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# Uncovering the origin of enhanced field emission properties of rGO–MnO<sub>2</sub> heterostructures: a synergistic experimental and computational investigation†

 Sachin R. Rondiya,<sup>a</sup> Indrapal Karbhal,<sup>b</sup> Chandradip D. Jadhav,<sup>c</sup> Mamta P. Nasane,<sup>d</sup> Thomas E. Davies,<sup>a</sup> Manjusha V. Shelke,<sup>b</sup> Sandesh R. Jadhkar,<sup>d</sup> Padmakar G. Chavan<sup>e</sup> and Nelson Y. Dzade<sup>\*a</sup>

The unique structural merits of heterostructured nanomaterials including the electronic interaction, interfacial bonding and synergistic effects make them attractive for fabricating highly efficient optoelectronic devices. Herein, we report the synthesis of MnO<sub>2</sub> nanorods and a rGO/MnO<sub>2</sub> nano-heterostructure using low-cost hydrothermal and modified Hummers' methods, respectively. Detailed characterization and confirmation of the structural and morphological properties are done *via* X-ray Diffraction (XRD), Field Emission Scanning Electron Microscopy (FESEM) and Transmission Electron Microscopy (TEM). Compared to the isolated MnO<sub>2</sub> nanorods, the rGO/MnO<sub>2</sub> nano-heterostructure exhibits impressive field emission (FE) performance in terms of the low turn-on field of 1.4 V μm<sup>-1</sup> for an emission current density of 10 μA cm<sup>-2</sup> and a high current density of 600 μA cm<sup>-2</sup> at a relatively very low applied electric field of 3.1 V μm<sup>-1</sup>. The isolated MnO<sub>2</sub> nanorods display a high turn-on field of 7.1 for an emission current density of 10 μA cm<sup>-2</sup> and a low current density of 221 μA cm<sup>-2</sup> at an applied field of 8.1 V μm<sup>-1</sup>. Besides the superior FE characteristics of the rGO/MnO<sub>2</sub> nano-heterostructure, the emission current remains quite stable over the continuous 2 h period of measurement. The improvement of the FE characteristics of the rGO/MnO<sub>2</sub> nano-heterostructure can be ascribed to the nanometric features and the lower work function (6.01 and 6.12 eV for the rGO with 8% and 16% oxygen content) compared to the isolated α-MnO<sub>2</sub>(100) surface ( $\Phi = 7.22$  eV) as predicted from complementary first-principles electronic structure calculations based on density functional theory (DFT) methods. These results suggest that an appropriate coupling of rGO with MnO<sub>2</sub> nanorods would have a synergistic effect of lowering the electronic work function, resulting in a beneficial tuning of the FE characteristics.

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## 1. Introduction

Nanoscale heterostructure design comprising different material compositions is emerging as an attractive strategy and essential building block for functional devices to achieve improved performance. The desired physicochemical properties of the

participating nanomaterials in nanostructured hybrids/composites complement each other by tuning their electronic properties to meet the requirements for the fabrication of efficient electronic devices.<sup>1</sup> Reduced Graphene Oxide (rGO) is an attractive and ideal nanomaterial to be paired with another suitable semiconductor for the development of multifunctional heterostructures because of its unique electronic properties, high electrical conductivity ( $5 \times 10^{-3}$  S cm<sup>-1</sup>), flexible structure, high aspect ratio, and high specific surface area (2630 m<sup>2</sup> g<sup>-1</sup>).<sup>2–5</sup> Owing to its unique physicochemical properties, rGO is being recognized as a material of great interest for potential applications in nanoelectronics,<sup>6</sup> nanoelectromechanical systems,<sup>7</sup> sensors,<sup>8</sup> catalysis,<sup>9</sup> energy storage devices,<sup>10,11</sup> optics,<sup>12</sup> and field emission (FE).<sup>13–15</sup> There exist several reports of the successful synthesis of rGO or modified graphene heterostructures with various semiconducting nanomaterials such as TiO<sub>2</sub>, SnO<sub>2</sub>, ZnO, Si, CdSe, *etc.* in the literature.<sup>16–18</sup>

<sup>a</sup>School of Chemistry, Cardiff University, Main Building, Park Place, Cardiff, CF10 3AT, Wales, UK. E-mail: RondiyaS@cardiff.ac.uk; DzadeNY@cardiff.ac.uk

<sup>b</sup>Physical and Materials Chemistry Division, CSIR-National Chemical Laboratory, Pune 411008, MH, India

<sup>c</sup>The State Key Laboratory of Refractories and Metallurgy, Institute of Advanced Materials and Nanotechnology, College of Materials and Metallurgy, Wuhan University of Science and Technology, Wuhan 430081, P. R. China

<sup>d</sup>Department of Physics, Savitribai Phule Pune University, Pune 411007, India

<sup>e</sup>Department of Physics, School of Physical Sciences, Kavayitri Bahinabai Chaudhari North Maharashtra University, Jalgaon 425001, India

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For field emission applications, where electrons are extracted from the surface of a metal/semiconductor by an electrostatic field through quantum mechanical tunneling, rGO-based nanocomposites such as rGO-Bi<sub>2</sub>S<sub>3</sub><sup>19</sup> and WS<sub>2</sub>-RGO,<sup>20</sup> have demonstrated superior field emission properties. Among transition metal oxides, manganese dioxide (MnO<sub>2</sub>) has attracted increasing interest for field emission applications, owing to their wide structural diversity combined with unique chemical and physical properties.<sup>21,22</sup> The advantages of MnO<sub>2</sub> as field emitter are the lower cost for raw materials and the fact that manganese is more environmentally friendly than other metal oxide.<sup>23,24</sup> MnO<sub>2</sub> has also attracted a lot of attention as an electrochemical pseudocapacitor material due to its high theoretical capacitance (1370 F g<sup>-1</sup>).<sup>25-27</sup> Wu *et al.* reported inspiring results such as low turn-on field value of 8.4 V μm<sup>-1</sup> at current density of 1 μA cm<sup>-2</sup> and maximum emission current density of 160 μA cm<sup>-2</sup> at an applied field 18 V μm<sup>-1</sup>.<sup>21</sup> The field emission applications of MnO<sub>2</sub> is, however, limited by its low specific surface area and poor electrical conductivity (10<sup>-5</sup> to 10<sup>-6</sup> S cm<sup>-1</sup>). Compared to its flat films, by fabricating rGO/MnO<sub>2</sub> nanocomposite the interface area can be significantly enlarged, which is desirable for field emission application.<sup>28</sup> Besides, rGO is solution-processable and thus can be deposited in large areas onto different kinds of substrates enabling simple and cost-effective fabrication of field electron emitters for display applications. The formation of rGO-MnO<sub>2</sub> nanostructures and their electrochemical performance have been extensively investigated and it was demonstrated that compared to the single metal-oxide, rGO/MnO<sub>2</sub> nanocomposites show superior electric conductivity, electric capacity and charge/discharge efficiency for supercapacitor performance.<sup>29-34</sup> These characteristics make rGO/metal-oxide nanocomposites promising materials for energy applications. Considering that field emission is geometry (shape, size, aspect ratio, alignment, and

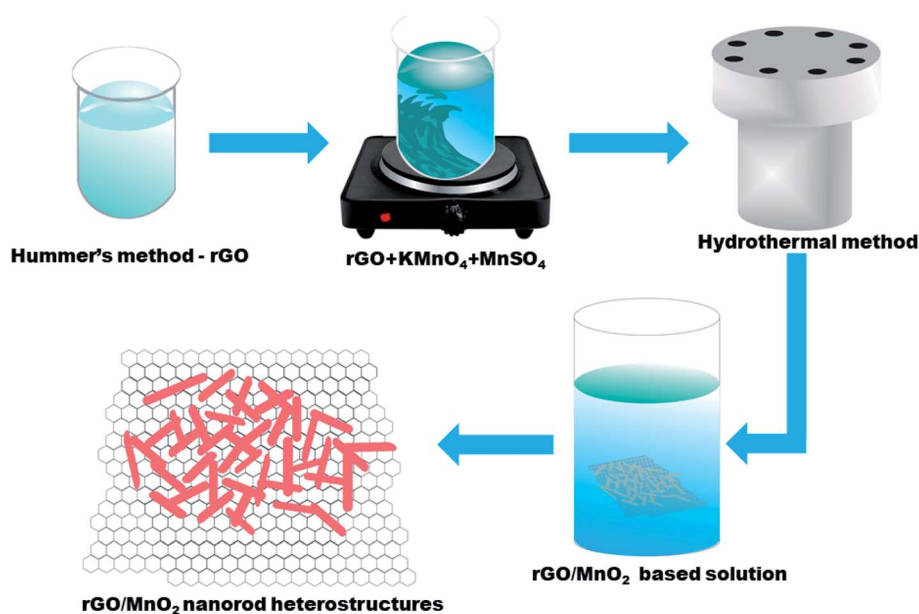
areal density of the nanostructure) and work function dependent phenomenon, well-aligned rGO/MnO<sub>2</sub> nanostructures is promising for enhancement of field emission characteristics.<sup>35</sup>

Herein, we report a simple and cost-effective solution-based method to prepared MnO<sub>2</sub> nanorods and rGO/MnO<sub>2</sub> nano-heterostructure. The structural and morphological verifications have been done by using X-ray diffraction (XRD), Field Emission Scanning Electron Microscope (FESEM), and Transmission Electron Microscopy (TEM). Finally, the field emission properties of the as-prepared MnO<sub>2</sub> nanorods and rGO/MnO<sub>2</sub> nano-heterostructure was systematically characterized and compared. The rGO/MnO<sub>2</sub> nano-heterostructure exhibits superior field emission characteristics compared to the MnO<sub>2</sub> nanorod. The rGO/MnO<sub>2</sub> nano-heterostructure demonstrates a low turn-on field of 1.4 V μm<sup>-1</sup> for an emission current density of 10 μA cm<sup>-2</sup> compared to 7.1 V μm<sup>-1</sup> for MnO<sub>2</sub> nanorod. The combined contribution of the sharp edges of the thin rGO sheets and high aspect ratio of the MnO<sub>2</sub> nanorods, coupled with synergetic effect in the rGO/MnO<sub>2</sub> nano-heterostructure are responsible for the observed enhanced field emission behavior. Consistent with the experimental data, our complementary first-principles DFT calculations predict lower work function for the rGO/MnO<sub>2</sub> nano-heterostructure compared to the isolated MnO<sub>2</sub> as the primary origin for improved field emission.

## 2. Experimental, characterization, and computational methods

### 2.1 Synthesis of MnO<sub>2</sub> and rGO/MnO<sub>2</sub>

The rGO has been synthesized by modified Hummer's method<sup>36</sup> whereas the MnO<sub>2</sub> nanorods has been synthesized by the hydrothermal method.<sup>37</sup> For the preparation of the rGO/MnO<sub>2</sub> composite, 1 mg ml<sup>-1</sup> rGO was dispersed in the 100 ml of DI



Scheme 1 The facile synthesis of MnO<sub>2</sub> nanorods and rGO/MnO<sub>2</sub> heterostructure.



water in a beaker. Later, 10 mM of  $\text{KMnO}_4$  and 10 mM of  $\text{MnSO}_4$  were added into the rGO solution and stirred for 30 min form a homogenous solution. The prepared solution was transferred into a stainless steel autoclave and kept at 160 °C for 24 h. After cooling to room temperature, the material was filter washed with DI water and ethanol to obtain rGO/ $\text{MnO}_2$  composite, which was dried in an oven at 80 °C for 12 hours and used for various characterizations presented next. Scheme 1 represents the synthesis steps which were followed for synthesis of rGO,  $\text{MnO}_2$  and rGO/ $\text{MnO}_2$  heterostructure.

## 2.2 Materials characterization

The  $\text{MnO}_2$  nanorods and rGO/ $\text{MnO}_2$  nano-heterostructure were characterized by various complementary experimental methods. The XRD patterns were obtained with a Bruker D8 Advance X-ray diffractometer using the Cu K $\alpha$  line ( $\lambda = 1.54 \text{ \AA}$ ) at 1° grazing angle. The HR-TEM micrographs and selected area electron diffraction (SAED) patterns were obtained with JEOL-JEM 2100 microscope operating at 200 kV. The samples were dry dispersed over 300 mesh copper grids coated with holey carbon film. A Field Emission Scanning Electron Microscope (FEG-SEM Model – Tescan MAIA3) was used to examine the morphology and surface topography of the  $\text{MnO}_2$  nanorods and rGOs/ $\text{MnO}_2$  composite. The accelerating voltage was 15 kV. X-ray spectroscopy (EDS) measurements were done using Oxford Instruments X-Max<sup>N</sup> 80 detector and analyzed using Aztec software. X-ray Photoelectron Spectroscopy (XPS) was carried on the samples using a Kratos Axis Ultra DLD photoelectron spectrometer utilizing monochromatic AlK $\alpha$  radiation operating at an energy of 120 W (10 × 12 kV). Data were analyzed using Casa XPS and modified Wagner sensitivity factors as supplied by the instrument manufacturer after subtraction of a Shirley background. All spectra were calibrated to the C(1s) line taken to be 284.8 eV.

## 2.3 Field emission

The field emission studies of the  $\text{MnO}_2$  nanorods and rGO/ $\text{MnO}_2$  heterostructure were carried out in the Ultra-High Vacuum (UHV) chamber at a base pressure of  $\sim 1 \times 10^{-8}$  mbar (Excel Instruments model: I-100). Detail experimental procedure may found in our earlier paper.<sup>38</sup> The configuration of the field emission experiment steps up is shown in ESI (scheme S1).† The distance between inter-electrode was maintained at 1 mm. The area of both specimens ( $\text{MnO}_2$  nanorods and rGOs/ $\text{MnO}_2$  heterostructure) was 0.25 cm<sup>2</sup>.

## 2.4 Computational methods

The first-principles spin polarized density functional theory (DFT) calculations were performed using the Vienna *Ab initio* Simulation Package (VASP),<sup>39–41</sup> a periodic plane wave DFT code which includes the interactions between the core and valence electrons using the Project Augmented Wave (PAW) method.<sup>42</sup> An energy cut-off of 600 eV, and Monkhorst–Pack<sup>43</sup> *k*-point mesh of  $7 \times 7 \times 3$  was used to sample the Brillouin zone of bulk  $\alpha$ - $\text{MnO}_2$ . Geometry optimizations were performed based on the conjugate-gradient algorithm until the residual

Hellmann–Feynman forces on all relaxed atoms reached  $10^{-3} \text{ eV \AA}^{-1}$ . The electronic exchange–correlation potential was calculated using the Perdew–Burke–Ernzerhof (PBE) generalized gradient approximation (GGA) functional.<sup>44</sup> To accurately reproduce the experimentally known band gaps and density of states features of  $\alpha$ - $\text{MnO}_2$  and rGO, the screened hybrid functional HSE06<sup>45</sup> was used with the exchange value of 25%. The projected density of states (PDOS) was calculated using tetrahedron method with Bloch correction.<sup>46</sup>

The most stable  $\alpha$ - $\text{MnO}_2$  (100) surface<sup>47</sup> was employed to form the nano-heterostructure with rGO (rGO/ $\alpha$ - $\text{MnO}_2$ ). The  $\alpha$ - $\text{MnO}_2$  (100) surface was created from the optimized bulk material using the METADISE code, which ensures the creation of surfaces with zero dipole moment perpendicular to the surface plane. The rGO/ $\alpha$ - $\text{MnO}_2$  nano-heterostructure was constructed with  $(2 \times 4)$ - $\alpha$ - $\text{MnO}_2$ (100) and  $(5 \times 5)$ -rGO supercells. We used *k*-point meshes of  $9 \times 9 \times 1$  for the rGO monolayer,  $5 \times 5 \times 1$  for the  $\alpha$ - $\text{MnO}_2$  (100) surface, and  $5 \times 5 \times 1$  for the rGO/ $\alpha$ - $\text{MnO}_2$  composite. In each simulation cell, a vacuum region of length 20 Å was added perpendicular to the surface to avoid interactions between periodic slabs. The electrostatic potential of each surface was averaged along the *c*-direction, using the Macro Density package.<sup>48–50</sup> The work function ( $\Phi$ ) was calculated as  $\Phi = V_{\text{vacuum}} - E_{\text{F}}$ , where  $V_{\text{vacuum}}$  and  $E_{\text{F}}$  are the vacuum and Fermi level, respectively. Dipole correction perpendicular to all surfaces was accounted for, which ensured that there is no net dipole perpendicular to the surfaces that may affect the potential in the vacuum level.<sup>51–53</sup>

# 3. Results and discussions

## 3.1 Characterization of $\text{MnO}_2$ and rGO/ $\text{MnO}_2$

The crystalline structures of the  $\text{MnO}_2$  nanorods and rGO/ $\text{MnO}_2$  nano-heterostructure were confirmed by XRD and the corresponding results are presented in Fig. S1.† All the diffraction peaks in Fig. S1† can be indexed to the tetragonal crystal structure of  $\text{MnO}_2$  (ICDD card no. 72-1982) with lattice constant  $a = b = 9.815 \text{ \AA}$  and  $c = 2.847 \text{ \AA}$ . We have observed the highest growth of  $\alpha$ - $\text{MnO}_2$  in the (211) plane.<sup>54</sup> The XRD diffraction pattern of the rGO- $\text{MnO}_2$  nanostructure is shown in inset of Fig. S1.† The broad peak at  $2\theta$  around 26° corresponds to the (002) plane of the reduced graphene oxide.<sup>55</sup> The field emission scanning electron microscope (FESEM) images in Fig. 1, reveal the morphological properties of the  $\text{MnO}_2$  nanorods and rGO/ $\text{MnO}_2$  nano-heterostructure. The FESEM images recorded at different magnifications (panels a–c of Fig. 1) show the formation of randomly distributed  $\text{MnO}_2$  nanorods. The low magnification image shown in Fig. 1a depicts large coverage of the  $\text{MnO}_2$  nanorods. The diameter of the formed ultra-long nanorods is estimated in the range of 140–150 nm as revealed by the high magnification FESEM images analyses (Fig. 1b and c). It is evident from the FESEM image of the rGO/ $\text{MnO}_2$  nano-heterostructure that the  $\text{MnO}_2$  nanorods are embedded in the rGO network (Fig. 1b–d). The high magnification image in Fig. 1f reveals the enormous coverage of the rGO/ $\text{MnO}_2$  nano-heterostructure.





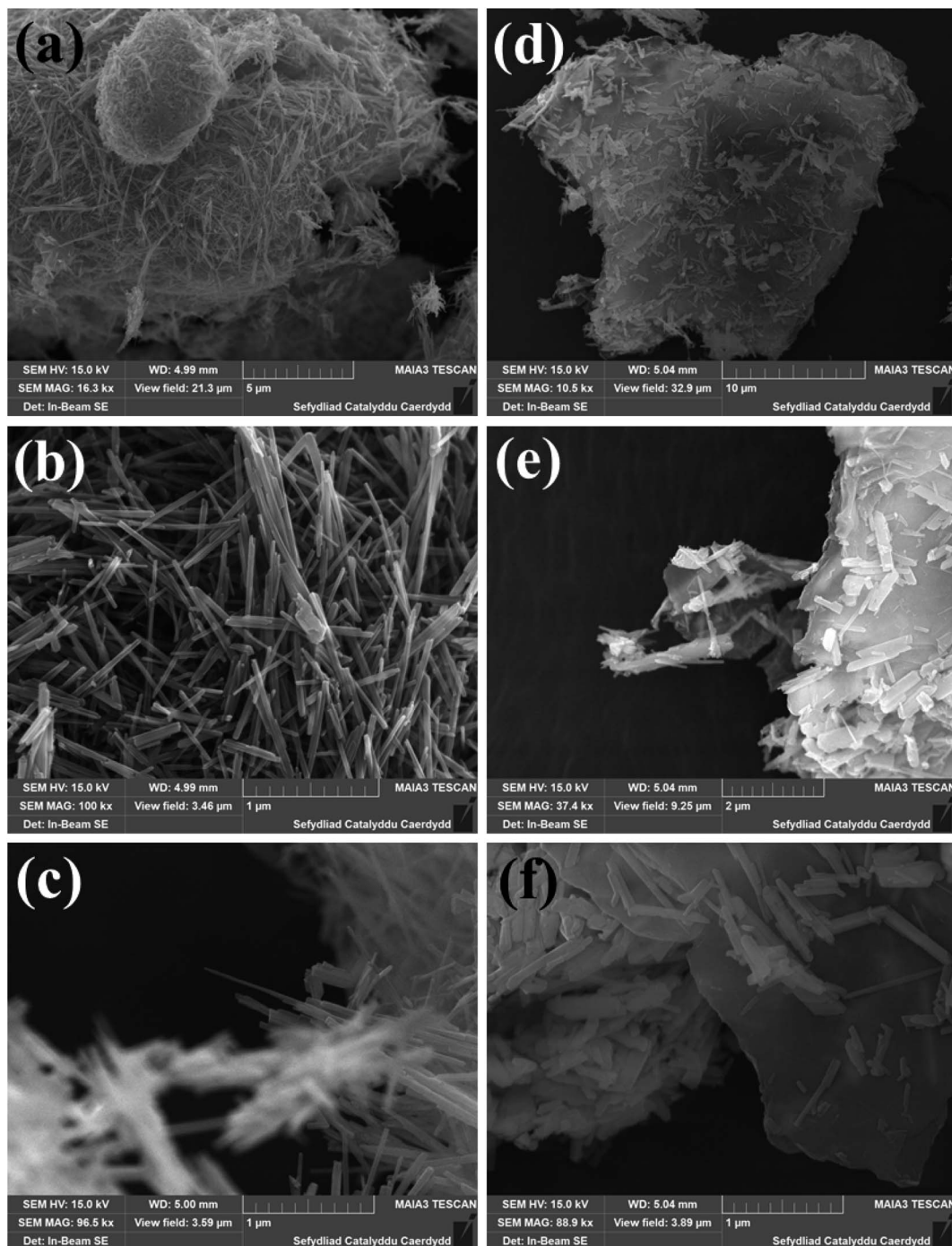


Fig. 1 FESEM images of MnO<sub>2</sub> nanorods (a–c) and rGO–MnO<sub>2</sub> heterostructures (d–f) recorded at different magnification.

Energy dispersive X-ray spectroscopy (EDX) composition analysis in the 0–10 keV energy range (ESI, Fig. S2<sup>†</sup>) confirmed that the MnO<sub>2</sub> nanorods and rGO/MnO<sub>2</sub> nano-heterostructure have the optimal stoichiometric atomic Mn : O and Mn : O : C ratios, respectively. Moreover, the FESEM-EDS elemental mapping (ESI, Fig. S2<sup>†</sup>) confirms an even distribution of the chemical constituents (Mn, O, and C). For detail structural analysis of the as-synthesized MnO<sub>2</sub> nanorods and rGO/MnO<sub>2</sub> nano-heterostructure, TEM studies were carried out. The TEM

micrograph (Fig. 2a) reveals the morphology of the MnO<sub>2</sub> nanorods, with sizes ranging between 50 and 70 nm in width and the average length of 1 μm. The selected area electron diffraction (SAED) pattern, depicted in Fig. 2b, confirms the polycrystalline nature of the MnO<sub>2</sub> nanorods. The lattice-resolved high-resolution transmission electron microscopy (HRTEM) image of the MnO<sub>2</sub> nanorod (Fig. 2c) clearly reveals its crystalline nature. The lattice fringes are clearly observed in the HR-TEM image (Fig. 2c and d) and a 0.69 nm interplanar



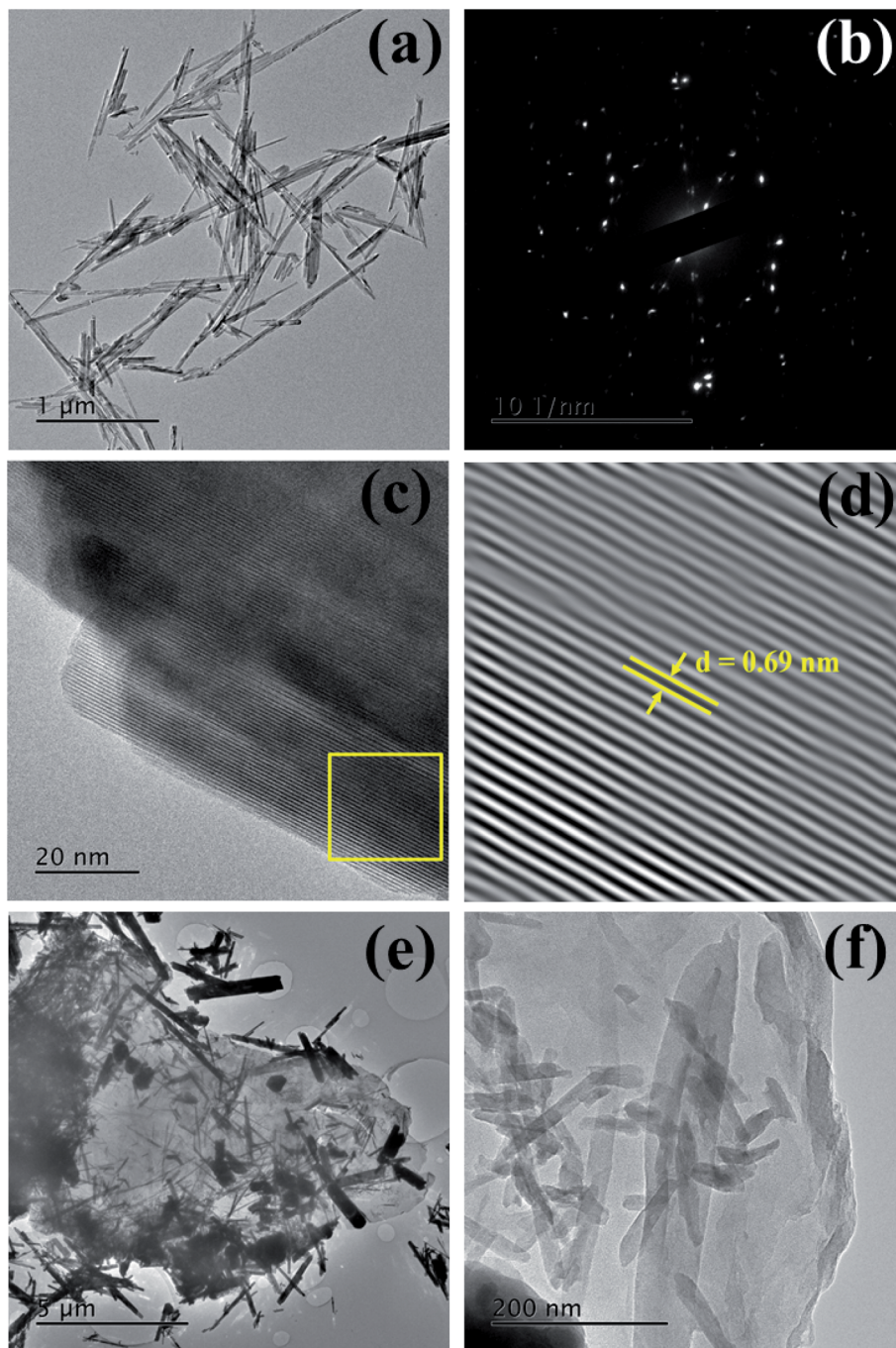


Fig. 2 (a) TEM images of MnO<sub>2</sub> nanorods (b) selected area electron diffraction (SAED) pattern of MnO<sub>2</sub> nanorods (c) HR-TEM image of MnO<sub>2</sub> nanorods with clear lattice resolution (d) Inverse Fourier transform of area shown in (c) with interplanar spacing of MnO<sub>2</sub> (e and f) TEM images of rGO-MnO<sub>2</sub> heterostructures.

distance indicates these planes to be of the (110) character.<sup>56</sup> The inverse FFT HR-TEM image and the corresponding profile plot of the MnO<sub>2</sub> nanorod are shown in Fig. S3(a and b).<sup>†</sup> A number of firmly attached  $\alpha$ -MnO<sub>2</sub> nanorods onto the rGO sheets can be clearly seen from Fig. 2e and f. X-ray photoelectron spectroscopy (XPS) analysis was applied to determine the oxidation state and elemental composition of the prepared MnO<sub>2</sub> nanorods and rGO/MnO<sub>2</sub> nano-heterostructure. The XPS

results for the MnO<sub>2</sub> nanorods are shown in Fig. S4(a and b)<sup>†</sup> and the rGO/MnO<sub>2</sub> nano-heterostructure in Fig. S4(c–e).<sup>†</sup> The peaks centered at 642.63 and 654.53 eV in the high-resolution spectrum of Mn 2p (Fig. S4a<sup>†</sup>) can be assigned to the Mn 2p<sub>3/2</sub> and Mn 2p<sub>1/2</sub> peaks, respectively, confirming the presence of MnO<sub>2</sub>.<sup>57–60</sup> The deconvolution peaks of the O 1s spectrum (Fig. S4b<sup>†</sup>) can be divided into three peaks which correspond to O–Mn bonding. The high-resolution spectrum of Mn 2p and O



1s shown in Fig. S4c and S4d,<sup>†</sup> respectively, for rGO/MnO<sub>2</sub> nano-heterostructure, confirm the presence of MnO<sub>2</sub>, whereas the C (1s) spectra (Fig. S4e<sup>†</sup>) exhibit peaks that originates from the rGO sheets.

### 3.2 Field emission investigations

Fig. 3(a and b) shows the emission current density as a function of the applied electrical field ( $J$ - $E$  curves) for the MnO<sub>2</sub> nanorods and rGO/MnO<sub>2</sub> nano-heterostructure. In this work, the turn-on field defined as the field required to draw an emission current density ( $J$ ) of 10  $\mu\text{A cm}^{-2}$  is found to be 7.1 and 1.4  $\text{V } \mu\text{m}^{-1}$  for

the MnO<sub>2</sub> nanorods and the rGO/MnO<sub>2</sub> nano-heterostructure, respectively. The rGO/MnO<sub>2</sub> nano-heterostructure also attains an impressive current density of 600  $\mu\text{A cm}^{-2}$  at an applied field of 3.1  $\text{V } \mu\text{m}^{-1}$  compared to the isolated MnO<sub>2</sub> which displayed a lower current density of 221  $\mu\text{A cm}^{-2}$  at a relatively higher applied field of 8.1  $\text{V } \mu\text{m}^{-1}$ . A comparison between the turn-on field values obtained in the present study and previously reported MnO<sub>2</sub> nanostructures and the MnO<sub>2</sub>/rGO nano-heterostructures is provided in Table 1.<sup>21,22,61-66</sup>

The Fowler-Nordheim (F-N) plot for MnO<sub>2</sub> nanorods and rGO/MnO<sub>2</sub> nano-heterostructure obtained from  $\ln(J/E^2)$  verses

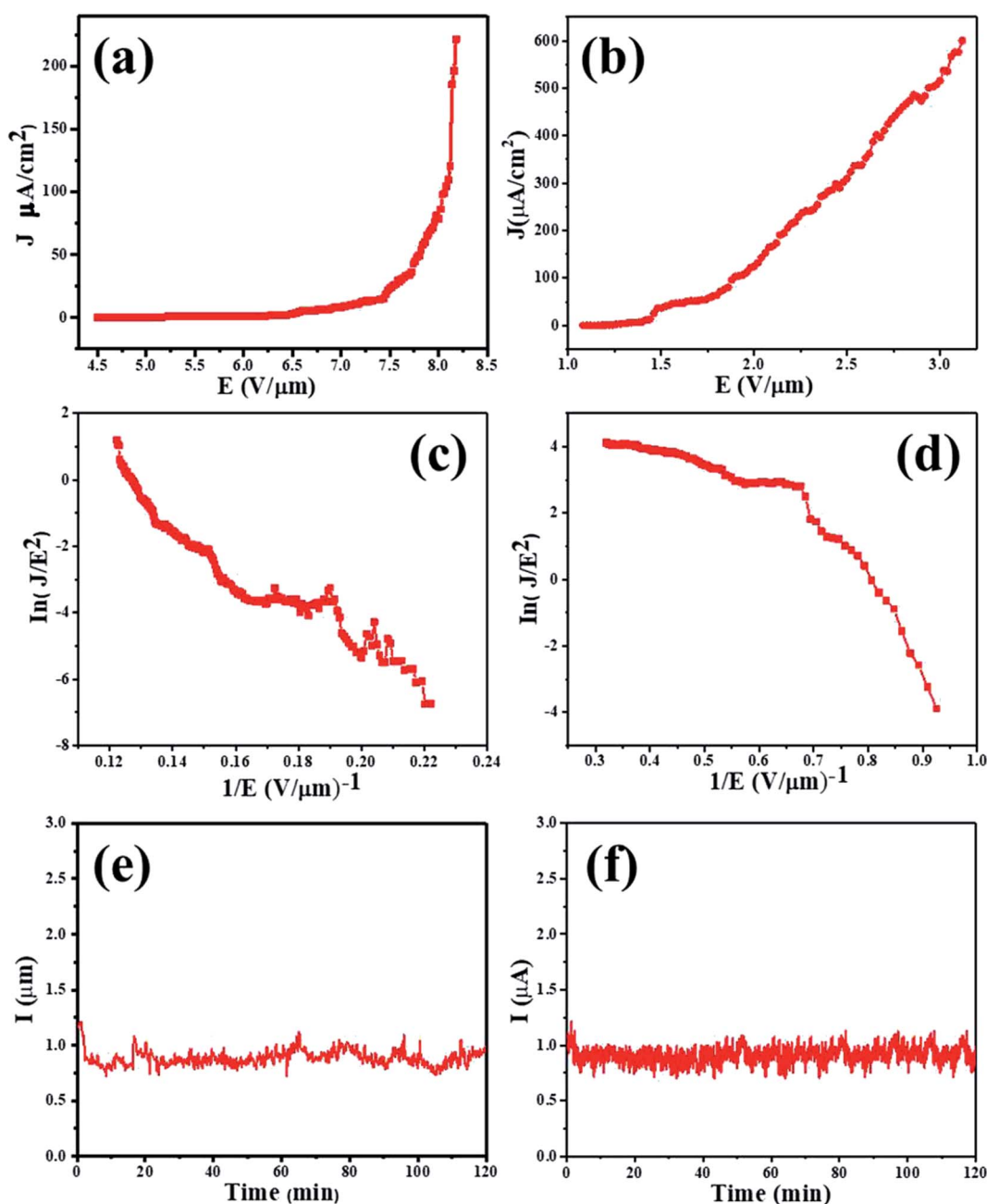


Fig. 3 FE characteristics of the MnO<sub>2</sub> nanorods and rGO-MnO<sub>2</sub> heterostructures: (a and b) emission current density versus applied electric field ( $J$ - $E$ ) curve, (c and d) F-N plot, and (e and f) emission current versus time ( $I$ - $t$ ) plot.





Table 1 Comparison of the turn-on field with reported MnO<sub>2</sub> nanomaterials and the MnO<sub>2</sub>/rGO nano-heterostructures in the present study

Sr. no.	Specimen	Turn-on field (V $\mu\text{m}^{-1}$ ) (for $J = 10 \mu\text{A cm}^{-2}$ )	Reference
1	MnO <sub>2</sub> nanotubes	8.4 (1 $\mu\text{A cm}^{-2}$ )	21
2	MnO <sub>2</sub> nanorods	5.90 (1 $\mu\text{A cm}^{-2}$ )	22
3	MnO <sub>2</sub> /rGO nanocomposite	3.6 (1 $\mu\text{A cm}^{-2}$ )	
4	rGO/TiO <sub>2</sub> nanocomposite	2.6	61
5	T-ZnO/rGO nanocomposite	1.54 (1 $\mu\text{A cm}^{-2}$ )	62
6	SnO <sub>2</sub> /rGO nanocomposite	1.8 (1 $\mu\text{A cm}^{-2}$ )	63
7	ZnO/graphene nanocomposite	1.3 (1 $\mu\text{A cm}^{-2}$ )	64
8	ZnO/graphene nanocomposite	2.1	65
9	SnO <sub>2</sub> /graphene nanocomposite	3.85 (1 $\mu\text{A cm}^{-2}$ )	66
10	MnO <sub>2</sub> nanorods	7.1	Present work
11	MnO <sub>2</sub> /rGO nano-heterostructure	1.4	

( $1/E$ ) is shown in Fig. 3(c and d). Consistent previous reports,<sup>21,22</sup> the FN plots of the MnO<sub>2</sub> nanorods and rGO/MnO<sub>2</sub> nano-heterostructure exhibit linear behavior in good agreement with their semiconducting nature. The nanometric features of the rGO/MnO<sub>2</sub> nano-heterostructure coupled with the high electrical conductivity of rGO are suggested as the key factors

behind the observed superior field emission properties.<sup>2-5</sup> Apart from the improved field emission performance, the electron emission current stability is an important parameter for device fabrication considerations. We have therefore measured the emission current as a function of time in order to ascertain the robustness of the rGO/MnO<sub>2</sub> nano-heterostructure. The

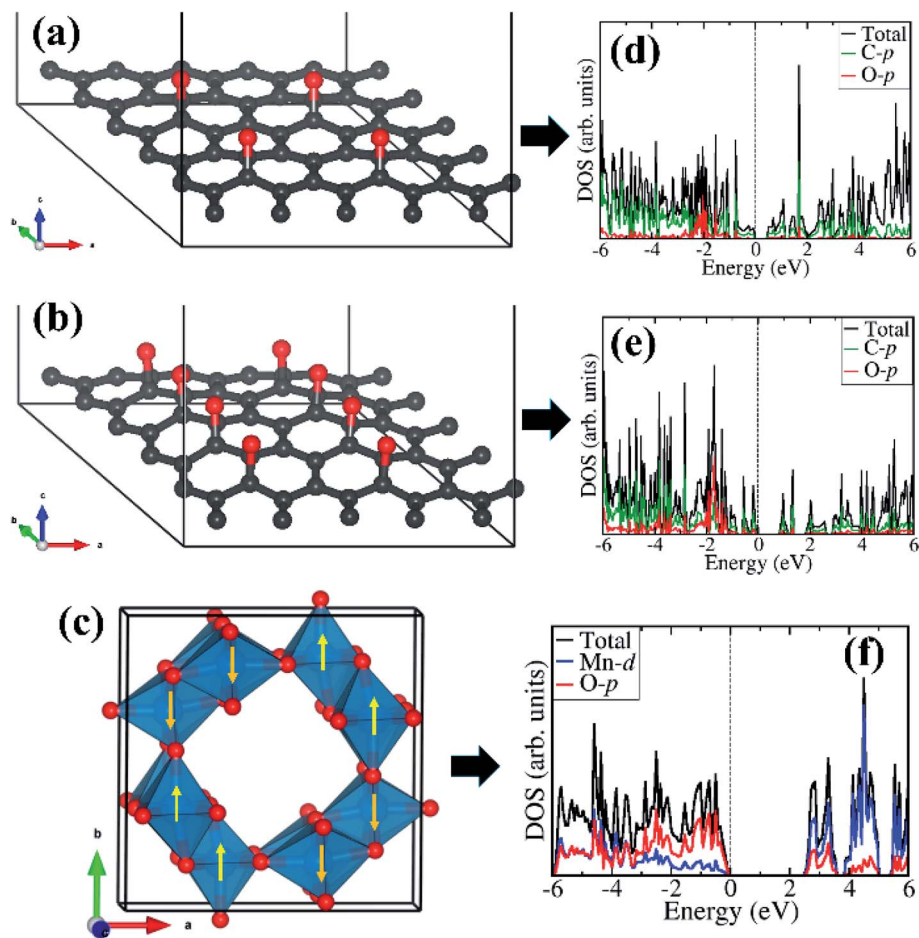


Fig. 4 Optimized structure of rGO with epoxide functional group with (a) 8% and (b) 16% oxygen content. The tetragonal structure of  $\alpha$ -MnO<sub>2</sub> showing the AFM spin ordering in (c). The corresponding partial density of states are shown in the right column (d–f). Atomic color: C = grey, O = oxygen, Mn = blue.



emission current *versus* time ( $I-t$ ) plot of the  $\text{MnO}_2$  nanorods and  $\text{rGO}/\text{MnO}_2$  nano-heterostructure were recorded continuously for 2 h at a preset value of emission current of  $\sim 1 \mu\text{A}$  as shown in Fig. 3(e and f). Generally, the results show that the emission current remains quite stable without showing any sign of diminishing over the 2 h period of continuous testing. Instabilities in the form of spikes can be attributed to the presence of residual gas molecules across the emitter surface.

### 3.3 Density functional theory analyses

Considering that field emission is a geometry and work function ( $\Phi$ ) dependent phenomenon, with lower  $\Phi$  enhancing the field emission characteristics, we have carried out first-principles density functional theory calculations to gain atomic-level insight into the electronic structure and work function of the isolated  $\text{rGO}$ ,  $\alpha\text{-MnO}_2$  (100) surface, and the  $\text{rGO}/\alpha\text{-MnO}_2$  (100) nanocomposite. As the electronic band gap of  $\text{rGO}$  vary depending on the degree of reduction, the  $\text{rGO}$  monolayer with two concentrations of epoxide functional groups with 8% and 16% oxygen contents were modelled as

shown in Fig. 4(a and b). The band gap of the  $\text{rGO}$  with 8% and 16% oxygen contents are predicted at 0.48 eV (Fig. 4d) and 0.81 eV (Fig. 4e). This is consistent with experimental data that showed that the band gap of  $\text{rGO}$  can be tuned from 0.264–0.786 eV by controlling the surface concentration of epoxide groups.<sup>67</sup> The bulk  $\alpha\text{-MnO}_2$  was modelled tetragonal crystal structure (space group  $-I4/m$ , no. 87) in the antiferromagnetic AFM-C2 configuration<sup>68</sup> as shown in Fig. 4c. A full unit cell relaxation yielded a strain-free  $\alpha\text{-MnO}_2$  with lattice parameters  $a = b = 9.763 \text{ \AA}$ ,  $c = 2.872 \text{ \AA}$ , which compares closely with known experimental data ( $a = b = 9.71 \text{ \AA}$  and  $c = 2.88 \text{ \AA}$ ).<sup>69,70</sup> The electronic band gap of  $\alpha\text{-MnO}_2$  is predicted at 2.42 eV (Fig. 4f), in close agreement with an experimental estimate of 2.23 eV<sup>71</sup> and previous theoretical prediction of 2.7 eV.<sup>72</sup> The valence band edge of  $\alpha\text{-MnO}_2$  is demonstrated to consist mainly of O-p states whereas the conduction band edge is dominated by Mn-d states. The smaller band gap of the  $\text{rGO}$  compared to the  $\text{MnO}_2$  suggests that the  $\text{rGO}$  has better electrical conductivity than the metal oxide.

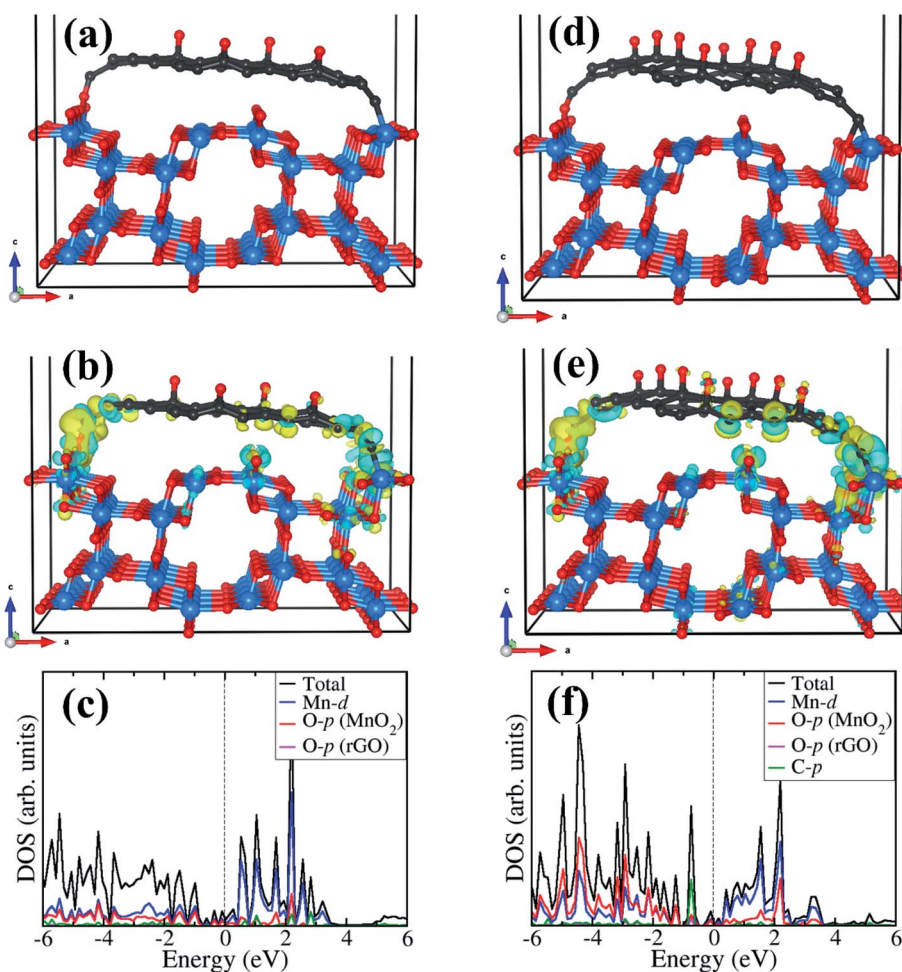


Fig. 5 Optimized structures of the  $\text{rGO}/\alpha\text{-MnO}_2(100)$  nanocomposite formed by  $\text{rGO}$  with epoxide functional group (a) 8% and (d) 16% oxygen content. The corresponding differential charge density iso surface contours and partial density of states are shown in (b) and (c) for 8% oxygen content, and in (e) and (f) for 16% oxygen content. The yellow and cyan regions indicate electron density depletion and accumulation by  $0.003 \text{ e \AA}^{-3}$ , respectively. Atomic color: C = grey, O = oxygen, Mn = blue.





The optimized structures of the rGO/ $\alpha$ -MnO<sub>2</sub>(100) nanocomposite formed by rGO with epoxide functional group with 8% and 16% oxygen contents are shown in Fig. 5a and d, respectively. The rGO is stabilized on the MnO<sub>2</sub> surface *via* C–O and C–Mn chemical bonds through the terminal C atoms. The interactions of the rGO with the MnO<sub>2</sub> surface gave rise to electron density redistribution within the rGO/ $\alpha$ -MnO<sub>2</sub>(100) nanocomposite, which was analyzed by determining the three-dimensional-charge density difference iso surface contours as shown in Fig. 5b for the rGO with 8% oxygen content and Fig. 5e for rGO with 16% oxygen content. The yellow and cyan regions represent charge depletion and accumulation in the space, respectively. We observe electron density accumulations mainly in the interfacial bonding regions in the rGO/ $\alpha$ -MnO<sub>2</sub>(100) nanocomposite, suggesting strong interactions between rGO and MnO<sub>2</sub>(100) surface. Consistent with the strong interaction, we observe strong hybridization between the C-p orbitals of the rGO with the Mn-d and O-p of the MnO<sub>2</sub> surface as shown in Fig. 5c and f, resulting in metallic conductivity of the composite systems.

The analysis of the work function ( $\Phi$ ) for the rGO monolayer,  $\alpha$ -MnO<sub>2</sub>(100) and the rGO/ $\alpha$ -MnO<sub>2</sub>(100) nanocomposite can

help us to understand the origin/direction of charge transfer at the rGO/ $\alpha$ -MnO<sub>2</sub>(100) interface. The  $\Phi$  for the isolated rGO is predicted at 5.21 and 5.85 eV for the 8% and 16% oxygen contents, respectively (Fig. 6a and b). This is consistent with a previous theoretical investigation which predicted a work function of rGO with epoxy groups to be 4.35 eV for 1.5% oxygen content and 5.6 eV for 20% oxygen content.<sup>73</sup> The  $\Phi$  for the  $\alpha$ -MnO<sub>2</sub>(100) surface is predicted at 7.22 eV (Fig. 6c), also in good agreement with earlier theoretical pH-corrected work function of 7.7 eV for the  $\alpha$ -MnO<sub>2</sub>(110) surface.<sup>72</sup> The work function of the rGO/ $\alpha$ -MnO<sub>2</sub>(100) nanocomposite is predicted at 6.01 and 6.12 eV for the rGO with 8% and 16% oxygen contents (Fig. 6d and e), both of which are lower than that of the isolated  $\alpha$ -MnO<sub>2</sub>(100) surface (Fig. 6c). The reduction in the work function of rGO/ $\alpha$ -MnO<sub>2</sub>(100) nanocomposite relative to the isolated  $\alpha$ -MnO<sub>2</sub>(100) surface can be ascribed to the interfacial bonding, electronic interaction and synergistic effects. Considering that the electron emission capability of a material is dictated by its work function, the observed superior field emission characteristics of the rGO/ $\alpha$ -MnO<sub>2</sub> nanocomposite compared to the isolated  $\alpha$ -MnO<sub>2</sub> material can be attributed to the predicted lower work function for the rGO/ $\alpha$ -MnO<sub>2</sub>(100) nanocomposite.

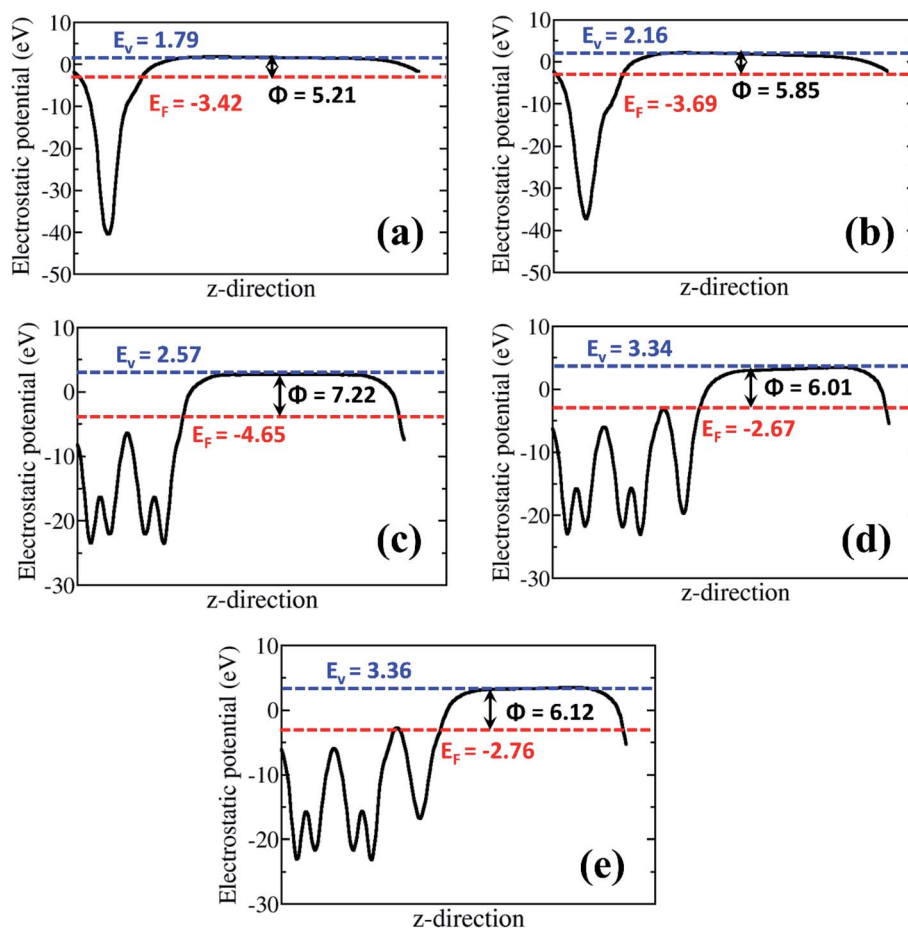


Fig. 6 The electrostatic potentials for isolated epoxide-rGO with (a) 8% and (b) 16% oxygen content; (c) isolated  $\alpha$ -MnO<sub>2</sub>(100) surface, and rGO/ $\alpha$ -MnO<sub>2</sub>(100) nanocomposite formed rGO with (d) 8% and (e) 16% oxygen content. The blue and red dashed lines represent the vacuum level ( $E_{vac}$ ) and the Fermi level ( $E_F$ ), respectively.  $\Phi$  denotes the work function.



Reduction in the work function has been observed in other composite materials compared to the isolated materials.<sup>52,53</sup> For instance, Susaki *et al.*, have shown that the deposition of a single unit cell of MgO on an Nb:SrTiO<sub>3</sub> substrate reduces the work function by about 0.8 eV.<sup>74</sup> Similarly, by decorating SnSe nanosheets with Au nanoparticles (Au/SnSe) and porous ZnO nanosheets with CuSCN nanocoins (CuSCN/ZnO) resulted in significant improvements in the FE characteristics owing to predicted lower functions.<sup>52,53</sup> The higher work function predicted for the  $\alpha$ -MnO<sub>2</sub>(100) surface compared to rGO monolayer (Fig. 6) suggests that spontaneous electrons transfer will flow from the rGO monolayer to the  $\alpha$ -MnO<sub>2</sub>(100) after the two are coupled together. Besides the reduction of the work function, the formation of nano-protrusions (denoted by red circles) in the rGO/MnO<sub>2</sub> nano-heterostructure (Fig. S5†) can act as effective emission sites. In addition, the better electrical conductivity ( $5 \times 10^{-3} \text{ S cm}^{-1}$ ) of rGO is expected to play an important role in electron transportation. The rGO as backbone may result in easy and high percolation of electrons from the rGO to MnO<sub>2</sub> nanorods giving rise to the observed superior field emission behavior of the rGO/MnO<sub>2</sub> nano-heterostructure.

## 4. Conclusion

In summary, we report the successful synthesis of MnO<sub>2</sub> nanorods and rGO/MnO<sub>2</sub> nano-heterostructure using cost effective hydrothermal and modified Hummer's methods, respectively. The coupling of rGO sheets with MnO<sub>2</sub> nanorods is demonstrated to have a synergistic effect in improving the FE characteristics of formed rGO/MnO<sub>2</sub> nano-heterostructure. The dramatic reduction of the turn-on field by  $5.7 \text{ V } \mu\text{m}^{-1}$  for an emission current density of  $10 \text{ } \mu\text{A cm}^{-2}$  and the achieved high current density of  $600 \text{ } \mu\text{A cm}^{-2}$  with an applied field of  $3.1 \text{ V } \mu\text{m}^{-1}$  demonstrate the superior FE characteristics of the rGO/MnO<sub>2</sub> nano-heterostructure compared to the isolated porous MnO<sub>2</sub> nanorods. The results are corroborated by first principles DFT calculations, which predict lower work function for the rGO/MnO<sub>2</sub> nano-heterostructure (6.01 and 6.12 eV for the rGO with 8% and 16% oxygen contents, respectively) compared to the isolated MnO<sub>2</sub> (7.22 eV) as the primary origin for the improved field emission of the rGO/MnO<sub>2</sub> nanocomposite. The controlled nanofabrication of rGO/MnO<sub>2</sub> heterostructure reported here provides a promising approach for designing highly efficient MnO<sub>2</sub>-based next generation FE electron sources and extend their practical applications in micro/nano electronic devices.

## Conflicts of interest

There are no conflicts to declare.

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## References

- 1 D. J. Milliron, S. M. Hughes, Y. Cui, L. Manna, J. Li, L. W. Wang and A. P. Alivisatos, *Nature*, 2004, **430**, 190–195.
- 2 M. D. Stoller, S. Park, Y. Zhu, J. An and R. S. Ruoff, *Nano Lett.*, 2008, **8**(10), 3498–3502.
- 3 Y. Sun, Q. Wu and G. Shi, *Energy Environ. Sci.*, 2011, **4**, 1113–1132.
- 4 G. A. M. Ali, S. A. Makhlof, M. M. Yusoff and K. F. Chong, *Rev. Adv. Mater. Sci.*, 2015, **41**, 35–43.
- 5 G. A. M. Ali, M. M. Yusoff and K. F. Chong, *ARPN J. Eng. Appl. Sci.*, 2016, **16**, 9712–9717.
- 6 G. Eda, G. Fanchini and M. Chhowalla, *Nat. Nanotechnol.*, 2008, **3**, 270–274.
- 7 J. T. Robinson, M. Zalalutdinov, J. W. Baldwin, E. S. Snow, Z. Wei, P. Sheehan and B. H. Houston, *Nano Lett.*, 2008, **8**(10), 3441–3445.
- 8 J. T. Robinson, F. K. Perkins, E. S. Snow, Z. Q. Wei and P. E. Sheehan, *Nano Lett.*, 2008, **8**, 3137–3140.
- 9 D. R. Dreyer, H. Jia and C. W. Bielawski, *Angew. Chem. Int. Ed.*, 2010, **49**, 6813–6816.
- 10 D. Wang, R. Kou, D. Choi, Z. Yang, Z. Nie, J. Li, L. V. Saraf, D. Hu, J. Zhang, G. L. Graff, J. Liu, M. A. Pope and I. A. Aksay, *ACS Nano*, 2010, **4**, 1587–1595.
- 11 W. Gao, N. Singh, L. L. Song, A. L. Reddy, L. Ci, R. Vajtai, Q. Zhang, B. Wei and P. M. Ajayan, *Nat. Nanotechnol.*, 2011, **6**, 496–500.
- 12 K. P. Loh, Q. Bao, G. Eda and M. Chhowalla, *Nat. Chem.*, 2010, **2**, 1015–1024.
- 13 H. Yamaguchi, K. Murakami, G. Eda, T. Fujita, P. Guan, W. Wang, C. Gong, J. Boisse, S. Miller, M. Acik, K. Cho, Y. Chabal, M. Chen, F. Wakaya, M. Takai and M. Chhowalla, *ACS Nano*, 2011, **5**(6), 4945–4952.
- 14 R. Kumar, R. Singh, A. Vaz, R. Yadav, C. Rout and S. Moshkalev, *New J. Chem.*, 2017, **41**, 8431–8436.
- 15 R. Roy, A. Jha, D. Sen, D. Banerjee and K. K. Chattopadhyay, *J. Mater. Chem. C*, 2014, **2**, 7608–7613.
- 16 C. Chen, W. Cai, M. Long, B. Zhou, Y. Wu, D. Wu and Y. Feng, *ACS Nano*, 2010, **4**(11), 6425–6432.
- 17 C. Zhu, S. Guo, P. Wang, L. Xing, X. Y. Fang, Y. M. Zhai and S. J. Dong, *Chem. Commun.*, 2010, **46**, 7148–7150.
- 18 Z. Rujia, Z. Zhang, L. Jiang, K. Xu, Q. Tian, S. Xue, J. Hu, Y. Bando and D. Golberg, *Mater. Chem.*, 2012, **22**, 19196–19201.
- 19 G. H. Gote, S. R. Bhopale, M. A. More and D. J. Late, *Phys. Status Solidi A*, 2019, **216**, 1900121.



- 20 C. S. Rout, P. D. Joshi, R. V. Kashid, D. S. Joag, M. A. More, A. J. Simbeck, M. Washington, S. K. Nayak and D. J. Late, *Sci. Rep.*, 2013, **3**, 3282.
- 21 M. Wu, J. Lee, Y. Wang and C. Wan, *J. Phys. Chem. B*, 2004, **108**(42), 16331–16333.
- 22 K. Hareesh, S. R. Suryawanshi, A. B. Phatangare, M. A. More, V. N. Bhoraskar and S. D. Dhole, *Mater. Lett.*, 2016, **185**, 472–475.
- 23 J. Wu, H. Huang, L. Yu and J. Hu, *Adv. Mater. Phys. Chem.*, 2013, **3**, 201–205.
- 24 W. Wei, X. Cui, W. Chen and D. G. Ivey, *Chem. Soc. Rev.*, 2011, **40**, 1697–1721.
- 25 T. Bordjiba and D. Belanger, *Electrochim. Acta*, 2010, **55**, 3428–3433.
- 26 K. Sun, H. Wang, H. Peng, Y. Wu, G. Ma and Z. Lei, *Int. J. Electrochem. Sci.*, 2015, **10**, 2000–2013.
- 27 S. R. Majid Rusi, *Solid State Ionics*, 2014, **262**, 220–225.
- 28 G. A. M. Alia, M. M. Yusoff, H. Algarnic and K. F. Chonga, *Ceram. Int.*, 2018, **44**, 7799–7807.
- 29 S. Ghasemi, S. R. Hosseini and O. Boore-talari, *Ultrason. Sonochem.*, 2018, **40**, 675–685.
- 30 R. Rajagopal and K. Ryu, *ChemElectroChem*, 2018, **5**, 2218.
- 31 S. Ghosal and P. Bhattacharyya, *IEEE Trans. Electron. Dev.*, 2019, **66**, 3882.
- 32 J. Hong, T. T. Mengesha, S. Hong, H. Kim and Y. Hwang, *J. Korean Phys. Soc.*, 2020, **76**, 264–272.
- 33 O. Sadak, W. Wang, J. Guan, A. K. Sundramoorthy and S. Gunasekaran, *ACS Appl. Nano Mater.*, 2019, **2**, 4386–4394.
- 34 J. G. Radich and P. V. Kamat, *ACS Catal.*, 2012, **2**, 807–816.
- 35 B. Q. Zeng, G. Y. Xiong, S. Chen, H. Jo, W. Z. Wang, D. Z. Wang and Z. F. Ren, *Appl. Phys. Lett.*, 2006, **88**, 213108.
- 36 D. C. Marcano, D. V. Kosynkin, J. M. Berlin, A. Sinitskii, Z. Sun, A. Slesarev, L. B. Alemany, W. Lu and J. M. Tour, *ACS Nano*, 2010, **4**(8), 4806–4814.
- 37 S. Santhanagopalan, A. Balram and D. D. Meng, *ACS Nano*, 2013, **7**(3), 2114–2125.
- 38 G. P. Patil, P. K. Baviskar, V. S. Bagal, R. D. Ladhe, A. B. Deore, M. A. More, B. R. Sankapal and P. G. Chavan, *RSC Adv.*, 2016, **6**, 71958.
- 39 G. Kresse and J. Hafner, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1993, **47**, 558.
- 40 G. Kresse and D. Joubert, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1999, **59**, 1758.
- 41 G. Kresse and G. J. Furthmüller, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1996, **54**, 11169–11189.
- 42 P. E. Blöchl, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1994, **50**, 17953–17979.
- 43 H. J. Monkhorst and J. D. Pack, *Phys. Rev. B: Solid State*, 1976, **13**, 5188–5192.
- 44 I. Biaggio, R. W. Hellwarth and J. P. Partanen, *Phys. Rev. Lett.*, 1997, **78**, 891.
- 45 V. Krukau, O. A. Vydrov, A. F. Izmaylov and G. E. Scuseria, *J. Chem. Phys.*, 2006, **125**, 224106-1–224106-5.
- 46 P. E. Blöchl, O. Jepsen and O. K. Andersen, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1994, **49**, 16223–16233.
- 47 D. A. Tompsett, S. C. Parker and M. S. Islam, *J. Mater. Chem. A*, 2014, **2**, 15509–15518.
- 48 A. Walsh and K. T. Butler, *Acc. Chem. Res.*, 2013, **47**, 364–372.
- 49 M. Ganose, K. T. Butler, A. Walsh and D. O. Scanlon, *J. Mater. Chem. A*, 2016, **4**, 2060–2068.
- 50 L. Wu, N. Y. Dzade, L. Gao, D. O. Scanlon, Z. Öztürk, N. Hollingsworth, B. M. Weckhuysen, E. J. M. Hensen, N. H. L. Leeuw and J. P. Hofmann, *J. Adv. Mater.*, 2017, **28**, 9602–9607.
- 51 N. Y. Dzade and N. H. Leeuw, *Phys. Chem. Chem. Phys.*, 2017, **19**, 27478–27488.
- 52 P. K. Baviskar, S. R. Rondiya, G. P. Patil, B. R. Sankapal, H. M. Pathan, P. G. Chavan and N. Y. Dzade, *ACS Omega*, 2020, **5**(12), 6715–6724.
- 53 S. R. Rondiya, C. D. Jadhav, P. G. Chavan and N. Y. Dzade, *Sci. Rep.*, 2020, **10**(1), 1–10.
- 54 B. He, G. Cheng, S. Zhao, X. Zeng, Y. Li, R. Yang, M. Sun and L. Yu, *J. Solid State Chem.*, 2019, **269**, 305–311.
- 55 P. K. S. Mural, M. Sharma, A. Shukla, S. Bhadra, B. Padmanabhan, G. Madras and S. Bose, *RSC Adv.*, 2015, **5**, 32441–32451.
- 56 H. Wang, Z. Lu, D. Qian, Y. Li and W. Zhang, *Nanotechnology*, 2007, **18**, 115616.
- 57 B. J. Tan, K. J. Klabunde and P. M. A. Sherwood, *J. Am. Chem. Soc.*, 1991, **113**, 855–861.
- 58 V. Subramanian, H. Zhu, R. Vajtai, P. M. Ajayan and B. Wei, *J. Phys. Chem. B*, 2005, **109**, 20207–20214.
- 59 B. Li, G. Rong, Y. Xie, L. Huang and C. Feng, *Inorg. Chem.*, 2006, **45**, 6404–6410.
- 60 H. Huang and X. Wang, *Nanoscale*, 2011, **3**, 3185–3191.
- 61 G. P. Patil, V. S. Bagal, C. R. Mahajan, V. R. Chaudhari, S. R. Suryawanshi, M. A. More and P. G. Chavan, *Vacuum*, 2016, **123**, 167–174.
- 62 C. Wu, T. W. Kim and T. Guo F. Li, *Sci. Rep.*, 2016, **6**, 38613.
- 63 S. R. Bansode, R. T. Khare, K. K. Jagtap, M. A. More and P. Koinkar, *Mater. Sci. Semicond. Process.*, 2017, **63**, 90–96.
- 64 W. T. Zheng, Y. M. Ho, H. W. Tian, M. Wen, J. L. Qi and Y. A. Li, *J. Phys. Chem. C*, 2009, **113**, 9164.
- 65 R. Zou, G. He, K. Xu, Q. Liu, Z. Zhang and J. Hu, *J. Mater. Chem. A*, 2013, **1**, 8445.
- 66 J. Ding, X. Yan, J. Li, B. Shen, J. Yang, J. Chen and Q. Xue, *ACS Appl. Mater. Interfaces*, 2011, **3**, 4299.
- 67 Y. Jin, Y. Zheng, S. G. Podkolzin and W. Lee, *J. Mater. Chem. C*, 2020, **8**, 4885–4894.
- 68 G. Kruthika, J. Karthikeyan and P. Murugan, *Phys. Chem. Chem. Phys.*, 2017, **19**, 3770–3776.
- 69 F. Cheng, T. Zhang, Y. Zhang, J. Du, X. Han and J. Chen, *Angew. Chem. Int. Ed.*, 2013, **52**, 2474–2477.
- 70 C. S. Johnson, D. W. Dees, M. F. Mansuetto, M. M. Thackeray, D. R. Vissers, D. Argyriou, C. K. Loong and L. Christensen, *J. Power Sources*, 1997, **68**, 570–577.
- 71 N. Sakai, Y. Ebina, K. Takada and T. Sasaki, *J. Phys. Chem. B*, 2005, **109**, 9651–9655.
- 72 M. J. Young, A. M. Holder, S. M. George and C. B. Musgrave, *Chem. Mater.*, 2015, **27**, 1172–1180.
- 73 P. V. Kumar, M. Bernardi and J. C. Grossman, *ACS Nano*, 2013, **7**, 1638–1645.
- 74 T. Susaki, A. Makishima and H. Hosono, *Phys. Rev.*, 2011, **83**, 115435.

