TABLES FOR SOLUTION OF THE HEAT-CONDUCTION EQUATION WITH A TIME-DEPENDENT HEATING RATE

ARTHUR E. BERGLES

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Arthur E. Bergles

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Heat Transfer Laboratory Engineering Projects Laboratory Department of Mechanical Engineering Massachusetts Institute of Technology Cambridge, Massachusetts

ABSTRACT

Tables are presented for the solution of the transient onedimensional heat flow in a solid body of constant material properties with the heating rate at one boundary dependent on time. These tables allow convenient and rapid estimation of the temperature distribution in the many practical cases where the mathematical model applies. Examples illustrating use of the tables are given.

INTRODUCTION

An important case of the heat-conduction equation concerns the largeslab geometry where one boundary is insulated and the heat input at the other boundary is an arbitrary function of time. This mathematical model approximates, for example, the aerodynamic heating of a flight vehicle re-entering the earth's atmosphere. Complex problems of this nature are normally handled by numerical integration and employ large-scale digital computers. However, the design engineer needs to determine rapidly the transient non-uniform temperature distribution in the skin of such a vehicle for any specified flight condition. In this respect the analytical solution offers considerable advantage over the numerical procedure as the entire heating history need not be considered to obtain the temperature distribution at any instant of time.

There has been considerable interest in this problem in recent years, and various solutions have been presented for the transient one-dimensional heat flow in a solid body of constant material properties with the heating rate at one boundary dependent on time.

This is a linear boundary-value problem which readily yields its solution to any of the elegant methods available for treating such problems. The solution for a constant flux into the solid is given in the standard work of Carslaw and Jaeger[1]^{*}. A small-scale chart gives this infinite series solution over a range of parameters. Carslaw and Jaeger also give the solution for a time-dependent flux which can be expressed

*Numbers in brackets denote References at end of paper.

as a single term. These single-term-flux solutions are, however, limited in their usefulness as most actual heating rates are more complex functions of time.

Sutton [2, 3] and Chen [4] presented the solution in integral form for an arbitrary heating rate. However, these integrals can be very difficult to evaluate depending on the form of the heating rate. This type of solution is only of academic interest to the engineer who seeks a convenient and rapid solution. Even the solutions given in reference [1] for certain time variations of the heat input are of limited practical usefulness as the computation labor involved in evaluating the infinite series is a major burden.

Solutions for a more general variation of the heat flux have been presented. Sutton [2] give a solution for a polynomial variation of the heat input with time. Certain terms in the solution were neglected so that the temperature distribution could be given as a polynomial in time and position. This reference includes an involved table of coefficients which can be used to calculate the desired temperatures. The solution is, however, accurate only for large times due to simplification of the solution.

A polynomial variation of the heat input was also considered by Bergles and Kaye [5]. The exact solution was given in terms of the infinite series of the repeated integrals of the error functions. The design charts given by Kaye and Yeh [6] can be used to rapidly estimate these infinite series. Such graphical solutions conserve on space; however, a limited range of parameters is considered, and their accuracy is limited to two significant figures.

-2-

The polynomial time variation of the heating rate is sufficiently general to be of considerable practical interest. In view of the fact that available graphical and tabular solutions are inaccurate and inconvenient to use, accurate tabular solutions were prepared based on the analysis of reference [5].

ANALYSIS

Consider transient one-dimensional heat flow in a solid body of constant material properties with the heating rate at one boundary dependent on time for the slab shown in Fig. 1. The following solution is presented in quite general terms by consideration of heat generation and initial temperature distribution.

The general differential equation, assuming a heat-generation term of the form

$$P(N) = K + M N^2$$
 (1)

is given by

$$\partial^2 \tau (\mathbf{N}, \mathbf{X}) / \partial \mathbf{N}^2 - \partial \tau (\mathbf{N}, \mathbf{X}) / \partial \mathbf{X} + \mathbf{K} + \mathbf{M} \mathbf{N}^2 = 0$$
 (2)

The initial temperature distribution is assumed to be represented by an even order polynomial of the form

$$\mathcal{T}(\mathbf{N},\mathbf{0}) = \mathbf{FN}^2 + \mathbf{GN}^4 \tag{3}$$

The boundary conditions are as follows:

$$\partial \tau (0, \mathbf{X}) / \partial \mathbf{N} = 0 \tag{4}$$

$$\partial \tau(\mathbf{1},\mathbf{x})/\partial \mathbf{N} = \mathbf{Q}(\mathbf{x})$$
 (5)

where the heat flux in Eq. (5) can be expressed as a polynomial of (s + 1) terms

$$Q(X) = \sum_{s=0}^{s} H_{s} X^{s}$$
(6)

Equations (2) to (6) were solved by means of the Laplace transformation. The solution for the temperature distribution is given by

$$\tau(N,X) = \sum_{B=0}^{B} X^{B} Z_{2S+1} - (2F + 4G) Z_{1} - 24G X Z_{3} + g_{2S+1} - (2F + K) Z_{1} - 24G Z_{1} - 24$$

where

$$Z_{2s+1} = 2^{2s+1} s! x^{1/2} \sum_{r=0}^{\infty} \left\{ i^{2s+1} erfc \quad \left[(2r+1-N)/2x^{1/2} \right] + i^{2s+1} erfc \quad \left[(2r+1+N)/2x^{1/2} \right] \right\}$$
(8)

It is seen that the solution to Eq. (2) can be simply represented as a polynomial in the functions of time and position defined by Eq. (8). The transient temperature distribution can be readily computed if the various temperature distribution functions are computed and arranged in tabular form.

COMPUTATION OF TEMPERATURE DISTRIBUTION FUNCTIONS

The infinite series of the repeated integrals of the error function, which comprises the temperature distribution function given by Eq. (8), was evaluated using the IEM 704 digital computer.

The repeated integrals of the error function were taken from the tables of Kaye [7]. The tables, together with appropriate differences, were stored in the machine. The Everett central-difference interpolation

method was chosen as it is more accurate and converges more rapidly than the forward-difference methods such as those of Newton or Gauss.

Values of the index s were chosen from 0 to 5, and the position ratio N varied from 0 to 1.0 in steps of 0.2. The values of relative time X were chosen over a range of 0 - 40 so that linear interpolation could be used in the tables. Linear interpolation is valid throughout all but two per cent of the final tables which are presented in the Appendix.

APPLICATIONS

The problem of the slab with prescribed heat flux at its surface is of increasing importance in technical applications. The general requirement for employment of the present model is that the heat flux be independent of the temperature of the body. Heat can be supplied, for example, by a flat heater embedded in the solid; in this case there is no loss of heat at the surface, and the boundary condition is accurately satisfied if the thermal capacity of the heater is negligible. The boundary condition is also satisfied for a flight vehicle re-entering the earth's atmosphere where the allowable surface temperature is small compared with the gas temperature and can be neglected.

A special case of the present solution was derived in reference [6] for use in estimating the transient temperature distribution in a wedgeshaped wing flying at supersonic speeds. The flux variation was obtained from specified time variation of the surface coefficient of heat transfer and of the temperature difference between the adiabatic wall and the wall. The assumption of a time variation of the temperature difference is,

-5-

however, equivalent to specifying the solution to the equation. The close agreement noted in that investigation between the analytical and numerical results is, therefore, to be expected since information from the numerical solution was used as a boundary condition for the analytical solution. An analytical solution for aerodynamic heating using a heat-transfer coefficient and adiabatic-wall temperature which are time dependent is given in reference [8]. The complexity of this solution is so great, however, that it is apparent that the numerical approach is more desirable for this type of aerodynamic heating problem.

The present solution is applicable to situations where the temperature gradient in all but one direction can be neglected. The temperature distribution for certain simple two- and three-dimensional geometries, such as the brick-shaped solid, can be treated using the tables and the standard product solution technique of Newman.

The tables have sufficiently fine intervals at low values of the relative time so that the very thick slabs or semi-infinite solids can be readily considered. Only the large values of the position ratio are used in this case.

The inverse problem arises when the surface heat flux versus time is sought from knowledge of an interior temperature versus time. Stolz [9] presented one of the few general treatments of the subject and developed a numerical inversion method. The present analytical method is applicable to all cases where the heat flux can be expressed as a polynomial and where the initial temperature is uniform.

The interior temperature at a location near the surface is monotored as a function of time. The values of N and X are readily calculated if

-6-

the properties are known and the corresponding Z are obtained from the tables. The coefficients of the surface heat flux are then solved from the series of simultaneous equations derived from the data.

$$H_{0} Z_{1} + H_{1} Z_{3} X + \dots = t - t_{b}$$

$$H_{0} Z_{1} + H_{1} Z_{3} X' + \dots = t' - t_{b}$$
etc.

EXAMPLES

Satellite Re-entry

The temperature distribution in the skin of a flight vehicle can be readily estimated by the present method for an important case of aerodynamic heating. The solution given by Eq. (7) is valid for a situation where the heat input to a body can be determined independent of the surface temperature. This occurs, for example, when a vehicle re-enters the earth's atmosphere as its allowable surface temperature is small compared with the gas temperature and can be neglected [10]. For hypersonic re-entry velocities the heat input is determined primarily by the solution to the equation of motion.

Scala [11] presents the aerodynamic heating rate for the ballistic re-entry of a satellite from an initial orbit of 900,000 feet. The ballistic parameter W/C_DA_D is chosen to be 200 lb/ft², and the re-entry velocity is approximately 24,000 ft/sec at a path angle of 92.5 degrees from the local vertical. The heating rate obtained in that study is presented in Fig. 2. Figure 3 includes one-dimensional transient heatconduction calculations used in the heat sink section of Scala's work. The temperature profiles were obtained by numerical analysis involving use of a large digital computer. The temperature distribution will now be approximated by means of the present method.

The heating rate to the satellite as given in Fig. 2 is first expressed as a polynomial in time. A simple four-point curve fit suffices to give a good representation of the temporal variation of the heating rate.

 $q/A = 1.76 \times 10^{-4} e^3 - 9.887 \times 10^{-3} e^2 + 0.5711e$ (9) The thickness of the material is five inches. The product of ρc_p is 44 Btu/ft³ °F, and the thermal conductivity is assumed to be 43.2 Btu/hr ft °F[12]. Using these properties the heating rate is expressed in the form of Eq. (6).

 $Q = 15.77 \times 10^5 X^3 - 13.91 \times 10^4 X^2 + 12.63 \times 10^3 X$ (10) The solution as given by Eq. (7) for a uniform initial temperature of

t - 170 =
$$15.77 \times 10^5 X^3 Z_7 - 13.91 \times 10^4 X^2 Z_5 + 12.63 \times 10^3 X Z_3$$
 (11)

Temperature profiles were calculated for the same flight times as used in the computer solution. For a time of 100 sec, for example, the value of $X = \alpha \theta/1^2 = 0.1571$. At N = 1.0 the value of Z_7 using linear interpolation is 0.2044. The temperature profiles as calculated by the present method are shown to be in close agreement with the computer solution in Fig. 3. The small deviations in the profiles are due to consideration of gas-cap radiation and use of a temperature-dependent thermal conductivity in the computer solution. Approximately two hours of desk-calculator computations were necessary to obtain the four temperature profiles. The rest of the temperature history can be calculated by fitting a curve to the remaining portion of the heat flux versus time plot and proceeding as above. The initial temperature distribution must, however, be included. An even-order-polynomial curve can be fitted to the temperature distribution at the time of application of the new heat input. Several such steps should serve in most cases to solve the problem if the entire heat-input history cannot be expressed accurately as a single polynomial.

Evaluation of Diffusivity

The present solution can be used to rapidly evaluate the thermal diffusivity of a material. Diffusivity experiments achieve one-dimensional heat flow by a. use of a large slab sample where only a small central portion is used for measurements, b. insulation of the sides so as to prevent heat losses, and c. use of a guard heater to prevent radial heat loss.

An experiment to determine the diffusivity of copper at 932 ${}^{\circ}F$ (500 ${}^{\circ}C$) is devised using a large-slab sample. A heat source, such as an electric resistance element, is uniformly applied to the surface; whereas, the back face is insulated. Thermocouples are installed at two interior locations, x = 1.0 and 2.0 inches. The block is initially at a uniform temperature $t_b = 72 {}^{\circ}F$. The heating, with a heat source producing a heat flux $q/A \sim 10^6$ Btu/hr ft², is commenced at time $\Theta = 0$, and the temperatures at the interior locations are monotored. The time required for each of these points to reach 932 ${}^{\circ}F$ is then recorded.

-9-

x = 2.0 in., N' = 0.8, Θ^{i} = 31.3 sec. x = 1.0 in., N" = 0.4, Θ^{n} = 41.8 sec.

The analytical solution as obtained from Eq. (7) is given as

$$\mathcal{T}(N, X) = t - t_b = H_0 Z_1(N, X) = (q/A) (1/k) Z_1$$
 (12)

It is readily seen that the temperature distribution functions must be equal for the two points. Thus

$$Z_{1}^{\prime}(N^{\prime}, X^{\prime}) = Z_{1}^{\prime\prime}(N^{\prime\prime}, X^{\prime\prime})$$
 (13)

subject to the condition that

$$X'/X'' = \Theta'/\Theta'' = 31.3/41.8$$
 (14)

The tables for Z_1 , N = 0.4, and 0.8 give the desired solution

$$Z_1^{i} = Z_1^{i} = 0.8684$$

 $X^{i} = 0.7152, \quad X^{i} = 0.9551$

from which the diffusivity is calculated.

$$\alpha = 1^2 \text{ X/}\Theta = (2.5/12)^2 0.7152/(31.3/3600) = 3.57 \text{ ft}^2/\text{hr}.$$

It is noted that the evaluation of the diffusivity requires no measurement of the heat flux. If the density and specific heat are known from other simple experiments, this experimental determination also yields the thermal conductivity indirectly. The thermal conductivity can, however, be directly evaluated from Eq. (12) if the heat flux is measured.

The analytical solution to Eq. (7) for constant heat flux has been modified and used with success in diffusivity experiments by Butler and Inn [13] and Sheer, et al. [14] among others. These investigators made use of the linearity of the temperature distribution at values of the relative time X > 0.5. The present method is more general as the data can be taken for any value of the relative time.

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SYMBOLS

A	=	surface area
k	=	thermal conductivity
1	=	thickness of slab
N	=	position ratio = x/l
P	=	heat-generation term = q^1/k
Q,	=	heat-flux term = $(q/A)(1/k)$
đ	=	rate of heat transfer
g *	=	heat generation per unit volume
r,s	=	integers, 0, 1, 2,
t	=	temperature
t _o	=	constant initial temperature
x	=	relative time = $\alpha 0/1^2$
x	=	normal distance
Z	=	defined by Eq. (8)
F, G, H, K, M	8	constants
α	=	thermal diffusivity
au	=	temperature difference = $(t - t_b)$
•	1	time
Subscripts and Supers	cri	pts
w	=	wall position
S	=	integer



FIG. I ONE - DIMENSIONAL HEAT CONDUCTION IN A SLAB



THE TEMPERATURE DISTRIBUTION FUNCTION - Z1

х	$\mathbf{N} = \mathbf{O}$	N = .2	$\mathbf{N} = .4$	N = .6	N = .8	N = 1.0
0.000	0.00000	0.00000	0.00000	0.0000	0.0000	0.0000
.002	.00000	.00000	.00000	.0000	.0000	.0505
.004	.00000	.00000	.00000	•0000 :	.0008	.0714
.006	.00000	.00000	.00000	.0000	.0029	.0874
.008	.00000	.00000	.00000	.0001	.0061	.1109
.010	.00000	.00000	.00000	.0002	.0101	.1128
.012	.00000	.00000	.00000	.0005	.0144	.1236
.014	.00000	.00000	.00001	.0009	.0190	.1335
.016	.00000	.00000	.00004	.0016	.0237	.1427
.018	.00000	.00000	.00008	.0024	.0285	.1514
.020	.00000	.00000	.00015	.0034	.0333	.1596
.022	.00000	.00001	.00026	.0046	.0382	.1674
.024	.00000	.00001	.00041	.0059	.0430	.1748
.026	.00000	.00003	.00060	.0073	.0478	.1820
.028	.00000	.00004	.00085	•008 9	.0525	.1888
.030	.00000	.00007	.00115	.0105	.0572	.1954
.032	.00001	.00011	.00150	.0123	.0618	.2019
.034	.00002	.00016	.00190	.0141	.0664	.2081
.036	.00002	.00022	.00237	.0161	.0710	.2141
.038	.00004	.00030	.00288	.0181	.0754	.2200
.040	.00006	.00039	.00345	.0201	.0799	.2256
.042	.00008	.00051	.00407	.0222	.0842	.231 2
.044	.00011	.00064	.00474	.0243	.0885	.2367
.046	.00016	.00079	.00546	.0265	.0928	.2420
.048	.00021	.00097	.00623	.0288	.0970	.247 2
.050	.0003	.0012	.0070	.0310	.1012	.2523
.055	.0005	.0018	.0093	•0368	.1113	.2646
.060	.0008	.0025	.0117	.0426	.1212	.2764
.065	.0012	.0034	.0144	.0486	.1308	.2877
.070	.0017	.0045	.0173	.0546	.1402	•2985
.075	.0024	.0057	.0203	.0606	.1493	.3090
.080	.0032	.0071	.0235	.0667	.1582	•3192
.085	.0042	.0086	.0268	.0727	.1669	.3290
.090	.0052	.0103	•0303	.0788	.1754	•3385
.095	•0065	.0122	.0338	.0848	.1838	• 34 (8
.100	.0079	.0141	•0375	.0908	.1919	.3568
.105	.0094	.0162	.0412	.0968	.1999	•3656
.110	.0111	.0185	.0450	.1027	.2078	•3(42
.115	.0129	.0209	.0489	.1086	.2155	.3027
.120	.0149	.0234	.0529	.1145	.2230	•3909
.125	.0170	.0260	.0569	.1204	.2305	•3990
.130	.0192	.0287	.0609	.1262	.2378	.4069
.135	.0216	.0315	.0651	.1320	.2450	.4146
.140	.0240	.0344	.0692	.1378	.2521	.4222

-16-

x	$\mathbf{N} = \mathbf{O}$	N = .2	N = .4	N = .6	N = .8	N = 1.0
.145	.0266	.0375	.0734	.1435	.2591	.4297
.150	.0293	.0400 ·	.0///	•1492	.2000	.43(1
•155	.0321	.0438	.0820	.1549	.2(28	.4443
.100	.0350	.0471	.0003	.1005	-2(95	•4515 hc9c
.107	.0300	.0707	.0907	.1001	.2001	•4202 ·
175	00411	-0739	-0951	•1(1)	•2921	•4074
180	.0443	•07(7	-0995	1808	•2992	•4/23
185	.0410	-0011 06h7	1090	1882	• 30 70	•4(90 h857
100	.0509	0685	1120	1005	.3119	.4071
105	·V)44	.0005	·1130	1000	· 3102	·4922
•195	•0719	.0122	•112	•1772	• 3244	•490(
.200	.0615	.0761	.1220	.2046	.3306	.5052
.210	.0688	.0840	.1312	.2155	.3427	-5178
.220	.0764	.0920	-1405	.2262	•3546	.5302
•230	.0843	.1003	.1498	-2369	.3664	•5424
.240	.0923	.1087	.1592	•24(5	.3780	•5544
.250	.1005	.1172	.1000	.2580	.3894	.5001
.200	.1009	-1259 John	-1(OL	•2005	.4007	•5770
.210	•11(4	•134(.10((.2(90	.4119	•5092
.200	.1201	•143(.19/3	•2094	-4230	.0000
.290	•1349	.1521	.2009	•2990	.4340	0110
. 300	•1430	.1010	•2100	• 3101	.4440 hee6	.0220
• 310	1610	1803	.2203	• 3204	•4770 b66b	•0330 6hhm
• 320	.1019	1806	•2300 ohrz	- 3301	·4004	•044 (6555
• 330	1801	1001	•2471	• 3409	.4/10	•0777
250	1807	2085	•2777	- 3511	1082	6760
.360	.1001	.2180	.2055	.3715	5086	.6875
.370	.2086	.2276	-2850	.3817	.5101	.6981
.380	.2181	.2372	2018	.3010	.5205	7086
.390	.2276	2468	3047	1050	.5398	.7190
	10	12100			•/3/0	•1200
.400	.2372	. 2565	-3145	.4121	•5502	.7294
.420	.2566	.2770	•3343	•4323	.5707	.7501
.440	.2760	• 29 55	•3542	•4525	.5912	.7707
.460	•295 5	.3151	.3740	.4727	.6116	.7911
.480	.3151	•3348	•3939	•4928	.6319	.8115
.500	•3348	• 3545	.4138	•5129	.6522	.8319
.520	• 3545	•3743	•4337	•5330	.6724	.8521
.540	•3743	.3941	.4536	•5530	.6925	.8723
.560	.3941	.4140	.4736	.5731	.7127	.8925
.580	.4140	.4339	.4935	.5931	.7328	.9127

X	$\mathbf{N} = \mathbf{O}$	N = .2	N = .4	N = .6	N = .8	N = 1.0
.600	.4339	.4538	.5135	.6132	.7529	.9328
.650	.4837	.5036	.5634	.6632	.8031	.9830
.700	.5335	.5535	.6134	.7133	.8532	1.0331
.750	.5835	.6034	.6634	.7633	.9032	1.0832
.800	.6334	.6534	.7134	.8133	.9533	1.1333
.850	.6834	.7034	.7633	.8633	1.0033	1.1833
.900	.7334	.7534	.8133	.9133	1.0533	1.2333
.950	.7834	.8033	.8633	.9633	1.1033	1.2833
1.0 3.4 5.6 7.8 9.0 1.2 3.4 5.6 7.8 9.0 1.2 3.4 5.6 7.8 9.0 1.2 3.4 5.6 7.8 9.0 1.2 1.2 1.2 1.2 1.2 1.2 1.2 1.2	$\begin{array}{c} .105^{+}\\ .833\\ 1.833\\ 2.833\\ 3.833\\ 3.833\\ 5.833\\ 5.833\\ 5.833\\ 5.833\\ 5.833\\ 3.833\\ 11.8333\\ 12.8333\\ 11.8333\\ 12.8333\\ 11.8333\\ 13.8333\\ 14.8333\\ 17.8333\\ 23.8333\\ 23.8333\\ 23.8333\\ 24.8333\\ 23.8333\\ 24.8333\\ 23.8333\\ 24.8333\\ 23.8333\\ 25.8333\\ 33.8333\\ 35.8333\\ $	$\begin{array}{c}$	$\begin{array}{c} .913\\ 1.913\\ 2.913\\ 3.9$	1.013 2.013 3.013 4.013 5.013 6.013 9.013 10.013 11.013 12.013 13.013 14.013 15.013 14.013 15.013 14.013 15.013 19.013 20.013 21.013 22.013 24.013 25.013 24.013 25.013 24.013 25.013 24.013 25.013 29.013 31.013 31.013 32.013 33.013 34.013 35.013 36.013 37.013 38.013	1.153 2.153 3.153 4.153 5.153 4.153 5.153 9.153 9.153 10.153 12.153 13.153 14.153 15.153 14.153 15.153 14.153 15.153 16.153 20.153 21.153 22.153 24.153 24.153 24.153 24.153 24.153 25.153 24.153 25.153 24.153 25.153 26.153 26.153 26.153 26.153 26.153 26.153 27.153 26.153 26.153 26.153 26.153 27.153 26.153 27.153 26.153 27.153 26.153 27.153 26.153 27.153 26.153 27.153 26.153 27.153	$\begin{array}{c} 1.333\\ 2.333\\ 3.333\\ 4.333\\ 5.333\\ 4.333\\ 5.333\\ 5.333\\ 9.333\\ 10.333\\ 12.333\\ 12.333\\ 14.333\\ 14.333\\ 15.333\\ 14.333\\ 15.333\\ 19.333\\ 21.333\\ 21.333\\ 22.333\\ 24.333\\ 25.333\\ 24.333\\ 29.333\\ 29.333\\ 31.333\\ 32.333\\ 31.333\\ 32.333\\ 33.333\\ $
39.0	38.833	38.853	38.913	39.013	39.153	39•333
40.0	39.833	39.853	39.913	40.013	40.153	40•333

.

THE TEMPERATURE DISTRIBUTION FUNCTION - Z1

THE TEMPERATURE DISTRIBUTION FUNCTION - Z₃

X	$\mathbf{N} = \mathbf{O}$	N = .2	N = .4	N = .6	N = .8	N = 1.0
0.000	0.00000	0.00000	0.00000	0.00000	0.0000	0.0000
.002	•00000	.00000	•00000	.00000	.0000	.0336
.004	.00000	.00000	.00000	.00000	.0002	.0476
.006	.00000	.00000	.00000	•00000	.0007	.0583
.008	.00000	•00000	•00000	.00001	.0016	.0673
.010	.00000	+ 00000 +	.00000	.00003	.0029	.0752
.012	.00000	.00000	.00000	•00008	.0045	.0824
.014	.00000	.00000	.00000	.00017	.0062	.0890
.016	.00000	.00000	.00000	•00030	.0081	.0952
.018	•00000	.00000	.00001	.00049	.0101	.1009
.020	.00000	.00000	.00002	.00073	.0122	.1064
.022	.00000	.00000	.00004	.00102	.0143	.1116
.024	.00000	.00000	.00006	.00137	.0165	.1165
.026	•00000	.00000	.00010	.00177	.0187	.1213
.028	.00000	.00001	.00014	.00222	.0210	.1259
.030	.00000	.00001	.00020	.00272	.0232	.1303
.032	.00000	.00001	.00027	.00326	.0255	.1346
•034	.00000	.00002	.00035	.00384	.0278	.1387
.036	•00000	.00003	.00045	.00447	.0300	.1427
•038	.00000	.00004	.00056	.00513	.0323	.1466
.040	.00001	.00006	.00069	.00583	.0346	.1505
.042	.00001	•00008	.00084	. 00656	.0368	.1542
.044	.00001	.00010	.001.00	.00732	.0391	.1578
.046	.00002	.00012	.00118	.00811	.0413	.1613
•048	.00003	.00016	.00137	.00892	.0436	.1648
.050	.00003	.00019	.0016	.0098	.0458	.1682
.055	•00006	.00031	.0022	.0120	.0513	.1764
.060	.00011	.00046	.0029	.0143	.0567	.1843
•065	.00018	.00065	.0037	.0167	.0620	.1918
.070	.00027	.00089	.0045	.0192	.0673	.1990
.075	.00039	.00117	.0055	.0217	.0725	.2060
•080	.00054	.00149	.0065	.0243	.0775	.2128
.085	.00072	.00187	.0076	.0270	.0825	.2193
.090	.00094	.00229	.0088	.0297	.0875	•2257
-095	.00120	.00276	.0100	.0325	.0923	•2319
.100	.00150	.00320	.0113	.0352	.09(1	·2319
.105	.00104	.00304	.0120	.0300	.1010	·2430
.110	.00222	.00440	.0140	.0400	.1004	•2495
100	.00207	.00512	.0154	.0430	.1110	·2771
105	00362	•00703	•070A	.0407	1000	2000
120	·W303	.00070	•0104	• • • • • • • • • • • • • • • • • • • •	2 OP	.2000
125	•004TO	.00130	•0200 0014	·V722	• 1244 1027	0761
•132	+004(0	.00022	.0210	·U77V	10201	0975
• 140	•00743	*00AT0	·0252	.0519	•1330	·2012

-19-

x	$\mathbf{N} = \mathbf{O}$	N = .2	N = .4	N = .6	N = .8	N = 1.0
.145	.00611	.01003	.0249	.0607	.1372	.2865
.150	.00684	.01099	.0265	.0636	.1414	.2914
.155	.00761	.01200	.0283	.0664	.1455	.2962
.160	.00842	.01305	.0300	.0693	.1496	.3009
.165	.00927	.01413	.0318	.0721	.1536	.3056
.170	.01016	.01525	.0336	.0750	.1576	.3102
.175	.01109	.01640	.0354	.0778	.1616	.3147
.180	.01206	.01759	.0372	.0807	.1655	.3192
.185	.01306	.01882	.0391	.0835	.1694	.3236
.190	.01411	.02008	.0410	.0863	.1732	•3219
•195	.01518	.02137	.0429	.0891	.1770	.3322
.200	.0163	.0227	.0448	.0920	.1808	• 3365
.210	.0186	.0254	.0487	.0976	.1882	.3448
.220	.0211	.0283	.0527	.1032	.1955	• 3530
.230	.0236	.0312	.0567	.1088	.2027	.3609
.240	.0263	.0343	•0608	.1143	.2097	• 366'(
•250	.0291	.0374	.0649	.1199	.2167	•3(64
.260	.0321	.0406	.0691	.1254	•2235	.3039
.270	.0351	.0440	.0733	.1309	.2303	• 3913
.280	.0382	.0474	.0(75	•1303	.2310	• 3900 hose
.290	.0413	.0500	.0010	•1410	•2430	.4070
.300	.0440	.0744	.0002	•14/2	.2701	.4120
- 310	.04(9	.0500	•0905 oolio	•1521	·2500	1266
. 320	•0714 05h0	.0011	•0949	1601	•2030	h22h
• <u>3 30</u>	.0549 058h	.0074	1028	1688	2756	.+ <u>.</u>
250	.0504	.0092	1082	1749	.2818	.hh57
- 360	.0657	.0769	.1128	.1795	.2879	4532
.370	.0604	.0809	.1173	1848	.2940	4597
.380	.0732	.0849	.1218	.1902	.3001	.4661
.390	.0771	.0889	.1264	.1955	.3061	.4725
.400	.0810	.0930	.1310	.2007	.3121	.4788
.420	.0889	.1012	.1402	.2113	.3239	.4912
.440	.0970	.1096	.1495	. 2218	.3356	.5034
.460	.1052	.1181	.1588	.2323	.3471	.5155
.480	.1135	.1267	.1682	.2427	.3586	.5274
.500	.1219	.1354	.1776	.2531	• 3699	•5398
.520	.1305	.1443	.1871	·2635	.3812	.5508
.540	.1392	. 1532	.1966	.2738	.3923	.5623
.560	.1479	.1622	.2061	.2842	.4034	•5738
.580	.1567	.1712	.2157	.2945	.4144	.5851

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THE TEMPERATURE DISTRIBUTION FUNCTION - Z₃

х	N = O	N = .2	N = .4	N = .6	N = .8	N = 1.0
.600	.1656	.1802	.2253	.3048	.4254	.5964
.650	.1882	.2032	.2494	.3304	.4525	.6242
.700	.2111	.2264	.2736	.3560	.4793	.6516
.750	.2342	.2499	.2979	.3815	.5059	.6787
.800	.2576	.2735	.3223	.4069	.5323	.7056
.850	.2812	.2974	• 3468	.4323	.5585	.7322
.900	.3049	.3213	.3713	.4576	.5846	.7586
.950	. 3288	• 3454	• 3959	.4829	.6106	.7849
1.0	• 353	.370	.421	.508	.637	.811
2.0	.843	.861	•917	1.011	1.145	1.322
3.0	1.340	1.359	1.416	1.512	1.648	1.826
4.0	1.838	1.857	1.915	2.012	2.149	2.520
5.0	2.337	2.357	2.415	2.512	2.050	2.029
6.0	2.837	2.856	2.915	3.012	3.151	3.334
7.0	3.336	3.356	3.414	3.513	3.051	3.030
8.0	3.836	3.855	3.914	4.013	4.171	4.331 h 831
9.0	4.335	4.300	4.414 h 01h	4.713 5 013	4.071 5 150	4.0JL 5 221
10.0	4.035	4.077	4.914 5 bab	5.013 5.512	5.650	5 821
11.0	2.332	2•322 5 855	2.414 5 01/1	5.713	6 152	6 331
12.0	5.035	5.055	5.914 6 h1h	6 513	6.652	6.832
12.0	6 835	6 854	6.014	7.013	7,152	7,332
15.0	7.335	7.354	7.414	7,513	7.652	7.832
16.0	7.835	7.854	7.914	8.013	8.152	8.332
17.0	8,334	8.354	8.414	8.513	8.652	8.832
18.0	8.834	8.854	8.914	9.013	9.152	9.332
19.0	9.334	9.354	9.414	9.513	9.652	9.832
20.0	9.834	9.854	9.914	10.013	10.152	10.332
21.0	10.334	10.354	10.414	10.513	10.652	10.832
22.0	10.834	10.854	10.914	11.013	11.153	11.332
23.0	11.334	11.354	11.414	11.513	11.653	11.832
24.0	11.834	11.854	11.914	12.013	12.153	12.332
25.0	12.334	12.354	12.414	12.513	12.653	12.832
26.0	12.834	12.854	12.914	13.013	13.153	13.332
27.0	13.334	13.354	13.414	13.513	13.073	13.032
28.0	13.034	13.054	13.914	14.013	14.173	14.556
29.0	14.334	14.304		14.713	14.073	15 222
30.0	14.034	14.074	14.914 15 hih	15 513	15 652	15 822
31.0	17•334	15 954	15 01	16 012	16 152	16 222
32.0	15.034	16 251	15.914	16 512	16 652	16 833
33.0 ah 0	16 824	10.394 16.85h	16 01	17 013	17,153	17,333
25 0	17 224	17 25h	10.914 17 h1h	17.513	17.653	17.833
36.0	17.82年	17.854	17.014	18.013	18.153	18.333
37.0	18.337	18.354	18.414	18.513	18.653	18.833
38.0	18.82	18.854	18.014	19.013	19.153	19.333
39.0	19, 334	19,354	19.414	19.513	19.653	19.833
40.0	19.834	19.854	19.914	20.013	20.153	20.333
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THE TEMPERATURE DISTRIBTION FUNCTION - Z5

X	$\mathbf{N} = \mathbf{O}$	N = .2	N = .4	N = .6	N = .8	N = 1.0
.145	.00213	.00399	.0119	03hh	.0023	2000
.150	.00241	.00442	.0128	.0362	.0054	•2292
.155	.00272	.00487	.0137	.0380	.008/	•2351
.160	.00304	.00534	.0147	.0399	1014	·2309
.165	.00339	.00583	.0156	.0417	1014	·2401
.170	.003 76	.00634	.0166	.0436	1074	•2447
.175	.00414	•0068 8	.0176	0454	.1102	•2401
.180	.00455	.00743	.0186	.0473	.1122	•2710
.185	.00498	.00801	.0197	.0491	.1161	•2555
.190	.00543	.00860	.0207	.0510	.1100	•2709
.195	.00 589	.00921	.0218	.0528	.1218	.2658
.200	.0064	.0098	.0229	.0547	.1247	.2601
.210	.0074	.0112	.0251	.0584	.1302	2758
.220	•008 5	.0126	.0274	.0622	.1357	.2823
.230	.0097	.0140	.0297	.0659	.1411	2887
.240	.0109	.01 55	.0321	·0696	.1464	2010
.250	.0123	.0171	.0345	.0733	.1516	.3010
.260	.0136	.0188	.0369	.0771	1568	.3060
.2(0	.0151	.0205	.0394	.0808	.1619	.3128
.280	.0166	.0223	.0420	.0845	.1670	3186
.290	.0182	.0241	.0445	.0882	.1719	3242
• 300	.0198	.0260	.0471	.0919	.1768	3298
• 310	.0215	.0279	.0497	•0956	.1817	.3353
• 320	.0232	•0298	.0524	.0992	.1865	.3407
· 330	.0250	.0319	. 0550	.1029	.1913	.3460
. 340	•0268	•0339	.0577	.1066	.1960	.3513
• 370	.0201	.0360	.0605	.1102	.2006	3565
· 300	.0306	.0381	.0632	.1139	.2052	.3616
· 3(V	.0326	.0403	.0660	.1175	•0298	.3667
• 300	.0346	.0425	•0688	.1212	.2143	.3717
• 390	.0301	•0448	.0715	.1248	.2188	.3766
-400	.0388	.0471	.0744	.1284	.2233	. 3815
-420	.0431	.0517	.0801	.1356	.2321	.3911
160	.0475	•0565	•0858	.1428	.2408	.4005
·400	.0521	.0614	.0 936	.1500	.2493	.4098
•400 500	.0567	.0628	.0975	.1571	.2578	.4189
• 500	.0015	.0714	.1034	.1643	.2661	4279
510	.0004	•0766	.1094	.1713	.2744	.4367
560	.0(14	.0819	.1154	.1784	.2826	.4455
580	.0(04	.0872	.1214	.1855	.2907	.4541
•)00	*00T0	.0926	.1275	.1925	.2987	.4626

х	N = O	N = .2	N = .4	N = .6	N = .8	N = 1.0
.600	.0868	.0980	.1336	.1996	.3066	.4710
.650	.1001	.1119	.1489	.2170	.3262	.4916
.700	.1139	.1261	.1645	.2344	.3455	.5118
.750	.1279	.1405	.1801	.2518	.3645	.5316
.800	.1422	.1552	.1959	.2690	.3832	.5511
.850	.1567	.1701	.2117	.2862	.4018	.5702
.900	.1715	.1852	.2276	.3034	.4201	.5892
.950	.1864	.2004	.2436	.3205	.4383	.6079
$\begin{array}{c} .950\\ 1.0\\ 2.0\\ 3.0\\ 4.0\\ 5.0\\ 6.0\\ 7.0\\ 8.0\\ 9.0\\ 10.0\\ 11.0\\ 12.0\\ 13.0\\ 14.0\\ 15.0\\ 14.0\\ 15.0\\ 14.0\\ 15.0\\ 14.0\\ 15.0\\ 14.0\\ 15.0\\ 22.0\\ 23.0\\ 24.0\\ 25.0\\ 24.0\\ 25.0\\ 26.0\\ 27.0\\ 28.0\\ 29.0\\ 31.0\\ 32.0\\ 31.0\\ 35.0\\ 37.0\\ 37.0\end{array}$.1864 $.201$ $.518$ $.846$ 1.176 1.508 1.840 2.172 2.505 2.838 3.171 3.503 3.837 4.170 4.503 3.837 4.170 4.503 5.169 5.502 5.835 6.502 5.835 6.502 5.835 7.168 8.835 9.168 8.835 9.168 8.835 9.168 10.835 10.168 10.835 11.168 11.501 11.835 12.168	.2004 .216 .535 .864 1.195 1.526 1.859 2.191 2.524 2.857 3.523 3.428 4.522 4.855 5.189 5.522 5.855 6.8555 6.8555 8.884 7.521 7.8555 8.884 9.521 10.8544 10.8544 12.521 12.854 12.884	.2436 .260 .587 .918 1.250 1.583 1.916 2.249 2.582 2.915 3.581 4.914 4.914 5.5814 4.914 5.5814 5.5814 5.9147 7.5814 8.5804 9.247 9.5804 10.5804 11.580 11.247 11.580 12.247	.3205 .338 .675 1.010 1.344 1.678 2.345 2.679 3.345 2.679 3.346 3.679 4.3479 4.679 5.013 5.679 6.013 6.679 7.013 6.680 3.346 9.013 8.680 9.013 10.680 11.013 11.036 12.013 12.013 12.013 12.013	.4383 .456 .804 1.142 1.478 1.813 2.482 2.8160 3.483 1.513 2.8160 3.483 1.513 4.4815 5.8181 5.4858 6.8182 7.485 8.8152 9.4859 9.4859 10.15259 12.485 12.152 12.485	.6079 .626 .9719 1.319 1.6561 2.22233660 3.6604 3.66933 3.66933637 5.6697314 6.993158 8.993258 9.993258 10.9982598 10.999825998 10.9999825998 10.99998599859985999999999999999999999999
38.0	12.501	12.521	12.580	12.679	12.819	12.998
39.0	12.835	12.854	12.914	13.013	13.152	13.332
40.0	13.168	13.188	13.247	13.346	13.485	13.665

0567.	6600.	£0020.	10900.	0/.00.	87000.	0#T*
6601	+F00.	04670.	<u><u> </u></u>	05100	79000.	SET.
3081	6000.	#2010.	<0<00.	26100.	1.4000*	•T30
+201.	+0(0+	TTITO	76400.	ΥΤΤΟΟ'	84000*	५टा∙
1081	0220	TACTO	ZT+00'	66000.	0+000*	.120
2821	5250.	Z/#T0*	002001	49000*	*00033	STT.
042L TT1T0	1000	024C0	003200.	2,000.	12000*	OTT*
TIOT	2030	StrZTO*	10200	09000*	12000.	ζ ΟΤ•
TEOT.	+2+0.	021010	05200.	05000*	1000.	00T*
06CT.	12+0.	ZZOTO*	ST200.	T#000*	£1000.	<u> 560°</u>
1462.	T0+0*	97600	ESTOO.	SE000.	00000	060*
#05T*	+120.	#T800*	#ST00.	\$2000°	70000.	≤80 •
654T.	1.420.	971.00	LETOO.	61000.	60000.	080.
ET#T.	6120.	.00622	EOTOO.	†T000 *	£0000°	510.
492T"	2620.	52500+	T8000.	0T000*	20000.	070.
ζ Τ Σ Τ•	5920.	84400*	£9000*	70000.	τοοοο.	≤90 •
+021.	1220.	69200*	1.17000*	50000*	T0000*	090*
156	•0510	16200.	-00033	£0000*	00000.	550.
ECTT.	EQTO.	.0023T	• 00053	S0000.	00000*	020.
0322	-010					
0611.	STTO.	70200.	61000.	τοοοο.	00000.	840.
90TT ·	τ9το.	48100.	9 1000*	T0000*	00000*	9 †0°
5801.	τέτο·	-00162	Et000.	T0000°	00000*	ή ηΟ·
Τζοι.	07TO .	τητοο.	TT000*	T0000*	00000*	240°
.103S	OETO	.00122	80000.	00000*	00000*	040.
9 00T *	6770.	†0T00 *	70000.	00000*	00000*	8£0 .
6160.	6070.	88000.	≤0 000 °	00000*	00000*	920*
τ <u>≤60</u> •	66 00 °	£7000.	†0000*	00000*	00000.	4E0°
•0923	68 0 0 •	09 000 *	£0000°	00000*	00000*	•035
£680.	0800.	84000°	50000.	00000*	00000*	050.
£980*	0700.	75000.	T0000*	00000*	00000*	820.
S80.	τ900'	82000.	τοοοο.	00000*	00000*	•056
6610°	2 200.	.00021	00000*	00000*	00000.	• 05ф
59 70.	ήή00°	ήΤΟΟΟ *	00000.	00000.	00000*	. 022
9270.	9E00°	0T000°	00000*	00000*	00000*	•020
3 690 •	8500.	9 0000 °	00000.	•00000	00000.	810 .
-0652	.0022	•00003	00000*	00000*	00000*	970.
0T90°	9700.	20000.	•00000	00000*	00000*	ηΤΟ •
5950.	0T00 ·	τοοοο.	00000.	00000*	00000.	SIO.
9150.	9 000 '	00000.	00000	00000.	00000*	0T0 *
T970.	£000.	00000*	•00000	00000*	00000*	800.
0010	τοοο.	00000	00000.	00000*	00000.	900*
93 50.	0000	00000.	00000.	00000.	00000.	†00°
.023L	0000.	00000.	00000.	00000.	00000*	.002
0000.0	00000.0	000000*0	00000.0	00000.0	0000010	000.0
N = 1.0	8. = N	9° = N	ή· = Ν	S. = N	0 = N	x

x	$\mathbf{N} = \mathbf{O}$	N = .2	N = .4	N = .6	N = .8	N = 1.0
.145	.00090	.00190	.00660	.02210	.0684	.1964
.150	.00104	.00213	.00716	.02337	.0709	.1998
.155	.00118	.00236	.00773	.02466	.0733	.2031
.160	.00133	.00261	.00831	.02597	.0757	.2063
.165	.00150	.00287	.00892	.02727	.0781	.2095
.170	.00168	.00315	.00953	.02859	.0805	.2127
.175	.00187	.00344	.01016	.02992	.0829	-2158
.180	.00207	.00374	.01081	.03125	.0852	.2189
.185	.00228	.00405	.01147	.03259	.0876	.2219
.190	.00251	.00438	.01214	.03393	.0899	.2248
.195	.00274	.00472	.01283	.03528	.0921	.2278
.200	.0030	.0051	.01350	.0366	.0944	.2307
.210	.0035	.0058	.0150	•0393	•0989	.2364
.220	.0041	.0066	.0164	.0421	.1034	.2420
.230	.0047	.0074	.0179	.0448	.1077	.2474
.240	.0054	.0083	.0195	.0476	.1121	.2527
.250	.0061	.0092	.0211	.0503	.1163	.2579
.260	.0069	.0102	.0227	.0531	.1206	.2630
.270	.0077	.0112	.0244	.0559	.1247	.2681
.280	•0085	.0123	.0261	.0586	.1288	.2730
.290	.0094	.0134	.0278	.0614	.1329	.2778
. 300	.0103	.0145	.0295	.0642	.1369	•2826
- 310	.0113	.0157	.0313	.0009	·1409	.2873
• 320	.0123	.0109	.033L	•0697	•1440 •1907	•2919
• 3 30	.0133 0145	.0101	.0349	.0(25	•1407	.2964
- 350	.0144	•0194	.0300	.0752	•1525	.3009
- 350	.0155	.0207	-0300 0105	.0/00	•1563	.3053
- 300	.0107	.0220	.0405	.0000	.1001	.3097
-380	.0103	·0233	0424	.0035	1676	• <u>3140</u> 9180
.300	.0203	0241	.0443	.0003	.1010	• 3102
•		.0202	.0403	.0090	•117	• 5224
.400	.0216	.0276	.0482	.0918	.1749	.3266
.420	.0242	.0306	.0522	.0973	.1821	• 3,347
.440	.0270	.0336	.0562	.1027	, .1892	• 3427
•460	.0298	•0368	.0603	.1082	.1962	.3505
.480	.0328	.0401	.0644	.1137	.2031	.3582
•500	.0358	.0434	.0685	.1191	.2099	•3657
.520	•0389	•0468	.0717	.1245	.2166	.3731
•540	.0422	·0503	.0770	.1299	.2224	. 3804
•560	.0454	.0538	.0813	.1353	.2298	.3876
.580	•0488	.0574	. 0856	.1407	.2362	• 3946

Х	N = O	N = .2	N = .4	N = .6	N = .8	N = 1.0
.600 .650 .700 .750 .800 .850 .900 .950	.0522 .0610 .0702 .0797 .0895 .0995 .1097 .1201	.0611 .0705 .0802 .0901 .1003 .1107 .1213 .1320	.0899 .1009 .1121 .1234 .1348 .1463 .1579 .1695	.1460 .1594 .1726 .1858 .1990 .2121 .2252 .2382	.2427 .2585 .2740 .2892 .3041 .3188 .3334 .3478	.4016 .4187 .4353 .4515 .4674 .4829 .4982 .5132
$\begin{array}{c} 1.0\\ 2.0\\ 3.0\\ 5.0\\ 7.0\\ 9.0\\ 10.0\\ 12.0\\ 9.0\\ 112.0\\ 15.0\\ 10.0\\$	$\begin{array}{c} \textbf{.131}\\ \textbf{.360}\\ \textbf{.647}\\ \textbf{.8943}\\ \textbf{.943}\\ \textbf{.9399}\\ \textbf{.8943}\\ \textbf{.9399}\\ \textbf{.89939}\\ \textbf{.9399}\\ \textbf{.89939}\\ \textbf{.9399}\\ \textbf{.89939}\\ \textbf{.9399}\\ \textbf{.89939}\\ \textbf{.9399}\\ \textbf{.89999}\\ \textbf{.899999}\\ \textbf{.89999}\\ \textbf{.89999}\\ \textbf{.899999}\\ \textbf{.899999}\\ \textbf{.89999}\\ .899$.143 .375 .618 .865 1.13 1.610 1.85992.358 2.608773.3.6076 4.85052.35555555 5.6055555555555555555555555555555555555	$\begin{array}{c} .181\\ .423\\ .6708\\ .111\\ .46516\\ .9167\\ .111\\ .46516\\ .91655\\ .9167\\ .111\\ .46515\\ .91655\\ .1416\\ .46516\\ .1655\\ .1416\\ .4664\\ .1644\\ .4644\\ .4644\\ .4644\\ .4644\\ .4664\\ .1644\\ .4644\\ .4664\\ .1644\\ .4644$.4644\\ .4644\\ .4644\\ .4644 .4644\\ .4644 .4644\\ .4644 .4644\\ .4644 .4644\\ .4644 .4644\\ .4644 .4644\\ .4644 .4644 .4644\\ .4644 .	$\begin{array}{c} .251\\ .506\\ .759\\ 1.010\\ 1.260\\ 1.511\\ 2.011\\ 2.262\\ 2.512\\ 2.512\\ 2.512\\ 2.512\\ 2.512\\ 2.512\\ 3.262\\ 2.512\\ 3.262\\ 2.513\\ 3.262\\ 4.512\\ 4.513\\ 5.263\\ 5.513\\ 5.263\\ 5.513\\ 7.263\\ 7.513\\ 7.513\\ 7.513\\ 8.263\\ 8.513\\ 8.263\\ 8.513\\ 9.263\\ 9.513\\ 9.263\\ 9.763\end{array}$.362 .631 .888 1.141 1.394 1.645 1.896 2.147 2.398 2.648 2.147 2.398 2.648 2.399 3.399 3.399 3.399 3.399 3.650 4.150 4.651 5.401 5.151 5.651 5.651 5.651 5.651 5.651 5.651 5.651 5.651 5.651 5.651 5.651 5.651 5.651 5.651 5.651 5.651 5.651 5.651 5.652 7.652 7.652 8.15	.528 .803 1.062 1.317 1.570 1.823 2.325 2.576 2.325 2.576 2.325 2.576 2.325 2.576 2.325 2.576 2.325 2.576 2.325 2.576 2.327 3.328 3.579 4.579 4.579 4.530 5.330 5.5830 6.3311 7.5811 7.5811 8.3311 7.5811 8.3311 9.5811 9.5811 9.5811 9.5811 9.6
40.0	9.835	9.855	9.914	10.013	10.152	10.331

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THE TEMPERATURE DISTRIBUTION FUNCTION - Z9

$\begin{array}{cccccccccccccccccccccccccccccccccccc$	x	$\mathbf{N} = \mathbf{O}$	N = .2	N = .4	N = .6	N = .8	N = 1.0
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	0.000	0.00000	0.00000	0,0000	0 00000	0.0000	• • • • • •
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	•002	•00000	.00000	.00000	00000	0.0000	0.0000
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.004	•00000	.00000	.00000	-00000	.0000	.0205
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	•006	.00000	.00000	.00000	.00000	.0000	.0290
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.008	.00000	.00000	.00000	•00000	.0001	.0355
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.010	•00000	.00000	.00000	.00000	.0002	.0410
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.012	.00000	.00000	.00000	•00000	.0003	.0459
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.014	.00000	.00000	.00000	.00000	.0006	.0502
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.016	.00000	.00000	.00000	.00001	.0009	.0543
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	•018	.00000	.00000	.00000	.00002	.0013	.0580
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	•020	.00000	.00000	.00000	-00003	•0010	.0615
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	•022	.00000	•00000	.00000	00007	.0023	•0648
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.024	.00000	.00000	.00000	.00010	•0020	.0680
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.026	.00000	.00000	.00000	.0001/	•0034 00ka	.0710
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.028	•00000	.00000	.00000	.00010	.0040	.0739
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	•030	•00000	.00000	,00001	.00025	.004	.0767
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	•032	.00000	.00000	.00001	.00023	•0054	.0794
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.034	.00000	.00000	-00002	.00031	•0061	.0820
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.036	.00000	•00000	.00002	-000/17	•0000	•0845
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	•038	•00000	.00000	.00003	.00057	.00/5	.0870
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	•040	.00000	.00000	.00004	.00067	.0003	.0094
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.042	.00000	•00000	.00005	.00079	.009I	.0917
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.044	.00000	•00000	.00006	.00001	.0099	.0940
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.046	•00000	•00000	.00007	.00105	.0101	-0902
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.048	.00000	•00000	.00009	.00119	.0122	.0903
$\begin{array}{cccccccccccccccccccccccccccccccccccc$						• 9160	•1005
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.050	.00000	.00001	.00011	.00134	.0121	1005
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.055	.00000	.00001	.00016	.00175	.0152	1027
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.060	•00000	.00002	.00023	.00222	.0173	.10()
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.065	.00001	.00003	.00032	.00273	.0105	1160
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.070	.00001	•00005	.00042	.00329	.0216	109
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.075	.00001	•00006	.00054	.00389	.0238	.1256
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	•000	.00002	•00009	•00068	.00453	.0260	.1207
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.005	•00003	.00012	•00084	.00520	.0281	.1337
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.090	•00004	.00015	.00101	.00591	.0303	.1376
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	100	.00005	.00020	.00120	.00665	.0325	.1413
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.105	•00007	.00024	.00141	.00741	.0346	.1450
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.110	.0009	.00030	.00164	.00820	.0368	.1486
$\begin{array}{cccccccccccccccccccccccccccccccccccc$.115	.00012	•00036	.00188	.00902	.0389	.1521
.125 .00010 .00050 .00242 .01071 .0432 .1588 .125 .00022 .00059 .00271 .01159 .0453 .1621 .130 .00026 .00068 .00301 .01248 .0474 .1653 .135 .00031 .00078 .00333 .01339 .0494 .1685 .140 .00037 .00089 .00367 .01431 .0515 .1716	.120	.00018	.00043	.00214	•00986	.0410	.1555
.130 .00026 .00059 .00271 .01159 .0453 .1621 .135 .00031 .00078 .00333 .01248 .0474 .1653 .140 .00037 .00089 .00367 .01431 .0515 .1716	125	•00003	.00050	.00242	.01071	.0432	.1588
.135 .00031 .00301 .01248 .0474 .1653 .135 .00031 .00078 .00333 .01339 .0494 .1685 .140 .00037 .00089 .00367 .01431 .0515 .1716	.130	•00022 00002	•00059	.00271	. 01159	.0453	.1621
.140 $.00037$ $.00089$ $.00333$ $.01339$ $.0494$ $.1685$ $.140$ $.00037$ $.00089$ $.00367$ $.01431$ $.0515$ $.1716$.135	•00020 00031	.00068	.00301	.01248	.0474	.1653
.00009 .00367 .01431 .0515 .1716	.140	•00031 00037	•000/0	.00333	•01,339	.0494	.1685
		•••••	•00009	•00367	.01431	.0515	.1716

-28-

THE TEMPERATURE DISTRIBUTION FUNCTION - 29

X	$\mathbf{N} = \mathbf{O}$	N = .2	N = .4	N = .6	N = .8	N = 1.0
.145	.00043	.00101	.00402	.01525	.0536	.1746
.150	.00050	.00114	.00438	.01620	.0556	.1776
.155	.00058	.00128	.00476	.01716	.0576	.1805
.160	.000 66	.00142	.00515	.01814	.0597	.1834
.165	.00074	.00157	.00555	.01912	.0617	.1862
.170	.00084	.00174	.00596	.02011	.0636	.1891
.175	.00094	.00191	.00639	.02111	.0656	.1918
.180	.00105	.00209	.00683	.02212	.0676	.1945
.185	.00116	.00227	.00727	.02314	.0695	.1972
.190	.00129	.00247	.00773	.02417	.0715	.1999
.195	.00142	.00268	.00820	.02520	.0734	.2025
.200	.00156	.00289	.0087	.02624	.0753	.2051
.210	.00186	·00335	.0097	.02832	.0791	.2101
.220	.00219	.00383	.0107	.03043	.0829	.2151
.230	.00255	.00435	.0118	.03256	.08 66	.2199
.240	.00294	.00490	.0128	.03470	.0902	. 2246
•250	.00337	.00549	.0140	.03685	.0939	·2293
•260	.00383	.00610	.0151	.03901	.0974	.2338
.270	.00431	.00675	.0163	.04119	.1010	.2383
.280	•00483	.00742	.0175	•04336	.1045	.2426
•290	-00537	.00813	.0187	•04555	.1079	•2469
• 300	•00594	.00007	.0200	.04774	.1114	.2512
• 310	.00054	.00903	.0213	•04993	.1147	•2553
• 320	.00(1)	.01042	.0226	.052E3	•1101	•2594
• 330	.00103	.01124	.0239	.07433	·1214	•2034
• <u>)</u>	*000001	.01200	.0273	·U2023	.1247	.20(4
.360	.00922	01295	0280	.05073	1200	.2(13
.370	.01071	01476	.0200	.00094	1 212	•2()2
.380	.01149	.01570	.0308	.06535	• 1376	-2190
• 390	.01230	.01666	.0322	.06755	.1407	.2865
.400	.0131	.0177	.0337	06975	1428	.2001
.420	.0149	.0197	.0366	.07416	.1500	.2073
.440	.0167	.0218	.0396	.07856	.1560	30/13
.460	.0186	.0240	.0427	.08295	.1620	3112
.480	.0206	.0263	.0457	.08734	.1679	.3180
.500	.0227	.0286	0489	.09172	.1738	3246
.520	.0248	.0310	.0520	.09609	.1795	.3311
.540	.0270	.0334	.0552	.10045	.1852	3375
.560	.0293	.0360	.0585	.10481	.1908	3438
.580	.0316	.0385	.0617	.10916	.1963	.3501
				-		

х	$\mathbb{N} = \mathbb{O}$	N = .2	N = .4	N = .6	N = .8	N = 1.0
.600	.0340	.0412 .0479	•0650 •0734	.1135 .1243	.2018 .2152	•3562 •3711
.700	0468	.0550	.0820	.1251	2284	3855
750	0526	0623	0006	1458	2112	· JOJJ 2005
800	0607	.0608	000/1	1564	0520 	• 3777 h120
850	0680	.0090	1082	1671	•275 9	1266
.000	0755	085)	1172	1776	2004	1207
.900	•0[55	.0074	1062	•T[]O	.2100	•4391
•970	.0052	·0935	•1203	• 1002	•2901	•4720
1.0	.091	.102	.135	.199	• 303	.465
2.0	.267	.281	• 326	.405	•524	.695
3.0	•457	•473	•522	.607	•733	.906
4.0	.651	•668	.720	.809	•938	1.113
5.0	.848	.865	.919	1.009	1.141	1.317
6.0	1.046	1.064	1.118	1.210	1.343	1.519
7.0	1.244	1.262	1.317	1.411	1.544	1.721
8.0	1.443	1.461	1.517	1.611	1.745	1.923
9.0	1.642	1.660	1.716	1.811	1.946	2.124
10.0	1.841	1.860	1.916	2.011	2.147	2.325
11.0	2,040	2.059	2.116	2.211	2.347	2.525
12.0	2.240	2.259	2.316	2.412	2.548	2.726
13.0	2.439	2.458	2.515	2.612	2.748	2.927
14.0	2.639	2.658	2.715	2.812	2.949	3.127
15.0	2.838	2.857	2.915	3.012	3.149	3.327
10.0	3.038	3.057	3.115	3.212	3.349	3.528
17.0	3.238	3.257	3.312	3.412	3.549	3.728
10.0	3.438	3.457	3.515	3.612	3.750	3.928
19.0	3.031	3.057	3.715	3.012	3.950	4.129
20.0	3.031	3.070	3.915	4.012	4.170	4.329
21.0	4.031	4.050	4,117	4.212	4.370	4.529
22.0	4.23(4.270	4.317 h 57h	4.412	4.750	4. (29
23.0	4.451	4.470	4.714	4.012 h 810	4.150	4.929
25.0	1. 826	h 856	4. (14 h Olh	4.012 5 010	4.970 5 151	5 220
25.0	5.026	5 056	4•914 5 11h	5 010	2°1211 5 259	5 520
27.0	5.236	5.256	5.31h	5.439	5,551	5.730
28.0	5 436	5.456	5.514	5.612	5 751	5.030
29.0	5.636	5.655	5.714	5.812	5.051	6.130
30.0	5.836	5-855	5.014	6.012	6.151	6.330
31.0	6.036	6.055	6.114	6.979	6.351	6.530
32.0	6.236	6.255	6.314	6.412	6.551	6.730
33.0	6.436	6.455	6.514	6.612	6.751	6.030
34.0	6.636	6.655	6.714	6.812	6.951	7,131
35.0	6.835	6.855	6.914	7.013	7,151	7,221
36.0	7.035	7.055	7.11h	7.213	7.251	7.531
37.0	7,235	7,255	7,214	7.412	7,557	7,721
38.0	7.435	7.455	7.514	7.613	7,751	7.931
39.0	7.635	7.655	7.714	7.813	7,951	8,121
40.0	7.835	7.855	7.011	8.013	8,151	8,331
			1+24-	~~~	مل <i>ع کر</i> مانه ۵ ۷۰	سرر و ر

THE TEMPERATURE DISTRIBUTION FUNCTION - Z

х	N = O	N = .2	N = .4	N = .6	N = .8	N = 1.0
0.000	0.00000	0.00000	0.00000	0.00000	0.00000	0.0000
.002	.00000	.00000	.00000	.00000	.00000	.0186
.004	.00000	.00000	.00000	.00000	.00000	.0264
.006	.00000	.00000	.00000	.00000	.00003	.0323
.008	.00000	.00000	.00000	.00000	.00009	.0373
.010	.00000	.00000	.00000	.00000	.00020	.0417
.012	.00000	.00000	.00000	.00000	.00036	.0457
.014	.00000	.00000	.00000	.00000	.00058	.0493
.016	.00000	.00000	.00000	.00001	.00085	.0527
.018	.00000	.00000	.00000	.00001	.00116	.0559
.020	.00000	.00000	.00000	.00002	.00152	.0589
.022	.00000	.00000	.00000	.00004	.00191	.0618
.024	.00000	.00000	.00000	.00005	.00234	.0646
.026	.00000	.00000	.00000	.00007	.00279	.0672
.028	.00000	.00000	.00000	.00010	.00328	.0697
.030	.00000	.00000	.00000	.00014	.00379	.0722
.032	.00000	.00000	.00001	.00018	.00432	.0746
.034	.00000	.00000	.00001	.00022	.00487	.0769
.036	.00000	.00000	.00001	.00027	.00545	.0791
•038	.00000	.00000	.00001	.00033	.00603	.0813
.040	•00000	.00000	.00002	.00040	.00664	. €834
.042	.00000	.00000	.00002	.00047	.00725	.0854
.044	.00000	.00000	.00003	.00055	.00788	.0874
.046	.00000	.00000	.00004	.00064	.00851	·0894
.048	.00000	.00000	.00005	.00073	•00916	.0913
.050	.00000	.00000	.00006	.00083	.0098	.0932
.055	.00000	•00000	.00008	.00110	.0115	.0978
.060	.00000	.00001	.00012	.00142	.0132	.1021
•065	.00000	.00002	.00017	.00177	.0149	.10 63
.070	.00000	.00002	.00024	.00216	.0167	.1103
.075	.00000	.00003	.00031	.00258	.0185	.1142
•080	.00000	.00004	•00039	.00303	.0202	.1179
•085	.00001	.00006	.00049	.00352	.0220	.1215
.090	.00002	.00008	•00060	.00403	.0238	.1251
.095	.00003	.00010	.00072	.00457	.0256	.1285
.100	.00004	.00012	.00085	.00512	.0274	.1318
.105	.00005	.00015	.00100	.00571	.0292	.1351
.110	• 0000 6	.00019	.00116	.00631	.0310	.1382
.115	80000	.00023	.00133	.00 693	.0328	.1414
.120	.00009	.00027	.00151	.00758	.0 346	.1444
.125	.00011	.00033	.00170	.00824	.0364	.1474
•130	.00013	.00039	.00191	.00892	.0382	.1503
•135	.00016	.00045	.00213	.00961	•0399	.1532
.140	.00019	.00051	.00236	.01032	.0417	.1560

-31-

THE TEMPERATURE DISTRIBUTION FUNCTION - Z

X	$\mathbf{N} = \mathbf{O}$	N = .2	N = .4	N = .6	N = .8	N = 1.0
.145	.00022	.00058	.00260	.01104	.0434	.1587
.150	.00026	.000 66	.00284	.01177	.0452	.1614
.155	.00030	.00075	.00310	.01252	.0469	.1641
.160	.00035	.00084	.00337	.01327	.0486	.1667
.165	.00039		.00366	.01404	.0504	.1693
.170	.00045	.00103	.00395	.01481	.0521	.1719
.175	.00051	.00114	.00425	.01 559	.0538	.1744
.180	.00057	.00125	.00456	.01639	.0555	.1768
.185	.00065	.00137	.00488	.01719	.0571	.1793
.190	.00072	.00150	.00521	.01799	0588	.1817
.195	.00080	.00163	.00 555	.01881	.0605	.1841
.200	.00089	.00177	.00590	.0196	.0621	.1864
.210	.00107	.00206	.00661	.0213	.0654	.1910
.220	.00127	.00237	.00736	.0230	.0687	. 1955
.230	.00149	.00271	.00813	.0247	.0719	.1999
.240	.00173	.00308	.00893	.0264	.0751	.2042
.250	.00199	.00348	.00976	.0281	.0782	.2084
.260	.00228	.00389	.01061	.0299	.0813	.2125
.270	.00259	.00433	.01149	.0316	.0844	.2166
.280	.00292	.00479	.01239	. 0334	.0875	.2206
.290	.00328	.00527	.01332	.0352	.0905	.2245
.300	-00365	.00578	.01426	.0370	.0935	.2283
- 310	.00404	.00631	.01523	.0388	•0964	.2321
• 320	200445	°00002	.01055	.0406	.0994	.2358
· 330	•00400	.00/42	.01(22	.0424	.1023	•2395 alim
• 340	.00533	.00001	.01024	.0442	.1052	·2431
· 570	.00500	.00002	•01929	.0400 01/78	.1000	.2400
• <u>500</u>	00682	.00925	.02034	.0410	.1109	.2501
380	00736	01057	·V2142	.0491	1151	.2730
.300	.00702	01125	02251	.0717	1109	2510
• • • • • •	100192	.0112)	~~JOT	•0222	° 77775	.2004
.400	.0085	.0120	.0247	.0551	.1220	.2637
.420	.0097	.0134	.0 270	.0588	.1274	.2702
.440	.0110	.0149	.0293	.0624	.1327	.2766
.460	.0123	.0165	.0317	.0661	.1380	.2828
.480	.0137	.0182	.0341	.0697	.1432	.2890
.500	.0152	.0199	.0365	.0734	.1483	.2950
.520	.0168	.0217	.0390	.0771	.1534	• 3008
•540	.0183	.0235	.0415	.0808	.1584	.3066
•500	.0199	.0254	.0441	.0844	.1633	.3123
• 500	.0216	.0273	.0467	.0881	.1682	.3179

х	N = O	N = .2	N = .4	N = .6	N = .8	N = 1.0
.600	.0234	.0292	.0493	.0917	.1730	• 3234
.650	.0279	.0343	.0559	.1008	.1848	• 3368
.700	.0328	.0397	.0627	.1098	.1964	• 3498
.750	.0379	.0453	.0697	.1188	.2077	• 3624
.800	.0433	.0510	.0767	.1278	.2188	• 3746
.850	.0488	.0570	.0838	.1367	.2297	• 3865
.900	.0546	.0631	.0911	.1456	.2404	• 3982
.950	.0605	.0694	.0984	.1545	.2509	• 4096
1.0 2.0 3.0 5.0 7.0 9.0 11.2 3.0 9.0 11.2 13.0 0.0 0.0 0.0 0.0 0.0 0.0 0.0 0.0 0.0	.067 .206 .362 .685 .685 .685 1.013 1.178 1.349 1.675 1.841 2.173 2.506 2.672 2.673 3.504 1.73 3.671 3.670 4.670 4.670 4.670 5.569 5.669 5.602 6.002 6.169	.076 .220 .377 .538 1.0396 1.121.52840 1.5284006 2.222.5695 3.33.52907 3.6852396 3.52896 3.52896 3.52907 3.52906 3.52896 3.5286	.106 $.262$ $.428$ $.753$ $.9195$ 1.2517 1.5830 1.2517 1.5830 1.2918 2.2416 2.2416 2.2416 2.2415 3.57916 3.5916 5	.163 $.336$ $.506$ $.674$ $.849$ 1.0077 1.341 1.5118 1.0178 2.3452 2.6795 3.31795 2.6796 4.0178 5.1796 3.5129 3.5129 3.6796 4.5129 4.5129 5.5139 5.5139 5.5139 5.5139 5.5139	$\begin{array}{c} .261\\ .452\\ .629\\ .801\\ 1.1409\\ 1.409\\ 1.4108\\ 1.4108\\ $	$\begin{array}{c} .421\\ .620\\ .800\\ .975\\ 1.146\\ 1.316\\ 1.485\\ 1.653\\ 1.985\\ 1.653\\ 1.989\\ 2.157\\ 4.2.329\\ 2.499\\ 2.3249\\ 2.6526\\ 3.167\\ 3.4961\\ 3.825\\ 4.3962\\ 4.396\\ 4.829\\ 5.1630\\ 5.830\\ 5.997\\ 6.163\\ 5.997\\ 6.1630\\ 6.497\\ 6.1630\\ 6.1600\\ 6.1$
39.0	6.336	6.355	6.414	6.513	6.651	6.831
40.0	6.502	6.522	6.581	6.679	6.818	6.998

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