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Phonon anomalies and superconductivity in the Heusler compound YPd_2Sn

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We have studied the structural and electronic properties of YPd_2Sn in the Heusler structure using a generalized gradient approximation of the density functional theory and the *ab initio* pseudopotential method. The electronic results indicate that the density of states at the Fermi level is primarily derived from Pd d states, which hybridize with Y d and Sn p states. Using our structural and electronic results, phonons and electron-phonon interactions have been studied by employing a linear response approach based on the density functional theory. Phonon anomalies have been observed for transverse acoustic branches along the [110] direction. This anomalous dispersion is merely a consequence of the strong coupling. By integrating the Eliashberg spectral function, the average electron-phonon coupling parameter is found to be $\lambda = 0.99$. Using this value, the superconducting critical temperature is calculated to be 4.12 K, in good accordance with the recent experimental value of 4.7 K. © 2014 AIP Publishing LLC. [<http://dx.doi.org/10.1063/1.4887355>]

I. INTRODUCTION

Heusler compounds have attracted a great deal of attention in recent years since they are covenanting for several handy applications in the field of smart materials and spintronics.^{1–7} It should be mentioned that the area of spintronics is one of the most quickly broadening areas of nanoscience and technology since the incorporation of the electron's spin provides an additional degree of freedom to be used for information processing in nanodevices.⁸ Outside of spintronics, Heusler compounds with high spin polarization, high Curie temperature and concurrently a low saturation magnetization are convenient for ultrahigh density magnetic memory storage devices.⁷ In addition to their magnetic properties, superconductivity was discovered for several Heusler compounds. In 1982, superconductivity was discovered for the cubic Heusler compound YPd_2Sn with the transition temperature of 3.72 in the experimental work of Ishikawa.⁹ This discovery has motivated experimentalists to investigate its structural and superconducting properties. In 1984, Johnson and Shelton¹⁰ observed that the application of hydrostatic pressure linearly depressed the superconducting transition temperature with a typical value of $\frac{dT_c}{dP} \sim 10^{-5}$ K/bar for the Heusler compounds RPd_2X (R = Sc, Y, Tm, Yb, and Lu and X = Sn, Pb). Following this study, Malik and co-workers showed that the Heusler compound YPd_2Sn is superconducting with a transition temperature of 4.55 K.^{11,12} Heat-capacity measurement¹³ indicated that the transition temperature of this Heusler compound is between 5.04 and 5.26 K. Recently, the lattice parameter of YPd_2Sn has been measured in the experimental work of Höting and co-workers.¹⁴ Superconducting properties of YPd_2Sn have been presented in the experimental work of Klimczuk *et al.*¹⁵ Following this work, Saadaoui and co-workers¹⁶ have reported on muon-spin rotation and relaxation (μSR) and ¹¹⁹Sn nuclear magnetic resonance (NMR) measurements to study the microscopic superconducting and magnetic properties of the Heusler compound YPd_2Sn . Although

considerable improvement has been done in experimental explanation of the superconducting properties of the Heusler compound YPd_2Sn , no systematic theoretical attempt has been made to investigate structural, electronic, and superconducting properties of this material so far. Such studies must be done, since phonons play the role of bringing about the coupling between electrons to form Cooper pairs, which are amenable for superconductivity, as explained in the BCS theory.

In this paper, we have carried out *ab initio* calculations of the ground state and electronic properties of YPd_2Sn in the Heusler structure. The electronic properties are investigated by using the plane wave pseudopotential method within the generalized gradient approximation of the density functional scheme. This is followed by the application of a linear response scheme¹⁷ for the calculation of phonon dispersion curves and phonon density of states. Atomic displacement patterns and electron phonon coupling parameters for zone-centre phonon modes have been presented and discussed in detail. Furthermore, the Eliashberg spectral function and the average electron-phonon coupling constant λ are calculated. We have compared the Eliashberg spectral function with the corresponding phonon density of states (DOS) in detail. This comparison shows that Pd-related phonon modes make a large contribution to the average electron-phonon coupling constant as a result of the considerable existence of Pd d electrons at the Fermi energy. Using the calculated value of λ , the superconducting critical temperature and the electronic specific heat coefficient γ are calculated. Finally, the calculated superconducting parameters are favourably compared with their corresponding values in the recent experimental work of Klimczuk and co-workers.¹⁵

II. THEORY

The QUANTUM ESPRESSO code¹⁷ employs a plane-wave basis set of the expansion of the Kohn-Sham orbitals.¹⁸ The exchange-correlation energy of electrons is defined in the generalized gradient density functional approximation (GGA)

using the Perdew-Burke-Erzerhof expression.¹⁹ The electron-ion interaction is described by using norm-conserving pseudo-potentials.^{20,21} The maximum plane wave cutoff energy is taken as 60 Ry while the electronic charge density is expanded in a basis cut off up to 240 Ry. Self-consistent solutions of the Kohn-Sham equations¹⁸ were obtained by employing a set of Monkhorst-Pack special \mathbf{k} points²² within the Brillouin zone. The numerical integration of the Brillouin zone (BZ) for the total-energy calculations is performed by using a $16 \times 16 \times 16$ Monkhorst-Pack \mathbf{k} -point sampling.²² On the other hand, the \mathbf{k} -point mesh that is used for the electronic calculations in the BZ is chosen as $32 \times 32 \times 32$.

The Monkhorst-Pack $16 \times 16 \times 16$ \mathbf{k} -points grid within the BZ is used for the phonon calculations. We calculate the phonon frequencies and atomic displacements using the linear response method,¹⁷ which avoids the use of supercells and allows calculation of the dynamical matrix at arbitrary phonon \mathbf{q} vectors. A static linear response of the valence electrons is taken into account in terms of the variation of the external potential corresponding to periodic displacements of the atoms in the unit cell. The screening of the electronic system in response to the displacement of the atoms is considered in a self consistent manner. Twenty nine dynamical matrices are calculated using a mesh of (8, 8, 8) reciprocal lattice divisions in order to obtain full phonon dispersion relations. These dynamical matrices are Fourier transformed to obtain the full phonon spectrum and density of states. We estimate that the phonon frequencies are accurate to within 0.1 THz for the present choice of the kinetic energy cutoff and the special \mathbf{k} points.

The GGA electronic calculations together with the linear response method for the determination of the phonon dispersion relations, electron-phonon coupling, and solution of the Eliashberg equations have made it possible to obtain transition temperatures and study in detail electron-phonon interaction. The Eliashberg electron-phonon spectral function $\alpha^2F(\omega)$ can be presented as^{23–25}

$$\alpha^2F(\omega) = \frac{1}{2\pi N(E_F)} \sum_{\mathbf{qj}} \frac{\gamma_{\mathbf{qj}}}{\hbar\omega_{\mathbf{qj}}} \delta(\omega - \omega_{\mathbf{qj}}), \quad (1)$$

where $N(E_F)$ is the electronic density of states at the Fermi level and $\gamma_{\mathbf{qj}}$ indicates the phonon line-width. When the electron energies around the Fermi level are linear in the range of phonon energies, the phonon line-width is given by the Fermi's "golden rule" formula^{24,25}

$$\gamma_{\mathbf{qj}} = 2\pi\omega_{\mathbf{qj}} \sum_{\mathbf{k}nm} |g_{(\mathbf{k}+\mathbf{q})m;\mathbf{k}n}^{\mathbf{qj}}|^2 \delta(\varepsilon_{\mathbf{k}n} - \varepsilon_F) \delta(\varepsilon_{(\mathbf{k}+\mathbf{q})m} - \varepsilon_F), \quad (2)$$

where the Dirac delta functions express energy conservation conditions, and g is the electron-phonon matrix element. Finally, we are able to determine the electron-phonon coupling parameter (λ) from $\alpha^2F(\omega)$, which is a good measure of the overall strength of the electron-phonon interaction; it is given by

$$\lambda = \int \lambda(\omega) = 2 \int \frac{\alpha^2F(\omega)}{\omega} d\omega. \quad (3)$$

The summations in Eqs. (1) and (2) are made using a dense mesh ($(32 \times 32 \times 32)$ Monkhorst-Pack mesh) of \mathbf{k} points in the BZ. The Dirac delta functions in this equation are replaced with a Gaussian function of width 0.01 Ry.

III. RESULTS

A. Structural and electronic properties

YPd₂Sn adopts the cubic Heusler structure with the space group $Fm\bar{3}m$ (see Fig. 1(a)), which includes one formula unit in the primitive unit cell with Wyckoff positions of atoms at Y (4b) $(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$, Pd (8c) $(\frac{1}{4}, \frac{1}{4}, \frac{1}{4})$, and Sn (4a) (0, 0, 0). For the calculation of the ground state properties such as the equilibrium lattice constant a , the bulk modulus B , and its pressure derivative B' , we have calculated the total energy as a function of lattice constant. The obtained relationship between total energy and lattice constant is displayed in Fig. 1(b). The calculated total energies are fitted to the Murnaghan's equation of state²⁶ to determine the values of the ground state properties. The calculated bulk static properties from our total energy calculations are presented in Table I together with available experimental results. In comparison with the experimental data, we have found that the

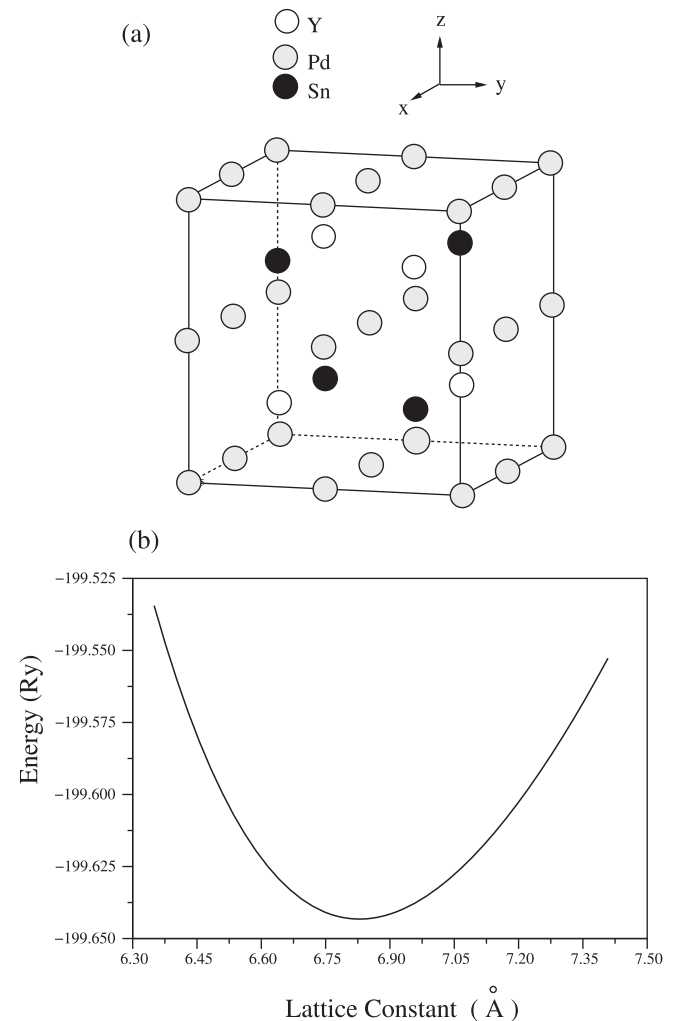


FIG. 1. (a) The cubic Heusler structure of YPd₂Sn. (b) Calculated total energies as a function of lattice constant.

TABLE I. Static properties of the superconducting Heusler compound YPd_2Sn and their comparison with previous experimental and theoretical results.

Source	a (Å)	B (Mbar)	B'
This study(GGA)	6.8289	1.004	4.70
Experimental ^{9,10}	6.7190	1.073	
Experimental ^{11,12}	6.7160		
Experimental ¹³	6.7165		
Experimental ¹⁴	6.7144		
Experimental ¹⁵	6.7160		

calculated lattice constant is slightly larger than the measured ones with around 1.7%. The calculated bulk modulus (B) deviates from its experimental value of 1.073 Mbar within 6%. To our knowledge, there are no other theoretical calculations for the ground state properties of YPd_2Sn in the literature, so more theoretical calculations should be made for comparison.

The calculated energy band structure, at the equilibrium lattice constant for YPd_2Sn , along the high symmetry directions in the Brillouin zone is shown in Fig. 2. The analysis of the electronic structure indicates the metallic nature of this material because conduction and valence bands cross the Fermi level. The lowest band in this figure lying between 9.0 and 7.5 eV below the Fermi level is dominated by the hybridization of Sn 5s and Pd 4p states. This band is well separated from other bands by a gap of 2.5 eV. This band has a parabolic shape along the main symmetry directions of Γ -K, Γ -X, and Γ -L while it is totally flat along other symmetry directions. The next group of occupied bands is from a mixture of Pd, Y, and Sn orbitals. In particular, we have observed a flat conduction band at approximately 0.15 eV above the Fermi level along the Γ -X and Γ -L symmetry directions. This feature is very important for the superconducting properties of YPd_2Sn since electrons near the Fermi surface are strongly involved in the formation of the superconducting state.

The total and partial DOS for YPd_2Sn are provided in Fig. 3. The total DOS has remarkable features with energies of -7.8, -4.8, -4.1, -3.3, -2.8, and -1.1 eV. The first

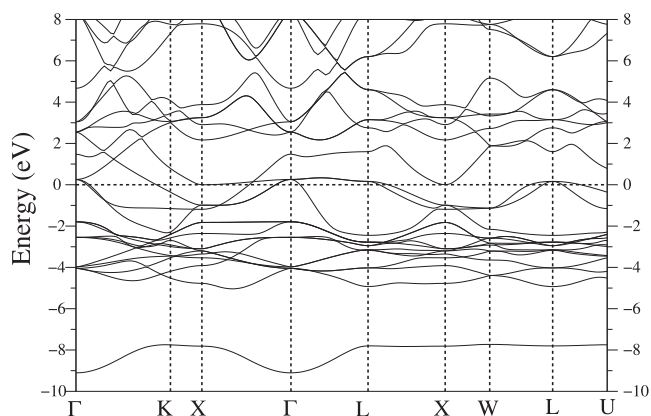


FIG. 2. Electronic band structure for YPd_2Sn along the high symmetry lines of the Brillouin zone. Zero energy represents the Fermi level.

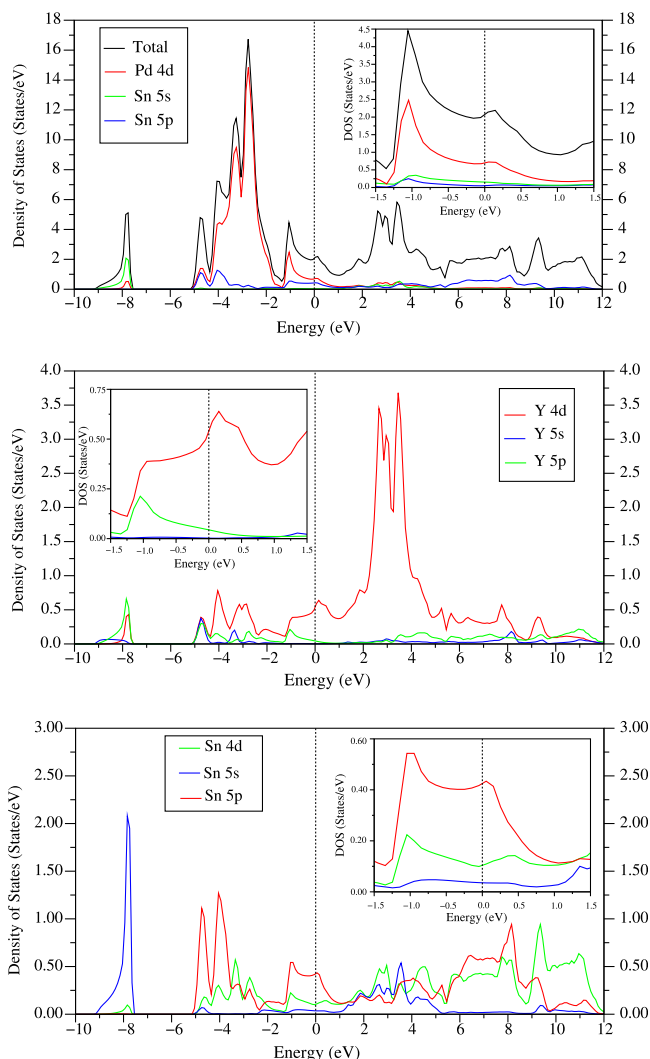


FIG. 3. Total and partial density of states for YPd_2Sn . The Fermi energy corresponds to 0 eV. The insets show results in a small energy range around the Fermi energy.

peak is created by the 5s states of Sn with a significant contribution from Pd 4p states. The second one originates mainly from the mixed Pd 4d and Sn 5p states with some contributions from Pd 4s and Pd 4p states. The main contribution to the peak at -4.1 eV comes from the 4d electrons of Pd atoms with a small hybridization of Sn 5p states. The peak at -3.3 eV is mainly characterized by Pd 4d states with much lesser contributions coming from Y 4d, Sn 4d, Sn 4p, and Pd 4p states. The sharpest peak at -2.8 eV arises from the 4d electrons of Pd with negligible contributions from other electronic states. The last peak below the Fermi level is located at -1.1 eV. This peak is mainly due to Pd 4d states with smaller contributions coming from other electronic states. The density of states at the Fermi level ($N(E_F)$) amounts to be 2.06 States/eV which can be compared with its experimental value of 2.23 State/eV in the work of Klimczuk *et al.*¹⁵ The orbital analysis of the DOS shows that mainly the Pd 4d (34%), Y 4d (25%), and Sn 5p (20%) states are contributing to the total value of DOS at the Fermi level. Finally, the unoccupied DOS are mainly dominated by the 4d electrons of Y atom.

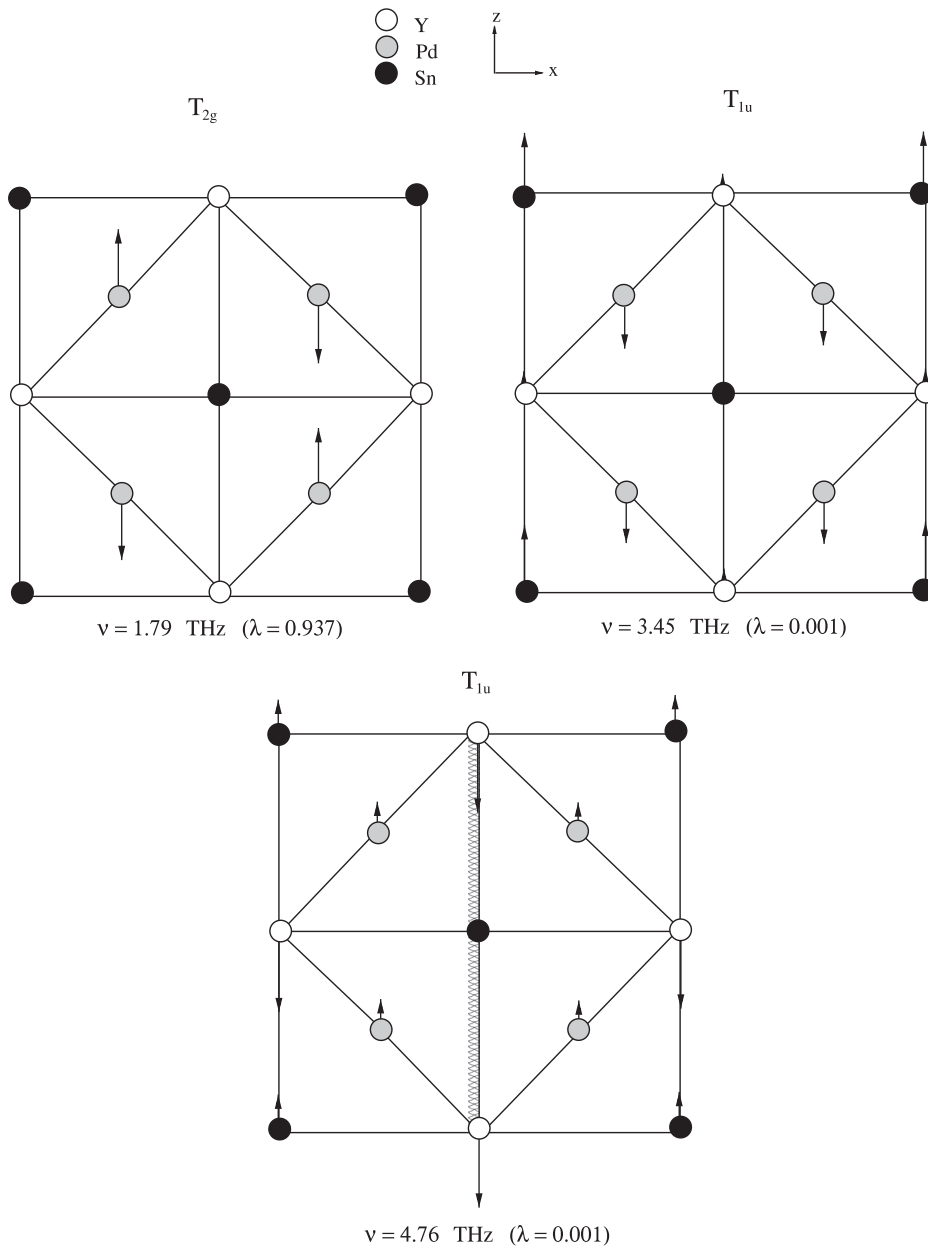


FIG. 4. Eigenvectors representations of zone-centre phonon modes for YPd_2Sn . The electron-phonon coupling parameter for these phonon modes is also given.

B. Electron-phonon interaction and superconductivity

The primitive unit cell of the YPd_2Sn Heusler structure has one formula unit (total four atoms), giving a total of 12 phonon branches, including three acoustic branches and nine optical branches. The optical phonon modes correspond to the following irreducible representations at the zone-centre:

$$\Gamma = T_{2g} + 2T_{1u}.$$

Each T vibration is threefold degenerate. The T_{2g} mode is Raman active while the T_{1u} modes are infrared active. We have illustrated the atomic displacement pattern of these zone-centre phonon modes together with their electron-phonon coupling parameters in Fig. 4. The opposing vibrations of Pd atoms generate the T_{2g} phonon mode at 1.79 THz. It is important to note that the largest electron-phonon coupling parameter ($\lambda = 0.937$) is found for this phonon mode. This observation is not surprising since the largest contribution to

the $N(E_F)$ comes from the d states of Pd atoms. We have to mention that the total electron-phonon coupling parameter for this phonon mode is three times of this value due to mode degeneracy. The first T_{1u} mode is due the opposing motion of Sn and Pd atoms with negligible contribution coming from Y atoms. The second one includes atomic vibrations of Y atoms against Pd and Sn atoms with maximum contribution coming from Y atoms due to their light mass.

In Fig. 5(a), the calculated phonon-dispersion relations are illustrated along several high-symmetry directions in the face-centered cubic zone. As can be seen, all phonon modes have positive frequencies, and there are no phonon branches with dispersions that dip towards the zero frequency. This picture confirms that the cubic Heusler YPd_2Sn is dynamically stable. The phonon dispersion curves can be separated into two apparent regions: a low-frequency region between 0 and 3.7 THz and a high-frequency region between 4.5 and 5.3 THz. A narrow gap of about 0.8 THz between these

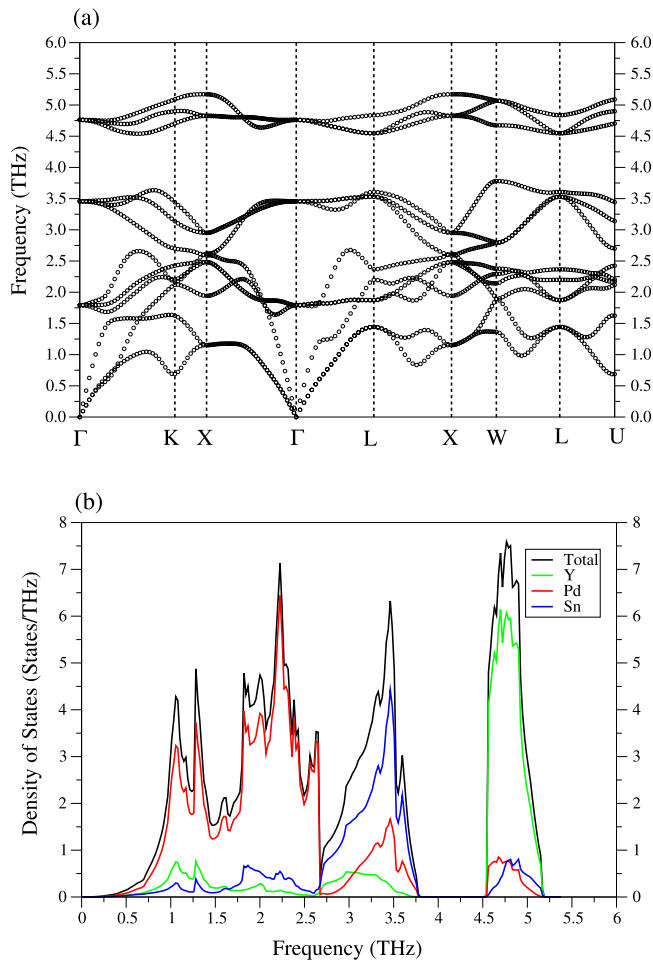


FIG. 5. (a) Calculated phonon dispersion relations for the Heusler compound YPd_2Sn along symmetry lines of the face-centered cubic Brillouin zone. (b) The calculated phonon total and atom-projected density of states.

regions comes from the mass difference between the different types of atoms in the unit cell. There are nine phonon branches extending up to 3.7 THz in the low-frequency region. All phonon branches in this region exhibit considerable amount of dispersion. The most striking feature of our phonon spectrum is the pronounced minimum of transverse acoustic branch (TA) along Γ -K, L-X, W-L, and L-U symmetry directions. We strongly believe that these phonon anomalies play an important role in developing superconductivity in the Heusler compound YPd_2Sn since a similar observation has been made for several superconductors.²⁷⁻³¹ There are three optical branches in the high-frequency region of the phonon spectrum. These phonon branches are less dispersive than the optical phonon branches in the low-frequency region. The nature of the phonon dispersion relations can be understood more clearly by the total and partial phonon DOS in Fig. 5(b). The partial DOS illustrates a dominance of Pd atoms in the low energy range up to 2.7 THz. The intermediate frequency region between 2.7 and 3.7 THz is mainly dominated by the vibrations of Sn atoms with smaller contributions coming from the vibrations of remaining atoms. Above the phonon band gap region, the vibrations of Y atoms are very dominant due to their lighter atom mass than the other two atoms.

As we have pointed out above, the most interesting feature of the phonon dispersion relations for YPd_2Sn is the anomalous behaviour exhibited by the lower acoustic branch along the Γ -K symmetry direction. In order to find a relationship between the anomalous phonon dispersion and electron-phonon interaction, we have calculated the mode-dependent electron-phonon coupling parameter $\lambda_{\mathbf{q},j}$ for the low-frequency phonon branches. Exciting features can be noted along the two main symmetry directions [110] and [111] in Fig. 6. First, along the [110] direction, the T_{2g} mode splits into three phonon branches: two transverse optical (TO₁ and TO₂) and one longitudinal optical (LO) branch. Along [111] direction, transverse branches are doubly degenerate. The electron-phonon coupling parameters of these optical branches decreases rapidly with increasing \mathbf{q} wave vector along both the symmetry directions. Thus, these phonon branches make a large contribution to electron-phonon interaction in YPd_2Sn only close to zone-centre. Second, the TA_2 branch lies below the TA_1 branch in the some part of Γ -K direction. This phonon anomaly is found to occur between $\mathbf{q} = \frac{2\pi}{a}(0.10, 0.10, 0.00)$ and $\mathbf{q} = \frac{2\pi}{a}(0.25, 0.25, 0.00)$. As a consequence of this phonon anomaly, the electron-phonon coupling parameter of this branch hardens rapidly with increasing wave vector up to $\mathbf{q} = \frac{2\pi}{a}(0.25, 0.25, 0.00)$. In particular, the most notable feature is that the electron-phonon coupling parameter of the TA_2 branch reaches a peak value of 1.20 at $\mathbf{q} = \frac{2\pi}{a}(0.225, 0.225, 0.000)$. Third, the branch TA_1 starts to show a negative dispersion (phonon anomaly) beyond $\mathbf{q}_0 = \frac{2\pi}{a}(0.575, 0.575, 0.00)$. Thus, this branch acquires a clear “dip” close to the K zone-boundary. Due to this phonon anomaly, the electron-phonon coupling parameter of this phonon branch increases with increasing \mathbf{q} wave vector beyond \mathbf{q}_0 up to the K point. The peak value of electron-phonon coupling parameter for this phonon branch is found to be 1.12 at the K point. Thus, it will be interesting to analyze the atomic displacement pattern of TA_1 phonon mode at the K point. Fig. 7 presents the eigenvectors representation of this phonon mode. As can be seen from this figure, this phonon mode is mainly dominated by the vibrations of Pd atoms with some contributions coming from other atoms in the unit cell. Different from the [110] direction, the doubly degenerate TA branch behaves normally and shows a positive dispersion with increasing \mathbf{q} wave vector along the [111] direction. However, the LA branch makes a larger contribution to the electron-phonon interaction than the TA branches. In particular, the electron-phonon coupling parameter of LA branch takes its largest value of 0.63 at $\mathbf{q} = \frac{2\pi}{a}(0.20, 0.20, 0.20)$.

In order to sight the strengths with which different modes of the atomic motion couple to electrons, and thus are skillful in affecting superconductivity in YPd_2Sn most, we have displayed the Eliashberg spectral function together with the variation of the electron-phonon coupling parameter with increasing frequency $\lambda(\omega)$ in Fig. 8. By integrating the Eliashberg spectral function using Eq. (3), the average electron-phonon coupling parameter is found to be $\lambda = 0.99$. A comparison of the Eliashberg spectral function $\alpha^2F(\omega)$ with the phonon density of states (see Fig. 5(b)) shows that the features in the electron-phonon spectral function are

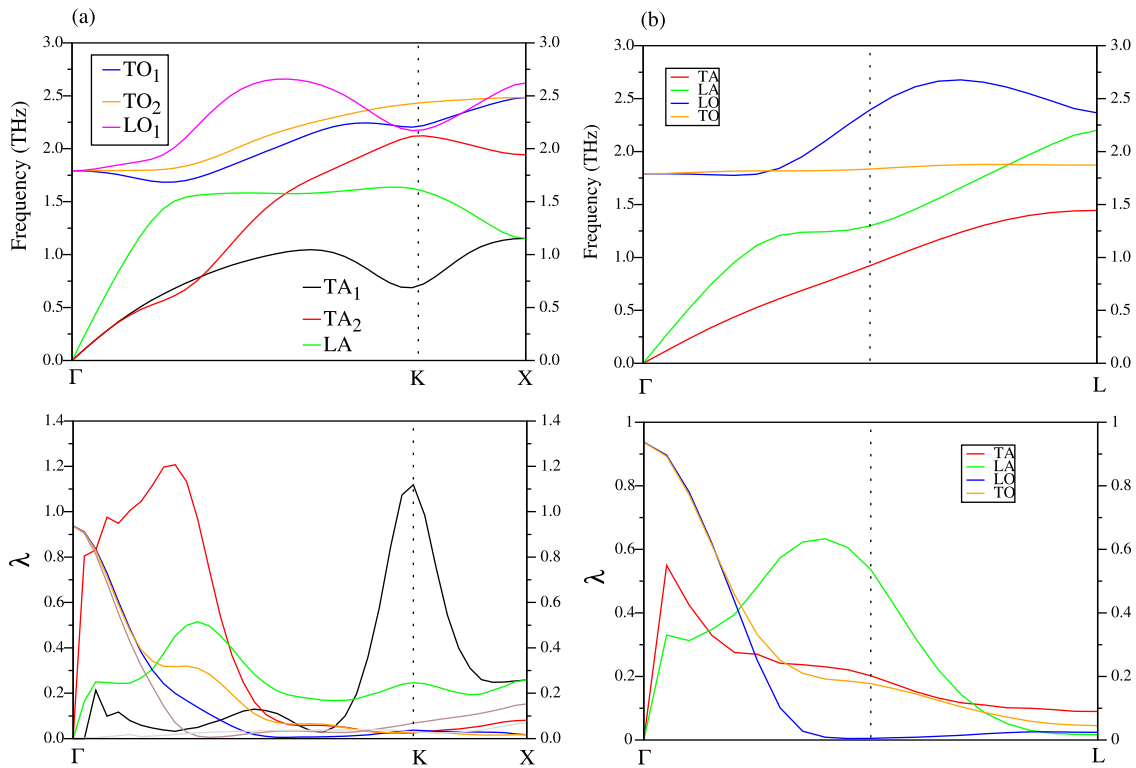


FIG. 6. The dispersion of the lowest six phonon branches with their corresponding electron-phonon coupling parameter (a) along the [110] direction and (b) along the [111] direction.

considerably decreased relative to the corresponding features in the phonon density of states above the phonon band gap region. Thus, we can say that phonon branches above the gap region make very small contribution to the average electron-phonon coupling parameter (their contribution is within 4%). Such small contribution from high frequency phonons can be connected with the factor $\frac{1}{\omega}$ in Eq. (3).

However, phonon modes below 2.7 THz contribute about 88% towards λ . Our partial phonon density of states shows that the Pd-related phonon modes exist below 2.7 THz. This is an expected result due to considerable contribution of Pd d states to $N(E_F)$.

With λ computed, the transition temperature (T_C) for YPd₂Sn can be calculated from the Allen-Dynes formula^{24,25}

$$T_C = \frac{\omega_{ln}}{1.2} \exp\left(-\frac{1.04(1 + \lambda)}{\lambda - \mu^*(1 + 0.62\lambda)}\right), \quad (4)$$

where the logarithmically averaged phonon frequency ω_{ln} is given as

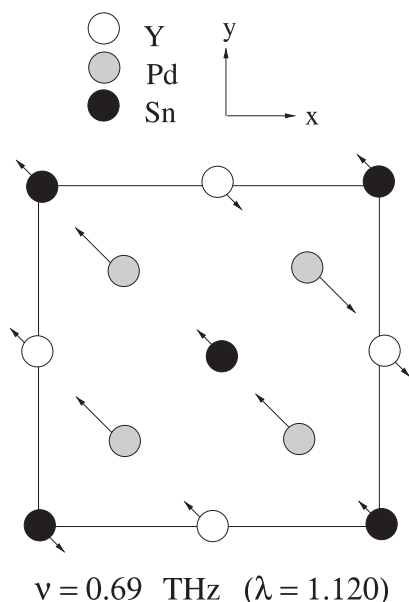


FIG. 7. Eigenvectors representations of first transverse acoustic phonon mode (TA₁) at the K point for YPd₂Sn. The electron-phonon coupling parameter for this phonon is also presented.

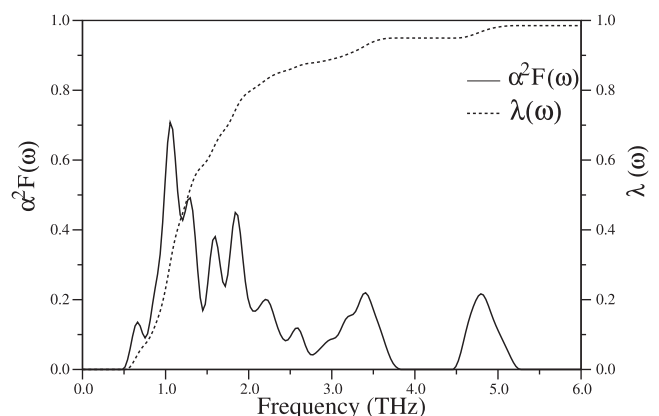


FIG. 8. The calculated electron-phonon spectral function $\alpha^2F(\omega)$ (left axis) and the variation of the electron-phonon coupling parameter ($\lambda(\omega)$) with increasing frequency $\lambda(\omega)$.

$$\omega_{\text{ln}} = \exp\left(2\lambda^{-1} \int_0^{\infty} \frac{d\omega}{\omega} \alpha^2 F(\omega) \ln \omega\right). \quad (5)$$

Inserting the calculated value of λ into Eq. (5), the value of ω_{ln} is found to be 1.47 THz. μ^* in Eq. (4) is the screened Coulomb pseudopotential parameter, which takes a value between 0.1 and 0.16.^{24,25} Using the Allen-Dynes formula and taking typical values of $\mu^* = 0.10, 0.11, 0.12, 0.13, 0.14, 0.15,$ and 0.16 , we obtain $T_C = 4.81, 4.57, 4.32, 4.11, 3.89, 3.66,$ and 3.45 , respectively. The average of these values is 4.12 K, which is in acceptable agreement with the experimental values of 3.74,⁹ 4.55,¹² 5.04,¹³ and 4.7 K.¹⁵

The calculated values of the superconducting state parameters for YPd₂Sn are compared with their corresponding experimental values in Table II. In general, our results are in qualitative agreement with available experimental results.^{9,10,12,13,15} In particular, the calculated value of $N(E_F)$ is slightly lower than its experimental value of 2.23 (Ref. 15) while the calculated value of λ is higher than its experimental value of 0.70.¹⁵ When making this comparison, it should be noted that the computed value of $N(E_F)$ comes directly from our band structure results, while the experimentally quoted value is estimated from the measured Sommerfeld coefficient of the specific heat and the estimate of λ . Also, our calculated value of $\lambda = 0.96$ is higher than the experimental value of 0.70. We have used the *quasi-static approximation*³² to calculate the Debye temperature (Θ) of YPd₂Sn.^{33,34} The value of Θ is found to be 199 K, which is in an excellent agreement with the recent experimental value of 210 K in the work of Klimczuk *et al.*¹⁵ The electron-phonon coupling constant λ also enters the electronic specific heat coefficient γ , which is given as

$$\gamma = \frac{1}{3} \pi^2 k_B^2 N(E_F) (1 + \lambda). \quad (6)$$

From the above equation, the value of γ is obtained to be 9.6 mJ mol⁻¹ K⁻², which compares very well with its experimental value of 9.2 mJ mol⁻¹ K⁻².¹⁵ Finally, the BCS energy gap ($\Delta = 1.76 k_B T_C$) is found to be 0.49 K, which can be compared with the corresponding experimental value of 0.83 K in the work of Klimczuk *et al.*¹⁵ The agreement between our results and recent experimental results¹⁵ signs the good quality of our GGA calculations. As a result, our calculations indicate that the Heusler compound YPd₂Sn is a conventional phonon-mediated superconductor with the strong electron-phonon coupling strength of 0.99.

TABLE II. The comparison of superconducting state parameters for YPd₂Sn. We have used $\Delta = 1.76 k_B T_C$.

Source	$N(E_F)$	λ	ω_{ln} (K)	Θ (K)	T_C (K)	$\left(\frac{\gamma}{\text{mJ mol K}^2}\right)$	Δ (meV)
This work	2.06	0.96	70.96	199	4.12	9.5	0.49
Experimental ^{9,10}				165	3.72		
Experimental ¹²					4.55		
Experimental ¹³					5.04		
Experimental ¹⁵	2.23	0.70		210	4.7	9.2	0.83

IV. SUMMARY

We have studied the static and electronic properties of the Heusler compound YPd₂Sn applying density functional theory within a generalised gradient approximation. Our results compare very well with existing experimental results for lattice constant and bulk modulus. Our electronic calculations validate the metallic nature of this compound. From the analysis of the total and partial electronic density of states, we have observed that the states around the Fermi level are contributed by Pd 4d, Y 4d, and Sn 5p states with much smaller contribution from other electronic states.

Using our calculated lattice constant and electronic structure, lattice dynamics of YPd₂Sn has been studied by employing a linear response approach based on density functional perturbation theory. This compound is dynamically stable, as no instabilities in the phonon dispersion curves have been found. However, a striking feature in the phonon spectrum of this material is the pronounced minimum of transverse acoustic branches along the symmetry direction [110]. From the analysis of phonon density of states, we have found that Pd-related vibrations are dominant in the low frequency region up to 2.7 THz. On the other hand, Sn-related vibrations are strong in the middle frequency region between 2.7 and 3.7 THz. Above the phonon band gap region, Y-related vibrations are dominant due to the smaller mass of Y atom.

We have presented a numerically accurate correlation between phonon anomalies and the electron-phonon coupling parameter by studying electron-phonon interaction in this material. Our electron-phonon interaction results show that the electron phonon coupling parameters of the acoustic branches (TA₁ and TA₂) reach their peak values of 1.12 and 1.20, respectively, along the Γ -K direction where phonon anomalies are observed.

The phonon spectrum and the electron-phonon matrix elements are used to calculate the Eliashberg function from which the superconducting properties can be obtained. The deviation of the shape of Eliashberg function from that of phonon density of states confirms that the phonon modes related to Pd-vibrations play a very important role in the process of scattering of electrons. From the Eliashberg function, the average electron-phonon coupling parameter is calculated to be 0.99, indicating reasonably strong interaction between electrons and phonons in this compound. Using the calculated the values of λ and ω_{ln} , the transition temperature (T_C) calculated from the Allen-Dynes formula is found to be 4.12 K, which compares well with the recent experimental value of 4.7 K.

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