

# SMOS INSTRUMENT PERFORMANCE AND CALIBRATION AFTER 5 YEARS IN ORBIT

M. Martín-Neira (1), R. Oliva (2), I. Corbella (3), F. Torres (3), N. Duffo (3), I. Durán (3), J. Kainulainen (4), J. Closa (5), A. Zurita (5), F. Cabot (6), A. Khazaal (6), E. Anterrieu (7), J. Barbosa (8), G. Lopes (8), J. Tenerelli (9), R. Díez-García (10), J. Fauste (2), F. Martín-Porqueras (11), V. González-Gambau (12), A. Turiel (12), S. Delwart (13), R. Crapolicchio (13), M. Suess (1)

## Corresponding author

M. Martín-Neira and M. Suess are with ESA-ESTEC, Keplerlaan 1, 2200 AG Noordwijk, The Netherlands. (e-mail: [manuel.Martín-Neira@esa.int](mailto:manuel.Martín-Neira@esa.int))

## Authors' affiliations

- (1) European Space Agency at ESTEC, Noordwijk, The Netherlands.
- (2) European Space Agency at ESAC, Villanueva de la Cañada, Spain.
- (3) Polytechnic University of Catalonia (UPC), Barcelona, Spain.
- (4) Harp Technologies Ltd., Espoo, Finland.
- (5) EADS-CASA Espacio, Madrid, Spain.
- (6) Centre d'Etudes Spatiales de la Biosphère (CESBIO), Toulouse, France.
- (7) Institut de Recherche en Astrophysique et Planétologie (IRAP), Toulouse, France.
- (8) DEIMOS, Lisbon, Portugal.
- (9) OceanDataLab, Brest, France.
- (10) IDEAS, ESAC, Villanueva de la Cañada, Spain.
- (11) VEGA Telespazio, ESAC, Villanueva de la Cañada, Spain.
- (12) SMOS Barcelona Expert Centre, Barcelona, Spain.
- (13) European Space Agency at ESRIN, Frascati, Italy.

## Keywords

Soil Moisture and Ocean Salinity (SMOS) Mission, soil moisture, sea surface salinity, L-Band radiometry, aperture synthesis, MIRAS

40 **1 Abstract**

41  
42 ESA's Soil Moisture and Ocean Salinity (SMOS) mission, launched 2-Nov-2009, has been in orbit for over 5  
43 years, and its Microwave Imaging Radiometer with Aperture Synthesis (MIRAS) in two dimensions keeps  
44 working well. The calibration strategy remains overall as established after the commissioning phase, with a few  
45 improvements. The data for this whole period has been reprocessed with a new fully polarimetric version of the  
46 Level-1 processor which includes a refined calibration schema for the antenna losses. This reprocessing has  
47 allowed the assessment of an improved performance benchmark. An overview of the results and the progress  
48 achieved in both calibration and image reconstruction is presented in this contribution.

49 **2 INTRODUCTION**

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51 With an experience of over 5 years of in-orbit operation, much has been learnt on how MIRAS works  
52 and how its observations can be improved through better calibration and image reconstruction  
53 techniques. The purpose of this paper is to update the reader with the latest results on the payload  
54 performance and data processing of the SMOS mission (Mecklenburg et al., 2012). SMOS is currently  
55 delivering several products, some of them used by operational systems, others only for scientific  
56 research (Mecklenburg et al., in press). MIRAS is a Microwave Imaging Radiometer with two-  
57 dimensional Aperture Synthesis, which remains being the first and so far, the only one of its kind, in  
58 space. The main feature of MIRAS is that it obtains two-dimensional images at every snapshot without  
59 needing any mechanical scanning of its antenna, a very distinct capability when compared with  
60 traditional scanners or push-broom radiometers. A detailed description of the instrumental aspects of  
61 MIRAS can be found in (McMullan et al., 2008) while the on-board Calibration System and respective  
62 in-flight calibration strategy are described in (Brown et al., 2008) and (Martín-Neira et al., 2008). One  
63 year after launch the calibration approach was slightly modified with the initial flight experience, and  
64 the first SMOS instrument in-orbit performance was assessed in (Oliva et al., 2013), including the  
65 effect of the unexpectedly severe Radio Frequency Interference from ground transmitters (Oliva et al.,  
66 2012). The present paper will then follow the same structure as (Oliva et al., 2013), with important

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67 additions brought by the accumulated experience of over 5 years: Section 3 provides an overview of the  
68 main sources of error and the current mitigation strategies used to overcome them; Section 4  
69 summarizes the current status of calibration activities, including all latest modifications to the initial  
70 calibration plan; Section 5 presents the in-orbit behaviour of the most critical instrument parameters;  
71 Section 6 gives the performance obtained with the latest version of the Level-1 processor, through the  
72 spatial and temporal analysis of brightness temperature images, and finally, Section 7 includes a view  
73 on the current investigations that should lead to the next version of the Level-1 processor with a hint on  
74 the expected improvements.

75 It is worth mentioning that, at the time of the writing of this paper, the running version of the  
76 operational SMOS Level-1 data processor is V620, that a new version, V700, has been delivered and is  
77 under assessment, and that the entire data record of the SMOS mission (from January 2010 onwards)  
78 has been reprocessed with V620 and is available to the whole SMOS user community.

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## 80 **3 ERROR SOURCES AND MITIGATION TECHNIQUES**

### 81 **3.1 Error Sources**

82 Different error sources cause different effects on the SMOS brightness temperature images. Therefore  
83 in this section the error sources will be presented according to the effect they produce in the images.

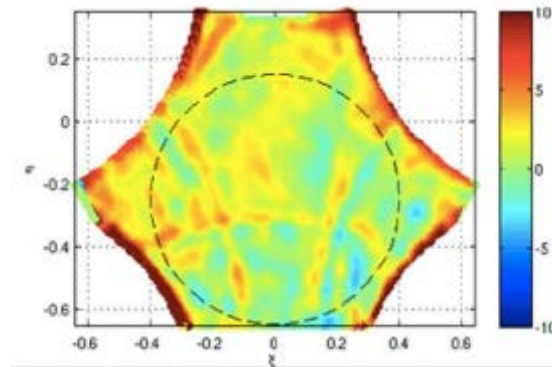
#### 84 **3.1.1 Systematic Spatial Ripple**

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86 Figure 1 presents the deviation, with respect to a forward model, of an image of the brightness  
87 temperature measured by SMOS over a portion of the South-Eastern Pacific Ocean in X-polarization  
88 (X-polarization refers to the image formed with the signal collected by the horizontal probe of MIRAS  
89 antenna elements). The comparison is performed after averaging many snapshots so that random errors  
90 induced by the radiometric resolution can be neglected, and only systematic errors remain. The most  
91 prominent features of such deviation image are a +0.96 K bias and a 1.5 K rms spatial ripple, both  
92 statistics evaluated within the dashed circle shown in Figure 1. Similar statistics can be computed for

93 the Y-polarization as well as for the Stokes-3 and Stokes-4 parameters, obtaining, in general, different  
94 values for the different parameters, values which, in turn, depend on the particular image reconstruction  
95 approach being applied, that is, on the Level-1 processor version. Furthermore, and although it is not as  
96 easy to show as with measurements of the relatively uniform ocean, bias and ripples also appear in  
97 images acquired over any region of the Earth, be it land, ice or coastlines, and over the Cold Sky,  
98 exhibiting a magnitude which is scene-dependent. Interpreting the bias as a spatial ripple of an infinite  
99 spatial wavelength, both bias and spatial ripple shall be understood as comprised within the ‘spatial  
100 ripple’ referred to in what follows.

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104 Figure 1: Example of bias and spatial ripples of SMOS images when compared to a **radiative transfer**  
105 **ocean model. The axes are the director cosines and the colour scale is in Kelvin.**

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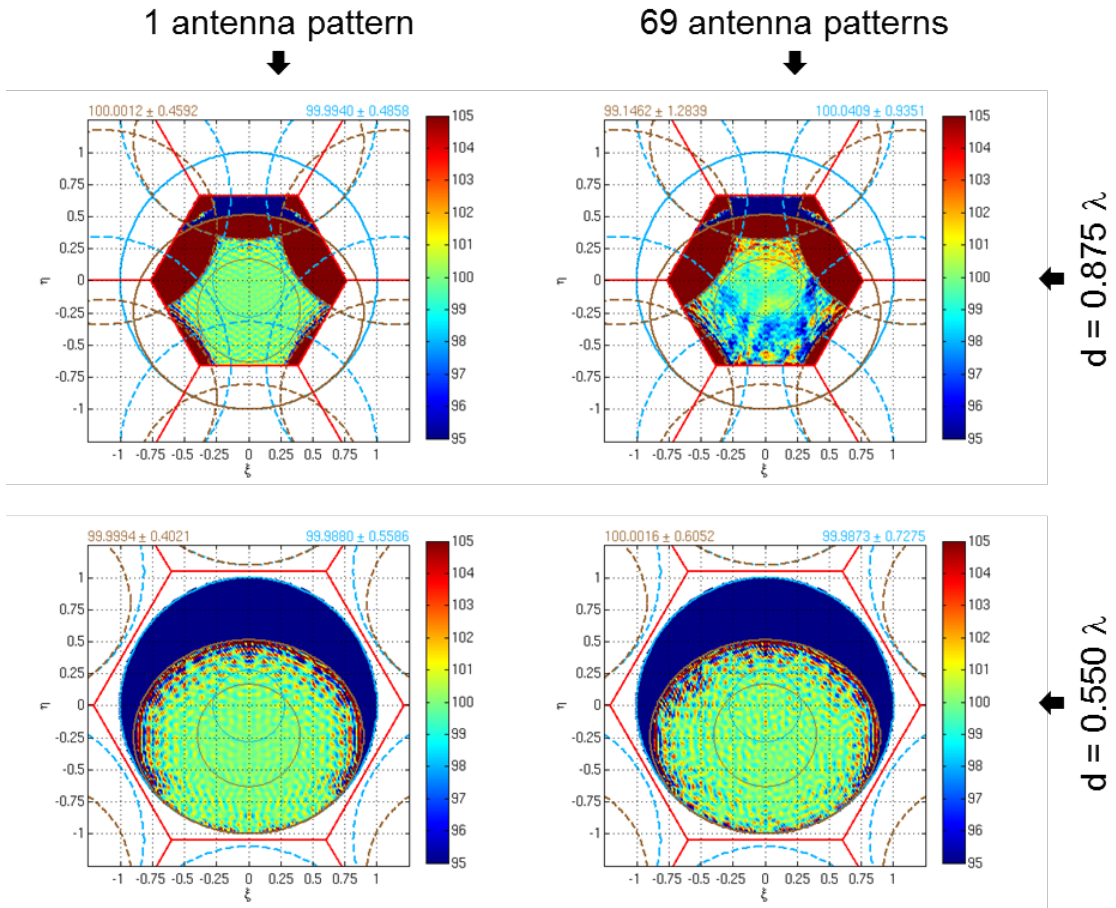
107 The cause and existence of the systematic spatial ripple underlying all SMOS images was already  
108 studied before SMOS launch (Camps et al., 2005)(Anterrieu, 2007). Thanks to the investigations  
109 conducted during the last 5 years in several directions, using flight data, an important conclusion has  
110 been consolidated: the spatial ripple results mostly from the combination of having different antenna  
111 element patterns and imaging in alias conditions (that is, using a spatial sampling which leads to having  
112 aliased images in some parts of the real space). This is illustrated in Figure 2, which shows a similar  
113 deviation image to Figure 1 obtained simulating different conditions: the left and right columns assume  
114 identical and different –perfectly known– antenna patterns respectively, while the rows correspond to  
115 different antenna element spacings, the one of MIRAS ( $0.875\lambda$ ) in the top, and another one which is  
116 alias-free ( $0.55\lambda$ ) in the bottom. As it is evident, the spatial ripple appears only in the top right plot, that

117 is, for alias condition and different antenna patterns. If there is no aliasing, or/and if the antenna  
118 patterns are all identical, then there is no significant spatial ripple.

119 As a consequence of this, two further results have now been well established: first, even in the ideal  
120 case where the antenna pattern of every antenna element were known perfectly well, a non-negligible  
121 systematic spatial error would still be present in SMOS images, dubbed ‘noise floor’; second, the part  
122 of the scene outside the alias-free area does have an impact on the scene recovered in the alias-free  
123 region, or in other words, the spatial ripple in the alias-free area depends on the scene in the aliased  
124 regions (Corbella et al., 2014).

125 The first result might be the most striking one, since for long, it had been expected that the accurate  
126 knowledge of the antenna patterns, very carefully characterized on ground and used in the image  
127 reconstruction, would have enabled the acquisition of images with negligible ripple. It also emphasizes  
128 the need of having the interferometric instrument with either an alias-free element spacing or as similar  
129 antenna patterns as possible, to suppress the spatial ripple from the images. The second result, on the  
130 other hand, has been the basis to build ripple reduction methods to improve the current SMOS images,  
131 as will be seen in the section devoted to mitigation techniques below. Finally the contribution to the  
132 spatial ripple due to the limited knowledge of the antenna patterns and residual calibration errors shall  
133 not be forgotten.

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Figure 2: Illustration that the spatial ripple results from the combination of different antenna patterns and alias condition. **The image shows the Earth and the sky at a spatially uniform but different brightness temperature, viewed with the nominal SMOS geometry. The axes are the director cosines and the colour bar represents the retrieved brightness temperature in Kelvin.**

### 142 3.1.2 Sun and RFI Tails

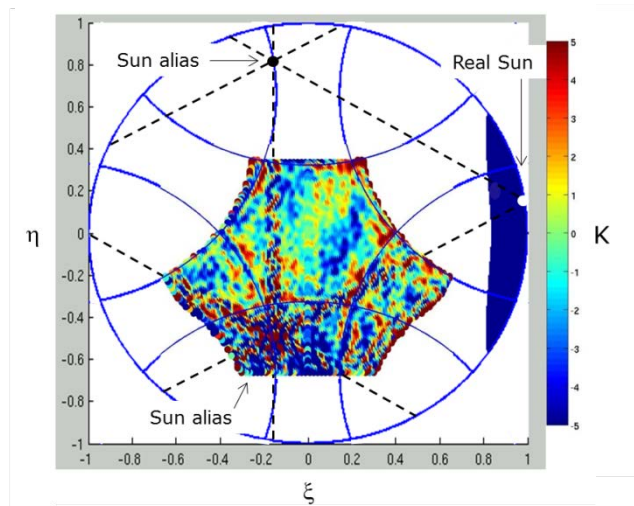
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144 Figure 3 is a deviation image as Figure 1 but acquired at a time when the Sun is in front of the antenna.

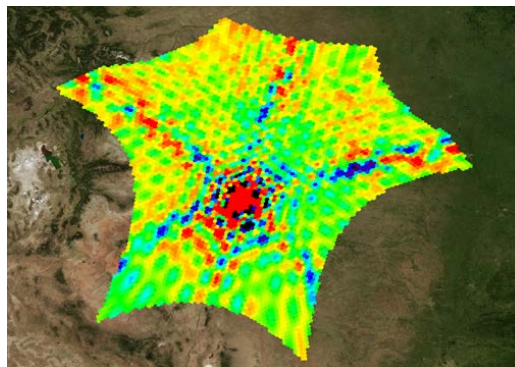
145 The Sun can be seen as a white circle towards the right side of the unity circle, which represents the  
146 front hemisphere of the antenna. The area shaded in blue is the locus of possible positions of the Sun as  
147 seen from the SMOS orbit around the year. The Sun locus is sufficiently far away from the extended  
148 field of view of SMOS that it would not cause any ripples if it were not because of the aliasing and the  
149 side lobes. Indeed, a replica of the Sun is clearly visible inside the field of view. Moreover, the real Sun  
150 and its replica appear connected by lines of side lobes which cross the entire field of view, including  
151 the alias free region. Extending the lines of side lobes through the black dashed lines it is possible to

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152 located a second alias in the upper part of the unity circle, outside the extended alias-free field of view.  
153 The cause of these side lobes is the same as that of the spatial ripple: the differences across antenna  
154 patterns enhance the side lobes joining the Sun aliases, which are generated by the element spacing.  
155 Patterns of side lobes are also excited by Radio Frequency Interference (RFI) transmitters, which  
156 behave within the image reconstruction process, much in the same way as the Sun. This is illustrated in  
157 Figure 4, where the extended alias-free field of view is projected on ground at a location of an RFI  
158 source: an hexagonal pattern of side lobes is clearly visible. RFI sources remain being an important  
159 error source in SMOS (Oliva et al., in press).



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Figure 3: Sun tails and alias affecting an ocean image



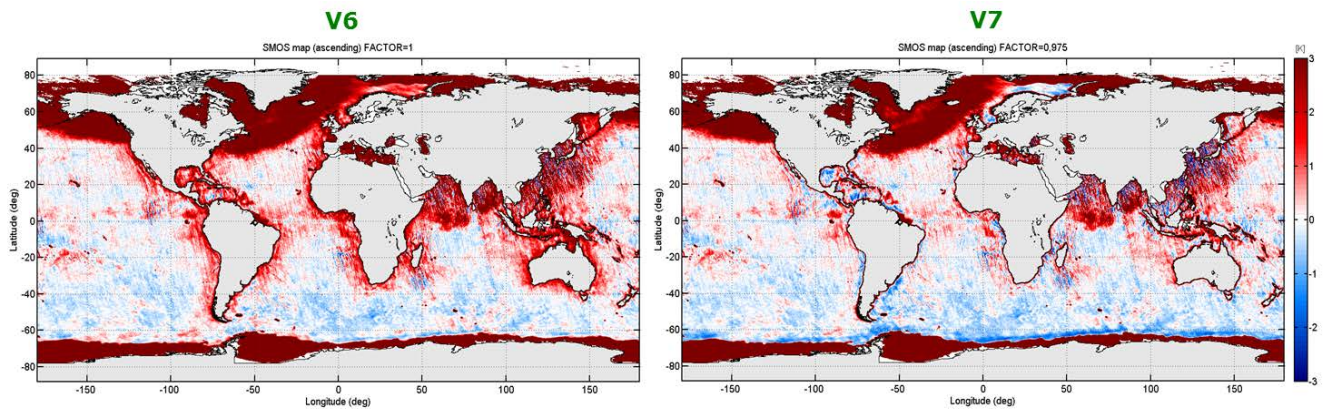
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Figure 4: Hexagonal pattern of side lobes excited by a strong Radio Frequency Interference source

### 168 3.1.3 Land-Sea Contamination

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170 The left plot of Figure 5 shows a global view of the oceans with the accumulated deviation of SMOS  
171 measurements from a forward model. Brightness temperature warmer than the model are **seen** around  
172 all continental masses which wrongly lead to fresher water retrievals. This feature of SMOS images is  
173 referred to as ‘land-sea contamination’ and is of concern among the Sea Surface Salinity community  
174 because it can extend several hundreds of kilometers into the open ocean. Much effort has been devoted  
175 to understand the reason behind this land-sea contamination. The current understanding is that it is  
176 caused by several contributions. The most important one seems to be a calibration error, of the order of  
177 2%, in the amplitude of the correlations, the so-called  $G_{kj}$  correlator efficiency coefficients (Corbella et  
178 al., in press). Although this problem with the correlation coefficients has been identified, to date, the  
179 root cause has not been found yet and the search within the instrumental details related to it continues.  
180 The second contributor is the spatial ripple described above, generated by the warmer land and  
181 extending into the ocean.

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184 **Figure 5: Stokes-1/2 residual images against a radiative transfer ocean model using present**  
185 **correlation efficiency factors (left), where the land-sea contamination is clearly observed around**  
186 **the continental masses, and using corrected values (right), with significantly reduced errors. The**  
187 **warm areas around Alaska, Greenland, Arabian Sea, Gulf of Bengal and Sea of China are due to**  
188 **Radio Frequency Interference, and the blue rim around Antarctica is due to un-modelled sea ice.**  
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### 190 3.1.4 Seasonal Variations

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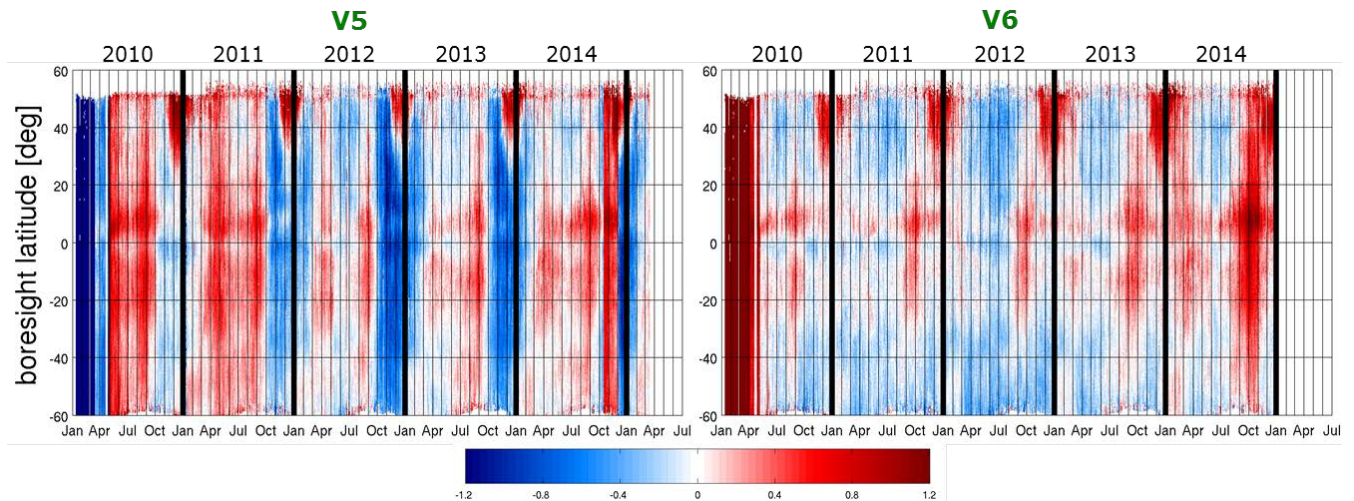
192 The right plot of Figure 6 shows the current deviation of the Stokes-1/2 parameter over the Pacific  
193 Ocean, averaged within the alias-free field of view, with respect to an ocean model, along the  
194 descending passes. The plot spans 5 years, from 2010 till 2014, and from **60° South to 60° North** in



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195 latitude, with a brightness temperature scale of  $\pm 1.2$  K. This **Hovmöller** plot constitutes a powerful tool  
196 to analyze any seasonal (and latitudinal) variations. Besides the red stripe during the Commissioning  
197 Phase in early 2010, the variations are contained within  $\pm 0.4$  K except for the eclipse periods (**mid-**  
198 **November to mid-February**) and a region around October where some warm signatures are observed.  
199 During an eclipse, the Sun, which can be as high as  $31^\circ$  above the antenna horizon, is suddenly  
200 occulted by the Earth. The antenna skin temperature falls down a couple of tens of degrees, from  
201 around  $28^\circ\text{C}$  to some  $5^\circ\text{C}$ . When the satellite exits the Earth shadow, the Sun warms up the antenna  
202 again until it reaches the temperature it would have if there had been no eclipse. This post-eclipse  
203 transient causes, in every descending orbit, a warm anomaly in the brightness temperature which  
204 extends to latitudes as low as  $30^\circ$  North. On the other hand, the reason for the October warm anomaly  
205 has not yet been uncovered. Attempts to correlate it with residual galactic noise or other geophysical  
206 signatures have failed and hence, an instrumental origin should be assumed. Furthermore this anomaly  
207 seems more intense in 2014.

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Figure 6: Latitude-Time **Hovmöller** plot of the descending pass Stokes-1/2 deviation from an ocean forward model, averaged in the alias-free area, with V505 (left) and V620 (right) Level 1 processor versions. **Colour bar is given in Kelvin.**

### 214 3.1.5 Orbital Variations

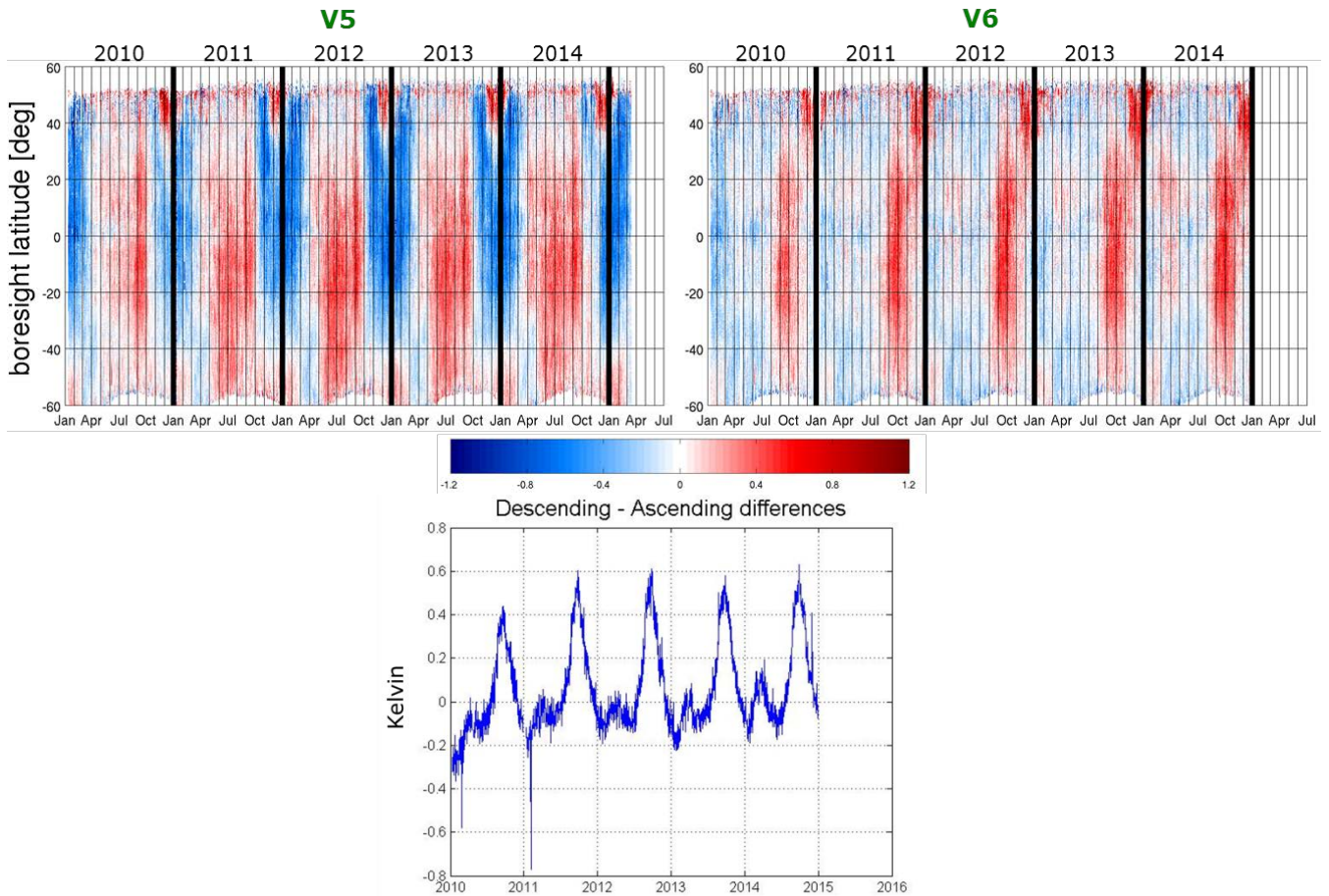
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216 The **top** right plot of Figure 7 shows the current difference between descending and ascending passes  
217 of the Stokes-1/2 parameter over the Pacific Ocean, averaged within the alias-free field of view. **The**  
218 **Stokes-1/2 parameters of each pass are first corrected for significant forward model**  
219 **contributions which range from tens of Kelvin (sea surface physical temperature and emissivity)**  
220 **to just a few Kelvin (reflected cold sky, galactic glint, direct and reflected Sun, atmospheric up**  
221 **and down welling). This correction is needed to separate geophysical from instrumental effects as**  
222 **much as possible.** The plot span in latitude and time, as well as the brightness temperature scale, are  
223 the same as those of Figure 6. The difference between **the two** passes, **shown in the bottom panel, is**  
224 within **0.8** K peak-to-peak, the maximum departure, of about +0.5 K, happening around October,  
225 **which is the period when the galactic glint is the strongest. There is evidence that such deviation**  
226 **in October is very likely related to a mis-modeling of the galactic glint.** It is worth mentioning that  
227 the descending minus ascending difference does not show any particular increase in 2014 as seen in the  
228 descending pass of Figure 6, meaning that both passes follow each other's variations also in 2014. **The**  
229 **top right panel of Figure 7 also shows a clear impact of the eclipse period in latitudes above 30°**  
230 **North.**

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Figure 7: Descending-minus-Ascending pass Stokes-1/2 parameter over the Pacific Ocean, averaged in the alias-free area, with V505 (top left) and V620 (top right) Level 1 processor versions. **Colour bar is given in Kelvin. Same parameter averaged from 40° South to 5° North to avoid RFI and eclipse effects (bottom)**

### 240 3.2 Error Mitigation Techniques

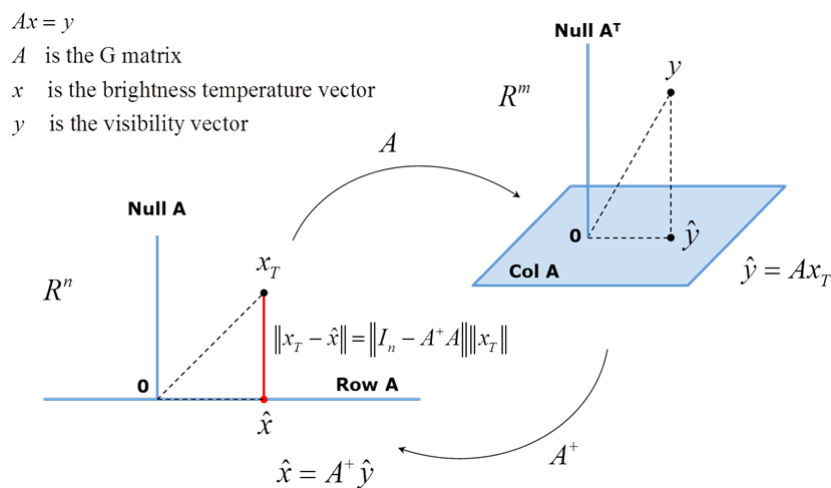
241 This section presents a summary of the techniques which have been attempted to mitigate, with more or  
242 less success, the various error sources described in the previous section. The overall approach to  
243 mitigate any of the errors has been to first understand the mechanism causing them and to then build a  
244 new calibration or image reconstruction approach to reduce it. Empirical corrections have been avoided  
245 as much as possible, and when accepted, they have been applied just once for all 5 year data set. In this  
246 section the focus shall be in image reconstruction based solutions, leaving for a later section those  
247 improvements brought into the new in-orbit calibration plan.

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### 3.2.1 Correction of Systematic Spatial Ripple

As explained earlier, given an antenna element spacing and a level of dissimilarity between the patterns of these, there is a minimum spatial ripple, the noise floor, which cannot be avoided. From a purely linear algebra point of view to the image reconstruction process, the noise floor can be understood with the help of Figure 8. A true brightness temperature distribution  $x_T$  in  $R^n$  space (with  $n$  being large) maps into  $y$  in the visibility domain in  $R^m$  ( $m$  also being large). However only a limited set of visibilities are actually measured, which in turn defines a fundamental hexagonal region in the physical space, smaller than the unity circle. Therefore only a visibility vector  $\hat{y}$  projected onto the column space of the –assumed perfectly known–  $G$  matrix is available for inversion. Using the pseudo-inverse matrix of  $G$ ,  $G^+$ , a least squares solution  $\hat{x}$  is found, which belongs to the row space of the  $G$  matrix, at some unavoidable distance from the true brightness temperature distribution  $x_T$ , this distance defining the noise floor. The noise floor is therefore the component of the true brightness temperature vector in the direction of the null space of the  $G$  matrix, and is therefore proportional to the brightness temperature of the scene.



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Figure 8: Noise floor definition

The following methods have been attempted, at Level-1 data processing, to mitigate the noise floor:

- a) Amplitude Mask

276 This technique is based on assuming that the spatial ripple can be mitigated by the use of a  
277 multiplicative mask built over uniform targets like the Southern Pacific Ocean or Antarctica (Lin et  
278 al., 2011, 2012)(Torres et al., 2012).

279 b) Floor Error Mask (FEM)

280 A difference brightness temperature map is built between the measurements and what an ideal  
281 instrument would reconstruct. Such map is produced over a full year, and split into a scene  
282 dependent and a scene independent components. The latter is used to correct SMOS images for the  
283 scene independent spatial ripple (Anterrieu et al., 2015).

284 c) Pre-Distorted G-matrix (PDG)

285 Using the Binomial Inverse Theorem the G-matrix is pre-distorted a priori in order to reduce the  
286 error contribution due to antenna pattern differences (Díez-García et al., 2014).

287 d) Average Pattern Reconstruction (APR)

288 The visibilities are decomposed in two components, one corresponding to an instrument with  
289 identical antenna patterns (and receiver responses) and a residual. The residual is then removed from  
290 the visibilities and an image reconstruction using an instrument with identical antenna and receivers  
291 is performed.

292 e) Initial guess based techniques

293 These methods are based on the second result stated in the section devoted to the spatial ripple, i.e.  
294 that the scene content outside the alias free field of view has an impact on the recovered scene inside  
295 it. According to this result, it is desirable to reduce the brightness temperature content outside the  
296 alias free field of view, which can be achieved if a first guess of the scene is subtracted prior to the  
297 image reconstruction. The set of methods implementing this technique is referred to as “Gibbs”  
298 methods, as they try to reduce the contrast within the scene, hence the Gibbs ringing, and by  
299 extension, the spatial ripple (Camps et al., 2008)(Khazaal et al., 2009) (Corbella et al., 2014). There  
300 are different levels of Gibbs techniques with increased implementation complexity: Gibbs-0 assigns

301 a single value to the whole unity circle; Gibbs-1 differentiates between the sky and the Earth disk,  
302 giving a single value of brightness temperature to the whole Earth disk, and another one to the sky;  
303 Gibbs-2 introduces, in addition, a differentiation within the Earth disk between land and ocean  
304 masses, assigning a constant but different value to each part; an option of Gibbs-2 includes a Fresnel  
305 variation over incidence angle over the ocean, instead of using a single value; in Gibbs-3 the first  
306 guess is obtained from the long record of measured brightness temperatures as opposed to the use of  
307 models, as done in the previous cases.

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309 The Gibbs-n techniques have been, among all, the most successful ones in reducing the spatial ripple.  
310 The success of the Gibbs-n techniques depends on the amplitude of any residual calibration errors.  
311 Nonetheless, because of the existence of the noise floor, it is unlikely that the current level of spatial  
312 ripple be significantly reduced in the future. The current Level-1 processor implements a Gibbs-1  
313 technique, the remaining spatial ripple being removed at Level-2 through the Ocean Target  
314 Transformation (Gourion et al., 2013), to enable the retrieval of Sea Surface Salinity.

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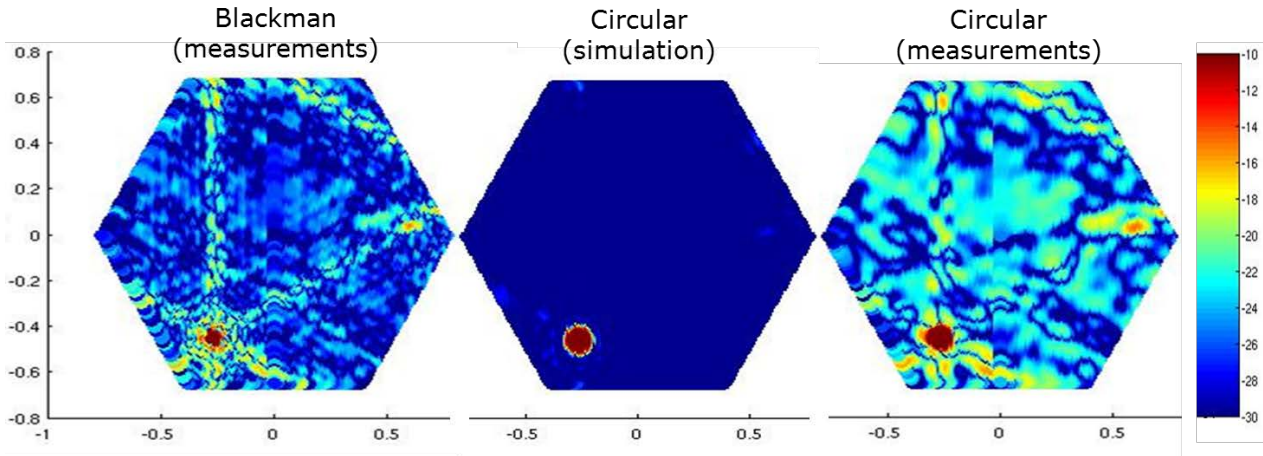
### 316 **3.2.2 Correction of Sun and RFI Tails**

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318 Given the element spacing in SMOS, the aliases of the Sun are unavoidable. The mitigation of the side  
319 lobes joining the real Sun and its aliases, visible in the left panel of Figure 9 with peak amplitudes of  
320 the order of  $-18$  dB, should be achievable through the application of a stronger apodization window in  
321 the spatial frequency domain than the nominal Blackman window (Camps, 1996) (Anterrieu et al.,  
322 2002). When a Circular window with much increased tapering than the nominal Blackman is applied to  
323 the spatial frequencies, the simulated response of the center panel should be obtained. Instead, the  
324 measurements, shown in the right panel, reveal that the side lobe level remains roughly at the same  
325 level (Torres et al., 2014). This puzzling result is a consequence, again, of the combination of the alias  
326 condition and the dissimilarities across the antenna patterns, as was the spatial ripple. In other words,

327 the noise floor sets a limit in the control that can be reached on the side lobes through windowing in the  
328 spatial frequency domain.

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Figure 9: Measured Sun response with standard Blackman window (left), expected image with a stronger circular apodization window (center) and the corresponding measurements (right). **Colour scale gives the brightness temperature in logarithmic scale (dB-K).**

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The following other techniques have been attempted to remove the Sun alias and its side lobes from the

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SMOS images:

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a) Use of a measured Sun response

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During one orbit of the 9<sup>th</sup>, 10<sup>th</sup>, 11<sup>th</sup> and 13<sup>th</sup> of January 2013, SMOS was pointed towards the

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galactic pole at a time when the Sun elevation was 21.3°, 17.1°, 10.1° and 21.6°, respectively,

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above the antenna plane. Four Sun responses were acquired over the uniform low cold sky

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background. These responses were later translated in position and scaled in amplitude, to

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correct for the Sun during measurements in nominal pointing. This method is still under

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research.

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b) Use of an estimation of the Sun response

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From an image contaminated by the Sun, the position and amplitude of the Sun are estimated,

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an artificial response is built based on them, and then subtracted from the original image

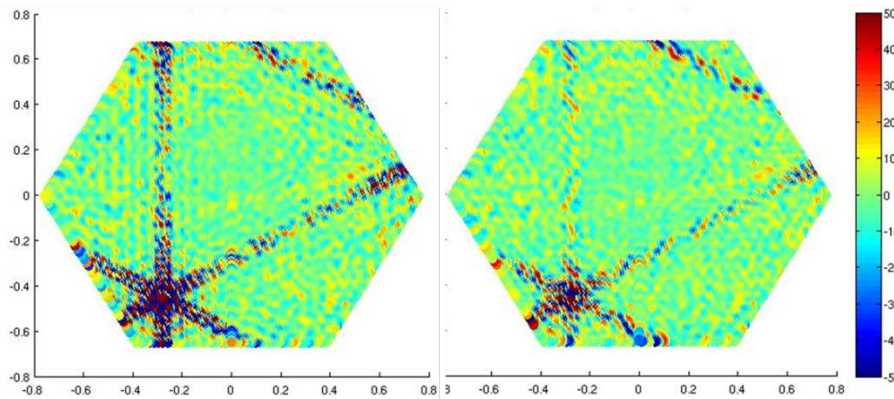
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(Camps et al., 2004). In this correction technique the most critical parameter is the accuracy in

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the localization of the Sun, which depends on the number of solar spots (Chiuderi Drago et al.

1977). An iterative process to estimate the position and brightness temperature of the Sun has shown promising results, but the increase in computational time is critical and is still under evaluation. A simpler method has been implemented in the Level-1 data processor of SMOS, with limited improvement, and has been complemented by flagging. Figure 10 shows an example of the Sun correction.



**Figure 10: SMOS image of the Sun in cold sky pointing mode before (left) and after (right) Sun correction. Colour scale is in Kelvin.**

In principle, the methods for the correction of the Sun alias and its tails can also be applied to remove RFI sources effects. However, the population of RFI sources is irregular, clustering in some regions of the Earth, with several interferers appearing at the same time inside the field of view of SMOS. In this situation a correction for the RFI sources of the type described for the Sun is difficult. Nonetheless, techniques to better detect, flag and correct for RFI sources keep being developed and assessed (Khazaal et al., 2014). In parallel, a technique called the Nodal Sampling (González-Gambau et al., 2015) has been proposed to image RFI polluted areas and is under assessment, showing some promising results.

### 3.2.3 Correction of Land-Sea Contamination

During the investigation of the land-sea contamination error it has become clear that a mismatch between the amplitude of the visibility at the origin  $V(0,0)$  and the rest of the samples,  $V(k,j)$ , generates this kind of degradation, as shown in the left plot of Figure 5. Empirically it has been proven that



373 affecting the correlator efficiency coefficients  $G_{kj}$  of the visibility samples  $V(k,j)$  outside the origin by a  
374 factor near 0.98 removes significantly the land-sea contamination. The right plot of Figure 5 shows the  
375 improvement when this correction is applied. The warm brightness temperature halos surrounding the  
376 continents have mostly disappeared. The possibility of correcting the correlation efficiency coefficients  
377  $G_{kj}$  has been implemented in the latest version of the SMOS Level-1 processor (V700) and will be  
378 subject of validation before it is accepted to enter into operation.

379 It is worth mentioning that a parallel empirical correction of the land-sea contamination is being  
380 prepared at Level-2 based on a mask built with the 5 year long record of SMOS data (SMOS Level-2  
381 Ocean Salinity Team, 2015).

382

### 383 **3.2.4 Correction of Seasonal and Orbital Variations**

384 The seasonal and orbital variations are observed in the right plots of Figures 6 and 7 respectively as  
385 warm anomalies around October and in the eclipse season every year. The current strategy to correct  
386 for these fluctuations is to simplify the calibration approach of the instrument as much as possible by  
387 using the All-LICEF mode (Torres et al., 2006) explained below and then attempt new corrections to  
388 mitigate them.

390

## 391 **4 IN FLIGHT CALIBRATION PLAN**

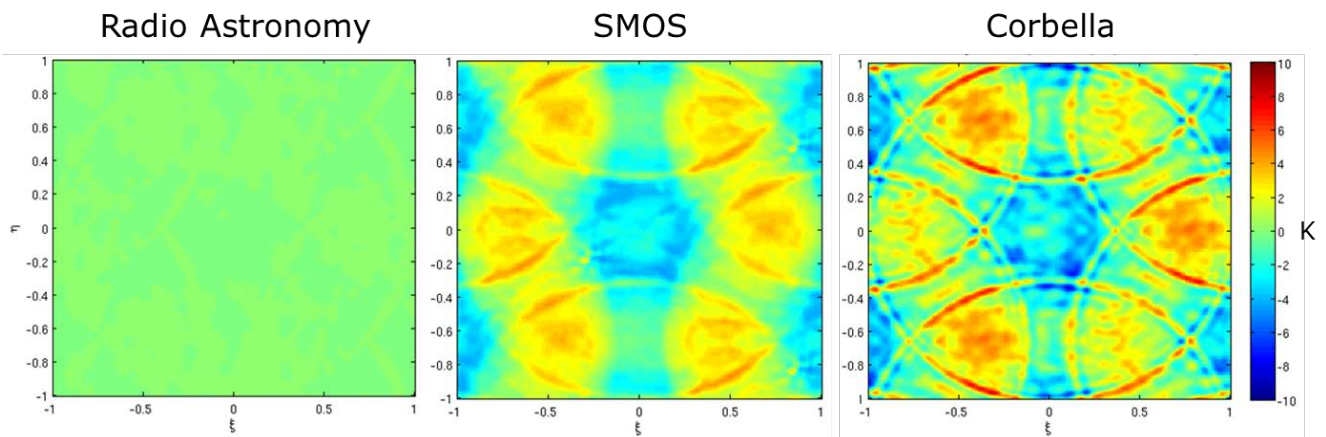
### 392 **4.1 The Corbella Equation**

393 The Corbella equation, introduced in 2003 (Corbella et al., 2004a), involves a fundamental  
394 modification to the formulation of interferometry, as used in radio-astronomy, that is necessary to  
395 describe the way an aperture synthesis radiometer of the type of MIRAS works. The calibration of  
396 SMOS is based on the Corbella equation, and hence, one of the first and most important tasks  
397 undertaken in the frame of the calibration of the instrument was its verification (Martín-Porqueras et  
398 al., 2010). Such exercise would ideally involve the imaging of two perfectly uniform black body targets

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399 at two different physical temperatures. Since the Cold Sky near the galactic pole is the only reasonable  
400 realization to such uniform target, the validation of the Corbella equation focused on the Cold Sky.  
401 With the help of a simulator, two sets of visibility samples of the Cold Sky were produced using the  
402 radio astronomy and the Corbella equations. Then, to improve contrast, the visibility samples at the  
403 origin were set to zero, and a simple Fourier Transform was applied to provide the images of the  
404 modified brightness temperature of the Cold Sky in each case. The image obtained using the radio  
405 astronomy formulation is shown in the left panel of Figure 11, while the one resulting from the use of  
406 the Corbella equation is in the right panel. These images were then compared to the one measured by  
407 SMOS, shown in the center panel of Figure 11, which was generated in the same way, i.e. through a  
408 Fourier Transform of the measured visibility samples setting the one at the origin to zero. The image  
409 using the Corbella equation is very similar to that measured by SMOS, while the Cold Sky retrieved  
410 with the radio astronomy equation does not capture any of the features present in the measurements.  
411 This test validated the Corbella equation.

412



413  
414 Figure 11: Expected modified brightness temperature of the Cold Sky using the radio astronomy (left)  
415 and the Corbella (right) equation; center is the SMOS measurement (note: the visibility sample at the  
416 origin has been set to zero to improve contrast)

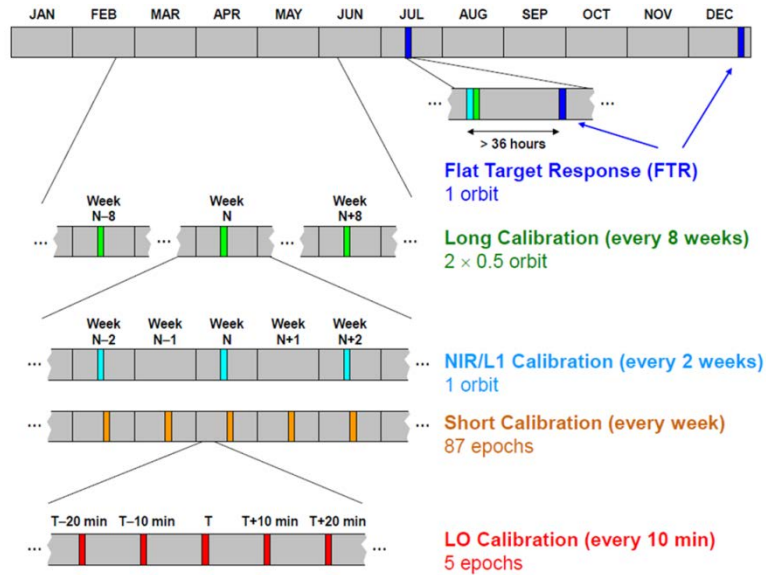
#### 417 4.2 Routine In-orbit Calibration Plan

418 The routine in-orbit calibration plan was established at the end of the Commissioning Phase, in May  
419 2010. With the experience of the first year of the operational phase (Oliva et al., 2013), weekly Short  
420 Calibrations while flying over Antarctica were added as from March 2011 to track the temporal

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421 variation of the voltage offset of the receivers, leading to the calibration plan shown in Figure 12,  
 422 which is the one currently used in SMOS.

423



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Figure 12: SMOS routine in-orbit calibration plan

### 428 4.3 In-flight Calibration Improvements

429 Two improvements have been made based on the in-flight experience over these 5 years: the “warm”  
 430 external calibrations and the addition of an RFI check to validate the external calibrations.

#### 431 4.3.1 Warm Calibrations

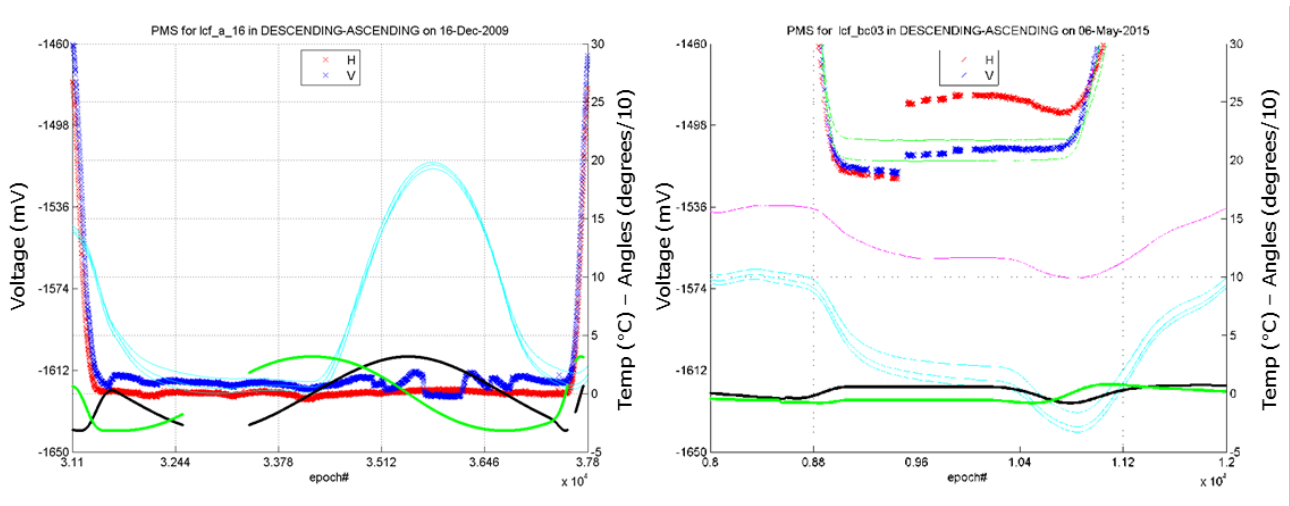
432  
 433 Detailed analysis of the external calibrations revealed that a few LICEF receivers of MIRAS exhibited  
 434 small and smooth unexpected jumps in their PMS (Power Monitoring System) detector voltages. These  
 435 jumps seemed to correlate well with the skin temperature of the antenna, happening more frequently for  
 436 colder skin temperatures, and appeared to be reversible in the sense that, for warmer skin temperatures,  
 437 the usual values were again obtained. To illustrate this, refer to the left panel in Figure 13, which spans  
 438 one full orbit flown pointing zenith during the Commissioning Phase, **including the transitions from**  
 439 **and to Earth pointing at the beginning and end of the plot, respectively**. The 3 cyan curves provide  
 440 the skin temperature of the antenna measured by 3 thermistors (named Tp7). The black and green lines  
 441 give, respectively, the elevation of the Sun over the antenna plane and its azimuth, in decadegrees as

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442 read from the scale on the right. The Sun elevation is negative (the Sun is behind the antenna) except  
443 for a portion in the right half of the plot, where it reaches an elevation of about  $30^\circ$ . As soon as the Sun  
444 appears in front of the antenna the skin temperature (cyan lines) increases from near  $0^\circ\text{C}$  till some  
445  $20^\circ\text{C}$ , to return back to just  $1^\circ\text{C}$  or  $2^\circ\text{C}$  as the Sun sets behind the antenna horizon. The dark blue and  
446 dark red crosses correspond to the detected voltages in the vertical and horizontal polarizations. The  
447 vertical polarization shows some fluctuations. The first anomaly appears towards the left of the plot  
448 when the physical temperature drops below some  $10^\circ\text{C}$ . At this moment the detected voltage jumps up  
449 a few millivolts, which is unexpected because as the instrument is pointed towards cold sky the  
450 detected voltage is expected to constantly decrease towards a minimum level. The anomalous higher  
451 value is maintained until the temperature rises again above some  $12^\circ\text{C}$ . Then several fluctuations  
452 happen centered around the maximum of temperature in the right half of the plot, showing a high  
453 degree of symmetry and correlation with the temperature evolution. The detected voltage attains a right  
454 value only in the center of these fluctuations, coinciding with the warmest temperature interval around  
455 the peak. It is plausible that these fluctuations could be related to a change in the electrical phase of  
456 some Teflon pieces of the antenna at some physical temperature range. Another example is given in the  
457 right panel of Figure 13, which corresponds to another of the affected receivers, this time during a  
458 typical external calibration manoeuvre: as the antenna cools down and its temperature reaches about  
459  $2^\circ\text{C}$ , the detected voltages at the two polarizations experience jumps of 10 and 40 mV about. The  
460 purple line in the right panel is the estimated physical temperature at which the Teflon parts of the  
461 antenna could be. To avoid these voltage jumps, the external calibration manoeuvres are planned, since  
462 October 2014, **at a modified time to have** the Sun at some positive elevation angle over the antenna  
463 plane. The Sun illumination on the antenna keeps it warm, avoiding the skin temperature to fall too low  
464 and the PMS detector voltage fluctuations. The positive Sun elevation is however kept below a limit of  
465  $10^\circ$  to ensure that its presence does not degrade the external calibration acquisitions. Careful checks  
466 were carried out to detect any effect from the direct signal of the Sun that could compromise the quality

467 of the external calibration. In addition the warm calibrations perform the external calibration with the  
468 antenna at a more similar temperature to the measurement mode, which is desirable, and reduce thermal  
469 excursion on the antenna, improving reliability. Warm calibrations can be planned any time along the  
470 year except around the equinoxes, when the Sun elevation is just too low, in which case, the usual cold  
471 calibrations are performed instead.

472



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476

Figure 13: Example of PMS detector voltage fluctuations in a zenith pointing orbit (left) and during a typical external calibration (right) that led to the introduction of the external ‘warm’ calibrations

#### 477 4.3.2 RFI Check in Validation of External Calibrations

478

479 External calibrations, where SMOS is pointed towards the Cold Sky, are executed only over the Pacific  
480 Ocean to avoid picking up signals from strong RFI transmitters on ground through the back lobes of the  
481 antennas. However, in one instance, an external calibration carried out 3 June 2015 in the North-  
482 Eastern Pacific Ocean, near Alaska, appeared contaminated by some ground interference. This caused  
483 some disturbance in the data production chain as the calibration file had been ingested before the  
484 problem was discovered. To avoid this, since then, every external calibration (these are performed once  
485 every 2 weeks) is manually checked for RFI degradation before being accepted for use in the Level-1  
486 data processor. An automatic procedure is being built up to replace the manual check.

487

## 488 **5 IN FLIGHT INSTRUMENT MONITORING**

489

490 The values of some of the key instrument parameters which are monitored or calibrated in flight are  
491 presented next.

### 492 **5.1 Physical Temperature**

#### 493 **5.1.1 Skin Antenna Temperature**

494

495 MIRAS carries a thermistor (labeled Tp7) inside the head of the central screw of the antenna of the 3  
496 Noise Injection Radiometers whose readings are representative of the physical skin temperature of any  
497 of the antennas (Rubiales et al., 2015). This temperature is important because it affects the amount of  
498 noise emitted by the front end equivalent resistor. It also provides an indication on how different the  
499 thermal conditions of the antenna are between an external calibration and the nominal measurement  
500 mode.

501 The evolution of the temperature readings from the 3 Tp7 thermistors is shown in Figure 14. The skin  
502 temperature goes through its largest excursion (from about 6°C to 28°C) during **every** boreal winter  
503 solstice, when the Sun reaches maximum elevation above the antenna plane (around 31°) and is  
504 eclipsed by the Earth. There is a second period of large thermal excursion (from 8°C to 18°C  
505 approximately) around every boreal summer solstice where the Sun elevation reaches up to 15°  
506 elevation above the antenna horizon. During the equinoxes the temperature excursion is the smallest  
507 (between 5°C to 12°C) and the lowest skin temperatures are recorded, except for the external  
508 calibration events. The latter correspond to the individual spikes that drop below 0°C in Figure 14. The  
509 Tp7 temperatures went through an initial cooling transient, clearly observed during the first half of  
510 2010, to then flatten out into a very small long term residual cooling trend.

511

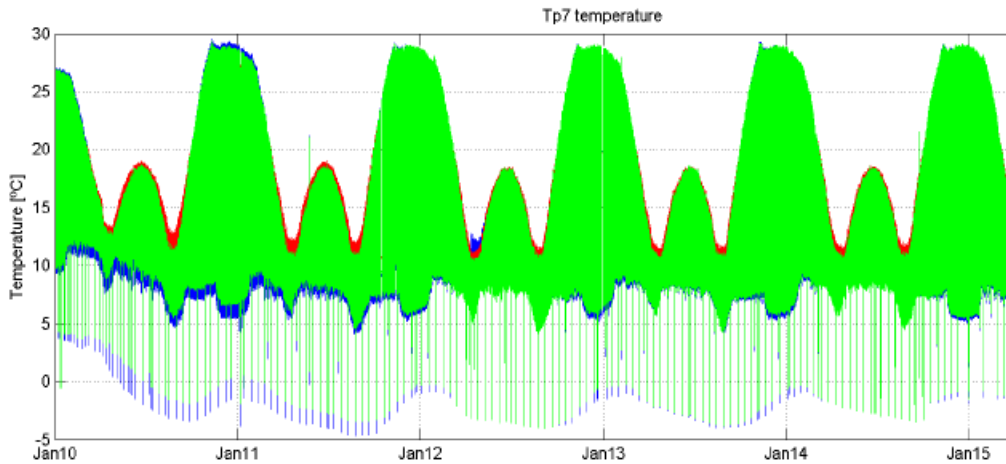


Figure 14: Evolution of the skin antenna temperature measured by the 3 Tp7 thermistors

### 5.1.2 Inner Receiver Temperature

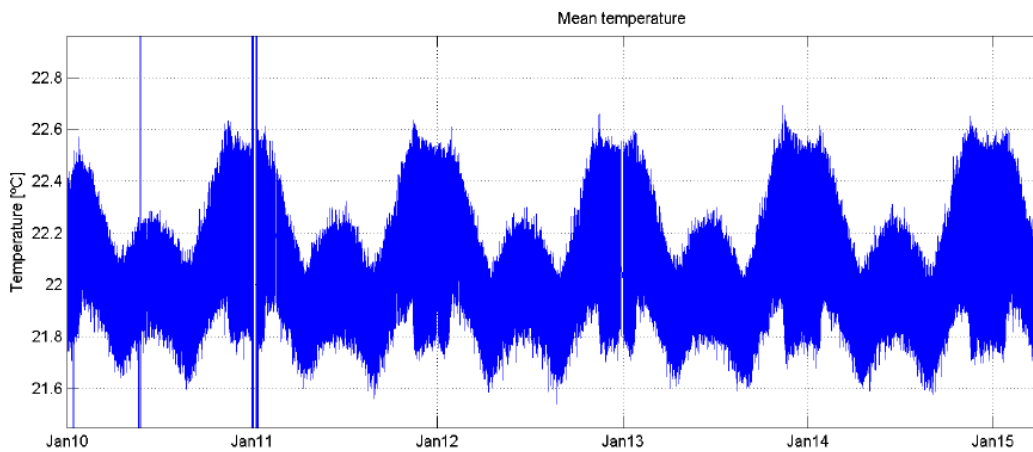
Every one of the 72 LICEFF receivers of MIRAS has a thermistor (labelled Tp6) next to an internal matched load in the front-end electronics used as warm point in the amplitude calibration. This thermistor senses the inner temperature of the receiver. The average value of Tp6 across all LICEF receivers is shown in Figure 15. The physical temperature of the receivers is **seen** to be quite stable along the mission, centered around 22°C with a global peak to peak variation of about 1°C. As for Tp7, the Tp6 readings present larger excursions during the solstices, and narrower variation around the equinoxes, where its lowest values are attained.

## 5.2 Receiver Parameters

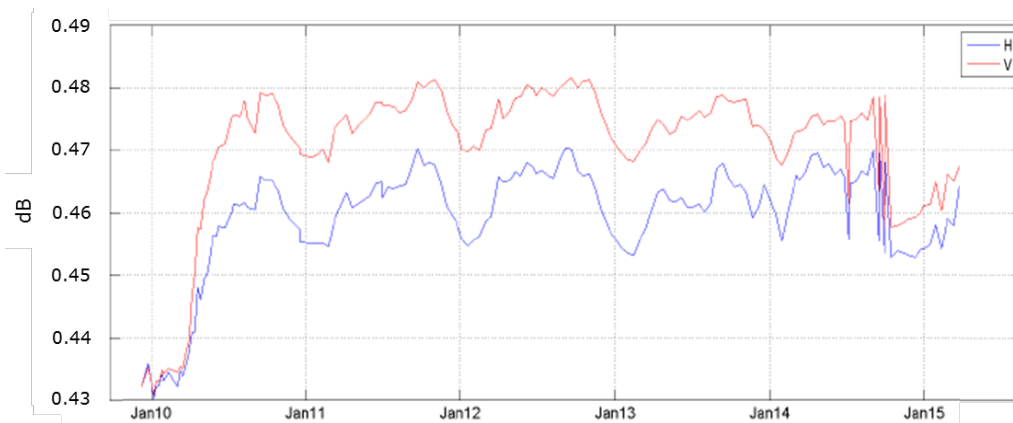
### 5.2.1 Antenna Losses

The antenna has two distinct loss components: one due to the radiating resonant cavity, and the other due to the intermediate layer circuit that combines the signal from the pair of probes of each polarization. The first component is tiny and difficult to measure on ground. It was estimated to be of about  $L1=0.05$  dB, by calculations based on the geometry and materials of the antenna design. On the other hand, the losses of the intermediate layer circuit, of about  $L2=0.25$  dB, was measured on ground. The total antenna losses are then expected to be around 0.30 dB. During the in orbit calibration, the

535 antenna loss are directly measured using the Cold Sky and the internal matched load (Corbella et al.,  
536 2012). The average value across all LICEF receivers for each polarization is shown in Figure 16. The  
537 in-flight measured antenna losses are about 0.17 dB larger than their pre-launch estimated value. It is  
538 worth noticing the rapid evolution exhibited during the first 6 months of the mission, as well as the  
539 seasonal fluctuations, the latter being partly driven by the PMS detector voltage fluctuations described  
540 earlier. The antenna losses present a different evolution after October 2014, reflecting the introduction  
541 of the warm external calibrations to avoid the mentioned PMS fluctuations.



543  
544 Figure 15: Average inner LICEF receiver temperature Tp6. **The spikes in early Jan'10, May'10 and**  
545 **Jan'11 are due to 3 anomalies occurred in the instrument.**  
546



547  
548 Figure 16: Evolution of the antenna losses as measured in flight  
549  
550

### 551 5.2.2 Receiver Detector Gains

552 The average detector gain across all LICEF receivers is presented in Figure 17 for each polarization.

553 The absolute gain is shown in the left panel. Similar features to those found in the evolution of the  
554

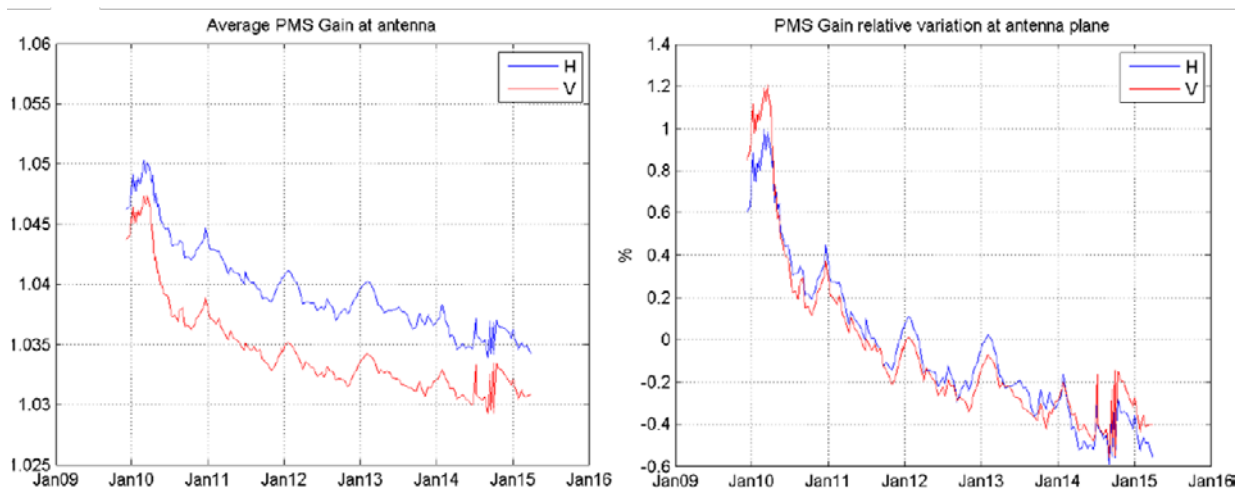


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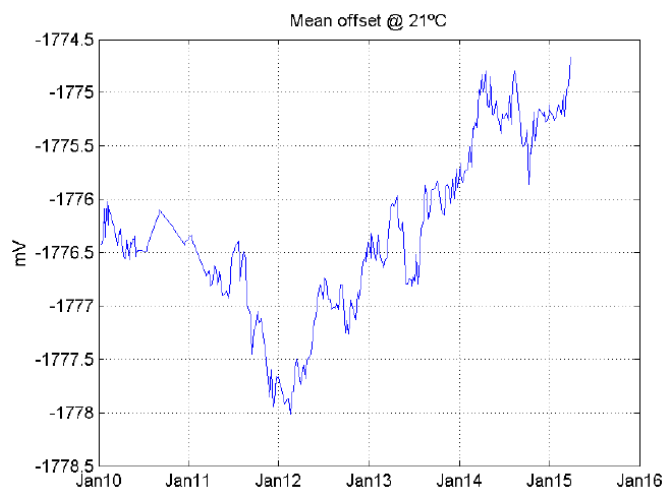
555 antenna losses are repeated here: an initial rapid transient followed by seasonal variations. In addition  
556 the receiver gains seem to be undergoing an exponential decay which, according to the relative gain  
557 variation shown in the right panel, has accumulated a total decrease of about 1.5%. **The reason for this**  
558 **decay is unknown, but could be caused by the overall thermal stabilization over mission life time.**

559  
560 **5.2.3 Receiver Detector Voltage Offsets**  
561

562 The average voltage offset across all LICEF receivers is shown in Figure 18. The behavior is somewhat  
563 erratic, without any clear trend, with rapid fluctuations that led, in March 2011, to the introduction of  
564 weekly short calibrations as from March 2011 to track them. The voltage offset is therefore well  
565 followed with a weekly refresh rate and **calibrated** out.



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568 Figure 17: Evolution of the end-to-end average receiver gain in mV/K (left) and in percentage variation  
569 taking June 2011 as reference value (right)  
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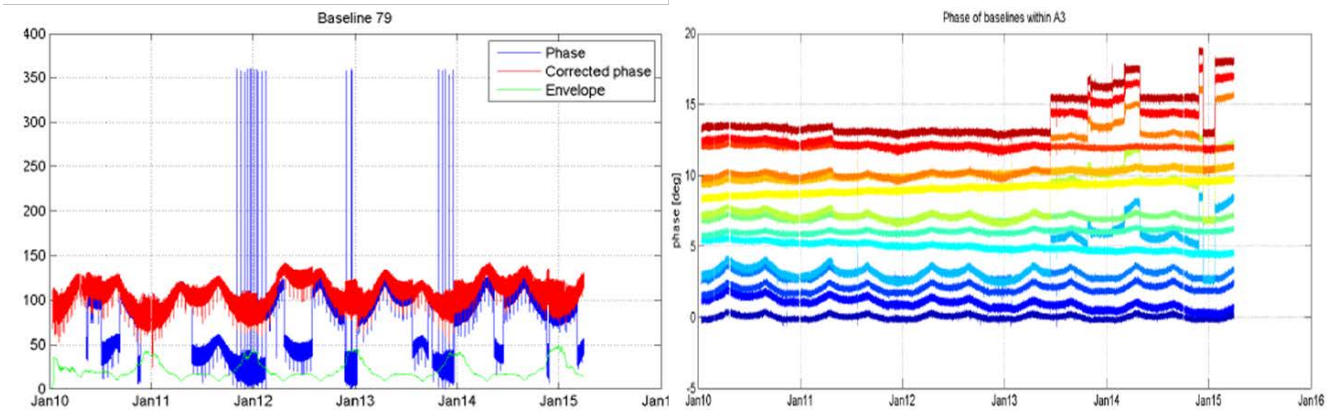
573  
574 Figure 18: Evolution of the average detector offset voltage (in mV) across all LICEF receivers  
575

### 576 **5.3 Baseline Parameters**

#### 577 **5.3.1 Relative Phase across Receivers**

578 The relative phase across receivers is measured once every 10 minutes through the injection of a burst  
579 of correlated noise into all LICEF receivers. Phases between receivers sitting in different arm segments,  
580 hence fed by physically different local oscillators, present the strongest temporal variations. **In**  
581 **addition some of the 12 phase-lock-loop circuits available on board lose lock from time to time**  
582 **for reasons still unknown, but probably due to temperature (in total 9 unlocks per year, on**  
583 **average), generating phase jumps.** As an example, the phase between one receiver in the hub and  
584 other receiver in the first segment of one of the arms of MIRAS is shown by the blue line in the left  
585 panel of Figure 19. After correcting for any unlocks (blue spikes), the corrected phase in red is  
586 obtained. The fluctuations of the corrected phase are correlated with the physical temperature, which  
587 causes slight differential changes in the two local oscillators involved in the particular baseline.

589 On the other hand, the relative phase across receivers fed by the same local oscillator, i.e. sitting in the  
590 same arm segment, is rather stable, **and are not affected by the unlocks.** The right panel of Figure 19  
591 brings the evolution of the relative phase across the 6 receivers of the third segment of one of the arms,  
592 shifted vertically only for presentation purposes. The peak to peak phase fluctuations are below  $1^\circ$   
593 across the entire mission except for a few jumps that can be attributed to one specific receiver, probably  
594 caused by thermal expansion in one cable connection, but for which there is no proven explanation.



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Figure 19: Relative phase between 2 LICEF receivers which do not share the same local oscillator (left) - spikes correspond to  $360^\circ$  wrappings - and across each of the 15 possible pairs from the set of 6 LICEF receivers of one arm segment connected to the same local oscillator (right)

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### 5.3.2 Correlator Amplitude Coefficients

The average value of all the 2556 efficiency coefficients  $G_{kj}$  of the correlator for each baseline formed by receivers  $k$  and  $j$  is shown in Figure 20. **These  $G_{kj}$  coefficients represent the value of the so-called Fringe-Washing Function (FWF in short) at zero delay. They provide the correlation losses due to the on-board calibration network and the differences in the frequency responses of the receivers. The  $G_{kj}$  coefficients are used to denormalize the value of the raw correlations obtained from the 1-bit sampled signals. In theory the  $G_{kj}$  cannot be larger than 1, but because they are obtained through a combination of measurements of correlated noise injection from common Noise Sources (NS) and closure equations, some of them reach values slightly above 1. The  $G_{kj}$  coefficients are very stable along the mission. The average of all  $G_{kj}$  is about 0.99, without any significant temporal trend.**

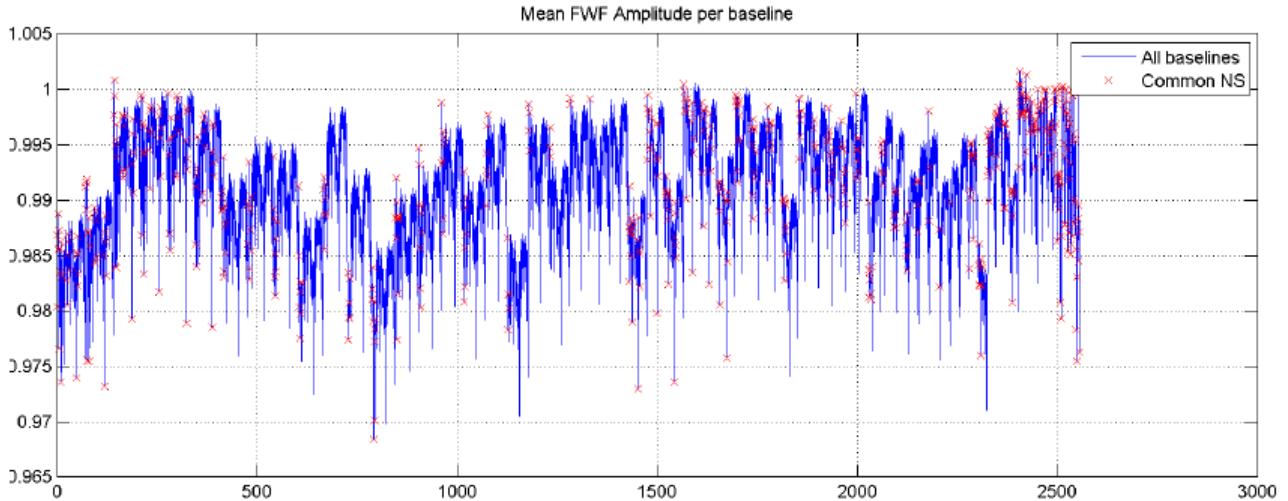


Figure 20: Average value of the  $G_{kj}$  correlator efficiency coefficients

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#### 620 5.4 Noise Injection Radiometer Parameters

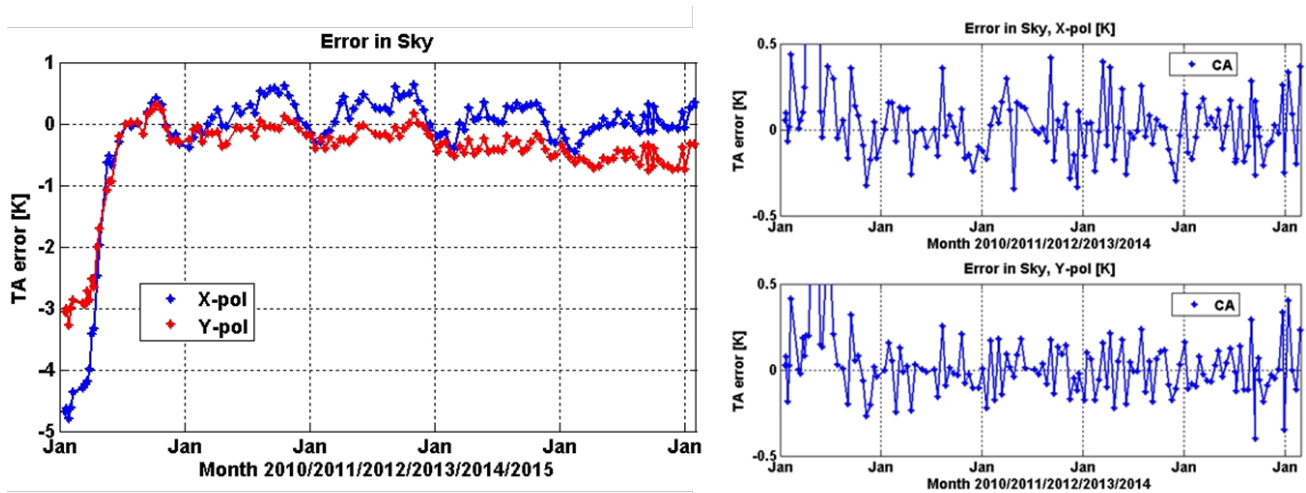
621 The 3 Noise Injection Radiometers (NIR) embarked on SMOS are the reference radiometers of MIRAS  
622 (Colliander et al., 2007). They serve several purposes: (a) they measure the visibility sample at the  
623 origin; (b) they provide the reference antenna temperature to de-normalize the visibility samples  
624 obtained through 1-bit correlations; (c) they measure the amplitude of the noise diodes of the on-board  
625 Calibration System (Lemmetynen et al., 2007); and (d) they are used as an additional receiver to  
626 generate visibility samples. A summary of their in orbit performance is presented next.

##### 627 5.4.1 NIR Long and Short Term Stability

628  
629 The NIR long term stability is illustrated in the left panel of Figure 21, which gives the error in each of  
630 the two polarizations against the Cold Sky for all external calibrations manoeuvres, performed once  
631 every 2 weeks, when a single calibration point, at the beginning of the Operational Phase (June 2010) is  
632 used. As it is shown, after the transient of the first 6 months of the mission (Kainulainen et al., 2010,  
633 2012), the stability of the NIR units is remarkable: even when being calibrated only once in 4 years, the  
634 long term drift is only of about  $-0.1$  K/year. The right panel of Figure 21 shows the short term stability  
635 of the NIR units by providing the Cold Sky error just before the NIR parameters are refreshed at every  
636 external calibration. The accumulated error over 2 weeks is within  $\pm 0.4$  K. The bi-weekly external

637 calibrations remove completely any long term drift. The annual behavior of the NIR error, that is  
638 clearly seen in Figure 21 (left) is investigated a lot during the mission. Trials have been made to  
639 relate the drift to internal instrument properties like physical temperatures, and to external  
640 conditions, like Sun position or other components that contribute to the total antenna  
641 temperature. Contributions of such external sources are reviewed e.g. in (Colliander, 2015), in  
642 which forward geophysical model simulations were used to simulate NIR antenna temperature  
643 measurements.  
644 The observed 0.13 K standard deviation error of the sky measurements in Figure 21 scales down  
645 to about 0.1 K error in measurement pointing mode (Kainulainen et al, 2012), a smaller error  
646 than the 0.2 K radiometric resolution of the NIR units, and thus, not affecting the retrieval of the  
647 geophysical parameters.

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Figure 21: NIR-CA stability when calibrated only once (left) or every 2 weeks (right)

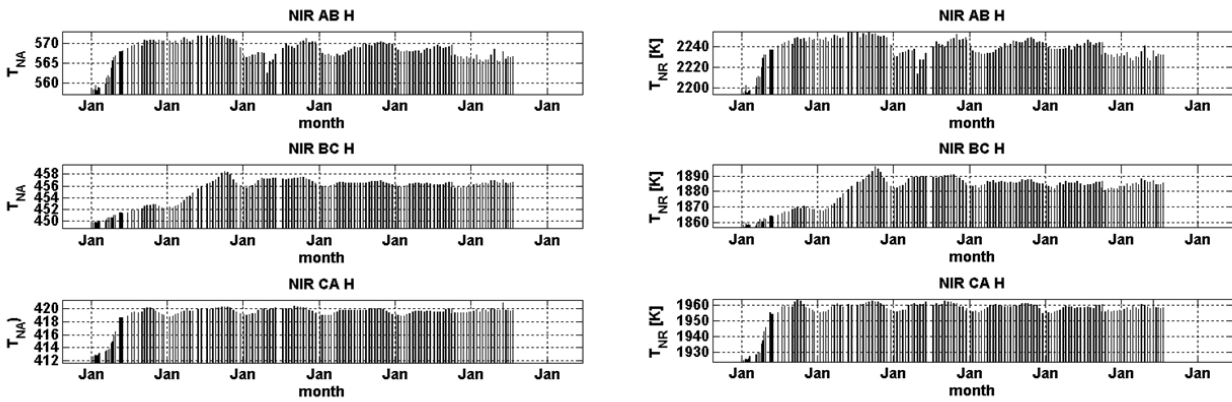
#### 652 5.4.2 Noise Injection Temperatures

653

654 NIR uses some internal noise diodes to operate. The stability of these diodes and their injection circuits  
655 are critical to establish the performance of the NIR units and the whole MIRAS instrument. Two levels  
656 of noise injection temperatures are used, a low one  $T_{na}$  to measure the antenna temperature, and a high  
657 one,  $T_{nr}$  to measure the diodes of the on-board Calibration System. Figure 22 shows the 5 year  
658 evolution of  $T_{na}$  and  $T_{nr}$  in the left and right panels respectively, for the horizontal polarization (similar

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659 results are obtained in vertical polarization). Besides the initial transient, the variations over 5 years are  
660 within 0.2% peak to peak for the most stable unit, NIR-CA.



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Figure 22: Temporal evolution of the noise injection temperatures of the NIR units

## 664 **6 SMOS LEVEL-1 DATA PROCESSOR PERFORMANCE**

665

666 Figure 23 shows the evolution of the SMOS Level-1 data processor that transforms the raw visibility  
667 samples into calibrated brightness temperature records. The first version used during the  
668 Commissioning Phase (November 2009 till May 2010) was V324 while the first processor supporting  
669 the Operational Phase was V344. During 2011 version V500 was deployed and a full reprocessing of  
670 the mission data set between January 2010 till October 2011 was carried out with a slightly modified  
671 version V505. A new version V600 was ready by end of 2012, which, after further enhancements,  
672 became V620 by early 2014. A second full mission reprocessing of the data between January 2010 and  
673 May 2015 took place using V620, version which was deployed in early 2015. This section will focus  
674 on the improvements brought in by V620, used in the second reprocessing, with respect to the earlier  
675 V505 of the first reprocessing, as well as in the performance assessment of the new processor version.

676

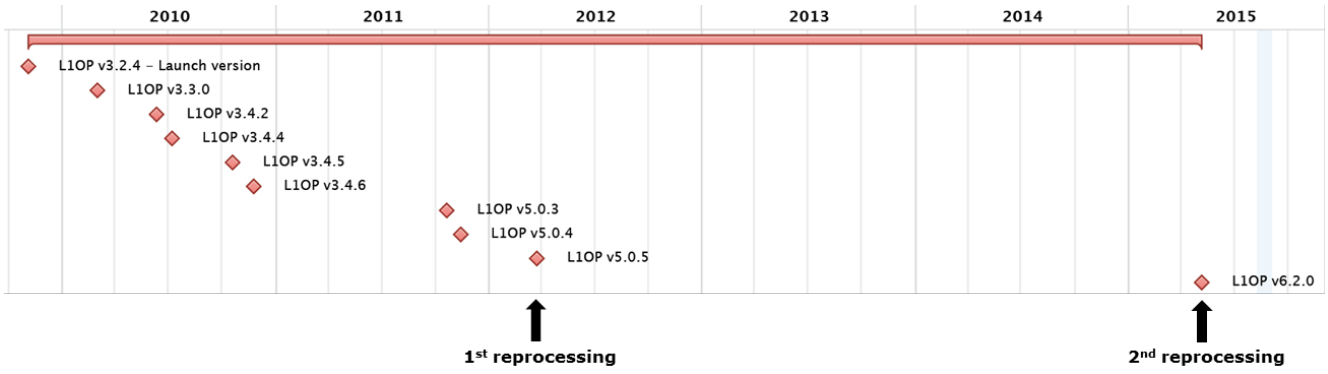


Figure 23: SMOS Level-1 data processor evolution

## 6.1 Enhancements in Level-1 Processor Version V620 over V505

### 6.1.1 Fully Polarimetric Operation

The equation below gives the matrix relation between the brightness temperature and the visibility vectors through the G observation matrix (Martín-Neira et al., 2002) (Corbella et al., 2004b):

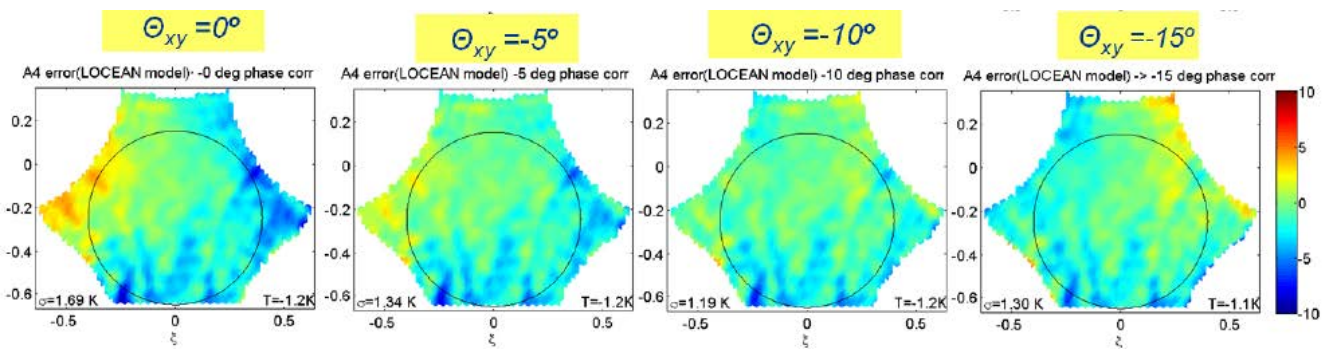
$$\begin{bmatrix} \Delta V_{kj}^{HH} \\ \Delta V_{kj}^{VV} \\ \Delta V_{kj}^{HV} \end{bmatrix} = \begin{bmatrix} G_{T_{HH}}^{HH} & G_{T_{VV}}^{HH} & G_{T_{HV}}^{HH} \\ G_{T_{HH}}^{VV} & G_{T_{VV}}^{VV} & G_{T_{HV}}^{VV} \\ G_{T_{HH}}^{HV} & G_{T_{VV}}^{HV} & G_{T_{HV}}^{HV} \end{bmatrix} \times \begin{bmatrix} \Delta T_B^{HH}(\xi, \eta) \\ \Delta T_B^{VV}(\xi, \eta) \\ \Delta T_B^{HV}(\xi, \eta) \end{bmatrix}$$

The elements along the main diagonal of the G matrix involve co-polar antenna patterns. The elements outside the main diagonal include one or two cross-polar patterns. V505 was built using only co-polar patterns, that is, the G matrix was block-diagonal. The new version V620 makes use of the cross-polar patterns, which were measured on ground, and hence, implements a fully populated G matrix with all the blocks shown above. The first attempts to use the cross polar patterns were not successful though, because there was a sign inconsistency between the ground measurements and the Level-1 data processor. Once this issue was identified and solved, the expected performance improvement was achieved (Torres et al., 2015). The G matrix is used inside the Level-1 processor in the forward modelling as well as in the inverse process of image reconstruction (Khazaal et al., 2015). V620 is the first fully polarimetric SMOS Level-1 data processor. **Among all enhancements, this one providing fully polarized data is perhaps the most significant one. In fact, this is the first time that fully polarized interferometric images from space are processed.**

699 **6.1.2 Use of the Relative Phase between Polarizations Measured on Ground**

700

701 During the Image Validation Test (IVT) of MIRAS that was carried out on ground before launch, a set  
702 of 4 probes was placed in the ceiling of the Maxwell Electromagnetic Compatibility chamber of  
703 ESTEC to measure the relative phase between all LICEF receivers (Corbella et al., 2009). The  
704 instrument was operated in both dual and full polarization. For version V505 of the Level-1 processor,  
705 2 separate sets of relative phases were retrieved from the IVT test: one with the relative phases for the  
706 horizontal polarization and another one for the vertical polarization. When preparing the next version,  
707 V620, of the processor, it became clear that the two sets of relative phases could have an offset between  
708 them which had to be corrected. Such phase bias across the two polarizations was in fact causing  
709 distortions in the Stokes-3 and, most clearly, Stokes-4 parameters. Therefore, the IVT data set was re-  
710 analyzed to determine the missing phase offset between the two polarizations, which was found to be of  
711  $-6.8^\circ$ . This value was verified using Stokes-4 images over the ocean: as shown in Figure 24, the error  
712 (sigma displayed in the lower left corner) was indeed minimized for a phase offset close to the retrieved  
713 one. All phases corresponding to vertical polarization were then reduced by that amount in version  
714 V620 of the processor.



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Figure 24: Validation of the relative phase between polarizations using Stokes-4 over ocean.

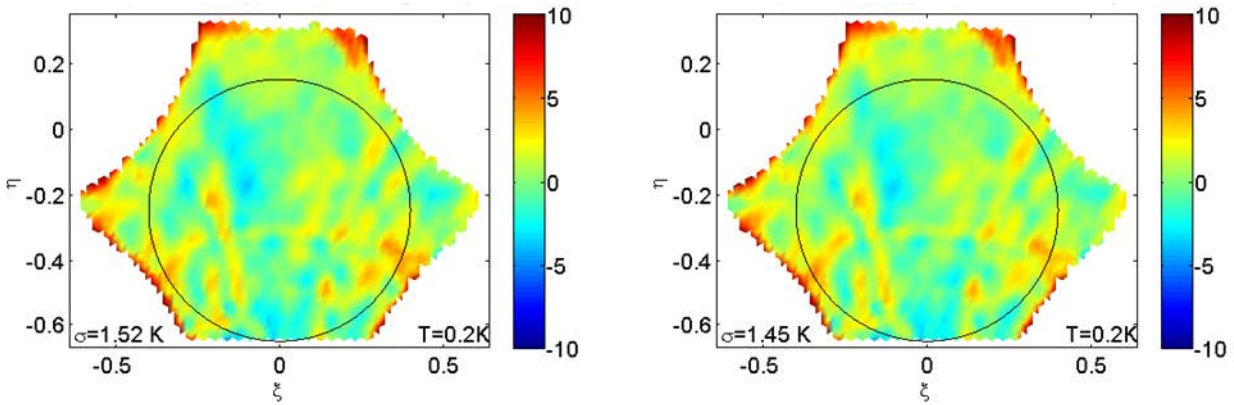
719 **6.1.3 Use of Average Antenna Patterns across 3 Frequencies Measured on Ground**

720

721 The antenna patterns of all and every element of MIRAS embedded in the array were carefully  
722 measured in an antenna test range. The measurements were performed at the center frequency, 1413.5  
723 MHz, as well as at the band edges, 1404 and 1423 MHz. In version V505 of the Level-1 processor only



724 the patterns at the center frequency were used. In version V620, the average antenna pattern across the  
725 3 measured frequencies is used, which reduces slightly the spatial ripple, as shown in the left lower  
726 corner of each panel of Figure 25.



727  
728 Figure 25: Deviation images from an ocean model in X-polarization using 1 pattern (left) or the  
729 average of the 3 measured patterns on ground (right) –similar results were obtained in Y-polarization–.  
730

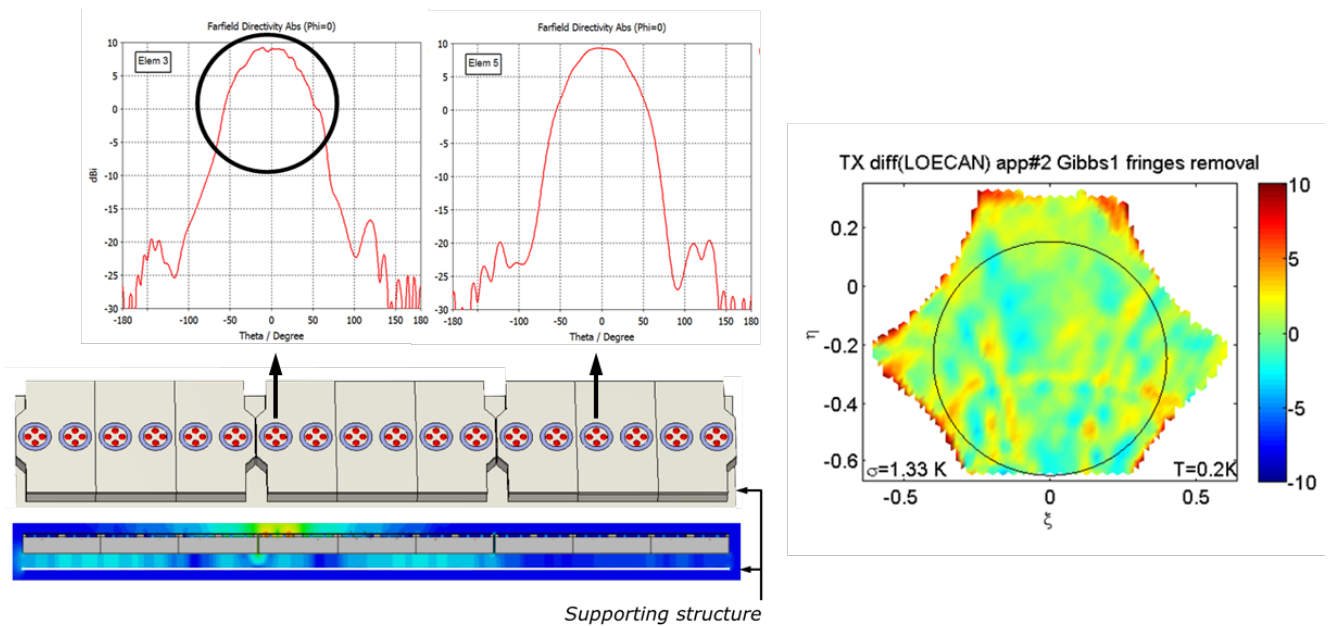
#### 731 **6.1.4 Replacement of Antenna Pattern of Hinge Elements by their Neighbour's**

732  
733 Analyzing in detail the measured antenna patterns, it was noticed that those corresponding to elements  
734 next to one of the hinges of the deployed arms (upper left panel of Figure 26), presented some ripples  
735 which the elements in the center of the arm segments did not exhibit (upper center panel). Some  
736 research led to the conclusion that these ripples were caused by travelling waves between the arm and  
737 the supporting structure that was used to hold it during the measurements, as shown in the lower left  
738 panel of Figure 26, which would leak out through the next hinge and segment contour causing a typical  
739 interference pattern. Since the instrument in flight configuration does not have any supporting structure  
740 in the back, the real patterns of the hinge elements should be free of fringes and similar to those of the  
741 central elements in each arm segment. For this reason, in the Level-1 processor version V620, the  
742 antenna patterns of the hinge elements was replaced by that of its inner neighbours, which reduces  
743 slightly the spatial ripple of the images, as illustrated in the right panel of Figure 26, to be compared to  
744 the left panel of Figure 25.

#### 745 **6.1.5 Use of Only the Most Stable NIR Unit**

746

747 From Figure 22, it is seen that the NIR-CA unit (bottom panels) is the most stable of the 3 Noise  
748 Injection Radiometers of MIRAS. To assure the best temporal stability, it was decided to use, in V620,  
749 only the antenna temperature measured by NIR-CA, and not those from NIR-AB and NIR-BC, for the  
750 visibility sample at the origin  $V(0,0)$ . All 3 NIR units are still employed in the measurements of other  
751 visibility samples outside the origin as well as in the measurement of the noise diodes of the on-board  
752 Calibration System.



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Figure 26: Antenna patterns of elements next to a hinge (top left) and away from it (top center); electromagnetic simulation showing leakage between arm and back supporting structure (left bottom); image obtained by replacing hinge patterns by their neighbour's (right).

### 759 6.1.6 Use of In-Orbit Antenna Loss

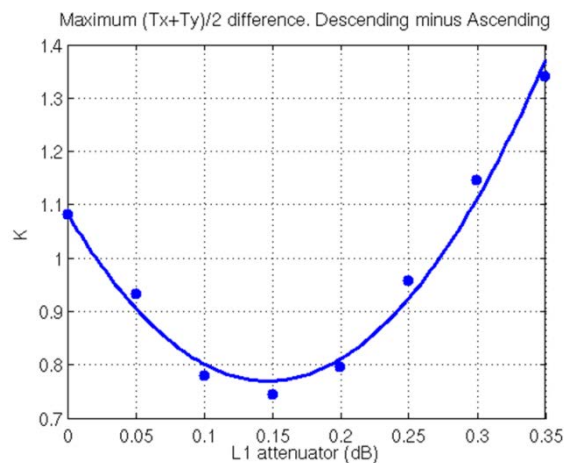
760  
761

As mentioned earlier, antenna losses  $L1+L2$  could not be accurately measured on ground due to set up  
762 uncertainties, and instead, were characterized in orbit, for V620 of the Level-1 processor, thanks to the  
763 more favourable external calibration manoeuvres (refer to Figure 16). The final estimation of the  
764 antenna loss was carried out after removing those instances with detector voltage fluctuations as the  
765 examples in Figure 13. Once the total loss had been measured, a second step, critical for the temporal  
766 stability, was performed, that of determining the  $L1$  and  $L2$  separately for the Noise Injection  
767 Radiometers. As the external calibration could only provide the ensemble loss, the split into its two

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768 contributions was done by minimizing the orbital variations, **taking as metrics the descending minus**  
769 **ascending deviations** over the ocean with respect to a forward model. Such exercise resulted in a  
770 minimum of orbital variation for  $L1=0.15$  dB as illustrated in Figure 27, the remaining of the total loss  
771 measured in orbit being assigned to L2. The values used in V620 for the L2 antenna loss of each  
772 receiver and polarization are constant (do not change over time) and equal to their average value  
773 obtained across all external calibrations, using the optimized  $L1=0.15$  dB. **This procedure to split the**  
774 **antenna losses in its two contributors L1 and L2 was based on the much stronger relationship of**  
775 **L1 with orbital variations than L2, L1 being tightly influenced by the skin temperature of the**  
776 **antenna.**

777



778

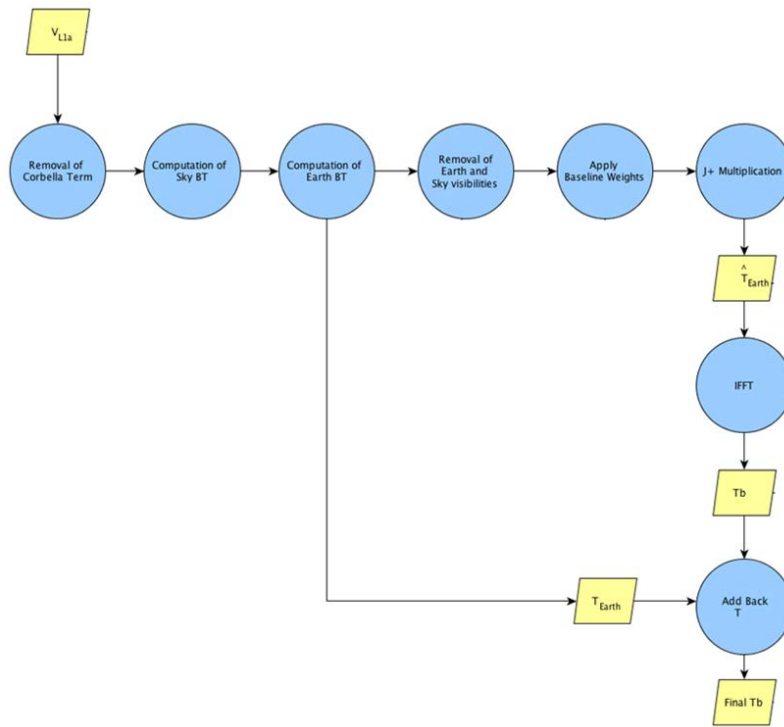
779 Figure 27: Split of NIR antenna losses achieved by minimizing the descending minus ascending pass  
780 difference of the deviations of the first Stokes parameter (divided by 2) from a model of the brightness  
781 temperature of the ocean (in Kelvin)  
782

### 783 6.1.7 Improved Gibbs-1 Image Reconstruction

784

785 In addition to being fully polarimetric, as commented earlier, the image reconstruction of V620,  
786 depicted in Figure 28, brings two other main improvements over V505. The first one is the  
787 implementation of the whole processing on hexagonal grids, as opposed to rectangular ones, to avoid  
788 interpolation errors. The second consists of a new Gibbs-1 ‘model approach’ by which the image  
789 reconstruction is performed over residual visibilities resulting from subtracting the estimated  
790 contribution of the Corbella term, the sky and the Earth, from the measured visibilities.

791



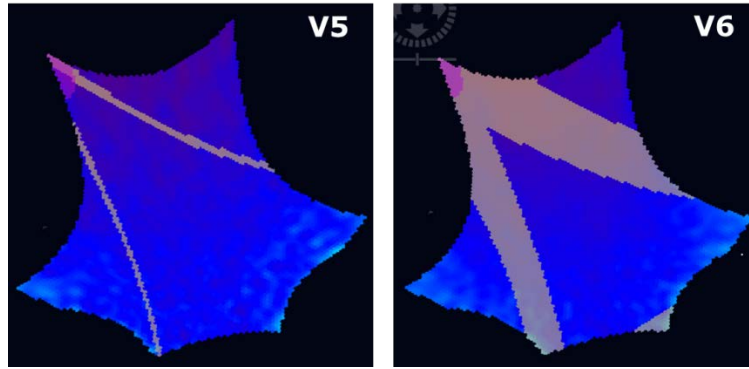
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Figure 28: Improved Gibbs-1 model approach for image reconstruction as implemented in V620.

### 795 6.1.8 Improved Removal and Flagging of Sun and RFI sources

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798 Since version V505, the algorithm for the Sun removal has been continually improved. In V620 the  
799 Sun brightness temperature is estimated through a 4-point **spatial** interpolation instead of taking the  
800 single closest neighbour point, this allowing a much finer positioning of the energy and thus, a  
801 reduction of the Sun tails. Several enhancements were also done to the flagging of the pixels affected  
802 by the Sun brightness temperature, in particular, the dynamic adjustment of the width of the tails of the  
803 real Sun and its aliases for every snapshot, and the correct flagging of all tails (one of the tails was not  
804 properly flagged in the previous version), as shown in Figure 29. As for the flagging and treatment of  
805 RFI sources and their impact, the improvements are reported in (Oliva et al., in press) (Khazaal et al.,  
2014) (Daganzo-Eusebio et al., 2013).



806  
807

808 **Figure 29: SMOS extended alias-free field of view showing, in grey, the pixels which are flagged**  
809 **due to the Sun tails (real Sun is on the right, with 2 alias on left top and left bottom). The width of**  
810 **the tails with flagged pixels has been increased with the improved Sun tails dynamic flagging of**  
811 **V620, and all tails are now flagged (note the lower right corner tail in V6).**  
812

## 813 **6.2 Performance of Level-1 Processor Version V620**

814 This section is devoted to present the performance of the currently operational version of the Level-1  
815 processor, V620, in comparison with the earlier version V505. The performance has been assessed over  
816 the entire data set of the two reprocessing campaigns, one with each processor version. The  
817 quantification of the performance has been carried out following some defined metrics, as depicted in  
818 Figure 30, comprising: calibration parameters, temporal stability (orbital, seasonal and yearly) of both  
819 the antenna and the brightness temperatures, systematic spatial errors in the images (bias and ripple),  
820 Sun correction, land-sea contamination and random noise. The Cold Sky, the Pacific Ocean and  
821 Antarctica are the targets where the metrics have been evaluated, **by comparing observations to a**  
822 **surveyed map of the sky, a radiative transfer model of the ocean (Ocean Target Transformation**  
823 **or OTT) or simply their average value over Antarctica (Ice Target Transformation or ITT)**  
824 **respectively.**

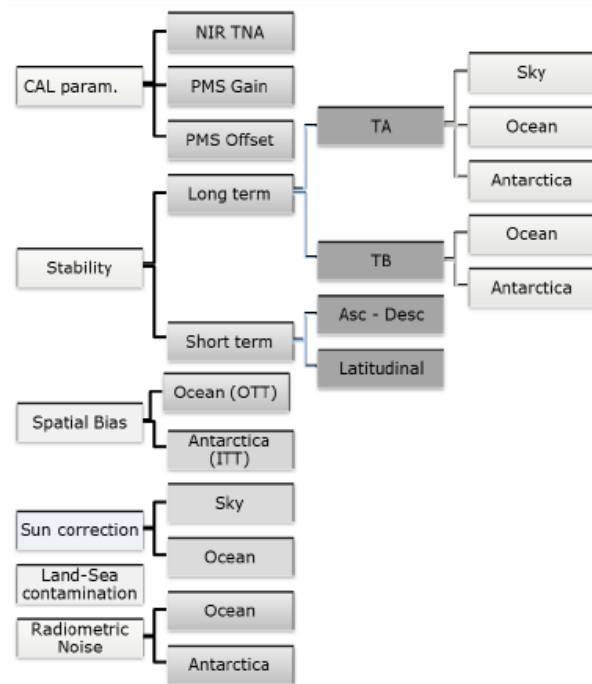
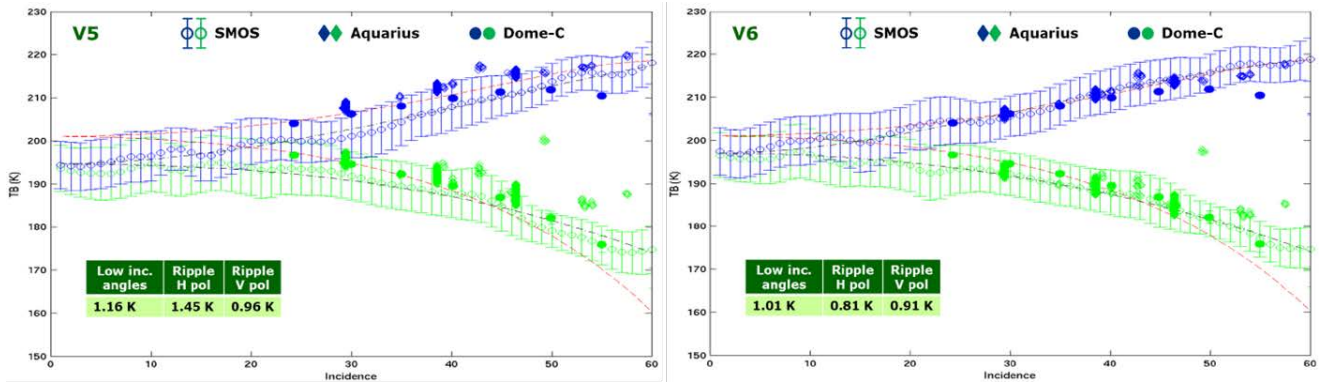


Figure 30: Overview of the metrics used to evaluated the performance of the Level-1 processor.

### 6.2.1 Removed Negative Slope at High Elevation Angles

V505 suffered from a negative slope in the dependence of the brightness temperature with elevation angle, in a way that measurements at low incidence were too cold when compared against modeled values. This is illustrated in the left panel of Figure 31 which shows the brightness temperature measured by SMOS, in horizontal and vertical polarizations, as a function of incidence angle over Dome-C station in Antarctica (Macelloni et al., 2013). As reference, the red dashed line represents a model, the diamonds are Aquarius observations and the solid circles ground measurements. Besides the discrepancy at high incidence angles which is expected due to the unavoidable sky contamination in the ground observations (through the part of the main lobe above the ice horizon), the SMOS brightness temperatures are clearly colder than the model towards 0° incidence. The enhancements implemented in V620 described earlier (in particular in Section 6.1.7) resulted in a substantial reduction of such cold trend at low incidence, as can be seen in the right panel of Figure 31. It is also worth noting that, for V620, the ripples along incidence angle are smaller and the match with Aquarius and Dome-C ground radiometer is better.

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Figure 31: Performance of V505 (left) and V620 (right) with incidence angle over Antarctica.

847

### 6.2.2 Lower Spatial Ripple

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Although there is a limit to how much the spatial ripple can be removed as explained before, Figure 32

850

shows that V620 achieves about 0.2 K lower spatial ripple in both polarizations than V505, **thanks to**

851

**the improvements in Sections 6.1.1 through 6.1.4 and 6.1.7.** The root mean square spatial ripple of

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V620 over most of the Extended Alias-Free Field of View is therefore of about 1.5 and 2.0 K for X and

853

Y polarizations respectively, evaluated over the ocean. It has to be noted that the bias of V620 is

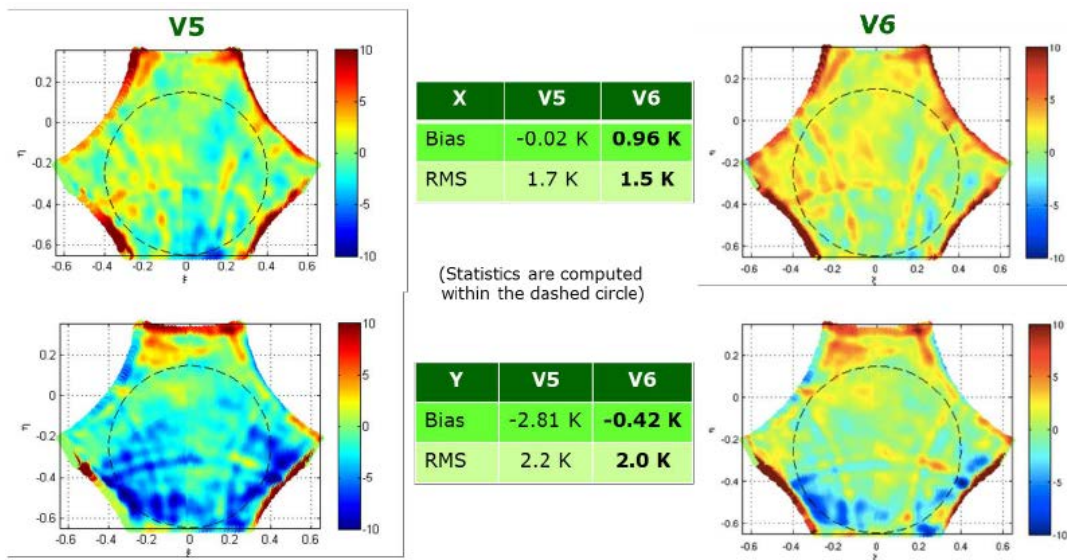
854

warmer than that of V505, overshooting almost 1 K in X polarization above the forward ocean model.

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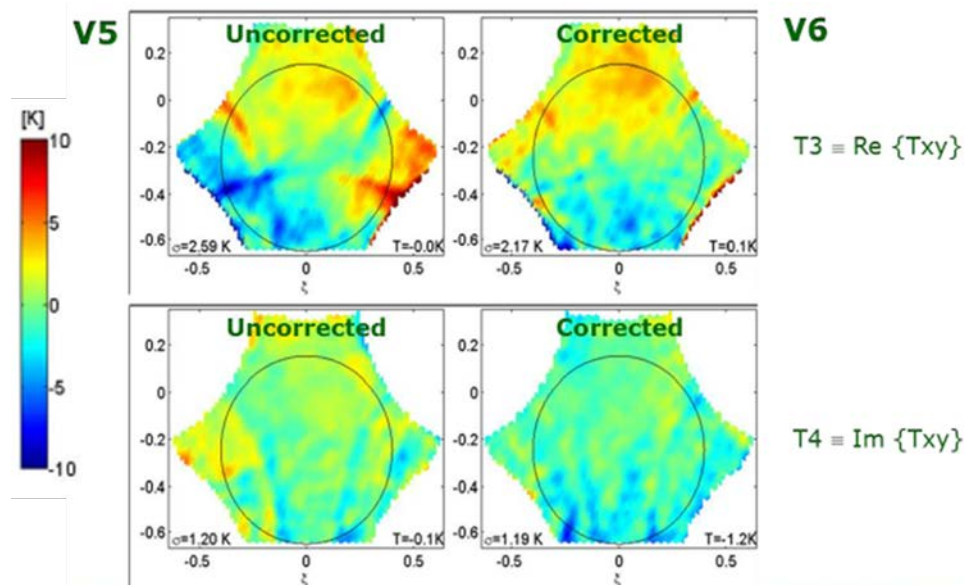
### 6.2.3 Improved Stokes-3 and Stokes-4 Parameters

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860 The fully polarimetry operation of V620 including the cross-polar patterns with corrected sign  
861 convention (**Section 6.1.1**) and relative phase between polarizations (**Section 6.1.2**) yield significant  
862 improvements in the Stokes-3 and Stokes-4 parameters (Lin et al., 2013). This is clearly shown in  
863 Figure 33, where more uniform residuals against the ocean model are obtained with V620 (right  
864 column) than with V505 (left column). The cleaner polarimetric brightness temperatures provided by  
865 V620 allows accurate estimation of the ionospheric total electron content and the Faraday rotation  
866 angle directly from SMOS observations (Corbella et al., 2015).

867



868

869

870 Figure 33: Residual Stokes-3 (top) and Stokes-r4 (bottom) of V505 (left) and V620 (right) over Ocean

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(colour scale in Kelvin)

#### 872 6.2.4 Removed Latitudinal Drift

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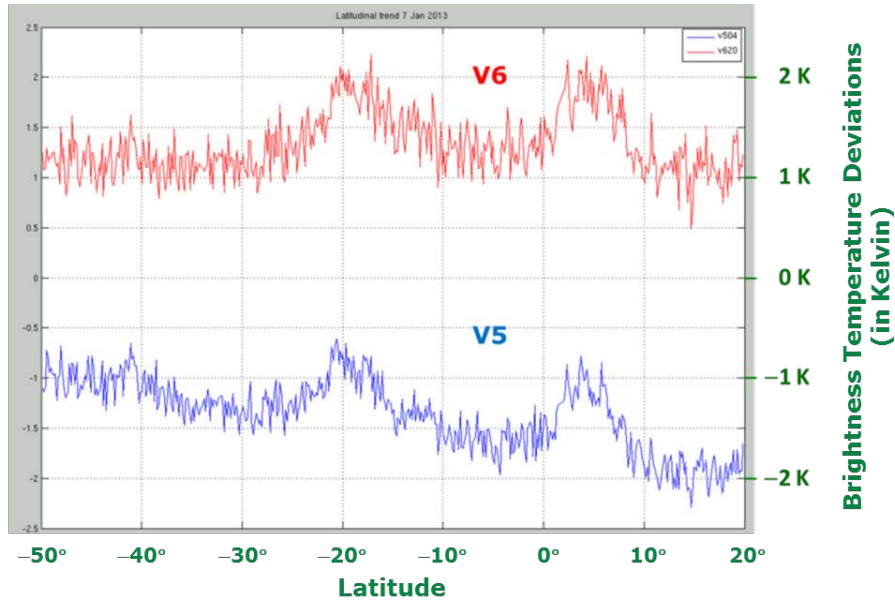
874 The red line in Figure 34 shows the deviation of the Stokes-1/2 parameter (average of the vertical and  
875 horizontal brightness temperatures) against a forward model along a strip of the Pacific Ocean from 50°  
876 South to 20° North, for version V620. **The plot is for one orbit, but the general behaviour is**  
877 **systematic, with some seasonal variations as shown in the Hövmoller plot in the right panel of**  
878 **Figures 6.** Besides some variations **on each side** of the equator which can be related to geophysical  
879 signals (**unmodeled rain and surface roughness effects**), there is no tendency with latitude as  
880 expected, **thanks mainly to the improvement in Section 6.1.6.** This is an important improvement



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881 over V505 processor, for which the Stokes-1/2 parameter, shown by a blue line on the same figure, did  
882 have a significant slope over latitude.

883



884

885

886 **Figure 34: Latitudinal trend of the brightness temperature deviation from an ocean model,**  
887 **averaged in the alias-free field of view, of V505 (blue line), and corrected value in V620 (red line).**  
888

889

### 889 **6.2.5 Reduced Orbital Variations**

890

891 The reduced orbital variations of V620 in the Stokes-1/2 parameter over the ocean have already been  
892 introduced and shown in the right panel of Figure 7. What is presented in the left panel of the same  
893 figure is the strong orbital variations, of about 2 K peak to peak, that V505 exhibited between  
894 ascending and descending passes in the same Stokes-1/2 parameter. These pronounced variations have  
895 been reduced in V620 mostly thanks to the optimization and use of the antenna losses measured in orbit  
896 **(Section 6.1.6) and the use of only the most stable NIR unit (Section 6.1.5).**

### 897 **6.2.6 Improved Yearly and Seasonal Stability**

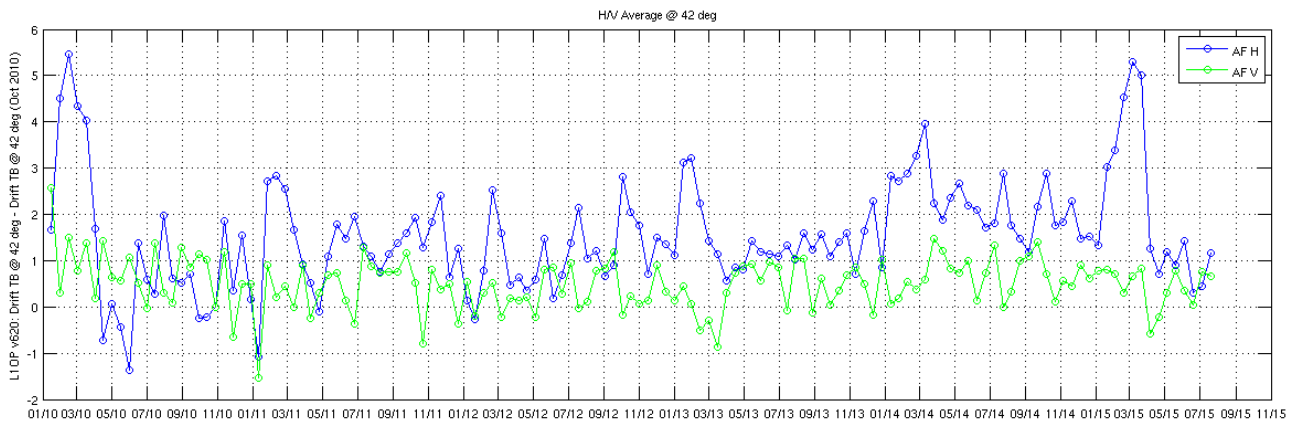
898

899 Figure 6 shows the latitudinal, seasonal and yearly variations of the Stokes-1/2 parameter over the  
900 ocean. The left plot, for V505, has strong variations with latitude which, as discussed above, have been  
901 removed in V620, shown in the right panel. V505 presents also significant seasonal variations, seen as  
902  $\pm 2$  K alternating blue and red bands in the corresponding **Hovmöller** plot. These strong seasonal

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903 fluctuations have been greatly reduced in V620, down to  $\pm 0.4$  K, **thanks to the improvement in**  
904 **Section 6.1.6.** In addition, a  $-0.18$  K/year cooling yearly drift is present in V505 data, which has  
905 disappeared in V620, for which the trend is of only  $-0.003$  K/year. The long term stability of V620 has  
906 further been verified over Antarctica, as presented in Figure 35, which shows the vertical and  
907 horizontal brightness temperature at  $42^\circ$  incidence angle measured by SMOS for a period of over 5  
908 years. The fluctuations in horizontal polarization are known to be due to surface roughness fluctuations,  
909 being the vertical polarization the best indicator to check for the instrument long term stability. As it  
910 can be seen, the yearly drift is indeed negligible.

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Figure 35: Over 5-year record of vertical and horizontal SMOS brightness temperatures at  $42^\circ$  incidence angle within the Alias-Free Field of View over Antarctica.

## 916 7 CURRENT INVESTIGATIONS

917

918 Despite the major improvements brought in by V620 version of the SMOS Level-1 processor over the  
919 previous version V505 which have been presented, several remaining challenges require further work,  
920 in particular, the residual orbital and seasonal variations, including the eclipse season, the further  
921 reduction of the spatial ripple and the Sun tails, not to forget the detection and mitigation of RFI  
922 effects. This section presents some of the on-going investigations to advance in these directions.

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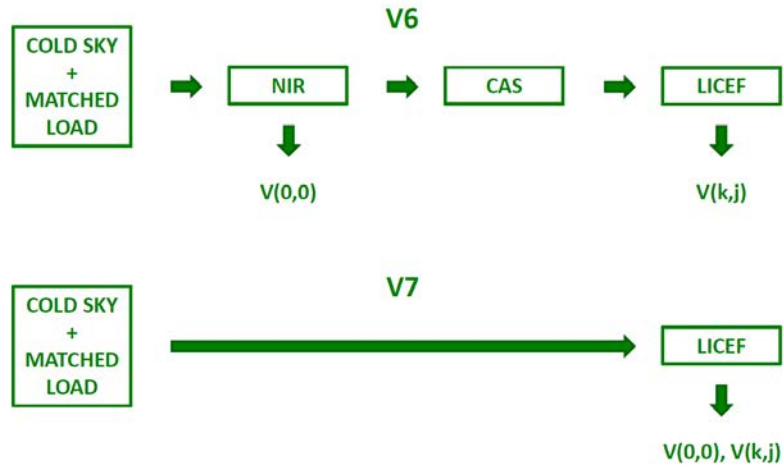
### 924 7.1 Level-1 Processor Version V700

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925 As mentioned in the introduction a new version of the processor, V700, has already been delivered,  
926 whose major change with respect to the currently operational version V620 is that it can run a so-called  
927 ‘All-LICEF’ mode. The implementation of the All-LICEF mode has, as objective, the simplification of  
928 the overall calibration procedure of the MIRAS instrument, to achieve better calibrated measurements  
929 and to make faster progress in correcting some of the remaining spatial errors and temporal variations.  
930 The concept behind the All-LICEF mode is depicted in Figure 36. The upper part shows the man  
931 calibration steps followed in V620: first the NIR radiometers are calibrated using the Cold Sky and an  
932 internal matched load, to provide the visibility sample at the origin  $V(0,0)$ ; then NIR is used to calibrate  
933 the noise diodes of the on-board Calibration System (CAS); finally the CAS system calibrates all the  
934 LICEF receivers of MIRAS. The lower part of Figure 36 shows instead, the much simpler and direct  
935 calibration approach of the All-LICEF mode: the LICEF receivers are directly calibrated using the Cold  
936 Sky and an internal matched load, and hence, the visibility sample at the origin can be measured by  
937 averaging all the calibrated antenna temperatures of the LICEF receivers. The straight advantage of the  
938 All-LICEF mode is therefore its extreme simplicity. This mode of operation was already envisaged  
939 well before launch, but only the accumulation of 5 years of flight experience allows to properly  
940 assessed its performance. It has to be noted that while the Noise Injection Radiometers are intrinsically  
941 more stable than the LICEF total power radiometers, the fact that there are 69 of the latter and only 3 of  
942 the former (in fact only the most stable NIR unit is currently used), offsets the final result in favour of  
943 the All-LICEF solution. To show one example, Figure 37 shows the temporal variations for V620 (left  
944 panel) and V700 in All-LICEF mode (right panel), and it can be noticed that the warming trend of 2014  
945 for V620 is not so pronounced in V700. Another illustration of the benefits of the All-LICEF mode is  
946 given in Figure 5, where the land-sea contamination correction based on the adjustment of the  
947 correlator efficiency coefficients explained earlier has been successfully tested with V700 (right panel)  
948 by comparison with the current V620 (left panel) where a significant coastal signature is still

949 observable. The All-LICEF branch of V700 is undergoing through a detailed assessment before it is  
 950 accepted for operational use.

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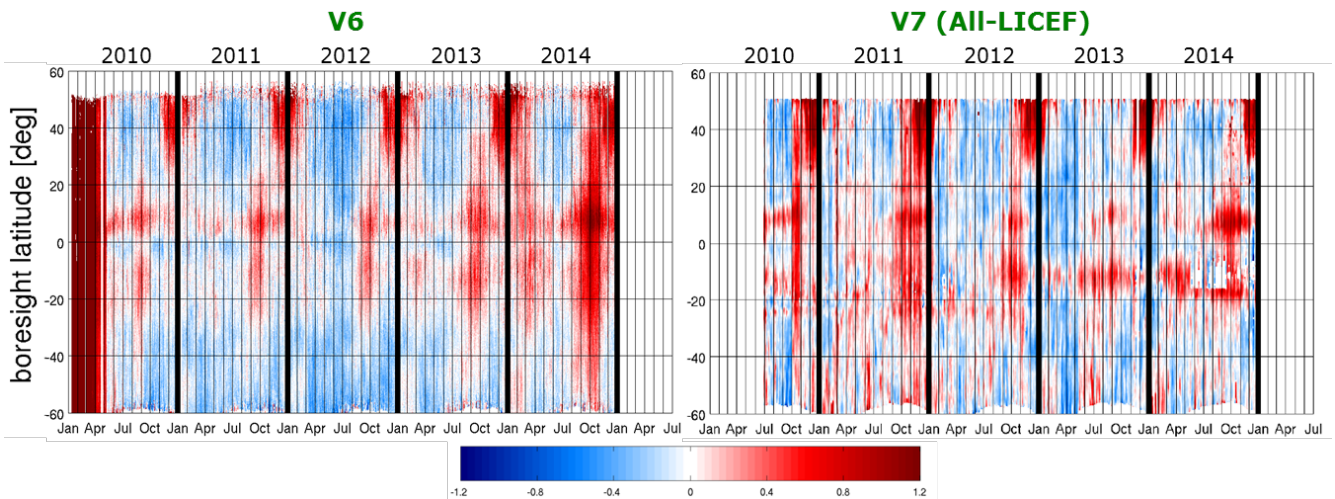
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Figure 36: All-LICEF mode

955 **7.2 Beyond V700**

956 Other improvements beyond version V700 are in the pipeline, namely the implementation of an  
 957 improved Gibbs-2 based image reconstruction algorithm, the further elaboration of improved spatial  
 958 ripple and Sun mitigation techniques, as well the reduction of residual temporal fluctuations.

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Figure 37: Descending-minus-Ascending pass Stokes-1/2 parameter over the Pacific Ocean, averaged in the alias-free area, with V620 (left) and V700 in All-LICEF mode (right) –colour scale in Kelvin–

965 **8 CONCLUSIONS**

966

967 SMOS mission has passed its extended life time of 5 years and continues to deliver good quality data,  
968 with both the MIRAS payload and the PROTEUS platform being in good health. The accumulated data  
969 record has allowed continued insight and improvement, from launch till today, of the MIRAS  
970 instrument and the Level-1 processor which transforms the visibility measurements into brightness  
971 temperature images. This paper has described the results of the detailed investigation of the calibration  
972 data and images, in various aspects, including instrumental behaviour and image reconstruction.  
973 Special focus has been given to the presentation of the performance of the currently operational version  
974 of the Level-1 processor, V620, and its improvements by comparison to the earlier version V505. V620  
975 is a fully polarimetric processor accounting for cross polar terms, with an enhanced image  
976 reconstruction technique and based on an improved in-orbit calibration approach which has resulted in  
977 overall cleaner and more stable brightness temperature images. Some hints have also been given about  
978 on-going investigations with a new processor version already delivered, V700, featuring the All-LICEF  
979 mode with a much simpler calibration strategy with potential to reduce the land-sea contamination.  
980 Finally, some future avenues, as the Gibbs-2 method, have been mentioned, addressing the still  
981 remaining challenges as the residual orbital and seasonal fluctuations, spatial ripple and Sun correction.

982 **9 ACKNOWLEDGEMENTS**

983

984 The authors would like to express their thanks to Mariano Kornberg with ESA, for providing the long  
985 series of instrument parameter values shown in the paper.

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