# Journal Pre-proofs

# Research paper

The long duration cryogenic system of the OLIMPO balloon–borne experiment: design and in–flight performance

A. Coppolecchia, L. Lamagna, S. Masi, P.A.R. Ade, G. Amico, E.S.Battistelli, P. Bernardis de, F. Columbro, L. Conversi, G. D'Alessandro, M.De Petris, M. Gervasi, F. Nati, L. Nati, A. Paiella, F. Piacentini, G. Pisano, G.Presta, A. Schillaci, C. Tucker, M. Zannoni

PII:	S0011-2275(20)30131-4
DOI:	https://doi.org/10.1016/j.cryogenics.2020.103129
Reference:	JCRY 103129
To appear in:	Cryogenics
Received Date:	7 February 2020
Accepted Date:	7 July 2020

Please cite this article as: Coppolecchia, A., Lamagna, L., Masi, S., Ade, P.A.R., Amico, G., Battistelli, E.S., Bernardis, P. de, Columbro, F., Conversi, L., D'Alessandro, G., Petris, M.D., Gervasi, M., Nati, F., Nati, L., Paiella, A., Piacentini, F., Pisano, G., Presta, G., Schillaci, A., Tucker, C., Zannoni, M., The long duration cryogenic system of the OLIMPO balloon–borne experiment: design and in–flight performance, *Cryogenics* (2020), doi: https://doi.org/10.1016/j.cryogenics.2020.103129

This is a PDF file of an article that has undergone enhancements after acceptance, such as the addition of a cover page and metadata, and formatting for readability, but it is not yet the definitive version of record. This version will undergo additional copyediting, typesetting and review before it is published in its final form, but we are providing this version to give early visibility of the article. Please note that, during the production process, errors may be discovered which could affect the content, and all legal disclaimers that apply to the journal pertain.

© 2020 Published by Elsevier Ltd.

ELSEVIER		ISSN 0011-2275
Ever tempe Boytemp Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting Constanting	rature engineering reconductivity onics	

# The long duration cryogenic system of the OLIMPO balloon–borne experiment: design and in–flight performance

A. Coppolecchia<sup>a,b,\*</sup>, L. Lamagna<sup>a,b</sup>, S. Masi<sup>a,b</sup>, P.A.R. Ade<sup>d</sup>,
G. Amico<sup>a</sup>, E.S. Battistelli<sup>a,b</sup>, P. de Bernardis<sup>a,b</sup>, F. Columbro<sup>a,b</sup>,
L. Conversi<sup>a,f</sup>, G. D'Alessandro<sup>a,b</sup>, M. De Petris<sup>a,b</sup>, M. Gervasi<sup>c</sup>,
F. Nati<sup>a,c</sup>, L. Nati<sup>a</sup>, A. Paiella<sup>a,b</sup>, F. Piacentini<sup>a,b</sup>, G. Pisano<sup>d</sup>,
G. Presta<sup>a,b</sup>, A. Schillaci<sup>a,e</sup>, C. Tucker<sup>d</sup>, M. Zannoni<sup>c</sup>

<sup>a</sup> Physics Department, University of Rome Sapienza, P.le Aldo Moro 5, 00185, Roma, Italy

<sup>b</sup>Istituto Nazionale di Fisica Nucleare, Sezione di Roma, P.le A. Moro 2, 00185 Roma, Italy

<sup>c</sup>Physics Department, University of Milano-Bicocca, Piazza della Scienza 3, 20126 Milano, Italy

<sup>d</sup>School of Physics and Astronomy, Cardiff University, Cardiff CF24 3YB, UK <sup>e</sup>California Institute of Technology, Pasadena, CA 91125, USA <sup>f</sup>ESA/ESRIN, Largo Galileo Galilei, 1, 00044 Frascati (RM), Italy

#### Abstract

We describe the design and in–flight performance of the cryostat and the self–contained <sup>3</sup>He refrigerator for the OLIMPO balloon–borne experiment, a spectrophotometer to measure the Sunyaev–Zel'dovich effect in clusters of galaxies.

The <sup>3</sup>He refrigerator provides the 0.3 K operation temperature for the four arrays of kinetic inductance detectors working in 4 bands centered at 150, 250, 350 and 460 GHz. The cryostat provides the 1.65 K base temperature for the <sup>3</sup>He refrigerator, and cools down the reimaging optics and the filters chain at about 2 K.

The integrated system was designed for a hold time of about 15 days, to achieve the sensitivity required by the planned OLIMPO observations, and successfully operated during the first long–duration stratospheric flight of OLIMPO in July 2018.

The cryostat features two tanks, one for liquid nitrogen and the other one for liquid helium. The long hold time has been achieved by means of custom stiff G10 fiberglass tubes support, which ensures low thermal conductivity and remarkable structural stiffness; multi-layer superinsulation,

Preprint submitted to Elsevier

<sup>\*</sup>Corresponding author

 $<sup>{\</sup>it Email~address:~alessandro.coppolecchia@roma1.infn.it~(A.~Coppolecchia~)}$ 

and a vapour cooled shield, all reducing the heat load on the liquid helium tank.

The system was tested in the lab, with more than 15 days of unmanned operations, and then in the long duration balloon flight in the stratosphere. In both cases, the detector temperature was below 300 mK, with thermal stability better than  $\pm$  0.5 mK.

The system also operated successfully in the long duration stratospheric balloon flight.

Keywords: Wet cryostat, long duration  ${}^{3}\text{He}$  sorption fridge, superinsulation, Helium vapor cooled shield, balloon-borne experiment.

# 1. Introduction

5

20

25

30

Precision measurements of the Cosmic Microwave Background (CMB) require very sensitive detectors, arranged in arrays, operating at mm wavelengths. Photon-noise-limited performance [1] can been achieved by cooling the optical system at cryogenic temperatures, and the detectors at sub-K temperatures, for the entire duration of the measurements.

The CMB exhibits a thermal spectrum, characterized by a single temperature,  $T_{rad} = (2.72548 \pm 0.00057)$  K [2], proving that the universe underwent a dense, hot phase, as described by the standard cosmological model. This radiation is very close to be spatially isotropic. Its low level of anisotropy, as measured by many experiments in the last decades, is determined both by the fluctuations of the photon density across the Last Scattering Surface (primary anisotropies), and by the structure formation processes, leaving their imprint on the CMB along its optical path to us (secondary anisotropies). Precision measurements of CMB anisotropy at different angular scales pro-

vide invaluable information about these processes [3].

Foreground emissions (galactic and extragalactic) must be minimized and removed by proper selection of the observation bands and sky regions. So, in addition to high sensitivity detectors, simultaneous measurements at different frequencies are needed. This allows to monitor and to efficiently separate different signal components, which feature different spectral shapes.

The best way to explore the whole spectral region of interest for the CMB is to operate instruments with no hindrance from the atmospheric opacity in the microwaves, either by flying them on satellites, or by keeping them at very high altitude on sub–orbital platforms, like rockets and stratospheric balloons [4].

OLIMPO [5, 6] is the largest telescope ever launched on a stratospheric balloon (2.6 m optical aperture). It was designed to observe the sky in the mm and sub-mm bands, using innovative Kinetic Inductance Detectors (KIDs)[7, 8, 9] and a plug-in Differential Fourier Transform Spectrometer

(DFTS) [10, 11]. The main goal of the experiment is to measure the secondary CMB anisotropy known as Sunyaev–Zel'dovich (SZ) [12] effect, in the direction of selected clusters of galaxies. To this purpose, it covers four frequency bands, centered at 150, 250, 350 and 460 GHz: these frequencies match the negative, zero, and positive regimes of the SZ spectral regions.

35

45

65

70

The limited electrical power available on the payload, and the extreme temperature and pressure of the stratospheric environment, do not permit the use of a dry mechanical cryocooler. So we opted for a wet configuration with liquid cryogens. Moreover this configuration allows to minimize the

- vibrations and consequent microphonic noise in the detectors. The chosen configuration provides the needed lifetime for the cryostat, without requiring any maintenance, nor electrical power. With careful design, it matches the forecast experiment duration at float: about 15 days for a polar stratospheric long duration balloon (LDB).
  - In this paper we describe the optimization and performance of the wet cryostat and the closed cycle <sup>3</sup>He refrigerator of OLIMPO.

Starting from the experience of BOOMERanG [13], the OLIMPO cryostat has been designed to have four different temperature stages: the external shell, the nitrogen tank, the helium vapor cooled shield, and the helium

tank. These stages are at gradually decreasing temperatures from the outside to the innermost insert, and they are thermally–insulated by structural supports made of fiberglass.

The room-temperature external shell has been designed to sustain a very low pressure in the internal volume and to be the first shield to the radiative load on the system.

The nitrogen tank, which operates at 77 K at standard pressure, has been designed with a bell shape, and plays an essential role for the inner stages of the cryostat providing heat–sinking for the mechanical and electrical connections from the outside, and an envelope against radiation loading.

Radiation shielding is further improved through multi-layer superinsulation<sup>1</sup>.

Vapors from the <sup>4</sup>He boiloff flow through a serpentine welded on a copper shield placed between the two tanks, cooling it at an intermediate temperature, close to 35 K (with a temperature of the nitrogen tank at 77 K) and close to about 21 K in flight (when the temperature of the nitrogen tank is around 52 K).

The <sup>4</sup>He stage has been built with a cylindrical shape, where the available liquid reduces to about 50 % of the full volume after the pump-down operation needed to reach a temperature of 1.65 K. The main purpose of the <sup>4</sup>He tank in the OLIMPO cryostat is to cool down the reimaging optics and to provide the thermal reference stage for the condensation point of the

<sup>&</sup>lt;sup>1</sup>Cryolam NRC–2, Metallized Products, Inc., Winchester, MA, USA.

 $^{3}$ He evaporative sorption refrigerator. This refrigerator has been customdesigned to work on the stratospheric balloon payload and to cool to 300 mK the focal planes of the experiment for about two weeks without being recycled.

# 2. Cryostat design

75

80



Figure 1: Exploded sections of the cryostat tanks (supports removed). We can see: the external shell with the window, the toroidal nitrogen tank, the helium vapor–shield, the cylindrical helium tank, the fiberglass support in green, the *cold optics* and the focal planes with the arrays. The fill and vent lines have been removed for clarity, and are discussed in §2.1 and §2.4.

A general drawing of the cryostat is shown in figure 1.

## 2.1. External shell and liquid nitrogen tank

The external shell and toroidal liquid nitrogen (LN2) tank, with a volume of 74 liters, are the same as those used for the cryostat [13] of the LDB BOOMERanG experiment [14, 15]. Both the external shell and the nitrogen tank are made of 6061 aluminum to limit their mass. The external shell has been modified changing the position of the window for radiation input, since in OLIMPO the window is located on the side wall to couple the cold reimaging optics to the focal plane of the Ritchey–Chretien room temperature telescope. The room–temperature shell is a vacuum–tight aluminum vessel able to maintain a pressure below  $10^{-7}$  mbar in the internal volume. It represents the first shield against the radiative load on the receiver, thanks to the low emissivity of the inner surface.

The LN2 tank uses the high latent heat of evaporation  $(199 \frac{kJ}{kg} \text{ at p about}$ 1 bar, versus  $21 \frac{kJ}{kg}$  for <sup>4</sup>He) to intercept the radiation loading from room temperature, through the evaporation of a comparatively limited amount of cryogen. Two stainless steel tubes with corrugated inserts are used as the fill and vent lines to the nitrogen tank. They have an outer diameter of 22 mm and 0.2 mm thickness, to minimize the conductive thermal loading. The corrugation segments increase the length thus reducing the heat load from the outer shell.

Differently from the BOOMERanG cryostat, where kevlar cords were used, in the OLIMPO cryostat the LN2 tank is connected to the room temperature shell and supported by a single fiberglass (G10) cylinder, 4 mm 100 thick, 800 mm long and 674 mm in diameter, as shown in figure 1. This support provides increased stiffness, which is needed due to the telescope elevation range between  $0^{\circ}$  and  $55^{\circ}$ , while preserving a reasonably low heat load of around 0.5 W during the flight. Since the atmospheric pressure at stratospheric altitude is a few mbar, liquid nitrogen in the tank cools further 105 down to about 52 K, changing from liquid to solid state and possibly loosing thermal contact with the tank walls. To cope with this effect, we added several internal copper braids, to ensure the thermal contact between the solid nitrogen and the nitrogen tank wall. The performance of this solution is briefly discussed in  $\S6$ . 110

## 2.2. Superinsulation

Multi-layer insulation (*superinsulation*) is implemented between the external shell and the nitrogen tank in order to reduce the radiative load. For *n*-layers of *superinsulation*, where each layer is made of highly reflective material on one side and moderately emitting material on the opposite side, the theoretical radiative power load on the innermost surface is:

$$\dot{Q}_{rad} = \left(\frac{1}{n+1}\right) \left(\frac{\epsilon}{2}\right) \sigma A \left(T_{n+1}^4 - T_0^4\right) \tag{1}$$

in the approximation  $A_i = A$  and  $\epsilon_i = \epsilon \ll 1$ , where A is the surface area,  $\epsilon$  is the surface emissivity and  $\sigma = 5.67 \cdot 10^{-8} W m^{-2} K^{-4}$  is the Stefan-Boltzmann constant.

120

However, in practice there is an extra conductive thermal load, due to a shallow contact between the layers, and an extra convective possible thermal load, due to the density of the layers which makes more difficult to remove the gas trapped among them. The number of layers has been chosen as a tradeoff between reduction of radiative load and available volume. We installed 30 layers of *superinsulation*: 15 between the external shell and the fiberglass support tube, and 15 between the fiberglass support and the LN2 tank, for a total thickness of the superinsulation blankets of ~ 15 mm; 20 layers covered the top and bottom bases of the LN2 tank, with a thickness of ~ 10 mm.

130

140

145

150

155

125

By measuring the boiling off rate of the liquid nitrogen before and after the assembly of the *superinsulation*, we have observed a heat load reduction from 52 W to 8.5 W. In the flight configuration, the boiloff rate is about 0.2 l/h, which guarantees a nitrogen hold time of about 15 days for a 74 l reservoir.

## 135 2.3. Helium vapor cooled shield (VCS)

The Nitrogen tank supports a VCS, cooled by the cold gas evaporating from the liquid <sup>4</sup>He (L<sup>4</sup>He) tank. The fiberglass tube supporting the VCS from the LN2 tank is 2.5 mm thick, 465 mm long, and 550 mm inner diameter (see figure 2 on the left); the fiberglass tube supporting the L<sup>4</sup>He tank from the VCS is 2.5 mm thick, 440 mm long and 478 mm inner diameter.

The main purpose of this shield is to absorb radiated (mainly) and conducted heat coming from the nitrogen tank, reducing the total heat load on the L<sup>4</sup>He tank. The VCS is a self regulating system: if the shield is too warm, the evaporation rate of helium in the L<sup>4</sup>He tank increases, due to higher loading from the shield, thus producing a larger flux of cold vapors and therefore providing additional cooling for the shield itself.

This shield is split in two copper cylinders, interrupted by a center flange where both fiberglass cylinders are bolted on. A copper serpentine, collecting the flow of cold gas evaporated from the L<sup>4</sup>He tank, is hard soldered around the cylinders. The length of the serpentine is ~ 11 m and the internal diameter is 8 mm. Since its high impedance might produce pressure oscillations during the liquid helium pump–down process, a lower impedance heat exchanger was assembled in parallel to the serpentine, with a motorized valve regulating the fraction of <sup>4</sup>He vapors flux flowing along the two paths. The shield and the valve with the heat exchanger are visible in figure 2.

The valve is a *Swagelok* model SS–6BW–TW<sup>2</sup>, made of stainless steel. The valve was degreased with ultrasonic cleaning to prevent it from getting stuck when cooled at cryogenic temperatures. The valve is rotated by a *GlobeMotors* DC motor with bronze bushings, which have been rebored (+100  $\mu$ m on the diameter) for cryogenic temperature operation. In §4 we shortly report on the performance of this system.

<sup>&</sup>lt;sup>2</sup>https://www.swagelok.com



Figure 2: Pictures: on the left the Helium vapor cooled shield, the fiberglass tube support, the serpentine and the exchanger. On the right detail of the exchanger, the valve and the mechanism used to operate the valve, regulating the flux of the vapors.

## 2.4. Liquid Helium Tank

The L<sup>4</sup>He tank cap volume is 64 liters. The boiling temperature of helium drops from 4.2 K at standard pressure to below 1.65 K when the pressure of the vapors is reduced to a few mbar in the stratosphere.

The innermost and coldest stages of the OLIMPO cryostat (including the refrigerator, the detectors, the cold optics and the optical filter chain) have been thermally and mechanically designed to take advantage of the stratospheric environment, and thus cannot be properly operated when the pressure over the helium reservoir is higher than a few mbar. Therefore the 170  $L^4$ He reservoir must be pumped in the laboratory to operate the <sup>3</sup>He refrigerator, allowing for pre-flight detectors validation and calibration. Just before launch, the mechanical rotary pump used for maintaining the low pressure on the  $L^4$ He bath is disconnected, after having shut the vent line by means of a remotely operated valve. In this way the pressure and tem-175 perature of the L<sup>4</sup>He bath drift slightly up while, in a few hours, the payload reaches the floating altitude. There, the vent valve is re-opened, and the low ( $\sim 300$  Pa) outer pressure maintains the liquid at a steady temperature of 1.65 K. This procedure avoids a large loss of cryogen which would be unavoidable if pumping the He bath in uncontrolled conditions during the ascent. The efficiency of this procedure has been confirmed in flight, as described in  $\S6$ .

Even after an optimized pump-down (about 1.5 mbar/min) and tuning of the pumping line impedance,  $\sim 50$  % of the L<sup>4</sup>He is used to cooldown the remaining  $\sim 50$  % and the attached equipment. The final temperature of the L<sup>4</sup>He bath is below 1.65 K, enough to efficiently condensate the <sup>3</sup>He in the refrigerator and to cool the optics box and the coldest stages of the optical filtering chain.

In order to interface the bottom of the helium tank directly to the cold optics we decided to make both the top and bottom flanges flat. The thickness t of the tank was chosen according to the following formula [16]:

$$t = 0.866 \cdot r \cdot \sqrt{\frac{P \cdot X}{R}} \tag{2}$$

where r is the radius of the tank, P is the pressure applies from the outer side, X, the safety factor, and R, the breaking strength. For a standard pressure of ~  $10^5$  Pa, radius r = 0.225 m, X = 5 and  $R \sim 3 \times 10^8$  Pa for the aluminum we obtain a thickness of about 8 mm. For the fabrication of the flange we started from a 30 mm thick flat and lightened it by milling out most of the material, leaving a network of thick ribs, defining many areas with typical side of 100 mm, and reinforcing a 5 mm thick flat skin.

We decided to add some ribs around the tank and some ribs on the internal side of the bottom flange of the tank, to optimize the thermal contact between the liquid and the fridge.



Figure 3: Schematic diagram of  $L^4$ He tank connections. Two steel tubes connect the tank to the exterior: one for Fill, one for Vent. The former is directly welded to the helium tank, while the latter is split in two parallel paths. One is a serpentine hard soldered on the VCS, and one is a bypass heat exchanger on the top cover of the VCS. The flow through the heat exchanger can be regulated by means of motorized valve 1.

To fill with liquid Helium and to provide a vent for the vapors, we use two tubes, made of stainless steel with a thickness of 0.2 mm and a diameter of 12.7 mm. The evaporation tube is split into two directions, one towards the serpentine and the other towards the valve, then towards the exchanger flowing into the Helium tank. The filling tube is connected straight to the tank, and is also connected to one of the two serpentine outputs, as shown in fig. 3. The fill and vent tubes are interrupted with a 200 mm long section

where 185 mm  $\emptyset$  stainless steel belows is inserted. This reduces the heat load and allows for small shifts of the two ends, due to thermal contraction 210 effects.

The geometry of the system and the heat loads on the helium bath, mainly from cables harness and fiberglass tube support, are such that the equilibrium of the VCS is expected to be reached at a temperature of 35 K (with the LN tank at 77 K). The fiberglass tube connected between the VCS

and  $L^4$ He is 2.5 mm thick, 45 mm long and 476 mm in diameter.

2.4.1. External Plumbing



Figure 4: System design of external Plumbing for helium tank. The vapors flow through the vent line while the fill line is closed. When the instrument is in pre-flight configuration and during ascent the values 1, 3 and motorized value 2 are closed, while value 2 is open. When the system is at float in the stratosphere, the motorized valve 2 is opened and pumping is restored, due to the external pressure of about 3 mbar. Valve 2 is used only in the laboratory, during pump-down (see text).

The filling and venting tubes for the helium tank are externally connected through a plug-in network of valves and tubes, which allows to adjust onthe-fly the line impedance of the pumping system used to lower the pressure 220 over the liquid (see figure 4). This system also features a remotely-operated motorized valve, used to seal the tank before launch and open it to the outside pressure when the cryostat is at float in the stratosphere. The motorized valve is a critical subsystem, and potentially a single-point failure. For this reason we avoided unnecessary complexity in the motor control 225 electronics, which allows only for two states of the value: fully open or completely shut, without any control of the intermediate positions. However, in the laboratory, during the L<sup>4</sup>He pump-down process, we have to finetune dynamically the pumping speed, to avoid Taconis oscillations. For this reason we have added a manually operated valve (valve 2 in figure 4) in series 230 to the motorized valve. The manually operated valve 2 is used only during the pump-down and left wide-open when the pump-down is completed.

9

### 3. Window and Filters

245

250

The OLIMPO focal plane features a large throughput (~2 cm<sup>2</sup>sr). This means that the stack of filters must limit out-of-band radiation which would represent a very large heat load. The radiation enters the cryostat through the window and encounters the first thermal shader, located right behind it. The window is 8 mm thick and it is made of HDPE (High Density Polyethylene), a material with high transmittance and low reflectance at mm wavelengths [17], thus ensuring minimum excess background on the detectors, while preserving excellent mechanical strength and vacuum tightness.

In order to avoid interference due to internal reflections, an anti-reflection coating made of porous Teflon on a low-density polyethylene substrate was applied on the HDPE surfaces. The 130 mm aperture of the window is determined by the configuration of the cold optics and by the size of the radiation beam.



Figure 5: Transmission spectra of the low-pass filters mounted in the OLIMPO cryostat.

Then, the radiation passes through the second thermal shader, the 1 THz low-pass filters, and the third thermal shader, all three anchored to the LN stage. After passing through the fourth thermal shader and the 775 GHz low-pass filter mounted on the vapor <sup>4</sup>He shield at 35 K (in-flight the temperature is about 21 K), the incoming radiation encounters the 630 GHz low-pass filter mounted on the *optics box* shield at about 2 K. Fig. 5 shows the transmission spectra of the three low-pass filters mounted in the OLIMPO cryostat.



Figure 6: Measured He evaporation rate (black dots) and empirical fit (red solid line) versus vapor cooled shield temperature.

## 255 4. Laboratory performance of the cryostat

The cryostat has been assembled and tested in the lab in more than 30 complete cooling cycles. A 40 m<sup>3</sup>/h rotary pump has been used to pump on the L<sup>4</sup>He bath, reaching a steady state temperature, after a  $\sim$ 12 hours long pumpdown, of 1.65 K.

We have experimented different settings for the gas flow in the long VCS serpentine and in the top low-impedance heat exchanger described in section §2.3. The best configuration (no oscillations, temperature gradient of a few K across the shield) was found with the valve in a midway position, allowing for a non-negligible gas flux in the long VCS serpentine, while most of the gas ( $\sim 70$  %) flows through the top low-impedance heat exchanger.

The total heat load on the L<sup>4</sup>He in steady-state conditions (and with LN at 77 K) is of the order of 80 mW, as estimated from the measured L<sup>4</sup>He evaporation rate (0.10 l/h). The resulting hold-time for both L<sup>4</sup>He and LN is about 15 days, as confirmed by several complete operation cycles in laboratory.

The best fit of the data results with an empirical relation  $(y[l/h] = 6.58 \cdot e^{-\frac{x[K]}{3.96}} + 0.1)$  between the shield temperature and the L<sup>4</sup>He evaporation rate is shown in fig.6, derived from measurements taken during the pump-down process (with LN<sub>2</sub> at 77 K). This relation is useful to approximately estimate the evaporation rate during the operations of the cryostat at balloon altitude, in the absence of a flux meter, since the VCS temperature is monitored continuously during the flight (see discussion in §6).

a

270

275

260

# 5. <sup>3</sup>He Refrigerator

We have designed a simple single-shot <sup>3</sup>He refrigerator, composed of a cryopump, a heat exchanger in thermal contact with the <sup>4</sup>He tank, and an 280 evaporator. The fridge is very similar to the one used in the BOOMERanG experiment [18], with the exception of the gas–gap heat switch, used this time to connect the cryopump to the <sup>4</sup>He bath when needed. The fridge contains about 40  $l_{STP}$  of <sup>3</sup>He, required to keep the four detector arrays in their aluminum holders and their horn arrays at 300 mK for about 16 days. 285 The total geometrical volume available inside the fridge is  $1000 \text{ cm}^3$ , but the available volume for the gas is about 80 % of that, after accounting for the volume of the carbon grains inside the cryopump. The internal gas pressure at room temperature is thus of the order of 4 MPa, so all the parts of the evaporator must be designed to withstand high internal pressure with 290 a reasonable safety factor.

spheres, with a diameter of 120 mm, and is enclosed in a gold-plated copper shield, to prevent radiation to warm the optics box when the active carbon is heated for fridge cycling. The thickness of the hemispheres (T = 2 mm)295 has been calculated from the equation [16]:

$$T = \frac{PDX}{5R\Psi} \tag{3}$$

where P is the pressure, D is the diameter, X is the safety factor (which is about 5), R is the breaking strength of material,  $\Psi$  is a form factor (23/22) for a sphere).

The spherical cryopump is composed of two welded stainless steel half-

In order to increase the temperature uniformity of the carbon grains, we have inserted a copper finger, crossing the cryopump along a diameter. One end of this tube is connected to the cryopump heater resistor, and to the cryopump–exchanger thermal link, operated through the gas–gap heat switch.

Two fiberglass supports, above and below the cryopump, connect the cryopump to the copper support of the exchanger, taking care of protection against radial accelerations, as shown in figure 7. Further protection against radial accelerations is provided for the evaporator chamber, which is secured coaxially with the main axis of the cryostat by three insulating radial spikes (in vespel, with a heat load lower than  $\sim 1 \ \mu W$ ) supported by an aluminum 310 ring structure anchored on three Vespel tubes. Moreover, thanks to the optical scheme and the consequent arrangement of the 4 detector arrays, the thermal link structure connecting the evaporator to the detector arrays ensures additional structural stiffness along three axes.

315

300

305

The cryopump and the evaporator are connected to the condensation section by means of thin–wall (t ~ 150  $\mu m$ ) stainless steel tubes (9.5 mm in diameter for the evaporator-condenser section, 10 mm in diameter for

#### Journal Pre-proofs



Figure 7: Left panel: picture of <sup>3</sup>He refrigerator (upside down in this picture). Note in the bottom part of the picture the spherical cryopump, with the copper link that connected the inner copper tube of cryopump with heater and heat switch, and the two fiberglass supports; in the middle the thin wall stainless steel tube connecting the cryopump to the heat exchanger at  $L^4$ He temperature, in the top part the aluminum structure of the evaporator, which is supported by three Vespel tubes.

 $Right\ panel:$  drawing of  $^{3}{\rm He}$  refrigerator. GGHSW is the gas–gap heat switch , CP is the cryopump and EXCH is the heat exchanger.

the condenser–cryopump section), which prevent damage from axial accelerations, while ensuring good thermal insulation from the heat exchanger.

The safety factor to withstand the internal pressure is  $X \sim 5$ . The evaporator is a 60 mm diameter stainless steel cylinder (D), with a thickness (T) of 2 mm, as computed from the equation [16]:

$$T = \frac{PDX}{2R\alpha} \tag{4}$$

where  $\alpha$  is a coefficient depending on the quality of weldings, in our case is  $\alpha \sim 0.8$ , P is the pressure and R is the breaking strength of material.

While the thickness of the copper baseplate, 10 mm, has been computed from equation 2, in this case the safety factor is  $X \sim 4$  and R is smaller than the breaking strength of copper due to welding.

The entire cryogenic system was installed at the telescope and operates

325



Figure 8: Picture of the cryogenic instrument insert (shown upside down with respect to its positioning inside the  $L^4$ He cryostat). We can see the four arrays of detectors, the thermal link, and coaxial cables between arrays and stage at 2 K. On the right we can see the Refrigerator with its cryopump shield and its evaporator with the aluminum structural support.

in a wide range of inclinations from vertical, due to the elevation range of the telescope boresight, ranging from  $e = 0^{\circ}$  for calibration to  $e = 55^{\circ}$  at the maximum useful elevation. This means that the symmetry axis of the cryostat will change from vertical to  $55^{\circ}$  offset from vertical. These movements can change the exchange surface between the L<sup>3</sup>He and the bottom flange of the evaporator, thus changing the temperature gradient between the liquid and the detectors. To minimize this effect, the inner surface of the copper baseplate has been increased by adding a walled structure, shaped to maintain a more uniform distribution of the liquid even when tilted from vertical.

The thermal contact between the copper flange of the evaporator and the array holders is obtained by means of a gold-plated electrolytic tough pitch (ETP) copper link.

In figure 8 we can see a picture of the Cold Optics with the  ${}^{3}\text{He}$  re-

frigerator, the four arrays of detectors, supported by Vespel tubes for the High Frequency arrays, and by thin perforated–layers of Vespel for the Low Frequency arrays, and the thermal link connecting the refrigerator to the arrays.

345

In figure 9 we plot the temperatures of the cryopump, heat exhanger and evaporator during a typical cooling cycle. The cycle can start when the condensation point has been cooled by the main <sup>4</sup>He bath below 1.7 K, the cryopump is also cold, and all the gas is adsorbed. Then electrical 350 power is applied to the cryopump heater. The cryopump heater power is not constant, it is adjusted dynamically to optimize the desorption of the gas, while keeping the condensation point reasonably cold. In the initial phase a significant power  $\sim 0.2$  W is applied to the cryopump heater, to warm it up and initiate desorption. This results in a significant temperature 355 increase of the condensation point. After this initial phase, when most of the gas is desorbed, the heating power can be gradually reduced, while the cryopump continues to warm-up, so that the condensation point cools down. This improves the efficiency of gas condensation. The cryopump is maintained at a temperature higher than 30 K, until the temperature of the 360 condensation point returns to 1.65 K, which ensures the condensation of the full charge of  ${}^{3}$ He. Finally (30.5 h in the plot), the heater is switched off, and the heat switch between the cryopump and the main <sup>4</sup>He bath is closed, thus activating the cryosorption and reducing the pressure on the <sup>3</sup>He bath, reaching the asymptotic temperature. 365



Figure 9: Temperature of the cryopump, exchanger and evaporator during a cooling cycle.

## 5.1. Thermometry

We equipped the cryostat with several cryogenic thermometers, both silicon diodes and germanium resistance temperature sensors, (as specified in table 1) which are needed for cryo-operation monitoring and to exclude or decorrelate systematic effects in the measurements of the CMB, due to temperature drifts.

Stages	Diode	germanium
77 K	1	0
$35~{ m K}$	2	0
2 K	2	0
Refrigerator	3	2

Table 1: Type and number of thermometers of the OLIMPO cryostat, mounted at different stages.

The silicon diodes used into the cryostat are DT-670 provided by Lake Shore Cryotronics, Inc.<sup>3</sup>. For the <sup>3</sup>He refrigerator temperatures, we used germanium resistance temperature sensors, also from Lake Shore Cryotronics, Inc, model GR-200A-1500, for the range of temperatures  $6 \div 1.4$  K and 375 GR-200A-100 for the range of temperatures  $2.1 \div 0.259$  K. These are biased with an ultrastable 1 nA sinusoidal current (AC at 3.33 Hz) and readout by means of synchronous demodulation of the voltage signal across the thermistor. Custom readout boards provide temperature measurements with absolute accuracy of  $\pm 0.5$  mK at 300 mK and  $\pm 0.1$  K at 1 K for the germa-380 nium resistance temperature sensors, and  $\pm 0.5$  K in the 4.2 $\div$ 77 K range for the silicon diodes. The performance of the KID detectors depends strongly on their operating temperature and the temperature of the radiative environment. The accuracy of the thermometry in flight is thus necessary to confirm that the detectors operate in conditions within the range of temper-385 atures where they have have been calibrated in the laboratory.

## 5.2. Refrigerator load curve

In order to characterize the performance of the refrigerator we performed a specific cooling cycle, without the four arrays and their cables (in terms of thermal load, without 30  $\mu$ W). We mounted a stable resistor (100 k $\Omega$ ) on the evaporator to dissipate controlled power levels, and measured the equilibrium temperature for different loads. For the measurement of the load curve, see figure 10, the maximum temperature variation of the exchanger was about 10 % of its base temperature (1.65 K).

<sup>&</sup>lt;sup>3</sup>https://www.lakeshore.com



Figure 10: Load Curve of the OLIMPO refrigerator: evaporator temperature vs external heat load power. With the nominal heat load of 30  $\mu$ W the temperature of the evaporator is close to 300 mK.

## <sup>395</sup> 6. In–flight performance

At 05:57 GMT on July  $14^{th}$  2018, after regular fill and pump-down procedures, the motorized valve of the <sup>4</sup>He tank was closed by command in preparation of the launch of the experiment. The vapor pressure of the <sup>4</sup>He bath started to increase slowly, and, as a consequence, the temperatures of the <sup>4</sup>He bath and of the <sup>3</sup>He evaporator increased up to 1.71 K and 307 mK, respectively.

At 07:07 GMT on July  $14^{th}$  the OLIMPO payload was launched from Longyearbyen (Svalbard, Norway.  $78^{\circ}13' \text{ N } 15^{\circ}33' \text{ E}$ ).

In figure 11 we plot the liquid nitrogen temperature from before the launch until floatation, where the temperature stabilized at  $\sim 52$  K, and nitrogen is solid. This low temperature of the LN stage is beneficial for the heat load of the LHe stage and for the duration of the mission. Note that the temperature of the nitrogen stage during the flight is measured to be less than 2 K above the sublimation temperature of solid nitrogen at the floating pressure. Since the heat load on the nitrogen tank is estimated to be  $\sim 8$  W, the effective heat conductance of the Cu braids in the LN tank is > 4 W/K.

At 11:37 GMT, at an altitude of about 35 km, the temperatures of the  ${}^{4}$ He bath and of the  ${}^{3}$ He evaporator were 1.90 K and 320 mK respectively, and the motorized valve was opened by telecommand. In this configuration the  ${}^{4}$ He tank was connected directly to the stratospheric pressure (about

415



Figure 11: Temperature of LN tank before launch (first 29 hours, with the liquid evaporating at standard pressure), during ascent, and at the beginning of floatation at 37.8 km (subsequent hours). The passage through the triple point of nitrogen is evident at hour 32.5.

3 mbar) and the  $^4\mathrm{He}$  bath vapour pressure and its temperature started to decrease.

At 12:30 GMT on July 14<sup>th</sup> the OLIMPO payload reached the float altitude of 37.8 km and both temperatures, of <sup>4</sup>He bath and of the <sup>3</sup>He evaporator, drifted down to 1.65 K and 295 mK, respectively.

In figure 12 we plot the temperatures of the <sup>4</sup>He bath and of the <sup>3</sup>He evaporator during the operations described above. A detailed view of an hour of flight data shows that temperature drifts are smaller than 0.5 mK/h. In figure 13 we plot the frequency spectrum of the temperature fluctuations shown in the zoom of figure 12.

Communications with the OLIMPO payload were lost 24 hours after the launch, therefore we do not have temperature records afterwards. However, the temperature of the VCS (as well as those of the other stages) was already stabilized before the connection was lost. Moreover, integrating the evaporation rate computed from the temperature profile of the VCS before communication interruption, we estimate that 33.5 liters of  $L^4$ He were still present in the tank at the moment of communications black-out.

At that point the temperature of the VCS was 21 K. With the LN2 tank at 77 K, this temperature would correspond to an evaporation rate of 0.134 l/h. However, at float the LN2 tank sets to a reduced pressure, and its temperature was 52 K, thus reducing the heat load on the VCS with respect to the laboratory tests. This means that an evaporation rate of 0.134 l/h is

425

435



Figure 12: The temperatures of the <sup>4</sup>He bath and of the <sup>3</sup>He evaporator during the operations in preparation of the launch, during the launch and in flight. The zoom in the top panel shows that the temperature drifts are smaller than 0.5 mK/h. In the bottom panel the *red-line* is the time when the motorized valve was closed while the *green-line* is the time when it was opened.



Figure 13: Power spectrum of the temperature fluctuations measured at the 0.3 K  $^3{\rm He}$  evaporator during one hour of flight data.

conservative, since this low temperature of the VCS is also due to the lower thermal load from the LN2 stage. Anyway, assuming this evaporation rate and the estimated volume of  $L^4$ He remaining in the tank, we estimate a minimum residual hold time in flight of about 11 days. Assuming the same heat load as in the laboratory (which is also a conservative assumption due to the lower radiative load on the <sup>4</sup>He tank from the shield) we estimate a residual hold time of 14 days.

445

The payload was separated from the balloon on July  $19^{th}$ , 2018, at 02:28 GMT. The cryostat has been subsequently recovered, and the vacuum jacket was still evacuated when arrived in our laboratory 6 months later.

## 7. Conclusions

We have developed a wet cryostat and a self contained <sup>3</sup>He refrigera-450 tor, suitable for long hold time (about 15 days) in the stratosphere. The <sup>3</sup>He refrigerator has been optimized to be robust, reliable and suitable for unmanned operation. The cryostat and the <sup>3</sup>He refrigerator have been qualified in a stratospheric balloon flight. The <sup>4</sup>He tank was sealed before the launch to maintain the  $L^4$ He pumped during ascent. Once at float, the 455 sealing valve was succesfully re-opened towards the stratospheric vacuum (3 mbar) to resume pumping. After a few hours the system stabilized at the nominal evaporation rate, fullfilling the requirement to cover the nominal duration of the flight (14 days). The <sup>3</sup>He refrigerator temperature during the flight has been 296 mK with temperature drifts smaller than 0.5 mK/h, 460 providing an excellent environment for the operation of 4 arrays of Kinetic Inductance Detectors [19, 20].

### Acknowledgements

We acknowledge the Italian Space Agency (ASI) for funding the OLIMPO 465 experiment, and organizing the 2018 launch campaign.

## References

- J. M. Lamarre, Photon noise in photometric instruments at farinfrared and submillimeter wavelengths, Applied Optics 25 (1986) 870–876. doi:10.1364/AO.25.000870.
- [2] D. Fixen, The temperature of the cosmic microwave background, The Astrophysical Journal 707 (2009) 916–920. doi:10.1088/0004-637X/707/2/916.
  - [3] E. Di Valentino et al., Exploring cosmic origins with CORE: Cosmological parameters, Journal of Cosmology and Astroparticle Physics 4 (2018) 017. doi:10.1088/1475-7516/2018/04/017.

- [4] S. Masi, A. Coppolecchia, E. S. Battistelli, P. de Bernardis, F. Columbro, G. D'Alessandro, L. Lamagna, A. Paiella, F. Piacentini, G. Presta, Balloon-borne Cosmic Microwave Background experiments, European Physical Journal Web of Conferences 209. doi:10.1051/epjconf/201920901046.
- [5] A. Coppolecchia, G. Amico, E. Battistelli, P. D. Bernardis, A. Cruciani, A. D'Addabbo, G. D'Alessandro, S. D. Gregori, M. D. Petris, R. Gualtieri, L. Lamagna, S. Masi, L. Nati, F. Nati, L. Pagano, F. Piacentini, R. Puddu, D. Romano, M. Salatino, A. Schillaci, P. Ade,
- P. Mauskpf, D. Morozov, C. Tucker, R. Tucker, A. Boscaleri, P. Camus,
  S. Colafrancesco, P. Marchegiani, G. D. Stefano, G. Romeo, M. Gervasi, M. Zannoni, K. Irwin, C. Reintsema, C. P. Novaes, A. Wuensche, D. Yvon, V. Vdovin, OLIMPO: A 4-bands imaging spectro-photometer for balloon-borne observations of the Sunyaev Zel'dovich effect, Proceedings of the International School of Physics "Enrico Fermi" 186 (2014) 257–264. doi:10.3254/978-1-61499-476-3-257.
- [6] F. Nati, P. Ade, A. Boscaleri, D. Brienza, M. Calvo, S. Colafrancesco, L. Conversi, P. de Bernardis, M. de Petris, A. Delbart, P. Hargrave, A. Iacoangeli, J. Juin, C. Magneville, C. M. Bettolo, S. Masi, P. Mauskopf, F. Melchiorri, L. Nati, A. Orlando, Y. Rephaeli, G. Romeo, J. Pansart, S. Peterzen, F. Piacentini, G. Pisano, G. Polenta, L. Salvaterra, G. Savini, C. Tucker, D. Yvon, **The OLIMPO experiment**, New Astronomy Reviews 51 (2007) 385– 389. doi:10.1016/j.newar.2006.11.066.
- [7] A. Paiella, E. S. Battistelli, F. Columbro, A. Coppolecchia, G. D'Alessandro, P. de Bernardis, , L. Lamagna, S. Masi, F. Piacentini, M. Castellano, I. Colantoni, S. Gordon, P. Mauskopf, Design and Electrical Performance of the Kinetic Inductance Detectors of the OLIMPO Experiment, 2017 16th International Superconductive Electronics Conference, ISEC 2017 2018–January (2018) 1–3. doi:10.1109/ISEC.2017.8314223.
- [8] A. Paiella, A. Coppolecchia, L. Lamagna, P. A. R. Ade, E. S. Battistelli, M. G. Castellano, I. Colantoni, F. Columbro, G. D'Alessandro, P. de Bernardis, S. Gordon, S. Masi, P. Mauskopf, G. Pettinari, F. Piacentini, G. Pisan, G. Presta, C. Tucker, Kinetic Inductance Detectors for the OLIMPO experiment: design and pre-flight characterization, Journal of Cosmology and Astroparticle Physics 2019 (01) (2019) 039. doi:10.1088/1475-7516/2019/01/039.
- 515 [9] A. Paiella, E. Battistelli, M. Castellano, I. Colantoni, F. Columbro, A. Coppolecchia, G. D'Alessandro, P. de Bernardis, S. Gordon, L. La-

480

490

485

495

500

505

magna, H. Mani, S. Masi, P. Mauskopf, G. Pettinari, F. Piacentini, G. Presta, **Kinetic Inductance Detectors and readout electronics for the OLIMPO experiment**, Journal of Physics: Conference Series 1182 (2019) 012005. doi:10.1088/1742-6596/1182/1/012005. URL https://iopscience.iop.org/article/10.1088/1742-6596 /1182/1/012005

- [10] A. Schillaci, G. D'Alessandro, P. de Bernardis, S. Masi, C. P. Novaes, M. Gervasi, M. Zannoni, Efficient differential Fouriertransform spectrometer for precision Sunyaev-Zel'dovich effect measurements, Astronomy and Astrophysics 565 (2014) A125. doi:10.1051/0004-6361/201423631.
- [11] P. de Bernardis, S. Colafrancesco, G. D'Alessandro, L. Lamagna, P. Marchegiani, S. Masi, A. Schillaci, Low-resolution spectroscopy of the Sunyaev-Zel'dovich effect and estimates of cluster parameters, Astronomy and Astrophysics 538 (2012) A86. doi:10.1051/0004-6361/201118062.
- [12] R. A. Sunyaev, Y. B. Zeldovich, The Observations of Relic Radiation as a Test of the Nature of X-Ray Radiation from the Clusters of Galaxies, Comments on Astrophysics and Space Physics 4 (1972) 173.
- [13] S. Masi, P. Cardoni, P. de Bernardis, F. Piacentini, A. Raccanelli, F. Scaramuzzi, A long duration cryostat suitable for balloon borne photometry, Cryogenics 39 (1999) 217–224. doi:10.1016/s0011-2275(99)00018-1.
- B. P. Crill, P. A. R. Ade, D. R. Artusa, R. S. Bhatia, J. J. Bock, A. Boscaleri, P. Cardoni, S. E. Church, K. Coble, P. de Bernardis, G. de Troia, P. Farese, K. M. Ganga, M. Giacometti, C. V. Haynes, E. Hivon, V. V. Hristov, A. Iacoangeli, W. C. Jones, A. E. Lange, L. Martinis, S. Masi, P. V. Mason, P. D. Mauskopf, L. Miglio, T. Montroy, C. B. Netterfield, C. G. Paine, E. Pascale, F. Piacentini, G. Polenta, F. Pongetti, G. Romeo, J. E. Ruhl, F. Scaramuzzi, D. Sforna, A. D. Turner, BOOMERANG: A Balloon-borne Millimeter-Wave Telescope and Total Power Receiver for Mapping Anisotropy in the Cosmic Microwave Background, The Astrophysical Journal Supplement Series 148 (2003) 527–541. doi:10.1086/376894.
- [15] S. Masi, P. A. R. Ade, J. J. Bock, J. R. Bond, J. Borrill, A. Boscaleri, P. Cabella, C. R. Contaldi, B. P. Crill, P. de Bernardis, G. de Gasperis, A. de Oliveira-Costa, G. de Troia, G. di Stefano, P. Ehlers, E. Hivon, V. Hristov, A. Iacoangeli, A. H. Jaffe, W. C. Jones, T. S. Kisner, A. E. Lange, C. J. MacTavish, C. Marini Bettolo, P. Mason, P. D. Mauskopf,

520

525

530

540

535

545

550

T. E. Montroy, F. Nati, L. Nati, P. Natoli, C. B. Netterfield, E. Pascale, F. Piacentini, D. Pogosyan, G. Polenta, S. Prunet, S. Ricciardi, G. Romeo, J. E. Ruhl, P. Santini, M. Tegmark, E. Torbet, M. Veneziani, N. Vittorio, Instrument, method, brightness, and polarization maps from the 2003 flight of BOOMERanG, Astronomy and Astrophysics 458 (2006) 687–716. doi:10.1051/0004-6361:20053891.

[16] R. R. Conte, Éléments de Cryogénie, Masson & Cie (1970) 73.

565

570

560

[17] G. D'Alessandro, A. Paiella, A. Coppolecchia, M. G. Castellano, I. Colantoni, P. de Bernardis, L. Lamagna, S. Masi, Ultra high molecular weight polyethylene: Optical features at millimeter wavelengths, Infrared Physics & Technology 90 (2018) 59–65. doi:10.1016/j.infrared.2018.02.008.

[18] S. Masi, E. Aquilini, P. Cardoni, P. de Bernardis, L. Martinis, F. Scaramuzzi, D. Sforna, A self-contained <sup>3</sup>He refrigerator suitable for long duration balloon experiments, Cryogenics 38 (1998) 319–324. doi:10.1016/S0011-2275(97)00174-4.

[19] S. Masi, P. de Bernardis, A. Paiella, F. Piacentini, L. Lamagna, A. Coppolecchia, P. A. R. Ade, E. S. Battistelli, M. G. Castellano, I. Colantoni,

F. Columbro, G. D'Alessandro, M. De Petris, S. Gordon, C. Magneville,
P. Mauskopf, G. Pettinari, G. Pisano, G. Polenta, G. Presta, E. Tommasi, C. Tucker, V. Vdovin, A. Volpe, D. Yvon, Kinetic Inductance
Detectors for the OLIMPO experiment: in-flight operation and performance, Journal of Cosmology and Astroparticle Physics 2019 (07) (2019) 003. doi:10.1088/1475-7516/2019/07/003.

[20] A. Paiella, P. A. R. Ade, E. S. Battistelli, M. G. Castellano, I. Colantoni, F. Columbro, A. Coppolecchia, G. D'Alessandro, P. de Bernardis, M. De Petris, S. Gordon, L. Lamagna, C. Magneville, S. Masi, P. Mauskopf, G. Pettinari, F. Piacentini, G. Pisano, G. Polenta, G. Presta, E. Tommasi, C. Tucker, V. Vdovin, A. Volpe, D. Yvon, Inflight performance of the LEKIDs of the OLIMPO experiment, J. Low Temp. Physics. 199 (1-2) (2020) 491–501. doi:10.1007/s10909-020-02372-y.

575

580

# **Declaration of interests**

 $\boxtimes$  The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

The authors declare the following financial interests/personal relationships which may be considered as potential competing interests:



The long duration cryogenic system of the OLIMPO balloon-borne experiment: design and infight performance:

A. Coppolecchia, L. Lamagna, S. Masi, P.A.R. Ade, G. Amico, E.S. Battistelli, P. de Bernardis, F. Columbro, L. Conversi, G. D'Alessandro, M. De Petris, M. Gervasi, F. Nati, L. Nati, A. Paiella, F. Piacentini, G. Pisano, G. Presta, A. Schillaci, C. Tucker, M. Zannoni.

# **Highlights:**

The paper describes the performance of the cryogenic system of the OLIMPO balloon-borne telescope. The cryo system consists of a large wet cryostat (LN and LH<sub>e</sub>) hosting a self-contained  $3H_e$  Refrigerator. The cold volume available for instrumentation (optics and detection system) is more than 60 liters. Special care has been taken to ensure high stability and stiffness of the instrument inserted for cryostat tilt up to 55° from vertical. The cryogenic system has been designed to operate continuously and autonomously in the stratosphere for two weeks. The system has been validated in the lab and flown successfully at 37.8 Km of altitude in July 2018. The  $3H_e$  refrigerator temperature during the flight was 296 mK with a thermal stability better than 0.5 mK/h, providing an excellent environment for the operation of the detectors.

## Authors statement

This work is the result of a collaboration effort within the OLIMPO experiment. All the Authors of this paper have contributed significantly to several activities. In particular, the team at the Physics Department of Sapienza University of Rome ( A. Coppolecchia, L. Lamagna, S. Masi, G. Amico, E.S. Battistelli, P. de Bernardis, F. Columbro, L. Conversi, G. D'Alessandro, M. De Petris, F. Nati, L. Nati, A. Paiella, F. Piacentini, G. Presta, A. Schillaci ) has contributed mainly to the design, fabrication, test and optimization of the cryogenic system; the team of Milan Bicocca ( M. Gervasi, M. Zannoni ) has contributed mainly to the tests, the team at Cardiff University ( P.A.R. Ade, G. Pisano, C. Tucker ) contributed mainly to the filters chain design, fabrication, test, validation. The first three authors also took care of paper writing and paper-specific data analysis.