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A MUSIC-type Algorithm for Detecting Internal Corrosion from Electrostatic Boundary Measurements ^{*}

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Abstract

We establish an asymptotic representation formula for the steady state current perturbations caused by internal corrosive boundary parts of small surface measure. Based on this formula we design a non-iterative method of MUSIC (multiple signal classification) type for localizing the corrosive parts from voltage-to-current observations. We perform numerical experiments to test the viability of the algorithm and the results clearly demonstrate that the algorithm works well even in the presence of relatively high noise ratios.

Mathematics subject classification (MSC2000): 35R30

Keywords: corrosion, asymptotic representation formula, reconstruction, layer potentials

1 Introduction

Hidden corrosion, for instance in pipes, is a serious problem that is responsible for significant economic losses and represents a dangerous threat to safety. In the field of nondestructive evaluation, new and improved techniques are therefore constantly being sought to facilitate the detection of such hidden corrosion.

Corrosion occurs in many different forms, and several mathematical models for electrostatic fields in the presence of corrosion have been studied in the literature (see, for example, Kaup and Santosa [16], Kaup et al. [17], Vogelius and Xu [21], Inglese [12], Luong and Santosa [13], Banks et al. [3] and references therein).

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In this paper, we adopt the potential model: Laplace's equation holds in the metal pipe, and the effect of corrosion is described by means of a linear Robin boundary condition. A model like this is motivated by a number of observations. On the one hand, corrosion tends to roughen a surface: this effect can be modelled by the introduction of a thin coating characterized by rapid oscillations. In the limit as the thickness of the coating goes to zero and the rapidity of the oscillations diverges, the emergence of a Robin boundary conditions has been shown by Buttazo and Kohn [5]. On the other hand, the study of active electrochemical corrosion processes can be based on "first principles", which conjecture a certain relation between the number of "free electrons" and the energy level (typically an exponential Boltzmann-type distribution). Potential models with exponential boundary conditions that arise from such considerations are often associated with the names of Butler and Volmer. Vogelius and Xu [21] study a potential model of this kind. If we linearize the boundary condition of this model, we get the linear Robin boundary condition studied here.

To set up the problem mathematically, consider a simply connected, bounded $C^{2,\alpha}$ domain U in \mathbb{R}^2 for some $0 < \alpha \leq 1$, and a simply connected $C^{2,\alpha}$ domain D , compactly contained in U . $\Omega = U \setminus \overline{D}$ represents the specimen to be inspected, *e.g.*, the cross section of the pipe. We define $\Gamma_e = \partial U$ and $\Gamma_i = \partial D$, so that $\partial\Omega = \Gamma_i \cup \Gamma_e$. Suppose that the inaccessible boundary Γ_i contains some corrosive, connected parts I_s , $s = 1, \dots, m$. The open curves I_s are well-separated, *i.e.*, there is a constant $c_0 > 0$ such that $\text{dist}(I_i, I_j) > c_0$ for $i \neq j$, and the corrosion coefficient, $\gamma_s \in C^1(I_s)$, of each I_s , $s = 1, \dots, m$, satisfies

$$0 \leq \gamma_s \leq C_0 \quad ,$$

for some constant C_0 , and is not identically zero. Let

$$\gamma(x) = \sum_{s=1}^m \gamma_s \chi_s(x), \quad x \in \Gamma_i \quad , \quad (1.1)$$

where χ_s denotes the characteristic function of I_s . A typical shape of Ω would be an annulus. We assume that the curve length of each I_s is small, to be specific,

$$d_0 \epsilon \leq |I_s|_{\Gamma_i} \leq D_0 \epsilon, \quad s = 1, \dots, m \quad , \quad (1.2)$$

where ϵ is a small parameter, representing the common order of magnitude of I_s , and d_0, D_0 are positive constants. Each curve I_s is centered around a fixed point z_s :

$$I_s = \{z \in \Gamma_i : d_{\Gamma_i}(z, z_s) < |I_s|/2\} \quad . \quad (1.3)$$

Here $|\cdot|_{\Gamma_i}$ and $d_{\Gamma_i}(\cdot, \cdot)$ refer to the one-dimensional curve length on Γ_i . The internal voltage potential u_ϵ , generated by a voltage f applied to Γ_e , satisfies

$$\begin{cases} \Delta u_\epsilon = 0 & \text{in } \Omega \quad , \\ -\frac{\partial u_\epsilon}{\partial \nu} + \gamma u_\epsilon = 0 & \text{on } \Gamma_i \quad , \\ u_\epsilon = f & \text{on } \Gamma_e \quad , \end{cases} \quad (1.4)$$

where ν is the **inward** normal to Ω on Γ_i (*i.e.*, the outward normal to D).

A realistic inverse problem in corrosion detection consists of the determination of the corrosion damage on the inaccessible boundary part, Γ_i , when the available data are a finite

number of voltage-to-current pairs on the accessible part Γ_e . The difficulties of this inverse problem stem from its inherent ill-posedness and nonlinearity. Most of the techniques for detecting the corrosion are based on iterative algorithms: least-square-fit algorithms and Newton-type iteration schemes. These methods typically incur tremendous computational costs to get close to the true corrosion damage, but more troubling: they frequently get “stuck” near local minima.

The purpose of this work is to design a direct (non-iterative) technique for electrostatic corrosion detection. This algorithm is of MUSIC-type (multiple signal classification) and is based on an accurate asymptotic representation formula for the steady state boundary currents. In many ways it is closely related to the algorithm developed in [4] for the purpose of detection of small internal defects.

The Cauchy data continuation technique, as studied in the paper by Yang, Choulli, and Cheng [22] for thermal detection of corrosion, is also a direct method. For reasons of instability of the solution of the Cauchy problem it is less accurate than the method described here. Furthermore it makes use of the formula $\gamma = (1/u_\epsilon) \partial u_\epsilon / \partial \nu$ which, in the electrostatic case is somewhat “unfortunate”, due to potential sign changes (and corresponding zeroes) of u_ϵ . A significant (“signal-to noise”) advantage of the present method is that it directly uses the boundary current perturbation data $(\frac{\partial u_\epsilon}{\partial \nu} - \frac{\partial u_0}{\partial \nu})|_{\Gamma_e}$ (as opposed to $\frac{\partial u_\epsilon}{\partial \nu}|_{\Gamma_e}$). Here u_0 is the voltage potential in the absence of corrosion, *i.e.*, the solution to

$$\begin{cases} \Delta u_0 = 0 & \text{in } \Omega \text{ ,} \\ -\frac{\partial u_0}{\partial \nu} = 0 & \text{on } \Gamma_i \text{ ,} \\ u_0 = f & \text{on } \Gamma_e \text{ .} \end{cases} \quad (1.5)$$

The present paper is organized as follows: In Section 2 we review some basic facts about layer potentials. In the following section we establish a representation formula for the unique solution to (1.4). This formula generalizes the formula proved by Kang and Seo in [14, 15]. The main goal of Section 4 is to rigorously derive an asymptotic expansion for $\frac{\partial u_\epsilon}{\partial \nu} - \frac{\partial u_0}{\partial \nu}$ on Γ_e . Section 5 is then devoted to the design of the numerical algorithm, and several computational experiments with it.

2 Layer Potentials

In this section we review some well-known properties of layer potentials for the Laplacian, and establish some useful identities. Let $\Phi(x)$ be the fundamental solution associated with the Laplacian Δ ,

$$\Phi(x) = \frac{1}{2\pi} \ln |x| \text{ , } x \neq 0 \text{ .}$$

Let D be a simply connected, bounded $C^{2,\alpha}$ domain in \mathbb{R}^2 , and let $\Gamma := \partial D$. Let $H^1(D)$ denote the set of functions $f \in L^2(D)$ such that $\nabla f \in L^2(D)$. Furthermore, we define $H^2(D)$ as the space of functions $u \in H^1(D)$ such that $\partial^2 u \in L^2(D)$, and the space $H^{3/2}(D)$ as the interpolation space $[H^1(D), H^2(D)]_{1/2}$. Let $H^{1/2}(\partial D)$ be the set of traces of functions in $H^1(D)$ and $H^{-1/2}(\partial D) = (H^{1/2}(\partial D))^*$. Finally, let $H^1(\partial D)$ denote the set of functions $f \in L^2(\partial D)$ such that $\partial f / \partial \tau \in L^2(\partial D)$, where $\partial / \partial \tau$ is the tangential derivative.

The single-layer and double-layer potentials of a function $\varphi \in L^2(\Gamma)$, denoted $\mathcal{S}_\Gamma[\varphi]$ and $\mathcal{D}_\Gamma[\varphi]$, respectively, are defined by

$$\mathcal{S}_\Gamma[\varphi](x) := \int_\Gamma \Phi(x-y)\varphi(y) d\sigma_y, \quad x \in \mathbb{R}^2, \quad (2.1)$$

$$\mathcal{D}_\Gamma[\varphi](x) := \int_\Gamma \frac{\partial}{\partial \nu_y} \Phi(x-y)\varphi(y) d\sigma_y, \quad x \in \mathbb{R}^2 \setminus \Gamma. \quad (2.2)$$

For a function u , defined in $\mathbb{R}^2 \setminus \Gamma$, we denote

$$\frac{\partial u}{\partial \nu} \Big|_{\pm}(x) := \lim_{t \rightarrow 0^+} \langle \nabla u(x \pm t\nu_x), \nu_x \rangle, \quad x \in \Gamma,$$

if the limit exists. Here ν_x is the outward unit normal to $\Gamma = \partial D$ at x , and $\langle \cdot, \cdot \rangle$ denotes the scalar product in \mathbb{R}^2 . $u|_{\pm}$ are defined likewise.

It is well-known (see for example [11]) that for $\varphi \in L^2(\Gamma)$

$$\frac{\partial(\mathcal{S}_\Gamma[\varphi])}{\partial \nu} \Big|_{\pm}(x) = \left(\pm \frac{1}{2}I + \mathcal{K}_\Gamma^* \right) [\varphi](x), \quad x \in \Gamma, \quad (2.3)$$

$$\mathcal{D}_\Gamma[\varphi]|_{\pm}(x) = \left(\mp \frac{1}{2}I + \mathcal{K}_\Gamma \right) [\varphi](x), \quad x \in \Gamma, \quad (2.4)$$

where \mathcal{K}_Γ is defined by

$$\mathcal{K}_\Gamma[\varphi](x) = \frac{1}{2\pi} \int_\Gamma \frac{\langle y-x, \nu_y \rangle}{|x-y|^2} \varphi(y) d\sigma_y,$$

and \mathcal{K}_Γ^* is the L^2 -adjoint of \mathcal{K}_Γ , i.e.,

$$\mathcal{K}_\Gamma^*[\varphi](x) = \frac{1}{2\pi} \int_\Gamma \frac{\langle x-y, \nu_x \rangle}{|x-y|^2} \varphi(y) d\sigma_y.$$

Since Γ is a $\mathcal{C}^{2,\alpha}$ curve, there is a constant C such that

$$\frac{|\langle x-y, \nu_x \rangle|}{|x-y|^2} \leq C, \quad x, y \in \Gamma. \quad (2.5)$$

Thus it immediately follows that \mathcal{K}_Γ is bounded on $L^p(\Gamma)$ for $1 \leq p \leq \infty$. Moreover, the operator $\varphi \mapsto \nabla_T \mathcal{K}_\Gamma[\varphi]$, where ∇_T denotes the tangential derivative on Γ , is a Calderón-Zygmund operator on Γ which is known to be bounded on $L^p(\Gamma)$ ($1 < p < \infty$), $\mathcal{C}^\beta(\Gamma)$ and $\mathcal{C}^{1,\beta}(\Gamma)$ ($0 < \beta < 1$) (see [19]). Therefore \mathcal{K}_Γ and \mathcal{K}_Γ^* map continuously $L^p(\Gamma)$ into $H^{1,p}(\Gamma)$, $\mathcal{C}^\beta(\Gamma)$ into $\mathcal{C}^{1,\beta}(\Gamma)$ ($0 < \beta < 1$), and $\mathcal{C}^{1,\beta}(\Gamma)$ into $\mathcal{C}^{2,\beta}(\Gamma)$ ($0 < \beta \leq \alpha$) since Γ is $\mathcal{C}^{2,\alpha}$. Here $H^{1,p}(\Gamma)$ is the collection of functions in $L^p(\Gamma)$ whose derivatives also belong to $L^p(\Gamma)$.

Lemma 2.1 *The operators $\frac{1}{2}I + \mathcal{K}_\Gamma$ and $\frac{1}{2}I + \mathcal{K}_\Gamma^*$ are invertible on $L^p(\Gamma)$, $1 < p$, as well as on $\mathcal{C}^k(\Gamma)$ for $k = 1, 2$.*

Proofs of Lemma 2.1 on the spaces $L^2(\Gamma)$ and $\mathcal{C}(\Gamma)$ can be found in [11]. The invertibility on the other spaces can be seen as follows: consider the equation

$$\left(\frac{1}{2}I + \mathcal{K}_\Gamma \right) [u] = f \quad \text{or} \quad \left(\frac{1}{2}I + \mathcal{K}_\Gamma^* \right) [u] = f.$$

If $f \in L^p(\Gamma)$ for $p > 2$ and $u \in L^2(\Gamma)$, then $\mathcal{K}_\Gamma u \in L^p(\Gamma)$ and hence $u \in L^p(\Gamma)$. This shows that $\frac{1}{2}I + \mathcal{K}_\Gamma$ and $\frac{1}{2}I + \mathcal{K}_\Gamma^*$ are invertible on $L^p(\Gamma)$ for $p > 2$. Invertibility on $L^p(\Gamma)$ for $1 < p < 2$ can be seen by duality. Similarly one can prove invertibility on $\mathcal{C}^k(\Gamma)$ for $k = 1, 2$.

In the case where Γ is a circle of radius r

$$\frac{\langle \nu_x, x - y \rangle}{|x - y|^2} = \frac{1}{2r} \quad , \quad x, y \in \Gamma \quad ,$$

and hence

$$\mathcal{K}_\Gamma[\varphi](x) = \mathcal{K}_\Gamma^*[\varphi](x) = \frac{1}{4\pi r} \int_\Gamma \varphi(y) d\sigma_y \quad . \quad (2.6)$$

It now follows from (2.3) and (2.6) that

$$\mathcal{S}_\Gamma[1](x) = \begin{cases} r \ln r & \text{if } |x| \leq r \quad , \\ r \ln |x| & \text{if } |x| \geq r \quad , \end{cases} \quad (2.7)$$

in the case where Γ is a circle of radius r .

Finally, we recall some boundedness properties of single layer potentials for later use. If $p > 1$, there is a constant C_p such that

$$\|\mathcal{S}_\Gamma[\varphi]\|_{L^\infty(\Gamma)} \leq C_p \|\varphi\|_{L^p(\Gamma)} \quad (2.8)$$

for all $\varphi \in L^p(\Gamma)$. In fact, (2.8) can be proved using Hölder's inequality. It is also well-known that

$$\|\mathcal{S}_\Gamma[\varphi]\|_{H^1(D)} \leq C \|\varphi\|_{L^2(\Gamma)} \quad (2.9)$$

for any $\varphi \in L^2(\Gamma)$.

3 Representation Formula

Recall that $\Omega = U \setminus \overline{D}$, where U and D are simply connected, bounded \mathcal{C}^2 domains. Let $\Gamma_e = \partial U$ and $\Gamma_i = \partial D$. For $f \in H^1(\Gamma_e)$, let u_0 be the solution in the absence of corrosion, *i.e.*, the solution to the problem (1.5). We seek to represent u_0 in the following form

$$u_0 = \mathcal{D}_{\Gamma_e}[\varphi_0] + \mathcal{S}_{\Gamma_i}[\psi_0] \quad \text{in } \Omega$$

for some functions $\varphi_0 \in H^1(\Gamma_e)$ and $\psi_0 \in L^2(\Gamma_i)$. The boundary conditions in (1.5) are, due to the relations (2.4) and (2.3), equivalent to

$$\begin{pmatrix} \frac{1}{2}I + \mathcal{K}_{\Gamma_e} & \mathcal{S}_{\Gamma_i} \\ -\frac{\partial}{\partial \nu_i} \mathcal{D}_{\Gamma_e} & -\frac{1}{2}I - \mathcal{K}_{\Gamma_i}^* \end{pmatrix} \begin{pmatrix} \varphi_0 \\ \psi_0 \end{pmatrix} = \begin{pmatrix} f \\ 0 \end{pmatrix} \in H^1(\Gamma_e) \times L^2(\Gamma_i) \quad , \quad (3.1)$$

where ν_e and ν_i denote outward normal ν to Ω on Γ_e and Γ_i (the subscripts e and i emphasize that they are defined on Γ_e and Γ_i , respectively). For example, $\frac{\partial}{\partial \nu_i} \mathcal{D}_{\Gamma_e}[\varphi] = \nu_i \cdot \nabla \mathcal{D}_{\Gamma_e}[\varphi]$.

Lemma 3.1 For $p > 1$, let $X_p := L^p(\Gamma_e) \times L^p(\Gamma_i)$ and

$$A_0 := \begin{pmatrix} \frac{1}{2}I + \mathcal{K}_{\Gamma_e} & \mathcal{S}_{\Gamma_i} \\ -\frac{\partial}{\partial \nu_i} \mathcal{D}_{\Gamma_e} & -\frac{1}{2}I - \mathcal{K}_{\Gamma_i}^* \end{pmatrix}. \quad (3.2)$$

Then A_0 is invertible on X_p , as well as on $H^1(\Gamma_e) \times L^2(\Gamma_i)$.

Proof. Since there is a distance between Γ_e and Γ_i , the operator A_0 is a compact perturbation of

$$\begin{pmatrix} \frac{1}{2}I + \mathcal{K}_{\Gamma_e} & 0 \\ 0 & -\frac{1}{2}I - \mathcal{K}_{\Gamma_i}^* \end{pmatrix},$$

which is known to be invertible on X_p (Lemma 2.1). Therefore, it suffices, by applying the Fredholm alternative, to show that the operator A_0 is injective. Suppose that $(\varphi, \psi) \in X_p$ satisfies

$$A_0 \begin{pmatrix} \varphi \\ \psi \end{pmatrix} = 0.$$

Since $(\frac{1}{2}I + \mathcal{K}_{\Gamma_e})[\varphi] = -\mathcal{S}_{\Gamma_i}[\psi]$ on Γ_e and $\mathcal{S}_{\Gamma_i}[\psi]$ is \mathcal{C}^2 on Γ_e , we obtain that $\varphi \in \mathcal{C}^2(\Gamma_e)$. Likewise it follows that $\psi \in \mathcal{C}^2(\Gamma_i)$. Therefore, the function u defined by $u = \mathcal{D}_{\Gamma_e}[\varphi] + \mathcal{S}_{\Gamma_i}[\psi]$ in Ω is a classical solution to (1.5), with $f = 0$. Since this solution is unique, we get

$$\mathcal{D}_{\Gamma_e}[\varphi] + \mathcal{S}_{\Gamma_i}[\psi] = 0 \quad \text{in } \Omega,$$

and hence

$$\mathcal{D}_{\Gamma_e}[\varphi] + \mathcal{S}_{\Gamma_i}[\psi] = 0 \quad \text{in } U.$$

The ‘‘continuity’’ of the normal derivative of a double-layer potential, in combination with (2.3) now gives

$$\begin{aligned} \psi &= \frac{\partial}{\partial \nu} \mathcal{S}_{\Gamma_i}[\psi] \Big|_+ - \frac{\partial}{\partial \nu} \mathcal{S}_{\Gamma_i}[\psi] \Big|_- \\ &= \frac{\partial}{\partial \nu} (\mathcal{D}_{\Gamma_e}[\varphi] + \mathcal{S}_{\Gamma_i}[\psi]) \Big|_+ - \frac{\partial}{\partial \nu} (\mathcal{D}_{\Gamma_e}[\varphi] + \mathcal{S}_{\Gamma_i}[\psi]) \Big|_- = 0 \quad \text{on } \Gamma_i. \end{aligned}$$

In other words $\psi = 0$, and $\mathcal{D}_{\Gamma_e}[\varphi] = 0$ in U . It follows from (2.4) that $(\frac{1}{2}I + \mathcal{K}_{\Gamma_e})[\varphi] = 0$, and the invertibility of $\frac{1}{2}I + \mathcal{K}_{\Gamma_e}$ now insures that $\varphi = 0$.

Invertibility of A_0 on $H^1(\Gamma_e) \times L^2(\Gamma_i)$ follows easily by the fact that if $(\varphi, \psi) \in L^2(\Gamma_e) \times L^2(\Gamma_i)$ and $A_0 \begin{pmatrix} \varphi \\ \psi \end{pmatrix} \in H^1(\Gamma_e) \times L^2(\Gamma_i)$, then $\mathcal{K}_{\Gamma_e}[\varphi] \in H^1(\Gamma_e)$, and hence $\varphi \in H^1(\Gamma_e)$. This completes the proof. \square

As an immediate consequence of Lemma 3.1, we obtain the following theorem.

Theorem 3.2 The solution u_0 to (1.5) can be represented as

$$u_0 = \mathcal{D}_{\Gamma_e}[\varphi_0] + \mathcal{S}_{\Gamma_i}[\psi_0] \quad \text{in } \Omega,$$

where (φ_0, ψ_0) is the unique solution to (3.1).

Using the exact same techniques we may also derive a representation formula for the unique solution to the boundary value problem

$$\begin{cases} \Delta v_0 = 0 & \text{in } \Omega \text{ ,} \\ -\frac{\partial v_0}{\partial \nu} = g & \text{on } \Gamma_i \text{ ,} \\ v_0 = 0 & \text{on } \Gamma_e \text{ ,} \end{cases} \quad (3.3)$$

where $g \in L^2(\Gamma_i)$. The formula simply reads

$$v_0 = \mathcal{D}_{\Gamma_e}[\tilde{\varphi}_0] + \mathcal{S}_{\Gamma_i}[\tilde{\psi}_0] \quad \text{in } \Omega \text{ ,}$$

where $(\tilde{\varphi}_0, \tilde{\psi}_0)$ is the solution to

$$A_0 \begin{pmatrix} \tilde{\varphi}_0 \\ \tilde{\psi}_0 \end{pmatrix} = \begin{pmatrix} 0 \\ g \end{pmatrix} .$$

We proceed to derive a representation formula for u_ϵ , the solution to (1.4), valid for ϵ sufficiently small. In this case the relevant integral operator is

$$A_\gamma := \begin{pmatrix} \frac{1}{2}I + \mathcal{K}_{\Gamma_e} & \mathcal{S}_{\Gamma_i} \\ -\frac{\partial}{\partial \nu_i} \mathcal{D}_{\Gamma_e} + \gamma \mathcal{D}_{\Gamma_e} & -\frac{1}{2}I - \mathcal{K}_{\Gamma_i}^* + \gamma \mathcal{S}_{\Gamma_i} \end{pmatrix} .$$

Observe that

$$A_\gamma = A_0 + M_\gamma B \text{ ,} \quad (3.4)$$

where M_γ represents multiplication by γ , and

$$B = \begin{pmatrix} 0 & 0 \\ \mathcal{D}_{\Gamma_e} & \mathcal{S}_{\Gamma_i} \end{pmatrix} .$$

Since

$$\|\gamma\|_{L^p(\Gamma_i)} \leq C\epsilon^{1/p} \text{ ,}$$

we have

$$\begin{aligned} \left\| M_\gamma B \begin{pmatrix} \varphi \\ \psi \end{pmatrix} \right\|_{X_p} &\leq \|\gamma\|_{L^p(\Gamma_i)} \|\mathcal{D}_{\Gamma_e}[\varphi] + \mathcal{S}_{\Gamma_i}[\psi]\|_{L^\infty(\Gamma_i)} \\ &\leq C_p \epsilon^{1/p} \left\| \begin{pmatrix} \varphi \\ \psi \end{pmatrix} \right\|_{X_p} \text{ ,} \end{aligned}$$

for any $p > 1$. Here $\|(\varphi, \psi)\|_{X_p} := \|\varphi\|_{L^p(\Gamma_e)} + \|\psi\|_{L^p(\Gamma_i)}$ and C_p is a constant depending only on p . In other words, we have a bound for the operator norm, $\|M_\gamma B A_0^{-1}\|_p$, of $M_\gamma B A_0^{-1}$ on X_p :

$$\|M_\gamma B A_0^{-1}\|_p \leq C_p \epsilon^{1/p} \text{ .} \quad (3.5)$$

Likewise, one can show that

$$\left\| M_\gamma B A_0^{-1} \begin{pmatrix} \varphi \\ \psi \end{pmatrix} \right\|_{H^1(\Gamma_e) \times L^2(\Gamma_i)} \leq C_2 \epsilon^{1/2} \left\| \begin{pmatrix} \varphi \\ \psi \end{pmatrix} \right\|_{H^1(\Gamma_e) \times L^2(\Gamma_i)} \quad (3.6)$$

Since

$$A_\gamma = (I + M_\gamma B A_0^{-1}) A_0 \quad ,$$

we arrive at the following theorem.

Theorem 3.3 *There exists a positive number ϵ_0 such that for any $0 < \epsilon < \epsilon_0$ A_γ is invertible on X_p and on $H^1(\Gamma_e) \times L^2(\Gamma_i)$. Moreover, for any $f \in H^1(\Gamma_e)$, the solution u_ϵ to (1.4) can be represented as*

$$u_\epsilon = \mathcal{D}_{\Gamma_e}[\varphi_\epsilon] + \mathcal{S}_{\Gamma_i}[\psi_\epsilon] \quad \text{in } \Omega \quad ,$$

where $(\varphi_\epsilon, \psi_\epsilon) \in H^1(\Gamma_e) \times L^2(\Gamma_i)$ is the unique solution to

$$A_\gamma \begin{pmatrix} \varphi_\epsilon \\ \psi_\epsilon \end{pmatrix} = \begin{pmatrix} f \\ 0 \end{pmatrix} \quad . \quad (3.7)$$

The last result in this section gives a complete expansion for u_ϵ in terms of iterates easily constructed from u_0 . We can only prove convergence of this expansion for ϵ sufficiently small. First note that

$$A_\gamma^{-1} = A_0^{-1} (I + M_\gamma B A_0^{-1})^{-1} \quad , \quad (3.8)$$

and

$$M_\gamma B \begin{pmatrix} \varphi \\ \psi \end{pmatrix} = \begin{pmatrix} 0 \\ \gamma (\mathcal{D}_{\Gamma_e}[\varphi] + \mathcal{S}_{\Gamma_i}[\psi]) \end{pmatrix} \quad . \quad (3.9)$$

By expanding (3.8) in a Neumann series, we get

$$A_\gamma^{-1} = A_0^{-1} + \sum_{n=1}^{+\infty} (-1)^n A_0^{-1} (M_\gamma B A_0^{-1})^n \quad .$$

Now let

$$\begin{pmatrix} \varphi_n \\ \psi_n \end{pmatrix} := A_0^{-1} (M_\gamma B A_0^{-1})^n \begin{pmatrix} f \\ 0 \end{pmatrix}, \quad n = 0, 1, 2, \dots \quad , \quad (3.10)$$

then

$$\begin{pmatrix} \varphi_\epsilon \\ \psi_\epsilon \end{pmatrix} = A_\gamma^{-1} \begin{pmatrix} f \\ 0 \end{pmatrix} = \sum_{n=0}^{+\infty} (-1)^n \begin{pmatrix} \varphi_n \\ \psi_n \end{pmatrix} \quad . \quad (3.11)$$

Moreover, the following recursive relation holds:

$$\begin{pmatrix} \varphi_{n+1} \\ \psi_{n+1} \end{pmatrix} = A_0^{-1} M_\gamma B \begin{pmatrix} \varphi_n \\ \psi_n \end{pmatrix}, \quad n = 0, 1, 2, \dots \quad . \quad (3.12)$$

It then follows from (3.6) and (3.10) that

$$\|\varphi_n\|_{H^1(\Gamma_e)} + \|\psi_n\|_{L^2(\Gamma_i)} \leq C_0 C^n \epsilon^{n/2} \|f\|_{H^1(\Gamma_e)}, \quad n = 0, 1, 2, \dots \quad , \quad (3.13)$$

for some constant C independent of ϵ . Here C_0 is a bound for A_0^{-1} .

If we define

$$u_n := \mathcal{D}_{\Gamma_e}[\varphi_n] + \mathcal{S}_{\Gamma_i}[\psi_n] \quad \text{in } \Omega, \quad n = 0, 1, 2, \dots, \quad (3.14)$$

which as far as u_0 is concerned is consistent with being the solution to (1.5), then the identity (3.11) immediately yields

$$u_\epsilon = \sum_{n=0}^{+\infty} (-1)^n u_n \quad \text{in } \Omega.$$

It is not difficult to see that this series converges in $H^{3/2}(\Omega)$ (for ϵ sufficiently small). In fact, thanks to elliptic estimates, we have

$$\begin{aligned} \|u_n\|_{H^{3/2}(\Omega)} &\leq C' (\|\mathcal{D}_{\Gamma_e}[\varphi_n]\|_{H^1(\Gamma_e)} + \|\mathcal{S}_{\Gamma_i}[\psi_n]\|_{H^1(\Gamma_i)}) \\ &\leq C' (\|(\frac{1}{2}I + \mathcal{K}_{\Gamma_e})[\varphi_n]\|_{H^1(\Gamma_e)} + \|\mathcal{S}_{\Gamma_i}[\psi_n]\|_{H^1(\Gamma_i)}) \\ &\leq C_1 (\|\varphi_n\|_{H^1(\Gamma_e)} + \|\psi_n\|_{L^2(\Gamma_i)}) \end{aligned}$$

for some constant C' and C_1 , and hence, by (3.13),

$$\|u_n\|_{H^{3/2}(\Omega)} \leq C_0 C_1 C^n \epsilon^{n/2} \|f\|_{H^1(\Gamma_e)}. \quad (3.15)$$

Thus the series for u_ϵ converges strongly in $H^{3/2}(\Omega)$ as long as ϵ is sufficiently small.

By (3.9) and (3.12),

$$\begin{pmatrix} \varphi_n \\ \psi_n \end{pmatrix} = A_0^{-1} \begin{pmatrix} 0 \\ \gamma(\mathcal{D}_{\Gamma_e}[\varphi_{n-1}] + \mathcal{S}_{\Gamma_i}[\psi_{n-1}]) \end{pmatrix} = A_0^{-1} \begin{pmatrix} 0 \\ \gamma u_{n-1}|_{\Gamma_i} \end{pmatrix}, \quad n \geq 1,$$

and so u_n , $n \geq 1$, is the solution to

$$\begin{cases} \Delta u_n = 0 & \text{in } \Omega, \\ -\frac{\partial u_n}{\partial \nu} = \gamma u_{n-1} & \text{on } \Gamma_i, \\ u_n = 0 & \text{on } \Gamma_e. \end{cases} \quad (3.16)$$

Here we have used the observation immediately following Theorem 3.2. In summary we have proven:

Theorem 3.4 *There exists a positive number ϵ_0 such that, for any $0 < \epsilon < \epsilon_0$, and any $f \in H^1(\Gamma_e)$, the solution u_ϵ to (1.4) has the expansion*

$$u_\epsilon = \sum_{n=0}^{+\infty} (-1)^n u_n \quad \text{in } H^{3/2}(\Omega), \quad (3.17)$$

where u_0 is the solution to (1.5) and u_n , $n \geq 1$, are the solutions to (3.16).

4 Asymptotic Expansion

We now derive a formula for the principal term of $\frac{\partial}{\partial \nu}(u_\epsilon - u_0)$ (on Γ_e) as $\epsilon \rightarrow 0$. This formula forms the basis for our corrosion detection algorithm. In order to derive our formula we need an elliptic regularity estimate for the solution to the boundary value problem

$$\begin{cases} \Delta v = 0 & \text{in } \Omega \text{ ,} \\ -\frac{\partial v}{\partial \nu} = \gamma g & \text{on } \Gamma_i \text{ ,} \\ v = 0 & \text{on } \Gamma_e \text{ .} \end{cases} \quad (4.1)$$

In particular we need that

$$\left\| \frac{\partial v}{\partial \nu} \right\|_{\mathcal{C}^1(\Gamma_e)} \leq C_p \|\gamma g\|_{L^p(\Gamma_i)} \text{ ,} \quad (4.2)$$

for any $p > 1$. This regularity estimate is a consequence of well-known Schauder estimates, and could just be stated without proof – however, for completeness we show how it also immediately follows from our integral representation of the solution. We have seen that the solution v to (4.1) can be represented as

$$v = \mathcal{D}_{\Gamma_e}[\varphi] + \mathcal{S}_{\Gamma_i}[\psi] \quad \text{in } \Omega \text{ ,}$$

where (φ, ψ) is the solution to

$$A_0 \begin{pmatrix} \varphi \\ \psi \end{pmatrix} = \begin{pmatrix} 0 \\ \gamma g \end{pmatrix} \text{ ,}$$

i.e. ,

$$\begin{cases} \left(\frac{1}{2}I + \mathcal{K}_{\Gamma_e} \right) [\varphi] + \mathcal{S}_{\Gamma_i}[\psi] = 0 & \text{on } \Gamma_e \text{ ,} \\ -\frac{\partial}{\partial \nu} \mathcal{D}_{\Gamma_e}[\varphi] - \left(\frac{1}{2}I + \mathcal{K}_{\Gamma_i}^* \right) [\psi] = \gamma g & \text{on } \Gamma_i \text{ .} \end{cases} \quad (4.3)$$

Since A_0 is invertible on X_p , we have

$$\|\psi\|_{L^p(\Gamma_i)} \leq C_p \|\gamma g\|_{L^p(\Gamma_i)} \text{ .} \quad (4.4)$$

Let $2\delta = \text{dist}(\Gamma_i, \Gamma_e)$, and define $\Omega_\delta := \{x \in \Omega : \text{dist}(x, \Gamma_e) \leq \delta\}$. Then

$$\|\mathcal{S}_{\Gamma_i}[\psi]\|_{\mathcal{C}^{2,\alpha}(\Omega_\delta)} \leq C \|\psi\|_{L^p(\Gamma_i)} \leq C_p \|\gamma g\|_{L^p(\Gamma_i)} \text{ ,}$$

where constants depend on δ . It now follows from the first equation in (4.3) and Lemma 2.1 that

$$\begin{aligned} \|\varphi\|_{\mathcal{C}^{2,\alpha}(\Gamma_e)} &= \left\| \left(\frac{1}{2}I + \mathcal{K}_{\Gamma_e} \right)^{-1} [(\mathcal{S}_{\Gamma_i}[\psi])|_{\Gamma_i}] \right\|_{\mathcal{C}^{2,\alpha}(\Gamma_e)} \\ &\leq C \|\mathcal{S}_{\Gamma_i}[\psi]\|_{\mathcal{C}^{2,\alpha}(\Gamma_e)} \\ &\leq C_p \|\gamma g\|_{L^p(\Gamma_i)} \text{ .} \end{aligned} \quad (4.5)$$

Note that if $\varphi \in \mathcal{C}^{2,\alpha}(\Gamma_\epsilon)$, then $\mathcal{D}_{\Gamma_\epsilon}[\varphi] \in \mathcal{C}^{2,\alpha}(U)$. In fact, if $\varphi \in \mathcal{C}^{2,\alpha}(\Gamma_\epsilon)$, then $\mathcal{K}_{\Gamma_\epsilon}[\varphi] \in \mathcal{C}^{2,\alpha}(\Gamma_\epsilon)$. Since $u := \mathcal{D}_{\Gamma_\epsilon}[\varphi]$ is the unique solution to the Dirichlet problem $\Delta u = 0$ in U and $u|_{\Gamma_\epsilon} = \frac{1}{2}\varphi + \mathcal{K}_{\Gamma_\epsilon}[\varphi]$, we have $\mathcal{D}_{\Gamma_\epsilon}[\varphi] \in \mathcal{C}^{2,\alpha}(U)$. Therefore, we have in particular

$$\left\| \frac{\partial v}{\partial \nu} \right\|_{\mathcal{C}^1(\Gamma_\epsilon)} \leq \|\mathcal{D}_{\Gamma_\epsilon}[\varphi] + \mathcal{S}_{\Gamma_i}[\psi]\|_{\mathcal{C}^2(\Omega_\delta)} \leq C_p \|\gamma g\|_{L^p(\Gamma_i)} ,$$

which is exactly the estimate (4.2). By a combination of (4.2) with the estimate

$$\|\gamma g\|_{L^p(\Gamma_i)} \leq C\epsilon^{1/p} \|g\|_{L^\infty(\Gamma_i)} ,$$

we obtain the following lemma.

Lemma 4.1 *Let v be the solution to (4.1). Then, for each $p > 1$, there is a constant C_p , independent of g , such that*

$$\left\| \frac{\partial v}{\partial \nu} \right\|_{\mathcal{C}^1(\Gamma_\epsilon)} \leq C_p \epsilon^{1/p} \|g\|_{L^\infty(\Gamma_i)}. \quad (4.6)$$

Based on (3.16), (3.17), and (4.6)

$$\sum_{n=k}^{+\infty} \left\| \frac{\partial u_n}{\partial \nu} \right\|_{\mathcal{C}^1(\Gamma_\epsilon)} \leq C_p \epsilon^{1/p} \sum_{n=k}^{+\infty} \|u_{n-1}\|_{L^\infty(\Gamma_i)} , \quad (4.7)$$

for any $k \geq 1$. From (3.14) we have

$$\|u_{n-1}\|_{L^\infty(\Gamma_i)} = \|\mathcal{D}_{\Gamma_\epsilon}[\varphi_{n-1}] + \mathcal{S}_{\Gamma_i}[\psi_{n-1}]\|_{L^\infty(\Gamma_i)} \leq C_p \left\| \begin{pmatrix} \varphi_{n-1} \\ \psi_{n-1} \end{pmatrix} \right\|_{X_p} ,$$

which in combination with (3.5) and (3.10) yields

$$\|u_{n-1}\|_{L^\infty(\Gamma_i)} \leq C_p \left(C\epsilon^{1/p} \right)^{n-1} \|f\|_{L^p(\Gamma_\epsilon)} .$$

Insertion into (4.7) now gives

$$\sum_{n=k}^{+\infty} \left\| \frac{\partial u_n}{\partial \nu} \right\|_{\mathcal{C}^1(\Gamma_\epsilon)} \leq C_p \epsilon^{1/p} \sum_{n=k}^{+\infty} \left(C\epsilon^{1/p} \right)^{n-1} \|f\|_{L^p(\Gamma_\epsilon)} \leq C_{k,p} \frac{\epsilon^{k/p}}{1 - C\epsilon^{1/p}} \|f\|_{L^p(\Gamma_\epsilon)} ,$$

for any $k \geq 1$, provided $C\epsilon^{1/p} < 1$. The preceding analysis immediately leads to the following theorem.

Theorem 4.2 *For all $p > 1$, there is a constant C_p such that*

$$\left\| \frac{\partial u_\epsilon}{\partial \nu} - \frac{\partial u_0}{\partial \nu} \right\|_{\mathcal{C}^1(\Gamma_\epsilon)} \leq C_p \epsilon^{1/p} \|f\|_{L^p(\Gamma_\epsilon)} . \quad (4.8)$$

We now derive the desired formula for the principal term of $\frac{\partial u_\epsilon}{\partial \nu} - \frac{\partial u_0}{\partial \nu}$. Using the notation from before we easily obtain that

$$\frac{\partial u_\epsilon}{\partial \nu}(x) - \frac{\partial u_0}{\partial \nu}(x) = -\frac{\partial u_1}{\partial \nu}(x) + O(\epsilon^{2/p}) ,$$

uniformly on Γ_e . We thus need to derive a formula for the principal term of $\frac{\partial u_1}{\partial \nu}$. For that purpose, let $G(x, y)$ denote the Green's function for the problem (3.3), *i.e.*, for each $x \in \Omega$, $G(x, y)$ is the solution to

$$\begin{cases} \Delta_y G(x, y) = -\delta_x(y) & \text{in } \Omega \text{ ,} \\ \frac{\partial}{\partial \nu_y} G(x, y) = 0, & y \in \Gamma_i \text{ ,} \\ G(x, y) = 0, & y \in \Gamma_e \text{ .} \end{cases} \quad (4.9)$$

By continuity $G(\cdot, \cdot)$ may be extended to all of $\bar{\Omega} \times \bar{\Omega} \setminus \{(x, y) : x = y\}$, and it satisfies $G(x, y) = G(y, x)$. In terms of G the solution v to (4.1) can be written

$$v(x) = \int_{\Gamma_i} G(x, y) \gamma(y) g(y) d\sigma_y \text{ , } x \in \Omega \text{ .}$$

Remember: ν is the inward normal to Ω on Γ_i . With $g = u_0$ this yields

$$u_1(x) = \int_{\Gamma_i} G(x, y) \gamma(y) u_0(y) d\sigma_y \text{ , } x \in \Omega \text{ ,}$$

and, by taking the normal derivative on Γ_e we now obtain

$$\frac{\partial u_1}{\partial \nu}(x) = \int_{\Gamma_i} \frac{\partial}{\partial \nu_x} G(x, y) \gamma(y) u_0(y) d\sigma_y \text{ , } x \in \Gamma_e \text{ .}$$

Note that $\frac{\partial}{\partial \nu_x} G(x, \cdot)$ is a \mathcal{C}^2 function on Γ_i for any $x \in \Gamma_e$ and that u_0 is also \mathcal{C}^2 on Γ_i since Γ_i is away from Γ_e . Thus we have

$$\frac{\partial}{\partial \nu_x} G(x, y) \gamma(y) u_0(y) = \sum_{s=1}^m \left[\frac{\partial}{\partial \nu_x} G(x, z_s) u_0(z_s) + O(\epsilon) \right] \gamma_s(y) \chi_s(y) \text{ ,}$$

where χ_s is the characteristic function of I_s . We thus finally get

$$\frac{\partial u_1}{\partial \nu}(x) = \sum_{s=1}^m \frac{\partial}{\partial \nu_x} G(x, z_s) u_0(z_s) \int_{I_s} \gamma_s d\sigma + O(\epsilon^2) \text{ .}$$

Define

$$\langle \gamma \rangle_s := \int_{I_s} \gamma_s d\sigma \text{ .}$$

Then we have arrived at the following formula for the principal term of $\frac{\partial}{\partial \nu} u_\epsilon - \frac{\partial}{\partial \nu} u_0$.

Theorem 4.3 *For $x \in \Gamma_e$:*

$$\frac{\partial u_\epsilon}{\partial \nu}(x) - \frac{\partial u_0}{\partial \nu}(x) = - \sum_{s=1}^m \langle \gamma \rangle_s u_0(z_s) \frac{\partial}{\partial \nu_x} G(x, z_s) + O(\epsilon^{1+\beta}) \quad (4.10)$$

for any $0 < \beta < 1$.

5 MUSIC type algorithm for reconstruction

We now apply the asymptotic formula (4.10) to design a MUSIC (Multiple Signal Classification) algorithm for locating small internal corrosive parts from boundary measurements. MUSIC is generally used in signal processing problems as a method for estimating the individual frequencies from multiple-harmonic signals [20]. In the present context the MUSIC algorithm uses information about the spectral structure of the (exterior Dirichlet-to-Neumann) boundary map for the Laplace operator with the “internal” corrosion (Robin) boundary condition. The eigenvectors corresponding to the most significant eigenvalues span a “signal subspace” in the sense that they contain nearly all the information about the corrosive parts that can be extracted from the boundary map. The remaining eigenvectors span a sort of “noise subspace”. In order to be of any practical interest our algorithm uses spectral information about a particular discrete version of the boundary map.

It is worth mentioning that the present algorithm is also related to the linear sampling method of Colton and Kirsch [9]. We refer to Cheney [7], and Kirsch [18], for detailed discussions of the connection between MUSIC-type algorithms and the linear sampling method.

Define the (exterior Dirichlet-to-Neumann) map Λ_γ from $H^{1/2}(\Gamma_e)$ into $H^{-1/2}(\Gamma_e)$ by

$$\Lambda_\gamma(f) = \frac{\partial u_\epsilon}{\partial \nu} \Big|_{\Gamma_e}, \quad (5.1)$$

where u_ϵ is the solution to (1.4). Let Λ_0 be the exterior Dirichlet-to-Neumann map for the case when no corrosion is present. We proceed to show how partial spectral information about $\Lambda_\gamma - \Lambda_0$ may effectively be used to determine the corroded internal boundary parts. The estimate (4.8) immediately gives that

$$\|(\Lambda_\gamma - \Lambda_0)(f)\|_{C^1(\Gamma_e)} \leq C_p \epsilon^{1/p} \|f\|_{L^2(\Gamma_e)}, \quad (5.2)$$

for any $1 < p \leq 2$. In particular this shows that the operator $\Lambda_\gamma - \Lambda_0$, originally defined on $H^{1/2}(\Gamma_e)$, can be extended as an operator on all of $L^2(\Gamma_e)$. This operator is compact when regarded as an operator from $L^2(\Gamma_e)$ into itself. We also note that for $g \in H^{1/2}(\Gamma_e)$

$$\int_{\Gamma_e} (\Lambda_\gamma - \Lambda_0)(f) g \, d\sigma = \int_{\Omega} [\nabla u_\epsilon \cdot \nabla v_\epsilon - \nabla u_0 \cdot \nabla v_0] \, dx + \int_{\Gamma_i} \gamma u_\epsilon v_\epsilon \, d\sigma \quad (5.3)$$

where v_ϵ and v_0 are solutions to (1.4) and (1.5) with f replaced by g , respectively. Note that

$$\int_{\Omega} \nabla u_\epsilon \cdot \nabla u_0 \, dx = \int_{\Gamma_e} f \frac{\partial u_0}{\partial \nu} \, d\sigma = \int_{\Omega} \nabla u_0 \cdot \nabla u_0 \, dx.$$

Thus we obtain by substituting f in place of g in (5.3)

$$\int_{\Gamma_e} (\Lambda_\gamma - \Lambda_0)f f \, d\sigma = \int_{\Omega} |\nabla(u_\epsilon - u_0)|^2 \, dx + \int_{\Gamma_i} \gamma |u_\epsilon|^2 \, d\sigma.$$

In summary we have established the following lemma.

Lemma 5.1 $\Lambda_\gamma - \Lambda_0$ is self-adjoint, positive, and compact on $L^2(\Gamma_e)$.

The identity (4.10) shows that

$$(\Lambda_\gamma - \Lambda_0)(f)(x) = - \sum_{s=1}^m \langle \gamma \rangle_s u_0(z_s) \frac{\partial}{\partial \nu_x} G(x, z_s) + O(\epsilon^{1+\beta}), \quad x \in \Gamma_e, \quad (5.4)$$

for any $0 < \beta < 1$. It is not difficult to see that the remainder term $O(\epsilon^{1+\beta})$ is actually bounded by $C_\beta \epsilon^{1+\beta} \|f\|_{L^2(\Gamma_e)}$ with C_β independent of f . Define the operator T on $H^{1/2}(\Gamma_e)$ by

$$(Tf)(x) = - \sum_{s=1}^m \langle \gamma \rangle_s u_0(z_s) \frac{\partial}{\partial \nu_x} G(x, z_s), \quad x \in \Gamma_e . \quad (5.5)$$

Since u_0 depends linearly on f , T is a linear operator and can be extended as a bounded operator on $L^2(\Gamma_e)$. Furthermore the estimate we have just verified may easily be sharpened to show that

$$\Lambda_\gamma - \Lambda_0 = T + O(\epsilon^{1+\beta}) , \quad (5.6)$$

in the operator norm on $L^2(\Gamma_e)$.

Lemma 5.2 *The operator T is a compact, self-adjoint, positive semi-definite operator on $L^2(\Gamma_e)$.*

If $\langle \gamma \rangle_s = c_s \epsilon$ for some fixed constants c_s , then this follows immediately by multiplication of (5.4) by ϵ^{-1} , consideration of the $\epsilon \rightarrow 0$ limit, and use of Lemma 5.1. Here it is essential that the remainder term is bounded by $C_\beta \epsilon^{1+\beta} \|f\|_{L^2(\Gamma_e)}$, *i.e.*, that we have norm convergence. More generally the lemma can also be verified by noting that T is a finite-dimensional operator, and that u_0 is given by

$$u_0(y) = - \int_{\Gamma_e} \frac{\partial}{\partial \nu_x} G(x, y) f(x) d\sigma(x), \quad y \in \Omega , \quad (5.7)$$

so that

$$\int_{\Gamma_e} T(f)h = \sum_{s=1}^m \langle \gamma \rangle_s u_0(z_s) v_0(z_s) ,$$

where v_0 is the solution to (1.5), with f replaced by h .

We introduce the linear operator $\mathcal{G} : L^2(\Gamma_e) \rightarrow \mathbb{R}^m$

$$\mathcal{G}f = (u_0(z_1), \dots, u_0(z_m)) , \quad (5.8)$$

where u_0 , as before, is the solution to (1.5). If \mathbb{R}^m is endowed with the standard Euclidean inner product then

$$\langle \mathcal{G}f, a \rangle = \sum_{s=1}^m a_s u_0(z_s) = - \int_{\Gamma_e} \left(\sum_{s=1}^m a_s \frac{\partial}{\partial \nu_x} G(x, z_s) \right) f(x) d\sigma(x) ,$$

for arbitrary $a = (a_1, \dots, a_m) \in \mathbb{R}^m$. Therefore, the adjoint operator $\mathcal{G}^* : \mathbb{R}^m \rightarrow L^2(\Gamma_e)$ is given by

$$\mathcal{G}^* a = - \sum_{s=1}^m a_s \frac{\partial}{\partial \nu_x} G(\cdot, z_s) . \quad (5.9)$$

Exactly as in Brühl *et al.* [4] there is a simple factorization of the operator T , and a corresponding characterization of its range.

Lemma 5.3

- (i) \mathcal{G}^* is injective;
- (ii) \mathcal{G} is surjective;
- (iii) $T = \mathcal{G}^* \mathcal{M} \mathcal{G}$, where

$$\mathcal{M}a = \left(\langle \gamma \rangle_1 a_1, \dots, \langle \gamma \rangle_m a_m \right), \quad a = (a_1, \dots, a_m) \in \mathbb{R}^m;$$

- (iv) $\text{Range}(T) = \text{span} \left\{ \frac{\partial}{\partial \nu_x} G(x, z_s); s = 1, \dots, m \right\}$.

Proof. Suppose $\mathcal{G}^* a = 0$, and let w denote the function $w(x) = \sum_{s=1}^m a_s G(x, z_s)$. Then $w = \frac{\partial w}{\partial \nu} = 0$ on Γ_e , and therefore by unique continuation $w(x) = \sum_{s=1}^m a_s G(x, z_s) = 0$ in all of Ω . It immediately follows that $a_s = 0$ for $s = 1, \dots, m$. Assertion (ii) follows from (i) and the well-known relation between the range and the null space of adjoint finite-dimensional operators: $\text{Range}(\mathcal{G}) = \text{Ker}(\mathcal{G}^*)^\perp$. Using (5.5), (5.8), and (5.9), it is easy to see that (iii) holds. Now, according to (iii) we have $\text{Range}(T) = \text{Range}(\mathcal{G}^* \mathcal{M} \mathcal{G}) = \text{Range}(\mathcal{G}^*)$, since \mathcal{M} and \mathcal{G} are surjective. This yields (iv), and the proof is complete. \square

The following result forms the basis for our algorithm to identify the midpoint locations, z_s , of the small corroded boundary parts.

Theorem 5.4 *A point $z \in \Gamma_i$ belongs to the set $\{z_s : s = 1, \dots, m\}$ if and only if*

$$\frac{\partial}{\partial \nu_x} G(\cdot, z) \Big|_{\Gamma_e} \in \text{Range}(T) . \quad (5.10)$$

Proof. Assume $\frac{\partial}{\partial \nu_x} G(\cdot, z) \Big|_{\Gamma_e} \in \text{Range}(T)$. As a consequence of (iv) of Lemma 5.3 there exist coefficients a_s , $1 \leq s \leq m$ such that

$$\frac{\partial}{\partial \nu_x} G(x, z) = \sum_{s=1}^m a_s \frac{\partial}{\partial \nu_x} G(x, z_s) , \quad \text{for } x \in \Gamma_e . \quad (5.11)$$

Since $G(x, z) = 0$ for any $x \in \Gamma_e$ and $z \in \Gamma_i$, we also have

$$G(x, z) = \sum_{s=1}^m a_s G(x, z_s) , \quad \text{for } x \in \Gamma_e , \quad (5.12)$$

and by unique continuation it follows that $G(x, z) = \sum_{s=1}^m a_s G(x, z_s)$ for all $x \in \Omega$. Due to the singularity of $G(x, z)$ at $x = z$ this is only possible if $z \in \{z_s : s = 1, \dots, m\}$, and so we have established the sufficiency of the condition (5.10). The necessity follows immediately from (iv) of Lemma 5.3. \square

The finite-dimensional, self-adjoint operator T has a spectral decomposition

$$T = \sum_{p=1}^m \lambda_p v_p \otimes v_p , \quad \|v_p\|_{L^2(\Gamma_e)} = 1 ,$$

where λ_p , $1 \leq p \leq m$, are the (positive) eigenvalues of T , and v_p is the corresponding eigenfunction. Note that the rank of T , m , is exactly the number of corroded parts. We

assume that $\lambda_1 \geq \lambda_2 \geq \dots \geq \lambda_m > 0$. Let $P_k : L_0^2(\Gamma_e) \rightarrow \text{span} \{v_1, \dots, v_k\}$, $k = 1, \dots, m$, be the orthogonal projector $P_k = \sum_{p=1}^k v_p \otimes v_p$. It then follows from Theorem 5.4 that

$$z \in \{z_s : s = 1, \dots, m\} \quad \text{iff} \quad (I - P_m) \left(\frac{\partial}{\partial \nu_x} G(\cdot, z) \Big|_{\Gamma_e} \right) = 0 \quad ,$$

or equivalently,

$$z \in \{z_s : s = 1, \dots, m\} \quad \text{iff} \quad \cot \theta(z) = +\infty, \quad (5.13)$$

where the angle $\theta(z) \in [0, \pi/2)$ is defined by

$$\cot \theta(z) = \frac{\left\| P_m \left(\frac{\partial}{\partial \nu_x} G(\cdot, z) \Big|_{\Gamma_e} \right) \right\|_{L^2(\Gamma_e)}}{\left\| (I - P_m) \left(\frac{\partial}{\partial \nu_x} G(\cdot, z) \Big|_{\Gamma_e} \right) \right\|_{L^2(\Gamma_e)}} \quad . \quad (5.14)$$

Since $\Lambda_\gamma - \Lambda_0$ is a self-adjoint, positive, and compact operator on $L^2(\Gamma_e)$, it admits a spectral decomposition

$$\Lambda_\gamma - \Lambda_0 = \sum_{p=1}^{+\infty} \lambda_p^\epsilon v_p^\epsilon \otimes v_p^\epsilon \quad , \quad \|v_p^\epsilon\|_{L^2(\Gamma_e)} = 1 \quad , \quad (5.15)$$

with $\lambda_1^\epsilon \geq \lambda_2^\epsilon \geq \dots \geq \lambda_m^\epsilon \geq \dots > 0$. Let $P_k^\epsilon : L^2(\Gamma_e) \rightarrow \text{span} \{v_1^\epsilon, \dots, v_k^\epsilon\}$, $k = 1, 2, \dots$, be the orthogonal projector $P_k^\epsilon = \sum_{p=1}^k v_p^\epsilon \otimes v_p^\epsilon$.

The data for our identification algorithm consist of (an appropriate approximation to) the spectral decomposition of $\Lambda_\gamma - \Lambda_0$. Motivated by (5.10), we seek to find those points $z \in \Gamma_i$ such that

$$\frac{\partial}{\partial \nu_x} G(\cdot, z) \Big|_{\Gamma_e} \in \text{Range} (\Lambda_\gamma - \Lambda_0) \quad . \quad (5.16)$$

This condition is fulfilled approximately by $z \in \Gamma_i$ if the following quantity is sufficiently large:

$$\cot \theta_k^\epsilon(z) := \frac{\left\| P_k^\epsilon \left(\frac{\partial}{\partial \nu_x} G(\cdot, z) \Big|_{\Gamma_e} \right) \right\|_{L^2(\Gamma_e)}}{\left\| (I - P_k^\epsilon) \left(\frac{\partial}{\partial \nu_x} G(\cdot, z) \Big|_{\Gamma_e} \right) \right\|_{L^2(\Gamma_e)}} \quad . \quad (5.17)$$

Since we don't *a priori* know m , we plot the angles $\theta_k^\epsilon(z)$ (for $k = 1, 2, \dots$) as a function of z . When the plots "stabilize", *i.e.*, don't significantly change by increasing k , we have a good estimate for m , and we interpret locations where small angles are attained as good candidates for the z_s .

Once the locations z_s are approximately found, we can also estimate the integrated corrosion coefficients $\langle \gamma \rangle_s$. Our procedure for doing this is the following. Define w_s , $s = 1, \dots, m$, to be the solution to

$$\begin{cases} \Delta w_s = 0 & \text{in } \Omega \quad , \\ -\frac{\partial w_s}{\partial \nu} = 0 & \text{on } \Gamma_i \quad , \\ w_s(x) = \frac{\partial}{\partial \nu_x} G(x, z_s) & x \in \Gamma_e \quad . \end{cases} \quad (5.18)$$

It then follows from (5.4) and (5.15) that

$$\begin{aligned} -\sum_{s=1}^m \langle \gamma \rangle_s w_{s'}(z_s) \frac{\partial}{\partial \nu_x} G(\cdot, z_s) &\approx (\Lambda_\gamma - \Lambda_0) \left(\frac{\partial}{\partial \nu_x} G(\cdot, z_{s'}) \right) \\ &\approx \sum_{p=1}^m \lambda_p^\epsilon \left\langle v_p^\epsilon, \frac{\partial}{\partial \nu_x} G(\cdot, z_{s'}) \right\rangle v_p^\epsilon . \end{aligned}$$

By integrating both sides of the above formula against $v_{s'}$, we obtain

$$-\sum_{s=1}^m w_{s'}(z_s) \left\langle v_{s'}^\epsilon, \frac{\partial}{\partial \nu_x} G(\cdot, z_s) \right\rangle \langle \gamma \rangle_s \approx \lambda_{s'}^\epsilon \left\langle v_{s'}^\epsilon, \frac{\partial}{\partial \nu_x} G(\cdot, z_{s'}) \right\rangle, \quad (5.19)$$

for $s' = 1, \dots, m$. Therefore the values of $\langle \gamma \rangle_s$, $s = 1, \dots, m$, can be calculated by solving a linear system. For the simplest case, $m = 1$, the formula reads

$$\langle \gamma \rangle_1 \approx -\frac{\lambda_1^\epsilon}{w_1(z_1)} .$$

5.1 Numerical Results

In this section we present some numerical experiments with the MUSIC-type algorithm we have designed to find the internally corroded parts, $I_s \subset \mathbb{R}^2$, $s = 1, 2, \dots, m$. The domain $\Omega \subset \mathbb{R}^2$ is taken to be an annulus, *i.e.*, $\Omega = U \setminus \overline{D}$, where U and D are disks centered at $(0, 0)$ with radii r_e and r_i , respectively. The measurements used for the identification represent a discrete analogue of the Dirichlet-to-Neumann map $\Lambda_\gamma - \Lambda_0$ on Γ_e . In order to find a good approximation from the “identification point of view” the following observation will be useful.

Lemma 5.5 *The functions $\frac{\partial G}{\partial \nu_x}(\cdot, y)$, $y \in \Gamma_i$ span a dense subspace of $L^2(\Gamma_e)$.*

Proof. The solution to the boundary value problem (1.5), for a given $f \in L^2(\Gamma_e)$, is given by

$$u_0(y) = -\int_{\Gamma_e} \frac{\partial}{\partial \nu_x} G(x, y) f(x) d\sigma(x) .$$

If $f \in L^2(\Gamma_e)$ were orthogonal to $\{ \frac{\partial G}{\partial \nu_x}(\cdot, y) : y \in \Gamma_i \}$ it thus follows that

$$\frac{\partial u_0}{\partial \nu} = u_0 = 0 \quad \text{on } \Gamma_i .$$

By unique continuation this implies that $u_0 = 0$ in all of Ω , and so consequently $f = 0$. \square

With this observation in mind we approximate $\Lambda_\gamma - \Lambda_0$ by using Dirichlet data of the form $f(\cdot) = \frac{\partial G}{\partial \nu_x}(\cdot, y)$, $y \in \Gamma_i$, and solving the integral equations (3.1) and (3.7). Note that G is a quite particular fundamental solution, namely the solution to (4.9). Practically it is very important that these normal derivatives can be calculated efficiently. To understand how, let \tilde{G} be the Green’s function for U , that is, for each $x \in U$, let $\tilde{G}(x, y)$ be the solution to

$$\begin{cases} \Delta_y \tilde{G}(x, y) = -\delta_x(y) & \text{in } U, \\ \tilde{G}(x, y) = 0 & y \in \Gamma_e. \end{cases} \quad (5.20)$$

Since U is a disk, $\tilde{G}(x, y)$ has a simple explicit expression. In fact, \tilde{G} is given by

$$\tilde{G}(x, y) = -\frac{1}{2\pi} \ln|x-y| + \frac{1}{2\pi} \ln \left| \frac{r_e}{|x|} x - \frac{|x|}{r_e} y \right|. \quad (5.21)$$

For any $y \in \Gamma_i$ we may now compute $\frac{\partial G}{\partial \nu_x}(\cdot, y)$ on Γ_e by solving the integral equation (5.22) of the following lemma.

Lemma 5.6 *For any $y \in \Gamma_i$, let $G_y(x) := G(x, y)$, and $\tilde{G}_y(x) := \tilde{G}(x, y)$ for $x \in \Gamma_e$. Then*

$$\left(I + 2\mathcal{K}_{\Gamma_e}^* - 4 \frac{\partial \mathcal{S}_{\Gamma_i}}{\partial \nu_e} \frac{\partial \mathcal{S}_{\Gamma_e}}{\partial \nu_i} \right) \left[\frac{\partial G_y}{\partial \nu_e} \right] (x) = 2 \frac{\partial \tilde{G}_y}{\partial \nu_e}(x), \quad x \in \Gamma_e, \quad (5.22)$$

where ν_e and ν_i denote outward normal vectors to U on Γ_e , and to D on Γ_i , respectively.

Proof. Fix $y \in \Omega$ and let $u(x) = (G_y - \tilde{G}_y)(x)$ for $x \in \Omega$. Because of the symmetry of G and \tilde{G} , u satisfies $\Delta u = 0$ in Ω , $u = 0$ on Γ_e , and $\frac{\partial u}{\partial \nu_i} = -\frac{\partial \tilde{G}_y}{\partial \nu_i}$ on Γ_i . Thus by the Green theorem, we have

$$u(x) = -\mathcal{S}_{\Gamma_e} \left[\frac{\partial u}{\partial \nu_e} \right] (x) + \mathcal{D}_{\Gamma_i}[u](x) + \mathcal{S}_{\Gamma_i} \left[\frac{\partial u}{\partial \nu_i} \right] (x), \quad x \in \Omega .$$

However, in view of (2.3), we have

$$\mathcal{D}_{\Gamma_i}[u](x) = \mathcal{S}_{\Gamma_i} \left(\frac{1}{2}I + \mathcal{K}_{\Gamma_i}^* \right)^{-1} \left[\frac{\partial}{\partial \nu_i} \mathcal{D}_{\Gamma_i}[u] \Big|_{\Gamma_i} \right] (x)$$

for all $x \in \mathbb{R}^2 \setminus \bar{D}$. Therefore there is a density function ψ with $\int_{\Gamma_i} \psi = 0$ such that

$$(G_y - \tilde{G}_y)(x) = -\mathcal{S}_{\Gamma_e} \left[\frac{\partial G_y}{\partial \nu_e} - \frac{\partial \tilde{G}_y}{\partial \nu_e} \right] (x) + \mathcal{S}_{\Gamma_i}[\psi](x), \quad x \in \Omega . \quad (5.23)$$

Similarly to (3.1) we now obtain from (5.23)

$$\left(\frac{1}{2}I + \mathcal{K}_{\Gamma_i}^* \right) [\psi] = -\frac{\partial \tilde{G}_y}{\partial \nu_i} + \frac{\partial \mathcal{S}_{\Gamma_e}}{\partial \nu_i} \left[\frac{\partial G_y}{\partial \nu_e} - \frac{\partial \tilde{G}_y}{\partial \nu_e} \right] \quad \text{on } \Gamma_i .$$

Since Γ_i is a circle, $\mathcal{K}_{\Gamma_i}^*[\varphi] = 0$ when $\int_{\Gamma_i} \varphi = 0$ (cf. (2.6)) and hence

$$\psi = -2 \frac{\partial \tilde{G}_y}{\partial \nu_i} + 2 \frac{\partial \mathcal{S}_{\Gamma_e}}{\partial \nu_i} \left[\frac{\partial G_y}{\partial \nu_e} - \frac{\partial \tilde{G}_y}{\partial \nu_e} \right] .$$

Since $\tilde{G}_y(x) + \mathcal{S}_{\Gamma_e} \left[\frac{\partial \tilde{G}_y}{\partial \nu_e} \right] (x) = -\Phi(x, y)$ for $x \in U$ and $y \in U$ (cf. Equation (2.23) of [1]), it follows that

$$\psi = 2 \frac{\partial \mathcal{S}_{\Gamma_e}}{\partial \nu_i} \left[\frac{\partial G_y}{\partial \nu_e} \right] + 2 \frac{\partial \Phi_y}{\partial \nu_i} . \quad (5.24)$$

By substituting (5.24) into (5.23) and taking normal derivatives (on Γ_e) of both sides of (5.23), we obtain

$$\left(\frac{1}{2}I + \mathcal{K}_{\Gamma_e}^* - 2 \frac{\partial \mathcal{S}_{\Gamma_i}}{\partial \nu_e} \frac{\partial \mathcal{S}_{\Gamma_e}}{\partial \nu_i} \right) \left[\frac{\partial G_y}{\partial \nu_e} \right] = \left(\frac{1}{2}I + \mathcal{K}_{\Gamma_e}^* \right) \left[\frac{\partial \tilde{G}_y}{\partial \nu_e} \right] + 2 \frac{\partial \mathcal{S}_{\Gamma_i}}{\partial \nu_e} \left[\frac{\partial \Phi_y}{\partial \nu_i} \right] \quad \text{on } \Gamma_e ,$$

for any fixed $y \in \Omega$. We claim that in the limit as y approaches Γ_i from inside Ω

$$\lim \left(\left(\frac{1}{2}I + \mathcal{K}_{\Gamma_e}^* \right) \left[\frac{\partial \tilde{G}_y}{\partial \nu_e} \right] + 2 \frac{\partial \mathcal{S}_{\Gamma_i}}{\partial \nu_e} \left[\frac{\partial \Phi_y}{\partial \nu_i} \right] \right) = \frac{\partial \tilde{G}_y}{\partial \nu_e} , \quad (5.25)$$

and (5.22) immediately follows. It remains to verify (5.25). First of all, by (2.6),

$$\begin{aligned} \mathcal{K}_{\Gamma_e}^* \left[\frac{\partial \tilde{G}_y}{\partial \nu_e} \right] &= \frac{1}{4\pi r_e} \int_{\Gamma_e} \frac{\partial \tilde{G}_y}{\partial \nu_e}(x) ds \\ &= \frac{1}{4\pi r_e} \int_U \Delta \tilde{G}_y(x) dx = -\frac{1}{4\pi r_e} , \end{aligned}$$

for any $y \in \Omega$. Secondly, in the limit as y approaches Γ_i from inside Ω , we have

$$\lim \mathcal{S}_{\Gamma_i} \left[\frac{\partial \Phi_y}{\partial \nu_i} \right] (x) = -\frac{1}{4\pi} [\ln|x-y| - \ln|x|] , \quad \text{for } |x| > r_i . \quad (5.26)$$

In fact, we have

$$\mathcal{S}_{\Gamma_i} \left[\frac{\partial \Phi_y}{\partial \nu_i} \right] (x) = \frac{1}{4\pi^2} \int_{\Gamma_i} \ln|x-z| \frac{\partial}{\partial \nu_z} \ln|z-y| ds(z) = \mathcal{D}_{\Gamma_i}[\Phi_x](y) ,$$

and hence from (2.4) and (2.6)

$$\begin{aligned} \lim \mathcal{S}_{\Gamma_i} \left[\frac{\partial \Phi_y}{\partial \nu_i} \right] (x) &= -\frac{1}{2} \Phi_x(y) + \frac{1}{4\pi r_i} \int_{\Gamma_i} \Phi_x(y) d\sigma(y) \\ &= -\frac{1}{4\pi} [\ln|x-y| - \ln|x|] , \end{aligned}$$

which is the desired formula. Note that the last equality holds because of (2.7).

A similar statement holds for any derivative with respect to x , $|x| > r_i$. In view of (5.21), it is now a matter of straight-forward computation to check that in the limit as y approaches Γ_i from inside Ω

$$\lim \left(\mathcal{K}_{\Gamma_e}^* \left[\frac{\partial \tilde{G}_y}{\partial \nu_e} \right] + 2 \frac{\partial \mathcal{S}_{\Gamma_i}}{\partial \nu_e} \left[\frac{\partial \Phi_y}{\partial \nu_i} \right] \right) = \frac{1}{2} \frac{\partial \tilde{G}_y}{\partial \nu_e} \quad \text{on } \Gamma_e ,$$

and we obtain (5.25). This completes the proof. \square

For computation, we discretize Γ_e and Γ_i by choosing points

$$\{r_e(\cos \theta_n, \sin \theta_n) | \theta_n = 2\pi(n-1)/N, n = 0, 1, \dots, N-1\}$$

and

$$\{r_i(\cos \theta_n, \sin \theta_n) | \theta_n = 2\pi(n-1)/N, n = 0, 1, \dots, N-1\},$$

with $N = 256$. Put $x_n := r_e(\cos \theta_n, \sin \theta_n)$ and $y_n := r_i(\cos \theta_n, \sin \theta_n)$ for $n = 0, 1, \dots, 255$. The discrete data we use for our numerical algorithm consist of the 256×256 matrix $((\Lambda_\gamma - \Lambda_0)(\frac{\partial G}{\partial \nu_x}(\cdot, y_n))(x_m))_{m,n=1}^{256}$, or rather its SVD (singular value decomposition). For comparison we also calculate the 256×256 matrix $(T(\frac{\partial G}{\partial \nu}(\cdot, y_n))(x_m))_{m,n=1}^{256}$ and its SVD. Notice that none of these two matrices are symmetric.

Figure 1 displays the singular values of our discrete versions of $\Lambda_\gamma - \Lambda_0$ and T when there are two internally corroded parts. The singular values of our approximation to T exhibit a significant drop, to machine epsilon, after the first two (consistent with the fact that the range of T is two-dimensional) there is also a drop in the singular values of our approximation to $\Lambda_\gamma - \Lambda_0$, but it is of course not nearly as significant due to the discrepancy between $\Lambda_\gamma - \Lambda_0$ and T , and due to the very approximate nature of our discrete data. It is nonetheless still possible to estimate the location of the corroded parts very accurately from the SVD of our discrete version of $\Lambda_\gamma - \Lambda_0$.

Using the first k (left-)singular vectors of our data, we plot approximate values of $\theta_k^\epsilon(z) = \cot^{-1} L_k^\epsilon(z)$ where

$$L_k^\epsilon(z) := \frac{\|P_k^\epsilon(\frac{\partial G(\cdot, z)}{\partial \nu})\|_{L^2(\Gamma_\epsilon)}}{\|(I - P_k^\epsilon)(\frac{\partial G(\cdot, z)}{\partial \nu})\|_{L^2(\Gamma_\epsilon)}}. \quad (5.27)$$

We start this process with just one singular vector ($k = 1$) and continue with increasing k until the plot of $\theta_k^\epsilon(z)$ stabilizes. Values near zero correspond to locations that are likely to be close to the I_s (or z_s).

Example 1. In this example, the outer radius $r_e = 1$ and the inner one $r_i = 0.8$ and there are two corroded parts. The top part of Figure 1 shows the locations of the corrosion and the approximate SVD of $\Lambda_\gamma - \Lambda_0$ and T . The bottom of Figure 1 clearly shows that there are two corroded parts, since the minima of θ_k^ϵ stabilize already at $k = 2$. The minima of θ_2^ϵ are the locations of the corroded parts. Since we use the same 256 points for the calculation and plotting of θ_k^ϵ and for the potential location of the points z_s , the fact that $|z_s - z_s^c| = 0$ really asserts that z_s is accurately estimated up to the spacing between two adjacent mesh points, which is approximately 0.0172. By solving (5.19), we compute the integrated corrosion coefficients: $\langle \gamma \rangle_1^c = 0.3846$, $\langle \gamma \rangle_2^c = 0.1660$. The actual, integrated coefficients are $\langle \gamma \rangle_1 = 0.3927$, $\langle \gamma \rangle_2 = 0.1767$.

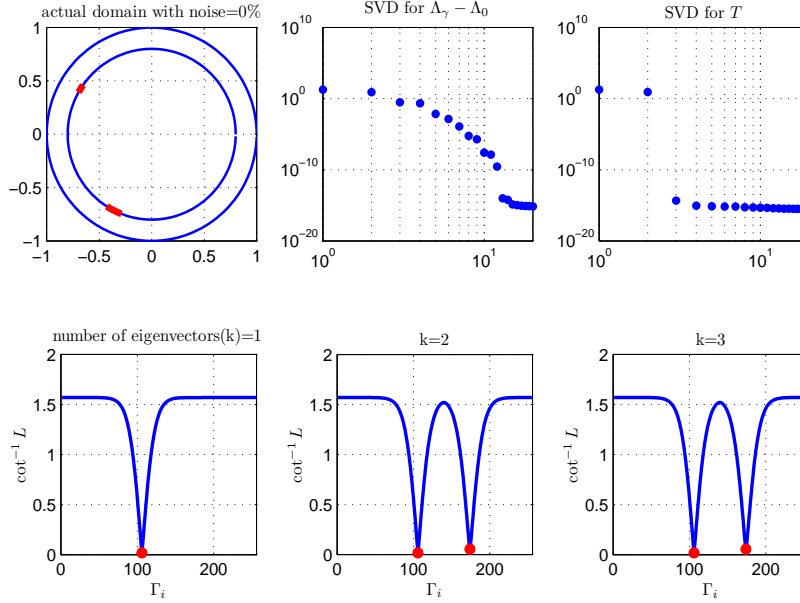
Example 2. In this example we consider a case with five corroded parts. The actual data of the configuration is summarized in the top of Table 1. The corroded parts are numbered counter-clock-wise starting with the smallest positive angle. Note that the first three corroded parts have low corrosion coefficients while the remaining two have relatively high ones. The computational results with 1% noise are summarized in Table 1 and Figure 2. It is interesting to note that the corroded parts 1 and 2 which are close to each other and have low corrosion coefficients, are detected as a single one. The actual midpoint locations of corroded parts 1 and 2 are $(0.0686, 0.6966)$ and $(-0.1197, 0.6897)$, and the single detected location is $(-0.1027, 0.6924)$. On the other hand, the corroded parts 4 and 5, which have high corrosion coefficients, are clearly detected already with $k = 3$.

actual data			
m	γ_s	z_s	$\langle \gamma \rangle_s$
5	0.1	(0.0686, 0.6966)	0.0120
	0.5	(-0.1197, 0.6897)	0.0430
	0.3	(-0.4307, 0.5518)	0.0206
	1.0	(0.2519, - 0.6531)	0.1203
	1.2	(0.5723, - 0.4031)	0.1031

detected data				
m^c	z_s^c	$\langle \gamma \rangle_s^c$	$ z_s - z_s^c $	$ \langle \gamma \rangle_s - \langle \gamma \rangle_s^c $
4	(-0.1027, 0.6924)	0.0514	0.0172	0.0084
	(-0.4170, 0.5622)	0.0212	0.0172	0.0006
	(0.2519, - 0.6531)	0.1181	0	0.0021
	(0.5723, - 0.4031)	0.1026	0	0.0005

result for each step					
number of sing. vectors(k)	m^c	z_s^c	$ z_s - z_s^c $	$\langle \gamma \rangle_s^c$	$ \langle \gamma \rangle_s - \langle \gamma \rangle_s^c $
1	1	(0.3889, - 0.5820)		0.1554	
2	2	(-0.1366, + 0.6865)		0.0579	
		(0.3889, - 0.5820)		0.1554	
3	3	(-0.1366, + 0.6865)		0.0579	
		(0.2519, - 0.6531)		0.1181	
		(0.5723, - 0.4031)		0.1026	
4		(-0.0686, + 0.6966)	0.0515	0.0486	0.0056
		(-0.4031, + 0.5723)	0.0344	0.0237	0.0030
		(0.2519, - 0.6531)	0	0.1181	0.0021
		(0.5723, - 0.4031)	0	0.1026	0.0005
5	4	(-0.1027, + 0.6924)	0.0172	0.0514	0.0084
		(-0.4170, + 0.5622)	0.0172	0.0212	0.0006
		(0.2519, - 0.6531)	0	0.1181	0.0021
		(0.5723, - 0.4031)	0	0.1026	0.0005
6	4	(-0.1027, + 0.6924)	0.0172	0.0514	0.0084
		(-0.4170, + 0.5622)	0.0172	0.0212	0.0006
		(0.2519, - 0.6531)	0	0.1181	0.0021
		(0.5723, - 0.4031)	0	0.1026	0.0005

Table 1: Summary of computational results with five corroded parts, and 1% noise. Two corroded parts with low corrosion coefficient are detected as a single one.



The results for steps

number of sing. vectors(k)	m^c	z_s^c	$ z_s - z_s^c $	$\langle \gamma \rangle_s^c$	$ \langle \gamma \rangle_s - \langle \gamma \rangle_s^c $
1	1	(-0.6759, 0.4280)		0.3846	
2	2	(-0.6759, 0.4280)	0	0.3846	0.0081
		(-0.3597, -0.7146)	0	0.1660	0.0107
3	2	(-0.6759, 0.4280)	0	0.3846	0.0081
		(-0.3597, -0.7146)	0	0.1660	0.0107

Figure 1: Pipe with two internal corroded parts, and “approximate” singular values of $\Lambda_\gamma - \Lambda_0$ and T in upper figures. The bottom figure shows θ_k^c for three different values of k . In the table, m^c is the number of estimated corroded parts, z_s and z_s^c are the actual midpoint locations and the estimated midpoint locations, and $\langle \gamma \rangle_s$, $\langle \gamma \rangle_s^c$ are the actual and the estimated, integrated, corrosion coefficients, for a particular k .

Example 3. Figure 3 shows the computational results with various degree of noise. The information about the actual locations and corrosion coefficients is summarized in Table 2. Observe that the first two corroded parts have low corrosion coefficients. They are detected as a single corroded part with 5% noise. The other two corroded parts, which have high corrosion coefficients, are detected very well even with a high ratio of noise.

6 Conclusion

We have designed a non-iterative algorithm of MUSIC-type for detecting small, internal, corroded boundary parts based on external, electrostatic boundary measurements. Our

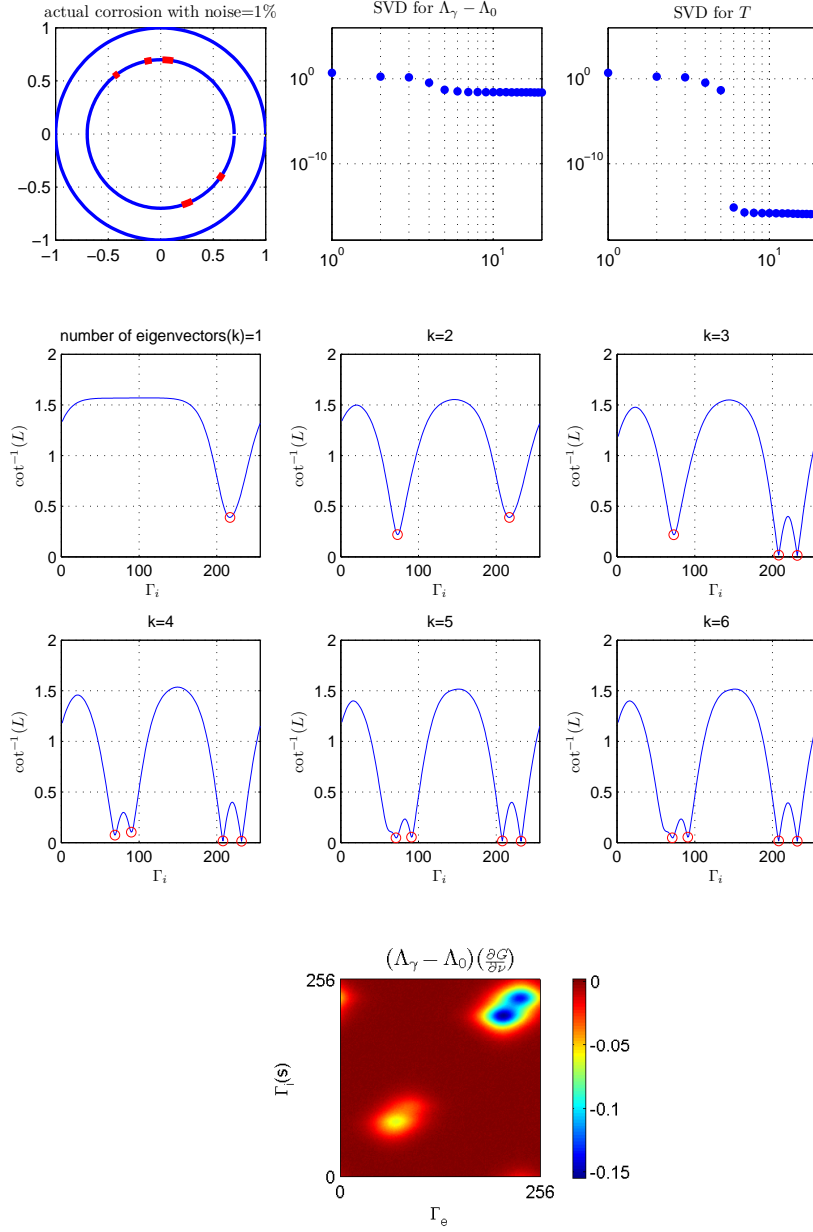


Figure 2: The computational results with 1% noise. The last plot shows the values of $(\Lambda_\gamma - \Lambda_0)\left(\frac{\partial G}{\partial \nu}\right)$ on $\Gamma_e \times \Gamma_i$.

method relies on an asymptotic representation formula for the steady state current perturbations caused by small, internal corroded boundary parts. We have performed numerical

m	γ_s	z_s	$\langle \gamma \rangle_s$
4	0.1	(0.5657, + 0.5657)	0.0137
	0.05	(0.3597, + 0.7146)	0.0049
	3.0	(0.6307, - 0.4922)	0.2945
	3.0	(0.7464, - 0.2879)	0.2945

Table 2: Pipe with four corrosive parts.

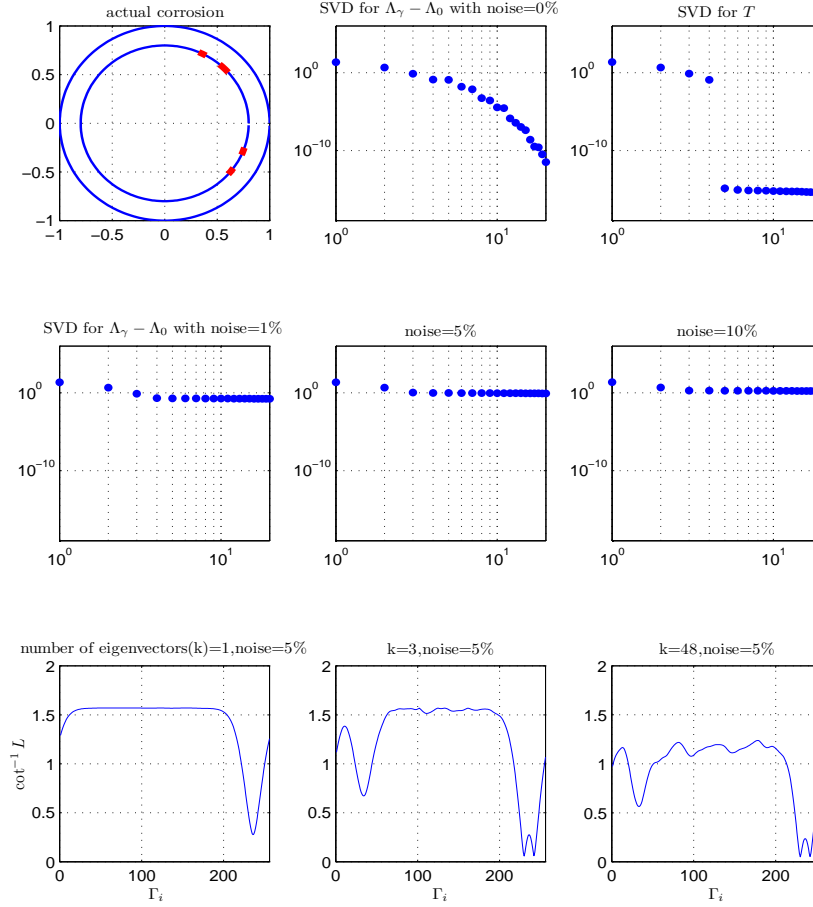


Figure 3: Computational results with varying degree of noise

experiments to test the viability of the algorithm using a discrete version of $\Lambda_\gamma - \Lambda_0$ as data. The numerical tests clearly demonstrate that the algorithm works well even in the presence of relatively high noise ratios, provided the corrosion coefficients are sufficiently large. The latter restriction is to be expected, since only those corroded parts with reasonably sized

corrosion coefficients can cause sufficient changes in currents. Furthermore, it is practically most important to detect the corroded parts with the highest coefficients, since they are likely to “develop” more rapidly, and therefore potentially cause the most serious damage (to the pipe). It should be mentioned that our algorithm cannot identify separately the size of the corroded parts (without assuming that we know the corrosion coefficients). We can only reconstruct the integral $\langle \gamma \rangle_s = \int_{I_s} \gamma \, d\sigma$ for $s = 1, \dots, m$. It is likely that a higher-order asymptotic expansion of the boundary perturbations would yield formulas that permit such identification.

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