

LETTER TO THE EDITOR

## Departure from the constant-period ephemeris for the transiting exoplanet WASP-12 b<sup>★,★★</sup>

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### ABSTRACT

**Aims.** Most hot Jupiters are expected to spiral in toward their host stars because the angular momentum of the orbital motion is transferred to the stellar spin. Their orbits can also precess as a result of planet-star interactions. Calculations show that both effects might be detected for the very-hot exoplanet WASP-12 b using the method of precise transit-timing over a time span of about 10 yr. **Methods.** We acquired new precise light curves for 29 transits of WASP-12 b, spanning four observing seasons from November 2012 to February 2016. New mid-transit times, together with those from the literature, were used to refine the transit ephemeris and analyze the timing residuals.

**Results.** We find that the transit times of WASP-12 b do not follow a linear ephemeris with a  $5\sigma$  confidence level. They may be approximated with a quadratic ephemeris that gives a change rate in the orbital period of  $(-2.56 \pm 0.40) \times 10^{-2} \text{ s yr}^{-1}$ . The tidal quality parameter of the host star was found to be equal to  $2.5 \times 10^5$ , which is similar to theoretical predictions for Sun-like stars. We also considered a model in which the observed timing residuals are interpreted as a result of the apsidal precession. We find, however, that this model is statistically less probable than the orbital decay.

**Key words.** stars: individual: WASP-12 – planets and satellites: individual: WASP-12 b – planet-star interactions

## 1. Introduction

With its orbital period of about 26 hours, the transiting planet WASP-12 b (Hebb et al. 2009) belongs to a group of hot Jupiters on the tightest orbits. It has a mass of  $M_b = 1.39 \pm 0.19 M_{\text{Jup}}$  (Knutson et al. 2014) and a radius  $R_b = 1.90 \pm 0.09 R_{\text{Jup}}$  (Maciejewski et al. 2013), which results in a mean density of

only 20% that of Jupiter. The planet is inflated; it fills ~60% of its Roche lobe (Li et al. 2010; Budaj 2011).

The planet shape departs from spherical symmetry, and the two bodies in the system – the planet and the host star – raise mutual tides. The nonspherical mass component of the gravitational field results in precession of the orbit (e.g., Ragozzine & Wolf 2009). This apsidal rotation might be observed through precise timing of transits for non-zero orbital eccentricities. The total apsidal precession is the sum of components such as tidal bulges, rotation bulges, and relativistic effects. With a precession period of 18 yr, WASP-12 b was found to be a promising candidate for detecting apsidal precession, mainly produced by tides risen on the planet (Ragozzine & Wolf 2009; Damiani & Lanza 2011). The precession rate may be used to determine the second-order Love number, which is related to the planet's internal density profile.

Planets on short-period orbits are expected to be unstable to tidal dissipation and finally spiral in toward the host star because they transfer the angular momentum of the orbital motion through tidal dissipation inside the star (e.g., Levrard et al. 2009;

\* Partly based on (1) data collected with the Nordic Optical Telescope, operated on the island of La Palma jointly by Denmark, Finland, Iceland, Norway, and Sweden, in the Spanish Observatorio del Roque de los Muchachos of the Instituto de Astrofísica de Canarias, (2) observations made at the Centro Astronómico Hispano Alemán (CAHA), operated jointly by the Max-Planck Institut für Astronomie and the Instituto de Astrofísica de Andalucía (CSIC), and (3) data collected with telescopes at the Rozhen National Astronomical Observatory.

\*\* The light curves are available at the CDS via anonymous ftp to [cdsarc.u-strasbg.fr](http://cdsarc.u-strasbg.fr) (130.79.128.5) or via <http://cdsarc.u-strasbg.fr/viz-bin/qcat?J/A+A/588/L6>

Essick & Weinberg 2016). The rate of this orbital decay can help determine the efficiency of the tide dissipation. The decaying orbital period is expected to be observed through transit timing. For some planets, the cumulative shift in transit times may be about  $10^2$  s after ten years (Birkby et al. 2014). Tentative detections of the orbital decay were reported for the systems OGLE-TR-113 and WASP-43 (Adams et al. 2010; Blecic et al. 2014), but have not been confirmed by further observations (Hoyer et al. 2016; Jiang et al. 2016).

In this study, we analyze new light curves for the transits of WASP-12 b and reanalyze some literature light curves to detect any longtime variations in the orbital period that may be attributed to one of the two mechanisms mentioned above.

## 2. Observations and data processing

We acquired 31 complete light curves for 29 transits between November 2012 and February 2016 using seven telescopes:

- the 2.56 m Nordic Optical Telescope (NOT) at the Observatorio del Roque de los Muchachos, La Palma (Spain) with the ALFOSC instrument in spectroscopic mode;
- the 2.2 m reflector (CA) at the Calar Alto Astronomical Observatory (Spain) with CAFOS in imaging mode;
- the 2.0 m *Ritchey-Chrétien-Coudé* telescope (ROZ) at the National Astronomical Observatory Rozhen (Bulgaria), equipped with a Roper Scientific VersArray 1300B CCD camera;
- the 1.8 m Bohyunsan optical telescope (BOAO) at the Bohyunsan Optical Astronomy Observatory (South Korea), equipped with a 4k CCD imaging instrument;
- the 1.5 m *Ritchey-Chrétien* telescope (OSN) at the Sierra Nevada Observatory (Spain) with a Roper Scientific VersArray 2048B CCD camera;
- the 1.2 m Trebur one-meter telescope (TRE) at the Michael Adrian Observatory in Trebur (Germany), equipped with an SBIG STL-6303 CCD camera;
- the 0.6 m Cassegrain telescope (PIW) at the Center for Astronomy of the Nicolaus Copernicus University in Piwnice (Poland), equipped with an SBIG STL-1001 CCD camera.

One transit was observed simultaneously with two telescopes and the other with a single telescope in two bands. Most of the data were acquired through *R*-band filters. To achieve a higher transit timing precision, some data were acquired without any filter. The list of observing runs is presented in Table A.1.

The data reduction was performed with the AstroImageJ package<sup>1</sup> (AIJ, Collins et al. 2016) following a standard procedure including de-biasing (or dark current removal) and flat-fielding. Photometric time series were obtained with differential aperture photometry. The aperture radius was allowed to vary to compensate for variable seeing. A set of comparison stars was optimized to achieve the lowest photometric scatter and to minimize trends caused by the differential atmospheric extinction. The light curves were examined for linear trends against airmass, position on the matrix, time, and seeing. The fluxes were normalized to unity for out-of-transit brightness. Time stamps were converted to barycentric Julian dates in barycentric dynamical time ( $\text{BJD}_{\text{TDB}}$ ).

The same procedure was applied to re-reduce data acquired with the 2.2 m telescope at Calar Alto, with the 2.0 m telescope at Rozhen, with the 1.8 m BOAO, and with the 1.2 m telescope in

Trebur. These data were published in Maciejewski et al. (2011) and Maciejewski et al. (2013). The re-reduction improved the photometric quality of the light curves by up to 15%.

The light curves from the NOT/ALFOSC were generated from series of low-resolution spectra obtained with a custom-built slit with a width of  $40''$ . We bias- and flat-field corrected the data and extracted spectra and corresponding calibration arcs. The wavelength calibration was based on He and Ne lamps using an IRAF script dedicated to NOT/ALFOSC long-slit data. On 2014 Jan. 31, grism 10, which covers the spectral range 3300–10550 Å, was used and the light curve was constructed from the whole spectral range. On 2015 Jan. 22 and 24, we used grism 4, which covers the spectral range 3200–9100 Å. The light curves were constructed from the 4950–6050 Å spectral range, which corresponds to a photometric *V*-band filter.

The set of new light curves was enhanced with photometric time series that are available in the literature. To limit the sample to the most reliable data, we considered only complete transit light curves that were acquired with telescopes with mirrors greater than 1 m. In addition to data from Maciejewski et al. (2011) and Maciejewski et al. (2013), we used photometric time series from Copperwheat et al. (2013), Chan et al. (2011), Cowan et al. (2012), and Stevenson et al. (2014).

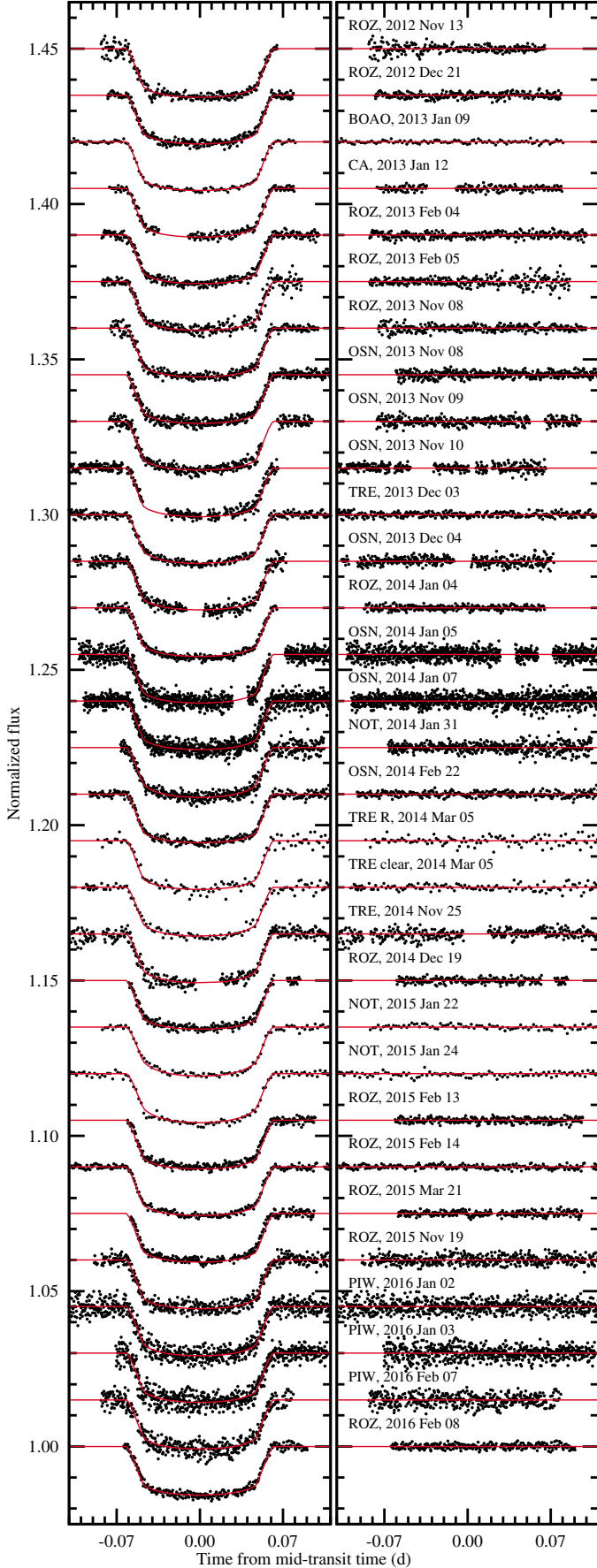
The Transit Analysis Package (Gazak et al. 2012) was employed to determine mid-transit times and their uncertainties for individual epochs. The transit parameters such as the orbital inclination, scaled semi-major axis, ratio of planetary to stellar radii, and coefficients of the quadratic limb-darkening (LD) law for *R*-band data were taken from Maciejewski et al. (2013). For the remaining bands, the LD coefficients were linearly interpolated from tables of Claret & Bloemen (2011). During the fitting procedure, the parameters were allowed to vary under Gaussian penalty determined by their uncertainties. The mid-transit time, airmass slope, and flux offsets were the free parameters. The median values of marginalized posteriori probability distributions of the ten Markov chain Monte Carlo values with  $10^5$  steps each and the 15.9 and 85.1 percentile values of these distributions were taken as the best-fit parameters and upper and lower  $1\sigma$  errors, respectively.

The mid-transit times for the new light curves, which are shown in Fig. 1, are given in Table A.2, those that we redetermined from literature data are presented in Table A.3. To extend the time covered by the observations, we also used the mid-transit time for epoch 0 from Hebb et al. (2009). It was obtained from a global fit and therefore represents an averaged value for the early epochs.

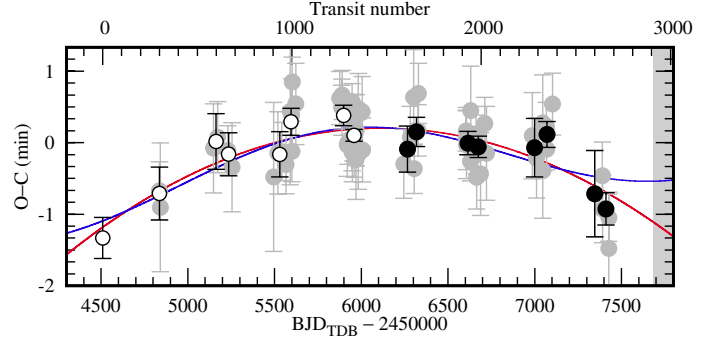
## 3. Results

The transit times were used to refine the transit ephemeris, and the timing residuals were examined for any long-term variations. Individual timing residuals were binned into half-season bins. Since complete transits of WASP-12 b are observable from the ground from October to March, the first bin of each observing season included data acquired between October and December, the second observations acquired between January and March. Uncertainties of individual mid-transit times were taken as weights, and the weighted standard deviations were calculated as errors. A fit of a linear ephemeris was found to be unsatisfying with the reduced  $\chi$ -squared ( $\chi_{\text{red}}^2$ ) equal to 4.02. This value corresponds to a *p*-value of  $5.3 \times 10^{-7}$ , which allows us to reject the null hypothesis with a  $5\sigma$  (99.9999%) confidence level. A Lomb-Scargle periodogram of the timing residuals reveals a

<sup>1</sup> <http://www.astro.louisville.edu/software/astroimagej>



**Fig. 1.** *Left:* new transit light curves for WASP-12 b, sorted by observation date. *Right:* the residuals from the transit model adopted from Maciejewski et al. (2013).



**Fig. 2.** Transit timing residuals from the linear ephemeris. The black open and filled circles denote literature and new data, respectively. Each point represents a half-season average of individual observations, which are marked with gray dots. The continuous line sketches the quadratic ephemeris discussed in Sect. 3. The dashed line illustrates the periodic signal discussed in Sect. 4. The gray area shows the next observing season, in which new data are expected from which we hope to be able to distinguish between both models.

long-term ( $\sim 3300$  d) signal that is similar to the time span covered by the observations. A quadratic ephemeris in the form of

$$T_{\text{mid}} = T_0 + P_b E + \frac{1}{2} \delta P_b E^2, \quad (1)$$

where  $E$  is the transit number from the cycle-zero epoch  $T_0$  and  $\delta P_b$  is the change in the orbital period between succeeding transits, yields a much better fit with  $\chi^2_{\text{red}} = 1.03$ . We obtained  $T_0 = 2\,454\,508.97696 \pm 0.00016$  BJD<sub>TDB</sub>,  $P_b = 1.09142162 \pm 0.00000021$  d, and  $\delta P_b = (-8.9 \pm 1.4) \times 10^{-10}$  days epoch<sup>-2</sup>. The latter quantity translates into the short-term change rate in the orbital period  $\dot{P}_b = \delta P_b / P_b = (-2.56 \pm 0.40) \times 10^{-2}$  s yr<sup>-1</sup>. The timing residuals from the linear ephemeris together with the quadratic model are plotted in Fig. 2.

#### 4. Discussion

The negative value of  $\delta P_b$  may be interpreted as evidence of an orbital decay that is driven by tidal dissipation in the host star. For a star-planet system in which the total angular momentum  $L_{\text{tot}}$  is conserved, but energy is dissipated by tides, no equilibrium state exists if  $L_{\text{tot}}$  is smaller than the critical angular momentum  $L_{\text{crit}}$  (see Eq. (2) in Levrard et al. 2009), which is required for the star-planet system to reach a state of dual synchronization. The spin of the star may be a significant, if not dominant, component of  $L_{\text{tot}}$ . It depends on the rotational velocity of the star  $\omega_*$ , which remains roughly constrained for the WASP-12 star by spectral observations (Albrecht et al. 2012). For the WASP-12 system, we obtained  $L_{\text{tot}}/L_{\text{crit}} = 0.4\text{--}0.7$ , depending on the rotation period of the star. Since this ratio is obviously lower than 1, the planet will unavoidably spiral inward<sup>2</sup>. Following Eq. (5) in Levrard et al. (2009), we find a relatively short in-spiral time of about  $10^6$  yr, which is very short compared to the age of the system of  $\sim 2 \times 10^9$  yr (Hebb et al. 2009).

The observed rate of the orbital decay may be used to determine the stellar tidal quality parameter  $Q'_*$ . This quantity is the ratio of the stellar tidal quality factor  $Q$ , which characterizes a

<sup>2</sup> In calculations we used a normalized moment of inertia (NMOI) equal to 0.04 for the star. The quantity was interpolated from tables of Claret & Gimenez (1989). For the planet, we used Jupiter's NMOI of 0.26 from Helled et al. (2011).

body's response to tides, to the second-order stellar tidal Love number  $k_2$ . Following [Blecic et al. \(2014\)](#), we adopted Eq. (3) of [Levrard et al. \(2009\)](#) for synchronous planetary rotation and negligible eccentricity and obliquity

$$Q'_* = 9P_b \dot{P}_b^{-1} \frac{M_b}{M_*} \left( \frac{R_*}{a_b} \right)^5 \left( \omega_* - \frac{2\pi}{P_b} \right). \quad (2)$$

We obtained  $Q'_*$  of about  $2.5 \times 10^5$ , which is on the same order as  $Q'_* = 4.3 \times 10^5$  calculated from models of [Essick & Weinberg \(2016\)](#) for solar-type host stars.

Alternatively, the observed departure from the linear ephemeris may be a part of the 3300 d periodic signal induced by star-planet tidal interactions. Using Eq. (86) of [Migaszewski \(2012\)](#), we found the timescale of the rotation of the pericenter caused by tides that have risen in the planet as a response to stellar gravity to be of about  $10^1$  yr. This value might correspond to the observed signal<sup>3</sup>. To explore this possibility, we employed the Systemic software (version 2.182, [Meschiari et al. 2009](#)). We used the transit-timing dataset enhanced with occultation times from [Campo et al. \(2011\)](#), [Croll et al. \(2011, 2015\)](#), [Crossfield et al. \(2012\)](#). We also included a homogenous set of precise radial velocity (RV) measurements from [Knutson et al. \(2014\)](#). This Doppler time series, free of any year-to-year instrumental trends, was acquired with the Keck/HIRES instrument between December 2009 and December 2013. The orbital period, planetary mass, mean anomaly for a given epoch, eccentricity  $e_b$ , longitude of periastron for a given epoch, and periastron precession rate  $\dot{\omega}$  were allowed to be fitted. The Levenberg-Marquardt algorithm was used to find the best-fit model with a Keplerian approach. The parameter uncertainties were estimated as median absolute deviations from a bootstrap run of  $10^5$  trails. We obtained  $e_b = 0.00110 \pm 0.00036$  and  $\dot{\omega} = 0.095 \pm 0.020$  deg per day, which corresponds to a period of the periastron precession  $\tau_\omega = 10.4 \pm 2.2$  yr.

The periastron precession with the very low value of  $e_b$  would have a marginal effect on transit parameters that are directly determined from light curves. In particular, the range of variations in transit duration is expected to be 2.7 s, much smaller than typical transit duration uncertainties of 1–3 min.

In Fig. 2 we also show the best-fit sinusoidal signal with a period equal to  $\tau_\omega$ . With  $\chi^2_{\text{red}} = 1.46$ , the goodness of fit is noticeably poorer than for the quadratic model. The Bayesian information criterion (BIC, [Schwarz 1978](#))

$$\text{BIC} = \chi^2 + k \ln N, \quad (3)$$

where  $k$  is the number of fit parameters and  $N$  is the number of data points, also favors the quadratic model with  $\text{BIC} = 21.8$  over the periodic model with  $\text{BIC} = 28.6$  with a probability ratio of  $e^{\Delta\text{BIC}/2} = 29$ . We note that observations in the upcoming season 2016/2017, which is represented by the gray area in Fig. 2, are expected to let us clearly distinguish between both models.

## 5. Conclusions

Our new precise observations, spread over four years, show that mid-transit times of the WASP-12 b planet do not follow a linear ephemeris. This phenomenon may be interpreted as the result of orbital decay, periastron precession due to planetary tides, or a

<sup>3</sup> Other contributions to the rotation of the pericenter can be neglected. The timescales of the rotation caused by rotational deformation of the star and the planet are of about  $10^4$  yr. The timescales for the tidal deformation of the star and relativistic effects are of about  $10^3$  yr.

combination of these two effects. The statistics formally favors the orbital decay scenario. The tidal quality parameter for the host star  $Q'_*$  was found to be slightly lower than the theoretical predictions. In turn, the periastron precession model is consistent with theoretical predictions and places tight constraints on the orbital eccentricity of the planet. It may be used to precisely determine the Love number of the planet, which is related to its internal structure. Future precise timing observations are expected to provide evidence in favor of one of the two scenarios. This will lead to a better understanding of the properties of the host star or the planet.

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**Appendix A: Supplementary materials**

Table A.1 presents details of new light curves acquired for the transits of WASP-12 b. Table A.2 lists the mid-transit times for the new light curves. Table A.3 shows the mid-transit times based on the literature data that were used in our study.

**Table A.1.** Details on new light curves obtained by us.

Date UT	Epoch	Telescope	Filter	$\Gamma$	$pnr$
2012 Nov. 13	1591	ROZ	$R_C$	1.22	1.05
2012 Dec. 21	1625	ROZ	$R_C$	1.22	0.72
2013 Jan. 09	1643	BOAO	$R_B$	0.38	0.95
2013 Jan. 12	1646	CA	$R_C$	1.22	0.72
2013 Feb. 04	1667	ROZ	$R_C$	1.22	0.74
2013 Feb. 05	1668	ROZ	$R_C$	1.22	1.17
2013 Nov. 08	1920	ROZ	$R_C$	1.22	0.82
		OSN	$R_C$	1.71	0.72
2013 Nov. 09	1921	OSN	$R_C$	1.69	0.80
2013 Nov. 10	1922	OSN	$R_C$	1.71	0.76
2013 Dec. 03	1943	TRE	none	1.03	0.72
2013 Dec. 04	1944	OSN	$R_C$	1.71	0.86
2014 Jan. 04	1973	ROZ	$R_C$	1.22	0.61
2014 Jan. 05	1974	OSN	$R_C$	3.75	0.80
2014 Jan. 07	1975	OSN	$R_C$	3.99	0.81
2014 Jan. 31	1997	NOT	none	1.99	0.97
2014 Feb. 22	2018	OSN	$R_C$	1.09	0.74
2014 Mar. 05	2028	TRE	$R_B$	0.39	2.09
		TRE	none	0.39	1.52
2014 Nov. 25	2270	TRE	none	1.25	1.24
2014 Dec. 19	2292	ROZ	$R_C$	1.22	0.67
2015 Jan. 22	2324	NOT	$V$	0.32	1.17
2015 Jan. 24	2325	NOT	$V$	0.32	1.36
2015 Feb. 13	2344	ROZ	$R_C$	1.22	0.65
2015 Feb. 14	2345	ROZ	$R_C$	0.76	0.73
2015 Mar. 21	2377	ROZ	$R_C$	1.22	0.63
2015 Nov. 19	2599	ROZ	$R_C$	1.76	1.00
2016 Jan. 02	2640	PIW	none	2.62	1.20
2016 Jan. 03	2641	PIW	none	2.40	1.18
2016 Feb. 07	2673	PIW	none	2.01	1.24
2016 Feb. 08	2674	ROZ	$R_C$	1.22	0.69

**Notes.** Date UT is given for mid-transit time. Epoch is the transit number from the initial ephemeris given in Hebb et al. (2009).  $R_C$  and  $R_B$  denote Cousins and Bessel  $R$ -band filters, respectively.  $\Gamma$  is the median number of exposures per minute.  $pnr$  is the photometric scatter in  $10^{-3}$  normalized flux per minute of the observation adopted from Fulton et al. (2011).

**Table A.2.** New mid-transit times for individual epochs.

Epoch	$T_{\text{mid}}$ (BJD <sub>TDB</sub> 2 450 000+)	$N_{\text{lc}}$	Telescope
1591	6245.42729 <sup>+0.00033</sup> <sub>-0.00033</sub>	1	ROZ
1625	6282.53584 <sup>+0.00030</sup> <sub>-0.00030</sub>	1	ROZ
1643	6302.18179 <sup>+0.00046</sup> <sub>-0.00046</sub>	1	BOAO
1646	6305.45536 <sup>+0.00024</sup> <sub>-0.00026</sub>	1	CA
1667	6328.37556 <sup>+0.00027</sup> <sub>-0.00025</sub>	1	ROZ
1668	6329.46733 <sup>+0.00029</sup> <sub>-0.00029</sub>	1	ROZ
1920	6604.50489 <sup>+0.00021</sup> <sub>-0.00020</sub>	2	ROZ, OSN
1921	6605.59624 <sup>+0.00030</sup> <sub>-0.00030</sub>	1	OSN
1922	6606.68760 <sup>+0.00033</sup> <sub>-0.00034</sub>	1	OSN
1943	6629.60726 <sup>+0.00019</sup> <sub>-0.00019</sub>	1	TRE
1944	6630.69917 <sup>+0.00043</sup> <sub>-0.00043</sub>	1	OSN
1973	6662.35014 <sup>+0.00019</sup> <sub>-0.00018</sub>	1	ROZ
1974	6663.44136 <sup>+0.00019</sup> <sub>-0.00019</sub>	1	OSN
1975	6664.53256 <sup>+0.00031</sup> <sub>-0.00032</sub>	1	OSN
1997	6688.54384 <sup>+0.00040</sup> <sub>-0.00041</sub>	1	NOT
2018	6711.46415 <sup>+0.00025</sup> <sub>-0.00026</sub>	1	OSN
2028	6722.37807 <sup>+0.00046</sup> <sub>-0.00047</sub>	2	TRE
2270	6986.50195 <sup>+0.00043</sup> <sub>-0.00042</sub>	1	TRE
2292	7010.51298 <sup>+0.00039</sup> <sub>-0.00039</sub>	1	ROZ
2324	7045.43831 <sup>+0.00046</sup> <sub>-0.00049</sub>	1	NOT
2325	7046.53019 <sup>+0.00049</sup> <sub>-0.00047</sub>	1	NOT
2344	7067.26715 <sup>+0.00022</sup> <sub>-0.00023</sub>	1	ROZ
2345	7068.35834 <sup>+0.00020</sup> <sub>-0.00021</sub>	1	ROZ
2377	7103.28423 <sup>+0.00031</sup> <sub>-0.00030</sub>	1	ROZ
2599	7345.57867 <sup>+0.00042</sup> <sub>-0.00040</sub>	1	ROZ
2640	7390.32708 <sup>+0.00033</sup> <sub>-0.00034</sub>	1	PIW
2641	7391.41818 <sup>+0.00033</sup> <sub>-0.00032</sub>	1	PIW
2673	7426.34324 <sup>+0.00055</sup> <sub>-0.00052</sub>	1	PIW
2674	7427.43496 <sup>+0.00023</sup> <sub>-0.00022</sub>	1	ROZ

**Notes.** Epoch is the transit number from the initial ephemeris given in Hebb et al. (2009).  $T_{\text{mid}}$  is the mid-transit time.  $N_{\text{lc}}$  is the number of individual light curves used simultaneously for a given epoch.

**Table A.3.** Mid-transit times based on literature data.

Epoch	$T_{\text{mid}}$ (BJD <sub>TDB</sub> 2 450 000+)	Data source
0	4508.97685 <sup>+0.00020</sup> <sub>-0.00020</sub>	1
300	4836.40340 <sup>+0.00028</sup> <sub>-0.00028</sub>	2
304	4840.76893 <sup>+0.00062</sup> <sub>-0.00060</sub>	3
585	5147.45861 <sup>+0.00043</sup> <sub>-0.00042</sub>	4
608	5172.56138 <sup>+0.00036</sup> <sub>-0.00035</sub>	3
661	5230.40653 <sup>+0.00024</sup> <sub>-0.00024</sub>	5
683	5254.41761 <sup>+0.00043</sup> <sub>-0.00042</sub>	5
903	5494.52999 <sup>+0.00072</sup> <sub>-0.00074</sub>	4
925	5518.54147 <sup>+0.00040</sup> <sub>-0.00040</sub>	6
947	5542.55273 <sup>+0.00028</sup> <sub>-0.00029</sub>	4
969	5566.56385 <sup>+0.00028</sup> <sub>-0.00027</sub>	4
991	5590.57561 <sup>+0.00068</sup> <sub>-0.00071</sub>	4
998	5598.21552 <sup>+0.00035</sup> <sub>-0.00035</sub>	4
1000	5600.39800 <sup>+0.00029</sup> <sub>-0.00030</sub>	4
1001	5601.49010 <sup>+0.00024</sup> <sub>-0.00024</sub>	4
1021	5623.31829 <sup>+0.00039</sup> <sub>-0.00039</sub>	4
1253	5876.52786 <sup>+0.00027</sup> <sub>-0.00026</sub>	4
1263	5887.44198 <sup>+0.00021</sup> <sub>-0.00021</sub>	4
1264	5888.53340 <sup>+0.00027</sup> <sub>-0.00027</sub>	4
1266	5890.71635 <sup>+0.00024</sup> <sub>-0.00024</sub>	4
1293	5920.18422 <sup>+0.00031</sup> <sub>-0.00030</sub>	4
1296	5923.45850 <sup>+0.00022</sup> <sub>-0.00021</sub>	4
1317	5946.37823 <sup>+0.00018</sup> <sub>-0.00018</sub>	4
1318	5947.47015 <sup>+0.00017</sup> <sub>-0.00017</sub>	4
1319	5948.56112 <sup>+0.00033</sup> <sub>-0.00034</sub>	4
1322	5951.83536 <sup>+0.00011</sup> <sub>-0.00011</sub>	7
1323	5952.92708 <sup>+0.00013</sup> <sub>-0.00013</sub>	7
1329	5959.47543 <sup>+0.00017</sup> <sub>-0.00017</sub>	4
1330	5960.56686 <sup>+0.00032</sup> <sub>-0.00040</sub>	4
1339	5970.38941 <sup>+0.00039</sup> <sub>-0.00040</sub>	4
1340	5971.48111 <sup>+0.00035</sup> <sub>-0.00035</sub>	4
1350	5982.39509 <sup>+0.00034</sup> <sub>-0.00033</sub>	4
1351	5983.48695 <sup>+0.00035</sup> <sub>-0.00034</sub>	4
1371	6005.31533 <sup>+0.00037</sup> <sub>-0.00034</sub>	4
1372	6006.40637 <sup>+0.00031</sup> <sub>-0.00033</sub>	4

**Notes.** Epoch is the transit number from the initial ephemeris given in Hebb et al. (2009).  $T_{\text{mid}}$  is the mid-transit time.

**References.** (1) Hebb et al. (2009); (2) Copperwheat et al. (2013); (3) Chan et al. (2011); (4) Maciejewski et al. (2013); (5) Maciejewski et al. (2011); (6) Cowan et al. (2012); (7) Stevenson et al. (2014).