Asymptotic Analysis of Non-self-adjoint Hill Operators

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Abstract

We obtain the uniform asymptotic formulas for the eigenvalues and eigenfunctions of the Sturm-Liouville operators $L_t(q)$ with a potential $q \in L_1[0, 1]$ and with t-periodic boundary conditions, $t \in (-\pi, \pi]$. Using these formulas, we find sufficient conditions on the potential q such that the number of spectral singularities in the spectrum of the Hill operator L(q) in $L_2(-\infty, \infty)$ is finite. Then we prove that the operator L(q)has no spectral singularities at infinity and it is an asymptotically spectral operator provided that the potential q satisfies the sufficient conditions.

Key Words: Asymptotic formulas, Hill operator, Spectral singularities, Spectral operator.

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1 Introduction and Preliminary Facts

Let L(q) be the Hill operator generated in $L_2(-\infty,\infty)$ by the expression

$$-y^{\prime\prime} + q(x)y,\tag{1}$$

where q(x) is a complex-valued summable function on [0,1] and q(x+1) = q(x) for a.e. $x \in (-\infty, \infty)$. It is well-known that (see [7], [23] for real and [5], [16]-[18] for complexvalued q) the spectrum S(L(q)) of the operator L(q) is the union of the spectra $S(L_t(q))$ of the Sturm-Liouville operators $L_t(q)$ for $t \in (-\pi, \pi]$, where $L_t(q)$ is the operator generated in $L_2[0,1]$ by (1) and by the boundary conditions

$$y(1) = e^{it}y(0), \ y'(1) = e^{it}y'(0).$$
 (2)

In this paper we obtain the asymptotic formulas, uniform with respect to $t \in (-\pi, \pi]$, for the eigenvalues and eigenfunctions of $L_t(q)$. (We recall that the formula f(k,t) = O(h(k)) is said to be uniform with respect to t in a set A if there exist positive constants M and N such that |f(k,t)| < M | h(k) | for all $t \in A$ and $|k| \ge N$.) Using these asymptotic formulas, we find sufficient conditions on the potential q such that the number of the spectral singularities in S(L(q)) is finite and in a certain sense L(q) is an asymptotically spectral operator.

The spectral expansion for the self-adjoint operator L(q) was constructed by Gelfand [7] and Titchmarsh [23]. Tkachenko [24] proved that the non-self-adjoint operator L(q) can be reduced to the triangular form if all eigenvalues of the operators $L_t(q)$ for all $t \in (-\pi, \pi]$ are simple. McGarvey [17] proved that L(q) is a spectral operator if and only if the projections of the operators $L_t(q)$ are bounded uniformly with respect to t in $(-\pi, \pi]$. However, in general, the eigenvalues of $L_t(q)$ are not simple and their projections are not uniformly bounded. For instance, Gasymov [6] investigated the operator L(q) with the potential q

which can be continued analytically onto upper half plane and proved that this operator (in particular L(q) with the simple potential $q(x) = e^{i2\pi x}$), has infinitely many spectral singularities. Note that the spectral singularities of the operator L(q) are the points of S(L(q)) in neighborhoods of which the projections of L(q) are not uniformly bounded. In [26] we proved that a number $\lambda = \lambda_n(t) \in S(L)$ is a spectral singularity of L(q) if and only if the operator $L_t(q)$ has an associated function at the point $\lambda_n(t)$. In [25] (see also [27]) we constructed the spectral expansion for the operator L(q) with a continuous and complex-valued potential. In [28], we obtained the asymptotic formulas for the eigenvalue and eigenfunction of $L_t(q)$ with $q \in L_1[0,1]$ and $t \neq 0, \pi$. Then using these formulas, we proved that the eigenfunctions and associated functions of L_t form a Riesz basis in $L_2[0,1]$ for $t \neq 0, \pi$ and constructed the spectral expansion for the operator L(q). (See also [13], [29], [30] for the spectral expansion of the differential operators with periodic coefficients). Recently, Gesztezy and Tkachenko [8], [9] proved two versions of a criterion for the Hill operator L(q) with $q \in L_2[0,1]$ to be a spectral operator of scalar type, one analytic and one geometric. The analytic version was stated in term of the solutions of the Hill's equation. The geometric version of the criterion used the algebraic and geometric properties of the spectra of the periodic/antiperiodic and Dirichlet boundary value problems.

Since the spectral property of L(q) is strongly connected with the operators $L_t(q)$ for $t \in (-\pi, \pi]$, let us discuss briefly the works devoted to $L_t(q)$. It is known that the operator $L_t(q)$ is Birkhoff regular [14]. In the case $t \neq 0, \pi$ it is strongly regular and the root functions of the operator $L_t(q)$ form a Riesz basis (this result was proved independently in [4], [12] and [19]). In the cases t = 0 and $t = \pi$, the operator $L_t(q)$ is not strongly regular. In the case when an operator is regular but not strongly regular the root functions generally do not form even usual basis. However, it is known [20], [21] that they can be combined in pairs, so that the corresponding 2-dimensional subspaces form a Riesz basis of subspaces.

Let us also briefly describe some historical developments related to the Riesz basis property of the root functions of the periodic and antiperiodic boundary value problems. We will focus only on the periodic problem. The antiperiodic problem is similar to the periodic one. In 1996 at a seminar in MSU Shkalikov formulated the following result. Assume that q(x) is a smooth potential,

$$q^{(k)}(0) = q^{(k)}(1), \quad \forall k = 0, 1, ..., s - 1$$
(3)

and $q^{(s)}(0) \neq q^{(s)}(1)$. Then the root functions of the operator $L_0(q)$ form a Riesz basis in $L_2[0,1]$. Kerimov and Mamedov [11] obtained the rigorous proof of this result in the case $q \in C^4[0,1], q(1) \neq q(0)$. Indeed, this result remains valid for an arbitrary $s \geq 0$ and it was obtained in Corollary 2 of [22].

Another approach is due to Dernek and Veliev [1]. The result was obtained in terms of the Fourier coefficients of the potential q. Namely, we proved that if the conditions

$$\lim_{n \to \infty} \frac{\ln |n|}{nq_{2n}} = 0,\tag{4}$$

$$q_{2n} \sim q_{-2n} \tag{5}$$

hold, then the root functions of $L_0(q)$ form a Riesz basis in $L_2[0,1]$, where $q_n =: (q, e^{i2\pi nx})$ is the Fourier coefficient of q and everywhere, without loss of generality, it is assumed that $q_0 = 0$. Here (.,.) denotes inner product in $L_2[0,1]$ and $a_n \sim b_n$ means that $a_n = O(b_n)$ and $b_n = O(a_n)$ as $n \to \infty$. Makin [15] improved this result. Using another method he proved that the assertion on the Riesz basis property remains valid if condition (5) holds, but condition (4) is replaced by a less restrictive one: $q \in W_1^s[0,1]$, (3) holds and $|q_{2n}| > c_0 n^{-s-1}$ with some $c_0 > 0$ for sufficiently large n, where s is a nonnegative integer. Besides, some

conditions which imply the absence of the Riesz basis property were presented in [15]. The results which we obtained in [22] are more general and cover all the results discussed above. Several theorems on the Riesz basis property of the root functions of the operator $L_0(q)$ were proved. One of the main results of [22] is the following:

Let $p \ge 0$ be an arbitrary integer, $q \in W_1^p[0,1]$ and (3) holds with some $s \le p$. Suppose that there is a number $\varepsilon > 0$ such that either the estimate

$$|q_{2n} - S_{2n} + 2Q_0 Q_{2n}| \ge \varepsilon n^{-s-2} \tag{6}$$

or the estimate

$$|q_{-2n} - S_{-2n} + 2Q_0 Q_{-2n}| \ge \varepsilon n^{-s-2} \tag{7}$$

hold, where $Q_k = (Q(x), e^{2\pi i kx})$ and $S_k = (S(x), e^{2\pi i kx})$ are the Fourier coefficients of

$$Q(x) = \int_0^x q(t) dt$$
 and $S(x) = Q^2(x)$.

Then the condition

$$q_{2n} - S_{2n} + 2Q_0 Q_{2n} \sim q_{-2n} - S_{-2n} + 2Q_0 Q_{-2n} \tag{8}$$

is necessary and sufficient for the root functions of $L_0(q)$ to form a Riesz basis. Moreover, if (6) (or (7)) and (8) hold then all the large eigenvalues of $L_0(q)$ are simple.

Some sharp results on the absence of the Riesz basis property were obtained by Djakov and Mitjagin [2]. Moreover, recently, Djakov and Mitjagin [3] obtained some interesting results about the Riesz basis property of the root functions of the operators $L_0(q)$ with trigonometric polynomial potentials. Here we do not formulate precisely the results of [2] and [3], since it will take some additional pages which are not related to our results. Very recently Gesztezy and Tkachenko [10] proved a criterion for the root functions of $L_0(q)$ to form a Riesz basis in terms of the spectra of the periodic and Dirichlet boundary value problems.

Next, we present some preliminary facts, from [28] and [1], which are needed in the following.

Result 1 (see [28]). The eigenvalues $\lambda_n(t)$ and the eigenfunctions $\Psi_{n,t}(x)$ of the operator $L_t(q)$ for $t \neq 0, \pi$, satisfy the following asymptotic formulas

$$\lambda_n(t) = (2\pi n + t)^2 + O(\frac{\ln|n|}{n}), \ \Psi_{n,t}(x) = e^{i(2\pi n + t)x} + O(\frac{1}{n}).$$
(9)

These asymptotic formulas are uniform with respect to t in $[\rho, \pi - \rho]$, where $\rho \in (0, \frac{\pi}{2})$ (see Theorem 2 of [9]). In other words, there exist positive numbers $N(\rho)$ and $M(\rho)$, independent of t, such that the eigenvalues $\lambda_n(t)$ for $t \in [\rho, \pi - \rho]$ and $|n| > N(\rho)$ are simple and the terms $O(\frac{1}{n})$, $O(\frac{|n||n|}{n})$ in (9) do not depend on t. **Result 2 (see [1])**. Let conditions (4) and (5) hold. Then:

(a) All sufficiently large eigenvalues of the operator $L_0(q)$ are simple. They consist of two sequences $\{\lambda_{n,1}: n > N_0\}$ and $\{\lambda_{n,2}: n > N_0\}$ satisfying

$$\lambda_{n,j} = (2\pi n)^2 + (-1)^j p_{2n} + O\left(\frac{\ln|n|}{n}\right)$$
(10)

for j = 1, 2, where $p_n = (q_n q_{-n})^{\frac{1}{2}}$. The corresponding eigenfunctions $\varphi_{n,j}(x)$ satisfy

$$\varphi_{n,j}(x) = e^{i2\pi nx} + \alpha_{n,j}e^{-i2\pi nx} + O(\frac{1}{n}),$$
(11)

where $\alpha_{n,j} \sim 1$, $\alpha_{n,j} = \frac{(-1)^j p_{2n}}{q_{2n}} + O\left(\frac{\ln|n|}{nq_{2n}}\right)$, j = 1, 2. (b) The root functions of $L_0(q)$ form a Riesz basis in $L_2[0, 1]$.

In [31] and [32] we generalized the results of [1] for the operators generated by the differential equation of order n > 2 and by the system of differential equations.

To summarize, in [1] and [22] we obtained the asymptotic formulas for the operators $L_t(q)$ with $t = 0, \pi$. In [28] we obtained the asymptotic formulas for the operators $L_t(q)$ which are uniform with respect to $t \in [\rho, \pi - \rho]$. In this paper, we obtain the uniform asymptotic formulas in much more complicated case of $t \in [0, \rho] \cup [\pi - \rho, \pi]$ (see Theorem 3 and 4). We note that in our description, some formulas of Section 2 are similar to those given in [1], [22] and [28], but here we wish to obtain the uniform, with respect to $t \in [0, \rho] \cup [\pi - \rho, \pi]$, formulas which are absent in these papers. We will focus only on the case $t \in [0, \rho]$. (The case $t \in [\pi - \rho, \pi]$ can be considered in the same way). Since the eigenvalues of $L_{-t}(q)$ coincide with those of $L_t(q)$, we obtain the uniform, with respect to t in $(-\pi, \pi]$, asymptotic formulas for the operators $L_t(q)$. These formulas imply that if the potential q satisfies certain conditions, then there exists a positive constant C independent of t such that all the eigenvalues of $L_t(q)$ lying outside the disk $\{\lambda \in \mathbb{C} : |\lambda| \leq C\}$ are simple for all the values of t in $(-\pi,\pi]$. Since the spectral singularities of the operator L(q) are contained in the set of multiple eigenvalues of $L_t(q)$, we obtain the sufficient conditions on q such that the Hill operator L(q) has at most finitely many spectral singularities. Moreover, we prove that if q satisfies these conditions then L(q) has no spectral singularity at infinity and in the sense of Definition 3 given in Section 3, the operator L(q) is an asymptotically spectral operator.

$\mathbf{2}$ Uniform Asymptotic Formulas for $L_t(q)$

It is well-known that the eigenvalues of $L_t(q)$ are the squares of the roots of the equation

$$F(\xi) = 2\cos t,\tag{12}$$

where $F(\xi) = \varphi'(1,\xi) + \theta(1,\xi)$, and $\varphi(x,\xi)$ and $\theta(x,\xi)$ are the solutions of the equation

$$-y^{''} + q(x)y = \xi^2 y$$

satisfying the initial conditions $\theta(0,\xi) = \varphi'(0,\xi) = 1$, $\theta'(0,\xi) = \varphi(0,\xi) = 0$. In [14] (see chapter 1, sec. 3) it was proved that

$$F(\xi) - 2\cos\xi = e^{|Im\xi|}\varepsilon(\xi), \quad \lim_{|\xi| \to \infty} \varepsilon(\xi) = 0.$$
(13)

Let us consider the functions $F(\xi) - 2\cos\xi$ and $2\cos\xi - \cos t$ on the circle

$$C(n,t,\rho) := \{\xi \in \mathbb{C} : |\xi - (2\pi n + t)| = 3\rho\},\tag{14}$$

where $t \in [0, \rho]$ and ρ is a sufficiently small fixed number. By (13) there exists a positive number $N(0,\rho)$ such that

$$|F(\xi) - 2\cos\xi| < \rho^2 \tag{15}$$

for $\xi \in C(n, t, \rho)$ whenever $n > N(0, \rho)$ and $t \in [0, \rho]$. On the other hand, using the Taylor formula of $\cos \xi$ at the point $2\pi n + t$ for $\xi = 2\pi n + t + 3\rho e^{i\alpha}$, where $\alpha \in (-\pi, \pi]$, and taking into account the inequalities $|\sin t| \le \rho$ and $|\cos t| > \frac{9}{10}$ for $t \in [0, \rho]$, we obtain

$$|2\cos\xi - 2\cos t| = 2|-3\rho e^{i\alpha}\sin t + \frac{9}{2}\rho^2 e^{2i\alpha}\cos t + O(\rho^4)| > 2\rho^2.$$
(16)

By the Rouche's theorem, it follows from (15) and (16) that equation (12) and

$$\cos\xi - \cos t = 0 \tag{17}$$

have the same number of the roots inside $C(n, t, \rho)$, where $n > N(0, \rho)$. Since equation (17) has 2 roots inside the circle $C(n, t, \rho)$, equation (12) has also 2 roots (counting multiplicity) inside this circle for $n > N(0, \rho)$. On the other hand, it is proved in [14] (see chapter 1, sec. 3) that the estimation

$$F(\xi) - 2\cos\xi = o(\cos\xi - \cos t)$$

holds on the boundaries of the admissible strip $K_n =: \{\xi : |\text{Re }\xi| < (2n+1)\pi\}$ for $t \in [0, \rho]$. Hence the number of the roots of equations (12) and (17) are the same in the strip K_n . Similarly, these equations have the same number of the roots in the set $K_{n+1}\setminus K_n$ for large n. The following remark follows from these arguments.

Remark 1 There exists a large number $N(0, \rho)$ such that the number of the roots of equations (12) lying in the strip K_N is 2N + 1. Denote these roots by $\xi_n(t)$ for

 $n = 0, \pm 1, \pm 2, ..., \pm N$. The roots of equation (12) lying outside K_N consist of the roots lying inside the contours $C(n, t, \rho)$, defined in (14), for $n > N(0, \rho)$. Moreover, (12) has two roots, denoted by $\xi_{n,1}(t)$ and $\xi_{n,2}(t)$, lying inside $C(n, t, \rho)$. Thus

$$|\xi_{n,j}(t) - (2\pi n + t)| < 3\rho, \ \forall |n| > N(0,\rho), \ t \in [0,\rho], \ j = 1,2.$$
(18)

Since the entire function $\frac{dF}{d\epsilon}$ has a finite number of zeros inside the circle

 $\{\xi \in \mathbb{C} : |\xi - 2\pi n| = 4\rho\}$ and this circle encloses $C(n,t,\rho)$ for all $t \in [0,\rho]$, there exist at most finite $t_1, t_2, ..., t_k$ from $(0,\rho)$ for which $\xi_n(t_k)$ is a double root of (12). Let $0 < t_1 < t_2 < ... < t_k < \rho$. By the implicit function theorem the functions $\xi_{n,1}(t)$ and $\xi_{n,2}(t)$ can be chosen as analytic in intervals $(0,t_1)$, (t_k,ρ) and (t_s,t_{s+1}) for s = 1,2,...,k-1. Let ξ be any limit point of $\xi_{n,j}(t)$ as $t \to t_s$. Since $F(\xi_{n,j}(t)) = 2 \cos t$ for j = 1,2 and F is continuous, we have $F(\xi) = 2 \cos t_s$. However, this equation has only one double root $\xi_{n,1}(t_s) = \xi_{n,2}(t_s)$ inside $C(n,t_s,\rho)$. Thus

$$\lim_{t \to t_s^-} \xi_{n,1}(t) = \lim_{t \to t_s^+} \xi_{n,1}(t) = \lim_{t \to t_s^-} \xi_{n,2}(t) = \lim_{t \to t_s^+} \xi_{n,1}(t) = \xi_{n,1}(t_s) = \xi_{n,2}(t_s)$$

for s = 1, 2, ..., k. This implies that the eigenvalues $\lambda_{n,1}(t) = \xi_{n,1}^2(t)$ and $\lambda_{n,2}(t) = \xi_{n,2}^2(t)$ of $L_t(q)$ can be chosen as continuous function on $(0, \rho)$. By the result of [28] (see introduction) $\lambda_{n,1}(\rho)$ and $\lambda_{n,2}(\rho)$ are the simple eigenvalues of $L_\rho(q)$ for $n > N(\rho)$. Moreover, if $q \in L_1[0,1]$ and (4), (5) hold then by the result of [1] $\lambda_{n,1}(0)$ and $\lambda_{n,2}(0)$ are the simple eigenvalues of L_0 for $n > N_0$. These arguments imply the continuity of the functions $\lambda_{n,1}(t)$, $\lambda_{n,2}(t)$ and

$$d_n(t) =: |\lambda_{n,1}(t) - \lambda_{n,2}(t)| \tag{19}$$

on $[0, \rho]$ for $n > N =: \max\{N(0, \rho), N(\rho), N_0\}$. By (18) we have

$$\left|\lambda_{n,i}(t) - (2\pi n + t)^2\right| < 15\pi n\rho \tag{20}$$

for $t \in [0, \rho]$, n > N and j = 1, 2. Thus for $t \in [0, \rho]$ and n > N the disk

$$D(n,t,\rho) =: \{\lambda \in \mathbb{C} : |\lambda - (2\pi n + t)^2| < 15\pi n\rho\}$$
(21)

contains two eigenvalues (counting multiplicity) $\lambda_{n,1}(t)$ and $\lambda_{n,2}(t)$ that are continuous function on the interval $[0, \rho]$. In addition to these eigenvalues, the operator $L_t(q)$ for $t \in [0, \rho]$ has only 2N + 1 eigenvalues.

Using (20), one can readily see that

$$\left|\lambda_{n,j}(t) - (2\pi(n-k) + t)^2\right| > |k| |2n-k|$$
(22)

for $k \neq 0, 2n$ and $t \in [0, \rho]$, where n > N and j = 1, 2. To obtain the uniform asymptotic formulas for the eigenvalues $\lambda_{n,j}(t)$ and normalized eigenfunctions $\Psi_{n,j,t}(x)$, we use (22) and the iteration of the formula

$$(\lambda_{n,j}(t) - (2\pi(n-k) + t)^2)(\Psi_{n,j,t}, e^{i(2\pi(n-k)+t)x}) = (q\Psi_{n,j,t}, e^{i(2\pi(n-k)+t)x})$$
(23)

which can be obtained obtained from $-\Psi_{n,j,t}^{''} + q\Psi_{n,j,t} = \lambda_{n,j}(t)\Psi_{n,j,t}$ by multiplying $e^{i(2\pi(n-k)+t)x}$. To iterate (23) we use the following lemma.

Lemma 1 For the right-hand side of (23) the following equality

$$(q\Psi_{n,j,t}, e^{i(2\pi(n-k)+t)x}) = \sum_{m=-\infty}^{\infty} q_m(\Psi_{n,j,t}, e^{i(2\pi(n-k-m)+t)x})$$
(24)

and inequality

$$\left| (q\Psi_{n,j,t}, e^{i(2\pi(n-k)+t)x}) \right| < 3M$$
 (25)

hold for all n > N, $k \in \mathbb{Z}$, j = 1, 2 and $t \in [0, \rho]$, where $M = \sup_{n \in \mathbb{Z}} |q_n|$, and N is defined in Remark 1. The eigenfunction $\Psi_{n,j,t}(x)$ satisfies the following, uniform with respect to $t \in [0, \rho]$, asymptotic formula

$$\Psi_{n,j,t}(x) = u_{n,j}(t)e^{i(2\pi n + t)x} + v_{n,j}(t)e^{i(-2\pi n + t)x} + h_{n,j,t}(x),$$
(26)

where $u_{n,j}(t) = (\Psi_{n,j,t}(x), e^{i(2\pi n + t)x}), v_{n,j}(t) = (\Psi_{n,j,t}(x), e^{i(-2\pi n + t)x}),$

$$(h_{n,j,t}, e^{i(\pm 2\pi n + t)x}) = 0, \quad ||h_{n,j,t}|| = O(\frac{1}{n}), \quad \sup_{x \in [0,1], \ t \in [0,\rho]} |h_{n,j,t}(x)| = O\left(\frac{\ln |n|}{n}\right), \quad (27)$$

$$|u_{n,j}(t)|^2 + |v_{n,j}(t)|^2 = 1 + O(\frac{1}{n^2}).$$
(28)

Proof. Equality (24) is obvious for $q \in L_2[0,1]$. For $q \in L_1[0,1]$ see Lemma 1 of [28]. Since $q\Psi_{n,j,t} \in L_1[0,1]$, we have

$$\lim_{|m|\to\infty} (q\Psi_{n,j,t}, e^{i(2\pi(n-k-m)+t)x}) = 0.$$

Therefore there exist C(t) and $k_0(t)$ such that

$$\max_{s\in\mathbb{Z}} \left| (q\Psi_{n,j,t}, e^{i(2\pi s+t)x}) \right| = \left| (q\Psi_{n,j,t}, e^{i(2\pi (n-k_0)+t)x}) \right| = C(t).$$

Now, using (22)-(24) and the obvious relations

$$|q_m| \le M, \quad \sum_{k \ne 0, 2n} \frac{1}{|k(2n-k)|} = O\left(\frac{\ln n}{n}\right)$$
 (29)

for $m \in \mathbb{Z}$, we obtain

$$C(t) = \left| (q\Psi_{n,j,t}, e^{i(2\pi(n-k_0)+t)x}) \right| = \left| \sum_{m=-\infty}^{\infty} q_m(\Psi_{n,j,t}, e^{i(2\pi(n-k_0-m)+t)x}) \right| = \left| q_{-k_0}(\Psi_{n,j,t}, e^{i(2\pi(n-k_0)+t)x}) + q_{2n-k_0}(\Psi_{n,j,t}, e^{i(-2\pi n+t)x}) \right| + \left| \sum_{\substack{m \neq -k_0, 2n-k_0}} q_m \frac{(q\Psi_{n,j,t}, e^{i(2\pi(n-k_0-m)+t)x})}{\lambda_{n,j}(t) - (2\pi(n-k_0-m)+t)^2} \right| \le 2M + \left| \sum_{\substack{m \neq -k_0, 2n-k_0}} \frac{MC(t)}{|\lambda_{n,j}(t) - (2\pi(n-k_0-m)+t)^2|} = 2M + C(t)O\left(\frac{\ln n}{n}\right) \right|$$

which implies that C(t) < 3M for all $t \in [0, \rho]$. Inequality (25) is proved. This with (23), (22) and (29) gives

$$\sum_{k \neq \pm n} \left| (\Psi_{n,j,t}, e^{i(2\pi k + t)x}) \right| = O\left(\frac{\ln n}{n}\right), \quad \sum_{k \neq \pm n} \left| (\Psi_{n,j,t}, e^{i(2\pi k + t)x}) \right|^2 = O\left(\frac{1}{n^2}\right).$$

Therefore decomposing $\Psi_{n,j,t}$ by basis $\{e^{i(2\pi k+t)x} : k \in \mathbb{Z}\}$ we get (26) and (27). The normalization condition $\|\Psi_{n,j,t}\| = 1$ with (26) and (27) implies (28)

Using (24) in (23), replacing k and m by 0 and n_1 respectively and then isolating the term containing the multiplicand $(\Psi_{n,j,t}, e^{i(-2\pi n+t)x})$ we obtain

$$(\lambda_{n,j}(t) - (2\pi n + t)^2)(\Psi_{n,j,t}, e^{i(2\pi n + t)x}) - q_{2n}(\Psi_{n,j,t}, e^{i(-2\pi n + t)x}) =$$
(30)
$$\sum_{\substack{n_1 \neq 0, 2n; \ n_1 = -\infty}}^{\infty} q_{n_1}(\Psi_{n,j,t}, e^{i(2\pi (n - n_1) + t)x}).$$

Now we iterate (30) by using the formula

$$(\Psi_{n,j,t}, e^{i(2\pi(n-n_1)+t)x}) = \sum_{n_2=-\infty}^{\infty} \frac{q_{n_2}(\Psi_{n,j,t}, e^{i(2\pi(n-n_1-n_2)+t)x})}{\lambda_{n,j}(t) - (2\pi(n-n_1)+t)^2}$$
(31)

obtained from (23) and (24). Taking into account that the denominator of the fraction in (31) is a large number for $n_1 \neq 0, 2n$ and $t \in [0, \rho]$ (see (22)), we iterate (30) as follows. Since this iteration is similar to that done in [1], here we give only the scheme of this iteration. First, we use (31) in (30), replacing the terms $(\Psi_{n,j,t}, e^{i(2\pi(n-n_1)+t)x})$ for $n_1 \neq 0, 2n$ in (30) by the right-hand side of (31) and get the summation with respect to n_1 and n_2 in the right-hand side of (30). We then isolate in this summation the terms containing one of the multiplicands $(\Psi_{n,j,t}, e^{i(2\pi n+t)x})$, $(\Psi_{n,j,t}, e^{i(-2\pi n+t)x})$ (i.e., terms with $n_1 + n_2 = 0, 2n$) and use (31) in the other terms. Repeating this process m-times, we obtain

$$(\lambda_{n,j}(t) - (2\pi n + t)^2 - A_m(\lambda_{n,j}(t), t))u_{n,j}(t) = (q_{2n} + B_m(\lambda_{n,j}(t), t))v_{n,j}(t) + R_m, \quad (32)$$

where

$$A_m(\lambda_{n,j}(t),t) = \sum_{k=1}^m a_k(\lambda_{n,j}(t),t), \ B_m(\lambda_{n,j}(t),t) = \sum_{k=1}^m b_k(\lambda_{n,j}(t),t),$$
(33)

$$a_{k}(\lambda_{n,j}(t),t) = \sum_{n_{1},n_{2},\dots,n_{k}} \frac{q_{n_{1}}q_{n_{2}}\dots q_{n_{k}}q_{-n_{1}-n_{2}-\dots-n_{k}}}{[\lambda_{n,j} - (2\pi(n-n_{1})+t)^{2}]\dots[\lambda_{n,j} - (2\pi(n-n_{1}-\dots-n_{k})+t)^{2}]},$$

$$b_{k}(\lambda_{n,j}(t),t) = \sum_{n_{1},n_{2},\dots,n_{k}} \frac{q_{n_{1}}q_{n_{2}}\dots q_{n_{k}}q_{2n-n_{1}-n_{2}-\dots-n_{k}}}{[\lambda_{n,j} - (2\pi(n-n_{1})+t)^{2}]\dots[\lambda_{n,j} - (2\pi(n-n_{1}-\dots-n_{k})+t)^{2}]},$$

$$R_{m} = \sum_{n_{1},n_{2},\dots,n_{m+1}} \frac{q_{n_{1}}q_{n_{2}}\dots q_{n_{m}}q_{n_{m+1}}(q\Psi_{n,j,t},e^{i(2\pi(n-n_{1}-\dots-n_{m+1})+t)x})}{[\lambda_{n,j} - (2\pi(n-n_{1})+t)^{2}]\dots[\lambda_{n,j} - (2\pi(n-n_{1}-\dots-n_{m+1})+t)^{2}]}.$$

Note that, here the sums are taken under conditions $n_s \neq 0$ and $n_1 + n_2 + ... + n_s \neq 0, 2n$ for s = 1, 2, ... Using (22), (25) and (29) one can easily verify that the equalities

$$a_{k} = O\left(\left(\frac{\ln|n|}{n}\right)^{k}\right), \ b_{k} = O\left(\left(\frac{\ln|n|}{n}\right)^{k}\right), \ R_{m} = O\left(\left(\frac{\ln|n|}{n}\right)^{m+1}\right)$$
(34)

hold uniformly with respect to t in $[0, \rho]$.

In the same way the relation

$$(\lambda_{n,j}(t) - (-2\pi n + t)^2 - A'_m(\lambda_{n,j}(t), t))v_{n,j}(t) = (q_{-2n} + B'_m(\lambda_{n,j}(t), t))u_{n,j}(t) = R'_m$$
(35)

can be obtained, where

$$A'_{m}(\lambda_{n,j}(t),t) = \sum_{k=1}^{m} a'_{k}(\lambda_{n,j}(t),t), \ B'_{m}(\lambda_{n,j}(t)) = \sum_{k=1}^{m} b'_{k}(\lambda_{n,j}(t),t),$$
$$a'_{k}(\lambda_{n,j}(t),t) = \sum_{n_{1},n_{2},\dots,n_{k}} \frac{q_{n_{1}}q_{n_{2}}\dots q_{n_{k}}q_{-n_{1}-n_{2}-\dots-n_{k}}}{[\lambda_{n,j} - (2\pi(n+n_{1})-t)^{2}]\dots[\lambda_{n,j} - (2\pi(n+n_{1}+\dots+n_{k})-t)^{2}]}$$

$$b'_{k}(\lambda_{n,j}(t),t) = \sum_{n_{1},n_{2},\dots,n_{k}} \frac{q_{n_{1}}q_{n_{2}}\dots q_{n_{k}}q_{-2n-n_{1}-n_{2}-\dots-n_{k}}}{[\lambda_{n,j} - (2\pi(n+n_{1})-t)^{2}]\dots[\lambda_{n,j} - (2\pi(n+n_{1}+\dots+n_{k}-t))^{2}]},$$

$$a'_{k} = O\left(\left(\frac{\ln|n|}{n}\right)^{k}\right), \ b'_{k} = O\left(\left(\frac{\ln|n|}{n}\right)^{k}\right), \ R'_{m} = O\left(\left(\frac{\ln|n|}{n}\right)^{m+1}\right),$$
(36)

 $n_s \neq 0, \, n_1 + n_2 + \ldots + n_s \neq 0, -2n \text{ for } s = 1, 2, \ldots, k.$

Now in (32) and (35) letting m tend to infinity, using (33), (34) and (36) we obtain

$$(\lambda_{n,j}(t) - (2\pi n + t)^2 - A(\lambda_{n,j}(t), t))u_{n,j}(t) = (q_{2n} + B(\lambda_{n,j}(t), t))v_{n,j}(t),$$
(37)

$$(\lambda_{n,j}(t) - (-2\pi n + t)^2 - A'(\lambda_{n,j}(t), t))v_{n,j}(t) = (q_{-2n} + B'(\lambda_{n,j}(t), t))u_{n,j}(t),$$
(38)

where

$$A(\lambda, t) = \sum_{k=1}^{\infty} a_k(\lambda, t), \ B(\lambda, t) = \sum_{k=1}^{\infty} b_k(\lambda, t), \ A' = \sum_{k=1}^{\infty} a'_k, \ B' = \sum_{k=1}^{\infty} b'_k.$$
(39)

The simplicity of the eigenvalues $\lambda_{n,j}(t)$ for $t \in [0, \rho]$, n > N, under some conditions on the potential q, are obtained from formulas (37) and (38) by the following scheme: First, we estimate the functions $A(\lambda, t), B(\lambda, t), A'(\lambda, t)$ and $B'(\lambda, t)$ (see Lemma 2). Next, using Lemma 2 we prove that the eigenvalues $\lambda_{n,j}(t)$ of $L_t(q)$ for $t \in [0, \rho]$ and n > N are the roots of the equation

$$(\lambda - (2\pi n + t)^2 - A(\lambda, t))(\lambda - (2\pi n - t)^2 - A'(\lambda, t)) = (q_{2n} + B(\lambda, t))(q_{-2n} + B'(\lambda, t))$$
(40)

$$\lambda = (2\pi n + t)^{2} + \frac{1}{2}(A(\lambda, t) + A'(\lambda, t)) - 4\pi nt - \sqrt{D(\lambda, t)}$$
(41)

and

$$\lambda = (2\pi n + t)^{2} + \frac{1}{2}(A(\lambda, t) + A'(\lambda, t)) - 4\pi nt + \sqrt{D(\lambda, t)},$$
(42)

where

$$D(\lambda, t) = (4\pi nt)^2 + q_{2n}q_{-2n} + D_1(\lambda, t) + D_2(\lambda, t),$$
(43)

$$D_{1}(\lambda, t) = 8\pi n t C(\lambda, t) + C^{2}(\lambda, t), \ C(\lambda, t) = \frac{1}{2} (A(\lambda, t) - A'(\lambda, t)),$$
(44)

$$D_{2}(\lambda, t) = q_{2n}B'(\lambda, t) + q_{-2n}B(\lambda, t) + B(\lambda, t)B'(\lambda, t).$$

To prove the simplicity of the eigenvalues $\lambda_{n,1}(t)$ and $\lambda_{n,2}(t)$ for $t \in [0, \rho]$ and n > N, we show that one of these eigenvalues satisfies (41) and the other one satisfies (42) and the roots of (41) and (42) are different. For this, we prove that the functions $B(\lambda, t)$, $B'(\lambda, t)$, $C(\lambda, t)$ (Lemma 3) and $\sqrt{D(\lambda, t)}$ (Lemma 4) satisfy some Lipschitz conditions. As a result, we find the conditions on q that guarantee the simplicity of those eigenvalues (Theorem 2).

Lemma 2 (a) The following equalities hold uniformly with respect to t in $[0, \rho]$:

$$A(\lambda_{n,j}(t),t) = O(n^{-1}), \ A'(\lambda_{n,j}(t),t) = O(n^{-1}).$$
(45)

(b) Let $q \in W_1^p[0,1]$ and (3) holds with some $s \leq p$. Then the equalities

$$B(\lambda_{n,j}(t),t) = o(n^{-s-1}), \ B'(\lambda_{n,j}(t),t) = o(n^{-s-1})$$
(46)

hold uniformly with respect to t in $[0, \rho]$.

Proof. (a) First, let us prove that

$$a_1(\lambda_{n,j}(t), t) = \frac{1}{4\pi^2} \sum_{k \neq 0, 2n} \frac{q_k q_{-k}}{k(2n-k)} + O\left(\frac{1}{n}\right).$$
(47)

Using (20) and the inequality $t < \rho$ and taking into account that ρ is a sufficiently small fixed number, one can see that if $|k| \leq 3 |n|$, then

$$|\lambda_{n,j} - (2\pi(n-k)+t)^2 - 4\pi^2 k(2n-k)| \le |n|.$$

Conversely, if |k| > 3 |n| then

$$|\lambda_{n,j} - (2\pi(n-k)+t)^2| > k^2 > n^2, |4\pi^2 k(2n-k)| > k^2 > n^2.$$

Therefore, taking into account the inequality in (29), we obtain

$$\sum_{k:|k|>3|n|} \frac{q_k q_{-k}}{\lambda_{n,j} - (2\pi(n-k)+t)^2} - \frac{1}{4\pi^2} \sum_{k:|k|>3|n|} \frac{q_k q_{-k}}{k(2n-k)} = O\left(\frac{1}{n}\right),$$
$$|\sum_{k:|k|\le3|n|,k\neq0,\,2n} \left(\frac{q_k q_{-k}}{\lambda_{n,j} - (2\pi(n-k)+t)^2} - \frac{q_k q_{-k}}{4\pi^2 k(2n-k)}\right)| \le$$

$$\sum_{k:|k|\leq 3|n|} \frac{M^2 |n|}{k^2 (2n-k)^2} \leq \sum_{k:|k|\leq |n|} \frac{M^2 |n|}{k^2 n^2} + \sum_{k:|n|<|k|\leq 3|n|} \frac{M^2 |n|}{n^2 (2n-k)^2} = O(\frac{1}{n}).$$

Thus (47) holds. In (47) grouping the terms $\frac{q_k q_{-k}}{k(2n-k)}$ and $\frac{q_{-k} q_k}{-k(2n+k)}$, we get

$$a_1(\lambda_{n,j}(t),t) = \frac{1}{4\pi^2} \sum_{k>0, \ k\neq 2n} \frac{q_k q_{-k}}{(2n+k)(2n-k)} + O\left(\frac{1}{n}\right).$$
(48)

To estimate the sum in (48) we consider, as done in [22], the function

$$G(x,n) = \int_0^x q(t)e^{-2\pi i(2n)t}dt - q_{2n}x$$

The Fourier coefficients $G_k(n) =: (G(x, n), e^{2\pi i kx})$ of G(x, n) are

$$G_k(n) = \frac{1}{2\pi i k} q_{2n+k}$$

for $k \neq 0$, and hence we have

$$G(x,n) = G_0(n) + \sum_{k \neq 2n} \frac{q_k}{2\pi i(k-2n)} e^{2\pi i(k-2n)x}.$$

Therefore, using the integration by parts and taking into account the obvious equalities G(1,n) = G(0,n) = 0, $G(x,n) - G_0(n) = O(1)$, we obtain

$$\frac{1}{4\pi^2} \sum_{k>0, \ k\neq 2n} \frac{q_k q_{-k}}{(2n+k)(2n-k)} = \int_0^1 (G(x,n) - G_0(n))^2 e^{2\pi i (4n)x} \, dx =$$
$$\frac{-1}{2\pi i (4n)} \int_0^1 2(G(x,n) - G_0(n))(q(x)e^{-2\pi i (2n)x} - q_{2n})e^{2\pi i (4n)x} \, dx = O\left(\frac{1}{n}\right).$$

This with (48) and (34) implies the first equality of (45). In the same way, we get the second equality of (45).

(b) If the assumptions of (b) hold, then

$$q_{2n} = o(n^{-s}), \ q_{n_1}q_{n_2}\cdots q_{n_k}q_{\pm 2n-n_1-n_2-\cdots-n_k} = o(n^{-s})$$
(49)

(see p. 655 of [22]). Using this and (22), in a standard way, we get

$$b_{k}(\lambda_{n,j}(t)) = o\left(\frac{\ln^{k} n}{n^{k+s}}\right) = o(n^{-s-1}), \quad b_{k}'(\lambda_{n,j}(t)) = o\left(\frac{\ln^{k} n}{n^{k+s}}\right) = o(n^{-s-1})$$
(50)

for $k \geq 2$. Now, it remains to prove that

$$b_1(\lambda_{n,j}(t)) = o(n^{-s-1}), \quad b'_1(\lambda_{n,j}(t)) = o(n^{-s-1}).$$
 (51)

Instead of the inequality in (29) using the equality $q_k q_{2n-k} = o(n^{-s})$ (see (49)) and arguing as in the proof of (47) we get

$$b_1(\lambda_{n,j}(t),t) = \frac{1}{4\pi^2} \sum_{k \neq 0, 2n} \frac{q_k q_{2n-k}}{k(2n-k)} + o(n^{-s-1})$$
(52)

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for all $t \in [0, \rho]$. In [22] (see p. 655) the summation in (52) is denoted by S_{2n} and it is proved that $S_{2n} = o(n^{-s-1})$ (see p. 658). Thus, from (52) we obtain the first equality of (51). In the same way, we get the second equality of (51). Now, (46) follows from (50) and (51)

Theorem 1 Let $q \in W_1^p[0,1]$ and (3) holds with some $s \leq p$. Suppose (5) holds and

$$|q_{2n}| > cn^{-s-1}$$
 (53)

for some c > 0. Then the eigenvalues $\lambda_{n,j}(t)$ of $L_t(q)$ for $t \in [0, \rho]$ and n > N are the roots of equation (40).

Proof. It follows from (53) and (46) that

$$q_{2n} + B(\lambda_{n,j}(t), t) \neq 0, \ q_{-2n} + B'(\lambda_{n,j}(t), t) \neq 0$$
(54)

for $t \in [0, \rho]$. Let us prove that this with formulas (37), (38) and (28) gives

$$u_{n,j}(t)v_{n,j}(t) \neq 0.$$
 (55)

If $u_{n,j}(t) = 0$ then by (28) $v_{n,j}(t) \neq 0$ and by (37) $q_{2n} + B(\lambda_{n,j}(t), t) = 0$ which contradicts (54). Similarly, if $v_{n,j}(t) = 0$ then by (28) and (38) $q_{-2n} + B'(\lambda_{n,j}(t), t) = 0$ which again contradicts (54). Therefore multiplying (37) and (38) side by side and then using (55), we get the proof of the theorem.

Now we consider some properties of the functions $A(\lambda, t)$, $B(\lambda, t)$, $A'(\lambda, t)$, $B'(\lambda, t)$ defined in (39) for $t \in [0, \rho]$ and $\lambda \in D(n, t, \rho)$, where $D(n, t, \rho)$ is the disk defined in (21).

Lemma 3 (a) There exists a constant K, independent of n > N and $t \in [0, \rho]$, such that

$$A(\lambda,t) - A(\mu,t) \mid < Kn^{-2} \mid \lambda - \mu \mid, \quad \mid A'(\lambda,t) - A'(\mu,t) \mid < Kn^{-2} \mid \lambda - \mu \mid,$$
 (56)

$$|C(\lambda_{n,j}(t),t)| < tKn^{-1}, |C(\lambda,t)| - C(\mu,t)| < tKn^{-2} |\lambda - \mu|$$
 (57)

for all $\lambda, \mu \in D(n, t, \rho)$, where N is defined in Remark 1.

(b) Let $q \in W_1^p[0,1]$, and (3) holds with some $s \leq p$. Then the functions $b_k(\lambda,t)$, $b'_k(\lambda,t)$, $B(\lambda,t)$, $B'(\lambda,t)$ for $\lambda, \mu \in D(n,t,\rho)$, k = 1, 2, ..., satisfy the following, uniform with respect to t in $[0,\rho]$, equality

$$f(\lambda, t) - f(\mu, t) = (\lambda - \mu)o(n^{-s-2}).$$
 (58)

Proof. (a) If $\lambda \in D(n, t, \rho)$, then

$$\left|\lambda - (2\pi(n-k)+t)^2\right| > |k| \left|2n-k\right|, \ \forall k \neq 0, 2n.$$
 (59)

To prove estimations (56) and (57) we use (59) and the following obvious equality

$$\sum_{k \neq 0, -2n} \frac{1}{\mid k^s (2n-k)^m \mid} = O\left(\frac{1}{n^p}\right)$$
(60)

if $\max\{s, m\} \ge 2$, where $p = \min\{s, m\} \ge 1$. By (60) and (59) the series in the formulas for the functions $a_k(\lambda, t)$, $a'_k(\lambda, t)$, $b_k(\lambda, t)$, $b'_k(\lambda, t)$ converge uniformly in a neighborhood of λ . It implies that these functions continuously depend on λ . Moreover, estimations (34) and (36) hold if we replace $\lambda_{n,j}(t)$ by λ . Therefore the series in the formulas for the functions $A(\lambda, t)$, $B(\lambda, t)$, $A'(\lambda, t)$ and $B'(\lambda, t)$ converge uniformly in a neighborhood of λ . Using (59) and (60) one can easily verify that these series can be differentiated, with respect to λ , term

by term. Moreover, taking into account the inequality

$$\mid \frac{d}{d\lambda} \big(\frac{1}{\lambda - (2\pi(n-k)+t)^2} \big) \mid \leq \frac{1}{k^2(2n-k)^2}$$

and (60), we see that the absolute values of the derivatives of $a_k(\lambda, t)$, $a'_k(\lambda, t)$, $b_k(\lambda, t)$, $b'_k(\lambda, t)$, with respect to λ is $O(n^{-k-1})$. Therefore, these functions satisfy the condition

$$g(\lambda, t) - g(\mu, t) = (\lambda - \mu)O(n^{-k-1}).$$
 (61)

Now (56) follows from (61).

To prove the first inequality of (57) we use substitutions $-n_1 - n_2 - \cdots - n_k = j_1$,

 $n_2 = j_k$, $n_3 = j_{k-1}, \ldots, n_k = j_2$ in the formula for the expression a'_k . Then the inequalities for the forbidden indices $n_p \neq 0$, $n_1 + n_2 + \cdots + n_p \neq 0$, -2n for $1 \leq p \leq k$ in the formula for a'_k take the form $j_p \neq 0$, $j_1 + j_2 + \ldots + j_p \neq 0$, 2n for $1 \leq p \leq k$, and

$$a_{k}'(\lambda_{n,j}(t)) = \sum_{n_{1},n_{2},\dots,n_{k}} \frac{q_{n_{1}}q_{n_{2}}\dots q_{n_{k}}q_{-n_{1}-n_{2}-\dots-n_{k}}}{[\lambda_{n,j} - (2\pi(n-n_{1})-t)^{2}]\dots[\lambda_{n,j} - (2\pi(n-n_{1}-\dots-n_{k})-t)^{2}]}.$$

Using (22) and (60) one can readily see that

$$\sum_{\substack{k=-\infty,\\k\neq 0,2n}}^{\infty} \mid \frac{1}{\lambda_{n,j}(t) - (2\pi(n-k) + t)^2} - \frac{1}{\lambda_{n,j}(t) - (2\pi(n-k) - t)^2} \mid = tO(\frac{1}{n}).$$

This with the inequality in (29) implies the first inequality in (57). Now arguing as in the proof of (56), we get the proof of the second inequality of (57).

(b) Using (49) and repeating the proof of (56) we get the proof of (b)

Lemma 4 Suppose the conditions of Theorem 1 hold. If at least one of the inequalities

$$\operatorname{Re} q_{2n} q_{-2n} \ge 0, \tag{62}$$

$$|\operatorname{Im} q_{2n}q_{-2n}| \ge \varepsilon |q_{2n}q_{-2n}| \tag{63}$$

are satisfied for some $\varepsilon > 0$ and for n > N, where N is defined in Remark 1, then there exists a constant c_1 such that

$$\left|\sqrt{D(\lambda_{n,1}(t),t)} - \sqrt{D(\lambda_{n,2}(t),t)}\right| < c_1 n^{-1} \mid \lambda_{n,1}(t) - \lambda_{n,2}(t) \mid$$
(64)

for n > N and $t \in [0, \rho]$.

Proof. First let us prove that

$$|D(\lambda_{n,j}(t),t)| > \frac{\varepsilon}{4} (|q_{-2n}q_{2n}| + (4\pi nt)^2),$$
(65)

$$D(\lambda_{n,j}(t),t) = ((4\pi nt)^2 + q_{2n}q_{-2n})(1+o(1))$$
(66)

for n > N and $t \in [0, \rho]$. If follows from (44), (57) and (53), (46), (5) that

$$D_1(\lambda_{n,j}(t),t) = t^2 O(1), \ D_2(\lambda_{n,j}(t),t) = o(q_{2n}n^{-s-1}) = o(q_{2n}q_{-2n}).$$
(67)

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Therefore, we have

$$D_1(\lambda_{n,j}(t), t) + D_2(\lambda_{n,j}(t), t) = o(|q_{-2n}q_{2n}| + (4\pi nt)^2).$$
(68)

Thus, by (43), to prove (65) and (66) it is enough to show that

$$\left|q_{-2n}q_{2n} + (4\pi nt)^{2}\right| > \frac{\varepsilon}{3}(\left|q_{-2n}q_{2n}\right| + (4\pi nt)^{2}).$$
(69)

For this we consider two cases. First case: $(4\pi nt)^2 \leq 2 |q_{2n}q_{-2n}|$). Then

$$|q_{2n}q_{-2n}| \ge \frac{1}{3}(|q_{2n}q_{-2n}| + (4\pi nt)^2).$$
(70)

If (62) holds then $|q_{2n}q_{-2n} + (4\pi nt)^2| \ge |q_{2n}q_{-2n}|$. Therefore, (69) follows from (70). If (63) holds then $|q_{2n}q_{-2n} + (4\pi nt)^2| \ge |\operatorname{Im} q_{2n}q_{-2n}| \ge \varepsilon |q_{2n}q_{-2n}|$ and again (69) follows from (70). Now let us consider the second case: $(4\pi nt)^2 > 2 |q_{2n}q_{-2n}|$). Then

$$\left|q_{2n}q_{-2n} + (4\pi nt)^2\right| > (4\pi nt)^2 - \left|q_{2n}q_{-2n}\right| > \frac{1}{3}(\left|q_{2n}q_{-2n}\right| + (4\pi nt)^2),$$

that is, (69) holds. Thus, (65) and (66) are proved.

Using (43), (57), (49), (46) and Lemma 3(b) one can easily verify that

$$|D(\lambda_{n,1}(t),t) - D(\lambda_{n,2}(t),t)| \le (5\pi t^2 K n^{-1} + n^{-2s-2}) |\lambda_{n,1}(t) - \lambda_{n,2}(t)|.$$
(71)

On the other hand, it follows from (67), (66) and (53), (5) that there exists a constant c_2 such that

$$\left|\sqrt{D(\lambda_{n,1}(t),t)} + \sqrt{D(\lambda_{n,2}(t),t)}\right| = \left|2\sqrt{D(\lambda_{n,1}(t),t)}(1+o(1))\right| > c_2(n^{-s-1}+nt).$$
(72)

Thus, from (71) and (72) we obtain (64). \blacksquare

Now using Lemma 2-4 and Theorem 1 we prove the following main result.

Theorem 2 Let $q \in W_1^p[0,1]$ and (3) holds with some $s \leq p$. Suppose (5) and (53) are satisfied. If at least one of the inequalities (62) and (63) hold, then the eigenvalues $\lambda_{n,j}(t)$ of $L_t(q)$ for n > N, j = 1, 2 and $t \in [0, \rho]$ are simple and

$$\lambda_{n,j}(t) = (2\pi n + t)^2 + \frac{1}{2} (A(\lambda_{n,j}(t), t) + A'(\lambda_{n,j}(t), t)) - 4\pi nt + (-1)^j \sqrt{D(\lambda_{n,j}(t), t)}.$$
 (73)

Proof. If both eigenvalues $\lambda_{n,1}(t)$ and $\lambda_{n,2}(t)$ of the operator $L_t(q)$ lying in the disk $D(n, t, \rho)$ (see (21)) satisfy equation (41), then

$$\lambda_{n,1}(t) - \lambda_{n,2}(t) = \left[\frac{1}{2}(A(\lambda_{n,1}(t), t) - A(\lambda_{n,2}(t), t)) + \frac{1}{2}(A'(\lambda_{n,1}(t), t) - A'(\lambda_{n,2}(t), t)] + (74) \left[\sqrt{D(\lambda_{n,1}(t), t)} - \sqrt{D(\lambda_{n,2}(t), t)}\right].$$

By (56) we have

$$|A(\lambda_{n,1},t) - A(\lambda_{n,2},t) + A'(\lambda_{n,1},t) - A'(\lambda_{n,2},t)| < 2Kn^{-2} |\lambda_{n,1}(t) - \lambda_{n,2}(t)|.$$
(75)

Thus, using this and (64) in (74) we get

$$\lambda_{n,1}(t) = \lambda_{n,2}(t). \tag{76}$$

In the same way, we prove that if both $\lambda_{n,1}(t)$ and $\lambda_{n,2}(t)$ satisfy (42) then (76) holds.

Now suppose that one of them, say $\lambda_{n,1}(t)$, satisfies (41) and the other $\lambda_{n,2}(t)$ satisfies (42). Then

$$\lambda_{n,1}(t) - \lambda_{n,2}(t) = \left[\frac{1}{2}(A(\lambda_{n,1},t) - A(\lambda_{n,2},t)) + \frac{1}{2}(A'(\lambda_{n,1},t) - A'(\lambda_{n,2},t)] + \left(\sqrt{D(\lambda_{n,1}(t),t)} + \sqrt{D(\lambda_{n,2}(t),t)}\right].$$
(77)

Therefore, by (75) and (72) there exists a constant c_3 such that

$$|\lambda_{n,1}(t) - \lambda_{n,2}(t)| > c_3(n^{-s-1} + nt).$$
(78)

Now it follows from (76) and (78) that the value $d_n(t)$ of the function d_n , defined in (19), for $t \in [0, \rho]$ belongs to the union of the disjoint sets $(c_3(n^{-s-1} + nt), \infty)$ and $\{0\}$. Moreover, as it is proved in Remark 1, d_n is a continuous function on $[0, \rho]$ which implies that the set $\{d_n(t) : t \in [0, \rho]\}$ is a connected set. Therefore, taking into account that $d_n(\rho) \in (c_3(n^{-s-1} + nt), \infty)$ (see (9)), we get $\{d_n(t) : t \in [0, \rho]\} \in (c_3(n^{-s-1} + nt), \infty)$. This means that $\lambda_{n,1}(t)$ and $\lambda_{n,2}(t)$ are different simple eigenvalues and one of them satisfies (41) and the other satisfies (42). Without loss of generality, it can be assumed that (73) holds where square root in (73) is taken with positive real part. Note that

$$\operatorname{Re}(\sqrt{D(\lambda_{n,j}(t),t)}) \neq 0 \tag{79}$$

due to the following reason. By (53), (62)) and (63) $\left| \arg(q_{-2n}q_{2n} + (4\pi nt)^2) \right| < \pi - \alpha$, where $\alpha \in (0, \pi)$. Thus, by (65) and (66) $\arg D(\lambda_{n,j}(t), t) \neq \pi$ and hence (79) holds

Now to prove the uniform asymptotic formulas for the eigenfunctions $\Psi_{n,j,t}(x)$ we need consider $v_{n,j}(t)$ and $u_{n,j}(t)$ (see (26)-(28)). Namely, we use the following

Lemma 5 Suppose that all conditions of Theorem 2 hold. Let $\lambda_{n,j}(t)$ be the eigenvalue of $L_t(q)$ satisfying (73), and $\Psi_{n,j,t}(x)$ be the corresponding eigenfunction. Then the relations

$$v_{n,1}(t) \sim 1, \ u_{n,2}(t) \sim 1$$
 (80)

hold uniformly for $t \in [0, \rho]$.

Proof. Multiplying (37) and (38) by $v_{n,j}(t)$ and by $u_{n,j}(t)$ respectively and then subtracting each other, we get

$$(-8\pi nt + A'(t) - A(t))u_{n,j}(t)v_{n,j}(t) = (q_{2n} + B(t))v_{n,j}^2(t) - (q_{-2n} + B'(t))u_{n,j}^2(t), \quad (81)$$

where, for brevity, $A(\lambda_{n,j}(t), t), A'(\lambda_{n,j}(t), t), B(\lambda_{n,j}(t), t)$ and $B'(\lambda_{n,j}(t), t)$ are denoted by A(t), A'(t), B(t) and B'(t) respectively.

First, suppose that $nt \leq |q_{2n}|$. Then it follows from (57) that $A'(t) - A(t) = o(q_{2n})$ and

$$|-8\pi nt + A'(t) - A(t)| < 9\pi |q_{2n}|.$$
(82)

On the other hand, relations (5), (46) and (53) imply that

$$q_{2n} + B(\lambda_{n,j}(t), t) \sim q_{-2n} + B'(\lambda_{n,j}(t), t) \sim q_{2n}.$$
(83)

Therefore, using (81)-(83) and taking into account that if the relation $u_{n,j}(t) \sim v_{n,j}(t)$ does not hold then $u_{n,j}(t)v_{n,j}(t) = o(1)$ (see (28)), we obtain $u_{n,j}(t) \sim v_{n,j}(t) \sim 1$ for j = 1, 2. Thus, (80) holds for the case $nt \leq |q_{2n}|$.

Now consider the case $nt > |q_{2n}|$. Using (73) in (37) and (38), we obtain

$$(C(t) - 4\pi nt + (-1)^j \sqrt{D(t)})u_{n,j}(t) = (q_{2n} + B(t))v_{n,j}(t),$$
(84)

$$(-C(t) + 4\pi nt + (-1)^{j}\sqrt{D(t)})v_{n,j}(t) = (q_{-2n} + B'(t))u_{n,j}(t).$$
(85)

Since $\operatorname{Re}(\sqrt{D(\lambda_{n,j}(t),t)}) > 0$, it follows from (84) for j = 1 and (57) that

$$|C(\lambda_{n,1}(t),t) - 4\pi nt - \sqrt{D(\lambda_{n,1}(t),t)}| \ge \operatorname{Re}(4\pi nt(1 + O(n^{-2})) + \sqrt{D(\lambda_{n,1}(t),t)}) > |q_{2n}|.$$

Using this and (83) in (84) for j = 1 we get $v_{n,1}(t) \sim 1$. In the same way we obtain the second relation of (80) from (85) for j = 2.

To obtain the asymptotic formulas of arbitrary accuracy we define successively the following functions

$$F_{n,j,1}(t) = (2\pi n + t)^2 - 4\pi nt + (-1)^j \sqrt{(4\pi nt)^2 + q_{2n}q_{-2n}},$$

$$F_{n,j,m+1}(t) = (2\pi n + t)^2 + \frac{1}{2} (A(F_{n,j,m},t) + A'(F_{n,j,m},t)) - 4\pi nt + (-1)^j \sqrt{D(F_{n,j,m},t)}$$

for m = 1, 2, ... Moreover we use the functions A^* , B^* which are obtained from A, B respectively by replacing q_{n_1} with $e^{i(2\pi(n-n_1)+t)x}$.

Theorem 3 (a) If the conditions of Theorem 2 hold, then the eigenvalue $\lambda_{n,j}(t)$ satisfies the following, uniform with respect to $t \in [0, \rho]$, formulas

$$\lambda_{n,j}(t) = (2\pi n + t)^2 - 4\pi nt + (-1)^j \sqrt{(4\pi nt)^2 + q_{2n}q_{-2n}} + O(\frac{1}{n}),$$
(86)

$$\lambda_{n,j}(t) = F_{n,j,m}(t) + O(\frac{1}{n^m}), \ m = 1, 2, \dots$$
(87)

(b) The normalized eigenfunction $\Psi_{n,j,t}(x)$ corresponding to $\lambda_{n,j}(t)$ is $\frac{\varphi_{n,j,t}(x)}{\|\varphi_{n,j,t}(x)\|}$, where $\varphi_{n,j,t}(x)$ satisfies the following, uniform with respect to $t \in [0, \rho]$, formulas

$$\begin{split} \varphi_{n,1,t}(x) &= e^{i(-2\pi n+t)x} + \alpha_{n,1}e^{i(2\pi n+t)x} + A^*(F_{n,1,m},t) + \alpha_{n,1}B^*(F_{n,1,m},t) + O(n^{-m-1}), \\ \varphi_{n,2,t}(x) &= e^{i(2\pi n+t)x} + \alpha_{n,2}e^{i(-2\pi n+t)x} + A^*(F_{n,2,m},t) + \alpha_{n,2}B^*(F_{n,2,m},t) + O(n^{-m-1}), \\ \alpha_{n,1}(t) &= \frac{C(F_{n,1,m},t) - 4\pi nt - \sqrt{D(F_{n,1,m},t)}}{q_{-2n} + B'(F_{n,1,m},t)} + O(\frac{1}{q_{2n}n^{m+1}}) = O(1), \\ \alpha_{n,2}(t) &= \frac{-C(F_{n,2,m},t) - 4\pi nt + \sqrt{D(F_{n,2,m},t)}}{q_{2n} + B(F_{n,2,m},t)} + O(\frac{1}{q_{2n}n^{m+1}}) = O(1). \end{split}$$

Proof. By (73) and (45) to prove (86) it is enough to show that

$$\sqrt{D(\lambda_{n,j}(t),t)} = \sqrt{(4\pi nt)^2 + q_{2n}q_{-2n}} + O(\frac{1}{n}).$$
(88)

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Using (66) and (69) one can easily verify that

$$|\sqrt{D(\lambda_{n,j}(t),t)} + \sqrt{(4\pi nt)^2 + q_{2n}q_{-2n}}| \ge \sqrt{\frac{\varepsilon}{6}}(4\pi nt + |\sqrt{q_{2n}q_{-2n}}|)$$

Therefore there exists a constant c_4 such that

$$\left|\sqrt{D(\lambda_{n,j}(t),t)} - \sqrt{(4\pi nt)^{2} + q_{2n}q_{-2n}}\right| = \left|\frac{D_{1}(\lambda_{n,j}(t),t) + D_{2}(\lambda_{n,j}(t),t)}{\sqrt{D(\lambda_{n,j}(t),t)} + \sqrt{(4\pi nt)^{2} + q_{2n}q_{-2n}}}\right| \leq c_{4}(\left|\frac{D_{1}(\lambda_{n,j}(t),t)}{4\pi nt}\right| + \left|\frac{D_{2}(\lambda_{n,j}(t),t)}{\sqrt{q_{2n}q_{-2n}}}\right|).$$

$$(89)$$

Moreover, from (67) and (5) we obtain

n

$$\frac{D_1(\lambda,t)}{4\pi nt} = O(\frac{1}{n}), \ \frac{D_2(\lambda,t)}{\sqrt{q_{2n}q_{-2n}}} = o(n^{-s-1}).$$
(90)

Hence, (88) follows from (89) and (90). Thus (86) is proved.

It follows from Lemma 3 and from the proof of (64) that the functions $A(\lambda, t)$, $A'(\lambda, t)$, $B(\lambda, t)$, $B'(\lambda, t)$ and $\sqrt{D(\lambda, t)}$ satisfy the equality

$$f(F_{n,j,k}(t) + O(n^{-k}), t) = f(F_{n,j,k}(t), t) + O(n^{-k-1}).$$
(91)

Now, we prove (87) by induction. It is proved for m = 1 (see (86) and the definition of $F_{n,j,1}(t)$). Assume that (87) is true for m = k. Substituting the value of $\lambda_{n,j}(t)$ given by (87) for m = k in the right-hand side of (73) and using (91) we get (87) for m = k + 1.

(b) Writing the decomposition of the normalized eigenfunction $\Psi_{n,j,t}(x)$ corresponding to the eigenvalue $\lambda_{n,j}(t)$ by the basis $\{e^{i(2\pi(n-n_1)+t)x} : n_1 \in \mathbb{Z}\}$, we obtain

$$\Psi_{n,j,t}(x) - u_{n,j}(t)e^{i(2\pi n+t)x} - v_{n,j}(t)e^{i(-2\pi n+t)x} =$$

$$\sum_{\substack{n \neq 0, 2\pi; n_1 = -\infty}}^{\infty} (\Psi_{n,j,t}(x), e^{i(2\pi (n-n_1)+t)x})e^{i(2\pi (n-n_1)+t)x}.$$
(92)

The right-hand side of (92) can be obtained from the right-hand side of (30) by replacing q_{n_1} with $e^{i2\pi(n-n_1)x}$. Since (37) is obtained from (30) by iteration, doing the same, we obtain

$$\Psi_{n,j,t}(x) = u_{n,j}(t)e^{i(2\pi n+t)x} + v_{n,j}(t)e^{i(-2\pi n+t)x} + u_{n,j}(t)A^*(\lambda_{n,j},t) + v_{n,j}(t)B^*(\lambda_{n,j},t)$$
(93)

from (92). First, let us consider the case j = 2. Using (87) and (91) in (37), taking into account (46) and (53), we get

$$\frac{v_{n,2}(t)}{u_{n,2}(t)} = \frac{-C(F_{n,2,m}(t),t) - 4\pi nt + \sqrt{D(F_{n,2,m}(t),t)}}{q_{2n} + B(F_{n,2,m}(t),t)} + O(\frac{1}{q_{2n}n^{m+1}}), \tag{94}$$

where m > s. Now, dividing both sides of (93) by $u_{n,2}(t)$, and denoting

$$\alpha_{n,2}(t) = \frac{v_{n,2}(t)}{u_{n,2}(t)}, \ \varphi_{n,2,t}(x) = \frac{\Psi_{n,2,t}(x)}{u_{n,2}(t)},$$

we obtain

$$\varphi_{n,2,t}(x) = e^{i(2\pi n + t)x} + \alpha_{n,2}(t)e^{i(-2\pi n + t)x} + A^*(\lambda_{n,2}(t), t) + \alpha_{n,2}(t)B^*(\lambda_{n,2}(t), t).$$
(95)

Here $\alpha_{n,2}(t) = O(1)$ due to (80). On the other hand, one can readily see that the functions $A^*(\lambda, t)$ and $B^*(\lambda, t)$ also satisfy (91). Therefore, from (95) we obtain the proof of (b) for j = 2. In the same way, we get the proof of (b) for j = 1.

To obtain the asymptotic formulas for the eigenvalues $\lambda_{n,j}(t)$ for $t \in [\pi - \rho, \pi]$ instead of (30) we use the formula

$$(\lambda_{n,j}(t) - (2\pi n + t)^2)(\Psi_{n,j,t}, e^{i(2\pi n + t)x}) - q_{2n+1}(\Psi_{n,j,t}, e^{i(-2\pi (n+1)+t)x}) =$$
(96)
$$\sum_{n_1 \neq 0, 2n+1; n_1 = -\infty}^{\infty} q_{n_1}(\Psi_{n,j,t}, e^{i(2\pi (n-n_1)+t)x}).$$

From (30) we obtained (37), (38). In the same way, from (96) we get

$$(\lambda_{n,j}(t) - (2\pi n + t)^2 - \widetilde{A}(\lambda_{n,j}(t), t))u_{n,j}(t) = (q_{2n+1} + \widetilde{B}(\lambda_{n,j}(t), t))v_{n,j}(t),$$

 $(\lambda_{n,j}(t) - (-2\pi(n+1) + t)^2 - \widetilde{A}'(\lambda_{n,j}(t), t))v_{n,j}(t) = (q_{-2n-1} + \widetilde{B}'(\lambda_{n,j}(t), t))u_{n,j}(t),$

where

$$\widetilde{A}(\lambda,t) = \sum_{k=1}^{\infty} \widetilde{a}_k(\lambda,t), \ \widetilde{B} = \sum_{k=1}^{\infty} \widetilde{b}_k, \ \widetilde{A}' = \sum_{k=1}^{\infty} \widetilde{a}'_k, \ \widetilde{B}' = \sum_{k=1}^{\infty} \widetilde{b}'_k.$$

Here $\tilde{a}_k, \tilde{a}'_k, \tilde{b}_k, \tilde{b}'_k$ differ from a_k, a'_k, b_k, b'_k respectively, in the following sense. The sums in $\tilde{a}_k, \tilde{a}'_k, \tilde{b}_k, \tilde{b}'_k$ are taken under conditions $n_1 + n_2 + \ldots + n_s \neq 0, \pm (2n + 1)$ instead of the condition $n_1 + n_2 + \ldots + n_s \neq 0, \pm 2n$ for $s = 1, 2, \ldots, k$. Besides in $\tilde{b}_k, \tilde{b}'_k$ the multiplicand $q_{\pm 2n - n_1 - n_2 - \ldots - n_k}$ of b_k, b'_k is replaced by $q_{\pm (2n+1) - n_1 - n_2 - \ldots - n_k}$. Moreover, instead of $F, \alpha_{n,j}, A^*, B^*$ we use $\tilde{F}, \tilde{\alpha}_{n,j}, \tilde{A}^*, \tilde{B}^*$ that can be defined in a similar way. Thus, instead of (5), (53), (62) and (63) using the relations

$$q_{2n+1} \sim q_{-2n-1}, \quad |q_{2n+1}| > cn^{-s-1},$$
(97)

$$\operatorname{Re} q_{2n+1} q_{-2n-1} \ge 0,$$
 (98)

$$|\operatorname{Im} q_{2n+1}q_{-2n-1}| \ge \varepsilon |q_{2n+1}q_{-2n-1}| \tag{99}$$

respectively and repeating the proof of Theorems 1-3 we get:

Theorem 4 Let $q \in W_1^p[0,1]$ and (3) holds with some $s \leq p$. Suppose (97) and at least one of the inequalities (98) and (99) hold. Then there exists $N(\pi, \rho)$ such that the eigenvalues $\lambda_{n,j}(t)$ for $n > N(\pi, \rho)$ and $t \in [\pi - \rho, \pi]$ are simple and satisfy the formulas

$$\lambda_{n,j}(t) = (2\pi n + t)^2 - 2\pi (2n+1)(t-\pi) + (-1)^j \sqrt{(2\pi (2n+1)(t-\pi))^2 + q_{2n+1}q_{-2n-1}} + O(\frac{1}{n}),$$
(100)

$$\lambda_{n,j}(t) = \widetilde{F}_{n,j,m}(t) + O(n^{-m}), \ m = 1, 2, \dots$$
(101)

The normalized eigenfunction $\Psi_{n,j,t}(x)$ corresponding to $\lambda_{n,j}(t)$ is $\frac{\varphi_{n,j,t}(x)}{\|\varphi_{n,j,t}(x)\|}$, where $\varphi_{n,j,t}(x)$ satisfies the following, uniform with respect to $t \in [\pi - \rho, \pi]$, formulas

$$\varphi_{n,1,t} = e^{i(-2\pi(n+1)+t)x} + \tilde{\alpha}_{n,1}(t)e^{i(2\pi n+t)x} + \tilde{A}^*(\tilde{F}_{n,1,m},t) + \tilde{\alpha}_{n,1}(t)\tilde{B}^*(\tilde{F}_{n,1,m},t) + O(n^{-m-1}),$$

$$\varphi_{n,2,t} = e^{i(2\pi n+t)x} + \tilde{\alpha}_{n,2}(t)e^{i(-2\pi(n+1)+t)x} + \tilde{A}^*(\tilde{F}_{n,2,m},t) + \alpha_{n,2}(t)\tilde{B}^*(\tilde{F}_{n,2,m},t) + O(n^{-m-1})$$

The following remark follows from Remark 1 and Theorems 1-4.

Remark 2 Suppose the conditions of Theorem 2 and Theorem 4 hold. One can readily see that (86) for t = 0 and $t = \rho$ gives formulas (10) and (9) respectively, if we use the notation: $\lambda_{-n}(t) =: \lambda_{n,1}(t)$ for n = 1, 2, ... and $\lambda_n(t) =: \lambda_{n,2}(t)$ for n = 0, 1, 2, ... Note that we use both notations $\lambda_n(t)$ and $\lambda_{n,j}(t)$. If the notation $\lambda_n(t)$ is used, then the corresponding eigenfunction and Fourier coefficients (see (26)) are denoted by $\Psi_{n,t}(x)$ and $u_n(t)$, $v_n(t)$. Similarly, (100) for $t = \pi$ and $t = \pi - \rho$ gives the formula obtained in [1] for $\lambda_{n,j}(\pi)$ and (9). Moreover, there is one-to-one correspondence between the eigenvalues (counting with multiplicities) and integers. Indeed, by (9) and Theorems 2-4 if $|n| > \max\{N, N(\pi, \rho)\}$ and $t \in [0, \pi]$ then the eigenvalues $\lambda_n(t)$ and $\lambda_{-n}(t)$ are simple and the number of the remaining eigenvalues of $L_t(q)$ is equal to $2 \max\{N, N(\pi, \rho)\} + 1$, where N and $N(\pi, \rho)$ are defined in Remark 1 and Theorem 4. Further, for simplicity of notation, $\max\{N, N(\pi, \rho)\}$ is denoted by N. Using the above notation, we see that the spectrum of $L_t(q)$ is

$$S(L_t) = \{\lambda_n(t) : n \in \mathbb{Z}\} = \{\lambda_{n,1}(t) : n = 1, 2, ...\} \cup \{\lambda_{n,2}(t) \ n = 0, 1, 2, ...\}.$$
 (102)

Since $\lambda_n(t)$ for |n| > N is a simple root of

$$F(\lambda) = 2\cos t,\tag{103}$$

where $F(\lambda)$ is the Hill's discriminant, we have

$$F(\lambda_n(t)) = 2\cos t, \ \frac{dF(\lambda_n(t))}{d\lambda} \neq 0, \ \frac{d\lambda_n(t)}{dt} = -(\frac{dF}{d\lambda})^{-1}2\sin t$$
(104)

for |n| > N, and $t \in [0, \pi]$. This implies that

$$\Gamma_n =: \{\lambda_n(t) : t \in [0,\pi]\}$$
(105)

is a simple (i.e. $\lambda_n : [0,\pi] \to \Gamma_n$ is injective) analytic arc with endpoints $\lambda_n(0)$ and $\lambda_n(\pi)$.

The eigenvalues of $L_{-t}(q)$ coincides with the eigenvalues of $L_t(q)$, because they are roots of equation (103) and $\cos(-t) = \cos t$. We define the eigenvalue $\lambda_n(-t)$ of $L_{-t}(q)$ by $\lambda_n(-t) = \lambda_n(t)$ for all $t \in (0, \pi)$. Then $\lambda_n(t)$ is an analytic function on $(-\pi, \pi]$.

Using Theorems 2-4 and taking into account Remark 2, we get

Theorem 5 Let $q \in W_1^p[0,1]$ and (3) holds with some $s \leq p$. If $q_n \sim q_{-n}$, $|q_n| > cn^{-s-1}$ and at least one of the following inequalities

$$\operatorname{Re} q_n q_{-n} \ge 0, \quad |\operatorname{Im} q_n q_{-n}| \ge \varepsilon \mid q_n q_{-n} \mid$$

hold for some c > 0 and $\varepsilon > 0$ and for n > N, where N is defined in Remark 2, then the eigenvalues $\lambda_n(t)$ of $L_t(q)$ for |n| > N and $t \in [0, \pi]$ are simple. The eigenvalues $\lambda_n(t)$ and the corresponding eigenfunctions $\Psi_{n,t}(x)$ satisfy the formulas (9) and the formulas obtained in Theorems 3 and 4.

3 Asymptotic Analysis of L(q)

Since the spectrum S(L(q)) of the operator L(q) is the union of the spectra $S(L_t(q))$ of the operators $L_t(q)$ for $t \in [0, \pi]$, it follows from (102) and (105) that

$$S(L(q)) = \bigcup_{n \in \mathbb{Z}} \Gamma_n.$$

By (104) and (105) the subset $\gamma =: \{\lambda_n(t) : t \in [\alpha, \beta]\}$, where $[\alpha, \beta] \subset [0, \pi]$, of Γ_n for |n| > N is a regular spectral arc of L(q) in sense of [9] (see Definition 2.4 of [9]). Following [24, 26, 9], we define the projection $P(\gamma)$ and the spectral singularities as follows

$$P(\gamma)f = \frac{1}{2\pi} \int_{\gamma} (\Phi_+(x,\lambda)F_-(\lambda,f) + \Phi_-(x,\lambda)F_+(\lambda,f)) \frac{\varphi(1,\lambda)}{p(\lambda)} d\lambda,$$
(106)

where $p(\lambda) = \sqrt{4 - F^2(\lambda)}$, $F(\lambda)$ is defined in (103),

$$\Phi_{\pm}(x,\lambda) =: \theta(x,\lambda) + (\varphi(1,\lambda))^{-1}(e^{\pm it} - \theta(1,\lambda))\varphi(x,\lambda)$$

is the Floquet solution and

$$F_{\pm}(\lambda, f) = \int_{\mathbb{R}} f(x) \Phi_{\pm}(x, \lambda) dx$$

The spectral singularities of the operator L(q) are the points of S(L(q)) in neighborhoods of which the projections $P(\gamma)$ of the operator L(q) are not uniformly bounded. In other words, we use the following definition.

Definition 1 We say that $\lambda \in S(L(q))$ is a spectral singularity of L(q) if for all sufficiently small $\varepsilon > 0$ there exists a sequence $\{\gamma_n\}$ of the regular spectral arcs $\gamma_n \subset \{z \in \mathbb{C} : | z - \lambda | < \varepsilon\}$ such that

$$\lim_{n \to \infty} \| P(\gamma_n) \| = \infty.$$
(107)

In the similar way we define the spectral singularity at infinity.

Definition 2 We say that the operator L(q) has a spectral singularity at infinity if there exists a sequence $\{\gamma_k\}$ of the regular spectral arcs such that $d(0, \gamma_k) \to \infty$ as $k \to \infty$ and (107) holds, where $d(0, \gamma_k)$ is the distance from the point (0, 0) to the arc γ_k .

To estimate the projections we use the following lemma of [16]:

Lemma 5.12 of [16] Let A' be in $L_{\infty}((0, 2\pi); B(L_2(0, 1)))$. Then for f in $L_2(-\infty, \infty)$ the limit in mean

$$Af = \lim_{N_i \to \infty} \frac{1}{2\pi} \sum_{-N_1 - N_3}^{N_2} \sum_{0}^{N_4} T_j^* \int_{0}^{2\pi} e^{it(j-k)} A'(t) T_k f dt$$
(108)

exists and defines a bounded operator in $L_2(-\infty,\infty)$ of norm $||A|| \le ||A'||_{\infty}$, where T_k is defined by $T_k(f(x)) = f(x+k)$ for $x \in [0,1)$, $T_k(f(x)) = 0$ for $x \ne [0,1)$ and

 $T_j^*(f(x)) = f(x-j)$ for $x \in [j, j+1), T_j^*(f(x)) = 0$ for $x \neq [j, j+1).$

Let $\{\chi_{n,t} : n \in \mathbb{Z}\}$ be biorthogonal to $\{\Psi_{n,t} : n \in \mathbb{Z}\}$ and $\Psi_{n,t}^*(x)$ be the normalized eigenfunction of $(L_t(q))^*$ corresponding to $\overline{\lambda_n(t)}$. Then

$$\chi_{n,t}(x)) = \frac{1}{\alpha_n(t)} \Psi_{n,t}^*(x), \ \alpha_n(t) = (\Psi_{n,t}(x), \Psi_{n,t}^*(x))_{(0,1)},$$
(109)

where $(.,.)_{(a,b)}$ denotes the inner product in $L_2(a,b)$. One can easily verify that

$$\Psi_{n,t}(x) = \frac{\Phi_+(x,\lambda_n(t))}{|\Phi_+(x,\lambda_n(t))|}, \ \chi_{n,t}(x) = \frac{1}{\overline{\alpha_n(t)}} \frac{\overline{\Phi_-(x,\lambda_n(t))}}{|\overline{\Phi_-(x,\lambda_n(t))}|},$$
(110)

$$\Psi_{n,t}(x+1) = e^{it}\Psi_{n,t}(x), \ \chi_{n,t}(x+1) = e^{it}\chi_{n,t}(x).$$
(111)

Now we are ready to prove the main results of this chapter:

Theorem 6 If all conditions of Theorem 5 hold, then

(a) The spectrum of the operator L(q) in a neighborhood of ∞ consist of the separated simple analytic arcs Γ_n for |n| > N with endpoints $\lambda_n(0)$ and $\lambda_n(\pi)$, where N is defined in Remark 2.

(b) The operator L(q) has at most finitely many spectral singularities.

(c) The projections $P(\gamma)$ of L(q) for all $\gamma \in \Gamma_n$ and |n| > N are uniformly bounded and the operator L(q) has no spectral singularity at infinity.

Proof. (a) Due to Remark 2 we need only to note that Γ_n for |n| > N are separated, that is, $\Gamma_n \cap \Gamma_k = \emptyset$ for $k \in \mathbb{Z} \setminus \{n\}$. This is true due to the following reason. The equality $\lambda_n(t) = \lambda_k(t)$ contradicts the simplicity of $\lambda_n(t)$. The equality $\lambda_n(t) = \lambda_k(t')$ for $t' \neq t$ and $t' \in [0, \pi]$ contradicts the first equality in (104).

(b) By Theorem 5 the equation $\frac{dF(\lambda)}{d\lambda} = 0$ has no zeros at Γ_n for |n| > N. Since $\frac{dF(\lambda)}{d\lambda}$ is an entire function, it has at most finite number of roots on the compact set $\bigcup_{|n| \le N} \Gamma_n$. Now the proof of (b) follows from the well-known fact that the spectral singularities of L(q) are contained in the set $\{\lambda : \frac{dF(\lambda)}{d\lambda} = 0, \lambda \in S(L(q))\}$ (see [9, 26]). (c) Changing the variable λ to the variable t in the integral in (106), using

$$d\lambda = -p(\lambda) \left(\frac{dF}{d\lambda}\right)^{-1} dt, \ \frac{dF(\lambda_n(t))}{d\lambda} = -\varphi(1,\lambda_n(t))(\Phi_+(x,\lambda_n(t)),\overline{\Phi_-(x,\lambda_n(t))})$$

and (110) by simple calculations we get

$$P(\gamma)f(x) = \frac{1}{2\pi} \int_{\delta} (f, \chi_{n,t})_{\mathbb{R}} \Psi_{n,t}(x) dt, \qquad (112)$$

where $\delta = \{t \in (-\pi, \pi] : \lambda_n(t) \in \gamma\}$. Let A'(t) be the operator defined by

$$A'(t)f = (f, \chi_{n,t})_{(0,1)}\Psi_{n,t}(x)$$
(113)

for $t \in \delta$ and A'(t) = 0 for $t \in (-\pi, \pi] \setminus \delta$. By (109) we have

$$\|A'(t)\| = |\alpha_n(t)|^{-1}, \ \forall t \in \delta,$$
(114)

where α_n is a continuous function and $\alpha_n(t) \neq 0$, since $\lambda_n(t)$ is a simple eigenvalue. Therefore $A' \in L_{\infty}((0, 2\pi); B(L_2(0, 1))).$

Let $f \in C_0$, where C_0 is the set of all compactly supported continuous functions, and A be the operator defined by (108). Then using Lemma 5.12 of [16] and (111)-(113) we get

$$A = \lim_{N_i \to \infty} \frac{1}{2\pi} \sum_{j=-N_1}^{N_2} \sum_{k=-N_3}^{N_4} T_j^* \int_0^{2\pi} e^{it(j-k)} A'(t) T_k f(x) dt =$$
$$\lim_{N_i \to \infty} \frac{1}{2\pi} \sum_{j=-N_1}^{N_2} \sum_{k=-N_3}^{N_4} T_j^* \int_0^{2\pi} e^{itj} (f(x+k)e^{-itk}, \chi_{n,t})_{(0,1)} \Psi_{n,t}(x) dt =$$
$$\lim_{N_i \to \infty} \frac{1}{2\pi} \sum_{j=-N_1}^{N_2} \int_0^{2\pi} (f, \chi_{n,t})_{\mathbb{R}} e^{itj} T_j^* \Psi_{n,t}(x) dt = P(\gamma) f(x).$$

Hence $Af = P(\gamma)f$ for all $f \in C_0$, where C_0 is dense in $L_2(-\infty,\infty)$. Moreover, A is bounded by Lemma 5.12 of [16] and $P(\gamma)$ is bounded, since $\gamma \subset \Gamma_n$ and Γ_n for |n| > N does not contain the spectral singularities. Therefore, we have $A = P(\gamma)$. Now Lemma 5.12 of [16]

with (114) implies that

$$\| P(\gamma) \| \leq \sup_{t \in \delta} | \alpha_n(t) |^{-1}$$

On the other hand, by (a), if γ_k is a regular spectral arcs such that $d(0, \gamma_k)$ is a sufficiently large number, then there exists n such that |n| > N and $\gamma_k \subset \Gamma_n$. Therefore, (c) follows from the following lemma.

Lemma 6 If all conditions of Theorem 5 hold, then there exists a constant d such that

$$|\alpha_n(t)|^{-1} < d \tag{115}$$

for all |n| > N and $t \in (-\pi, \pi]$.

Proof. For $t \in [\rho, \pi - \rho]$ inequality (115) follows from (9). Now we prove this for $t \in [0, \rho]$. The other cases are similar. Since the boundary condition (2) is self-adjoint we have $(L_t(q))^* = L_t(\overline{q})$. Moreover, the Fourier coefficients of \overline{q} has the form

$$(\overline{q}, e^{i2\pi nx}) = \overline{q_{-n}}.$$

Therefore, one can readily verify that if q satisfies the conditions of Theorem 5 then \overline{q} also satisfies these conditions. Thus all formulas and theorems obtained for L_t are true for L_t^* if we replace q_n with $\overline{q_{-n}}$. Hence, by formula (26), we have the following, uniform with respect to $t \in [0, \rho]$, formulas

$$\Psi_{n,j,t}^{*}(x) = u_{n,j}^{*}(t)e^{i(2\pi n+t)x} + v_{n,j}^{*}(t)e^{i(-2\pi n+t)x} + h_{n,j,t}^{*}(x),$$

$$(u_{n,j}^{*}(t))^{2} + (v_{n,j}^{*}(t))^{2} = 1 + O(n^{-1}), (h_{n,j,t}^{*}, e^{i(\pm 2\pi n+t)x}) = 0, \parallel h_{n,j,t}^{*} \parallel = O(n^{-1}).$$
(116)

Then

$$(\Psi_{n,j,t}(x),\Psi_{n,j,t}^*(x)) = u_{n,j}(t)\overline{u_{n,j}^*(t)} + v_{n,j}(t)\overline{v_{n,j}^*(t)} + O(n^{-1}).$$
(117)

By Lemma 5 we have $v_{n,1}^* \sim u_{n,2}^*(t) \sim 1$. Using this and (80) in (117) for j = 1, we get

$$(\Psi_{n,1,t}, \Psi_{n,1,t}^*) = v_{n,1}(t)\overline{v_{n,1}^*(t)} (1 + \frac{u_{n,1}(t)u_{n,1}^*(t)}{v_{n,1}(t)\overline{v_{n,1}^*(t)}}) + O(n^{-1}).$$
(118)

It follows from (84) and (85) that

$$\frac{u_{n,1}}{v_{n,1}} = \frac{(q_{2n} + B(t))}{(C(t) - 4\pi nt - \sqrt{D(t)})} = \frac{-C(t) + 4\pi nt - \sqrt{D(t)}}{(q_{-2n} + B'(t))}.$$
(119)

Then $\frac{u_{n,1}^*(t)}{v_{n,1}^*(t)}$ satisfies the formula obtained from (119) by replacing q_n with $\overline{q_{-n}}$. Hence, $\frac{\overline{u_{n,1}^*(t)}}{v_{n,1}^*(t)}$ satisfies the formula obtained from (119) by replacing q_n with q_{-n} . Thus, we have

$$\frac{u_{n,1}}{v_{n,1}}\frac{u_{n,1}^*(t)}{\overline{v_{n,1}^*(t)}} = \frac{-C(t) + 4\pi nt - \sqrt{D}}{(q_{-2n} + B^*(t))}\frac{(q_{-2n} + B^*(t))}{(C^*(t) - 4\pi nt - \sqrt{D^*(t)})},$$

where B^*, C^* and D^* are obtained from B, C and D by replacing q_n with q_{-n} . Since

$$(q_{-2n} + B'(t)) = q_{-2n}(1 + o(1)), \ (q_{-2n} + B^*(t)) = q_{-2n}(1 + o(1))$$

(see (46), (53)), the last equality can be written in the form

$$\frac{u_{n,1}}{v_{n,1}}\frac{\overline{u_{n,1}^*(t)}}{\overline{v_{n,1}^*(t)}} = \frac{-C(t) + 4\pi nt - \sqrt{D(t)}}{C^*(t) - 4\pi nt - \sqrt{D^*(t)}}(1 + o(1)).$$
(120)

Using (57) and (66) for L_t and L_t^* one can easily see that

$$|C(t)| + |\sqrt{D(t)}| + |C^{*}(t)| + |\sqrt{D^{*}(t)}| = O(f(n,t)),$$

where $f(n,t) = |4\pi nt| + |\sqrt{q_{2n}q_{-2n}}|$. Therefore, by (118), (120) and by $v_{n,1} \sim v_{n,1}^* \sim 1$, there exists a constant c_5 such that

$$\frac{1}{\mid (\Psi_{n,1,t}, \Psi_{n,1,t}^*) \mid} < c_5 \mid \frac{C^*(t) - 4\pi nt - \sqrt{D^*(t)}}{C^*(t) - \sqrt{D^*(t)} - C(t) - \sqrt{D(t)} + o(f(n,t))} \mid.$$
(121)

From (57) and (66) for L_t^* we get

$$|C^{*}(t) - 4\pi nt - \sqrt{D^{*}(t)}| < |9\pi nt| + 2 |\sqrt{q_{2n}q_{-2n}}| < 3f(n,t).$$
(122)

Similarly, by (57), (66) and (69) there exists a constant c_6 such that

$$|C^{*}(t) - \sqrt{D^{*}(t)} - C(t) - \sqrt{D(t)} + o(f(n,t))| > c_{6}f(n,t),$$
(123)

Thus, using (122) and (123) in (121) we get the proof of the lemma \blacksquare

It easily follows from this lemma the following result about the asymptotic spectrality of the operator L(q) defined as follows. Let $e(t, \gamma)$ be the spectral projection defined by contour integration of the resolvent of $L_t(q)$, where $\gamma \in R$ and R is the ring consisting of all sets which are the finite union of the half closed rectangles. In [17] it was proved Theorem 3.5 (for the differential operators of arbitrary order with periodic coefficients rather than for L(q)) which, for the operator L(q), can be written in the form:

L(q) is a spectral operator if and only if

$$\sup_{\gamma \in R} (\sup_{t \in (-\pi,\pi]} \| e(t,\gamma) \|) < \infty.$$

According to this theorem we give the following definition of the asymptotic spectrality.

Definition 3 The operator L(q) is said to be an asymptotically spectral operator if there exists a positive constant C such that

$$\sup_{\gamma \in R(C)} (\sup_{t \in (-\pi,\pi]} \| e(t,\gamma) \|) < \infty,$$
(124)

where R(C) is the ring consisting of all sets which are the finite union of the half closed rectangles lying in $\{\lambda \in \mathbb{C} : |\lambda| > C\}$.

Theorem 7 If all conditions of Theorem 5 hold, then the operator L(q) is an asymptotically spectral operator in sense of Definition 3.

Proof. Let C be a positive constant such that if $\lambda_n(t) \in \{\lambda \in \mathbb{C} : |\lambda| > C\}$, then |n| > N for all $t \in (-\pi, \pi]$, where N is defined in Remark 2. If $\gamma \in R(C)$, then γ contains in a finite number of the simple eigenvalues of $L_t(q)$. Thus, there exists a finite subset $J(t, \gamma)$ of $\{n \in \mathbb{Z} : |n| > N\}$ such that the eigenvalue $\lambda_k(t)$ lies in γ if and only if $k \in J(t, \gamma)$. It is

well-known that these eigenvalues are the simple poles of the Green function of $L_t(q)$ and the projection $e(t, \gamma)$ has the form

$$e(t,\gamma)f = \sum_{n \in J(t,\gamma)} \frac{1}{\alpha_n(t)} (f, \Psi_{n,t}^*) \Psi_{n,t}$$

Then by (26) $e(t,\gamma)f$ for $t \in [0,\rho]$ is the sum of $e^+(t,\gamma)f$, $e^-(t,\gamma)f$ and $e^h(t,\gamma)f$, where

$$e^{+}(t,\gamma)f = \sum_{n \in J(t,\gamma)} \frac{1}{\alpha_{n}(t)} (f, \Psi_{n,t}^{*}) u_{n}(t) e^{i(2\pi n + t)x},$$

$$e^{-}(t,\gamma)f = \sum_{n \in J(t,\gamma)} \frac{1}{\alpha_{n}(t)} (f, \Psi_{n,t}^{*}) v_{n}(t) e^{i(-2\pi n + t)x},$$

$$e^{h}(t,\gamma)f = \sum_{n \in J(t,\gamma)} \frac{1}{\alpha_{n}(t)} (f, \Psi_{n,t}^{*}) h_{n,t}, \quad || \ h_{n,t} \mid| = O(\frac{\ln |n|}{|n|})$$

It follows from (27), (28) and Lemma 6 that there exists a constant c_7 such that

$$\| e^{\pm}(t,\gamma)f \|^{2} < 2d^{2} \sum_{n:|n|>N} | (f,\Psi_{n,t}^{*}) |^{2}, \qquad (125)$$

$$\| e^{h}(t,\gamma)f \| < c_{7}d \sum_{n:|n|>N} | (f, \Psi_{n,t}^{*}) | \frac{\ln |n|}{|n|}$$
(126)

for all $t \in [0, \rho]$ and $\gamma \in R(C)$.

Now suppose that || f || = 1. By (116) there exists a constant c_8 such that

$$\sum_{n:|n|>N} |(f, \Psi_{n,t}^*)|^2 \le c_8.$$
(127)

This inequality with (125) gives

$$\| e^{\pm}(t,\gamma)f \|^2 < 2d^2c_8.$$

Thus, in (126), first using the Schwarz inequality for l_2 and then taking into account (127) we conclude that there exists a positive constant c_9 such that $|| e^h(t, \gamma) f || < c_9$ and

$$\parallel e(t,\gamma) \parallel < c_9 \tag{128}$$

for all $t \in [0, \rho]$ and $\gamma \in R(C)$. In the same way, one can prove inequality (128) for all $t \in (-\pi, \pi]$ and $\gamma \in R(C)$, that is, the theorem is proved.

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