



# Single $\pi^0$ production off neutrons bound in deuteron with linearly polarized photons

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Received: 12 March 2021 / Accepted: 10 June 2021

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Communicated by Patrizia Rossi

**Abstract** The quasifree  $\vec{\gamma}d \rightarrow \pi^0 n(p)$  photon beam asymmetry,  $\Sigma$ , has been measured at photon energies,  $E_\gamma$ , from 390 to 610 MeV, corresponding to center of mass energy from 1.271 to 1.424 GeV, for the first time. The data were collected in the A2 hall of the MAMI electron beam facility with the Crystal Ball and TAPS calorimeters covering pion center-of-mass angles from  $49^\circ$  to  $148^\circ$ . In this kinematic region, polarization observables are sensitive to contributions from the  $\Delta(1232)$  and  $N(1440)$  resonances. The extracted values of  $\Sigma$  have been compared to predictions based on partial-wave analyses (PWAs) of the existing pion photoproduction

database. Our comparison includes the SAID, MAID and Bonn–Gatchina analyses; while a revised SAID fit, including the new  $\Sigma$  measurements, has also been performed. In addition, isospin symmetry is examined as a way to predict  $\pi^0 n$  photoproduction observables, based on fits to published data in the channels  $\pi^0 p$ ,  $\pi^+ n$  and  $\pi^- p$ .

## 1 Introduction

Knowledge of the  $N^*$  and  $\Delta^*$  resonance decay couplings to nucleons and photons is largely restricted to charged states. Increasing the body of neutron-target measurements will

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allow a more highly constrained study of neutral states and their  $n\gamma$  couplings. The four charge channels ( $\pi^0 p$ ,  $\pi^+ n$ ,  $\pi^- p$  and  $\pi^0 n$ ) of pion photoproduction can be described in terms of three isospin amplitudes. This gives the possibility of predicting properties of one channel based on sufficiently detailed measurements of the other three. The  $\pi^0 n$  channel is the least-studied and was the subject of this experiment.

Most existing  $\gamma n$  data are unpolarized and provide around 1900  $\pi^0 n$  photoproduction data points [1] (Table 1). The  $\vec{\gamma} n \rightarrow \pi^0 n$  beam asymmetry,  $\Sigma$ , was previously measured by the GRAAL Collaboration [9]. The beam asymmetry measures the relative strength of the production with respect to the plane of photon linear polarization. Their measurements covered beam energy,  $E_\gamma$ , from 703 to 1475 MeV, corresponding to a centre-of-mass energy,  $W$ , range from 1.484 to 1.912 GeV, just above the current results.

Recently, the A2 Collaboration at MAMI published high-quality unpolarized measurements for  $\pi^0$  photoproduction off a neutron below  $E_\gamma = 813$  MeV [3]. The present data extend the range of the previous GRAAL polarized measurements [9], for  $\pi^0 n$  photoproduction down to masses where the  $\Delta(1232)$  and  $N(1440)$  resonances can contribute significantly to polarization observables. Further A2 Collaboration measurements of the  $\pi^0 n E$  asymmetry, with longitudinal polarized target and circularly polarized photons, for  $E_\gamma = 216\text{--}1606$  MeV [11], extend previous A2  $\pi^0 n E$  measurements [10]. These data will provide the basis for better-constrained  $\gamma n$  decay amplitudes in the near future.

Apart from lower-energy inverse reaction  $\pi^- p \rightarrow \gamma n$  measurements, the extraction of the two-body  $\gamma n \rightarrow \pi^- p$  and  $\gamma n \rightarrow \pi^0 n$  observables requires the use of a model-dependent nuclear correction, which mainly comes from final-state interaction (FSI) effects. In several papers, the GWU-ITEP group has shown that the FSI corrections on unpolarized cross sections are less than 20% (see, for instance, Refs. [12–14]). As polarization asymmetries measure ratios of cross sections, FSI effects are expected to have a considerably smaller effect on these, including  $\Sigma$ , and will be comparable, or less than our quoted systematic uncertainties from experimental sources. In this publication,  $\Sigma$  for the neutron bound in a deuteron is presented uncorrected for potential FSI effects so as not to add any model dependence to the results.

The organization for this paper is as follows. In Sect. 2, details of the A2 experiment and detectors are given; Sect. 3 outlines the event selection; Sect. 4 reviews the background subtraction; Sect. 5 covers the determination of the photon asymmetry; and Sect. 6 outlines the dominant sources of systematic uncertainty; Sect. 7 outlines the PWA methods used in the fits and predictions compared to data. Finally, Sect. 8 presents the results and interpretation of the present A2  $\Sigma$  data.

## 2 Experiment

The reaction  $\vec{\gamma} d \rightarrow \pi^0 n(p)$  was measured at the Mainzer Microtron (MAMI) electron accelerator facility, in August 2016. The 1.5 GeV MAMI electron beam, incident on an aligned diamond radiator, produced a photon beam via coherent bremsstrahlung, with significant linear polarization up to photon energies of 610 MeV. The energy of the photon beam was measured using the Glasgow-Mainz Tagged Photon Spectrometer with a resolution of around 4 MeV. This spectrometer measured the position of the degraded post-bremsstrahlung electron after traversing a 1.8 T magnetic dipole field [15]. The position was measured by a plastic scintillator focal plane consisting of 353 elements. The energy of the detected electron and therefore also the energy of the photon, was deduced from this position.

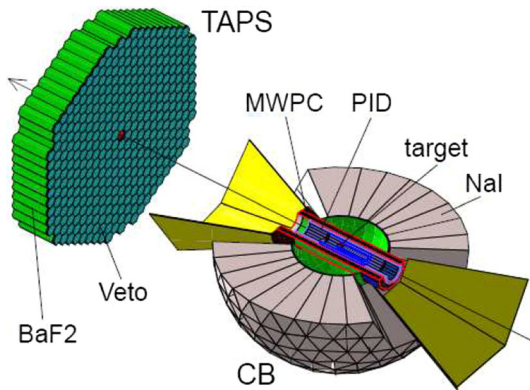
The photon beam interacted in a 10 cm long liquid deuterium target ( $LD_2$ ). The reaction products were detected in two calorimeters: the Crystal Ball (CB), a highly segmented array of 672 NaI(Tl) crystals arranged in a sphere centered on the target cell [16]; and the TAPS calorimeter, a forward wall of 366 BaF<sub>2</sub> and 72 PbWO<sub>4</sub> crystals arranged 1.5 m downstream from the CB center [17] (Fig. 1). The CB covered lab. frame angles  $21^\circ < \theta_{lab} < 159^\circ$  and TAPS approximately  $2^\circ < \theta_{lab} < 20^\circ$ . The  $LD_2$  target cell was surrounded by a Particle Identification Detector (PID) consisting of 24, 30 cm long plastic scintillators arranged in a cylindrical structure. This allowed separation of reactions with a scattered neutron from those with a proton. A Multi-Wire Proportional Chamber (MWPC) provided tracking information for charged particles which were not used for the all-neutral final state investigated here, however here they provided additional proton rejection. Charged particle identification was provided in the case of the TAPS detector by a thin plastic veto layer in front of each crystal. In addition, a 2.6 cm thick graphite cylinder was situated between the PID and the MWPC to be used as the analysing material for a nucleon polarimeter [18] and was not required for this analysis [19].

All simulations used in this analysis were performed with a full detector model using the Geant4 [20] toolkit.

The linear polarization of the photons was produced by coherent bremsstrahlung [21,22], with the electron beam scattered coherently from an aligned radiator. A thin diamond crystal (30  $\mu\text{m}$ ) with low mosaic-structure (i.e. few imperfections in the lattice) was used to minimize the energy smearing of the coherent spectrum arising from electron multiple scattering effects and crystal defects in the lattice [23]. The alignment of the diamond was carried out using the Stonehenge technique [24] with two orthogonal polarization plane orientations chosen to be at azimuthal angles of  $\pm 45^\circ$  with respect to the equatorial plane of the CB detector. To increase the degree of linear polarization a 2 mm diameter Pb-collimator was installed 2.5 m downstream of the radi-

**Table 1** Published data for  $\gamma n \rightarrow \pi^0 n$  reaction as given in the SAID database [33]: 1st column is the observable, 2nd column is the number of energy bins, 3rd column is the number of data points

Observable	Nexp	Ndata	$E_\gamma(\text{min})$ (MeV)	$E_\gamma(\text{max})$ (MeV)	$\theta(\text{min})$ ( $^\circ$ )	$\theta(\text{max})$ ( $^\circ$ )	Laboratory	Refs.
$d\sigma/d\Omega$	9	9	208	373	147	148	MAMI	[2]
	27	492	290	813	18	162	MAMI	[3]
	40	43	299	889	70	130	Tokyo	[4]
	49	931	446	1427	32	162	MAMI	[5]
	42	42	455	905	45	143	Tokyo	[6]
	35	35	462	784	60	135	Frascati	[7]
	3	28	911	1390	3	91	SLAC	[8]
	$\Sigma$	27	216	703	1475	53	164	GRAAL
$E$	17	151	446	1427	46	154	MAMI	[10]



**Fig. 1** Set-up of the A2 experiment. CB shows a NaI(Tl) calorimeter, TAPS shows a BaF<sub>2</sub> and PbWO<sub>4</sub> calorimeter, PID shows a plastic scintillator detector for particle identification, MWPC are two cylindrical multiwire proportional chambers, and target shows the liquid deuterium target LD<sub>2</sub> (see text for details)

ator, enhancing the ratio of coherently to incoherently scattered photons that reached the LD<sub>2</sub>. The degree of linear polarization was determined by calibrating against the linearly polarized photon beam asymmetry for  $\pi^0$  production off the proton. This was measured for each photon energy bin and compared to a recent SAID PWA solution including recent high statistics measurements in the same energy range [25]. The ratio of the measured asymmetry to the SAID values gave the photon polarization for each energy bin. The analysis of the proton asymmetry was performed in the same manner as the neutron asymmetry described here. The resulting photon polarization ranged from 15% at  $E_\gamma = 390$  MeV to a maximum of 55% at 610 MeV [19].

### 3 Data analysis

The photon asymmetry,  $\Sigma$ , for the reaction  $\vec{\gamma} d \rightarrow \pi^0 n(p)$  has been measured for beam energies in the range 390–

**Table 2** A summary of the loose cuts applied to the data before the sPlots fits

Variable	Cut range	Units
Tagged time	$-80 < t_{\pi^0} < 20$	ns
Coplanarity	$-50 < \Delta\phi < 50$	Degree
Missing mass	$1850 < M_{\text{miss}} < 2300$	MeV/c <sup>2</sup>
Cone angle	$0 < \theta_{\text{Cone}} < 0.5$	Radian
Invariant mass	$80 < M_{\text{inv}} < 200$	MeV/c <sup>2</sup>
Spectator momentum	$0 <  P_{\text{spec}}  < 200$	MeV/c

610 MeV and a center-of-mass (c.m.) production angle,  $\theta$ , range of 49°–148°. The semi-inclusive final state of interest included the recoiling neutron and  $\pi^0$ , and omitted the spectator proton. It was identified by detecting three neutral particles, two  $\gamma$ s stemming from the decay of the  $\pi^0$ , and a neutron as the third. The energy of the three particles was measured by the CB and TAPS calorimeters which, in coincidence with a tagged photon, allowed the reaction to be reconstructed. The classification of a neutral state was made if there were no hits in the PID or MWPC detectors for the CB or the TAPS veto layer. The spectator proton was not considered as it typically did not have sufficient energy to reach the calorimeters. The  $\pi^0$  was reconstructed from the combination of two of the three particles detected. All combinations were considered, and any incorrect  $2\gamma$  combinations were removed, either by subsequent cuts or by background subtraction.

Preliminary cuts were placed on a number of variables. These were guided by simulated signal and background channels to ensure no actual signal was lost and are given in Table 2. The missing mass was reconstructed using the mass of the missing 4-momentum defined as:

$$P_{\text{miss}} = P_{\text{beam}} + P_d - P_{\pi^0}, \tag{1}$$

where  $P_{\text{beam}}$  is the 4-momentum for the incident photon,  $P_d$  is the 4-momentum of the stationary deuterium target and

$P_{\pi^0}$  is the 4-momentum of the detected  $\pi^0$ . The mass of this missing 4-momentum gave a peak distributed around the summed mass of the two nucleons with some extra smearing from initial Fermi motion within the deuteron. When this initial momentum was low, as was generally the case, the resulting spectator proton momentum was also low and the reaction was approximately two-body with the  $\pi^0$  and participant nucleon being nearly coplanar in  $\phi$ :

$$\Delta\phi = \phi_{\pi^0} - \phi_{n-180^\circ} \sim 0^\circ, \quad (2)$$

where  $\phi_{\pi^0}$  is the reconstructed azimuthal angle of the  $\pi^0$  and  $\phi_{n-180^\circ}$  the azimuthal angle of the detected neutron after rotation by  $180^\circ$  around the z-axis.

The detection of the  $\pi^0$  and the neutron allowed the construction of the difference between the detected nucleon polar angle and the nucleon polar angle reconstructed from the  $\pi^0$  assuming a stationary initial state neutron. This gives the definition of the ‘‘Cone Angle’’,  $\theta_{\text{Cone}}$ .

The momentum of the spectator proton was also used to distinguish the quasi-free final state. It was calculated via:

$$P_{\text{spec}} = P_{\text{beam}} + P_d - P_{\pi^0} - P_n, \quad (3)$$

where  $P_n$  is the 4-momentum of the detected participant nucleon. The magnitude of momentum for the participant nucleon was calculated using conservation of momentum and energy in the three-body final state, using the measured  $\pi^0$  momentum and the neutron direction, as given by its cluster hit position in the CB.

## 4 Background subtractions

To further isolate the true  $\pi^0 n$  final state, several sources of background had to be subtracted from the selected events. The main sources of these backgrounds were: random electrons in the photon tagger; background to the two photon combination giving the  $\pi^0$ ; and other reactions producing the same detected particles as the  $\pi^0 n(p)$  reaction. The sPlot technique was used to remove these background events using a separate discriminatory variable for each source to produce event-by-event weights termed sWeights, for full details see [26]. The sWeights are normalized using the relative yields and covariance matrix of the signal and background derived from the fits. Weights corresponding to regions of high background are negative and effectively subtract off this contribution to the distribution. This is similar to how a ‘‘sideband subtraction’’ method works, but is more generally applicable. Consecutive fits were performed applying the sWeights from the previous fit. An sPlot fit to the discriminatory variables using appropriate probability density functions (PDFs) derived from simulated event samples, determined the yields of the different event species. The covariance matrix of this yield fit was then used to calculate the sWeights associated

with each event in the fitted sample. Including these sWeights in the subsequent observable fits allowed determination of the photon asymmetry for our signal. The fits to the discriminatory variables are described in the following sections.

### 4.1 Random Tagged Photons

Random coincidences with background electrons in the photon tagger were removed via the coincidence time between the  $\pi^0$  and the tagged beam photon as given by:

$$t_{\text{coin}} = \frac{t_{\gamma 1} + t_{\gamma 2}}{2} - t_{\text{tagger}}, \quad (4)$$

with the time of the electron in the tagger,  $t_{\text{tagger}}$  and the time in the calorimeters of the 2 photons  $t_{\gamma 1,2}$ . This resulted in a timing distribution strongly peaked at zero with a flat random background, as shown in Fig. 2 for the bin  $E_\gamma = 610$  MeV and  $\cos\theta_{CM} = 0.05$ . In this case, a Gaussian PDF was used for the signal with a uniform background function.

### 4.2 Background in $2\gamma$ invariant mass

Background to real  $\pi^0$  decays in the two  $\gamma$  invariant mass distribution can arise from a wrong combination of the three neutral clusters, or multiple clusters created by one actual particle. These background sources will not give a peaking structure in the invariant mass distribution and were thus subtracted using the sPlot technique. The  $\pi^0$  signal PDF was taken from a histogram template of simulated events, while the background was modelled by a third order Chebyshev polynomial. An example fit is shown for the bin  $E_\gamma = 610$  MeV and  $\cos\theta = 0.05$  in Fig. 3. In this mass range, the  $\pi^0$  signal was typically around 90% of the total events.

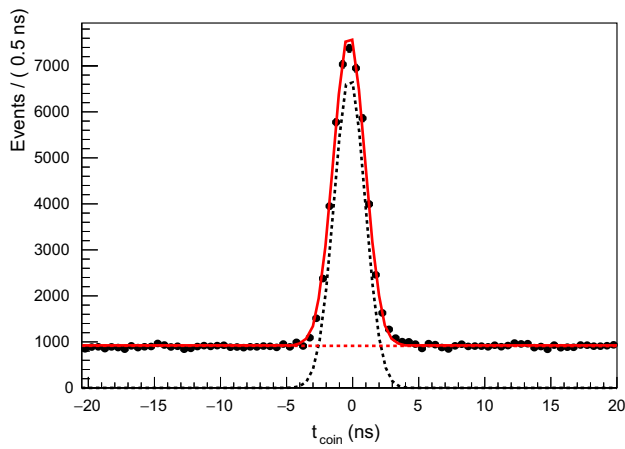
### 4.3 Background to the $\pi^0 n(p)$ final state

Background to the final state may come, for example, from events in which more than one pion is produced. This background is reduced with the loose cuts given in Table 2. To determine the sWeights for subtracting the residual background, the coplanarity given in Eq. 2 was used. The signal PDF shape was given by simulated data and the background by a second order Chebyshev polynomial. The resulting fit, for the bin  $E_\gamma = 610$  MeV and  $\cos\theta = 0.05$ , is shown in Fig. 4.

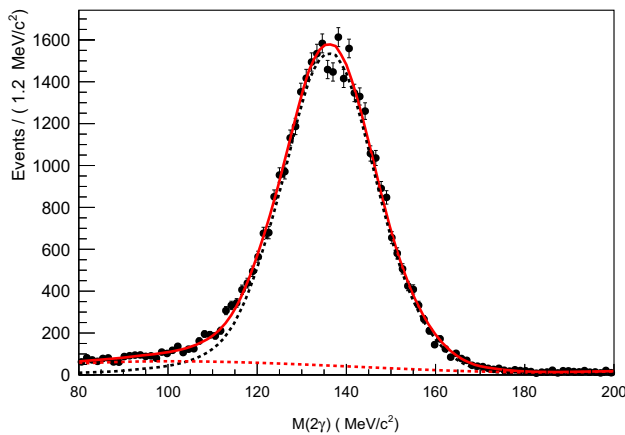
## 5 Determination of the photon asymmetry

The photon asymmetry,  $\Sigma$ , quantifies the effect of the polarization of the beam on the excitation of the neutron and its





**Fig. 2** The timing coincidence spectra between the photon beam tagger and the calorimeters. Black points are data; red solid line is full fit result; dashed black is signal Gaussian function; dashed red is flat background function. This fit was used to produce weights to subtract the random background events



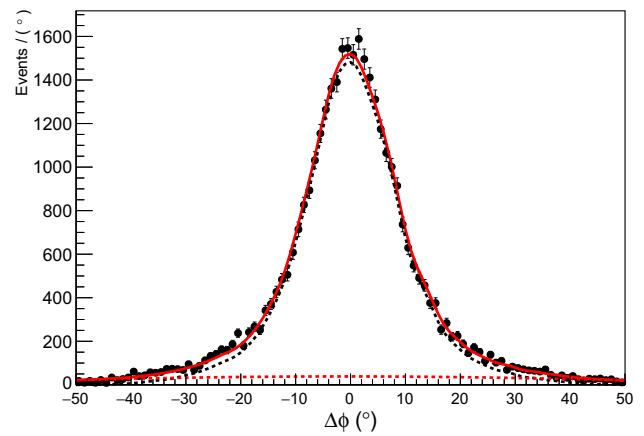
**Fig. 3** The invariant mass of the two detected  $\gamma$ s. Black points are data; red solid line is full fit result; dashed black is simulated signal function; dashed red is third degree Chebyshev polynomial. This fit was used to produce weights to subtract background events that did not have a  $\pi^0$

subsequent decay to a pion and nucleon. With a linearly polarized photon beam the differential cross section is

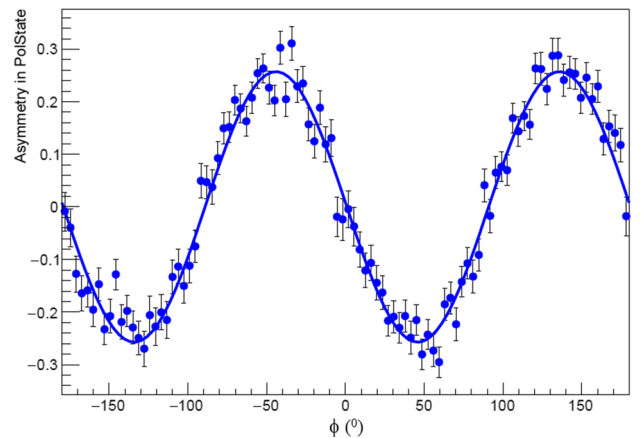
$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega}\right)_0 (1 + P_L \Sigma \cos 2\phi), \tag{5}$$

where  $\phi$  is the azimuthal angle of the meson production plane relative to the plane of linear polarization and  $P_L$  is the degree of linear polarization. Rotating the orientation of the diamond radiator allowed the plane of linear polarization to flip between  $-\frac{\pi}{4}$  and  $\frac{\pi}{4}$ . Shifting the plane by  $\frac{\pi}{2}$  effectively flips the sign of the asymmetry giving two polarization states  $P_S = \pm 1$ .

To extract  $\Sigma$  from the measured  $\phi$  distributions unbinned extended maximum likelihood fits were performed. The fit function was given by



**Fig. 4** The coplanarity between the  $\pi^0$  and detected neutron. Black points are data; red solid line is full fit result; dashed black is simulated signal function; dashed red is second degree Chebyshev polynomial. This fit was used to produce weights to subtract events that did not originate from the  $\vec{\gamma} d \rightarrow \pi^0 n(p)$  final state



**Fig. 5** Visualisation of asymmetry constructed from extracted  $\Sigma$  values as defined in Eq. 8. Blue points are the data and blue solid line is the result of the maximum likelihood fit

$$F(\Sigma : \phi, P_L, P_S) = 1 + \Sigma P_S P_L \cos(2\phi + \phi_0), \tag{6}$$

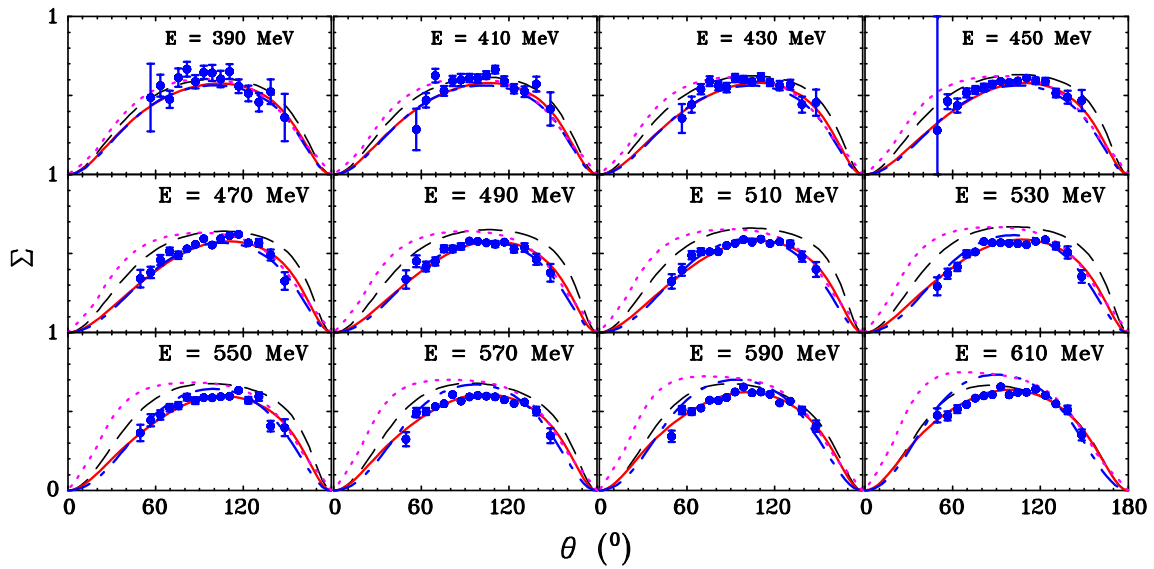
with  $\phi_0$  determined from fits to be  $95.8^\circ$ .

The negative log likelihood function given by

$$-\ln L = -\sum_i^N w_i \ln F(\Sigma : \phi_i, P_{L,i}, P_{S,i}) + B(\Sigma), \tag{7}$$

was minimized using Minuit as part of the ROOFIT [27] library. Here,  $N$  is the number of data events in the  $E_\gamma$  and  $\theta$  bin, while subscript  $i$  refers to the value of the variables for a given event. In particular,  $w_i$  represents the value of the sWeight used to subtract background events from the likelihood summation.

The PDF normalization integral term  $B(\Sigma)$  was determined by Monte-Carlo integration using simulated data. For this,  $P_S$  and  $P_L$  values were randomly chosen to match the fluxes and polarization degrees of the real data. This effec-



**Fig. 6**  $\Sigma$  for  $\vec{\gamma}n \rightarrow \pi^0n$  vs. pion production angle  $\theta$  in c.m. frame: A2 (blue filled circles); fit: SAID MU22 (red solid curves), SAID MA19 [3] (blue dash-dotted curves), Bonn–Gatchina BG2014-02 [28] (magenta dotted curves) and MAID2007 [29] (black dashed curves). Only angle-

dependent statistical uncertainties are shown for all data. Each plot corresponds to a 20 MeV wide bin in  $E_\gamma$  with the central value given in each. Only SAID MU22 has been fit to the new data

tively corrected for second order systematic effects due to differences in polarization state luminosity and degree of polarization.

Although the values for  $\Sigma$  were extracted using the event based maximum likelihood fits, we demonstrate how well the result matches the data by constructing an asymmetry in polarization state for our fitted PDF and compare it to the data asymmetry after background subtraction. For the PDF the plotted asymmetry is calculated as,

$$A(\Sigma_r : \phi) = \frac{F(\Sigma_r : \phi, P_L, +1) - F(\Sigma_r : \phi, P_L, -1)}{F(\Sigma_r : \phi, P_L, +1) + F(\Sigma_r : \phi, P_L, -1)} \tag{8}$$

and is shown in Fig. 5, for the bin  $E_\gamma = 610$  MeV and  $\cos \theta = 0.05$ , as an illustration.  $\Sigma_r$  represents the resulting value for the fitted parameter, while the  $\pm 1$  labels the polarization state.  $P_L$  is integrated over the  $E_\gamma$  bin.

The fits were performed in bins of 20 MeV for  $E_\gamma$  and 0.1 in  $\cos \theta$ . Results are shown in Fig. 6 alongside solutions of various PWAs described in Sect. 7.

### 6 Systematic uncertainties

The dominant sources of systematic uncertainty in the results came from the linear polarization calibration and the background subtraction. The fractional difference between a simple cuts-based analysis and the sPlot method gave an estimate of the systematic uncertainty in background subtraction

**Table 3** A summary of the tight cuts applied to the data to produce a low background event sample

Variable	Cut range	Units
Tagged time	$-5 < t_{\pi^0} < 5$	ns
Coplanarity	$-30 < \Delta\phi < 30$	Degree
Missing mass	$1850 < M_{\text{miss}} < 2100$	MeV/c <sup>2</sup>
Cone angle	$0 < \theta_{\text{Cone}} < 0.3$	Radian
Invariant mass	$110 < M_{\text{inv}} < 160$	MeV/c <sup>2</sup>
Spectator momentum	$0 <  P_{\text{spec}}  < 200$	MeV/c

method for each point. For the cuts-based analysis a similar procedure to the full extraction of  $\Sigma$ , described above, was performed but, rather than use a sPlot subtraction, tighter cuts were placed on the discriminatory variables to identify a cleaner sample of  $n\pi^0$  events with residual backgrounds estimated by simulations to be around 3.6%. These cuts are summarised in Table 3.

While a cuts-only analysis retains some small amount of background, as illustrated in Figs. 3 and 4, the sPlot method is expected to remove all the background. Any error in the subtraction of the background by the weights-based method is expected to be less than the effect of not subtracting the background. Hence, the difference between  $\Sigma$  extracted from the cuts and sPlot subtracted results is used as a conservative estimate of this systematic uncertainty [19]. Values for this systematic uncertainty were calculated for each  $E_\gamma$  and  $\theta$  bin with a mean value of 3%. For the systematic uncertainty in

the degree of linear polarization there are two factors, first the uncertainty in the SAID solution for  $\Sigma$  on  $p\pi^0$ , which was used to determine the polarization, and was estimated to be 2%. Second, the uncertainty on our extraction of the  $p\pi^0$  asymmetries which had a main contribution from the background subtraction which was estimated in a similar manner to the  $n\pi^0$  background subtraction and found to be 4%. Adding these two factors in quadrature gives an overall 4.5% systematic uncertainty in our  $\Sigma$  results due to the linear polarization. Other uncertainties were found to be much smaller than these sources: acceptance effects, such as the neutron detection efficiency, which vanishes due to the polarization flip; polarization degree and luminosity asymmetries were incorporated into the likelihood fit; and unbinned fits were used, meaning there were no artifacts from binning the data.

## 7 Multipole analysis

The MAID [29], SAID [30] and Bonn–Gatchina [31] analyses use different fit formalisms to extract the partial-wave (multipole) amplitudes underlying different data-sets. Comparing the different resulting amplitudes gives an estimate of the systematic errors inherent in the process.

For the MAID analysis, which was completed in 2007, the most recent data-sets have not been included and this must be considered when making comparisons [29]. The MAID unitary isobar approach applies a Breit–Wigner resonance plus background model, guaranteeing unitarity up to the two-pion production threshold.

The Bonn–Gatchina method fits a wider range of reactions utilizing elements of the K-matrix and P-vector approaches [31]. Both pion and photo-induced reactions are included in a multi-channel fit. Reactions with three-body final states are included using an event-based likelihood fit. The elastic pion-nucleon reaction is fitted based on existing amplitudes. The various data types are fitted with the possibility of renormalization and weighting.

The SAID method is an extension of the Chew–Mandelstam K-matrix approach used to fit pion-nucleon elastic scattering and  $\eta N$  production data. The resonance spectrum is fixed from this fit [32] and only the photo-couplings are allowed to vary. This differs from the MAID and Bonn–Gatchina analyses, which can add new resonances to improve the agreement with data. The formalism has built-in cuts associated with the  $\pi\Delta$ ,  $\rho N$  and  $\eta N$  thresholds but only single-pion photoproduction data are fitted. Data have been weighted and renormalized in previous fits. No weighting and only renormalization at the one percent level was utilized in fitting the present set of  $\Sigma$  data.

For each angular distribution, a normalization constant ( $X$ ) and its uncertainty ( $\epsilon_X$ ) were assigned. The quantity  $\epsilon_X$  is generally associated with the normalization uncertainty (if

**Table 4**  $\chi^2$  per data point for new A2 data. Predictions are from the SAID fit including present data (MU22), an older SAID fit to existing data (MA19 [3]), the new SAID fit MUXX (no world  $\pi^0 n$  data used) and MAID2007 [29]

Solution	$\chi^2/(\pi^0 n \text{ data})$
MU22	275/189 = 1.46
MA19	743/189 = 3.93
MUXX	624/189 = 3.30
MAID2007	1151/189 = 6.09

known). The modified  $\chi^2$  function to be minimized is given by

$$\chi^2 = \sum_i \left( \frac{X\eta_i - \eta_i^{exp}}{\epsilon_i} \right)^2 + \left( \frac{X - 1}{\epsilon_X} \right)^2, \quad (9)$$

where the subscript  $i$  labels the data points within the distribution,  $\eta_i^{exp}$  is an individual measurement,  $\eta_i$  is the corresponding calculated value, and  $\epsilon_i$  represents the angular-dependent statistical uncertainty. The total  $\chi^2$  is then found by summing over all measurements. This re-normalization freedom is often important in obtaining the best SAID fit results. For other data analyzed in the fit, such as the total cross sections and excitation data, the statistical and systematic uncertainties were combined in quadrature and no re-normalization was allowed.

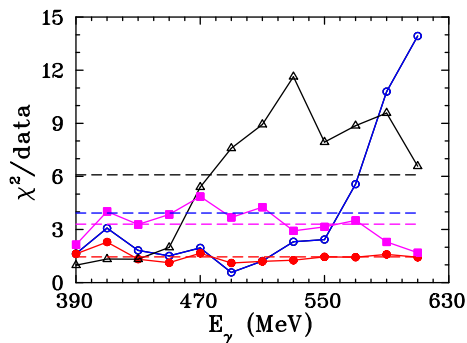
In fitting the present set of  $\Sigma$  data, an overall angle-independent systematic uncertainty of 1% was used for  $\epsilon_X$  in Eq. (9). The resulting values for  $X$  remained within  $\epsilon_X$  of unity on average.

A revised SAID multipole analysis has been completed, including the present set of  $\vec{\gamma}n \rightarrow \pi^0 n \Sigma$  data. This new global energy-dependent solution has been labeled as MU22. The overall fit quality of the present MU22 and previous MA19 [3] SAID fits is compared with the MAID2007 [29] solution in Tables 4 and 5. The inclusion of the present A2 data set provides a fit with significantly improved  $\chi^2/data$ , specifically at higher energies, in comparisons between the  $\pi^0 n$  fits and data ( $\chi^2/data$  for MA19 = 3.93 and 1.44 for MU22) as shown in Fig. 7 and Table 4. This demonstrates the influence of these asymmetry measurements with their small uncertainties. The overall comparison of the MA19 and MU22 solutions shows that the fit  $\chi^2/data$  values are essentially unchanged for  $\pi^0 p$  and  $\pi^+ n$  channels. The  $\chi^2$  per data point including all available data and the present A2 data for MA19 and MU22 (with MAID2007) is given in Table 5.

Additionally, an alternative MUXX solution was generated excluding all world  $\gamma n \rightarrow \pi^0 n$  data and show results in Fig. 8 and Table 4. The excellent comparison of the isospin-predicted  $\Sigma$  to the data strongly suggests the systematics in the new data are well under control. The comparisons

**Table 5**  $\chi^2$  per data point values for all charge channels covering the energy range from 155 to 1000 MeV. Fits as described in Table 4

Solution	$\chi^2/(\pi^0 p \text{ data})$	$\chi^2/(\pi^+ n \text{ data})$
MU22	13274/9534 = 1.39	7454/4039 = 1.85
MA19	12565/9534 = 1.32	7461/4039 = 1.85
MUXX	13171/9534 = 1.38	7259/4039 = 1.80
MAID2007	73638/9534 = 7.72	14599/4039 = 3.61
Solution	$\chi^2/(\pi^0 n \text{ data})$	$\chi^2/(\pi^- p \text{ data})$
MU22	2345/798 = 2.94	5879/3456 = 1.70
MA19	2649/798 = 3.32	5999/3456 = 1.74
MUXX	7639/798 = 9.57	5384/3456 = 1.56
MAID2007	4846/798 = 6.07	15365/3456 = 4.45

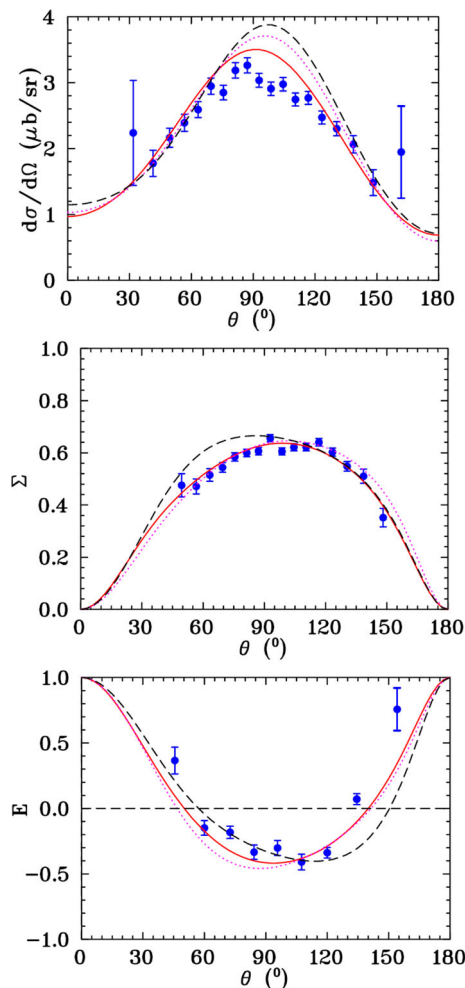
**Fig. 7** Comparison of  $\chi^2$  per data point for the previous SAID solution MA19 [3] (blue open circles) and MAID2007 [29] (black open triangles) applied to the present A2 data. Also shown are the new SAID solutions MU22 (red full circles) obtained after including the present A2 data and MUXX (magenta full squares) which did not use any  $\gamma n \rightarrow \pi^0 n$  data. In addition the fit  $\chi^2$  per data point values averaged over each energy bin  $E_\gamma$  are plotted, where the horizontal dashed lines are for the MU22 (red), MUXX (magenta), MA19 (blue), and MAID2007 (black) solutions. The solid lines connecting the points are included only to guide the eye

between the MUXX and the MAID solutions are interesting as both models use isospin symmetry to predict the  $\pi^0 n$  observables, based on the data from the other three charge channels which were available in 2007 and 2020.

## 8 Results and interpretation

A comprehensive set of  $\Sigma$  data for  $\vec{\gamma} d \rightarrow \pi^0 n(p)$  at 12 photon energies has been determined with the CB and TAPS spectrometers using a tagged photon beam at incident photon energies from 390 to 610 MeV. The present  $\Sigma$  data cover the resonance region from above the maximum of the  $\Delta$ -isobar to the Roper resonance.

The SAID MA19 [3], Bonn–Gatchina BG2014-02 [28] and MAID2007 [29] curves shown in Fig. 6 did not include the present A2 data in their fits. In addition, the MAID2007

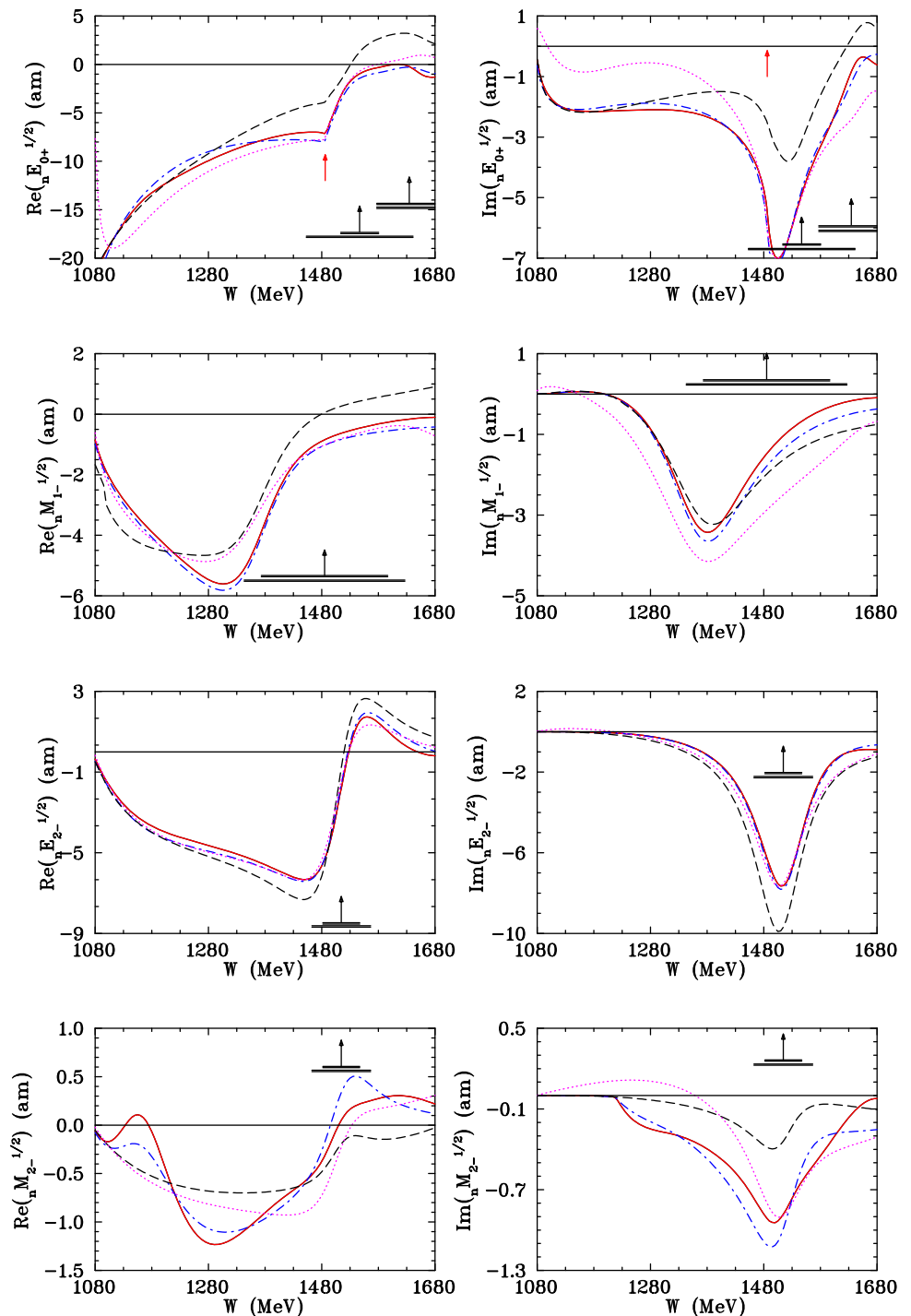
**Fig. 8**  $d\sigma/d\Omega$ ,  $\Sigma$  and  $E$  observables for  $\vec{\gamma} n \rightarrow \pi^0 n$  (blue full circles). All data were produced by the A2 Collaboration at MAMI:  $d\sigma/d\Omega$  at  $E_\gamma = 610$  MeV (top) are from Ref. [3],  $\Sigma$  at  $E_\gamma = 610$  MeV (middle) are present measurement, and  $E$  at  $E_\gamma = 603$  MeV (bottom) are from Ref. [10]. New SAID solutions MU22 (MUXX) are shown by red solid (magenta dotted) curves and MAID2007 [29] by black dashed curves

fit does not include measurements after 2007. MU22 includes all previous measurements and includes the A2 data. All fits and predictions agree well for the lowest energy where the  $\Delta$  resonance dominates. The angular distribution retains the shape of SAID MA19 up to the highest energies where some larger deviations become apparent.

In the fit, cross sections have larger angle-independent systematic uncertainties and the renormalization factor from Eq. (9) improves the description, but is not included in the plot. While the new data cover parts of the  $\Delta$  and Roper resonance regions, selected isospin multipoles are compared up to a photon energy of 1 GeV. The isospin 3/2 multipoles are taken as determined by the much larger proton-target database and only the isospin 1/2 neutron multipoles are shown in Fig. 9. Comparing MA19 to MU22, there are no sig-



**Fig. 9** Selected neutron multipole  $I = 1/2$  amplitude from threshold to  $W = 1.68$  GeV ( $E_\gamma = 1030$  MeV) for the states  $0^+$ ,  $1^-$  and  $2^-$ . The blue dash-dotted (red solid) curves correspond to the SAID MA19 [3] (new MU22 including present A2 data) solution. The magenta dotted (black dashed) curves give the Bonn–Gatchina BG2014-02 [28] (MAID2007 [29]). The vertical black arrows indicate Breit–Wigner (BW) mass ( $W_R$ ), and horizontal bars show full ( $\Gamma$ ) and partial ( $\Gamma_{\pi N}$ ) widths of resonances extracted by the BW fit of the  $\pi N$  data associated with the SAID solution SP06 [32]. Vertical red arrows show the  $\eta$  meson production threshold



nificant changes seen in the  $M_{1-}^{1/2}$  multipole which includes the  $N(1440)\frac{1}{2}^+ P_{11}$ . In addition, comparing imaginary parts of the multipoles connected to the nearby  $N(1535)\frac{1}{2}^- S_{11}$  ( $E_{0+}^{1/2}$ ) and  $N(1520)\frac{3}{2}^- D_{13}$  ( $E_{2-}^{1/2}$  and  $M_{2-}^{1/2}$ ), we see there is good agreement between the SAID and Bonn–Gatchina plots. The resonance couplings for the  $N(1440)$ ,  $N(1535)$

and  $N(1520)$  are expected to be in agreement with those reported in Ref. [3].

Exploring the effectiveness of isospin symmetry to predict  $\pi^0 n$  observables, in Fig. 8 the fit (MU22) is compared to predictions from MUXX and MAID2007 for  $d\sigma/d\Omega$ ,  $\Sigma$  and the double-polarization asymmetry  $E$ . At this energy, the qualitative features are generally reproduced, particularly for the current data.

As a final comment on the predictive ability of our fit, excluding all  $\pi^0 n$  data, a comparison of Tables 4 and 5 shows that fit MUXX is much less successful between the upper energy limit of the present experiment and 1 GeV in the photon energy. This change is due mainly to poor compatibility with GRAAL  $\Sigma$  data [9].

Our results for  $\Sigma$  for  $\vec{\gamma} d \rightarrow \pi^0 n(p)$  consist of 189 experimental points and are available from the SAID database [33], where systematic uncertainties for each bin have been added in quadrature.

**Acknowledgements** This work was supported in part by the UK Science and Technology Facilities Council (STFC Grants No. ST/L005719/1, ST/P004458/1, ST/T002077/1, ST/P004385/2, ST/V002570/1, ST/P004008/1 and ST/L00478X/2), the U. S. Department of Energy, Office of Science, Office of Nuclear Physics, under Awards no. DE-FG02-01ER41194, no. DE-SC0016583, no. DE-SC0016582, and no. DE-SC0014323. This work was supported by Schweizerischer Nationalfonds (Grants no. 200020-132799, no. 121781, no. 117601, and no. 113511), Deutsche Forschungsgemeinschaft (SFB Grant no. 443, No. SFB/TR 16, no. SFB 1044), DFG-RFBR (Grant no. 05-02-04014), European Community Research Infrastructure Activity (FP6), the U. S. DOE, U. S. NSF, and NSERC (Grant no. SAPPJ-2018-00020) Canada. This publication is part of a project that has received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement STRONG-2020-No 824093. We would like to thank all the technical and nontechnical staff of MAMI for their support.

**Data Availability Statement** This manuscript has associated data in a data repository. [Authors' comment: The data consisting of 189 experimental points and are available from the SAID database [33].]

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