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## FERROELECTRICITY OF THE NUCLEIC ACIDS

### AND SELECTED

NUCLEOTIDES AND POLYRIBONUCLEOTIDES

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A THESIS

Presented to

The Faculty of the Graduate Division

by

Richard Anthony Lorey

In Partial Fulfillment

of the Requirements for the Degree Doctor of Philosophy in the School of Physics

> Georgia Institute of Technology August, 1969

## FERROELECTRICITY OF THE NUCLEIC ACIDS

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Approved: Chailfman Date approved by Chairman <u>8-22-69</u>

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#### SUMMARY

The nucleic acids are known to be involved in the metabolism, reproduction and growth of living systems. Recent experiments have implied that ribonucleic acid (RNA) is intimately associated with the memory process. The present work was initiated as a search for a physical property of RNA that could relate RNA to this process.

A physical property which can readily account for information storage, and hence memory, is hysteresis. Ferromagnetic hysteresis is the mechanism responsible for information storage in such devices as tape recorders and computers. Since ferromagnetic materials exist in only trace amounts in neural tissue it is unreasonable to assume this mechanism could relate RNA and the memory process. Another type of hysteresis which can lead to an explanation of information storage is ferroelectric hysteresis. Hence, it is of interest to determine if RNA is ferroelectric.

A previous report indicated that deoxyribonucleic acid (DNA) may be ferroelectric. Therefore, for the present work it was considered desirable to study the dielectric properties and ferroelectric character of RNA and DNA, the several ribonucleotides and two synthetically produced polyribonucleotides.

The equipment used in this research consisted of apparatus used for hysteresis observation and apparatus used for measuring dielectric constants as functions of temperature. The circuit of Sawyer and Tower was used to display the ferroelectric hysteresis loops on an oscilloscope. Heating the samples permitted observation over the temperature range of 25°C to approximately 100°C. The dielectric constant measurements consisted of measuring capacitances dynamically by plotting capacitances on an X-Y recorder as a function of temperature.

The salts of the various acids were mixed with water and allowed to harden into relatively flat specimens. Electrodes were then painted on these samples and leads were attached to the electrodes. The material thus became the dielectric medium inside parallel plate capacitors. These capacitors were used in both the hysteresis observations and dielectric constant measurements.

The dielectric constants of the nucleic acids and other related materials exhibit the characteristic behavior of many ferroelectric materials as the temperature nears their Curie points. No high temperature Curie point was observed for any of the materials studied here. A low temperature irregularity in the dielectric constant of RNA may be a ferroelectric transition or simply a change in the value of the coercive field strength. No saturation of the hysteresis loops was observed for any of the materials studied here. Indications are that the nucleic acids and the associated biological molecules are ferroelectric in a

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temperature region between approximately  $30^{\circ}$ C and  $80^{\circ}$ C. The existence of ferroelectricity in the nucleic acids appears to be independent of the type of pentose sugar present, the types of bases present, the sequence of bases and the length of the molecule. Finally, it is suggested that the origin of the ferroelectricity of these materials may be in the phosphate group of the molecules.

Recommendations are made for ferroelectric, pyroelectric and crystallographic studies on the nucleic acids and several other related chemicals.

#### CHAPTER I

#### INTRODUCTION

Investigations of the chemical and physical properties of the nucleic acids have largely been concerned with the relationships of these properties to the role of the nucleic acids in the life processes. For example, it is known that the nucleic acids are of prime importance in controlling metabolism, reproduction and growth of living systems (1). Recently however, experiments have implied that ribonucleic acid (RNA) is intimately associated with the memory process of platyhelminths and mammals (2,3,4). It is therefore of interest to determine any relationship between the various properties of RNA and the memory process.

Any scientific correlation between the properties of RNA and the memory process must ultimately be based on chemical or physical principles. Furthermore, this correlation must take into account the distinction between long and short term memory. The two types of memory are defined here as follows: (1) long term memory is the reception and storage of information and the availability for recall long after the reception of the information, and (2) short term memory is the reception and storage of information, the immediate availability for recall after reception and the

possible short term retention of the information. As a mechanism to explain long term memory, it is quite reasonable to assume a chemical coding of the RNA. A chemical coding of genetic information is already known to exist in RNA and in deoxyribonucleic acid (DNA). The use of this mechanism to explain short term memory is not feasible since the chemical reaction of macromolecular synthesis, required by chemical coding, is too slow to account for the immediate recall required by the definition of short term memory. One is thus led to search for some other information storage mechanism to explain short term memory. From a physical standpoint an information storage mechanism based on hysteresis permits a ready explanation of the immediate recall of short term memory. Two common types of materials exhibiting hysteresis are ferromagnetic and ferroelectric materials. Ferromagnetic hysteresis accounts for the storage and recall of information of such mechanical devices as computers and tape recorders. However, an explanation based on a ferromagnetic hysteresis mechanism is unacceptable since ferromagnetic materials exist in only trace amounts in neural tissue. Hence, it is of interest to determine the existence of ferroelectricity in RNA.

#### History of Ferroelectricity

It is possible to divide all crystal types into 32 crystal classes, according to their rotational symmetries. Of these 32 crystal classes, or point groups as they are 2

often called, 11 have a center of symmetry (inversion center) and are called centrosymmetric. The remaining 21 classes lack an inversion center and are called noncentrosymmetric.

Crystals in 20 of the 21 noncentrosymmetric crystal classes exhibit the phenomenon of piezoelectricity (5). This is the ability of certain materials to polarize, producing an electric field, when subjected to a compression or tension. Conversely, when placed in an electric field, these materials expand or contract. It is possible then to predict piezoelectricity on the basis of crystallographic information alone, although the effect is not always large enough to detect.

Of those crystal classes which exhibit piezoelectricity, ten classes possess a polar axis (5). A polar axis is an axis along which exists a possible permanent (spontaneous) polarization. This spontaneous polarization is usually masked by surface charge. Heating these crystals will permit observation of the polarization by producing a change in the polarization. These crystals are thus said to be pyroelectric. The existence of pyroelectricity, like piezoelectricity, may be deduced once the point group symmetry of a crystal is known.

If the direction of the spontaneous polarization of a pyroelectric material can be reversed by application of an electric field without causing dielectric breakdown, the crystal is said to be ferroelectric. Ferroelectricity is a dielectric property and can be determined only from dielectric measurements, not from crystallographic measurements. The name is derived strictly from analogy between the dielectric behavior of these crystals and the magnetic properties of ferromagnetic materials. For example, it is found that ferroelectric materials exhibit a domain structure analogous to the magnetic domain structure of ferromagnetic crystals and exhibit electric hysteresis similar to ferromagnetic hysteresis. Furthermore, quite a few ferroelectric materials exhibit a transition temperature beyond which they are no longer polar and where their dielectric constants obey, again in analogy to ferromagnetism, a Curie-Weiss law

$$\epsilon = C/(T - T_0), \qquad (1)$$

where  $\epsilon$  is the dielectric constant, C is the Curie constant, T is the temperature and T<sub>o</sub> is a characteristic temperature (6).

Ferroelectricity was discovered in Rochelle salt by Valasek in 1921 (7). For some time thereafter the effect was not observed in any material except Rochelle salt. It is now known however, that ferroelectricity exists in many materials and an impressive list of known ferroelectric materials could be cited. It is possible to divide some well studied ferroelectric materials into four families: (1) the ferroelectric tartrate family represented by Rochelle salt, (2) the alkali metal dihydrogen phosphates and arsenates of which  $KH_2PO_4$  is a typical representative, (3) the oxygen-octahedra family which includes the subgroup with the perovskite (CaTiO<sub>3</sub>) structure, represented by BaTiO<sub>3</sub> and (4) the family represented by guanidine aluminum sulfate hexahydrate.

Several theories are known which attempt to explain the ferroelectricity, in large measure, of that class of ferroelectric materials represented by barium titanate. Barium titanate has a cubic structure in its nonpolar phase, which makes it a convenient structure for relatively simple theoretical treatment. Two of the theories which have been successful are those of Devonshire (1949) and Cochran (1960) (8,9,10). Devonshire's work is a thermodynamic, phenomenological theory, while Cochran's theory is based on lattice dynamics.

#### Physical Properties of Ferroelectric Materials

As mentioned above, ferroelectric materials exhibit a domain structure analogous to that in ferromagnetic materials. Since polarization is defined as the electric dipole moment per unit volume it is possible to associate with each domain a region in which the electric dipoles are all aligned. However, the association of these dipoles with specific pairs of opposite charges may not always be possible (see Appendix A).

The hysteresis exhibited by a ferroelectric material can easily be understood on the basis of domain structure. Figure 1 represents a hysteresis loop for a typical single crystal ferroelectric material. The exact shape of the

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hysteresis loop obtained for a given specimen will depend upon a number of factors. If the specimen is not a single crystal, one obtains a superposition of loops due to crystallites. A number of other factors which may affect the shape of the hysteresis loop are the temperature, humidity, strength and frequency of the switching field, conductivity of the specimen, physical dimensions and prior thermal and electrical history (7,11). Referring again to Figure 1, at point a the specimen has no net polarization, i.e., all the domains are randomly oriented. As the electric field is increased from zero (a to b) the polarization increases rapidly, this means that there is a net alignment of domains in the direction of the applied field. Upon further increase of the electric field (b to c) saturation is reached (c to d). At this point no new alignment takes place. The intersection of the polarization axis and the extrapolated portion of the linear part of this curve is referred to as the spontaneous polarization  $P_s$ . As the electric field is reduced to zero, a few of the domains aligned with the applied field switch back to random orientations and consequently at zero field a residual or remnant polarization Pr remains. As the electric field is increased in the opposite direction, switching of domains occurs until an equal number of domains are aligned in both directions and there is no net polarization. The value of the electric field at this point is called the coercive field strength  $E_c$ . This last value is

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not a well-defined quantity but rather depends on temperature, measuring frequency and waveform of the applied voltage (12). Upon further increase of the electric field in the negative direction, saturation is again reached and the remaining portion of the loop is analogous to that previously discussed. As already noted many ferroelectric materials exhibit a temperature at which the hysteresis loop disappears, in which case the loop degenerates to a straight line, i.e., the polarization is then a linear function of the applied field.

It is possible to define the dielectric constant or permittivity  $\epsilon$  of a ferroelectric material when in the polar phase in several ways (7). The susceptibility  $\chi$  is defined by

$$\chi = \mathbf{P}/\mathbf{E} \tag{2}$$

and is related to  $\epsilon$ , in the cgs system, by

$$\epsilon = 1 + 4 \pi \chi , \qquad (3)$$

and in the MKS system, by

$$\epsilon = \chi + 1. \tag{4}$$

Here, four distinct susceptibilities, and therefore four permittivities, may be defined: (1) the overall susceptibility P/E with values of P and E from point c of Figure 1, (2) the differential susceptibility  $\partial P/\partial E$  at any point, (3) the saturation susceptibility, which is the slope at point c of Figure 1, and (4) the initial or "small signal" susceptibility, which is the slope  $\partial P/\partial E$  at zero external field. In this thesis the dielectric constant will be taken to mean the initial or "small signal" dielectric constant.

The dielectric constants of ferroelectric materials are normally quite high. More accurately, one or more of the components of the dielectric constant tensor may be quite high. As the temperature is raised (or lowered in the case of a lower transition point) the dielectric constant increases, and at or near the transition point (or Curie point) may become exceedingly large. Above the Curie point in the nonferroelectric phase, the dielectric constant obeys a Curie-Weiss law (Equation 1).

### Chemical Structure of the Nucleic Acids

The nucleic acids are composed of pentose sugars, phosphoric acid, purines and pyrimidines. The pentose sugars involved are ribose in RNA and deoxyribose in DNA (1,13,14). The purines and pyrimidines are commonly referred to as bases. The purines adenine and guanine, and the pyrimidine cytosine, occur in both RNA and DNA. The pyrimidine uracil occurs only in RNA and the pyrimidine thymine (5-methyluracil) occurs only in DNA. The chemical structures of these purines and pyrimidines are illustrated in Figure 2. Figure 3 shows the chemical structure of the pentose sugars and phosphoric acid.

A purine or a pyrimidine base is bonded to a pentose



Pyrimidines







Phosphoric Acid



Ribose



Deoxyribose

Figure 3. Phosphoric Acid, Ribose Sugar and Deoxyribose Sugar

sugar and the combination is known as a nucleoside. Figure 4 illustrates the nucleoside of adenine, known as adenosine or adenine riboside, and of cytosine, known as cytidine or cytosine riboside. If the nucleoside is formed with deoxyribose sugar the combination is referred to as a deoxyriboside. The chemical structures of ribosides and deoxyribosides are identical insofar as bond positions are concerned and hence the illustration does not show any deoxyriboside.

The phosphoric ester of a nucleoside is known as a nucleotide. The phosphoric acid is attached to the pentose sugar as illustrated in Figure 5, where again only the case for ribose sugar is shown. The nucleotides illustrated are adenylic acid or adenosine-5'-monophosphate and cytidylic acid or cytidine-5'-monophosphate. The 5' refers to the particular sugar carbon atom bonded to the phosphate group. Other nucleotides are possible with the attachment of the phosphate group to a different position on the pentose sugar but these will not be considered here.

The nucleic acids are formed by the bonding of nucleotides into long chains. The nucleotides formed with adenine, guanine, cytosine and uracil make up RNA with bonding occurring between the phosphate group of one nucleotide and the pentose sugar of another nucleotide as shown in Figure 6. In DNA the uracil nucleotide is replaced by the thymine nucleotide. Furthermore, DNA consists of two of these nucleotide chains linked by hydrogen bonds between the purines





Adenylic Acid (adenosine-5'-monophosphate)



Cytidylic Acid (cytosine-5'-monophosphate)

# Figure 5. Adenylic Acid and Cytidylic Acid



and pyrimidines. In this connection adenine bonds with thymine, and guanine bonds with cytosine as shown in Figure 7.

It is possible to synthetically produce polyribonucleotides with only one particular type of base. The names given to the adenine and cytosine polyribonucleotides are polyadenylic acid and polycytidylic acid respectively.

The chemicals with which this thesis is concerned are listed in Table 1. As indicated in this table the salts of the various acids were used rather than the actual acids themselves. All these chemicals were obtained commercially.

#### Crystal Nature of the Chemicals Used in this Research

X-ray pictures (Laue photographs) were obtained of the four macromolecular materials used here, and of guanylic acid disodium and cytidylic acid disodium. Attempts to obtain Laue photographs of adenylic acid sodium and thymidylic acid sodium were unsuccessful. The necessarily long exposure times, of the order of seven or eight hours, permitted the absorbtion of water and subsequent dissolution of the latter two materials. The X-ray information indicated that the samples of guanylic acid disodium and cytidylic acid disodium were polycrystalline in nature. The photographs of the cytidylic acid disodium sample most clearly showed this crystal nature. It will be assumed by analogy that the adenylic acid sodium and thymidylic acid sodium samples were also polycrystalline. This analogy is reasonable because adenine and guanine are both purines and thymine and cytosine are both pyrimidines.

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Figure 7. Deoxyribose Nucleic Acid

Table 1. The Chemicals Used in this Research

- 1. Adenine
- 2. Adenosine
- 3. Adenylic Acid (Adenosine-5'-Monophosphate) Sodium
- 4. Cytidylic Acid (Cytidine-5'-Monophosphate) Disodium
- 5. Guanylic Acid (Guanosine-5'-Monophosphate) Disodium
- 6. Thymidylic Acid (Thymidine-5'-Monophosphate) Sodium
- 7. Uridylic Acid (Uridine-5'-Monophosphate) Disodium
- 8. Sodium Deoxyribonucleate
- 9. Sodium Ribonucleate
- 10. Polyadenylic Acid Potassium
- 11. Polycytidylic Acid Potassium

The photographs of the four macromolecular materials did not indicate crystalline structures. However, it is quite probable there exist some degree of intramolecular order with intermolecular disorder. Fibers of these materials drawn from solution with the macromolecules lying generally along the fiber axis are usually employed in X-ray work (15,16,17). The bulk form used here could be expected to yield little information concerning the crystal nature of these materials. All X-ray work here was done at room temperature.

In connection with the macromolecular structure it should be noted that DNA can exist both in a crystalline and in a paracrystalline form (15,16,18,19). The exact structure is determined by the relative humidity and the most crystalline form occurs at approximately 70 per cent relative humidity (18). At higher humidities the paracrystalline form results and at lower humidities the loss of water distorts the structure, causes irregularities, but leaves the basic structure intact (15). Figure 7 illustrates the chemical structure of DNA. The basic physical structure of DNA is a double helix with the sugar-phosphate chains forming the backbones and the hydrogen bonded purines and pyrimidines connecting these chains (1,13,14,18). The crystalline form of DNA has a face-centered monoclinic lattice, has a repeat distance along the helix axis of 28 Å, contains approximately 30 per cent water by weight and the bases are tilted about  $25^{\circ}$  from the helix axis (18,19). The paracrystalline form

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of DNA contains more water, has a repeat distance of 34 A and has all the bases parallel to each other and perpendicular to the helix axis but with random rotation about their own molecular axes (18,19). The structural changes which occur with changes in humidity are not part of the helixcoil transitions associated with heat denaturation which occurs in the range of  $70^{\circ}C$  to  $80^{\circ}C$  (1,20). In contrast to DNA and most of the other polyribonucleotides, the structure of polyadenylic acid is independent of humidity (17). Polvadenylic acid may form a double helix with the sugar-phosphate backbone chains held together by purine-purine hydrogen bonds and purine-phosphate hydrogen bonds (17). The bases are inclined about  $10^{\circ}$  from the perpendicular to the helix axis and the resulting structure is quite compact. Denaturation, helix-coil transition, occurs over a temperature range centered at 90°C (17).

#### Purpose of this Research

This work was stimulated by the search for a physical property of RNA which could relate RNA to the memory process. It is interesting to note that the results of this investigation are particularly pertinent to the search for possible mechanisms for information storage in neural tissue. However, the purpose of this research is the experimental investigation of the dielectric properties and ferroelectric character of the nucleic acids, RNA and DNA, and several other related compounds.

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#### Review of the Literature

A report on the dielectric properties and ferroelectric character of DNA was made in 1960 (21). The material used was the lyophilized sodium salt of DNA which was pressed into a wafer. Neither the lyophilized nor the wafer form of sodium DNA is used in this research. The only other known experimental work reported concerning the ferroelectricity of the nucleic acids was done on the sodium salt of RNA in 1968 and coauthored by the present investigator (22).

#### CHAPTER II

#### INSTRUMENTATION AND EQUIPMENT

The equipment used in this investigation consisted mainly of that designed to observe the dielectric hysteresis and that designed to measure dielectric constants. All of the measurement techniques used in this work are of standard nature and none of the equipment has been uniquely designed for this experiment.

The "classical" method of Sawyer-Tower was used to observe dielectric hysteresis (23). The actual circuit used is shown schematically in Figure 8, where  $C_0$  is a standard capacitor and  $C_x$  is the sample capacitor. Figure 9 shows a typical hysteresis loop obtained using this circuit with a sample crystal of barium titanate. The picture was taken with a polaroid camera and then reproduced from the photograph.

It is of value here to consider the Sawyer-Tower circuit in more detail. The capacitors  $C_0$  and  $C_X$  are related by the inequality

$$C_0 \gg C_X$$
. (5)

With this condition it is readily seen that the horizontal





Standard Potential Capacitor Divider

Potential Divider


input voltage is given by

$$\mathbf{V} = \mathbf{E}_{\mathbf{x}} + \mathbf{E}_{\mathbf{0}} \cong \mathbf{E}_{\mathbf{x}} , \qquad (6)$$

where V is the input voltage,  $E_x$  is the voltage across the sample and  $E_0$  is the voltage across  $C_0$ . This approximation follows from consideration of the equation relating  $E_0$  to the charge on either capacitor,

$$\mathbf{E}_{\mathbf{O}} = \mathbf{Q} / \mathbf{C}_{\mathbf{O}} \quad (7)$$

Since the applied electric field is equal to the sample voltage divided by the thickness of the sample which is a constant, the horizontal displacement on the oscilloscope is proportional to the electric field. Furthermore, since the polarization is equal to the charge on the sample divided by the area of the sample which is a constant, and the charge is related to the voltage as shown in Equation 7, the vertical displacement on the oscilloscope is proportional to the polarization.

If the sample  $C_x$  is replaced by a linear element such as a resistor, capacitor or a parallel combination of these two, the oscilloscope will display an ellipse or its degenerate case, a straight line. It may be shown (see Appendix B) that the coordinates of the oscilloscope trace are related by an equation of the following type

$$Ax^{2} + Bxy + Cy^{2} + Dx + Ey + F = 0$$
, (8)

where the constants are functions of the capacitance, resistance and frequency of the circuit and where not all constants necessarily differ from zero. This equation represents an ellipse or its degenerate case provided that the inequality

$$B^2 - 4AC < 0 \tag{9}$$

is satisfied. Thus, ideal linear elements will yield either an ellipse, a circle or a straight line.

Considering again the nonlinear element, in most cases the ferroelectric sample has a finite conductivity. The resulting hysteresis loop may be regarded as consisting of three contributions as seen in the equation

$$Q(V) = P(V) + C \cdot V + (1/R) \int V dt$$
 (10)

The first term in Equation 10 represents the spontaneous polarization contribution, the second term represents the capacitance contribution and third term represents the conductivity contribution (11). If the input voltage is sinusoidal and given by

$$\mathbf{V} = \mathbf{V}_{\mathbf{O}} \mathbf{sin} \ \omega \mathbf{t} , \tag{11}$$

Q may be expressed as

$$Q = P(V) + V_{o}(C \cdot \sin \omega t - 1 / \omega R \cdot \cos \omega t) , \qquad (12)$$

or

$$Q = P(V) + C/\cos\phi \cdot V_{o}\sin(\omega t - \phi) , \qquad (13)$$

where  $\phi$  is related to C by the equation

$$\tan\phi = 1 / \omega RC . \tag{14}$$

Thus the horizontal deflection voltage phase is shifted and the resulting hysteresis loop shows rounded ends.

The circuit used to measure the dielectric constant as a function of temperature is shown schematically in Figure 10. The thermocouple was made of chromel alumel and was used to drive the X-axis of an X-Y recorder. Figure 11 shows the calibration curve for the thermocouple using solid  $CO_2$ , ice and boiling water as reference points. The signal across  $C_0$  is amplified, rectified and made to drive the Y-axis of the X-Y recorder. The curve obtained on the X-Y recorder is then a measure of the capacitance (dielectric constant) as a function of temperature. Figure 12 shows a plot of peak input voltage versus rectified voltage showing a linear relation in the range utilized. To obtain a



# Figure 10. Capacitance Measuring Circuit





Input Voltage (volts)



capacitance scale on the X-Y recorder, standard capacitances were measured on a General Radio impedance bridge and then placed in the circuit in the position of  $C_x$ . This procedure was carried out with each sample measured.

The oven used to heat the samples was made by wrapping 18 loops of nichrome wire around a 100 ml pyrex beaker and surrounding this with aluminum oxide refractory brick. Access to the oven was provided by a teflon door through which leads to the sample and thermocouple were placed. A port was provided for observation of the sample with a microscope. Figure 13 shows a schematic diagram of the oven.

Aluminum cups were made on a lathe to the dimensions shown in Figure 13. These cups were used in the preparation of samples.

Other miscellaneous equipment used included a microscope used to measure the thickness of samples, a refrigerator used in drying samples and a hydraulic press used to make wafers of some of the materials.







Top View



# Figure 13. Sample Cups and Sample Oven

#### CHAPTER III

#### EXPERIMENTAL PROCEDURE

The experimental procedure in these experiments consisted of three phases; sample preparation, hysteresis observation and dielectric constant measurement.

# Sample Preparation

Two procedures were used to obtain specimens suitable for measurements. The first method involved dissolution of the sodium or potassium salt of the various acids in distilled water. The amount of each chemical used, with the corresponding amount of water, is listed in Appendix C. These do not necessarily correspond to saturated solutions but rather to amounts of solute and solvent which would readily result in dissolution. These solutions were then poured into the aluminum cups described in Chapter II. The cups were then placed in a refrigerator at approximately  $5^{\circ}C$ , to provide slow evaporation of the water. The evaporation of the water yielded generally flat samples whose physical appearance, size and shape are described briefly in Appendix C. In the case of cytidylic acid disodium it was necessary to allow evaporation to proceed at room temperature because the material obtained from the refrigerator redissolved when brought to room temperature. This dissolution procedure did

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not yield usable samples of adenine, adenosine or uridylic acid disodium. To obtain samples of these materials a second procedure was used. Wafers were made of the powdered material using a hydraulic press. The amount of each material used and the pressure applied are listed in Appendix C. For comparison purposes, wafers were also made of adenylic acid sodium.

Electrodes were painted on the samples using DuPont Silver Preparation-Electronic Grade 4929. Leads 0.010 inch in diameter were placed on the silver electrodes and became attached upon drying and hardening of the electrodes. The samples thus became the dielectric medium inside parallel plate capacitors. To insure complete hardening of the electrode material, the coated samples were allowed to remain in the refrigerator at approximately  $5^{\circ}C$  for 24 hours.

The thickness of each sample was measured using a stereomicroscope equipped with a reticle whose smallest division is 0.001 inch. The area of the electrodes were determined from the stencils used for painting the electrodes. Electrodes were painted on glass slides using the stencils and these areas were determined by using the stereomicroscope reticle.

In both the hysteresis observation and the dielectric constant measurement, the sample was placed in the oven described in Chapter II. Figure 14 shows the heating curves for this oven for a few settings of the Variac oven control.



Figure 14. Oven Heating Curves

The rate at which the samples were heated appears in the compilation in Appendix D.

# Hysteresis Observation

Attempts to observe hysteresis were made on all materials, some of which were in the dissolved and hardened form and some in wafer form. However, no hysteresis loops were observed for any of the wafers. The samples were placed in the Sawyer-Tower circuit and the loops observed on an oscilloscope.

In order not to cause excessive heating of the sample due to the applied alternating voltage, only small signals were used. In this way the heating of the sample was due almost entirely to the oven. In order to test whether the thermocouple reading represented the actual sample temperature, the following procedure was utilized. Several layers of the sodium RNA solution were applied to one junction of a thermocouple and allowed to harden. This junction was placed in the oven adjacent to another thermocouple and the oven heated at twice the fastest rate used in the experiment. Over the temperature range used, room temperature to approximately  $100^{\circ}C$ , the lag was never greater than three degrees.

The observed hysteresis loops were photographed with a Polaroid camera and subsequently were reproduced from the photographs.

# Dielectric Constant Measurement

The capacitances of each sample was measured as a

function of temperature using the apparatus described in Chapter II. Small signals were used to measure the capacitances and their values are listed in the compilations in Appendix D.

In all cases it was assumed that the samples were ideal parallel plate capacitors. The ratios of the thicknesses of the samples to the diameters of the circular electrodes are listed in Appendix D. With this assumption, computation of the dielectric constant from the value of the capacitance is quite straightforward. Tables and Graphs of these measurements are also in Appendix D.

# CHAPTER IV

#### RESULTS

#### Precision

It is of interest to examine the errors associated with the measurements in this research prior to discussing the results; this will hopefully provide a better understanding of the results and their interpretation. The quantities which are involved in the measurement of the dielectric constant that contribute to the error are: (1) the temperature, (2) the capacitance, (3) the sample electrode area, and (4) the sample thickness.

Figure 11, which is the calibration curve for the thermocouple used here, shows that the error in the temperature is less than three degrees between 25°C and 100°C. Furthermore, the greatest error occurs at the high temperature end of this region. As explained in Chapter III, the sample temperature lagged the oven temperature less than three degrees when the oven was heated at twice the fastest rate used during the actual experiment. Since the heating rate decreases during heating, the temperature lag is greatest at the lower portion of the above temperature scale. Hence, it is reasonable to assume that the temperature error is less than four or five degrees over the entire temperature scale utilized. The values of the capacitances obtained from the impedance bridge and used for calibration of the X-Y recorder are accurate to within one percent. However, a large error results due to the method used for measuring the sample capacitance. The finite width of the recorder pen point and the resulting trace contribute significantly to this error. When measuring small capacitances the trace width may represent a large percentage of the capacitance scale. A more sensitive scale might decrease this percentage provided any resulting increase in noise level did not broaden the trace sufficiently to negate the advantage of any scale change. Thus some optimum scale must be established in order to minimize the error when measuring small capacitances. Even when measuring larger capacitances, the error introduced by the trace width usually can not be neglected.

The electrode area was determined by the method described in Chapter III. Sufficient numbers of electrodes of each size used were measured in order to determine the reproducibility and spread in the values of the areas. This procedure was used to determine the error due to area in each sample.

The thickness of each sample was measured a minimum of four times by the method described in Chapter III. The average value of these measurements was taken to be the sample thickness. Any sample whose faces were nonparallel by two or more degrees was not used for measurement.

Using the parallel plate capacitance approximation, the dielectric constant  $\epsilon$  is obtained by using the equation

$$\epsilon = Cd / \epsilon_0 A$$
, (15)

where C is the sample capacitance, d is the sample thickness,  $\epsilon_0$  is the permittivity of free space and A is the electrode area. Thus the error in the dielectric constant is just the sum of the errors of the sample capacitance, thickness and area. This cumulative error is listed along with its associated dielectric constant in Appendix D.

Additionally, the water content of the samples leads to an error in the dielectric constant. This error becomes appreciable at higher temperatures where presumably the water molecules are more mobile and are able to contribute to the sample conductivity. Higher temperature is assumed here to be above  $70^{\circ}$ C or  $80^{\circ}$ C. Since an ideal conductor has an infinite dielectric constant, one would expect the exceedingly large dielectric constants that occur at high temperatures. The values of these dielectric constants are listed in Appendix D. In particular it is to be noted that the dielectric constant for thymidylic acid sodium is quite large even at room temperature. The water concentration in this chemical may have been extremely high and thus contributed to the conductivity at lower temperatures. Therefore, particular care must be exercised when evaluating these high temperature

constants.

Finally, expansion of the samples at the high temperatures defined here may lead to an appreciable error. This expansion is generally not uniform since the electrodes add local stresses to the surfaces and the departure of the shape of the samples from parallel plate capacitors may be significant.

# Interpretation

The interpretation of the hysteresis loops and the dielectric constants for each chemical is discussed separately. The hysteresis loops and at least one dielectric constant curve for each material are included in this chapter. These serve to illustrate many of the items to be discussed and the reader may refer to Appendix D for the remaining dielectric constant curves. All hysteresis loops shown were obtained at a frequency of 60 cps and unless otherwise stated the dielectric constants were obtained at a frequency of 1000 cps. On the hysteresis loop illustrations D is the electrode diameter in centimeters, d is the sample thickness in centimeters, HC is the rate of heating and refers to Figure 14 and the voltage scales shown are uncalibrated voltages from the measuring oscilloscope.

<u>Sodium RNA.</u> The most striking feature of the dielectric constant versus temperature curve for sodium RNA, shown in Figure 15, is the irregularity which occurs at approximately 53°C. Reference to Figures 37, 38, 39 and 40 in Appendix D shows



Figure 15. Dielectric Constant of Sodium RNA

that although the shape of this irregularity is not always the same, it does occur within a five degree temperature range between  $52^{\circ}C$  and  $57^{\circ}C$ . Figures 41 and 42 in Appendix D show respectively the dielectric constant versus temperature at 100 cps and at 500 cps. It is clear that the irregularity is decreased, less pronounced, at the lower frequencies. This set of seven curves was obtained from two separate sample preparations with measurements made at different times and with different ambient conditions. Figures 15, 37 and 38 refer to samples from one preparation, and Figures 39, 40, 41 and 42 refer to samples from another preparation. For later reference the dielectric constant curve may be divided into three temperature regions: (1) up to  $45^{\circ}C$  or  $50^{\circ}C$ , (2) from  $45^{\circ}C$  or  $50^{\circ}C$  to  $60^{\circ}C$ , and (3) above  $60^{\circ}C$ .

Figure 16 illustrates the hysteresis loops obtained for sodium RNA. As stated previously the samples used were not single crystals and were conducting, and a combination of these two factors is probably responsible for the general roundness and lack of saturation of the hysteresis loops. The loops below  $40^{\circ}$ C actually show very little evidence of hysteresis, those above  $58^{\circ}$ C show what is considered here to be quite good hysteresis and those loops between  $40^{\circ}$ C and  $58^{\circ}$ C show a gradual change between these two extremes. At temperatures above  $75^{\circ}$ C or  $80^{\circ}$ C the loops gradually disappeared as the sample denatured. All the loops appear to be very slightly asymmetric which may be indicative of an internal alignment





Figure 16. Hysteresis Loops of Sodium RNA

of dipoles during solidification.

It is quite evident that the three regions of the sodium RNA dielectric constant curve can be correlated with the three temperature divisions of the hysteresis loops. The gradual changes in the hysteresis loops correspond to the rather broad regions of the dielectric constant curve of 100 cps. The sharper regions of the dielectric constant curve of 1000 cps can not be correlated with hysteresis loops since it was not possible to obtain loops at a frequency of 1000 cps.

The five dielectric constant curves of 1000 cps were averaged and the result is plotted in Figure 17 for the temperature region of  $35^{\circ}$ C to  $50^{\circ}$ C. A Curie-Weiss law has been fitted to this curve and the result is also plotted in Figure 17. In the region between  $35^{\circ}$ C and  $50^{\circ}$ C the dielectric constant  $\epsilon$  obeys the following

$$\epsilon = 4.55 \times 10^{2} / (55.25 - T).$$
 (16)

Similar curves were obtained for the dielectric constant at 500 cps and 100 cps and are shown in Figures 18 and 19 respectively. In the temperature range between  $35^{\circ}C$  and  $55^{\circ}C$  the dielectric constant at 500 cps obeys

$$\epsilon = 5.15 \times 10^{5} / (61.2 - T),$$
 (17)

and in the same temperature range the dielectric constant at



Figure 17. Curie-Weiss Law for Sodium RNA at 1000 cps



Figure 18. Curie-Weiss Law for Sodium RNA at 500 cps



Figure 19. Curie-Weiss Law for Sodium RNA at 100 cps

100 cps obeys

$$\epsilon = 6.45 \times 10^5 / 58.2 - T$$
 (18)

These dielectric constant equations are not to be taken as proof that a ferroelectric transition exists. However, structural changes do take place in the nucleic acids as discussed in Chapter I. In view of this, it is reasonable to assume some sort of transition in the temperature range up to about 50°C. Such a transition may take place relatively slowly and cause a gradual onset of ferroelectricity. On the other hand, what appears to be the onset of ferroelectricity may actually be merely a decrease in the coercive field strength.

<u>Sodium DNA.</u> The dielectric constant curves for sodium DNA exhibit a slight irregularity in the temperature range between  $35^{\circ}C$  and  $45^{\circ}C$  (see Figures 20,43 and 44). The hysteresis loops shown in Figure 21 are asymmetric below about  $50^{\circ}C$  and symmetric above about  $70^{\circ}C$  with a transition region between  $50^{\circ}C$  and  $70^{\circ}C$ . The asymmetry in DNA is much greater than in similar samples of RNA (compare Figures 15 and 20). Insofar as it is possible to compare the hysteresis loops and dielectric constant of DNA in view of the difference in measuring frequencies, only a very slight correlation can be made. The irregularity in the dielectric constant curve can be compared with the change in hysteresis loops from asymmetric to symmet-



Figure 20. Dielectric Constant of Sodium DNA



ric even though the two temperature ranges do not overlap. Comparison of the dielectric constant curves of RNA at 1000 cps and 100 cps shows that the irregularity not only broadens but actually tends to shift to a slightly higher temperature. A similar shift here of only a few degrees centigrade would be sufficient to have the two temperature ranges overlap. Polyadenylic Acid Potassium. Figure 22 shows the dielectric constant curve for polyadenylic acid potassium. It is possible to broadly divide this curve into three regions: (1) the slowly increasing portion up to  $55^{\circ}C$  or  $60^{\circ}C$ , (2) the more rapidly increasing portion between  $60^{\circ}C$  and  $80^{\circ}C$ , and (3) the decreasing portion above  $80^{\circ}$ C. It is to be noted that the dielectric constant maximum value is approximately one forth of the maximum value of most other materials used here in the same temperature range. This is probably due to the greater stability of the double helix caused by the purine-phosphate bonding. The region above 80°C approximately corresponds to the helix-coil transition (denaturation) temperature. This portion of the curve was included here because the denaturing resulted in a smooth decrease in the dielectric constant. In all the other materials studied, when denaturing occurred, the dielectric constant change was rather abrupt and unsteady. Comparison of these temperature ranges with the hysteresis loops of Figure 23 shows some correlation. The middle temperature range of 60°C to 80°C corresponds to loops in the range between  $45^{\circ}$ C and  $75^{\circ}$ C. Above and below









\_ \_\_













75 °C





←\_\_\_\_↓ ↓ ↓ volt volts





Figure 23. Hysteresis Loops of Polyadenylic Acid Potassium

this range the loops have less definite hysteresis shape. <u>Polycytidylic Acid Potassium.</u> A very slight bulge can be observed in the dielectric constant curve of polycytidylic acid in the temperature range of about  $50^{\circ}$ C to  $70^{\circ}$ C (see Figure 24). Figure 25 shows the hysteresis loops for this material. Below about  $40^{\circ}$ C the loops are asymmetric, between  $40^{\circ}$ C and  $45^{\circ}$ C the loops become symmetric and above  $45^{\circ}$ C the loops are again asymmetric but in the opposite direction. This latter region probably results from a combination of the onset of denaturation and increased conductivity of the sample. Any correlation of the change in the shape of the loops to the slight bulge in the dielectric constant curve must be superficial.

Adenylic Acid Sodium. The slight bulge in the dielectric constant curve of adenylic acid at approximately  $40^{\circ}$ C as shown in Figure 26 may be related to the change in the hysteresis loops up to  $50^{\circ}$ C, as seen in Figure 27. The rapid increase in the dielectric constant curve above  $55^{\circ}$ C corresponds to the change in the symmetry of the hysteresis loops. These latter changes are probably due again to conductivity and denaturation.

<u>Guanylic Acid Disodium.</u> The bulge in the dielectric constant curve of this purine containing salt, as shown in Figure 28, and the temperature region in which change occurs in the hysteresis loops shown in Figure 29, compare favorably with the similar changes in adenylic acid sodium. Here the



Figure 24. Dielectric Constant of Polycytidylic Acid Potassium





33 °C



37 °C





47 °C

57 <sup>o</sup>C











66 <sup>o</sup>C











Figure 25. Hysteresis Loops of Polycytidylic Acid Potassium



Figure 26. Dielectric Constant of Adenylic Acid Sodium





Figure 27. Hysteresis Loops of Adenylic Acid Sodium



Figure 28, Dielectric Constant of Guanylic Acid Disodium




Figure 29. Hysteresis Loops of Guanylic Acid Disodium

bulge occurs closer to 45<sup>°</sup>C, the increase in the dielectric constant above the bulge is more rapid and the hysteresis loops become asymmetric rather than symmetric. However, in a general way guanylic acid disodium and adenylic acid sodium behave similarly.

Cytidylic Acid Disodium. The dielectric constant curve of cytidylic acid shown in Figure 30, shows a very smooth increase with no irregularities. The value of the dielectric constant does not become appreciable until approximately 65°C and then increases rapidly within a range of 15 degrees. This may be correlated to the hysteresis loops of Figure 31 which shows symmetric loops only in the range between  $70^{\circ}$ C and  $80^{\circ}$ C. Thymidylic Acid Sodium. Figure 32 is the dielectric constant curve for this material and the curve has the same general shape as that for cytidylic acid disodium. The hysteresis loops of Figure 33 are of poor quality but can be related to the dielectric constant curve in the same manner as was done with cytidylic acid disodium. Hence the pyrimidine containing salts, like the purine containing salts, behave similarly. Pressed Wafer Dielectric Constants. It was not possible to obtain hysteresis loops with any of the materials in the pressed wafer form. This was probably due to the extremely small crystallites and the masking effect of entrapped air. Figures 34 and 35 show the dielectric constant curves of adenylic acid sodium and uridylic acid disodium respectively. It was not possible to obtain similar curves for the pressed



Figure 30. Dielectric Constant of Cytidylic Acid Disodium



Figure 31. Hysteresis Loops of Cytidylic Acid Disodium



Figure 32. Dielectric Constant of Thymidylic Acid Sodium





Figure 33. Hysteresis Loops of Thymidylic Acid Sodium



Figure 34. Dielectric Constant of Adenylic Acid Sodium Wafer



Figure 35. Dielectric Constant of Uridylic Acid Disodium Wafer

wafer of either adenosine or adenine. Over the temperature range used, approximately  $25^{\circ}C$  to  $100^{\circ}C$ , the dielectric constants of adenosine and adenine remained extremely small.

The ferroelectric hysteresis of the materials studied here may arise from the permanent dipoles of the sugar, phosphate and bases, the various hydrogen bonds, or a combination of these sources. In the case of DNA the hydrogen bonds between the two strands may be responsible for the ferroelectricity, but if this is the sole source, single stranded RNA should not be ferroelectric (24). It has been proposed that the ferroelectric hysteresis of RNA is due to the stacking of bases but this explanation may not be adequate to explain the ferroelectricity of the smaller molecules such as adenylic acid (24). The ferroelectricity of the materials studied here probably arises from a combination of the permanent dipoles of the molecules and the hydrogen bonds. However, it is likely that the main contribution of the ferroelectricity originates in the phosphate group. This can be seen by considering the following: (1) the phosphate group is the most polar part of the sodium DNA molecule and hence also of the nucleotides (18), (2) the essential difference between adenosine and adenylic acid is the lack of a phosphate group in adenosine, as indicated by Figures 4 and 5, and (3) the dielectric constant of the pressed wafer of adenylic acid sodium is quite large and increases greatly with increasing temperature while the dielectric constant of the pressed wafer of adenosine

remains quite small in the same temperature range. Furthermore, it has been reported that the phosphate groups in these biological molecules have a tetrahedral arrangement (20,21,25). This information is pertinent since a partially successful theory by Slater (1941) bases the ferroelectricity of  $KH_2PO_4$ on the orientation of the  $(H_2PO_4)^-$  dipoles, where the oxygen atoms are located at the corners and the phosphorus atoms at the centers of tetrahedra which are hydrogen bonded to one another (26). The hydrogens do not lie midway between oxygen atoms but rather may exist in crystallographically non-equivalent positions with different energies. Thus an explanation of the ferroelectricity of these biological molecules based on the orientation of phosphate group dipoles is entirely plausible.

### CHAPTER V

## CONCLUSIONS AND RECOMMENDATIONS

It is advisable to briefly review the nature of the molecules studied here. The five nucleic acid bases consist of two purines, adenine and guanine, and three pyrimidines. cytosine, thymine and uracil. A nucleoside is the combination of any base and a pentose sugar. The nucleoside is a riboside or a deoxyriboside depending on whether the sugar is ribose or deoxyribose respectively. The nucleosides are adenosine, guanosine, cytidine, thymidine and uridine. The molecule formed from a bonding of a phosphoric acid molecule to the sugar molecule of a nucleoside is a nucleotide. The pentose sugar determines if these are ribonucleotides or deoxyribonucleotides. The nucleotides are adenylic acid, guanylic acid, cytidylic acid. thymidylic acid and uridylic acid. Of the nucleotides used here, only thymidylic acid is a deoxyribonucleotide. The bonding of nucleotides, phosphate group to sugar group, into chains result in either RNA or DNA. If only one type of nucleotide forms a chain, the result is a polyribonucleotide or a polydeoxyribonucleotide. In this work the sodium or potassium salts of the various acids were utilized.

Samples of a crystal nature were obtained for the salts of RNA, DNA, adenylic acid, guanylic acid, cytidylic acid

thymidylic acid, polyadenylic acid and polycytidylic acid. Evaporation of aqueous solutions of the salts produced these samples. Unsaturated hysteresis loops were observed for all these materials. Furthermore, they displayed increasing dielectric constants with increasing temperature which is characteristic of many ferroelectric materials.

A lower transition point seems to exist for sodium RNA although this effect may be due to a lowering of the value of the coercive field. Irrespective of this, the ferroelectric character of RNA has been demonstrated for the first time. The sequence of the nucleotides in the macromolecules does not affect their dielectric properties since RNA, DNA, polyadenylic acid and polycytidylic acid all exhibit this ferroelectric character. Ferroelectric behavior was observed in specimens containing: (1) ribose sugar, RNA, polyadenylic acid, polycytidylic acid, adenylic acid, guanylic acid and cytidylic acid, and (2) deoxyribose sugar, DNA and thymidylic acid. Hence, the type of pentose sugar present does not significantly affect the ferroelectricity of these materials. Since the purine nucleotides and pyrimidine nucleotides all appear to be ferroelectric, and since the dielectric properties of the purine nucleotides are similar and the pyrimidine nucleotides are similar, this property is independent of the type of base present. Comparison of the ferroelectric properties of adenylic acid and polyadenylic acid. and cytidylic acid and polycytidylic acid, indicates that the length of the

molecule is not a determining factor of the ferroelectricity of the molecule. In summary, the existance of ferroelectricity in the nucleic acids appears to be independent of the type of pentose sugar present, the types of bases present, the sequence of bases and the length of the molecule. Hence, one might speculate that in a living system, the ferroelectricity of the nucleic acids is not confined to any particular type of nucleic acid.

It was not possible to obtain samples of a crystal nature for adenine, adenosine or uridylic acid disodium. Wafers were made of the powdered material using a hydraulic press. For comparison purposes, wafers were also made of adenylic acid sodium. Hysteresis was not observed for any of these materials in the wafer form. The dielectric constants of the wafers of adenylic acid sodium and uridylic acid disodium both exhibit the characteristic behavior of many ferroelectric materials. However, it is not possible to say that uridylic acid disodium is ferroelectric. The dielectric constants of the wafers of adenine and adenosine did not show any characteristic ferroelectric behavior and were quite small. The results of the dielectric constant measurements of adenine and the adenine containing nucleoside and nucleotide, suggest that the phosphate group is the origin of the ferroelectricity of the materials studied here.

From the results of this work and from the conclusions reached, the following recommendations are made: (1) Ferro-

electric and pyroelectric studies should be made on the fiber form of the nucleic acids and polynucleotides, and on single crystals of the nucleotides and salts of these materials. These forms of the materials should result in better defined hysteresis loops which indicate saturation. (2) Studies should be made of the nucleosides and nucleotides to test whether or not the phosphate group is the origin of the ferroelectricity. Furthermore, deuterated nucleotides should be studied to determine the effect of the hydrogen bonds on the ferroelectricity. (3) Controlled humidity and temperature experiments should be performed to determine the effect of these on the electrical properties and on the crystal structure. Any observed effects may be useful in determining the nature of any transition which might occur.

#### APPENDIX A

## DIPOLES AND FERROELECTRICITY

In attempting to describe or determine the origin of the ferroelectricity of some material it is desirable to associate dipoles with the molecules of the material. This is true since many theoretical works on ferroelectricity have dipole orientation as their basis. However, it is not always possible to select pairs of equal and opposite charges with associated dipole moment given by the product of charge and separation. In the case of non-molecular crystals where atoms are linked in a three dimensional array it may be unreasonable to ascribe special associations between specific pairs of atoms. These pairs may not retain any special relationship with respect to their neighbors since all atoms may be symmetrically related. An example of this is sodium chloride where each atom of sodium is associated with six atoms of chlorine and each chlorine atom is associated with six sodium atoms.

Dipole identification in many instances is not only possible but also quite reasonable. Two cases where it might be possible to ascribe dipoles to a particular substance, not necessarily ferroelectric, are: (1) in an ionic crystal when application of an electric field causes separa-

tion of charges within each ion and (2) in a crystal which contains molecules which may be dipoles and are free to rotate with respect to the rest of the structure. In the case of some ferroelectrics, it may be possible to assign dipoles on the basis of both electrical and crystallographical information. Such an assignment is possible in the case of KH<sub>2</sub>PO<sub>4</sub>. The phosphorus and oxygen atoms form a tetrahedron with the oxygen atoms at the corners and the phosphorus atom at the center. Each tetrahedra is hydrogen bonded to four others. An assumption is made that each tetrahedra has only two hydrogens near it, the other two hydrogens being nearer other tetrahedra. This assumption makes the  $(H_2PO_4)^-$  group an electric dipole (12). Finally, in the case of BaTiO3 the dipole is assigned to the octahedron formed with oxygen at its corners and the off centered positioned titanium atom (12).

### APPENDIX B

## SAWYER-TOWER CIRCUIT WITH LINEAR ELEMENTS

Figure 36 shows three Sawyer-Tower circuits containing ideal linear elements.  $C_x$  and  $C_o$  are capacitances,  $R_x$  is a resistance,  $V_o$  is the voltage across  $C_o$  and  $V_{in}$  is the input voltage which is assumed to be sinusoidal and given by

$$V_{in} = E_o \sin \omega t$$
, (19)

where  $\mathbf{E}_{\mathbf{O}}$  is a magnitude,  $\boldsymbol{\omega}$  is the angular frequency and t is the time.

An oscilloscope is usually employed to display voltage when using the Sawyer-Tower circuit. The input voltage  $V_{in}$ of the circuit is the horizontal input to the oscilloscope. Thus, the x coordinate of the trace on the oscilloscope is just the value  $V_{in}$  and is given by the equation

$$x = E_0 \sin \omega t$$
. (20)

The vertical input to the oscilloscope is the voltage  $V_0$ , which is the y coordinate of the trace on the oscilloscope. The equation of the oscilloscope trace may be obtained by relating x and y. The voltage  $V_0$  is just the product of the input voltage  $V_{in}$  and the ratio of the impedance of  $C_0$ 



Figure 36. Sawyer-Tower Circuit with Linear Elements

to the total impedance.

Impedances may be written as complex exponentials and this practice will be followed here. In the following, the letter Z stands for impedance. The impedance of  $C_0$  may be written as

$$Z_{o} = (1/\omega C_{o}) \exp(-i\pi/2)$$
 (21)

The impedance of  $C_{\mathbf{x}}$  is

$$Z_{\mathbf{X}} = (1/\omega C_{\mathbf{X}}) \exp(-i\pi/2) . \qquad (22)$$

The impedance of  $R_x$  is just the value of the resistance, and the impedance of  $R_x$  in parallel with  $C_x$  is

$$Z_{x} = (1/R_{x}^{2} + \omega^{2}C^{2})^{-\frac{1}{2}}exp(-i\theta), \qquad (23)$$

where  $\theta$  is related to  $\omega$  ,  $R_x$  and  $C_x$  by the equation

$$\tan \theta = \omega R_{\mathbf{X}} C_{\mathbf{X}} . \tag{24}$$

For the circuit containing the two capacitances  $C_0$ and  $C_x$ , the total impedance is

$$Z = ((C_0 + C_x) / \omega C_0 C_x) \exp(-i \pi / 2).$$
(25)

The voltage  $V_0$ , and hence coordinate y, is then given by

$$\mathbf{V}_{\mathbf{O}} \approx (\mathbf{C}_{\mathbf{X}} / \mathbf{C}_{\mathbf{O}} + \mathbf{C}_{\mathbf{X}}) \mathbf{E}_{\mathbf{O}} \sin \omega t .$$
 (26)

Equation (20) may be used to eliminate  $\sin \omega t$  from Equation (26) and the result, relating y and x, is

$$y = (C_x (C_0 + C_x))x$$
 (27)

which is the equation of a straight line.

For the circuit containing  $R_{\mathbf{X}}$  and  $C_{\mathbf{O}}$ , the total impedance is

$$Z = (R_x^2 + 1/\omega^2 C_0^2)^{\frac{1}{2}} exp(-i\phi), \qquad (28)$$

where  $\phi$  is related to  $\omega$  ,  $R_x$  and  $C_0$  by the equation

$$\tan \phi = 1/\omega R_{\rm x} C_{\rm o} . \tag{29}$$

The voltage  $V_0$  is then given by the equation

$$\mathbf{V}_{\mathbf{O}} = \mathbf{K} \mathbf{E}_{\mathbf{O}} (\sin \omega t \cdot \sin \phi - \cos \omega t \cdot \cos \phi) , \qquad (30)$$

where K is a constant given by the equation

$$K = (1/\omega C_0) / (R_X^2 + 1/\omega^2 C_0^2)^{\frac{1}{2}}.$$
 (31)

Use of Equation (20) and the substitution of y for  $V_0$  yields the following expression

$$K^2 x^2 - 2K \sin \phi xy + y^2 - K^2 \cos^2 \phi E_0^2 = 0$$
. (32)

Equation (32) is of the form

$$Ax^2 + Bxy + Cy^2 + Dx + Ey + F = 0$$
, (33)

which is an equation of an ellipse or its degenerate case if the inequality

$$\mathbf{B}^2 - \mathbf{4AC} < \mathbf{0} \tag{34}$$

is satisfied (27). Substitutions from Equation (32) yields

$$B^{2}-4AC = 4K^{2}\sin^{2}\phi - 4K^{2} = 4K^{2}(\sin^{2}\phi - 1).$$
 (35)

From Equation (29),  $\sin \phi$  may be expressed as

$$\sin\phi = (1/\omega C_0)/(R_x^2 + 1/\omega^2 C_0^2)^{\frac{1}{2}}.$$
 (36)

Since the resistance is non zero by assumption,  $\sin \phi$  can not be equal to one. Hence, since K is positive, the inequality is satisfied and the oscilloscope trace is an ellipse or its degenerate case. For the circuit containing  $R_x$ ,  $C_x$  and  $C_o$ , the resulting equation relating x and y is identical in form to Equation (32) but with K given by

$$K = (1/\omega C_0)/(a^2 + b^2)^{\frac{1}{2}}$$
(37)

and with  $\phi$  expressed by the relation

$$\tan\phi = \mathbf{b}/\mathbf{a} , \qquad (38)$$

where b is given by

$$\mathbf{b} = (1 + \omega^2 R_x^2 C_x^2 + \omega^2 R_x^2 C_x C_0) / \omega C_0 (1 + \omega^2 R_x^2 C_x^2)$$
(39)

and a is given by

$$a = R_{x}^{2} / (1 + \omega^{2} R_{x}^{2} C_{x}^{2}) .$$
 (40)

The inequality given above is satisfied here just as in the case of the circuit containing only  $R_x$  and  $C_0$ . In this case the equation analogous to Equation (36) would require either  $\omega$ ,  $R_x$  or  $C_0$  to be zero in order that  $\sin \phi$  would equal one. Since none of these are zero by assumption, the resulting trace is an ellipse or its degenerate case.

## APPENDIX C

# SAMPLE CHARACTERISTICS

This appendix consists of tables describing the characteristics of the samples used in this research. It is to be noted here that the drying of cytidylic acid disodium occurred at room temperature, not at five degrees centigrade. The water content refers to the amount of water initially added to obtain the solution and not the final water content of the samples. In the cases of the polyribonucleotides and guanylic acid disodium, as water from the sample cup evaporated more solution was added in order to obtain samples of sufficient thickness. Table 2. Sample Characteristics

Material	Mass grams	Water Content cc	Drying Time days	Physical Appearance After Drying
Adenylic Acid Sodium	2.0	2.0	7-8	Transparent, colorless, multi-cracked, single piece
Cytidylic Acid Disodium	2.0	2.0	5-6	Transparent, colorless, single aggregate of needle shaped fibers
Guanylic Acid Disodium	2.0	36.0	1-2	Translucent, white, single piece, brittle, uncracked
Th <b>ym</b> idyl <b>ic</b> Acid Sodium	1.0	1.0	7-8	Transparent, slightly yellow, cracked into many small pieces
Sodium DNA	5.0	7.0	1-2	Transparent, brown, cracked into irregular sizes and shapes
Sodium RNA	5.0	8.0	1-2	Transparent, brown cracked into irregular sizes and shapes
Po <b>ly- adenylic</b> Acid Potassium Ly <b>oph</b> ilized	1.0	100	4-5	Translucent, white single piece, not brittle, uncracked
Poly- cytidylic Acid Potassium Lyophilized	0.50	35	4-5	Transparent, colorless, single piece, not brittle, uncracked

# Table 3. Wafer Characteristics

Material	Mass grams	Pressure Applied to Form Wafers psi
Adenine	1.0	45000
Adenosine	1.0	45000
Adenylic Acid Sodium	1.0	45000
Uridylic Acid Sodium	1.0	45000

### APPENDIX D

## DIELECTRIC CONSTANT COMPILATION

The tables in this appendix contain the data of the dielectric constant curves appearing both in this appendix and in Chapter IV. The heading in the tables labled "Figure Number" correlates each dielectric constant curve to a particular listing in the tables. The heading, "Heating Curve", refers to the curves of Figure 14. The measuring voltage V is the voltage applied across the series combination of  $C_o$  and the sample capacitance. The maximum electric field is the approximate field applied to the sample during the time when the sample capacitance is much less than the capacitance of  $C_o$ . Unless otherwise stated the measuring frequency is 1000 cycles per second.

Figure Number	15			37	38		
Heating Curve	II		II		II		
Electrode Area A cm <sup>2</sup>	0.340		0.340		0.186		
Area Diameter D cm	0.6	0.657		0.657		0.487	
Sample Thickness d cm	0.0	0.0456		0.0438		0.0394	
Ratio D/d	14.1		15.0		16.7		
Measuring Voltage V volts	0.63		0.63		0.63		
Maximum Electric Field V/d volts/	13.5 cm	ò	1	4.4	16	•0	
Temperature <sup>O</sup> C	€ Pe Er	rcent ror ±	¢	Percent Error ±	۶ ا	Percent Error ±	
25 30 35 40 45 50 51 52 53 54 56 58 62 55 58 62 50 72 79 80 85	8800 14500 23000 38700 61800 99000 107000 112000 114000 114000 107000 103000 101000 105000 109000 165000 200000	20 17 15 15 12 12 12 12 12 12 12 12 12 12 12 12	1100 1670 2700 4300 6800 10800 11300 11500 11400 10100 10100 10100 10500 11800 16400 20000 33600	0       20         0       17         0       15         0       15         0       15         0       12	12600 22000 35000 55500 89400 95300 100000 103000 102000 94800 90000 87500 87300 93000 118000	17 15 15 14 12 12 12 12 12 14 14 12	



Figure 37. Dielectric Constant of Sodium RNA



Figure 38. Dielectric Constant of Sodium RNA

Table	5.	Dielectric	Constant	of	Sodium	Ribonucleate
TOUTO	1.	DICICOULIO	oonsound	U.L	DOUTAW	TT DOUGOTCOUC

Figure Number	39	40	
Heating Curve	II	II	
Electrode Area A cm <sup>2</sup>	0.186	0.340	
Area Diameter D cm	0.487	0.657	
Sample Thickness d cm	0.0533	0.0483	
Ratio D/d	9.14	13.6	
Measuring Voltage V volts	0.60	0.60	
Maximum Electric Field V/d volts/cm	11.3	12.4	
Temperature <sup>O</sup> C	• Percent Error ±	<ul> <li>ℓ Percent</li> <li>Error ±</li> </ul>	
30 35 40 45 50 52 54 56 58 65 70 75 80	$\begin{array}{cccccc} 10700 & 25 \\ 17100 & 25 \\ 25300 & 20 \\ 37000 & 20 \\ 59000 & 20 \\ 67000 & 20 \\ 73700 & 20 \\ 75300 & 20 \\ 76600 & 20 \\ 76600 & 20 \\ 78600 & 20 \\ 89000 & 20 \\ 108000 & 15 \\ 147000 & 15 \\ 200000 & 15 \\ \end{array}$	$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	



Figure 39. Dielectric Constant of Sodium RNA



Figure 40. Dielectric Constant of Sodium RNA

i

Figure Number	41	42
Heating Curve	II	II
Electrode Area A cm <sup>2</sup>	0.186	0.186
Area Diameter D cm	0.487	0.487
Sample Thickness d cm	0.0445	0.0508
Ratio D/d	10.9	9.58
Measuring Voltage V volts	0.60	0.60
Maximum Electric Field V/d volts/cm	13.5	11.8
Frequency cps	100	500
Temperature <sup>O</sup> C	← Percent     Error     Err	t e Percent ± Error ±
30 35 40 45 50 55 60 65 69 70 71	28000 18 33000 15 46000 15 88000 12 160000 12 221000 12 289000 12 382000 12 382000 12	147002019600182540015330001547900128200012120000121470001217500012
72.5 75 80 82	523000 12 626000 12	200000 12 235000 12 324000 12 371000 12



Figure 41. Dielectric Constant of Sodium RNA at 100 cps



Figure 42. Dielectric Constant of Sodium RNA at 500 cps

i.

Table 7. Dielectric Constant of Sodium Deoxyribonucleate

Figure Number	r 20			43	44		
Heating Curve	IV		IV		II		
Electrode Area A cm <sup>2</sup>	0.340		0.340		0,186		
Area Diameter D cm	0.657		0.657		0,487		
Sample Thickness d cm	0.0	533	C	0.0724	0	.0546	
Ratio D/d	12.3	5	9.08		8.92		
Measuring Voltage V volts	0.6	0.63		0.63		0.63	
Maximum Electric Field V/d volts/	11.8 cm	•	8	3.71	11.	.5	
Temperature <sup>O</sup> C	€ Pe Er	rcent ror ±	E	Percent Error ±	د ] ۲	Percent Error ±	
25 30 35 37.5 40 42.5 45 50 55 60 55 60 65 70 72.5 74.5 76.5 77	20600 37000 61500 71000 74800 75700 89200 99800 109000 122000 143000 159000 176000	12 12 10 10 10 10 10 10 10 10 10 10	30000 46800 75800 83600 91900 96600 109000 124000 157000 178000 200000	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	33000 56500 71300 80700 85500 89000 103000 117000 134000 154000 154000 154000 200000	12 12 10 10 10 10 10 10 10 10	
80 83 85	233000	10	320000 363000	) 10 ) 10	345000	10	


Figure 43. Dielectric Constant of Sodium DNA



Figure 44. Dielectric Constant of Sodium DNA

Figure Number	22	45	46
Heating Curve	II	II	II
Electrode Area A cm <sup>2</sup>	0.340	0.340	0.340
Area Diameter D cm	0.657	0.657	0.657
Sample Thickness d cm	0.0013	0.0013	0.0013
Ratio D/d	506	506	506
Measuring Voltage V volts	0.25	0.25	0.25
Maximum Electric Field V/d volts/c	m 192	192	192
Temperature <sup>O</sup> C	€ Percent Error ±		<ul> <li>ℓ Percent</li> <li>Error ±</li> </ul>
30 35 40 45 50 55 60 62.5 67.5 70 75 80 82.5 80 82.5 85 90 95	$\begin{array}{ccccccc} 2000 & 20 \\ 4300 & 18 \\ 6100 & 15 \\ 7700 & 15 \\ 9500 & 15 \\ 11300 & 15 \\ 12500 & 15 \\ 13600 & 15 \\ 15600 & 15 \\ 19200 & 12 \\ 25300 & 12 \\ 37800 & 12 \\ 41700 & 12 \\ 41300 & 12 \\ 41300 & 12 \\ 37100 & 12 \\ 30600 & 12 \\ 20000 & 12 \end{array}$	$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	$\begin{array}{cccccc} 4500 & 18 \\ 5800 & 15 \\ 6900 & 15 \\ 7800 & 15 \\ 9000 & 15 \\ 10700 & 15 \\ 12500 & 15 \\ 15000 & 15 \\ 19400 & 12 \\ 26400 & 12 \\ 35400 & 12 \\ 47300 & 12 \\ 51300 & 12 \\ 51300 & 12 \\ 49600 & 12 \\ 44400 & 12 \\ 36400 & 12 \\ 23000 & 12 \end{array}$

Table 8. Dielectric Constant of Polyadenylic Acid Potassium

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Figure 45. Dielectric Constant of Polyadenylic Acid Potassium



Figure 46. Dielectric Constant of Polyadenylic Acid Potassium

		1(	)2

# Table 9. Dielectric Constant of Polycytidylic Acid Potassium

Figure Number	24	ł		47		48
Heating Curve	I	I		II		II
Electrode Area A cm <sup>2</sup>	0.(	340		0.340	С	.340
Area Diameter D cm	0.6	557		0.657	C	.657
Sample Thickness d cm	0.0	0140		0.0140	0	.0140
Ratio D/d	46.9	7	4	6.9	46	•9
Measuring Voltage V volts	0.2	25		0.25	0	.25
Maximum Electric Field V/d volts/d	17.8 cm	3	1	7.8	17	.8
Temperature <sup>O</sup> C	€ Pe Ei	ercent rror ±	€	Percent Error ±	ę	Percent Error ±
25 30 35 40 45 50 55 60 65 70 75 77.5 80 82.5 83.5	3000 6000 9700 15700 23700 36700 54500 76500 102000 129000 172000 200000	20 18 155 155 155 155 155 155 15	380 710 1900 2820 4080 5570 7750 10500 12800 14800 16000 17300	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	3400 9000 13000 19400 31000 50500 72000 89800 111000 137000 153000 192000 200000	20 18 15 15 15 15 15 15 15 15 15 15 15 15 15
84.5 85 92 97	251000	15	20000 25100 26000	0 15 0 15 0 15	251000	15



Figure 47. Dielectric Constant of Polycytidylic Acid Potassium



Figure 48. Dielectric Constant of Polycytidylic Acid Potassium

Table 10. Dielectric Constant of Adenylic Acid Sodium

Figure Number	26	49	50
Heating Curve	I	I	II
Electrode Area A cm <sup>2</sup>	0.840	0.840	0.840
Area Diameter D cm	1.03	1.03	1.03
Sample Thickness d cm	0.0406	0.0457	0,0488
Ratio D/d	25.4	22.6	21.1
Measuring Voltage V volts	0.63	0.63	0.63
Maximum Electric Field V/d volts/d	15.5 cm	13.8	12.9
Temperature <sup>O</sup> C	€ Percen Error	t e Percent ± Error ±	<ul> <li>Percent</li> <li>Error ±</li> </ul>
28 30 35 40 45 50 52.5 57.5 57.5 60 65 67.5 70 72.5 74 80 82	30001853001512000121750012225001229000123400012445001262000121400001215800012	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	46001512700121970012243001227100123140012390001255500121100001214000012169000122210001231900012
85 86		329000 12	354000 12



Figure 49. Dielectric Constant of Adenylic Acid Sodium



Figure 50. Dielectric Constant of Adenylic Acid Sodium

.

Figure Number	28	3	5	1	:	52
Heating Curve	II		I	I		II
Electrode Area A cm <sup>2</sup>	0.1	.86	0.	186	0	.186
Area Diameter D cm	0.4	187	0.	487	0	.487
Sample Thickness d cm	0.0	203	0.	0223	0	.0254
Ratio D/d	24.0	)	21.	8	19	.2
Measuring Voltage V volts	0.1	10	0.	40	0	.40
Maximum Electric Field V/d volts/	19.7 cm	*	17.	9	15	.7
Temperature <sup>O</sup> C	e Pe Er	rcent ror ±	€ P E	ercent rror ±	€ ]	Percent Error ±
30 34 36 38 40 42 44 46 48 49.5 50 52.5 57.5 57.5 59.5 59.5 60 61	7300 9000 10300 12000 14000 29000 45500 60000 63300 72400 88200 120000 200000 268000	22 22 22 22 22 22 20 20 20 20 20 20 20 2	8800 13600 18400 31000 51400 69000 85200 103000 143000 207000 378000	22 22 20 20 20 20 20 20 20 20 20	8900 12000 14800 19300 28000 46000 67300 84000 98400 111000 136000 200000 334000	22 22 22 20 20 20 20 20 20 20 20 20 20
62 64 69	344000 599000	20 20	976000	20	909000	20



Figure 51. Dielectric Constant of Guanylic Acid Disodium



Figure 52. Dielectric Constant of Guanylic Acid Disodium

Figure Number	30	53
Heating Curve	II	II
Electrode Area A cm <sup>2</sup>	0.186	0.340
Area Diameter D cm	0.487	0.657
Sample Thickness d cm	0.0635	0.0762
Ratio D/d	7.67	8.63
Measuring Voltage V volts	0.60	0.60
Maximum Electric Field V/d volts/cm	9.45	7.88

Temperature	oC	£	Percent Error ±	€	Percent Error ±
65		5600	22	2500	22
67.5		12800	20	5700	22
70		24500	20	11000	20
72		45300	20	17800	20
74		75000	20	30000	20
76	1	25000	20	50000	20
77	1	70000	20	64500	20
78	2	00000	20	82500	20
8 <b>0</b>	2	95000	20	127000	20
82	4	05000	20	179000	20
82.5				200000	20
84	4	94000	20	226000	20
87	5	87000	20	283000	20
91	7	45000	20	359000	20
98	ġ	15000	20	385000	20



Figure 53. Dielectric Constant of Cytidylic Acid Disodium

Figure Number	32	54
Heating Curve	II	II
Electrode Area A cm <sup>2</sup>	0.0410	0.0410
Area Diameter D cm	0.228	0.228
Sample Thickness d cm	0.0305	0.0267
Ratio D/d	7.48	8.54
Measuring Voltage V volts	0.30	0.30
Maximum Electric Field V/d volts/cm	9.84	11.2
Temperature <sup>O</sup> C	€ Percen Error	t c Percent ± Error ±
65 67.5 70 75 80 85 90 95 100	440002853000257000025138000223440002276500022119000022147000022163000022	70000258900025127000223000002270000022117000022152000022



Figure 54. Dielectric Constant of Thymidylic Acid Sodium

t of Adenylic	Acid	Sodium	Wafer
34		55	
III	]	II	
2.85	2	2.85	
1.91	]	.91	
0.216	C	.212	
3.85	ç	9.02	
<b>-</b> -			

Heating Curve	III	III
Electrode Area A cm <sup>2</sup>	2.85	2.85
Area Diameter D cm	1.91	1.91
Sample Thickness d cm	0.216	0.212
Ratio D/d	8.85	9.02
Measuring Voltage V volts	7.0	7.0
Maximum Electric	32.4	33.0

Field V/d volts/cm

Figure Number

Temperature	°C	E	Percent Error ±	E	Percent Error ±
50		1000	12	1000	12
55		2000	12	1700	12
60		3200	9	2700	9
65		7800	8	4800	9
70		28000	7	14000	7
75		71000	7	46200	7
80		153000	7	130000	7
82		200000	7	172000	7
83				200000	7
86		280000	7	275000	7
87				355000	7
87.5		362000	7		
90				441000	7
91		450000	7		
<b>9</b> 9		766000	7	752000	7

#### Table 14. Dielectric Constant r



Figure 55. Dielectric Constant of Adenylic Acid Sodium Wafer

с	Acid	Disodium	Waf
		56	
	-	гтт	

Table 15. Dielectric Constant of Uridylic Acid Disodium Wafer

Figure Number		35	56	
Heating Curve	נ	III	III	
Electrode Area A cm <sup>2</sup>	2	2.85	2.85	
Area Diameter D cm	נ	91	1.91	
Sample Thickness d cm	C	.216	0.216	
Ratio D/d	8	3.85	8.85	
Measuring Voltage V volts	7	.0	7.0	
Maximum Electric Field V/d volts/cm	32.4		32.4	
Temperature <sup>o</sup> C	ŧ	Percent Error ±	€ P E	ercent rror ±
30 35 40 45 50 55 60 63.5 67	$     1000 \\     4000 \\     10000 \\     21000 \\     43000 \\     83000 \\     144000 \\     200000 $	14 11 10 9 9 9 9	$     1000 \\     1800 \\     4300 \\     10000 \\     22000 \\     48000 \\     93000 \\     148000 \\     200000 $	14 13 11 10 9 9 9 9
70 71	386000	9	288000	9
75 77	514000	9	463000	9
80 85 80	703000 891000	9 9	788000	9
90	1550000	2	1220000	9



Figure 56. Dielectric Constant of Uridylic Acid Disodium Wafer

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VITA

Richard Anthony Lorey was born in Pittsburgh, Pennsylvania on December 28, 1939. He received his elementary and secondary education in the Pittsburgh public school system and graduated from Allegheny High School in 1957. He attended the University of Pittsburgh and received the degree Bachelor of Science in Physics in 1962. He served in the United States Army as a first lieutenant assigned to the Nuclear Defense Laboratory in Edgewood Arsenal, Maryland until 1965. It was during this time that he married Virginia Joan Longo on December 5, 1964. Their only son, Brandon Craig Lorey, was born in Atlanta, Georgia on November 10, 1968. Upon completion of his tour of active duty he entered graduate school in the School of Physics at the Georgia Institute of Technology. He received the degree Master of Science in Physics from the Georgia Institute of Technology in 1967. He was the recipient of a National Science Foundation Traineeship during 1967-1968 and a United States Steel Fellowship during 1968-1969.

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