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# CMS Conference Report

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## Determination of the Discovery Potential for Higgs Bosons in MSSM

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### Abstract

The CMS and ATLAS collaborations have performed detailed studies of the discovery potential for the Higgs boson in MSSM. Different benchmark scenarios have been studied both for the CP-conserving case and for the CP-violating one. Results are presented of the discovery potential in the parameter space of MSSM.

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# Determination of the Discovery Potential for Higgs Bosons in MSSM

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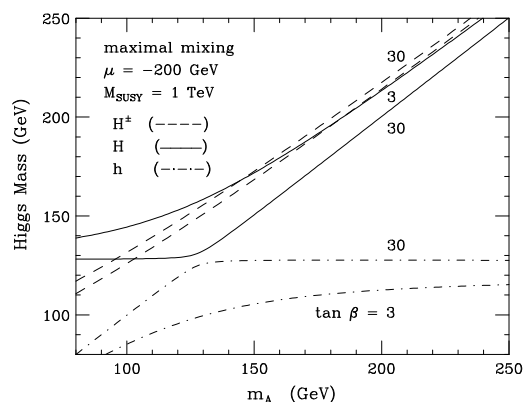
**Abstract.** The CMS and ATLAS collaborations have performed detailed studies of the discovery potential for the Higgs boson in MSSM. Different benchmark scenarios have been studied both for the CP-conserving case and for the CP-violating one. Results are presented of the discovery potential in the parameter space of MSSM.

**PACS.** 11.30.Pb Supersymmetry – 14.80.Cp Non-standard-model Higgs bosons

## 1 The MSSM Higgs Sector

Supersymmetry (SUSY) at the TeV scale provides an elegant solution to the hierarchy problem through introduction of superpartners to SM particles and cancellation of problematic loop corrections [1]. It allows light Higgs bosons in the context of GUT without fine tuning. In the Minimal Supersymmetric (MSSM) extension of the Standard Model (SM) two isospin Higgs boson doublets are introduced in order to preserve supersymmetry, one of which couples to the down-type fermions and the other one to the up-type. Eight degrees of freedom therefore exists, three of which are absorbed by the Z and W gauge bosons after the electroweak symmetry breaking. This leads to the existence of five elementary Higgs bosons in MSSM. These physical Higgs bosons are the two CP-even neutral scalar particles  $h$ ,  $H$ , one CP-odd neutral pseudoscalar particle  $A$ , and two charged particles  $H^\pm$ . Four masses of the elementary Higgs bosons are used to describe the MSSM Higgs sector:  $M_h$ ,  $M_H$ ,  $M_A$ ,  $M_{H^\pm}$ . In addition two parameters are needed which describe the properties of the scalar particles and their interactions with gauge bosons and fermions. The first of them is the mixing angle  $\beta$  which is related to the ratio of the two vacuum expectation values (vev) of the Higgs boson doublets:  $\tan\beta = v_2/v_1$ . The second parameter is the mixing angle,  $\alpha$ , in the neutral CP-even sector.

Several relations exist between the MSSM parameters introduced above. At tree-level only two of the parameters are independent and the others can be calculated in terms of them. In the case of CP-conservation the two parameters chosen to describe MSSM are usually  $M_A$ ,  $\tan\beta$ . The following hierarchies hold at tree-level too:  $M_h < M_Z$ ,  $M_A < M_H$  and  $M_W < M_{H^\pm}$ . Such limits would lead to a light scalar Higgs boson in a mass range already excluded by the LEP experiments. This tree-level bound on  $M_h$  receives large ra-



**Fig. 1.** The light, heavy and charged CP-even MSSM Higgs boson masses as a function of  $M_A$  for  $\tan\beta = 3$  and 30, including radiative corrections. The maximal mixing scenario is used with Higgs boson mass parameter  $\mu = -200$  GeV and SUSY mass scale  $M_{\text{SUSY}} = 1$  TeV [2].

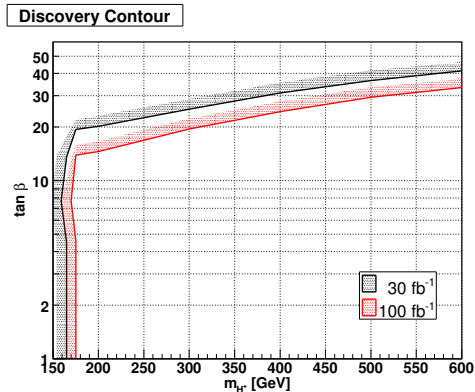
diative corrections from SUSY breaking effects in the Yukawa sector of the theory. The leading order corrections are of proportional to  $M_t^4$ , where  $M_t$  is the top quark mass. The upper mass bound for light scalar Higgs boson is reached at large  $M_A$  where  $h$  becomes SM like and  $M_h < 135$  GeV. The Higgs boson masses including the radiative corrections are shown in figure 1 [2].

## 2 Benchmark Scenarios

Beyond tree-level, the main corrections for Higgs boson masses stem from the  $t/\tilde{t}$  and  $b/\tilde{b}$  sector (the latter is important for large  $\tan\beta$ ). Sub-leading corrections come from all other sectors of MSSM. The Higgs sector phenomenology is thus connected via radiative corrections to the full spectrum on MSSM. In the unconstrained version of MSSM no particular SUSY breaking mechanism is assumed. A parametrization of all

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**Fig. 5.** Charged Higgs boson discovery contour in the MSSM. The regions above the curves are the part of the parameter space in which a  $5\sigma$ -discovery is feasible. Curves for two different integrated luminosities of simulated data with the ATLAS detector are shown.

for the light, neutral Higgs boson  $h$  from the inclusive  $pp \rightarrow h + X$  production with the  $h \rightarrow \gamma\gamma$  decay and for light and heavy scalar Higgs bosons,  $h$  and  $H$ , produced in the vector boson fusion (VBF)  $qq \rightarrow qqh(H)$  with  $h(H) \rightarrow \tau\tau \rightarrow l + j$  decay is shown in figure 3. The VBF channel dominates the discovery potential at low luminosity and covers most of the parameter space left over from the LEP exclusion limits. The discovery potential for the light neutral  $h$  for ATLAS  $30 \text{ fb}^{-1}$  simulated data is shown in figure 4 in the small  $\alpha$  scenario. The effect of suppressed branching ratios into  $\tau$  leptons is visible for  $\tan\beta > 20$  and  $20 < M_A < 300 \text{ GeV}$ . This hole in the discovery region is nicely complemented by  $h$  decays into gauge bosons from VBF or gluon-gluon fusion. Differences between the different scenarios are mainly due to the fact that in the same parameter space point the  $h$  mass differs, changing the sensitivity of the channels under consideration.

### 3.2 Charged Higgs Bosons

Since no charged scalar Higgs boson is predicted within the SM, search for such a Higgs at LHC is particularly interesting. At the LHC (and hadron colliders in general) single charged Higgs boson is produced through two main mechanisms:  $gg \rightarrow tH^\pm$  ( $2 \rightarrow 3$  process) and  $gb \rightarrow tH^\pm$  ( $2 \rightarrow 2$  process). These two processes are called twin processes [7] since they correspond to two different approximations describing the same basic process. The  $2 \rightarrow 2$  process is dominant at higher masses,  $M_{H^\pm} > M_t$ , due to resummation of potentially large logarithms by the  $b$  quark parton density. In this case the parton shower produces the outgoing  $b$  quark. For Higgs boson masses below the mass of the top,  $M_{H^\pm} < M_t$ , the  $2 \rightarrow 3$  process dominates since it incorporates on-shell top quark pair production with subsequent decay to charged Higgs boson. For charged Higgs boson masses range around the mass of the top quark both processes give comparable contributions. Theoretical treatment of this region is complicated

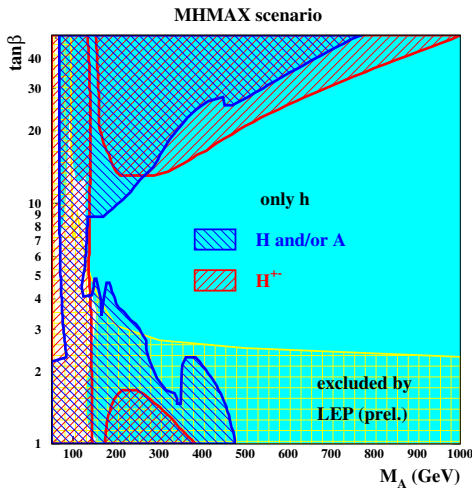
and a matching procedure must be applied in order to avoid double counting. The ATLAS collaboration has performed an analysis of simulated data for two different luminosities,  $30$  and  $300 \text{ fb}^{-1}$ . For the generation of events the MATCHIG [8] program was used that performs a matching of the above two processes allowing proper handling of the predictions and a consistent treatment of the transition region around the top quark mass. The decay channels used are  $H \rightarrow \tau\nu_\tau$ ,  $t \rightarrow bj\bar{j}$ . Results of the discovery contours obtained for the two luminosities are presented in figure 5. An improvement with respect to previous measurements is observed. With  $30 \text{ fb}^{-1}$  a discovery of charged Higgs boson with mass below  $160 \text{ GeV}$  is possible for all  $\tan\beta$ , whereas a discovery over the whole mass region is possible for  $\tan\beta > 40$ .

### 3.3 Overall CPC Discovery Potential

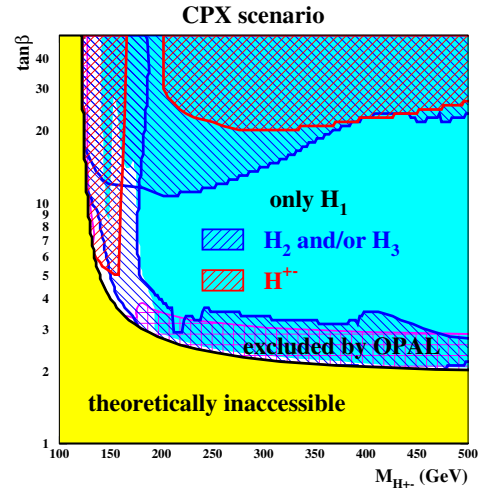
The overall discovery potential for Higgs boson in the  $M_h^{\text{max}}$  scenario with  $300 \text{ fb}^{-1}$  simulated data of the ATLAS detector is shown in figure 6. For all benchmark scenarios in a large part of the MSSM parameter space discovery is possible via several channels, which will allow a determination of parameters of the the MSSM Higgs sector. At least one Higgs boson can be discovered in the whole model parameter space and for a significant part of the parameter space more than one Higgs boson can be observed. This would allow to distinguish between SM and MSSM through direct observation. However in a large area only the observation of the light neutral Higgs boson  $h$  is possible. In order to discriminate in this case between a SM or MSSM origin of  $h$  the ratio of branching ratios ( $\text{BR}$ ) was calculated for the VBF production  $R = \text{BR}(h \rightarrow \tau\tau)/\text{BR}(h \rightarrow WW)$ . Then a discrimination variable is defined:  $\Delta = (R_{\text{MSSM}} - R_{\text{SM}})/\sigma_{\text{exp}}$ , where  $\sigma_{\text{exp}}$  denotes the expected error on  $R$ . Only statistical uncertainties have been taken into account. The sensitivity of this discrimination variable is shown in figure 7. Other studies have been performed for extracting the Higgs boson couplings from the LHC data and testing the sensitivity of the deviations from the SM [9]. They are based on global fits of the data.

### 3.4 CPX Scenario

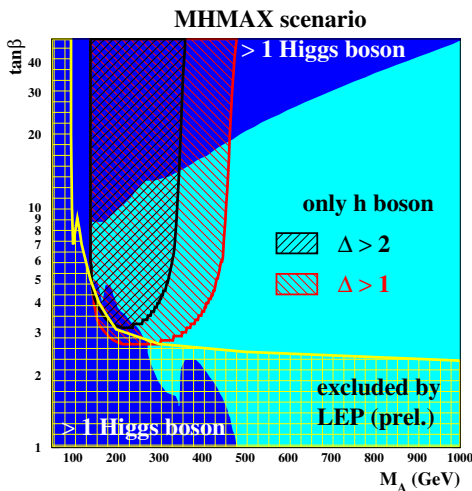
The overall discovery potential for neutral Higgs bosons in the CPX scenario is shown in figure 8 for a sample of simulated data of the ATLAS detector with a luminosity of  $300 \text{ fb}^{-1}$ . A coverage similar to the case of CPC scenarios is observed and in most of the parameter space at least one Higgs boson can be observed. It must be noted that the LEP exclusion limits for the CPX scenario are weaker and even low masses  $0 < M_{H_i} < 60 \text{ GeV}$  are not yet excluded [10]. A small region of low  $M_{H_i}$  and small  $\tan\beta$  remains where no discovery is possible for the channels and mass ranges investigated in the present studies. The mass of the three CP mass eigenstates  $H_1, H_2, H_3$  is respectively



**Fig. 6.** Overall discovery potential for Higgs boson in the  $M_h^{\max}$  scenario with  $300 \text{ fb}^{-1}$  as calculated with simulated data for the ATLAS detector.



**Fig. 8.** Overall discovery potential for Higgs bosons in the CPX scenario with  $300 \text{ fb}^{-1}$  of simulated data from the ATLAS detector.



**Fig. 7.** Sensitivity for discrimination between SM and MSSM origin of Higgs boson in the  $M_h^{\max}$  scenario with simulated data for the ATLAS detector. For more details see text.

105, 120 and 140 to 180 GeV. The ATLAS collaboration has started further studies in this particular region of the phase space.

## 4 Conclusions

The discovery potential of ATLAS and CMS Higgs boson searches based on the most recent theoretical calculations and analysis of simulated data has been discussed for four different CPC scenarios and the CP-violating CPX scenario. For the CPC MSSM Higgs boson all the parameter space is potentially covered with already  $30 \text{ fb}^{-1}$  luminosity of experimental data by discovery of at least one Higgs boson. The coverage observed on all four benchmark scenarios reflects probably most of the MSSM phase space. In large regions of

the phase space only the SM-like MSSM Higgs boson might be observed. Work is ongoing on strategies to distinguish between SM and MSSM origin of the Higgs boson in this case. For the CPV MSSM almost all parameter space is covered by observation of at least one Higgs boson in the CPX scenario. Studies are underway on small uncovered space ( $M_{H_1} < 50 \text{ GeV}$ ) not excluded by LEP searches.

## 5 Acknowledgements

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