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## Search for $\eta_b$ in two-photon collisions at LEP II with the DELPHI detector

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### Abstract

The pseudoscalar meson  $\eta_b$  has been searched for in two-photon interactions at LEP II. The data sample corresponds to a total integrated luminosity of  $617 \text{ pb}^{-1}$  at centre-of-mass energies ranging from 161 to 209 GeV. Upper limits at a confidence level of 95% on the product  $\Gamma_{\gamma\gamma}(\eta_b) \times \text{BR}(\eta_b)$  are 190, 470 and  $660 \text{ eV}/c^2$  for the  $\eta_b$  decaying into 4, 6 and 8 charged particles, respectively.

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# 1 Introduction

Two-photon collisions are very useful in searching for the formation of pseudoscalar mesons with  $J^{PC} = 0^{-+}$ . The high energy and high luminosity of LEP II are additional motivations to look for the  $b\bar{b}$  pseudoscalar quarkonium state  $\eta_b$  which has not yet been discovered [1, 2].

Its mass,  $m_{\eta_b}$ , is estimated by several theoretical models [3]. It should lie below that of the  $\Upsilon$  vector meson ( $m_\Upsilon=9.46$  GeV/ $c^2$ ) and the mass shift,  $\Delta m = m_\Upsilon - m_{\eta_b}$ , is estimated to be in the range 10 to 130 MeV/ $c^2$ .

The cross-section for two-photon resonance  $R$  formation with  $C=+1$

$$e^+e^- \rightarrow e^+e^-\gamma^*\gamma^* \rightarrow e^+e^-R$$

is given by [4]

$$\sigma(e^+e^- \rightarrow e^+e^-R) = \int \sigma_{\gamma\gamma \rightarrow \eta_b} dL_{\gamma\gamma}(W^2),$$

with the cross-section

$$\sigma_{\gamma\gamma \rightarrow \eta_b}(W^2, q_1^2, q_2^2) = 8\pi (2J_R + 1) \cdot \Gamma_{\gamma\gamma}(R) \cdot F^2(q_1^2, q_2^2) \cdot \frac{\Gamma_R}{(W^2 - m_R^2)^2 + m_R^2 \Gamma_R^2}.$$

Here  $L_{\gamma\gamma}(W^2)$  is the two-photon luminosity function,  $W$  is the two-photon centre-of-mass energy,  $q_1^2$  and  $q_2^2$  are the squares of the virtual-photon four-momenta. The resonance  $R$  is characterised by its spin  $J_R$ , mass  $m_R$ , total width  $\Gamma_R$  and its two-photon partial width  $\Gamma_{\gamma\gamma}(R)$ . In ‘‘quasi-real’’ ( $q^2 \sim 0$ ) photon interactions, the form factor  $F^2(q_1^2, q_2^2)$  is constant and can be taken to be unity.

To compute the  $\eta_b$  production cross-section, the partial width  $\Gamma_{\gamma\gamma}(\eta_b)$  must be known. Theoretical estimates [5] predict it to be in the range 260 to 580 eV/ $c^2$ . Setting  $m_{\eta_b}$  to 9.4 GeV/ $c^2$  leads to an expected production cross-section  $\sigma(e^+e^- \rightarrow e^+e^-\eta_b)$  of 0.14 to 0.32 pb at  $\sqrt{s}=200$  GeV.

Most of the observations of  $\eta_c$  decays have been to four charged particles, both pions and kaons [6]. Hence the  $\eta_b$  has been similarly searched for in 4, 6 and 8 charged particle final states. The expected backgrounds come from the  $\gamma\gamma \rightarrow q\bar{q}$  processes and the  $\gamma\gamma \rightarrow \tau^+\tau^-$  channel.

From the ALEPH experiment, upper limits on  $\Gamma_{\gamma\gamma}(\eta_b) \times \text{BR}(\eta_b)$  [1] are :

$$\begin{aligned} \Gamma_{\gamma\gamma}(\eta_b) \times \text{BR}(\eta_b \rightarrow 4 \text{ charged particles}) &< 48 \text{ eV}/c^2, \\ \Gamma_{\gamma\gamma}(\eta_b) \times \text{BR}(\eta_b \rightarrow 6 \text{ charged particles}) &< 132 \text{ eV}/c^2. \end{aligned}$$

The L3 experiment, looking for  $\eta_b$  in the decay modes  $\eta_b \rightarrow K^+K^-\pi^0$ ,  $\pi^+\pi^-\eta$ , 2, 4 and 6 charged particles (only or associated with one  $\pi^0$ ), observes 6 candidate events with 2.5 background events expected. This corresponds to a combined upper limit on  $\Gamma_{\gamma\gamma}(\eta_b) \times \text{BR}(\eta_b)$  [2]:

$$\Gamma_{\gamma\gamma}(\eta_b) \times \text{BR}(\eta_b \rightarrow \text{analysed channels}) < 200 \text{ eV}/c^2.$$

In this paper we report on the search for  $\eta_b$  in the reaction

$$e^+e^- \rightarrow e^+e^-\gamma^*\gamma^* \rightarrow e^+e^-\eta_b$$

with  $\eta_b$  decaying into the following final states:

$$\begin{aligned}\eta_b &\rightarrow 4\pi^\pm(K^\pm), \\ \eta_b &\rightarrow 6\pi^\pm(K^\pm), \\ \eta_b &\rightarrow 8\pi^\pm(K^\pm).\end{aligned}$$

Here the charged K's in parentheses indicate that a pair of pions may be replaced by a pair of kaons with net zero strangeness.

## 2 Experimental procedure

The analysis presented here is based on the data taken with the DELPHI detector [7, 8] in 1996-2000, covering a range of centre-of-mass energies from 161 to 209 GeV (average centre-of-mass energy: 195.7 GeV). The selected data set corresponds to the period when the Time Projection Chamber (TPC) was fully operational thus ensuring good particle reconstruction. This requirement reduces the integral luminosity for the analysis to  $617 \text{ pb}^{-1}$ .

For quasi-real photon interactions, the scattered  $e^\pm$  are emitted at very small polar angles. Hence there is no requirement on detecting them.

The  $e^+e^- \rightarrow e^+e^-\eta_b$  candidate events are selected by requiring final states with 4, 6 or 8 tracks with zero net charge. Charged-particle tracks in the detector are accepted if the following criteria are met:

- particle transverse momentum  $p_T > 150 \text{ MeV}/c$ ;
- impact parameter of a track transverse to the beam axis  $\Delta_{xy} < 0.5 \text{ cm}$ ;
- impact parameter of a track along the beam axis  $\Delta_z < 2 \text{ cm}$ ;
- polar angle of a track  $10^\circ < \theta < 170^\circ$ ;
- track length  $l > 30 \text{ cm}$ ;
- relative error of the track momentum  $\Delta p/p < 30\%$ .

No  $K_S^0$  reconstruction is attempted on each track pair. The identification of other neutral particles is made using calorimeter information. The calorimeter clusters which are not associated to charged-particle tracks are combined to form the signals from the neutral particles ( $\gamma$ ,  $\pi^0$ ,  $K_L^0$ , n). A minimum measured energy of 1 GeV for showers in the electromagnetic calorimeters and 2 GeV in the hadron calorimeters is required.

The selection of candidate events is achieved by applying the following criteria:

- no particle is identified as an electron or a muon by the standard lepton-identification algorithms [9];
- no particle is identified as a proton by the standard identification algorithm [9];
- there are no electromagnetic showers with energy  $E_{shower} > 1 \text{ GeV}$  or converted  $\gamma$ 's with energy  $E_\gamma > 0.2 \text{ GeV}$  in the event.

To ensure that no particle from the  $\eta_b$  decay has escaped detection, the square of the total transverse momentum of charged particles,  $(\sum \vec{p}_T)^2$ , is required to be small. The actual cut value is estimated from a Monte Carlo sample of  $\eta_b$  events produced in  $\gamma\gamma$  interactions. In this simulation the kinematical variables are generated using the algorithms developed by Krasemann et al. [10]. It is also assumed that the production amplitude factorizes into the quasi-real transverse photon flux and a covariant amplitude describing both the  $\eta_b$  production and decay [11]. The  $\eta_b \rightarrow (4, 6, 8)$  charged-particle decay processes are assumed to be described by the phase-space momenta distribution. The generated events are passed through the standard DELPHI detector simulation and reconstruction programs [8]. The same selection criteria are applied on the simulated events as on the data. Finally, an event is accepted on the basis of the trigger efficiency. Parametrized for a single track, as a function of its transverse momentum  $p_T$ , it ranges from 20% for  $p_T < 0.5$  GeV/ $c$  to about 95% at  $p_T > 2$  GeV/ $c$  [12]. Due to the high mass of the  $\eta_b$  resonant state and relatively large number of tracks in the final state, the overall trigger efficiency per event is about 93.6%, 94.5% and 94.6% for events with 4, 6 and 8 charged particles, respectively.

Fig.1 shows, in the visible invariant-mass interval  $8 \text{ GeV}/c^2 < W_{vis} < 10 \text{ GeV}/c^2$ , the fraction of remaining events as a function of a cut,  $P_T^2$ , on  $(\sum \vec{p}_T)^2$ , for the 4 charged-particle channel. It decreases rapidly for  $P_T^2 < 0.1 \text{ GeV}^2/c^2$ . Hence to preserve the statistics, 4, 6 charged-particle events with  $(\sum \vec{p}_T)^2$  up to  $0.08 \text{ GeV}^2/c^2$  and 8 charged-particle events with  $(\sum \vec{p}_T)^2$  up to  $0.06 \text{ GeV}^2/c^2$  were kept.

The  $\pi/K$  identification is based on the TPC  $dE/dx$  and RICH [13] measurements which are used both separately and combined, in order to check the consistency, in a neural network-based algorithm [14]. In the  $\eta_b$  search region defined as  $8 \text{ GeV}/c^2 < W_{vis} < 10 \text{ GeV}/c^2$ , the average  $K^\pm$  identification efficiency is about 54% and the purity is 82%. The misidentification of charged pions as kaons is about 1.5%. After application of the selection criteria and requiring  $W_{vis} > 5 \text{ GeV}/c^2$ , the 4, 6 and 8 charged-particle data samples contain 173, 328 and 113 events respectively.

The main background comes from inclusive  $\gamma\gamma \rightarrow q\bar{q}$  channels. This background is estimated using a Monte Carlo sample generated with the PYTHIA 6.143 program [15].

The possible contamination of the  $e^+e^- \rightarrow e^+e^-\tau^+\tau^-$  process is given special attention. To reduce it in the  $\gamma\gamma \rightarrow 4\pi$  channel where it is most important, events of topology 1-3 with respect to the hemispheres defined by the thrust axis computed in the  $4\pi$  centre-of-mass system and with an invariant mass, in each hemisphere, smaller than  $1.8 \text{ GeV}/c^2$ , are discarded. Only  $(1.0 \pm 0.3)\%$  of  $\eta_b$  events are eliminated by this cut.

The mass resolution in the search region, shown on Fig.2, has been estimated from the Monte Carlo sample of  $\gamma\gamma \rightarrow q\bar{q}$  interactions. It is about  $200 \text{ MeV}/c^2$  FWHM for all topologies for the 4 charged-particle events. Therefore we have chosen  $200 \text{ MeV}/c^2$  mass intervals to search for a possible signal.

### 3 Results

The visible invariant-mass spectra of the selected events are presented in Fig. 3. When an event has an odd number of  $K^\pm$ , the kaon mass is assigned sequentially to the other particles of opposite charge and the  $W_{vis}$  mass is simply taken as the average of the various mass combinations. The resulting mass shift, averaged over the 4, 6 and 8 particle samples,

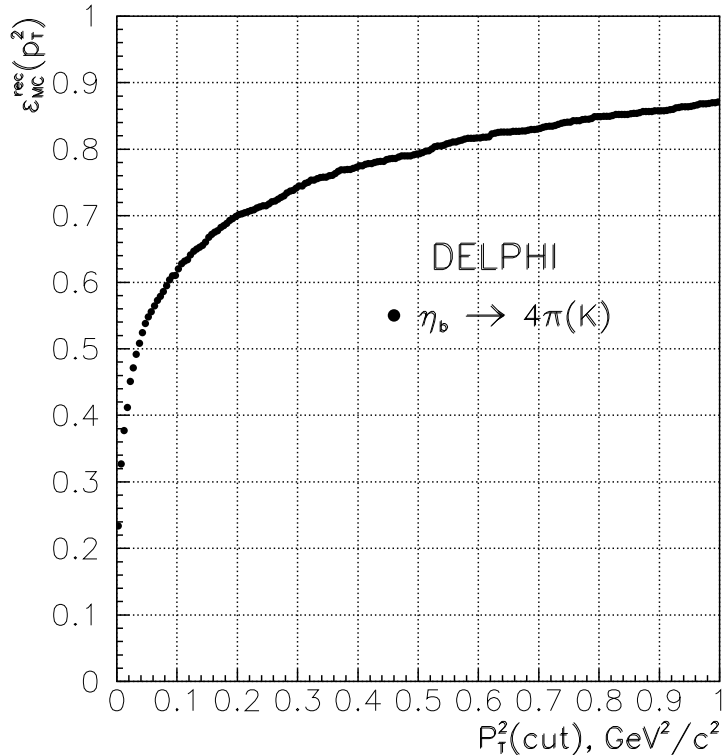


Figure 1: Efficiency of selected  $\eta_b$  Monte Carlo events of the 4 charged-particle channel, as a function of the cut  $(\sum \vec{p}_T)^2 < P_T^2$ , in the  $\eta_b$  search region:  $8 \text{ GeV}/c^2 < W_{vis} < 10 \text{ GeV}/c^2$ .

is about  $120 \text{ MeV}/c^2$  in the  $\eta_b$  search region.

The distributions are well reproduced by the  $\gamma\gamma \rightarrow q\bar{q}$  Monte Carlo simulation. The  $\eta_b$  candidates are expected to show up in the  $9.2$  to  $9.6 \text{ GeV}/c^2$  mass region.

Table 1 gives the number of 4, 6 and 8 charged-particle events in the  $9.2$  to  $9.4$  and  $9.4$  to  $9.6 \text{ GeV}/c^2$  mass regions, together with the number of expected background events computed taking into account the overall reconstruction and selection efficiency. Among the 3 observed  $\eta_b$  candidates only the event with 8 charged particles contains one identified kaon.

In the search for rare processes with a few observed events that may be compatible with background, an upper limit for the signal  $S$  can be derived considering a Poisson process with a background  $b$  and taking into account uncertainties in the background and efficiencies [16]

$$\text{CL} = 1 - \frac{\int g(b)f(\varepsilon) \sum_{k=0}^n P[k|(S\varepsilon + b)]d\varepsilon db}{\int g(b) \sum_{k=0}^n P[k|b]db}.$$

Here  $P(k|x)$  is the Poisson probability of  $k$  events being observed, when  $x$  are expected; CL is a confidence level,  $n$  is the number of observed events. The probability-density functions for the background  $g(b)$  and the efficiency  $f(\varepsilon)$  are assumed to be Gaussian and restricted to the range where  $b$  and  $\varepsilon$  are positive.

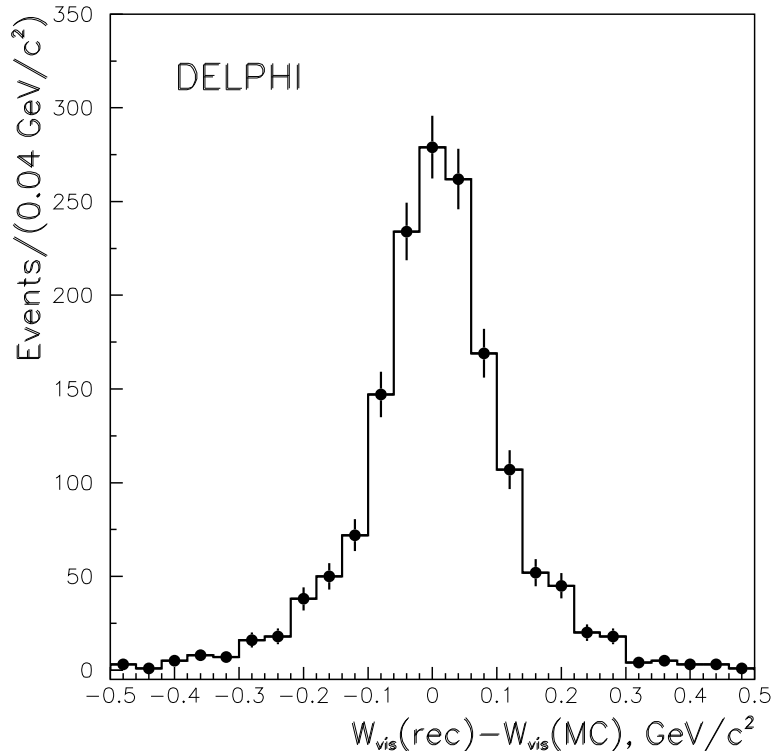


Figure 2: Difference between reconstructed and generated  $W_{vis}$  values for the selected 4 charged-particle events from the Monte Carlo  $\gamma\gamma \rightarrow q\bar{q}$  sample, in the  $\eta_b$  search region.

Upper limits at the 95% confidence level were calculated for each channel and a limit on  $\Gamma_{\gamma\gamma}(\eta_b) \times \text{BR}(\eta_b)$  could then be derived. The values are quoted in Table 1. Assuming that the  $\eta_b$  branching ratio for each channel is the same, the combined upper limit for  $\Gamma_{\gamma\gamma}(\eta_b) \times \text{BR}(\eta_b)$  is  $360 \text{ eV}/c^2$ ,  $\text{BR}(\eta_b)$  being the sum of the branching ratios of the decay modes considered [17].

We considered as main sources of systematic uncertainties: the statistical error of the background, the generator used for the  $\eta_b$  signal and the theoretical uncertainties of the  $\eta_b$  parameters. The limited statistics of our Monte Carlo event sample introduces relative uncertainties of 3%, 5%, 4% for the channels with 4, 6 and 8 charged particles respectively. To appreciate the influence of the generators, we have used PHOT02 [1, 18] which generates  $\eta_b$  events decaying into two gluon-jets. The relative differences in efficiency are of 24%, 11.4% and 6.1% for the 4, 6 and 8 charged particles channels. Varying the  $\eta_b$  mass within the range of  $9.33 - 9.45 \text{ GeV}/c^2$  generates a relative uncertainty of 2.5% on  $N_{ev}$ , for each considered  $\eta_b$  decay channel. The three kinds of uncertainties were added quadratically to obtain the upper limits quoted in Table 1.

## 4 Conclusions

The pseudoscalar meson  $\eta_b$  has been searched for through its decays to 4, 6 and 8 charged-particles in two-photon interactions at LEP II. The data sample corresponds to a total

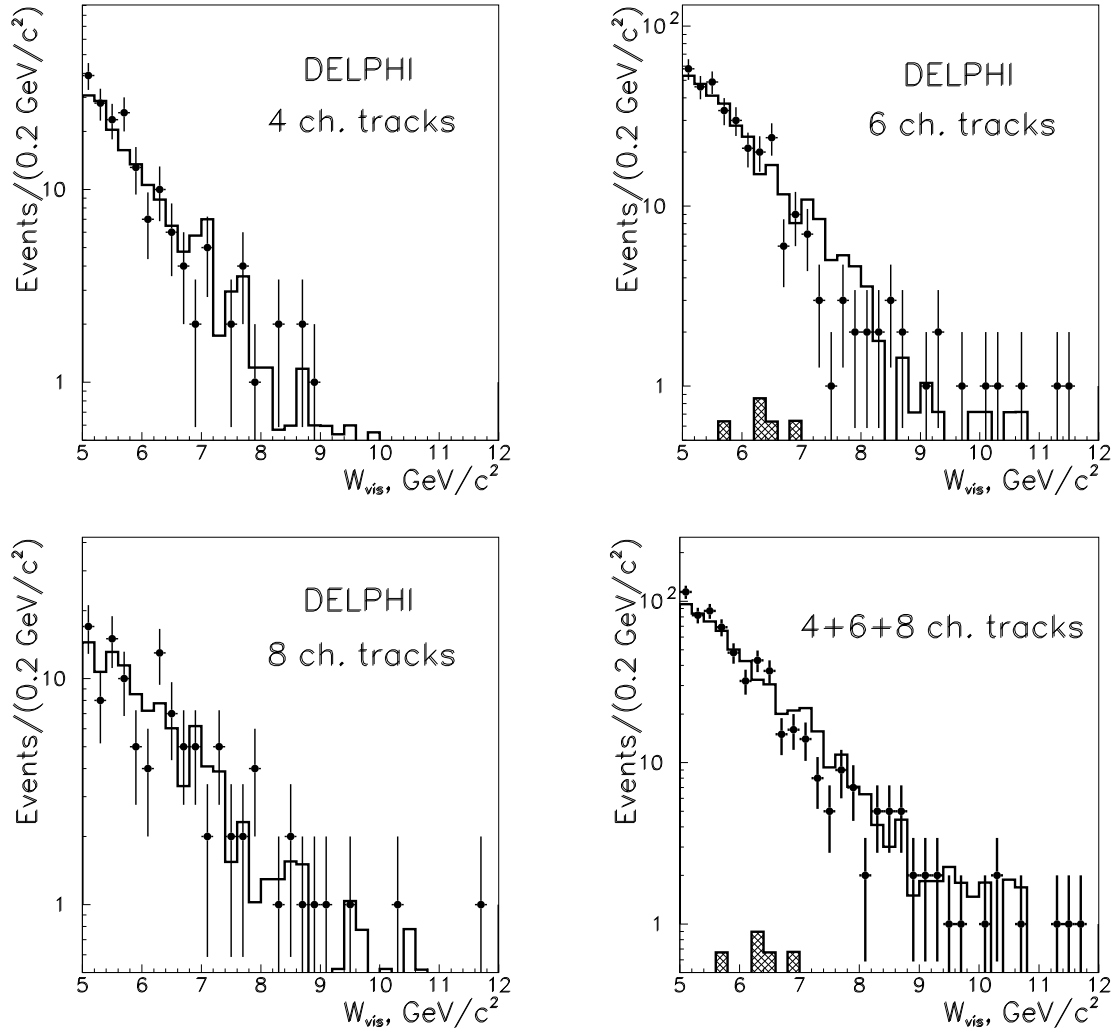


Figure 3: Invariant-mass distributions of selected events for 4, 6 and 8 charged-particle final states. Points with error bars are from the data; histograms present the expected number of background events from the  $\gamma\gamma \rightarrow q\bar{q}$  simulation; shaded histograms correspond to the expected  $\text{Isa } e^+e^- \rightarrow e^+e^-\tau^+\tau^-$  background.

integrated luminosity of  $617 \text{ pb}^{-1}$  collected at centre-of-mass energies ranging from 161 to 209 GeV.

Upper limits at a confidence level of 95% on the product  $\Gamma_{\gamma\gamma}(\eta_b) \times \text{BR}(\eta_b)$  are 190, 470 and 660  $\text{eV}/c^2$  for the  $\eta_b \rightarrow (4, 6, 8)$  charged particle decays, respectively. The combined upper limit for  $\Gamma_{\gamma\gamma}(\eta_b) \times \text{BR}(\eta_b)$ , for the sum of the three decay channels studied, assuming that the  $\eta_b$  branching ratio is the same for each channel, is 360  $\text{eV}/c^2$ , at the 95% confidence level.

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	$\eta_b$ decay modes		
	4 ch.tracks ( $N_{bkg}$ )	6 ch.tracks ( $N_{bkg}$ )	8 ch.tracks ( $N_{bkg}$ )
$N_{obs}$ ( $9.2 < W_{vis} < 9.4$ GeV/ $c^2$ )	0 (0.6)	2 (0.7)	0 (0.5)
$N_{obs}$ ( $9.4 < W_{vis} < 9.6$ GeV/ $c^2$ )	0 (0.6)	0 (0.4)	1 (1.0)
$N_{ev}$ (95% C.L. upper limit)	3.9	5.7	4.1
overall efficiency	5.9%	3.5%	1.8%
$\Gamma_{\gamma\gamma}(\eta_b) \times BR(\eta_b)$ , eV/ $c^2$ (95% C.L. upper limit)	190	470	660

Table 1: Number of observed 4, 6 and 8 charged-particle  $\eta_b$  candidates ( $N_{obs}$ ), expected background events ( $N_{bkg}$ ), 95% C.L. upper limits for signal events ( $N_{ev}$ ), overall efficiency and 95% C.L. upper limits on  $\Gamma_{\gamma\gamma}(\eta_b) \times BR(\eta_b)$ .

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