



GAP IN A SPIRALLY RIDGED POLE

D. W. Kerst

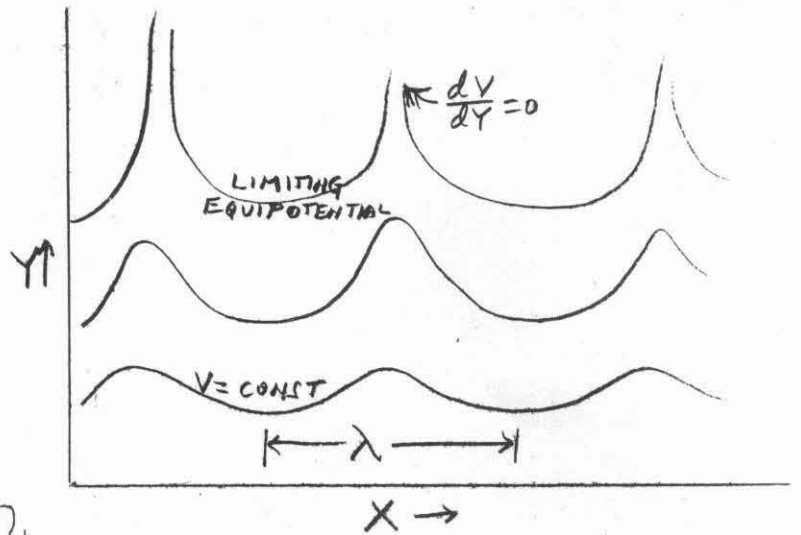
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This is a calculation similar to that made by Terwilliger in MURA/KMT-2 but for a sinusoidal distribution of magnetic field in the median plane instead of for a step distribution of field.

Terwilliger has pointed out that for $B = B_0 \gamma^k (1 + f \sin \frac{x}{\lambda})$, the magnetic potential for the case of $k = 0$ is

$$V = B_0 \left\{ Y + f \lambda \sin \frac{x}{\lambda} \sinh \frac{Y}{\lambda} \right\}$$

We want to calculate the maximum gap between opposing pole ridge tops, G , for different values of the fractional field variation, f , in the median plane. This maximum G is the separation between ridge tops for the limiting magnetic equipotential which develops an infinite crevice between ridges.



The crevices occur where $\frac{dV}{dY} = 0$ or $0 = 1 + f \sin \frac{x}{\lambda} \cosh \left(\frac{Y}{\lambda} \right)$ and the value of $\frac{x}{\lambda}$ at crevices is $\frac{x}{\lambda} = -\frac{\pi}{2}$. Thus $\frac{1}{f} = \cosh \left(\frac{Y_{\infty}}{\lambda} \right)$. For this equipotential $V_{\infty} = B_0 \left\{ \lambda \cosh^{-1} \frac{1}{f} - \lambda f \sqrt{\left(\frac{1}{f^2} - 1 \right)} \right\}$. On this same equipotential surface

$Y_{\min} = G/2$ occurs at $\sin x / \lambda = 1$ and it satisfies

$$V_{\infty} = B_0 \left\{ Y_{\min} + f \lambda \sin H \left(\frac{Y_{\min}}{\lambda} \right) \right\}$$

equating

$$\frac{Y_{\min}}{\lambda} + f \sin H \left(\frac{Y_{\min}}{\lambda} \right) = \cos H^{-1} \frac{1}{f} \sqrt{(1+f)(1-f)}.$$

If we now impose the condition that the smooth approxi-

mation $\nu_z^2 = \frac{f^2 r^2}{\lambda^2 N^2} = 2A.G.$ be a constant, then, since

in cases where $k \neq 0$ N is already chosen by radial motion

considerations, we are imposing the condition $\frac{f}{\lambda} = \text{constant}$.

This is also requiring that $\left. \frac{dB}{dx} \right|_{\max} = \text{constant}$ in the orbital

plane for different parameters tried.

The graphs show $\frac{G}{\lambda} = 2 \frac{Y_{\min}}{\lambda \left(\frac{r}{\lambda} \pi \right)}$ determined by the transcendental equation above and $\frac{G}{\lambda}$ which is proportional to the maximum gap attainable for $\frac{f}{\lambda} = \text{constant}$, $\nu_z = \text{constant}$, or $\left(\frac{Y}{H} \frac{dH}{dr} \right)_{\text{Peak}} = \text{constant}$.

The curves show that $f = \frac{1}{4}$ gives about the maximum gap if one is designing for a fixed ν_z , but only 10% of the gap is lost if $f = .35$ or $.15$ is used. The biggest usable value of G/λ is $.28$ which is a good design rule. This result is very closely the same as that given by the Illiac for the solution to the Cauchy problem with $k \sim 150$, that is for $B = B_0 (r/r_0)^k (1 + f \sin \frac{x}{\lambda})$ in the orbital plane. Terwilliger found a broad maximum at $f \doteq \frac{1}{3}$ for step shaped field variations.

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MURA/DWK IO

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PLOT OF MINIMUM GAP IN SPIRAL RIDGED
POLES (RIDGES SPACED λ) AS A FUNCTION
OF THE FLUTTER f .
 $V = H_0 (Y + f \lambda \sin X Y \lambda \sin H Y \lambda)$ IS POTENTIAL
FUNCTION.

$$B = B_0 (1 + f \sin X Y \lambda)$$

FIGURE II

