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Title: Effect of energy source, salt concentration and loading force on colloidal interactions between *Acidithiobacillus ferrooxidans* cells and mineral surfaces

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1 **Highlights:**

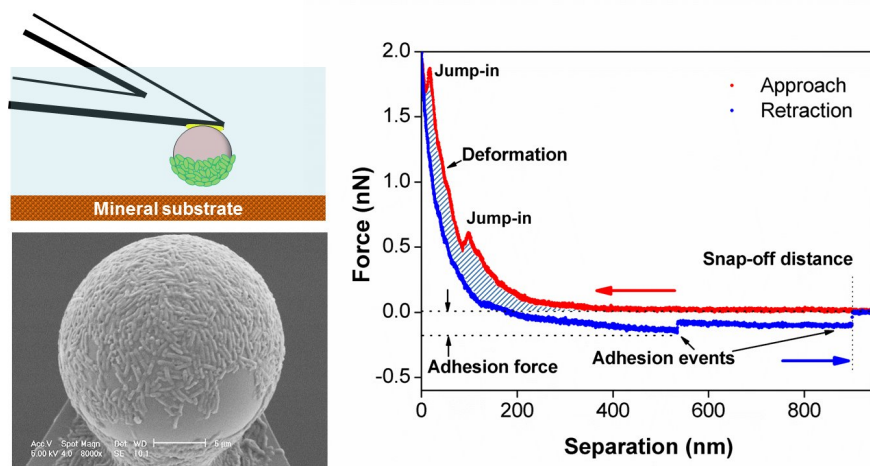
- 2 • Distinct retraction patterns of *A. f* grown with different energy sources were observed
- 3 • Interaction forces between *A. f* and minerals with bacterial probe were quantified
- 4 • The conformation of surface biopolymers was affected by salt concentration

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GRAPHICAL ABSTRACT



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A bacterial colloid probe and representative force-separation curves

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10 **Effect of energy source, salt concentration and loading force on colloidal**
11 **interactions between *Acidithiobacillus ferrooxidans* cells and mineral**
12 **surfaces**

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29 **Abstract**

30 The surface appendages and extracellular polymeric substances of cells play an
31 important role in the bacterial adhesion process. In this work, colloidal forces and
32 nanomechanical properties of *Acidithiobacillus ferrooxidans* (*A. f*) interacted with silicon
33 wafer and pyrite (FeS_2) surfaces in solutions of varying salt concentrations were
34 quantitatively examined using the bacterial probe technique with atomic force microscopy. *A.*
35 *f* cells were cultured with either ferrous sulfate or elemental sulfur as key energy sources. Our
36 results show that *A. f* cells grown with ferrous ion and elemental sulfur exhibit distinctive
37 retraction force vs separation distance curves with staircase and sawtooth shapes, respectively.
38 During the approach of bacterial probes to the substrate surfaces, surface appendages and
39 biopolymers of cells are sequentially compressed. The conformations of surface appendages
40 and biopolymers are significantly influenced by the salt concentrations.

41 **Keywords:** Atomic force microscopy, *Acidithiobacillus ferrooxidans*, bacterial adhesion,
42 nanomechanical properties

43 **1. Introduction**

44 Bacterial adhesion to mineral surfaces is of great importance to the growth of bacteria in
45 natural habitats and many industrial applications [1, 2]. In these processes, the bacterial
46 surface largely determines the adhesion process by the surface appendages (*e.g.* pili and
47 flagella) and extracellular polymeric substances. In the bioleaching process, *Acidithiobacillus*
48 *ferrooxidans* (*A. f*) was the first described metal sulfide oxidizing microorganism, which is
49 affiliated with the Gram-negative γ -Proteobacteria. It is one of the most important species in
50 the bioleaching of sulfide ores operating at temperature lower than 40 °C [3]. *A. f* is endowed

51 with a remarkably broad metabolic capacity, as it can live on the oxidation of ferrous salts,
52 elemental sulfur and a variety of sulfide minerals [4][5, 6].

53 Various growth substrates may induce physiological differences in the chemical
54 composition of cell surfaces, which reflects the response of cells in optimizing nutrient
55 uptake. Research associated with macroscopic assays of bacterial adhesion [7, 8], analysis of
56 chemical compositions for cell surface biopolymers [8, 9], and characterization of cell surface
57 structures [10, 11] has been well documented in the literature. However, the effect of
58 different energy sources on bacterial adhesion behavior and the correlation between surface
59 properties and fundamental interacting forces have not been fully resolved at nanoscale.

60 In the past decade, remarkable developments in atomic force microscopy (AFM) have
61 made it a versatile tool to determine the surface structures and specific interactions of
62 biological samples under near-physiological conditions [12, 13]. AFM is capable of sensing
63 picoNewton forces in aqueous solutions, and the obtained force-separation curves can
64 provide information on the adhesive and nanomechanical properties of biological samples
65 [14-18]. Tipped cantilevers have been extensively used as indenters to probe the elastic
66 properties of different bacterial cells such as *Escherichia coli* [19, 20], *Pseudomonas*
67 *aeruginosa* [21] and *Shewanella putrefaciens* [18]. Alternatively, a cell probe [20, 22] can be
68 used to measure the overall mechanical properties of the cell. A colloidal probe (a
69 microsphere glued onto the end of a cantilever) is often used to indent larger mammal cells
70 [17, 23]. However, the use of AFM to investigate the nanomechanical properties of
71 bioleaching bacteria has been rarely reported.

72 The goal of this study is to relate the adhesion behavior and nanomechanical interactions
73 to the biophysical responses of bacterial cells to the change of environmental conditions
74 (energy sources and salt concentrations). To this end, we performed AFM force

75 measurements using bacterial probes constructed with *A. f* cells grown with different energy
76 sources of Fe^{2+} and S^0 , and exposed the bacterial probes to solutions of various salt
77 concentrations. The shape of the retraction curves, adhesion forces and the Young's moduli
78 of cell surface biopolymers were compared for *A. f* cells grown in different energy sources.
79 Interesting features such as sequential "jump-in" events of approaching curves and distinct
80 retraction curve patterns of *A. f* cells grown with the energy sources were obtained. Overall,
81 our findings quantitatively describe the adhesion behaviors of *A. f* on mineral surfaces and the
82 nanomechanical properties may help the further understanding of responses of cell surface
83 appendages to environmental stimuli.

84 **2. Materials and methods**

85 **2.1 Microorganism and growth conditions**

86 *A. f* was kindly provided by Professor Guohua Gu (School of Mineral Processing and
87 Bioengineering, Central South University, China). Cells were cultured at 30 °C in 9K
88 medium (pH 2.0) [24]: $(\text{NH}_4)_2\text{SO}_4$ 3 g/L, KCl 0.1 g/L, $\text{K}_2\text{HPO}_4 \cdot 3\text{H}_2\text{O}$ 0.5 g/L, $\text{MgSO}_4 \cdot 7\text{H}_2\text{O}$
89 0.5 g/L, $\text{Ca}(\text{NO}_3)_2$ 0.01 g/L. Bacteria were grown with 4.47% (w/v) FeSO_4 and 3% (w/v)
90 elemental sulfur as energy source, respectively. *A. f* cells were incubated on a rotary shaker at
91 170 rpm to their mid-exponential growth phase.

92 **2.2 Substrate preparation**

93 The silicon wafers (100 oriented with a 100-nm thermal-oxide surface layer) were
94 purchased from Silicon Valley Microelectronics (USA). The silica surfaces were cleaned
95 using the RCA SC-1 process [25] and stored in Milli-Q deionized water (18.2 M Ω cm,
96 Millipore, USA) before use. A museum-grade pyrite (FeS_2) sample obtained from Ward's
97 Natural Science was embedded in an epoxy resin and cut off in a thickness of 3 mm. The
98 sample slice was manually polished with 3 μm and 1 μm diamond suspensions, respectively,

99 and thoroughly washed with acetone, ethanol and copious amount of deionized water. The
100 sample slice was UV treated for 30 min before exposure to the bacterial probe.

101 **2.3 Zeta potential measurements**

102 Cultures in the mid-exponential phase were filtered through Whatman filter paper to
103 remove suspended solid materials. Cells were harvested by centrifugation at acceleration of
104 10000 g-units for 15 min. The cell pellet was washed three times using sterilized H₂SO₄ (pH
105 1.5) and deionized water to remove trapped ions. The washed cell pellet was re-suspended in
106 0.001, 0.01, 0.1 and 1 M KCl solutions, respectively, to obtain a concentration of
107 approximately 1×10^7 cells/mL. The zeta potentials of *A. f* were measured using a zeta
108 potential analyzer (ZetaPLUS analyzer, Brookhaven Instruments Corp.) and calculated from
109 the electrophoretic mobility using the Smoluchowski equation embedded in the ZetaPlus
110 software. Measurements were conducted in triplets and the average value was used.

111 The zeta potential of planar silica surfaces was obtained using an EKA (electro kinetic
112 analyzer) with an asymmetric clamping cell (Anton Paar, GmbH, Austria). A piece of PMMA
113 was used as a supporting medium (back-plate) in the asymmetric clamping cell. The
114 streaming potential measurements were taken three times in each salt solution. The zeta
115 potential was calculated from the streaming potential according to the approach developed by
116 Fairbrother and Mastin embedded in the software [26]. The isoelectric point of pyrite surface
117 was reported to be about pH 2 in literatures [27].

118 **2.4 ATR FT-IR spectroscopy**

119 Cells in the mid-exponential phase were collected and washed three times with H₂SO₄
120 and deionized water. Cell pellets were dried at 50 °C for 30 min before acquiring the spectra.

121 The infrared spectra of bacterial surfaces were measured using a Perkin Elmer Spectrum 100
122 spectrometer. Spectra were the results of 43 scans with a resolution of 1 cm^{-1} in the range
123 $650\text{-}4000\text{ cm}^{-1}$.

124 **2.5 Preparation of bacterial probes**

125 Bacterial probes were prepared using a protocol described in the previous study [28]:
126 Tipless cantilevers (Veeco, model NP-OW) were cleaned with strong oxidizing Piranha
127 solution [29] for 30 min, and rinsed with copious amounts of deionized water before blow-
128 drying with high purity compressed nitrogen gas. A silica microsphere (20 μm in diameter,
129 Fuso Chemical Co., Japan) cleaned by the RCA SC-1 solution was glued to the end of a
130 tipless cantilever with a small amount of thermoplastic epoxy resin using a micromanipulator
131 under an optical microscope. The colloidal probe was functionalized with 1% (w/v)
132 polyethyleneimine (PEI, MW~1300, Sigma-Aldrich, Australia) solution for 2.5 h. The excess
133 solution was decanted and the probe was rinsed in deionized water and stored at $4\text{ }^{\circ}\text{C}$.

134 Bacterial pellets were washed and resuspended in a 3% (v/v) glutaraldehyde solution for
135 cell fixation at $4\text{ }^{\circ}\text{C}$ for 2.5 h [30]. After fixation the cells were washed with phosphate buffer
136 solution (PBS), and resuspended in PBS at $4\text{ }^{\circ}\text{C}$ overnight. The cell suspension was spread
137 onto a clean glass slide to allow the colloidal probe to touch the suspension by the means of a
138 micromanipulator. The bacterial probe was then gently rinsed with deionized water to remove
139 loosely attached cells and kept hydrated before force measurements. Scanning electron
140 microscopy (SEM) (Philips XL-30) was performed on all bacterial probes after AFM
141 measurements to verify the presence of cells on the microsphere.

142 2.6 AFM force measurements

143 Ultrahigh purity KCl (Sigma-Aldrich) was roasted at 500 °C for 12 h and used as the
144 supporting electrolyte. Force measurements were performed at room temperature in KCl
145 solutions at natural pH 5.6 ± 0.5 using a MFP-3D atomic force microscope (Asylum Research,
146 Santa Barbara, CA) equipped with a closed fluid cell. The actual spring constants of the
147 cantilevers were determined using the thermal noise method embedded in the Asylum
148 Research AFM software [31]. The cantilevers used in this study were found to have a spring
149 constant of 0.09 ± 0.02 N/m. After each probe being immersed in the solution for at least 20
150 min, force curves were recorded under a loading force of 2 nN at an approaching/retraction
151 velocity of 500 nm/s with a piezo movement of 6000 nm. At least three probes, as well as the
152 control probe (PEI-coated silica colloid probe), were used for each set of experiments at 3-5
153 contact locations per probe.

154 To investigate the effect of loading force on the nanomechanical properties of bacterial
155 cells, the force curves were measured at different loading forces from 0.5 nN to 2.8 nN. After
156 a series of consecutive force measurements applying increasing loading forces, the
157 measurement was conducted again at 2 nN. Once this force profile differed from the
158 previously measured ones at 2 nN, the bacterial probe was considered damaged and replaced
159 by a new one.

160 2.7 AFM data analysis

161 2.7.1 Raw data conversion

162 From the approaching and retraction curves, several useful parameters can be extracted,
163 for instance, the adhesion force, snap-off distance and Young's modulus. Due to the
164 deformable nature of biological samples, a correct conversion from raw data to the force-
165 separation curves is crucial to the investigation of bacteria-mineral interactions.

166 The direct results (raw data) recorded by AFM in a force measurement is a measure of
167 the cantilever deflection (d , V) versus the relative piezo displacement (z , nm). Deflection
168 (V) can be converted to cantilever deflection (nm) with a sensitivity of the cantilever obtained
169 by engaging the cantilever against a rigid surface. The raw data can then be converted to
170 force (nN or pN) according to Hooke's law: $F = k_c d$, where k_c is the spring constant of the
171 cantilever. For deformable bacterial cells, the conversion from raw data to force-separation
172 curves is not as simple as that for rigid solid surfaces. The cell deforms in response to the
173 interaction forces and loading forces as shown in Fig. 1A. When the bacterial probe is
174 brought toward the substrate surface, the soft cell is squeezed, reflecting the deformation of
175 the cell surface appendages, biopolymer brush and/or the cell body depending on the
176 magnitude of loading force.

177 The distance balance, $z_0 + 2R + s_0 = (z - d) + 2R + (s - \delta)$, gives the following equation
178 for the actual separation distance s :

$$179 \quad s = (z_0 + s_0) + \delta + d - z \quad (1)$$

180 The relative piezo position z and the cantilever deflection d are directly obtained from the
181 AFM measurement, while the values of sample deformation δ and $z_0 + s_0$ have to be
182 determined. Due to the deformable nature of cells and the presence of interaction forces

183 between cells and the substrate surfaces, we fit the portion of raw data at high loading force
 184 with the Hertz model to determine the required parameters (see Section 2.7.3 for more detail).

185 **2.7.2 Steric model**

186 Cell surface appendages and/or biopolymer chains can be considered as a layer of
 187 polymeric brush. The steric forces are often the dominant forces during the approach of cells
 188 to solid surfaces, particularly in high ionic strength solutions where the electrostatic forces
 189 between cell and substrate surfaces are largely compressed. The steric force per unit area
 190 between two surfaces, only one of which is coated with polymers is estimated by [32, 33]

$$191 \quad f = 50k_B T \Gamma^{3/2} e^{-2\pi s/L_0} \quad (2)$$

192 where k_B is the Boltzmann constant, T is the absolute temperature, s is the distance between
 193 the two surfaces, Γ is the density of grafted polymers in m^{-2} and L_0 is equilibrium length of
 194 the polymer brush.

195 In our case, the substrate surface is bare, while the cell surface is covered with
 196 biopolymers and considered as a microsphere with a radius R of 500 nm (Fig. 1B). We
 197 approximate the total force by integrating f over half of the cell surface as follows:

$$198 \quad F = \int_0^{\pi/2} f 2\pi r dr = \int_0^{\pi/2} f 2\pi R^2 \sin \theta d\theta = 50k_B T R \Gamma^{3/2} L_0 e^{-2\pi(s+R)/L_0} (1 - e^{-2\pi R/L_0}) \quad (3)$$

199 where $r = R \sin \theta$ is the radial distance from a random point on the cell surface to the vertical
 200 axis of symmetry.

201 **2.7.3 Hertz model**

202 Among different models describing the elastic response of soft samples, the Hertz model
 203 has been widely used to describe similar systems in AFM experiments [17, 34-36]. In the
 204 Hertz theory, the cell is assumed to be an isotropic material with a well-defined interface and

205 any surface interactions or adhesions are neglected [37]. At the region of Hertz contact, we
 206 have the following equations for the distance of cell deformation and the force of cell
 207 deformation in contact with a rigid plane to the AFM data [38]:

$$208 \quad \delta = Ad^{2/3} \quad (4)$$

$$209 \quad F = \frac{4}{3} \frac{E}{1-\nu^2} R^{1/2} \delta^{3/2} \quad (5)$$

210 where E is the Young's modulus, R is the radius of bacterial cells, taken as 500 nm, ν is
 211 the Poisson's ratio of the deformable bacterial cells (assumed to be 0.5) [34] and A
 212 represents relationships between the sample radius, the Poisson ratio and the Young's
 213 modulus of the sample according to different geometries of the systems. The Young's
 214 modulus of a bacterial cell is obtained from quantitative interpretation of the non-linear
 215 regime that follows the steric interaction portion.

216 The deformation δ is dependent on the contact point of the bacterial probe with the
 217 substrate. However, due to the influence of interaction forces (*e.g.* steric repulsion), it is
 218 difficult to determine the accurate contact point. Various approaches such as manual
 219 determination by visual inspection [36], semi-automated [39, 40] and automated approach
 220 with software [41, 42], have been used for determining the contact point. Here, we fit the
 221 region of raw curve at high loading force with the Hertz model to determine the separation
 222 between bacteria and substrate surfaces and the Young's modulus of the cell with Eq. (1) and
 223 (5).

224 **2.7.4 WLC model**

225 The wormlike chain (WLC) model is commonly used to describe the elasticity of the
 226 biopolymer chains such as proteins and DNA [43-45]:

$$227 \quad F(s) = \frac{k_B T}{L_p} \left[\frac{s}{L_c} + \frac{1}{4(1-s/L_c)^2} - \frac{1}{4} \right] \quad (6)$$

228 where L_p and L_c are the persistence length and contour length of the polymer chains. In this
 229 study, the WLC model was applied to interpret the sawtooth-shaped adhesion events of
 230 sulfur-grown cells.

231 **3. Results and discussion**

232 **3.1 Bacterial probes**

233 The SEM image in Fig. 2A shows a cell-coated colloidal probe which was used for a
 234 series of force measurements. The contact area of the microsphere was covered by bacterial
 235 cells, thus the measured force curves can reflect the true bacterial-mineral interactions. Fig.
 236 2C displays typical approach and retraction force curves recorded from a ferrous ion-grown
 237 cell probe. Far from the substrate surface, the bacterial probe senses no interaction forces
 238 between the surfaces. As the probe approaches the substrate, the approaching curve exhibits a
 239 repulsive force with several jump-in events. After the mutual contact, adhesins of cells adhere
 240 to the substrate surface and show multiple adhesion events during retraction. The deviation in
 241 the contact region between the approaching and retraction part of the force curve (the shaded
 242 area) is due to the deformation of the cell. Usually, the typical “loading-unloading hysteresis”
 243 can help to verify the quality of bacterial probes before conducting SEM tests.

244 **3.2 Surface characterization**

245 The zeta potentials of ferrous ion- and sulfur-grown *A. f* and the silica wafer in the
 246 presence of various salt concentrations are shown in Fig. 3A. At neutral pH, bacterial cells
 247 are negatively charged and the absolute value of zeta potential decreases with an increase of
 248 the salt concentration, which is in line with enhanced screening of the charges within the

249 surface appendage and/or biopolymer chains by ions present in the solution. The ferrous ion-
250 grown cells were slightly more negatively charged than sulfur-grown cells in all solutions,
251 which is in agreement with the findings of Sharma *et al.* [46]. The zeta potential results
252 demonstrate that different energy sources can affect the surface charge of *A. f.*

253 An early study [47] found that *A. f.* cells cultured with sulfur, pyrite (FeS_2) and
254 chalcopyrite (CuFeS_2) are more hydrophobic than ferrous ion-grown cells, indicating that
255 various energy sources can affect the cell surface properties. The depletion of soluble ferrous
256 salts renders the bacteria prone to attaching to the solid energy source such as elemental
257 sulfur and sulfide minerals for growth, which results in the synthesis of more proteinaceous
258 substances for the purpose of facilitating adhesion [8, 9, 46]. To analyze the difference of
259 functional groups on cell surfaces induced by various energy sources, ATR-FTIR tests were
260 performed.

261 As can be seen in Fig. 3B, most peaks in the spectra of ferrous ion- and sulfur-grown
262 cells are at similar positions. The assignments of the peaks indicate that *A. f.* surface consists
263 of extracellular polysaccharides, proteins and nucleic acids [48, 49]. Previous FTIR results
264 reported by Devasia *et al.* [47] suggested that a proteinaceous new cell surface appendage
265 was synthesized in sulfur-grown cells while such an appendage was found to be absent in
266 ferrous ion-grown cells. Our FTIR results differ from their results because cells cultured with
267 either ferrous ion or elemental sulfur show similar peaks which represent the proteins. The
268 only remarkable difference is in the range of $1345\sim 1403\text{ cm}^{-1}$. Ferrous ion-grown cells show
269 a small peak at 1388 cm^{-1} which is absent in the sulfur-grown ones. Sulfur-grown *A. f.* possess
270 two peaks at 1403 and 1345 cm^{-1} respectively, which are absent in the spectra of ferrous ion-
271 grown cells. The peaks in the range of $1388\sim 1403\text{ cm}^{-1}$ represent the C=O symmetric
272 stretching of COO- group in amino acids and fatty acids. The weak peak at 1345 cm^{-1}
273 represents sulfonic acid. Interestingly, the peak at 1388 cm^{-1} was observed in the FTIR

274 spectra of *L. ferrooxidans*, which is also a ferrous ion-oxidizer, while the peak at 1403 cm^{-1}
275 was observed in the spectra of *A. thiooxidans* which is a sulfur oxidizer [28]. It can be
276 inferred from the qualitative FTIR data that specific energy sources can likely induce the
277 differential expression of specific extracellular biopolymers and lead to the difference in the
278 ratio of various biopolymers.

279 3.3 Analysis of retraction curves

280 The retraction curves were examined to obtain information regarding the adhesion
281 behavior (*i.e.* adhesion force and snap-off distance) of the bacterial cells. Representative
282 force-separation curves recorded between *A. f* cells and substrates are shown in Fig. 4. A
283 notable proportion of the retraction curves obtained from ferrous ion-cultured cells are
284 characterized by a stair-step separation, while those recorded from sulfur-cultured cells are
285 featured by larger adhesion forces with sawtooth-shaped multiple peaks. Both separation
286 patterns indicate that the molecular bonds formed between bacterial cells and the substrate
287 surfaces break sequentially until they completely separated from each other [15].

288 A possible reason for the distinct retraction patterns is the energy source-induced
289 differential expression of biopolymers. The stair-step separation pattern likely results from
290 the extension or desorption of extracellular polysaccharides, which was reported by Sletmoen
291 *et al.* [50]. The sawtooth shaped adhesion peaks, which can be well fitted with WLC model,
292 are typical for the stretching of proteins. The fraction of the extracellular proteins of *A. f*
293 cultured with ferrous ion and sulfur was examined by chemical analysis [8] and two-
294 dimensional gel electrophoresis [51]. The results reveal higher amount of protein on sulfur-
295 grown cells compared to ferrous ion-grown cells and the various ratios of total
296 polysaccharides to proteins are responsible for the different attachment abilities of *A. f*, which
297 is also in line with our AFM measurements.

298 As expected for biological samples, the shape of each retraction curve, the number of
299 adhesion events, the magnitude of adhesion forces and the snap-off distances vary from
300 retraction to retraction (Fig. 4 and Table 1). This is attributed to the heterogeneity of the
301 bacterial surface. Different types of biopolymers exist on the cell surface and more than one
302 type can randomly adhere to the substrate surface on contact. In addition, a biopolymer chain
303 can adhere to the substrate at multiple sites on the chain.

304 The properties of retraction curves are summarized in Table 1. To determine the
305 statistical significance of the adhesiveness differences between ferrous ion- and sulfur-grown
306 *A. f.*, we plotted the adhesion forces in histograms (Fig. 4). The data display more than 800
307 retraction curves in total obtained from at least three independent experiments using different
308 bacterial probes and substrates under each salt concentration. As shown in Fig. 4, a wide
309 distribution of adhesion forces was observed due to the heterogeneous nature of the bacterial
310 surface. Salt concentration did not dramatically affect the mean adhesion forces of ferrous
311 ion- or sulfur-grown *A. f.* The adhesion forces of ferrous ion-grown cells are similar to that of
312 sulfur-grown cells on the pyrite surface. Comparing with other sets of experiments, sulfur-
313 grown *A. f.* exhibits a higher adhesion affinity to the silica surface. It is interesting to note that,
314 the stair-step unfolding force of ferrous ion-grown cells increases with an increase of the salt
315 concentration, which indicates that biopolymers become stiffer in solutions of higher salt
316 concentration and require larger unfolding forces during retraction.

317 A wide range of the snap-off distances between the bacterial probe and the substrate
318 surfaces was also observed in all solutions (Fig. 5), which again reflects the heterogeneity of
319 the biopolymers on the bacterial surface. The various snap-off distances also provide solid
320 evidence of different responses of the surface appendages and biopolymer chains to various
321 salt concentrations. The snap-off distances vary in a wide range up to 2.2 μm , demonstrating
322 the existence of pili and/or flagella. However, due to the use of different bacterial probes and

323 the small number of flagella, the frequency of the flagella adhesion events (larger than 1 μ m)
324 is low.

325 **3.4 Analysis of approaching curves**

326 The size of the biopolymers/surface appendages is sufficient to cause the steric repulsion
327 between the bacterial probe and the substrate surface. When the bacterial probes are gradually
328 lowered to touch the substrate surfaces, the steric interactions between the substrate and cell
329 surface biopolymers are first sensed by the probes, followed by mechanical contact and
330 deformation of the surface appendages and polymer brush of the cell envelope as a result of
331 the compression. The physical properties such as the net surface charge of cells and minerals
332 and the conformation of biopolymers change with increasing the salt concentration.

333 In this study, the steric model and Hertz model were applied to the force curves to
334 estimate the length and Young's modulus of the biopolymer layer. Although the electrical
335 double layer force can also demonstrate exponential repulsive behavior, the electrostatic
336 interaction at high salt concentration was weak and ignored because the spatial range where
337 the steric repulsions are operative is significantly larger than that of the electrostatic forces.
338 The Decay length in 0.001 M KCl is approximately 9.7 nm, while the repulsion starts from
339 approximately 600 nm (Fig. 6A). This suggests that the electrostatic model is not applicable
340 to these biopolymers, thus it is justified only to use the steric model and Hertz model to
341 interpret the approaching curves. For approach curves analysis, we only consider the force
342 curves between bacterial probes and the silica wafer, which is more homogeneous and
343 smoother than the pyrite surface.

344 **3.4.1 Effect of salt concentration on biopolymer**

345 A significant effect of salt concentration on the approaching curves was observed (Fig.
346 6A and B). The steric model (Eq. (3)) was fitted to the approaching curves at various salt

347 concentrations. For ferrous ion-grown *A. f.*, the distance of the repulsion force decreased with
348 an increase of the salt concentration (Fig. 6A). From 0.001 M to 1 M salt concentration, the
349 equilibrium length of the polymer brush L_0 was 675 ± 60 , 415 ± 23 , 383 ± 21 and 236 ± 4 nm,
350 respectively. The grafted polymer density Γ was between 6.9×10^{16} m⁻² to 13.9×10^{16} m⁻², the
351 magnitude of which was in agreement with other reports [2, 35, 52]. It is interesting to note
352 that the sulfur-grown cells showed an opposite trend. As can be seen from Fig. 6B, the
353 approaching curves of sulfur-grown *A. f.* showed more repulsive peaks and the distance of
354 repulsion gradually increased with an increase of the salt concentration.

355 The histograms presented in Fig. 6C and D were obtained by analyzing more than 50
356 approaching curves using Eq. 5 for each salt concentration, resulting in different Young's
357 modulus distributions. In our experiments, the applied forces (2 nN) were only sufficient to
358 indent less than 200 nm even in 0.001 M KCl solution (smaller than polymer brush thickness
359 estimated by the steric model). This indicates that the applied forces were mainly used to
360 compress the outer biopolymer brush of the cell wall.

361 Generally, the average values of Young's modulus of ferrous ion-grown cells are larger
362 than that of the sulfur-grown cells in all solutions. The Young's modulus of ferrous ion-
363 grown *A. f.* increased with increasing salt concentration. From 0.001 M to 1 M solution, the
364 average value of E was 30.2 ± 8.9 , 29.2 ± 5.1 , 46.5 ± 5.1 and 64.5 ± 9.1 kPa, respectively. The
365 Young's moduli of sulfur-grown *A. f.* were similar (about 22 kPa) at 0.001 M and 0.01 M KCl.
366 However, the peak distribution of E shifted toward smaller values as the salt concentration
367 increased, showing 14.9 ± 7.1 and 17.2 ± 9.4 kPa at 0.1 M and 1 M, respectively. The Young's
368 moduli obtained from our experimental data by interpretation with Hertz model are in
369 agreement with the magnitudes reported in the literature for E of bacterial biopolymer layers
370 [19, 22, 35, 53], generally ranging from 1 to 100 kPa.

371 In salt solutions with neutral pH, bacterial cells and substrate surface are negatively
372 charged due to the presence of anionic groups such as carboxyl and phosphate groups. By
373 increasing the salt concentration of the solution, the repulsive interactions between
374 neighboring charged chains of biopolymer are screened, leading to the collapse of the
375 biopolymer chains onto the cell membrane [54]. In low salt concentration solutions, the
376 biopolymer chains are more extended (larger L_0), resulting in softer bacterial cells (smaller
377 E). However, the biopolymer thickness of sulfur-grown cells increased with increasing salt
378 concentration. The E value of sulfur-grown cells is approximately 4 times softer than that
379 determined for ferrous ion-grown cells at 1 M. Although the reasons for the opposite trend
380 observed from sulfur-grown cells are as yet not well understood, the results seem to agree
381 with the surface characterization results and indicate that the presence of different external
382 polymeric appendages and/or biopolymers and the change in the ratio between different
383 biopolymers lead to the varied nanomechanical properties of *A. f.* Other surface-specific
384 equipment such as sum frequency generation may be helpful to reveal the molecular
385 differences in surface biopolymers on the cell surface in further researches.

386 The absolute determination of the Young's modulus for a biological sample may not be
387 accurate for the following reasons: (a) *A. f.* cells are rod-shaped instead of a spherical shape.
388 (b) The simple Hertz model ignores the cell-substrate adhesion. (c) The homogeneous
389 assumption for the cell surface in the Hertz model is not physically realistic. The
390 componential and structural complexity in the bacterial surface can lead to variation of the
391 Young's modulus depending on different indentation depth, *i.e.* depending on the layer which
392 is actually squeezed. Various surface components and appendages such as fimbria, flagellum,
393 pili and biopolymer chains can contribute to the overall stiffness measured with AFM.
394 However, a trend in the change of Young's modulus as a function of salt concentration may
395 provide some fundamental information for bacteria-mineral interactions.

396 3.4.2 Effect of loading force

397 Upon approaching, the bacterial probe senses a repulsive force due to the steric
398 interaction. A notable fraction of approaching curves exhibit non-monotonic discontinuities
399 referred as sequential “jump-in” events, which reflect large changes in loading force over
400 very small distances. We propose that as the loading force increases, the resistance levied by
401 the surface appendages (pili or flagella) and/or biopolymer chains of cells against the
402 substrate is suddenly relieved, allowing further compression of the cell with less applied force.
403 However, “jump-in” events are rarely observed for experiments between a tip cantilever and
404 bacterial cells in previous studies [2, 52, 55-57].

405 The reasons for using a colloidal bacterial probe instead of using a tip cantilever as an
406 indenter onto the cell surface are as follows: (a) A sharp tip of the cantilever can possibly
407 penetrate between the biopolymer chains, and deform the cell membrane with the tip apex
408 while deforming a small amount of the polymer chains with the tip sides [17]. (b) A sharp tip
409 with a small contact area can interact with various biopolymers during each approaching
410 process, thus largely enhancing the heterogeneity of the results. To obtain a global
411 nanomechanical property of the cell, a bacterial colloidal probe can provide a much larger
412 contact area, the results from which are equivalent to the average of many measurements with
413 the sharp tip (*i.e.* conducting a nanomechanical force mapping experiment).

414 The effect of loading force on the approaching curves is illustrated for the interactions
415 between ferrous ion-grown *A. f* and the silica surface in 0.001 M KCl solution. For the sake
416 of illustration, the approaching curves were plotted as loading force versus arbitrary
417 separation. As shown in Fig. 7, under a very small loading force of 0.5 nN, the approaching
418 curve displays a monotonic repulsion due to the steric interaction. As the loading force
419 increases, the approaching curves show a first “jump-in” event at a similar distance around
420 150 nm. Under higher loading forces, the approaching curves show several smaller “jump-in”

421 events following the first one. Although the number of “jump-in” events generally increases
 422 with an increase of the loading force, the number also varies from curve to curve, which
 423 again points to the heterogeneities and complexity of cell surfaces. These “jump-in” events
 424 likely reflect the sequential compression of the cell surface appendages and biopolymer
 425 chains. By fitting the “jump-in” events to the Hertz model, we obtained the E values from
 426 low to high loading forces are 6.9, 9.0, 18.4, 26.9, 34.8 and 34.5 kPa, respectively. As we can
 427 see in Fig. 7, under a loading force larger than 1.5 nN, the Young’s moduli of the first
 428 compression events are similar and close to the average value of 30.2 ± 8.9 kPa in 0.001 M
 429 KCl solution.

430 In order to further rule out the possibility that the “jump-in” events result from the
 431 compression of different cells, we used the representative curves in Fig. 7 as an example to
 432 estimate the contact area of the bacterial probe. The Hertz contact area can be estimated as
 433 follows:

$$434 \quad S = \pi a^2 = \pi \left[\frac{3FR(1-\gamma^2)}{4E} \right]^{2/3} \quad (7)$$

435 As the loading force increases from 0.5 to 2.8 nN, the estimated contact area increases from
 436 0.23 to $0.25 \mu\text{m}^2$, which is about half of the area of *A. f* cell (*ca.* $0.5 \mu\text{m}^2$). Similar results
 437 were reported by Zhang et al. [20] that under a loading force of 1.5 ± 0.2 nN, the contact
 438 radius of *E. coli* cells is approximately 45 ± 2 nm. Our previous results showed that when the
 439 bacterial probe was pressed extremely hard onto the silica surface, the area of destroyed cell
 440 layer contained about 10 cells (data not shown). Therefore, under a loading force of 2 nN, it
 441 is likely that only one cell is in contact with the substrate, which in turn supports our claim
 442 that the sequential compression events reflect the compression of biopolymer chains with
 443 different lengths instead of the compression of several cells. Furthermore, we observed that

444 most “jump-in” events of the approaching curves correspond one-to-one with the adhesion
445 peaks of the corresponding retraction curves for sulfur-grown cells. This also indicates the
446 “jump-in” events reflect the compression of fibrous biopolymers.

447 **4. Conclusions**

448 In this study, the fundamental interaction forces between *A. f* cells cultured with different
449 energy sources and the substrates were directly quantified with the bacterial probe technique.
450 Our results show that the conformational changes in biopolymers due to the salt
451 concentration are important factors in influencing the surface potentials, adhesion behavior
452 and the softness of the bacterial cells. This research provides fundamental understanding and
453 evidence that different energy sources and the salt concentration significantly influence the
454 adhesion behavior and cell nanomechanical properties.

455 Bacterial probe technique of AFM is advantageous in direct measurement of the
456 interacting forces between cells and the mineral surface, and is of great importance in the
457 investigation of bacteria-mineral interface research. However, due to a lengthy manual
458 analysis, one often had to compromise by reducing the number of force curves to lower
459 analysis time. For deformable samples with heterogeneous surface components, converting
460 raw data, calculating Young’s modulus, and extracting adhesion force and snap-off distance,
461 would have made the manual route very tedious. Commercial data analysis software suitable
462 for batched AFM data analysis of different geometrical systems will be beneficial for faster
463 and more accurate analysis of AFM data in future researches.

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598 **Captures to Figures**

599 **Figure 1.** (A) Schematic of the relative positions of a bacterial probe with cell deformation
600 (not to scale). The silica microsphere was omitted for simplification. z_0 is the initial piezo
601 position of the cantilever; z is the relative piezo position of the cantilever; d is the cantilever
602 deflection; δ is the deformation of the cell; s_0 is the initial separation between sample and
603 the substrate surface; s is the actual separation distance between the two surfaces. (B)
604 Schematic illustration of the abbreviations used to calculate the interaction force from the
605 pressure (force per unit area). The cell is considered as a microsphere.

606

607 **Figure 2.** (A) A representative SEM image of a bacterial probe of *A. f* (Fe^{2+}). (B) An AFM
608 height image of *A. f* (Fe^{2+}) immobilized on a glass slide. (C) Typical force curves of ferrous
609 ion-grown cells showing the adhesion events, jump-in events and deformation of the cell.

610

611 **Figure 3.** (A) Zeta potentials of *A. f* and silica wafers under various salt concentrations. (B)
612 FTIR spectra of *A. f* cells cultured with (a) ferrous ions and (b) elemental sulfur.

613

614 **Figure 4.** Histograms of adhesion forces and representative retraction curves (inset) obtained
615 from (A) *A. f* (Fe^{2+}) and (B) *A. f* (S^0) at various salt concentrations. Red lines on retraction
616 curves of sulfur-grown *A. f* show that the adhesion forces are well-fitted by the wormlike
617 chain model described by Eq. (6).

618

619 **Figure 5.** Scatter plots of the adhesion forces versus snap-off distances measured between the
620 substrates and (A) $A. f(\text{Fe}^{2+})$ or (B) $A. f(\text{S}^0)$. Data points were randomly collected from 30
621 different retraction curves measured for each salt concentration.

622

623 **Figure 6.** Representative approaching curves fitted with steric model showing the effect of
624 ionic strength on approaching curves of (A) $A. f(\text{Fe}^{2+})$ and (B) $A. f(\text{S}^0)$ on the silica surface.
625 Histograms (C and D) of the Young's modulus obtained at various salt concentrations.

626

627 **Figure 7.** Effect of loading forces on the approaching curves of $A. f(\text{Fe}^{2+})$ on silica surface in
628 0.001 M KCl solution.

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630 **Captures to Tables**

631 **Table 1.** Effect of salt concentration on the average properties (mean \pm SE) of adhesion events

632 between bacteria and substrates

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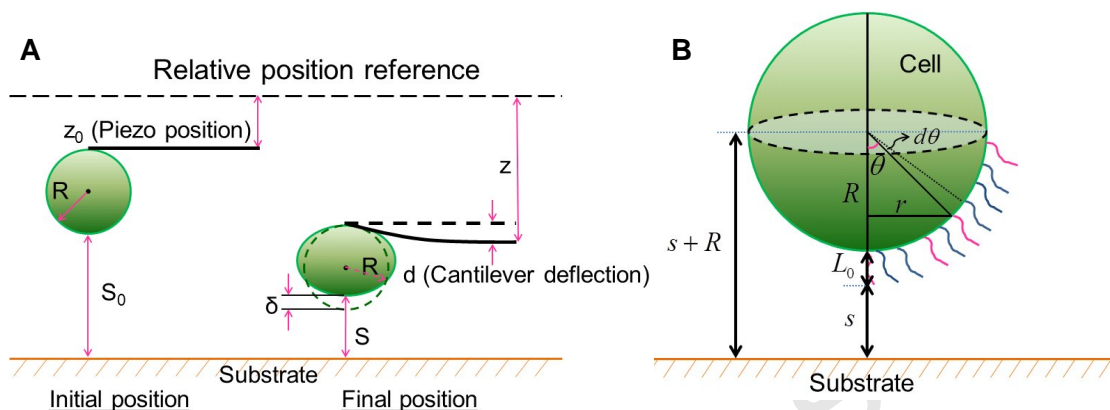
633 **Table 1.** Effect of salt concentration on the average properties (mean±SE) of adhesion events
 634 between bacteria and substrates

Bacteria	Substrates	Salt conc. (M)	Adhesion events (%)	Mean adhesion force (nN)	Stair-step unfolding force (nN)	Snap-off distance (nm)	
						Mean	Range
<i>A. f</i> (Fe ²⁺)	Silica	0.001	16	0.057±0.017	0.041±0.007	1498±332	308~2187
		0.01	28	0.074±0.023	0.065±0.016	558±275	40~1644
		0.1	42	0.026±0.013	0.048±0.005	254±62	78~627
		1	47	0.088±0.021	0.091±0.015	611±184	131~1679
	Pyrite	0.001	22	0.044±0.018	0.038±0.007	611±221	33~1625
		0.01	29	0.057±0.022	0.066±0.012	747±302	211~1983
		0.1	100	0.085±0.031	0.081±0.075	942±198	89~1840
		1	22	0.061±0.026	0.089±0.036	536±220	237~1853
<i>A. f</i> (S ⁰)	Silica	0.001	48	0.211±0.117	NA	267±36	76~401
		0.01	100	0.352±0.075	NA	222±44	58~351
		0.1	86	0.287±0.084	NA	326±102	73~800
		1	72	0.113±0.076	NA	537±176	102~ 1024
	Pyrite	0.001	53	0.069±0.036	NA	260±93	46~625
		0.01	75	0.053±0.021	NA	409±130	92~1154
		0.1	100	0.081±0.024	NA	572±182	34~1042
		1	86	0.082±0.035	NA	499±251	32~1688

635 NA: value was not estimated from this study.

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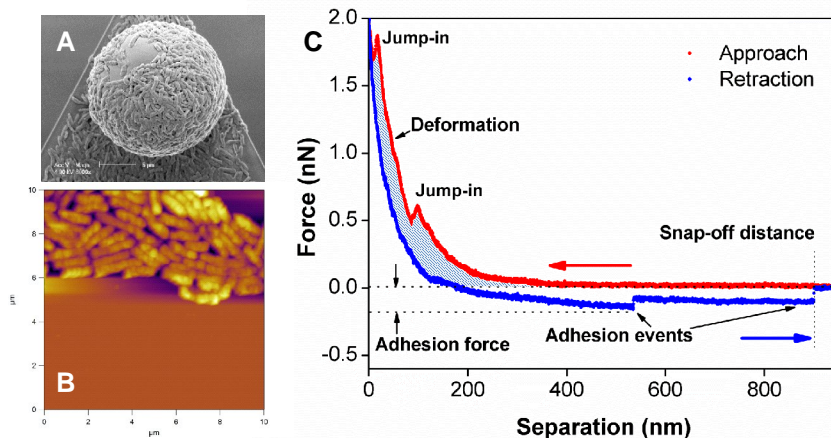
637

638 **Figure 1.** (A) Schematic of the relative positions of a bacterial probe with cell deformation
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 641 deflection; δ is the deformation of the cell; s_0 is the initial separation between sample and
 642 the substrate surface; s is the actual separation distance between the two surfaces. (B)
 643 Schematic illustration of the abbreviations used to calculate the interaction force from the
 644 pressure (force per unit area). The cell is considered as a microsphere.

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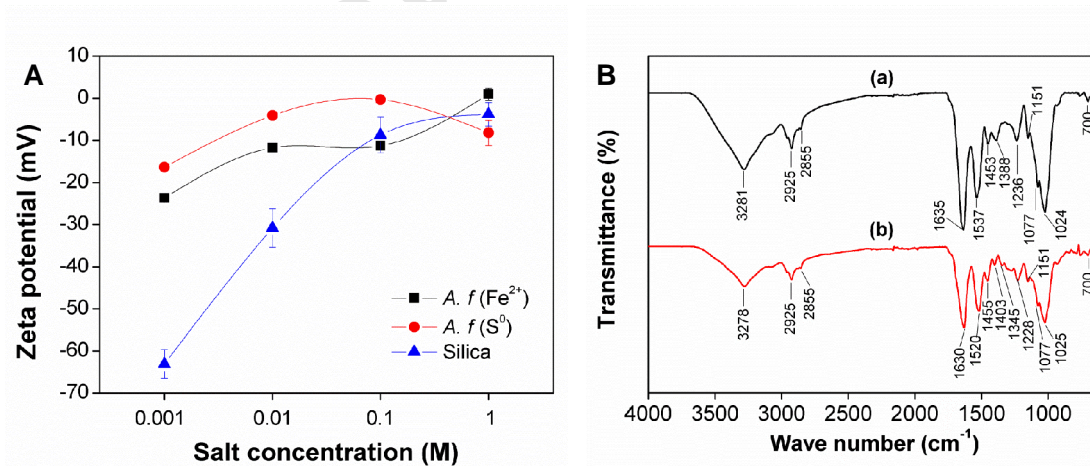
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647

648 **Figure 2.** (A) A representative SEM image of a bacterial probe of *A. f* (Fe^{2+}). (B) An AFM
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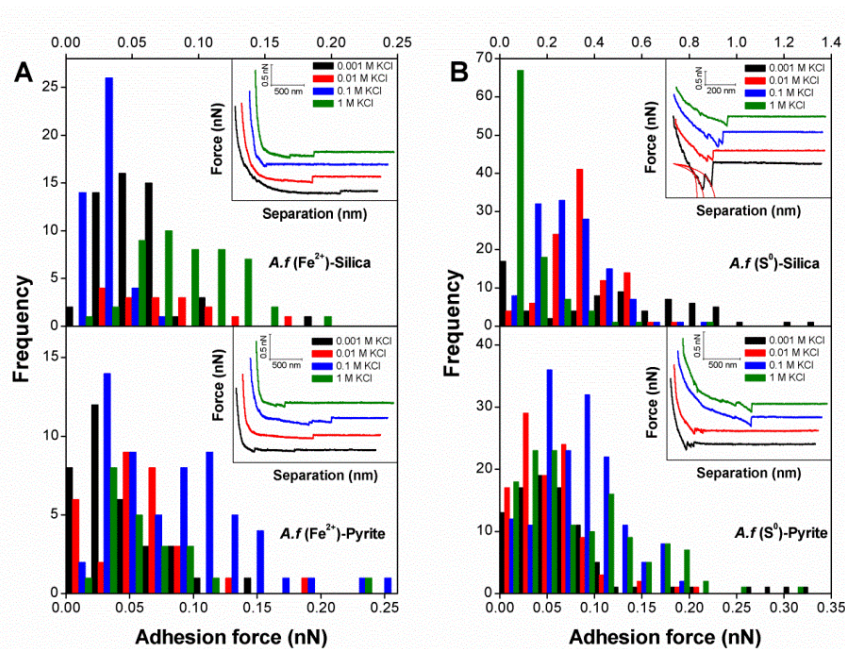


652

653 **Figure 3.** (A) Zeta potentials of *A. f* and silica wafers under various salt concentrations. (B)
 654 FTIR spectra of *A. f* cells cultured with (a) ferrous ions and (b) elemental sulfur.

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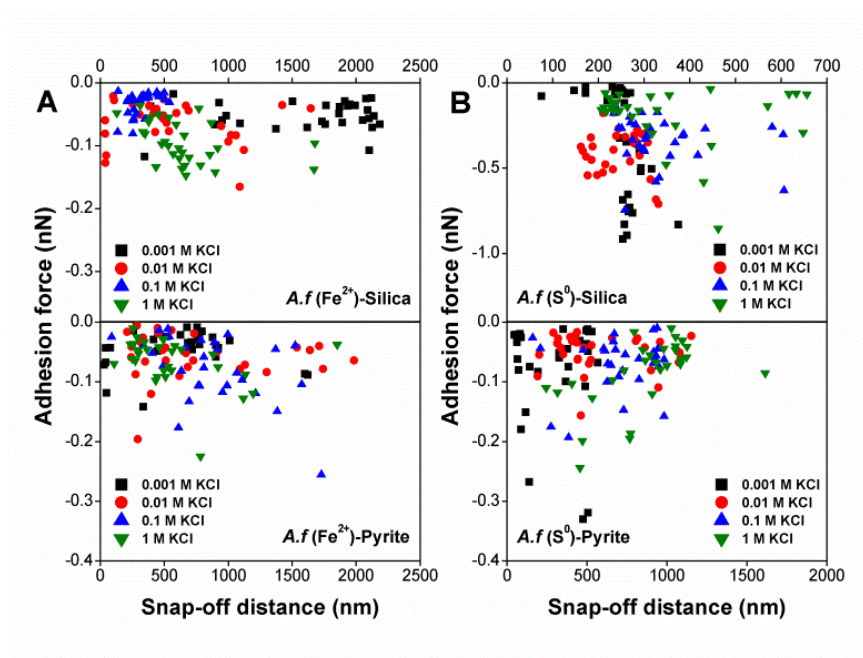
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657 **Figure 4.** Histograms of adhesion forces and representative retraction curves (inset) obtained
 658 from (A) *A. f* (Fe^{2+}) and (B) *A. f* (S^0) at various salt concentrations. Red lines on retraction
 659 curves of sulfur-grown *A. f* show that the adhesion forces are well-fitted by the wormlike
 660 chain model described by Eq. (6).

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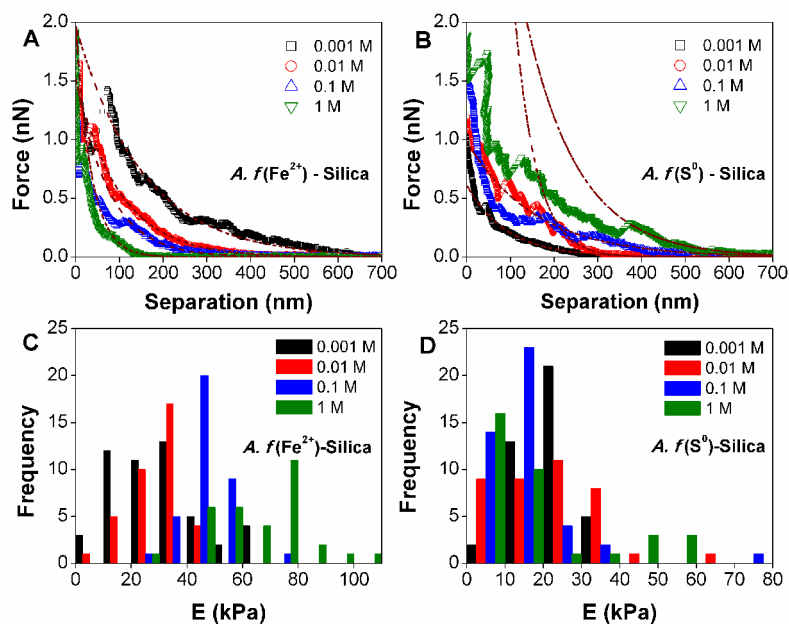
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663 **Figure 5.** Scatter plots of the adhesion forces versus snap-off distances measured between the
664 substrates and (A) *A. f* (Fe^{2+}) or (B) *A. f* (S^0). Data points were randomly collected from 30
665 different retraction curves measured for each salt concentration.

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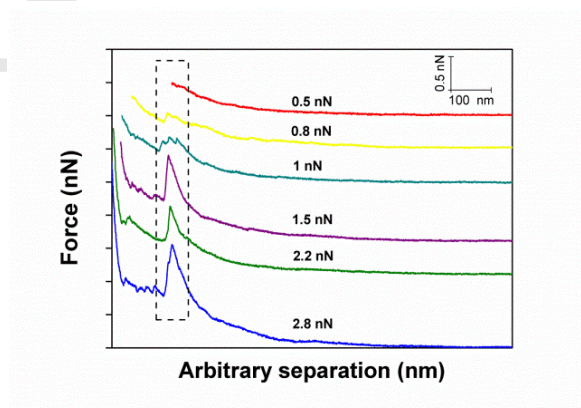
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669 **Figure 6.** Representative approaching curves fitted with steric model showing the effect of
 670 ionic strength on approaching curves of (A) $A. f(\text{Fe}^{2+})$ and (B) $A. f(\text{S}^0)$ on the silica surface.
 671 Histograms (C and D) of the Young's modulus obtained at various salt concentrations.

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674 **Figure 7.** Effect of loading forces on the approaching curves of *A. f* (Fe^{2+}) on silica surface in
675 0.001 M KCl solution.

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