MATHEMATICAL MODELLING AND ANALYSIS Volume 22 Number 4, July 2017, 441–463 https://doi.org/10.3846/13926292.2017.1319882 © Vilnius Gediminas Technical University, 2017 Publisher: Taylor&Francis and VGTU http://www.tandfonline.com/TMMA ISSN: 1392-6292 eISSN: 1648-3510

Global Existence and Asymptotic Behavior of Solutions for the Cauchy Problem of a Dissipative Boussinesq-Type Equation

Amin Esfahani and Hamideh B. Mohammadi

Damghan University

School of Mathematics and Computer Science, 36715-364 Damghan, Iran E-mail(*corresp.*): esfahani@du.ac.ir E-mail: amin@impa.br

Received December 15, 2016; revised April 11, 2017; published online July 15, 2017

Abstract. We consider the Cauchy problem for a Boussinesq-type equation modeling bidirectional surface waves in a convecting fluid. Under small condition on the initial value, the existence and asymptotic behavior of global solutions in some time weighted spaces are established by the contraction mapping principle.

Keywords: dissipative Boussinesq equation, asymptotic behavior, Sobolev spaces.

AMS Subject Classification: 35B30; 35Q55; 35Q72.

1 Introduction

The Kuramoto-Sivashinsky (KS) equation

$$u_t + \gamma u_{xxxx} + \alpha u_{xx} + u u_x = 0, \ u = u(x, t), \ (x, t) \in \mathbb{R} \times \mathbb{R}^+$$

is a well-known model of one-dimensional turbulence derived in various physical contexts such as chemical-reaction waves, propagation of combustion fronts in gases, surface waves in a film of a viscous liquid flowing along an inclined plane, patterns in thermal convection, rapid solidification (see e.g. [14, 19]), where α and γ are constant coefficients accounting for the long-wave instability (gain) and short-wave dissipation, respectively. By combining the dispersive effects of the KdV equation and the dissipative effects of the KS equation, the Kuramoto-Sivashinsky-Korteweg-de Vries (KS-KdV) equation

$$u_t + u_{xxx} + \gamma u_{xxxx} + \alpha u_{xx} + u u_x = 0, \ u = u(x,t), \ (x,t) \in \mathbb{R} \times \mathbb{R}^+$$

appears; which was first introduced by Benney [1]. This equation finds various applications in the study of unstable drift waves in plasmas [5], fluid flow along an inclined plane [1,16] convection in fluids with a free surface [7] the Eckhaus

instability of traveling waves [9], in solar dynamo wave [12], hydrodynamics and other fields [4].

The derivation of this equation in the physical situations mentioned above involves the assumption of unidirectional waves. The assumption of unidirectional waves for surface waves was removed in [10,13] and a modified Boussinesq system of equations was derived. One of these type of equation is the following dissipative Boussinesq equation:

$$u_{tt} - \Delta u + \Delta^2 u + \alpha \Delta u_t + \gamma \Delta^2 u_t = \Delta(\beta f(u_t) + g(u)).$$
(1.1)

Here u = u(x, t) is the unknown function of $x = (x_1, \dots, x_n) \in \mathbb{R}^n$, t > 0 and $\beta > 0$ and $\alpha \in \mathbb{R}$ are constants. The operator Δ is defined to be $\Delta = \sum_{j=1}^n \frac{\partial^2}{\partial x_j^2}$ and $\Delta^2 = \Delta \Delta$. The term u_t represents a frictional function dissipation, and the nonlinear term f(v) and g(v) are smooth functions of v under considerations and satisfies $f(v) = O(|v|^2)$ and $g(v) = O(|v|^2)$ for $v \to 0$. Equation (1.1) arises in the study of the stability of one-dimensional periodic patterns in systems with Galilean invariance and also the oscillations of elastic beams [3]. Ignoring the dissipation, (1.1) turns into the classical Boussinesq equation (see [2])

$$u_{tt} - \Delta u \pm \Delta^2 u = \Delta(u^2), u = u(x, t), (x, t) \in \mathbb{R}^n \times \mathbb{R}_+$$

appeared not only in the study of the dynamics of thin inviscid layers with free surface but also in the study of the nonlinear string, the shape-memory alloys, the propagation of waves in elastic rods and in the continuum limit of lattice dynamics or coupled electrical circuit. When $\gamma = \beta = 0$, the existence, uniqueness and long-time asymptotic of solutions to the Cauchy problem and the initial boundary value problem of equation (1.1) has been studied by several authors, see for instance [6, 11, 17, 18] and references therein.

In this paper we study the asymptotic behavior of solutions of the Cauchy problem associated to (1.1) with the initial values

$$u(x,0) = u_0(x), \quad u_t(x,0) = u_1(x).$$
 (1.2)

The article is organized as follows. In Section 2 we obtain the solution formula of (1.1) and study the decay property of the solution operators appearing in the solution formula. Then, in Section 3, we discuss the linear problem and show the decay estimates of the solutions in L^1 . We prove global existence and asymptotic behavior of solutions for the Cauchy problem (1.1) and (1.2) in L^2 in Section 4. Throughout this paper we assume $\gamma = 1 \leq -\alpha$.

Notations. Function \hat{f} denotes the Fourier transform of f(x), defined as

$$\hat{f}(\xi) = F(f)(\xi) = \int_{\mathbb{R}^n} f(x) \mathrm{e}^{\mathrm{i}x \cdot \xi} \mathrm{d}x.$$

We denote its inverse transform by F^{-1} . If a function $f \in L^r = L^r(\mathbb{R}^n)$, its usual norm is written as $||f||_{L^r}$. The usual Sobolev space of order s is defined by $W^{s,p} = W^{s,p}(\mathbb{R}^n) = (I - \Delta)^{s/2} L^p$ with the norm $||f||_{W^{s,p}} = ||(I - \Delta)^{s/2} f||_{L^p}$. The corresponding homogeneous Sobolev space of order s is defined by $\dot{W}^{s,p} = \dot{W}^{s,p}(\mathbb{R}^n) = (-\Delta)^{-s/2}L^p$ with the norm $||f||_{\dot{W}^{s,p}} = ||(-\Delta)^{s/2}||_{L^p}$. Note that $(I - \Delta)^{s/2}$ and $(-\Delta)^{s/2}$ are the Fourier multipliers defined by $(I - \Delta)^{s/2}f = F^{-1}((1 + |\xi|^2)^{s/2}\hat{f}(\xi))$ and $(-\Delta)^{s/2}f = F^{-1}(|\xi|^s\hat{f}(\xi))$. For p = 2, we write $H^s = H^s(\mathbb{R}^n) = W^{s,2}$ and $\dot{H}^s = \dot{W}^{s,2}$. For a nonnegative integer k, ∂_x^k denotes the totality of all the kth-order derivatives with respect to $x \in \mathbb{R}^n$. Also, for an interval I and a Banach space $X, C^k(I;X)$ denotes the space of k-times continuously differential functions on I with values in X.

Convention: Throughout the paper, we write sometimes function f(x, t) by f(t), where t is time variable.

2 Decay property of the linear part

The aim of this section is to derive the solution formula for the problem (1.1) and (1.2). First of all, we investigate the linear equation of (1.1).

$$u_{tt} - \Delta u + \Delta^2 u + \alpha \Delta u_t + \Delta^2 u_t = 0$$
(2.1)

with the initial data (1.2). By applying the Fourier transform to (2.1) we have

$$\hat{u}_{tt} + (|\xi|^4 - \alpha |\xi|^2)\hat{u}_t + (|\xi|^2 + |\xi|^4)\hat{u} = 0.$$
(2.2)

The corresponding initial values are given as

$$\hat{u}(\xi,0) = \hat{u}_0(\xi), \quad \hat{u}_t(\xi,0) = \hat{u}_1(\xi).$$
 (2.3)

The characteristic equation of (2.2) is

$$\lambda^{2} + (|\xi|^{4} - \alpha|\xi|^{2})\lambda + (|\xi|^{4} + |\xi|^{2}) = 0.$$

Let $\lambda = \lambda_{\pm}(\xi)$ be the corresponding eigenvalues, i.e

$$\lambda_{\pm}(\xi) = \frac{(\alpha|\xi|^2 - |\xi|^4) \pm \sqrt{(|\xi|^4 - \alpha|\xi|^2)^2 - 4(|\xi|^2 + |\xi|^4)}}{2}.$$
 (2.4)

The solution to the problem (2.2) and (2.3) is given in the form

$$\hat{u}(\xi,t) = \hat{G}(\xi,t)\hat{u}_1(\xi) + \hat{H}(\xi,t)\hat{u}_0(\xi), \qquad (2.5)$$

where

$$\hat{G}(\xi,t) = (e^{\lambda_{+}(\xi)t} - e^{\lambda_{-}(\xi)t}) / (\lambda_{+}(\xi) - \lambda_{-}(\xi)),$$
(2.6)

$$\hat{H}(\xi,t) = (\lambda_{+}(\xi)e^{\lambda_{-}(\xi)t}\lambda_{-}(\xi)e^{\lambda_{+}(\xi)t})/(\lambda_{+}(\xi) - \lambda_{-}(\xi)).$$
(2.7)

Let

$$G(x,t) = F^{-1}[\hat{G}(\xi,t)](x), \quad H(x,t) = F^{-1}[\hat{H}(\xi,t)](x).$$
(2.8)

With applying F^{-1} to (2.5), we obtain

$$u(x,t) = G(\cdot,t) * u_1 + H(\cdot,t) * u_0,$$
(2.9)

where * is the convolution in x. By our assumptions on the nonlinear terms of (1.1), we obtain the following solution formula to (1.1) and (1.2)

$$u(x,t) = G(\cdot,t) * u_1 + H(\cdot,t) * u_0 + \int_0^t G(\cdot,t-\tau) * \Delta(Z(u,u_t)(\tau))) d\tau, \quad (2.10)$$

where $Z(u, u_t)(t) = f(u(t)) + \beta g(u_t(t))$. This integral formula is derived by using the Duhamel principle. One can refer to [8, 15] to see how the Duhamel principle is applied for various nonlinear problems. Now we study the decay property of the linear equation (1.1). Our aim is to prove the following decay estimates of the solution operators G(t) and H(t) appearing in (2.9).

Lemma 1. The solution of (2.2) and (2.3) satisfies

$$\begin{aligned} |\xi|^2 (1+|\xi|^2) |\hat{u}(\xi,t)|^2 &+ |\hat{u}_t(\xi,t)|^2 \\ &\leq C e^{-c\omega(\xi)t} (|\xi|^2 (1+|\xi|^2) |\hat{u}_0(\xi)|^2 + |\hat{u}_1(\xi)|^2) \end{aligned}$$
(2.11)

for $\xi \in \mathbb{R}^n$ and $t \ge 0$, where $\omega(\xi) = |\xi|^2/(1+|\xi|^2)$.

Proof. By multiplying (2.2) by $\overline{\hat{u}}_t$ and taking the real part, we deduce that

$$\frac{1}{2}\frac{d}{dt}(|\hat{u}_t|^2 + (|\xi|^2 + |\xi|^4)|\hat{u}|^2) + (|\xi|^4 - \alpha|\xi|^2)|\hat{u}_t|^2 = 0.$$
(2.12)

Multiplying (2.2) by $\overline{\hat{u}}$ and take the real part yields

$$\frac{1}{2}\frac{d}{dt}((|\xi|^4 - \alpha|\xi|^2)|\hat{u}|^2 + 2Re(\hat{u}_t.\bar{\hat{u}})) + (|\xi|^2 + |\xi|^4|)|\hat{u}|^2 - |\hat{u}_t|^2 = 0.$$
(2.13)

Multiplying both sides of (2.12) and (2.13) by $(1 + |\xi|^2)$ and $|\xi|^2$ respectively, summing up the products yields

$$\frac{d}{dt}E + F = 0, \qquad (2.14)$$

where

$$E = (1+|\xi|^2)|\hat{u}_t|^2 + \{(1+|\xi|^2)(|\xi|^2+|\xi|^4)+|\xi|^2(|\xi|^4-\alpha|\xi|^2)\}|\hat{u}|^2 + 2|\xi|^2\Re(\hat{u}_t\bar{\hat{u}}),$$

$$F = \{2(1+|\xi|^2)(|\xi|^4-\alpha|\xi|^2)-2|\xi|^2\}|\hat{u}_t|^2 + 2|\xi|^2(|\xi|^2+|\xi|^4)|\hat{u}|^2,$$

where $\Re(\hat{u}_t \bar{\hat{u}})$ is the real part of $\hat{u}_t \bar{\hat{u}}$. It is easy to see that

$$C(1+|\xi|^2)E_0 \le E \le C(1+|\xi|^2)E_0, \qquad (2.15)$$

where

$$E_0 = |\hat{u}_t|^2 + |\xi|^2 (1 + |\xi|^2) |u|^2.$$

Noting that $F \ge |\xi|^2 E_0$ and with (2.15), we obtain

$$F \ge c \ \omega(\xi)E,\tag{2.16}$$

where $\omega(\xi) = |\xi|^2/(1+|\xi|^2)$. Using (2.14) and (2.16), we get

$$\frac{d}{dt}E + c \ \omega(\xi)E \le 0.$$

Thus

$$E(\xi, t) \le e^{-c \ \omega(\xi)t} E(\xi, 0),$$

which together with (2.15) proves the desired estimate (2.11).

Lemma 2. Assume that $\hat{G}(\xi, t)$ and $\hat{H}(\xi, t)$ are fundamental solutions of (2.1) in the Fourier space, which are given explicitly in (2.6) and (2.7). Then we have the pointwise estimates

$$|\xi|^2 (1+|\xi|^2) |\hat{G}(\xi,t)|^2 + |\hat{G}_t(\xi,t)|^2 \le C e^{-c\omega(\xi)t},$$
(2.17)

$$|\xi|^{2}(1+|\xi|^{2})|\hat{H}(\xi,t)|^{2} + |\hat{H}_{t}(\xi,t)|^{2} \le C|\xi|^{2}(1+|\xi|^{2})e^{-c\omega(\xi)t}, \qquad (2.18)$$

for $\xi \in \mathbb{R}^n$ and $t \ge 0$, where $\omega(\xi) = \frac{|\xi|^2}{1+|\xi|^2}$.

Proof. If $\hat{u}_0(\xi) = 0$, then from (2.5) we get

$$\hat{u}(\xi,t) = \hat{G}(\xi,t)\hat{u}_1(\xi), \quad \hat{u}_t(\xi,t) = \hat{G}_t(\xi,t)\hat{u}_1(\xi).$$

Substituting the equalities into (2.11) with $\hat{u}_0(\xi) = 0$ we obtain (2.17). In what follows, we consider $\hat{u}_1(\xi) = 0$. We have from (2.5) that

$$\hat{u}(\xi,t) = \hat{H}(\xi,t)\hat{u}_0(\xi), \quad \hat{u}_t(\xi,t) = \hat{H}_t(\xi,t)\hat{u}_0(\xi).$$

Substituting the equalities into (2.11) with $\hat{u}_1(\xi) = 0$, we obtain (2.18), which together with (2.17), we have completed the proof of the lemma. \Box

Lemma 3. Let l, k, j be nonnegative integers and assume that $1 \le p \le 2$. Then we have

$$\begin{aligned} \|\partial_x^k G(t) * \phi\|_{L^2} &\leq C(1+t)^{-\frac{n}{2}(\frac{1}{p}-\frac{1}{2})-\frac{k-j}{2}} \|\partial_x^j \phi\|_{\dot{W}^{-1,p}} \\ &+ Ce^{-ct} \|\partial_x^{k+l-2} \phi\|_{L^2}, \end{aligned}$$
(2.19)

$$\begin{aligned} \|\partial_x^k H(t) * \psi\|_{L^2} &\leq C(1+t)^{-\frac{n}{2}(\frac{1}{p}-\frac{1}{2})-\frac{k-j}{2}} \|\partial_x^j \psi\|_{L^p} \\ &+ Ce^{-ct} \|\partial_x^{k+l} \phi\|_{L^2}, \end{aligned}$$
(2.20)

for $0 \le j \le k$, where $k + l - 2 \ge 0$ in (2.19). Similarly, we have

$$\begin{aligned} \|\partial_x^k G_t(t) * \phi\|_{L^2} &\leq C(1+t)^{-\frac{n}{2}(\frac{1}{p}-\frac{1}{2})-\frac{k+1-j}{2}} \|\partial_x^j \phi\|_{\dot{W}^{-1,p}} \\ &+ Ce^{-ct} \|\partial_x^{k+l} \phi\|_{L^2}, \end{aligned}$$
(2.21)

$$\begin{aligned} \|\partial_x^k H_t(t) * \psi\|_{L^2} &\leq C(1+t)^{-\frac{n}{2}(\frac{1}{p}-\frac{1}{2})-\frac{k+1-j}{2}} \|\partial_x^j \psi\|_{L^p} \\ &+ Ce^{-ct} \|\partial_x^{k+l+2} \phi\|_{L^2}, \end{aligned}$$
(2.22)

for $0 \leq j \leq k+1$.

Proof. We only give a proof of (2.19). We apply the Plancherel theorem and use the pointwise estimate for \hat{G} in (2.17). This gives

$$\begin{split} \|\partial_x^k G_t(t) * \phi\|_{L^2}^2 &= \int_{\mathbb{R}^n} |\xi|^{2k} |\hat{G}(\xi,t)|^2 \|\hat{\phi}(\xi)|^2 d\xi \\ &= \int_{|\xi| \le 1} |\xi|^{2k} |\hat{G}(\xi,t)|^2 \hat{\phi}(\xi)|^2 d\xi + \int_{|\xi| \ge 1} |\xi|^{2k} |\hat{G}(\xi,t)|^2 \hat{\phi}(\xi)|^2 d\xi \\ &\le \int_{|\xi| \le 1} |\xi|^{2k-2} e^{-c|\xi|^2 t} |\hat{\phi}(\xi)|^2 d\xi \\ &+ C \int_{|\xi| \ge 1} e^{-c\omega(\xi)t} |\xi|^{2k} (|\xi|^2 (1+|\xi|^2))^{-1} |\hat{\phi}(\xi)|^2 d\xi \\ &\le C || \ |\xi|^{j-1} \hat{\phi}(\xi) \ \|_{L^p}^2 \ (\int_{|\xi| \le 1} |\xi|^{2(k-j)q} e^{-cq|\xi|^2 t} d\xi)^{\frac{1}{q}} \\ &+ C e^{-ct} \int_{|\xi| \ge 1} |\xi|^{2k-4} |\hat{\phi}(\xi)|^2 d\xi \\ &\le C || \ |\xi|^{j-1} \hat{\phi}(\xi) \ \|_{L^p}^2 \ (|| \ |\xi|^{2(k-j)} e^{-c|\xi|^2 t} ||_{L^q} \\ &+ C e^{-ct} \int_{|\xi| \ge 1} |\xi|^{2(k+l-2)} |\hat{\phi}(\xi)|^2 d\xi, \end{split}$$

where we used Hölder inequality with $\frac{1}{q} + \frac{2}{p} = 1$, $\frac{1}{p} + \frac{1}{p} = 1$. With a straight computation, we obtain

$$\||\xi|^{2(k-j)}e^{-c|\xi|^2t}\|_{L^q(|\xi|\le 1)} \le C(1+t)^{-n(\frac{1}{p}-\frac{1}{2})-(k-j)}.$$

It follows from the Hausdorff-Young inequality that

$$\||\xi|^{j-1}\hat{\phi}(\xi)\|_{L^{p}} \le \|\partial_{x}^{j}\phi\|_{\dot{W}^{-1,p}}.$$

Combining the above three inequalities yields (2.19). Similarly, we can prove (2.20)-(2.22). Thus the lemma is proved. \Box

Immediately we have from previous lemma the following corollary.

Corollary 1. Let $1 \le p \le 2$, and let k, j and l be nonnegative integers. Also, assume that G(x,t) and H(x,t) be the fundamental solution of (2.1) which are given in (2.6) and (2.7), respectively. Then we have

$$\|\partial_x^k G(t) * \Delta g\|_{L^2} \le C(1+t)^{-\frac{n}{2}(\frac{1}{p}-\frac{1}{2})-\frac{k+1-j}{2}} \|\partial_x^j g\|_{L^p} + Ce^{-ct} \|\partial_x^{k+l} g\|_{L^2},$$
(2.23)

for $0 \le k \le j+1$. It also for $0 \le k \le j+2$ holds that

$$\|\partial_x^k G_t(t) * \Delta g\|_{L^2} \le C(1+t)^{-\frac{n}{2}(\frac{1}{p}-\frac{1}{2})-\frac{k+2-j}{2}} \|\partial_x^j g\|_{L^p} + Ce^{-ct} \|\partial_x^{k+l+2} g\|_{L^2}.$$
(2.24)

3 Global existence and asymptotic behavior of solutions for L^1

The aim of this section is to prove the existence and asymptotic behavior of solutions to (1.1) and (1.2) with L^1 data. We first state the following lemma, which comes from [20].

Lemma 4. Assume that f = f(v) is smooth function, where $v = (v_1, \ldots, v_n)$ is a vector function. Suppose that $f(v) = O(|v|^{1+\theta})(\theta \ge 1$ is an integer) when $|v| \le v_0$. Then, for the integer $m \ge 0$, if $v, w \in W^{m,q}(\mathbb{R}^n) \cap L^p(\mathbb{R}^n) \cap L^\infty(\mathbb{R}^n)$ and $\|v\|_{L^\infty} \le v_0, \|w\|_{L^\infty} \le v_0$, then $f(v) - f(w) \in W^{m,r}(\mathbb{R}^n)$. Furthermore, the following inequalities hold:

$$\begin{aligned} \|\partial_x^m f(v)\|_{L^r} &\leq C \|v\|_{L^p} \|\partial_x^m v\|_{L^q} \|v\|_{L^{\infty}}^{\theta-1}, \\ \|\partial_x^m (f(v) - f(w))\|_{L^r} &\leq C \{ (\|\partial_x^m v\|_{L^q})\|v - w\|_{L^p} \\ &+ (\|v\|_{L^p} + \|w\|_{L^p} \|\partial_x^m (v - w)\|_{L^q}) \} (\|v\|_{L^{\infty}} + \|w\|_{L^{\infty}})^{\theta-1}, \end{aligned}$$
(3.1)

where $\frac{1}{r} = \frac{1}{p} + \frac{1}{q}$ and $1 \le p, q, r \le +\infty$.

Based on the decay estimates of solutions to the linear problem (2.1), we define the following solution space:

$$X = \{ u \in C([0,\infty); H^{s+2}(\mathbb{R}^n)) \bigcap C^1([0,\infty); H^s(\mathbb{R}^n)) : ||u||_X < \infty \},\$$

where

$$\|u\|_{X} = \sup_{t \ge 0} \bigg\{ \sum_{k \le s+2} (1+t)^{\frac{n}{4} + \frac{k}{2}} \|\partial_{x}^{k} u(t)\|_{L^{2}} + \sum_{k \le s} (1+t)^{\frac{n}{4} + \frac{k}{2}} \|\partial_{x}^{k} u_{t}(t)\|_{L^{2}} \bigg\}.$$

For R > 0, we define $X_R = \{ u \in X : ||u||_X \le R \}.$

Theorem 1. Let $n \ge 1$, $s \ge \max\{0, [\frac{n}{2}] - 1\}$ and suppose that functions $u_0 \in H^{s+2}(\mathbb{R}^n) \bigcap L^1(\mathbb{R}^n), u_1 \in H^s(\mathbb{R}^n) \bigcap \dot{W}^{-1,1}(\mathbb{R}^n)$ and functions f(v), g(v) are smooth and satisfy $f(v) = O(v^2), g(v) = O(v^2)$ for $v \to 0$. Put

$$E_0 := \|u_0\|_{L^1} + \|u_1\|_{\dot{W}^{-1,1}} + \|u_0\|_{H^{s+2}} + \|u_1\|_{H^s}.$$

If E_0 is suitably small, the Cauchy problem (1.1) and (1.2) has a unique global solution u(x,t) satisfying

$$X = u \in C([0,\infty); H^{s+2}(\mathbb{R}^n)) \bigcap C^1([0,\infty); H^s(\mathbb{R}^n))$$

Also, the solution satisfies the decay estimates

$$\|\partial_x^k u(t)\|_{L^2} \le CE_0(1+t)^{-\frac{n}{4}-\frac{k}{2}}, \quad \|\partial_x^l u_t(t)\|_{L^2} \le CE_0(1+t)^{-\frac{n}{4}-\frac{l+1}{2}}$$

for $0 \le k \le s+2$ and $0 \le l \le s$.

Proof. The Gagliardo-Nirenberg inequality gives

$$||u(t)||_{L^{\infty}} \le C ||\partial_x^{s_0} u||_{L^2}^{\theta} ||u||_{L^2}^{1-\theta} \le C(1+t)^{-\frac{n}{2}} ||u||_X,$$

where $s_0 = \frac{n}{2} + 1$, $\theta = \frac{n}{2s_0}$; i.e. $s \ge [\frac{n}{2}] - 1$. We define (see (2.10))

$$\Phi(u) = G(t) * u_1 + H(t) * u_0 + \int_0^t G(t-\tau) * \Delta(Z(u, u_t)(\tau))) d\tau,$$

where $Z(u, u_t)(t) = f(u(t)) - \beta g(u_t(t))$. We apply ∂_x^k to Φ and take the L^2 norm. We obtain

$$\begin{aligned} \|\partial_x^k \Phi(u)\|_{L^2} &\leq \|\partial_x^k G(t) * u_1\|_{L^2} + \|\partial_x^k H(t) * u_0\|_{L^2} \\ &+ C \int_0^t \|\partial_x^k G(t-\tau) * \Delta(Z(u,u_t)(\tau)))\|_{L^2} d\tau := I_1 + I_2 + J. \end{aligned}$$
(3.3)

First, we estimate I_1 . We apply (2.19) with p = 1, j = 0, l = 0 and get $I_1 \leq C(1+t)^{-\frac{n}{4}-\frac{k}{2}} \|u_1\|_{\dot{W}^{-1,1}} + Ce^{-ct} \|\partial_x^{(k-2)_+} u_1\|_{L^2} \leq CE_0(1+t)^{-\frac{n}{4}-\frac{k}{2}}$, where $(k-2)_+ = \max\{k-2, 0\}$.

For the term I_2 , we apply (2.20) with p = 1, j = 0 and l = 0. This yields

$$I_2 \le C(1+t)^{-\frac{n}{4}-\frac{k}{2}} \|u_0\|_{L^1} + Ce^{-ct} \|\partial_x^k u_0\|_{L^2} \le CE_0(1+t)^{-\frac{n}{4}-\frac{k}{2}}$$

Next, we estimate J. Let

$$J = \int_{0}^{t} G(t - \tau) * \Delta(Z(u, u_{t})(\tau))) d\tau$$

=
$$\int_{0}^{t/2} G(t - \tau) * \Delta(Z(u, u_{t})(\tau))) d\tau$$

+
$$\int_{t/2}^{t} G(t - \tau) * \Delta(Z(u, u_{t})(\tau))) d\tau =: J_{1} + J_{2}.$$

For the term J_1 , using (2.23) with p = 1, j = 0 and l = 0, we have

$$J_{1} \leq C \int_{0}^{t/2} (1+t-\tau)^{-\frac{n}{4}-\frac{k+1}{2}} \|Z(u,u_{t})(\tau))\|_{L^{1}} d\tau + C \int_{0}^{t/2} e^{-c(t-\tau)} \|\partial_{x}^{k}(Z(u,u_{t})(\tau)))\|_{L^{2}} d\tau =: J_{11} + J_{12}.$$

Note that by lemma (3.1) we have

$$\|f(u)\|_{L^1} \leq C \|u\|_{L^2}^2 \leq CR^2 (1+\tau)^{-\frac{n}{2}}, \|g(u_t)\|_{L^1} \leq C \|u_t\|_{L^2}^2 \leq CR^2 (1+\tau)^{-\frac{n}{2}}.$$

Therefore we have

$$J_{11} \leq CR^2 \int_0^{t/2} (1+t-\tau)^{-\frac{n}{4}-\frac{k+1}{2}} (1+\tau)^{-\frac{n}{2}} d\tau$$

$$\leq CR^2 (1+t)^{-\frac{n}{4}-\frac{k+1}{2}} \int_0^{t/2} (1+\tau)^{\frac{n}{2}} d\tau$$

$$\leq CR^2 (1+t)^{-\frac{n}{4}-\frac{k}{2}} \eta(t),$$

where

$$\eta(t) = \begin{cases} 1, & n = 1, \\ (1+t)^{-\frac{1}{2}} \ln(2+t), & n = 2, \\ (1+t)^{-\frac{1}{2}}, & n \ge 3. \end{cases}$$
(3.4)

We use (3.1) and obtain

$$\|\partial_x^k(Z(u,u_t)(\tau)))\|_{L^2} \le CR^2(1+t)^{-\frac{n}{4}-\frac{k}{2}-\frac{n}{2}}.$$
(3.5)

Consequently, we get

$$J_{12} \le CR^2 \int_0^{t/2} e^{-c(t-\tau)} (1+\tau)^{-\frac{n}{4} - \frac{k}{2} - \frac{n}{2}} d\tau \le CR^2 e^{-ct}$$

Finally, we estimate the term J_2 on the time interval [t/2, t]. Applying (2.23) with p = 2, j = k, l = 0 and using (3.5), we can estimate term J_2 as

$$J_{2} \leq C \int_{t/2}^{t} (1+t-\tau)^{-\frac{1}{2}} \|\partial_{x}^{k}(Z(u,u_{t})(\tau)))\|_{L^{2}} d\tau + C \int_{t/2}^{t} e^{-c(t-\tau)} \|\partial_{x}^{k}(Z(u,u_{t})(\tau)))\|_{L^{2}} \leq C R^{2} (1+t)^{-\frac{n}{4} - \frac{k}{2} - \frac{n-1}{2}}.$$
(3.6)

Thus we have shown that

$$J \le CR^2 (1+t)^{-\frac{n}{4} - \frac{k}{2}} \eta(t).$$

Substituting all these estimates into (3.3), we have

$$(1+t)^{\frac{n}{4}+\frac{k}{2}} \|\partial_x^k \Phi(u)\| \le CE_0 + CR^2,$$
(3.7)

for $0 \le k \le s + 2$. It follows from that (3.3)

$$\Phi(u)_t = G_t(t) * u_1 + H_t(t) * u_0 + \int_0^t G_t(t-\tau) * \Delta(Z(u,u_t)(\tau))) \|_{L^2} d\tau.$$
(3.8)

We use ∂_x^k to $\Phi(u)_t$ and take L^2 norm. This yields

$$\begin{aligned} \|\partial_x^k \Phi(u)_t\|_{L^2} &\leq \|\partial_x^k G_t(t) * u_1\|_{L^2} + \|\partial_x^k H_t(t) * u_0\|_{L^2} \\ &+ C \int_0^t \|\partial_x^k G_t(t-\tau) * \Delta(Z(u,u_t)(\tau)\|_{L^2} d\tau =: \acute{I}_1 + \acute{I}_2 + \acute{J}, \end{aligned}$$
(3.9)

for $0 \le k \le s$. For the term I_1 , we apply (2.21) with p = 1, j = 0 and l = 0 and obtain

$$\acute{I_1} \le C(1+t)^{-\frac{n}{4} - \frac{k+1}{2}} \|u_1\|_{\dot{W}^{-1,1}} + Ce^{-ct} \|\partial_x^k u_1\|_{L^2} \le CE_0(1+t)^{-\frac{n}{4} - \frac{k+1}{2}}.$$

Also, for the term I_2 , we apply (2.22) with p = 1, j = 0 and l = 0 and get

$$I_2 \le C(1+t)^{-\frac{n}{4} - \frac{k+1}{2}} \|u_0\|_{L^1} + Ce^{-ct} \|\partial_x^{k+2}u_0\|_{L^2} \le CE_0(1+t)^{-\frac{n}{4} - \frac{k+1}{2}}$$

To estimate the nonlinear term \hat{J} , we divide as $\hat{J} = \hat{J}_1 + \hat{J}_2$, where \hat{J}_1 and \hat{J}_2 correspond to the time intervals [0, t/2] and [t/2, t], respectively. By applying (2.24) with p = 1, j = 0 and l = 0, we have

$$\begin{split} \dot{J}_1 &\leq C \int_0^{t/2} (1+t-\tau)^{-\frac{n}{4} - \frac{k+2}{2}} \|Z(u, u_t)(\tau)\|_{L^1} d\tau \\ &+ C \int_0^{t/2} e^{-c(t-\tau)} \|\partial_x^{k+2} (Z(u, u_t)(\tau)\|_{L^2} d\tau =: J_{11} + J_{12}. \end{split}$$

By (3.1), we obtain

$$||Z(u, u_t)(\tau)||_{L^1} \le CR^2(1+\tau)^{-\frac{n}{2}}.$$

Therefore we get

$$\begin{split} J_{11} &\leq CR^2 \int_0^{t/2} (1+t-\tau)^{-\frac{n}{4}-\frac{k+2}{2}} (1+\tau)^{-\frac{n}{2}} d\tau \\ &\leq CR^2 (1+t)^{-\frac{n}{4}-\frac{k+1}{2}} \eta(t). \end{split}$$

Similarly as before, we can estimate J_{12} and obtain $J_{12} \leq CR^2 e^{-ct}$. Finally, we estimate the term J_2 by using (2.24) with p = 2, j = k + 2, l = 0 and get

$$\begin{aligned} \dot{J}_2 &\leq C \int_{t/2}^t \|\partial_x^{k+2}(Z(u,u_t)((\tau))\|_{L^2} d\tau \\ &+ C \int_{t/2}^t -\frac{n}{4} - \frac{k+1}{2} - \frac{n-1}{2} e^{-c(t-\tau)} \|\partial_x^{k+2}(Z(u,u_t)(\tau))\|_{L^2} \\ &\leq C R^2 \int_{t/2}^t (1+\tau)^{-\frac{n}{4} - \frac{k+1}{2} - \frac{n}{2}} d\tau \leq C R^2 (1+t)^{-\frac{n}{4} - \frac{k+1}{2} - \frac{n-1}{2}}. \end{aligned}$$

Consequently we have that

$$\dot{J} \le CR^2 (1+t)^{-\frac{n}{4} - \frac{k+1}{2}} \eta(t).$$

The above inequality implies

$$(1+t)^{\frac{n}{4}+\frac{k+1}{2}} \|\partial_x^k \Phi(u)_t\|_{L^2} \le CE_0 + CR^2.$$
(3.10)

Combining (3.10) and (3.7) and taking E_0 and R suitably small, we obtain $\|\Phi(u)\|_X \leq R$. For $u, \tilde{u} \in X_R$, (3.3) gives

$$\begin{split} \|\partial_x^k(\varPhi(u) - \varPhi(\tilde{u}))\|_{L^2} &= \int_0^t \|\partial_x^k G(t - \tau) * \varDelta(f(u) - f(\tilde{u}) - \beta(g(u_t)) \\ &- g(\tilde{u}_t))(\tau)\|_{L^2} d\tau = \int_0^{t/2} \|\partial_x^k G(t - \tau) * \varDelta(f(u) - f(\tilde{u}) - \beta(g(u_t)) \\ &- g(\tilde{u}_t))(\tau)\|_{L^2} d\tau + \int_{t/2}^t \|\partial_x^k G(t - \tau) * \varDelta(f(u) - f(\tilde{u})) \\ &- \beta(g(u_t) - g(\tilde{u}_t))(\tau)\|_{L^2} d\tau =: J_1 + J_2. \end{split}$$

For the term J_1 , we apply (2.23) with p = 1, j = 0 and l = 0, we arrive at

$$J_{1} \leq C \int_{0}^{t/2} (1+t-\tau)^{-\frac{n}{4}-\frac{k+1}{2}} \|(f(u)-f(\tilde{u})(\tau)-\beta(g(u_{t})-g(\tilde{u}))(\tau)\|_{L^{1}} d\tau +C \int_{0}^{t/2} e^{-c(t-\tau)} \|\partial_{x}^{k}(f(u)-f(\tilde{u})-\beta(g(u_{t}-g(u_{t})))(\tau)\|_{L^{2}} d\tau =: J_{11}+J_{12}.$$

By (3.2), we can estimate J_{11} as

$$J_{11} \leq CR \|u - \tilde{u}\|_X \int_0^{t/2} (1 + t - \tau)^{-\frac{n}{4} - \frac{k+1}{2}} (1 + \tau)^{-\frac{n}{2}} d\tau$$

$$\leq CR \|u - \tilde{u}\|_X (1 + t)^{-\frac{n}{4} - \frac{k}{2}} (1 + t)^{-\frac{n}{4} - \frac{k}{2}} \eta(t),$$

where η be defined in (3.4). It follows from the Gagliardo-Nirenberg inequality and (3.2) that

$$J_{12} \leq \int_{0}^{t/2} e^{-c(t-\tau)} \Big[(\|\partial_{x}^{k}u\|_{L^{2}} + \|\partial_{x}^{k}\tilde{u}\|_{L^{2}}) \|u - \tilde{u}\|_{L^{\infty}} \\ + (\|u\|_{L^{\infty}} + \|\tilde{u}\|_{L^{\infty}}) \|\partial_{x}^{k}(u - \tilde{u})\|_{L^{2}} + (\|\partial_{x}^{k}u_{t}\|_{L^{2}} + \|\partial_{x}^{k}\tilde{u}_{t}\|_{L^{2}}) \\ \times \|u_{t} - \tilde{u}_{t}\|_{L^{\infty}} + (\|u_{t}\|_{L^{\infty}} + \|\tilde{u}_{t}\|_{L^{\infty}}) \|\partial_{x}^{k}(u_{t} - \tilde{u}_{t})\|_{L^{2}} \Big] d\tau \\ \leq CR \int_{0}^{t/2} e^{-c(t-\tau)} (1+\tau)^{-\frac{n}{4} - \frac{k}{2} - \frac{n}{2}} \|u - \tilde{u}\|_{X} d\tau \\ \leq CR \|u - \tilde{u}\|_{X} e^{-ct}.$$

Finally, we estimate term J_2 on the time [t/2, t]. Applying (2.23) with p = 2, j = k, l = 0 and using (3.2), we obtain

$$J_{2} \leq C \int_{t/2}^{t} (1+t-\tau)^{-\frac{1}{2}} \|\partial_{x}^{k}(f(u)-f(\tilde{u})-\beta(g(u_{t})-g(\tilde{u}_{t}))(\tau))\|_{L^{2}} d\tau$$

+ $C \int_{t/2}^{t} e^{-c(t-\tau)} \|\partial_{x}^{k}(f(u)-f(\tilde{u})-\beta(g(u_{t})-g(\tilde{u}_{t}))(\tau))\|_{L^{2}}$
$$\leq CR \|u-\tilde{u}\|_{X} \int_{t/2}^{t} (1+t-\tau)^{-\frac{1}{2}} (1+t)^{-\frac{n}{4}-\frac{k}{2}-\frac{n}{2}} d\tau$$

$$\leq CR \|u-\tilde{u}\|_{X} (1+t)^{-\frac{n}{4}-\frac{k}{2}-\frac{n-1}{2}},$$

which implies

$$(1+t)^{\frac{n}{4}+\frac{k}{2}} \|\partial_x^k \Phi(u) - \Phi(\tilde{u})\|_{L^2} \le CR \|u - \tilde{u}\|_X.$$

Similarly, for $0 \le k \le s$ and $u, \tilde{u} \in X$ from (3.2) and (2.24), we deduce that

$$\begin{aligned} \|\partial_x^k (\Phi(u) - \Phi(\tilde{u})_t\|_{L^2} \\ &= \int_0^t \|\partial_x^k G_t(t-\tau) * \Delta(f(u) - f(\tilde{u}) - \beta(g(u_t) - g(\tilde{u}_t))(\tau)\|_{L^2} d\tau \end{aligned}$$

$$= \int_{0}^{t/2} \|\partial_{x}^{k} G_{t}(t-\tau) * \Delta(f(u) - f(\tilde{u}) - \beta(g(u_{t}) - g(\tilde{u}_{t}))(\tau)\|_{L^{2}} d\tau + \int_{t/2}^{t} \|\partial_{x}^{k} G_{t}(t-\tau) * \Delta(f(u) - f(\tilde{u}) - \beta(g(u_{t}) - g(\tilde{u}_{t}))(\tau)\|_{L^{2}} d\tau =: \hat{J}_{1} + \hat{J}_{2}.$$

For the term J_1 , we use (2.24) with p = 1, j = 0 and l = 2. We have

$$\begin{split} \dot{f_1} &\leq C \int_0^{t/2} (1 + t - \tau)^{-\frac{n}{4} - \frac{k+2}{2}} \|f(u) - f(\tilde{u})(\tau) - \beta(g(u_t) - g(\tilde{u}_t))(\tau)\|_{L^1} d\tau \\ &+ C \int_0^{t/2} e^{-c(t - \tau)} \|\partial_x^{k+2}(f(u) - f(\tilde{u}_t) - \beta(g(u_t) - g(\tilde{u}_t))(\tau)\|_{L^2} d\tau \\ &=: J_{11}' + J_{12}'. \end{split}$$

By (3.2), we have

$$\begin{split} J_{11} &\leq \int_{0}^{t/2} (1+t-\tau)^{-\frac{n}{4}-\frac{k+2}{2}} (\|u\|_{L^{2}}+\|\tilde{u}_{L^{2}})\|u-\tilde{u}\|_{L^{2}} \\ &+ (\|u_{t}\|_{L^{2}}+\|\tilde{u}_{t}\|_{L^{2}})(\|u_{t}-\tilde{u}_{t}\|_{L^{2}})d\tau \\ &\leq CR\|u-\tilde{u}\|_{X}(1+t)^{-\frac{n}{4}-\frac{k+1}{2}} \int_{0}^{t/2} (1+t-\tau)^{-\frac{n}{2}}d\tau \\ &\leq CR\|u-\tilde{u}\|_{X}(1+t)^{-\frac{n}{4}-\frac{k+1}{2}}\eta(t), \end{split}$$

where η be defined in (3.4). Also, the term J_{12} is estimated similarly as before and we can estimate the term J_{12} as

$$J_{12} \le CR \|u - \tilde{u}\|_X e^{-ct}.$$

By applying (2.24) with p = 2, j = k + 2, l = 0 and (3.1), we obtain

$$\begin{split} \dot{J_2} &\leq C \int_{t/2}^t \|\partial_x^{k+2} (f(u) - f(\tilde{u}) - \beta(g(u_t) - g(\tilde{u}_t)(\tau))\|_{L^2} d\tau \\ &+ C \int_{t/2}^t e^{-c(t-\tau)} \|\partial_x^{k+2} (f(u) - f(\tilde{\tilde{u}}) - \beta(g(u_t) - g(\tilde{u}_t)(\tau))\|_{L^2} d\tau \\ &\leq \int_{t/2}^t (\|\partial_x^{k+2}u\|_{L^2} + \|\partial_x^{k+2}\tilde{u}\|_{L^2})\|u - \tilde{u}\|_{L^\infty} \\ &+ (\|u\|_{L^\infty} + \|\tilde{u}\|_{L^\infty})\|\partial_x^{k+2} (u - \tilde{u})\|_{L^2} \\ &+ ((\|\partial_x^{k+2}u_t\|_{L^2} + \|\partial_x^{k+2}\tilde{u}_t\|_{L^2})\|u - \tilde{u}_t\|_{L^\infty} \\ &+ (\|u_t\|_{L^\infty} + \|\tilde{u}_t\|_{L^\infty})\|\partial_x^{k+2} (u - \tilde{u}_t)\|_{L^2})d\tau \\ &\leq CR \|u - \tilde{u}\|_X \int_{t/2}^t (1 + \tau)^{-\frac{n}{4} - \frac{k+2}{2} - \frac{n}{2}} d\tau \\ &\leq CR \|u - \tilde{u}\|_X (1 + t)^{-\frac{n}{4} - \frac{k+1}{2} - \frac{n-1}{2}} \\ &\leq CR \|u - \tilde{u}\|_X (1 + t)^{-\frac{n}{4} - \frac{k+1}{2}}. \end{split}$$

Substituting all these estimates together with the previous estimate and taking R suitably small, yields

$$\|\Phi(u) - \Phi(\tilde{u})\|_X \le \frac{1}{2} \|u - \tilde{u}\|_X.$$
(3.11)

From (3.11), we deduce that Φ is strictly contracting mapping. Then there exists a fixed point $u \in X_R$ of the mapping Φ , which is a solution to (1.1)–(1.2). The proof of the theorem is now complete. \Box

The proof of the previous theorem shows that when $n \ge 2$, the solution u to the integral equation (2.10) is asymptotic to the linear solution $u_L(t)$ given by the formula $u_L(t) = G(t) * u_1 + H(t) * u_0$ as $t \to \infty$. This result is stated as follows.

Lemma 5. Let $n \ge 2$ and assume the same conditions of Theorem 1. Then the solution u of problem (1.1)–(1.2) which is constructed in Theorem 1, can be approximated by the solution u_L to the linearized problem (2.1), (2.2) as $t \to \infty$. More precisely, we have

$$\begin{aligned} \|\partial_x^k(u-u_L)(t)\|_{L^2} &\leq CE_0^2(1+t)^{-\frac{n}{4}-\frac{k}{2}}\eta(t),\\ \|\partial_x^k(u-u_L)_t(t)\|_{L^2} &\leq CE_0^2(1+t)^{-\frac{n}{4}-\frac{k+1}{2}}\eta(t), \end{aligned}$$

for $0 \le k \le s+2$ and $0 \le k \le s$, respectively, where $u_L(t) := G(t) * u_1 + H(t) * u_0$ is the linear solution and $\eta(t)$ is defined in (3.4).

4 Decay estimates of solutions for L^2

In the previous section, we have proved global existence and asymptotic behavior of solutions to the Cauchy problem (1.1)–(1.2) with L^1 data.

In this section, we prove a similar decay estimate of solution with L^2 data for $n \ge 2$. Based on the decay estimates of solutions to the linear problem (2.1)-(2.2), we define the following solution space:

$$X = \{ u \in C([0,\infty); H^{s+2}(\mathbb{R}^n)) \bigcap C^1([0,\infty); H^s(\mathbb{R}^n)) : \|u\|_X < \infty \},\$$

where

$$\|u\|_{X} = \sup_{t \ge 0} \left\{ \sum_{k \le s+2} (1+t)^{\frac{k}{2}} \|\partial_{x}^{k} u(t)\|_{L^{2}} + \sum_{k \le s} (1+t)^{\frac{k}{2}} \|\partial_{x}^{k} u_{t}(t)\|_{L^{2}} \right\}.$$

For R > 0, we define

$$X_R = \{ u \in X : \|u\|_X \le R \}.$$

Note that from the Gagliardo-Nirenberg inequality for $u \in X_R$, we have

$$||u(t)||_{L^{\infty}} \le C(1+t)^{-\frac{n}{4}}.$$

Theorem 2. Suppose that $u_0 \in H^{s+2}, u_1 \in H^s(\mathbb{R}^n) \cap \dot{W}^{-1,2}(\mathbb{R}^n)$, such that $n \geq 1, s \geq max\{0, [\frac{n}{2}] - 1\}$, and the functions f(v) and g(v) are smooth and satisfies $f(v) = O(v^2), g(v) = O(v^2)$ for $v \to 0$. Let

$$E_1 := \|u_0\|_{L^2} + \|u_1\|_{\dot{W}^{-1,2}} + \|u_0\|_{H^{s+2}} + \|u_1\|_{H^s}.$$

If E_0 is suitably small, the Cauchy problem (1.1) and (1.2) has a unique global solution u(x,t) satisfying

$$X = u \in C([0,\infty); H^{s+2}(\mathbb{R}^n)) \bigcap C^1([0,\infty); H^s(\mathbb{R}^n)).$$

The solution u also satisfies the decay estimate

$$\|\partial_x^k u(t)\|_{L^2} \le CE_1(1+t)^{-\frac{k}{2}}, \quad \|\partial_x^h u_t(t)\|_{L^2} \le CE_1(1+t)^{-\frac{h+1}{2}}, \tag{4.1}$$

for $0 \le k \le s+2$ and $0 \le h \le s$.

Proof. Let the mapping Φ be defined in (3.3). Applying ∂_x^k to Φ and take L^2 norm. We have

$$\begin{aligned} \|\partial_x^k \Phi(u)\|_{L^2} &\leq \|\partial_x^k G(t) * u_1\|_{L^2} + \|\partial_x^k H(t) * u_0\|_{L^2} \\ &+ C \int_0^t \|\partial_x^k G(t-\tau) * \Delta(Z(u,u_t)(\tau))\|_{L^2} d\tau := I_1 + I_2 + J_2. \end{aligned}$$

We use (2.19) with p = 2 and j = l = 0 and get

$$I_1 \leq C(1+t)^{-\frac{k}{2}} \|u_1\|_{\dot{W}^{-1,2}} + Ce^{-ct} \|\partial_x^{(k-2)_+} u_1\|_{L^2} \leq CE_1(1+t)^{-\frac{k}{2}},$$

where $(k-2)_{+} = \max\{k-2, 0\}$. By applying (2.20) with p = 2 and j = l = 0, we get

$$I_2 \le C(1+t)^{-\frac{k}{2}} \|u_0\|_{L^2} + Ce^{-ct} \|\partial_x^k u_0\|_{H^{s+2}} \le CE_1(1+t)^{-\frac{k}{2}}.$$

To estimate the nonlinear J, as in the previous section, we divide as $J = J_1 + J_2$ where J_1 and J_2 correspond to the time intervals [0, t/2] and [t/2, t], respectively. For the term J_1 , we use (2.23) with p = 1 and j = l = 0 and deduce that

$$J_{1} \leq C \int_{0}^{t/2} (1+t-\tau)^{-\frac{n}{4}-\frac{k+1}{2}} \|Z(u,u_{t})(\tau)\|_{L^{1}} d\tau + C \int_{0}^{t/2} e^{-c(t-\tau)} \|\partial_{x}^{k}(Z(u,u_{t})(\tau)\|_{L^{2}} d\tau =: J_{11} + J_{12}.$$

By (3.1), we have $||Z(u, u_t)(\tau)||_{L^1} \leq CR^2$. Thus we can estimate the J_{11} as

$$J_{11} \leq CR^2 \int_0^{t/2} (1+t-\tau)^{-\frac{n}{4}-\frac{k+1}{2}} (1+\tau)^{-\frac{n}{2}} d\tau$$

$$\leq CR^2 (1+t)^{-\frac{k}{2}} \int_0^{t/2} (1+\tau)^{-\frac{n}{4}-\frac{1}{2}} d\tau \leq CR^2 (1+t)^{-\frac{k}{2}}.$$

By applying (3.1) and the Gagliardo-Nirenberg inequality, we get

$$\begin{aligned} \|\partial_x^k f(u)\|_{L^2} &\leq C \|u\|_{L^{\infty}} \|\partial_x^k u\|_{L^2} \leq C(1+t)^{\frac{n}{4}-\frac{k}{2}} R^2, \\ |\partial_x^k g(u_t)\|_{L^2} &\leq C \|u_t\|_{L^{\infty}} \|\partial_x^k u_t\|_{L^2} \leq C(1+t)^{\frac{n}{4}-\frac{k}{2}} R^2. \end{aligned}$$
(4.2)

455

Thus we have

$$J_{12} \le CR^2 \int_0^{t/2} e^{-c(t-\tau)} (1+\tau)^{-\frac{n}{4} - \frac{k}{2}} d\tau \le CR^2 e^{-ct}$$

It follows from (2.23) with p = 1, j = k and l = 2 that

$$J_{2} \leq C \int_{t/2}^{t} (1+t-\tau)^{-\frac{n}{4}-\frac{1}{2}} \|\partial_{x}^{k}(Z(u,u_{t})(\tau)\|_{L^{1}} d\tau + C \int_{t/2}^{t} e^{-c(t-\tau)} \|\partial_{x}^{k+2}(Z(u,u_{t})(\tau)\|_{L^{2}} =: J_{21} + J_{22}.$$

We have from (3.2) that

$$J_{21} \leq C \int_{t/2}^{t} (1+t-\tau)^{-\frac{n}{4}-\frac{1}{2}} (\|u\|_{L^{2}}\|\|\partial_{x}^{k}u\|_{L^{2}} + \|u_{t}\|_{L^{2}}\|\partial_{x}^{k}u_{t}\|_{L^{2}})d\tau$$

$$\leq CR^{2} \int_{t/2}^{t} (1+t-\tau)^{-\frac{n}{4}-\frac{1}{2}} (1+\tau)^{-\frac{k}{2}} d\tau$$

$$\leq CR^{2} (1+t)^{-\frac{k}{2}} \int_{t/2}^{t} (1+t-\tau)^{-\frac{n}{4}-\frac{1}{2}} d\tau \leq CR^{2} (1+t)^{-\frac{k}{2}}.$$

To estimate the term J_{22} , we have from (4.2) that

$$J_{22} \leq C \int_{t/2}^{t} e^{-c(t-\tau)} \|\partial_x^{k+2}(Z(u,u_t)(\tau)\|_{L^2} d\tau$$

$$\leq CR^2 \int_{t/2}^{t} e^{-c(t-\tau)} (1+\tau)^{-\frac{n}{4} - \frac{k+2}{2}} d\tau \leq CR^2 (1+t)^{-\frac{k}{2}}.$$

The above inequality shows that

$$(1+t)^{\frac{k}{2}} \|\partial_x^k \Phi(u)\| \le CE_1 + CR^2.$$
(4.3)

We deduce from (3.3) that

$$(\Phi(u))_t = G_t(t) * u_1 + H_t(t) * u_0 + \int_0^t G_t(t-\tau) * \Delta(Z(u,u_t)(\tau)) \|_{L^2} d\tau.$$
(4.4)

Applying ∂_x^k to $\varPhi(u)_t$ and taking $L^2\text{-norm}$ we have

$$\begin{aligned} \|\partial_x^k \Phi(u)_t\|_{L^2} &\leq & \|\partial_x^k G_t(t) * u_1\|_{L^2} + \|\partial_x^k H_t(t) * u_0\|_{L^2} \\ &+ C \int_0^t \|\partial_x^k G_t(t-\tau) * \Delta(Z(u,u_t)(\tau)\|_{L^2} d\tau \; =: \; \acute{I_1} + \acute{I_2} + \acute{J}. \end{aligned}$$

To estimate the term I_1 , apply (2.21) with p = 2, l = j = 0. It yields

$$I_1 \le C(1+t)^{-\frac{k+1}{2}} \|u_1\|_{\dot{W}^{-1,2}} + Ce^{-ct} \|\partial_x^k u_1\|_{L^2} \le CE_1(1+t)^{-\frac{k+1}{2}}.$$

Similarly, using (2.22) with p = 2, j = l = 0, we have

$$I_2 \le C(1+t)^{-\frac{k+1}{2}} \|u_0\|_{L^2} + Ce^{-ct} \|\partial_x^{k+2}u_0\|_{L^2} \le CE_1(1+t)^{-\frac{k+1}{2}}.$$

To estimate the nonlinear term \hat{J} , let

$$\begin{split} \dot{J} &= C \int_{0}^{t/2} \|\partial_{x}^{k} G_{t}(t-\tau) * \Delta(Z(u,u_{t})(\tau))\|_{L^{2}} d\tau \\ &+ C \int_{t/2}^{t} \|\partial_{x}^{k} G_{t}(t-\tau) * \Delta(Z(u,u_{t})(\tau))\|_{L^{2}} d\tau \; =: \; \acute{I_{1}} + \acute{I_{2}} + \acute{J}. \end{split}$$

It yields from (2.24) with p = 1 and j = l = 0 that

$$\begin{aligned} \dot{J}_1 &\leq C \int_0^{t/2} (1+t-\tau)^{-\frac{n}{4} - \frac{k+2}{2}} \|Z(u,u_t)(\tau)\|_{L^1} d\tau \\ &+ C \int_0^{t/2} e^{-c(t-\tau)} \|\partial_x^{k+2} (Z(u,u_t)(\tau)\|_{L^2} d\tau := J_{11} + J_{12}. \end{aligned}$$

We obtain from (3.1) that

$$||Z(u, u_t)(\tau)||_{L^1} \le CR^2.$$

Thus we can estimate J_{11} as

$$\begin{split} \dot{J_{11}} &\leq CR^2 \int_0^{t/2} (1+t-\tau)^{-\frac{n}{4} - \frac{k+2}{2}} d\tau \\ &\leq CR^2 (1+t)^{-\frac{k+1}{2}} \int_0^{t/2} (1+t-\tau)^{-\frac{n}{4} - \frac{1}{2}} d\tau \leq CR^2 (1+t)^{-\frac{k+1}{2}}. \end{split}$$

For the term J_{12} , we have from (4.2) that

$$J_{12} \leq CR^2 \int_0^{t/2} e^{-c(t-\tau)} (1+t)^{-\frac{n}{4} - \frac{k+2}{2}} d\tau \leq CR^2 e^{-ct}.$$

Applying (2.24) with p = 2, j = k + 2 and l = 0, we get

$$\begin{split} \dot{J_2} &\leq C \int_{t/2}^t \|\partial_x^{k+2}(Z(u,u_t)(\tau))\|_{L^2} d\tau \\ &+ C \int_{t/2}^t e^{-c(t-\tau)} \|\partial_x^{k+2}(Z(u,u_t)(\tau))\|_{L^2} d\tau \\ &\leq C R^2 \int_{t/2}^t (1+\tau)^{-\frac{n}{4} - \frac{k+2}{2}} d\tau \\ &\leq C R^2 (1+t)^{-\frac{k+1}{2}} \int_{t/2}^t (1+\tau)^{-\frac{n}{4} - \frac{1}{2}} d\tau \leq C R^2 (1+t)^{-\frac{k+1}{2}}. \end{split}$$

Thus we have

$$(1+t)^{-\frac{k+1}{2}} \|\partial_x^k \Phi(u)_t\|_{L^2} \le CE_1 + CR^2.$$
(4.5)

457

Combining (4.3) and (4.5) and taking E_0 and R suitably small, we obtain $\|\Phi(u)\|_X \leq R$.

For $u, \tilde{u} \in X_R$, by using (3.3) we obtain

$$\begin{split} &|\partial_x^k (\varPhi(u) - \varPhi(\tilde{u}))||_{L^2} \\ &= \int_0^t \|\partial_x^k G(t - \tau) * \varDelta(f(u) - f(\tilde{u}) - \beta(g(u_t)) - g(\tilde{u}_t))(\tau)\|_{L^2} d\tau \\ &= \int_0^{t/2} \|\partial_x^k G(t - \tau) * \varDelta(f(u) - f(\tilde{u}) - \beta(g(u_t)) - g(\tilde{u}_t))(\tau)\|_{L^2} d\tau \\ &+ \int_{t/2}^t \|\partial_x^k G(t - \tau) * \varDelta(f(u) - f(\tilde{u}) - \beta(g(u_t)) - g(\tilde{u}_t))(\tau)\|_{L^2} d\tau \\ &=: J_1 + J_2. \end{split}$$

By applying (2.23) with p = 1, j = 0 and l = 0, we have

$$J_{1} \leq C \int_{0}^{t/2} (1+t-\tau)^{-\frac{n}{4}-\frac{k+1}{2}} \| (f(u)-f(\tilde{u})(\tau)-\beta(g(u_{t})-g(\tilde{u}_{t}))(\tau) \|_{L^{1}} d\tau$$
$$+ C \int_{0}^{t/2} e^{-c(t-\tau)} \| \partial_{x}^{k}(f(u)-f(\tilde{u})-\beta(g(u_{t}-g(\tilde{u}_{t})))(\tau) \|_{L^{2}} d\tau$$
$$=: J_{11} + J_{12}.$$

Using (3.2), we get

$$J_{11} \leq CR \|u - \tilde{u}\|_X \int_0^{t/2} (1 + t - \tau)^{-\frac{n}{4} - \frac{k+1}{2}} d\tau$$

$$\leq CR \|u - \tilde{u}\|_X (1 + t)^{-\frac{k}{2}} \int_0^{t/2} (1 + t - \tau)^{-\frac{n}{4} - \frac{1}{2}} d\tau.$$

$$\leq CR \|u - \tilde{u}\|_X (1 + t)^{-\frac{k}{2}}.$$

Also, we have

$$J_{12} \leq \int_{0}^{t/2} e^{-c(t-\tau)} \Big[(\|\partial_{x}^{k}u\|_{L^{2}} + \|\partial_{x}^{k}\tilde{u}\|_{L^{2}}) \|u - \tilde{u}\|_{L^{\infty}} \\ + (\|u\|_{L^{\infty}} + \|\tilde{u}\|_{L^{\infty}}) \|\partial_{x}^{k}(u - \tilde{u})\|_{L^{2}} + (\|\partial_{x}^{k}u_{t}\|_{L^{2}} + \|\partial_{x}^{k}\tilde{u}_{t}\|_{L^{2}}) \|u_{t} - \tilde{u}_{t}\|_{L^{\infty}} \\ + (\|u_{t}\|_{L^{\infty}} + \|\tilde{u}_{t}\|_{L^{\infty}}) \|\partial_{x}^{k}(u_{t} - \tilde{u}_{t})\|_{L^{2}} \Big] d\tau \\ \leq CR \|u - \tilde{u}\|_{X} \int_{0}^{t/2} e^{-c(t-\tau)} (1+\tau)^{-\frac{n}{4} - \frac{k}{2}} d\tau \leq CR \|u - \tilde{u}\|_{X} e^{-ct}.$$

To estimate the term J_2 , apply (2.23) with p = 1, j = k, l = 2. We obtain

$$J_{2} \leq C \int_{t/2}^{t} (1+t-\tau)^{-\frac{n}{4}-\frac{1}{2}} \|\partial_{x}^{k}(f(u)-f(\tilde{u})-\beta(g(u_{t})-g(\tilde{u}_{t}))(\tau))\|_{L^{1}} d\tau + C \int_{t/2}^{t} e^{-c(t-\tau)} \|\partial_{x}^{k+2}(f(u)-f(\tilde{u})-\beta(g(u_{t})-g(\tilde{u}_{t}))(\tau))\|_{L^{2}} d\tau =: J_{21}+J_{22}.$$

By using (3.2), we get

$$\begin{aligned} J_{21} &\leq \int_{t/2}^{t} (1+t-\tau)^{-\frac{n}{4}-\frac{1}{2}} \Big[(\|\partial_{x}^{k}u\|_{L^{2}}+\|\partial_{x}^{k}\tilde{u}\|_{L^{2}}) \|u-\tilde{u}\|_{L^{2}} \\ &+ (\|u\|_{L^{2}}+\|\tilde{u}\|_{L^{2}})) \|\partial_{x}^{k}(u-\tilde{u})\|_{L^{2}} + (\|\partial_{x}^{k}u_{t}\|_{L^{2}}+\|\partial_{x}^{k}\tilde{u}_{t}\|_{L^{2}}) \|u_{t}-\tilde{u}_{t}\|_{L^{2}} \\ &+ (\|u_{t}\|_{L^{2}}+\|\tilde{u}_{t}\|_{L^{2}}) \|\partial_{x}^{k}(u_{t}-\tilde{u}_{t})\|_{L^{2}} \Big] d\tau \\ &\leq CR\|u-\tilde{u}\|_{X} \int_{t/2}^{t} (1+t-\tau)^{-\frac{n}{4}-\frac{1}{2}} (1+\tau)^{-\frac{k}{2}} d\tau \\ &\leq CR\|u-\tilde{u}\|_{X} (1+t)^{-\frac{k}{2}}. \end{aligned}$$

Finally, we estimate the term J_{22} as

$$\begin{aligned} J_{22} &\leq \int_{t/2}^{t} e^{-c(t-\tau)} \Big[(\|\partial_{x}^{k+2}u\|_{L^{2}} + \|\partial_{x}^{k+2}\tilde{u}\|_{L^{2}}) \|u - \tilde{u}\|_{L^{2}} \\ &+ (\|u\|_{L^{2}} + \|\tilde{u}\|_{L^{2}}) \|\partial_{x}^{k}(u - \tilde{u})\|_{L^{2}} + (\|\partial_{x}^{k}u_{t}\|_{L^{2}} + \|\partial_{x}^{k}\tilde{u}_{t}\|_{L^{2}}) \|u_{t} - \tilde{u}_{t}\|_{L^{2}} \\ &+ (\|u_{t}\|_{L^{2}} + \|\tilde{u}_{t}\|_{L^{2}}) \|\partial_{x}^{k}(u_{t} - \tilde{u}_{t})\|_{L^{2}} \Big] d\tau \\ &\leq CR \|u - \tilde{u}\|_{X} \int_{t/2}^{t} e^{-c(t-\tau)} (1+\tau)^{-\frac{n}{4} - \frac{k+2}{2}} d\tau \\ &\leq CR \|u - \tilde{u}\|_{X} (1+\tau)^{\frac{k}{2}}. \end{aligned}$$

Thus we have shown that

$$(1+t)^{\frac{k}{2}} \|\partial_x^k(\Phi(u) - \Phi(\tilde{u}))\|_{L^2} \le CR \|u - \tilde{u}\|_X.$$
(4.6)

Suppose that $u, \tilde{u} \in X_R$. It follows from (3.3) that

$$\begin{split} \|\partial_x^k(\varPhi(u) - \varPhi(\tilde{u})_t\|_{L^2} \\ &= \int_0^t \|\partial_x^k G_t(t-\tau) * \varDelta(f(u) - f(\tilde{u}) - \beta(g(u_t)) - g(\tilde{u}_t))(\tau)\|_{L^2} d\tau \\ &= \int_0^{t/2} \|\partial_x^k G_t(t-\tau) * \varDelta(f(u) - f(\tilde{u}) - \beta(g(u_t)) - g(\tilde{u}_t))(\tau)\|_{L^2} d\tau \\ &+ \int_{t/2}^t \|\partial_x^k G_t(t-\tau) * \varDelta(f(u) - f(\tilde{u}) - \beta(g(u_t)) - g(\tilde{u}_t))(\tau)\|_{L^2} d\tau \\ &=: j_1 + j_2. \end{split}$$

By using (2.24) with p = 1, j = 0, we have

$$\begin{split} \dot{J_1} &\leq C \int_0^{t/2} (1+t-\tau)^{-\frac{n}{4} - \frac{k+2}{2}} \|f(u) - f(\tilde{u})(\tau) - \beta(g(u_t) - g(\tilde{u}_t))(\tau)\|_{L^1} d\tau \\ &+ C \int_0^{t/2} e^{-c(t-\tau)} \|\partial_x^{k+2}(f(u) - f(\tilde{u}_t) - \beta(g(u_t) - g(\tilde{u}_t))(\tau)\|_{L^2} d\tau \\ &=: \ J_{11} + J_{12}. \end{split}$$

By (3.2), we obtain

$$\begin{aligned}
\dot{J_{11}} &\leq \int_{0}^{t/2} (1+t-\tau)^{-\frac{n}{4}-\frac{k+2}{2}} (\|u\|_{L^{2}}+\|\tilde{u}_{L^{2}})\|u-\tilde{u}\|_{L^{2}} \\
&\times (\|u_{t}\|_{L^{2}}+\|\tilde{u}_{t}\|_{L^{2}})(\|u_{t}-\tilde{u}_{t}\|_{L^{2}})d\tau \\
&\leq CR\|u-\tilde{u}\|_{X}(1+t)^{-\frac{k+1}{2}} \int_{0}^{t/2} (1+t\tau)^{-\frac{n}{4}-\frac{1}{2}}d\tau \\
&\leq CR\|u-\tilde{u}\|_{X}(1+t)^{-\frac{k+1}{2}}.
\end{aligned}$$

For the term J_{12} , by (3.2) we get

$$\begin{split} J_{12}^{'} &\leq \int_{0}^{t/2} e^{-c(t-\tau)} \left[(\|\partial_{x}^{k+2}u\|_{L^{2}} + \|\partial_{x}^{k+2}\tilde{u}\|_{L^{2}}) \|u-\tilde{u}\|_{L^{\infty}} \right. \\ &\quad + (\|u\|_{L^{\infty}} + \|\tilde{u}\|_{L^{\infty}})) \|\partial_{x}^{k+2}(u-\tilde{u})\|_{L^{2}} \\ &\quad + (\|\partial_{x}^{k+2}u_{t}\|_{L^{2}} + \|\partial_{x}^{k+2}\tilde{u}_{t}\|_{L^{2}}) \|u_{t} - \tilde{u}_{t}\|_{L^{\infty}} \\ &\quad + (\|u_{t}\|_{L^{\infty}} + \|\tilde{u}_{t}\|_{L^{\infty}}) \|\partial_{x}^{k+2}(u_{t} - \tilde{u}_{t})\|_{L^{2}} \right] d\tau \\ &\leq CR \|u-\tilde{u}\|_{X} \int_{0}^{t/2} e^{-c(t-\tau)} (1+\tau)^{-\frac{k+2}{2}} (1+\tau)^{-\frac{n}{4}} d\tau \\ &\leq CR \|u-\tilde{u}\|_{X} e^{-ct}. \end{split}$$

By applying (2.24) with p = 2, j = k + 2 and l = 0, we conclude that

$$\begin{split} \dot{J_2} &\leq C \int_{t/2}^t \|\partial_x^{k+2}(f(u) - f(\tilde{u}) - \beta(g(u_t) - g(\tilde{u}_t)(\tau)))\|_{L^2} d\tau \\ &\quad + C \int_{t/2}^t e^{-c(t-\tau)} \|\partial_x^{k+2}(f(u) - f(\tilde{u}) - \beta(g(u_t) - g(\tilde{u}_t)(\tau)))\|_{L^2} d\tau \\ &\leq \int_{t/2}^t \left[(\|\partial_x^{k+2}u\|_{L^2} + \|\partial_x^{k+2}\tilde{u}\|_{L^2})\|u - \tilde{u}\|_{L^\infty} \\ &\quad + (\|u\|_{L^\infty} + \|\tilde{u}\|_{L^\infty})\|\partial_x^{k+2}(u - \tilde{u})\|_{L^2} \\ &\quad + (\|\partial_x^{k+2}u_t\|_{L^2} + \|\partial_x^{k+2}\tilde{u}_t\|_{L^2})\|u - \tilde{u}_t\|_{L^\infty} \\ &\quad + (\|u_t\|_{L^\infty} + \|\tilde{u}_t\|_{L^\infty})\|\partial_x^{k+2}(u - \tilde{u}_t)\|_{L^2} \right] d\tau \\ &\leq CR\|u - \tilde{u}\|_X \int_{t/2}^t (1+\tau)^{-\frac{n}{4} - \frac{k+2}{2}} d\tau \\ &\leq CR\|u - \tilde{u}\|_X (1+t)^{-\frac{k+1}{2}} \int_{t/2}^t (1+\tau)^{-\frac{n}{4} - \frac{1}{2}} d\tau \\ &\leq CR\|u - \tilde{u}\|_X (1+t)^{-\frac{k+1}{2}}. \end{split}$$

Consequently, we have shown that

$$(1+t)^{\frac{k+1}{2}} \| \partial_x^k (\Phi(u) - \Phi(\tilde{u}))_t \|_X \le CR \|u - \tilde{u}\|_X.$$
(4.7)

Using (4.5) and (4.7) and taking R suitably small, it yields

$$\|\Phi(u) - \Phi(\tilde{u})\|_X \le \frac{1}{2} \|u - \tilde{u}\|_X.$$
(4.8)

From (4.8), we conclude that Φ is a contracting mapping. Then there exists a fixed point $u \in X_R$ of mapping Φ , which is a solution (1.1) and (1.2) and the proof is completed. \Box

Finally we study the asymptotic linear profile of the solution.

Suppose that u_L given by the formula $u_L(t) = G(t) * u_1 + H(t) * u_0$. In the previous two section, we have shown that the solution u to problem (1.1) and (1.2) can be approximated by the linear solution u_L . Now the aim is to derive a simpler asymptotic profile of the linear solution u_L .

In the Fourier space, we obtain $\hat{u}_L(\xi,t) = \hat{G}(\xi,t)\hat{u}_1(\xi,t) + \hat{H}(\xi,t)\hat{u}_0(\xi)$, where $\hat{G}(\xi,t)$ and $\hat{H}(\xi,t)$ are given explicitly in (2.6) and (2.7). First we give the asymptotic expansions of $\hat{G}(\xi,t)$ and $\hat{H}(\xi,t)$ for $\xi \to 0$. By using the Taylor expansion to (2.4), we obtain

$$\lambda_{\pm}(\xi) = \frac{1}{2}(\alpha|\xi|^2 - |\xi|^4) \pm \frac{|\xi|i}{2}(2 + |\xi|^2 - \frac{\alpha^2}{4}|\xi|^4 + O(|\xi|^4)$$
$$= \pm i|\xi| + \frac{\alpha}{2}|\xi|^2 + O(|\xi|^3)$$

and

$$\begin{aligned} \frac{1}{\lambda_{+} - \lambda_{-}} &= \frac{1}{i|\xi|\sqrt{4 + 4|\xi|^{2} - |\xi|^{6} - \alpha^{2}|\xi|^{4} + 2\alpha|\xi|^{5}}} \\ &= \frac{1}{2i|\xi|}(1 - \frac{1}{2}|\xi|^{2} + O(|\xi|^{4})). \end{aligned}$$

Substituting these expansions to (2.6) and (2.7), we obtain

$$\begin{aligned} \hat{G}(\xi,t) &= \frac{e^{\lambda_{+}t} - e^{\lambda_{-}t}}{\lambda_{+} - \lambda_{-}} \\ &= \frac{1}{2i|\xi|} \left(e^{\frac{\alpha}{2}|\xi|^{2}t} (e^{|\xi|ti} - e^{-|\xi|ti}) + e^{\frac{\alpha}{2}|\xi|^{2}t} (O(|\xi|^{2}) + O(|\xi|^{3}t)) \right) \end{aligned}$$

and

$$\begin{aligned} \hat{H}(\xi,t) &= \frac{\lambda_{+}e^{\lambda_{-}t} - \lambda_{-}e^{\lambda_{+}t}}{\lambda_{+} - \lambda_{-}} \\ &= \frac{1}{2}e^{\frac{\alpha}{2}|\xi|^{2}t}(e^{|\xi|ti} + e^{-|\xi|ti}) + e^{\frac{\alpha}{2}|\xi|^{2}t}(O(|\xi|^{2}) + O(|\xi|^{3}t)) \end{aligned}$$

for $\xi \to 0$. Let

$$\hat{G}_0(\xi,t) = \frac{1}{2i|\xi|} e^{\frac{\alpha}{2}|\xi|^2 t} (e^{i|\xi|t} - e^{-i|\xi|t}), \quad \hat{H}_0(\xi,t) = \frac{1}{2} e^{\frac{\alpha}{2}|\xi|^2 t} (e^{i|\xi|t} + e^{-i|\xi|t}).$$

Thus for $|\xi| \leq r_0$ we obtain

$$|(\hat{G} - \hat{G}_0)(\xi, t)| \le Ce^{-c|\xi|^2 t}, \quad |(\hat{H} - \hat{H}_0)(\xi, t)| \le C|\xi|e^{-c|\xi|^2 t},$$

where r_0 is a small positive constant. Now we define \overline{u}_L by

$$\overline{u}_L(t) = G_0(t) * u_1 + H_0(t) * u_0.$$
(4.9)

 \overline{u}_L gives an asymptotic profile of the linear solution u_L .

Theorem 3. Suppose that $n \geq 1$, $s \geq 0$ and $u_0 \in H^{s+2} \cap L^1$ and $u_1 \in H^s \cap \dot{W}^{-1,1}$. Put $E_0 = ||u_0||_{L^1} + ||u_1||_{\dot{W}^{-1,1}} + ||u_0||_{H^{s+2}} + ||u_1||_{H^s}$. Let u_L be the linear solution and \overline{u}_L be defined by (4.9). Thus we have

$$\|\partial_x^k (u_L - \overline{u}_L)(t)\|_{L^2} \le C E_0 (1+t)^{-\frac{n}{4} - \frac{k+1}{2}}$$

for $0 \le k \le s+2$.

Proof. It follows from definition that

$$(u_L - \overline{u}_L)(t) = (G - G_0)(t) * u_1 + (H - H_0)(t) * u_0.$$

So it suffices to show the following estimates:

$$\begin{aligned} \|\partial_x^k (G - G_0)(t) * u_1\|_{L^2} &\leq C(1+t)^{-\frac{n}{2}(\frac{1}{p} - \frac{1}{2}) - \frac{k+1-j}{2}} \|\partial_x^j u_1\|_{\dot{W}^{-1,p}} \\ &+ Ce^{-ct} \|\partial_x^{k+l-2} u_1\|_{L^2}, \\ \|\partial_x^k (H - H_0)(t) * u_0\|_{L^2} &\leq C(1+t)^{-\frac{n}{2}(\frac{1}{p} - \frac{1}{2}) - \frac{k+1-j}{2}} \|\partial_x^j u_0\|_{L^p} \\ &+ Ce^{-ct} \|\partial_x^{k+l} \phi\|_{L^2}, \end{aligned}$$

where $1 \le p \le 2$, and k, j and l are nonnegative integers such that $0 \le j \le k+1$. We assumed $k + l - 2 \ge 0$ in the first estimate. These estimates can be proved similarly as in the proof of Lemma 5 by using (4.1) for $|\xi| \le r_0$ and (2.17) and (2.17) and (4.4) for $|\xi| \ge r_0$. We omit the details. \Box

5 Conclusions

In the body of the paper, we have considered a dissipative Boussinesq-type equation. This equation (see Equation (1.1)) arises in the study of the stability of one-dimensional periodic patterns in systems with Galilean invariance and also the oscillations of elastic beams. Here we have shown that the Cauchy problem associated to this equation has a unique global solution in some suitable Sobolev space. The main difficulty was how to control the nonlinear term of the equation which is a combination of u and u_t . Under small condition on the initial value, we proves the existence by the contraction mapping principle due to small nonlinear terms. We also obtained some asymptotic behavior of global solutions in some time weighted spaces. In the presence of the dissipation and small nonlinearities, we established that the solution the nonlinear initial value problem behaves like the linear solution, and thus it can be approximated by the linear solution.

In the future, we plan to study this equation with more general nonlinearity. The sharpness of the decay estimates and low regularity of the solutions will be also very important and interesting issues.

Acknowledgements

The authors wish to thank the unknown referees for their careful reading.

References

- D.J. Benney. Long waves on liquid films. Journal of Mathematics and Physics, 45(1-4):150–155, 1966. https://doi.org/10.1002/sapm1966451150.
- [2] J. Boussinesq. Théorie des ondes et de remous qui se propagent le long d'un canal rectangulaire horizontal, en communiquant au liquide contene dans ce cannal des vitesses sensiblement pareilles de la surface au fond. Journal de Mathématiques Pures et Appliquées, 17:55–108, 1872.
- [3] C.I. Christov, G.A. Maugin and A.V. Porubov. On Boussinesq's paradigm in nonlinear wave propagation. *Comptes Rendus Mécanique*, **335**(9-10):521–535, 2007. https://doi.org/10.1016/j.crme.2007.08.006.
- [4] C.I. Christov and M.G. Velarde. Dissipative solitons. *Physica D: Nonlinear Phenomena*, 86(1-2):323–347, 1995. https://doi.org/10.1016/0167-2789(95)00111-G.
- [5] B.I. Cohen, J.A. Krommes, W.M. Tang and M.N. Rosenbluth. Non-linear saturation of the dissipative trapped-ion mode by mode coupling. *Nuclear Fusion*, 16(6):971–992, 1976.
- [6] A. Esfahani. Instability of the stationary solutions of generalized dissipative Boussinesq equation. Applications of Mathematics, 59(3):345–358, 2014. https://doi.org/10.1007/s10492-014-0059-1.
- [7] A.N. Garazo and M.G. Velarde. Dissipative Korteweg-de Vries description of Marangoni-Bénard oscillatory convection. *Physics of Fluids A: Fluid Dynamics*, 3(10):2295–2300, 1991. https://doi.org/10.1063/1.857868.
- [8] N. Hayashi, E.I. Kaikina, P. Naumkin and I.A. Shishmarev. Asymptotics for dissipative nonlinear equations. Springer-Verlag, Berlin, Heidelberg, 2006.
- [9] B. Janiaud, A. Pumir, D. Bensimon, V. Croquette, H. Richter and L. Kramer. The Eckhaus instability for travelling waves. *Physica D: Nonlinear Phenomena*, 55(3-4):269–286, 1992. https://doi.org/10.1016/0167-2789(92)90060-Z.
- [10] R.A. Kraenkel, S.M. Kurcbart, J.G. Pereira and M.A. Manna. Dissipative Boussinesq system of equations in the Bénard-Marangoni phenomenon. *Physical Review E*, 49(2):1759–1762, 1994. https://doi.org/10.1103/PhysRevE.49.1759.
- [11] M. Liu and W. Wang. Global existence and pointwise estimates of solutions for the multidimensional generalized Boussinesq-type equation. *Communications on Pure and Applied Analysis*, **13**(3):1203–1222, 2014. https://doi.org/10.3934/cpaa.2014.13.1203.
- [12] J. Mason and E. Knobloch. Long dynamo waves. *Physica D: Nonlinear Phenom*ena, **205**(1-4):100–124, 2005. https://doi.org/10.1016/j.physd.2005.01.006.
- [13] A.A. Nepomnyashchy and M.G. Velarde. A three dimensional description of solitary waves and their interaction in Marangoni-Bénard layers. *Physics of Fluids*, 6(1):187–198, 1994. https://doi.org/10.1063/1.868081.
- [14] G.I. Sivashinsky. Nonlinear analysis of hydrodynamic instability in laminar flames-I. derivation of basic equations. Acta Astronautica, 4(11-12):1177-1206, 1977. https://doi.org/10.1016/0094-5765(77)90096-0.
- [15] C.D. Sogge. Lectures on non-linear wave equations, Second edition. International Press, Boston, MA, 2008.
- [16] J. Topper and T. Kawahara. Approximate equations for long nonlinear waves on a viscous fluid. *Journal of the Physical Society of Japan*, 44(2):663–666, 1978. https://doi.org/10.1143/JPSJ.44.663.

- [17] Y. Wang. Asymptotic decay estimate of solutions to the generalized damped Bq equation. Journal of Inequalities and Applications, 2013(1):323, 2013. https://doi.org/10.1186/1029-242X-2013-323.
- [18] Y. Wang. On the Cauchy problem for one dimension generalized Boussinesq equation. International Journal of Mathematics, 26(03), 2015. https://doi.org/10.1142/S0129167X15500238.
- [19] T. Yamada and Y. Kuramoto. A reduced model showing chemical turbulence. Progress of Theoretical Physics, 56(2):681–683, 1976. https://doi.org/10.1143/PTP.56.681.
- [20] S. Zheng. Nonlinear evolution equations. Chapman & Hall/CRC, 2004.