

THE ULTRA-VIOLET ABSORPTION SPECTRA  
OF SOME HALOGEN COMPOUNDS IN  
VARIOUS SOLVENTS.

Thesis for the  
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Presented to the  
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by

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ponds/

## I. INTRODUCTION.

The modern theory regards the atom as made up of a positively charged nucleus surrounded by a shell of electrons, the number of electrons being equal to the positive charge on the nucleus. These electrons are supposed to exist in certain definite orbits or energy states, and energy may be absorbed or emitted as one quantum of monochromatic light when an electron passes from one state to another: the difference of the energies of the atom in the initial and final states is equal to the energy of the radiation absorbed or emitted.

For each electron there is a state of minimum energy known as the "ground state", and if the atom collides sufficiently violently with another atom (as at high temperatures) it may have one of its outer electrons raised to one of the states of higher energy, and the energy that the electron gives out in returning from this excited state to the ground state appears as monochromatic light. The light from a glowing gas or vapour is the sum of millions of such processes in the separate atoms.

Certain energy states will be more favoured than others, and some transitions will be frequent and others improbable. To each such transition corresponds/

corresponds a spectrum line, the intensity of which will depend on the above probabilities.

All spectra fall into two main classes: emission spectra and absorption spectra. Emission spectra can be excited in a number of ways (e.g. by violent collision). Absorption spectra are usually obtained by passing the light from a source of continuous spectrum through a cell containing the gas or liquid under investigation and examining the spectrum of the emerging light. From the bright continuous background of colour certain radiations are then missing - having been absorbed - and these constitute the absorption spectrum of the substance. The absorption spectrum of a substance is not necessarily the exact reverse of its emission spectrum, because chiefly those lines occur in the absorption spectrum which correspond to a transition from the ground state to some state of higher energy; in the emission spectrum however there can occur lines corresponding to a transition from one excited state to another.

Unlike the atom, which from a spectroscopic standpoint can only possess electronic energy, a molecule can have three kinds: electronic energy, vibrational energy (of the nuclei) and rotational energy. To each of these types of energy the quantum conditions must be applied. If only changes in the rotation energy take place, the radiation appears in the far infra-red (at/

(at wave-lengths in the neighbourhood of  $100\mu$  ).

If rotational and vibrational changes take place together, the magnitude of the latter raises the frequency of the emitted light into the near infra-red region. If an electronic change occurs simultaneously with these, the spectrum will be emitted in the visible or u. v. region. In a molecule of a substance in the gaseous state, to a first approximation, we may regard these energies as independent and write down the frequency of any band line as the sum of the contributions from these three separate sources. Thus:

$E = E_c + E_v + E_r$  (1) where the order of magnitude is  $E_c > E_v > E_r$

When the energy supplied to the molecule corresponds to wavelengths in the far infra-red the quanta of energy are not large enough to change the electronic or vibrational energy states but are only large enough to change the rotational energy states and so the spectrum of the substance consists of a series of lines the wavelengths of which correspond to the energy differences of the various rotational states of the molecule.

If however the energy supplied is of wavelengths in the near infra-red the quanta are large enough to affect the vibrational energy states. In this case a series of bands (instead of single lines) is obtained, each of which is made up of a group of lines; this/

this is due to the fact that rotational energy changes occur at the same time as the vibrational energy changes. If  $E$  is the energy supplied to the molecule then for each line

$E = E_v + E_r$  where  $E_v$  is the energy of a vibrational change and  $E_r$  is the energy of a rotational change, then for a definite value of  $E_v$  there will be series of values of  $E$  corresponding to different values of  $E_r$ , so that for each value of  $E_v$  possible to the molecule there will appear in the spectrum a band of lines each of which corresponds to a different value of  $E_r$ .

Wave-lengths in the visible and u.v. regions correspond to energy quanta large enough to produce changes in the electronic energy states; to each electronic energy change there corresponds a series of vibrational energy changes and to each vibrational energy change a series of rotational energy changes and the energy  $E$  supplied is given by equation (1). So that the spectrum of a substance in this region consists of a series of bands (each due to an electronic energy change) each made up of a series of smaller bands (due to vibrational energy changes), each of which, in its turn, is made up of a series of lines (due to rotational energy changes).

The foregoing remarks apply to the spectra of substances in the gaseous state. It is found that at constant temperature with increasing pressure of foreign/

foreign inert gas the rotational fine structure disappears and the vibrational bands gradually broaden and finally overlap and just appear as a series of maxima in the electronic band. For the symmetrical broadening of the rotational lines on the addition of an inert gas at high pressure to an absorbing molecular vapour at low partial pressure the Lorentz collision-damping hypothesis gives a good explanation up to a pressure of 10 atmospheres foreign gas. For still higher pressure of foreign gas there is a strong increase in the broadening, which as well as the shift in the edge of the band, and the increase in energy absorbed and the final appearance of new bands are due to effects other than the collision effect. These effects show analogies to the Stark effect although we cannot make definite assertions as to their nature.

The u. v. absorption spectrum of a substance dissolved in a chemically inert solvent consists of a series of maxima which in general show no fine structure. Teves (1) suggests that this non-appearance of fine structure may be explained by assuming that the solvent acts in part, as an inert gas at high pressure. But the problem involves much more than this because it is found that the absorption spectrum of a substance varies to a greater or less degree with the solvent; and the phenomenon of liquefaction is itself a result of intermolecular forces.



The first attempt to systematise the observed variation of an absorption spectrum of a solution with change of solvent was due to Kundt who propounded the rule which states "that the absorption band of a substance is shifted towards the red the greater is the index of refraction of the solvent in which the substance is dissolved".

In 1912 in a paper entitled "A Theory of Fluorescence" Baly and Krulla (2) brought forward the theory that every elementary atom possesses both primary and secondary valencies; in a compound formed by virtue of the primary valencies only the secondary valencies are unsatisfied, so that every atom is the centre of a field of force the nature and strength of which depends on the nature of the secondary valencies, and whose lines of force radiate in every direction. On the atoms coming together to form a molecule the fields of force condense together with the liberation of free energy, and a closed system of force lines is obtained and the free affinities of the atoms constituting the molecule will be considerably reduced; the whole of the free energy will only disappear when there is perfect equality between those of opposite types within the molecule. Where there is not perfect equality a "residual affinity" exists. These "condensed systems" may be unlocked by the absorption/

absorption of light, the light does work against the chemical forces which tend to produce the closed systems, and only those rays are absorbed which suit the particular system. In solutions the residual affinity of the solvent comes into play and this field of force will tend to penetrate the closed force field of the solute molecules, with the result that it will be partly or completely opened, and the amount of opening depends on the relation between the affinities of the solvent and solute and on their relative masses. The greater the difference in affinity the more will the closed system of the solute be opened, so that Kundt's rule should read "that the absorption band of a substance is shifted towards the red the greater the relative and residual affinity of the solvent." \*

This theory of Baly and Krulla contains the nucleus of the modern theory of association. The theory of the absorption of light by dissolved substances is very complicated because we have to consider the solute molecules each surrounded by a large number of solvent molecules. According to the theory of Fajans & Joos (3) if the solvent molecules are polar/

\* In modern quantitative work on the absorption of light by solutions it is usual to plot the log of the extinction coefficient (which is a function of the ratio  $\frac{I}{I_0}$ ) against the wave-length, and a curve with one or more maxima is obtained.

polar i.e. if the "centres of gravity" of the positive and negative charges of the molecule do not coincide they will exert a "deforming" influence on the shell of electrons of the solute molecules which they surround, the electrons being drawn towards the positive part of the solvent molecule and repelled by the negative part. It is very probable that the energy states of the deformed molecule are not the same as those of the undeformed molecule, so that a difference is to be expected between the spectrum of a substance in the gaseous state and in solution in a polar solvent. (3) It might be expected that the amount that the solute molecule is deformed will depend on the strength of the dipole of the solvent, and in many cases in practice a parallelism is found between the change in absorption (as compared with the gaseous state) and the dipole of the solvent. In non-polar solvents, such as hexane, the absorption spectrum is very similar to that in the gaseous state, so that in general the absorption in hexane is taken as that of the normal undeformed molecule, and is used as a standard for comparison of the spectra of the same substance in other solvents.

Scheibe (4) finds that the measured shift of an absorption maximum on change of solvent is not a simple function of the dipole of the solvent molecule, because the measured dipole is made up of constituent parts/

parts which are screened by non-charged groups; the measured dipole just gives the distance apart of the centres of gravity of the positive and negative charges. The effect of a dipole field on a neighbouring molecule depends on the possibility of approach, and for equal charges this can be quite different because of different screening groups.

Scheibe divides the effect of a change of solvent on the absorption spectrum of a dissolved substance into two parts:- (1) the change in height of the band and (2) the displacement of the centre of the band. He attributes these effects to two distinct processes.

Each molecule absorbs the frequency corresponding to the difference in energy between the original and final energy states of its electrons. This energy difference can vary from molecule to molecule in the solution because of (1) variation in the inner states of the molecules (e.g. rotation, vibration of atoms and groups) or (2) change in the value of the field strength of the neighbouring molecules by change in their distance from the absorbing molecule. The strength of absorption at a given wavelength is determined by the number of molecules with the corresponding energy difference in its two states. In liquid and dissolved substances at ordinary temperatures within certain limits there is a continuous variation/

variation in this energy difference. The absorption curve is an image of the frequency of occurrence of the different states for a dissolved substance and varies from solvent to solvent. This effect does not radically alter the shape of the curve, rather it only shows itself in variation of the position of the curve as measured by the position of the maximum absorption, i.e. by the energy difference which appears most often. By changing the field due to the solvent we can change the energy difference of the solute electron states so that the place of maximum absorption changes although the total area of the curve i.e. the total number of molecules absorbing remains the same; this shift is purely due to the effect of the neighbouring molecules i.e. to solvation.

In complicated organic molecules there may be an equilibrium between two tautomeric forms and the solvent may affect this equilibrium and therefore change the height of the curve. There are indications that two forms may exist in other molecules where ordinary conceptions of valency do not permit it. The pure solvent effect gives the shift in the position of the maximum, and the fall in the height is a result of the solvent effect on this equilibrium.

According to the quantum theory the height of the band gives a measure of the frequency of the absorption process, and a change in height can occur in/

in different ways according to the structure of the molecule. Scheibe (as described above) ascribes the change in the position of the band to compound formation between the solvent and solute, the variation being dependent on the dipole characters of the solvent and solute. If the chromophore group of the solution is positively charged the negative parts of the solvent dipoles are in its vicinity, so that the electron "spring" is made more difficult, i.e. requires more energy and the whole electron band shifts towards the ultra violet. If on the other hand, the chromophore group is negative the positive parts of the dipole are surrounding it, the electron spring is facilitated and the whole band shifts towards the red.

Wolf in 1929 in a paper on Dipole moments, Association and U.V. Absorption (5) developed a formula connecting the refractive index of a solution with the special vibrational frequencies of its constituent absorbing groups. But this formula only holds if no new forces are set up in the solution i.e. both solute and solvent are non-polar. If the solvent has a permanent dipole it will deform the electron shell of the solute and both the position and height of the absorption band will be changed. The polarisability of a molecule is a measure of the ease with which its electron shell can be deformed, and so the change in absorption/

absorption due to one particular group in a molecule by the introduction of another can be taken as a measure of the deforming power of the second group.

Wolf also emphasises the fact that if the solute possesses a dipole it will have a deforming effect on the electron shells of the solvent molecules. Owing to the small number of solute molecules compared with the number of solvent molecules this effect may be neglected if the solvent is polar; but this is not always possible in the case of a polar solute and a non-polar solvent. Fuchs & Donle (6) assert that it is possible to divide non-polar solvents into two groups, those with small, and those with large disturbing effect. The former group contains the saturated aliphatic and alicyclic hydrocarbons (e.g. hexane, cyclohexane, decalin), i.e. compounds having no polar groups and low polarisability. The second group, of non-polar solvents with a large influence on the solute, is composed of substances which contain polar groups whose influence is, however balanced out by symmetry, e.g.  $\text{CCl}_4$ ,  $\text{CS}_2$ , dioxane,  $\text{Cl}-\text{C}_6\text{H}_4-\text{Cl}$ ,  $\begin{array}{c} \text{Cl} \quad \text{Cl} \\ \diagdown \quad / \\ \text{C} = \text{C} \\ / \quad \diagdown \\ \text{Cl} \quad \text{Cl} \end{array}$ . These substances are all easily polarised, so that in the presence of a polar solute a dipole is induced in the solvent molecule which then acts as a polar solvent.

Another theory discusses the shift of absorption spectra in terms of energies of solvation rather than of/

of deformed orbits. In heteropolar compounds it is quite understandable that association should take place between the solute molecules, and it has also been shown to occur to some extent in non-polar hydrocarbons such as diphenyl and phenanthrene in solution. From this fact Berger (7) deduces that the permanent dipole in a molecule is not the only factor entering into association. Association of a dissolved substance is the resultant of two attractive forces (1) the mutual attraction between the dissolved molecules, and (2) the attraction between the solute and solvent molecules, and Berger develops formulae for these attractions depending on the polarisabilities of the molecules. Briegleb (8) (9) has shown that not only the mean polarisability but also the anisotropy of the polarisability is of importance in association.

So that for all substances in solution we may assume association between the solute and solvent molecules to a greater or less degree; and this association alters in some way the potential energies of the ground and excited states, which are intimately bound up with change in absorption of light. This is discussed by Iredale and Mills (10) who point out that the association need not always be of the same type. One form may have a greater difference in potential energy between the ground and excited states compared/



compared with the normal molecule, and another form may have associated molecules with a smaller difference in potential energy between the ground and excited states (again referred to the normal molecule). There is also the possibility that the difference in the potential energies of the ground and excited states is the same for the associate as for the normal molecule. These three types may be represented diagrammatically as follows:-

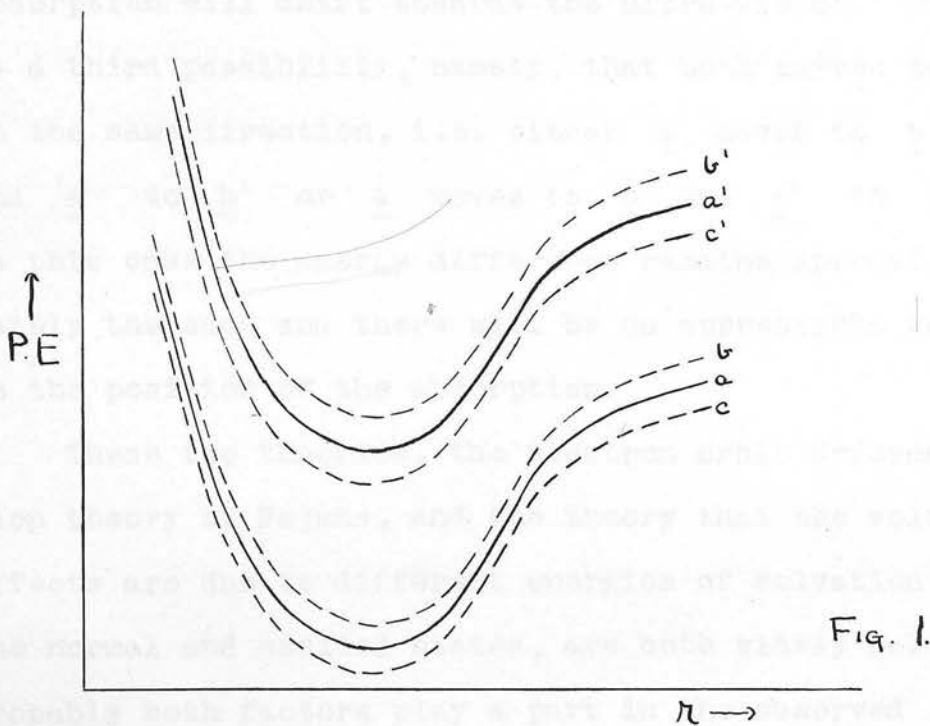


Fig. 1.

The curves a and a' represent the potential energy curves for the normal molecule in the ground and excited states. On dissolving the substance in a solvent the energies of both states will be altered (due to solvation) and the curves will move either up or down but/

but  $r$  for the equilibrium positions will not alter appreciably. They may or may not move in the same direction. If  $\underline{a}$  move to  $\underline{b}$  and  $\underline{a}'$  to  $\underline{c}'$  the difference in the energies of the ground and excited states will be less for this form than for the normal molecule, and less energy will be required to excite the molecule, so the absorption will move towards the red. If on the other hand  $\underline{a}$  move to  $\underline{c}$  and  $\underline{a}'$  to  $\underline{b}'$  the energy difference will increase, more energy will be required to excite the molecule and the absorption will shift towards the ultra-violet. There is a third possibility, namely, that both curves move in the same direction, i.e. either  $\underline{a}$  moves to  $\underline{b}$  and  $\underline{a}'$  to  $\underline{b}'$  or  $\underline{a}$  moves to  $\underline{c}$  and  $\underline{a}'$  to  $\underline{c}'$ . In this case the energy difference remains approximately the same and there will be no appreciable shift in the position of the absorption.

These two theories, the electron orbit deformation theory of Fajans, and the theory that the solvent effects are due to different energies of solvation in the normal and excited states, are both widely held. Probably both factors play a part in the observed shifts of the absorption bands.

The exact shape of the curves  $a\ b\ c$   $a'\ b'\ c'$  depends on the forces controlling the internal vibrations of the molecule in the normal and excited states, and these forces will be altered by the influence which the/  
the/

the solvent introduces. In general one could expect the effect of non-polar solvents to be slight, the effect of solvents containing strong but symmetrical polar groups to be observable, and for strongly polar solvents quite considerable. Their study will involve accurate determination of the shapes of absorption curves as well as of the position of the maximum, and perhaps, attention to the effect of temperature on the shape of the curves. Towards the end of the discussion of the results obtained in the work described in this essay, a tentative explanation of the effect of the solvents is advanced on the basis of the second of the two theories mentioned above.

## II. METHODS OF MEASUREMENT.

### 1. GENERAL REMARKS.

It is well known that when light passes through a medium a certain amount is absorbed. There are two principal laws which govern the amount of light which is absorbed. Lambert's law states that the proportion of light absorbed by a substance is independent of the intensity of the incident light. This law appears to be rigidly true in all experiments which have been made. (Except in such cases as the intensity of the light is sufficient to produce measureable changes in the stationary concentration of the absorbing molecule (59) ). Bunsen and Roscoe assumed only Lambert's law in their definition of "extinction coefficient":- "Let  $I_0$  be the intensity of the incident light which enters the medium,  $I$  the intensity remaining after its subsequent passage through a path length  $d$ , and  $\frac{1}{\alpha}$  the path length, passage through which reduces the light intensity to  $\frac{1}{10}$  th ; then

$$I = I_0 10^{-\alpha d}$$

where  $\alpha$ , a constant depending on the medium, is called by us the 'extinction coefficient'.  $d$  is to be taken in centimetres."

The Bunsen - Roscoe extinction coefficient is, then, the reciprocal of the thickness which is necessary/

necessary to weaken the light to  $\frac{1}{10}$  th of its incident value.

Beer's law states that if an absorbing substance is dissolved in a non-absorbing liquid its absorption of a beam of homogeneous light depends on the number of molecules of the absorbing substance which the beam of light passes through, that is, on the concentration of the solution. If  $I_0$  is the intensity of the light entering the liquid,  $I$  that after passing through  $d$  cm of the liquid, and  $c$  the concentration, Beer's law can be expressed  $I = I_0 \alpha^{-dc}$ . Here  $\alpha$  is a constant of the liquid and is obviously the ratio of the transmitted to the incident light for unit thickness and concentration. We can write this is another way. If  $\alpha_0$  is the extinction coefficient of a solution for unit thickness and concentration, then Beer's law says that  $I = I_0 \cdot 10^{-d\alpha_0 c}$ . In other words the extinction coefficient, for a given wave-length, of a solution of concentration "c" is  $c$  times that of the same solution of unit concentration.

Beer's law has been checked again and again and holds in many instances for dilute solutions with great exactness (11). When, as is very frequently the case, departures occur they indicate an alteration of the absorbing molecules, or the formation of fresh molecular aggregates, as a result of the alteration of/

of concentrations. At small concentrations the absorbing molecules will not be able to influence one another to any appreciable extent, but as the concentration increases and they approach close to one another polarisation and association effects will appear. Halban and Ebert (see above) find deviation from Beer's law for the  $\text{NO}_3^-$  ion at concentrations greater than <sup>about 0.1 molar</sup> and according to Weigert (12) the complex ions  $\text{Cr}_2\text{O}_7^{=}$ ,  $\text{CrO}_4^{=}$ ,  $\text{MnO}_4^-$ ,  $\text{PtCl}_6^{=}$  etc. show similar deviation at high concentrations. If the extinction coefficient of a crystal of known thickness and density be measured and then, successively, the extinction coefficients of solutions of the same substance with different concentrations, it is found that the extinction coefficient increases with dilution from the minimum value given by the solid to a maximum. It would seem that the relation between the maximum extinction coefficient  $K$  in a given solvent, and the extinction coefficient  $k$  at the concentration  $c$  in that solvent is given by

$$\frac{k}{K} = 1 - e^{-\frac{\alpha}{c}}$$

where  $\alpha$  is a constant characteristic of the solvent. For small values of  $c$  the right hand side of this equation approximates to 1 so that  $k = K$  i.e. Beer's law holds for dilute solutions. In all the solutions measured in this work the concentrations of the absorbing substances were small and in the region of/

of complete validity of Beer's law, so that any change in an absorption curve with concentration is assumed to be evidence for a corresponding change in the nature and concentration of the absorbing substance or substances present.

When the term "absorption spectrum" is used in a quantitative sense there is meant the curve showing the variation of the "extinction coefficient" with the wave-length (or frequency) where the extinction coefficient is defined by the Beer-Lambert laws for the absorption of light by an absorbing medium

$$I = I_0 e^{-\epsilon cd}$$

$$\text{or } \log \frac{I_0}{I} = \epsilon cd$$

Here  $I_0$  is the intensity of the beam of monochromatic light.

$I$  is the intensity of the same beam after passing through a length  $d$  of the absorbing medium.

$c$  is the concentration of the absorbing substance

$\epsilon$  is the extinction coefficient.

In practice,  $d$  is expressed in centimetres,  $c$  in mols per litre, and the more convenient  $\log_{10}$  are used instead of  $\log_e$ . This gives a different absolute value of the absorption coefficient, in this case denoted by  $K$  where  $K = 0.4343 \epsilon$ . Hence the determination of an "absorption spectrum" involves the measurement of values of  $K$  corresponding to a number/

number of different wave-lengths spaced at small intervals. It is usual in plotting the absorption curve to employ  $\log K$  instead of  $K$  as, since the values of  $K$  measured may range from  $10^{-1}$  to  $10^5$  in this way features of the curve occurring at widely different values of  $K$  are more equally brought out.

The older methods of measuring absorption spectra were carried out with the help of a Baly tube, and a spectroscope or spectrograph. The light source was one which gave either a continuous and uniform spectrum, or a line spectrum in which the lines are very close together and evenly distributed throughout every region of the spectrum to be explored. The method of making observations was as follows: the cell was filled with a solution of known concentration and placed in front of the slit of the spectroscope. The wave-length readings were then taken at both edges of the band or bands. The strength of the solution was then diminished to about  $\frac{1}{2}$  or  $\frac{1}{3}$  of the original concentration and the new readings of the edges of the bands were taken. The process was repeated until the dilution was such as to cause but a single very narrow band; <sup>the</sup> <sub>^</sub> mean of the readings for this narrow band gave the position of the "Maximum Absorption" or the "head" of the band. The experimental data were plotted in the form of curves the abscissae being wave-lengths or frequencies/



frequencies and the ordinates logarithms of equivalent thicknesses. This method was very laborious and did not give very dependable results.

Actually to measure  $K$ , the extinction coefficient, it would theoretically be simplest to pass a beam of light of known wave-length and known intensity through a known thickness of a known concentration of the absorbing substance and to measure the intensity of the emergent beam. This is the method used by V. Halban and Siedentopf (13) (14) who used a monochromator and a photo-electric cell and reached an accuracy of 0.1%. But if  $K$  is to be measured for a large number of different wave-lengths the method becomes exceedingly slow and laborious and a photographic method is enormously more convenient.

The later and more systematic methods of absorption spectroscopy depend on spectrophotometry i.e. on the science of measuring the intensity of the light of each wave-length in the various regions of the spectrum. Just as with ordinary photometry, the method is one of comparison by the method of rendering the known and unknown equal. In spectroscopy the standard rays are derived from the same source as are the unknown because one cannot rely on getting from two similar sources two identically equal spectra or on the one source remaining constant from second to second.

Hence/

Hence in most spectrophotometers a pencil of light from any suitable source is divided into two equal parts, one to pass directly into the spectroscope to function as a standard, and the other to suffer change by its passage through the given medium; the amount of change at any point can then be determined by reference to the respective standard.

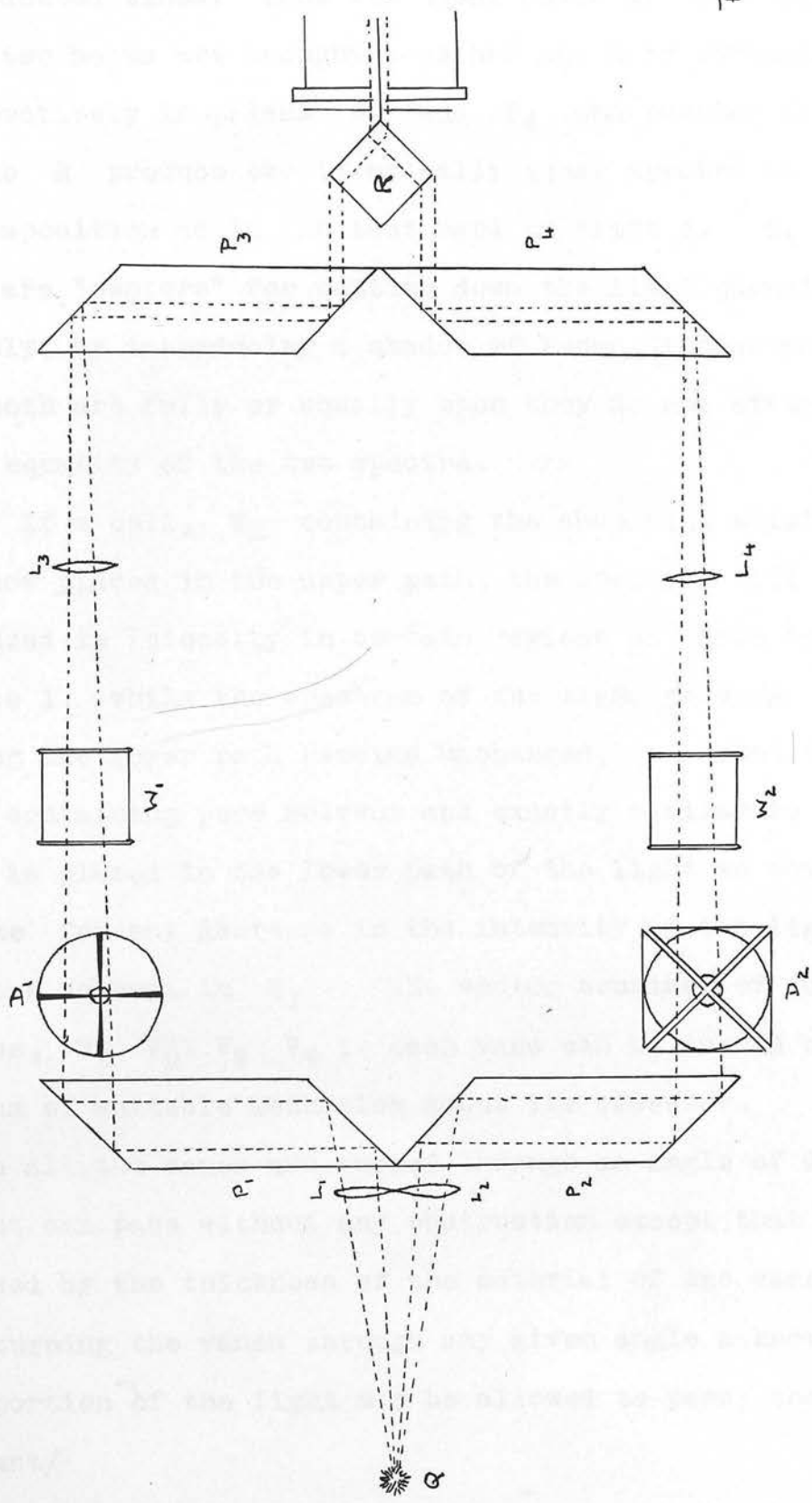
The first accurate method of spectrophotometry was devised by Vierordt in 1873 who may be regarded as a pioneer in applying measurements of absorption spectra to chemical problems.

Different methods have been used by various investigators for cutting down the intensity of the comparison beam. Houston (15) and Schaeffer (16) used two similar light sources and varied the intensity of the illumination from one by altering its position; Winther (17) used a calibrated blackened gauze for the reduction of intensity; V. Henri (18) used two equal sources and varied the time of exposure to one. The objection to these methods are all fairly obvious.

But the three most important methods for rapid measuring of the extinction coefficient over a series of wavelengths making use of the comparison of an unknown with a known intensity are: (1) the Judd Lewis sector photometer method (2) the Hilger rotating sector photometer (3) the Echelon cell.

In the Judd Lewis sector photometer method beams of light/

Fig. 2.



light from Q (Fig. 2) are reflected at the faces in prism  $P_1$  and at the faces in prism  $P_2$  as shown by the dotted lines. Thus two equal beams are produced. The two beams are brought together again by reflection respectively in prisms  $P_3$  and  $P_4$  and passing through rhomb R produce two identically equal spectra in juxtaposition as in the test band in Plate 1.  $D_1$  and  $D_2$  are "sectors" for cutting down the light quantitatively, by introducing a shadow of known dimensions, if both are fully or equally open they do not affect the equality of the two spectra.

If a cell,  $W_1$  containing the absorbing solution be now placed in the upper path, the spectrum will be reduced in intensity in certain regions as shown in plate 1. while the spectrum of the light passing along the lower path remains unchanged; a second cell  $W_2$  containing pure solvent and exactly similar to  $W_1$  is placed in the lower path of the light to compensate for any decrease in the intensity of the light by the solvent in  $W_1$ . The sector consists of four vanes,  $V_1 V_2 V_3 V_4$ ; each vane can be turned by means of suitable mechanism about its bisector. When all the vanes are turned through an angle of  $90^\circ$  light can pass without any obstruction except that caused by the thickness of the material of the vanes. By turning the vanes through any given angle a known proportion of the light may be allowed to pass; the amount/

amount of rotation of the vases is measured  
pointer.

In working the instrument the light is focused  
in the upper part of the vases so that the light is  
so that the light is focused in the upper part of the  
vases.

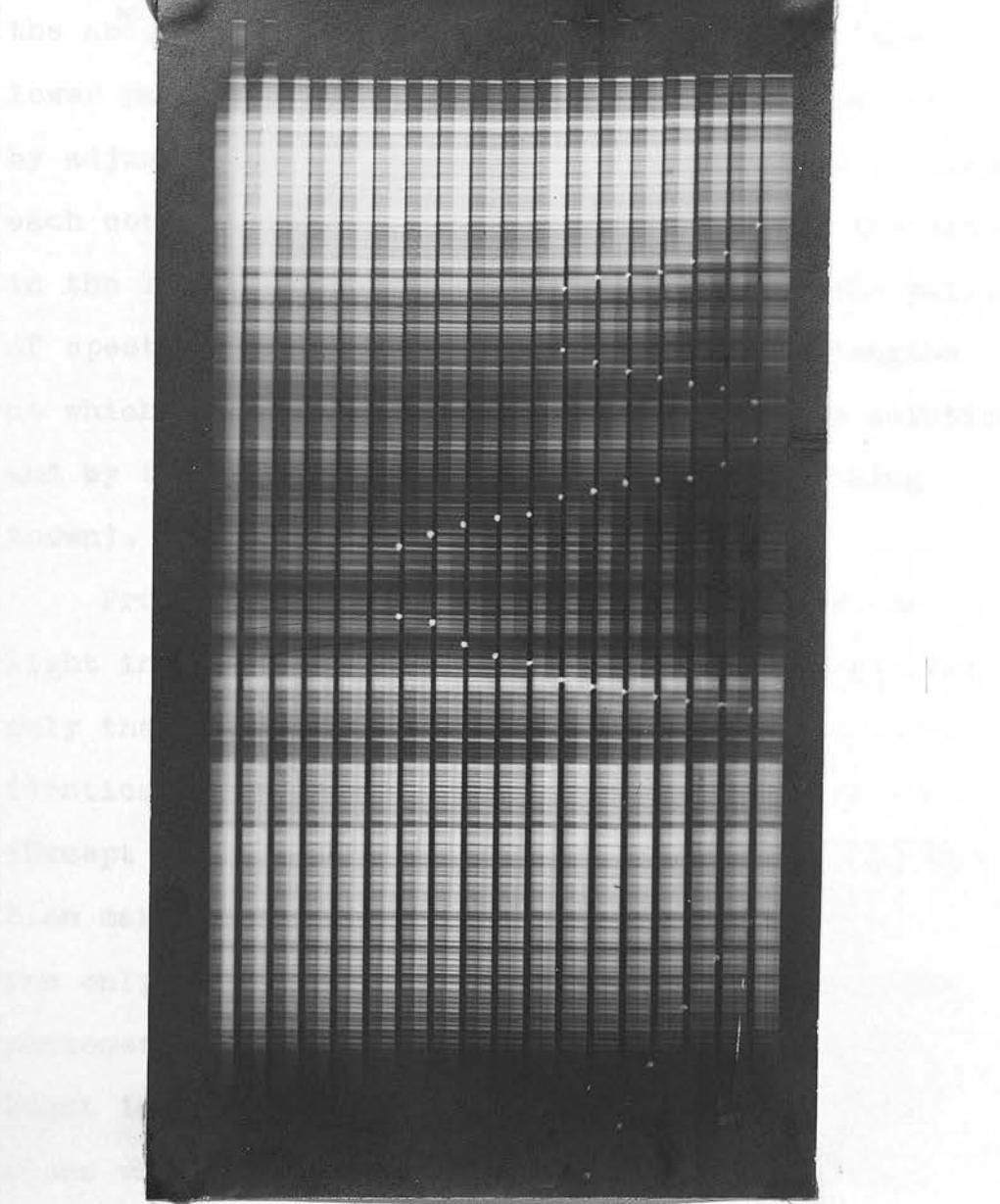


PLATE I.

amount of rotation of the vanes is measured by a pointer.

In working the photometer the vanes are arranged in the upper path so that the maximum light can pass, so that in that path the light is only cut down by the absorbing substance in the cell  $W_1$ . In the lower path the light is cut down by a known amount by adjusting the vanes; a series of exposures is taken, each corresponding to a different setting of the vanes in the lower path, and by an examination of the pairs of spectra obtained we can pick out the wavelengths at which the light is cut down equally by the solution and by the vanes (the latter "cutting down" being known).

Provided that the mechanism for reducing the light intensity is accurate, a method which involves only the comparison of equal intensities obtained by identical exposures cannot have theoretical errors. (Except when the intermittent nature of the illumination makes a difference, which seldom happens.) The only difficulty encountered with the Judd Lewis photometer is that the mechanism for reducing the light intensity involves somewhat elaborate corrections which are only valid if the mounting in the workshop has been accurately carried out.

In the Hilger rotating sector spectrograph the principle is similar to that in the Judd Lewis spectrograph/

spectrograph in that light from a single source is divided into two beams, one of which passes through the absorbing solution, and the other passes through the pure solvent and is then cut down by a known amount by mechanical means. The two beams eventually produce a pair of spectrum photographs in close juxtaposition, one of which is of reduced density throughout its whole length, the other - that which has passed through the material under test - being more dense than the first in certain parts and less so in others, there being certain wave-lengths where the two densities are equal. In this case the light is cut down in the comparison beam by means of a rotating sector the aperture of which can be varied. The beam passing through the absorbing solution also passes through a rotating sector of fixed opening. This is necessary so that both beams falling on the plate may be of an intermittent nature.

In both the Judd Lewis and Hilger Spectrographs the length of absorbing liquid is varied by using cells of different lengths.

Schwarzschild (19) found that when the product of intensity of illumination into time of exposure ( $I t$ ) is constant the density or blackening of a photographic plate was not necessarily constant. His experience is summarised in the equation

$$\frac{I_1}{I_2} = \left( \frac{t_1}{t_2} \right)^n \quad (\text{where } n \text{ is a constant.})$$

for constant density of the finished photographic plate.

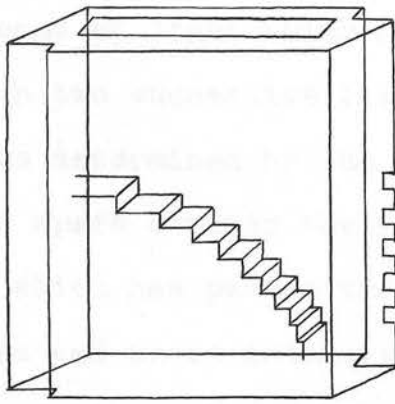
The significance of this Schwarzschild constant is obvious in all measurements of extinction coefficients by photographic methods, especially when a rotating sector is employed.

Howe and Gibson (20) stated that "within the limits of experimental error, and under the conditions obtaining in the use of the Hilger sector photometer, the index  $n$  may be taken as unity". This statement was later confirmed by Gibson, McNicholas, Tyndall, Frehafer and Matthewson (21). The same conclusion was reached by Scheibe, May and Fischer (22).

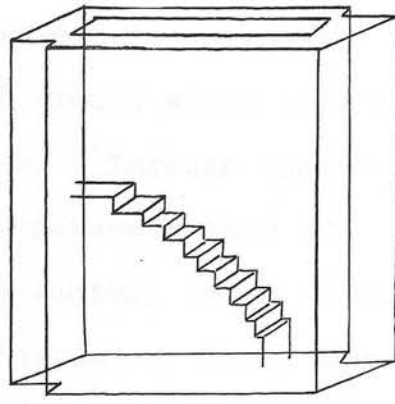
Baly, Morton, and Riding, (23) in a survey of the various methods of measuring the extinction coefficient of a substance come to the conclusion that "the sector photometer and the Judd Lewis sector spectrophotometer are both reliable instruments, and they give concordant results within a range of experimental error rarely if ever exceeding  $\pm 2\%$ ".

The latest device, an echelon cell, to study the absorption spectrum of a liquid was described by F. Twyman in a paper read before the Physical Society on 21st October 1932. Fig. 3. The ten steps in the echelon determine the thickness of the layer of the liquid to be examined. There are five notches along the/





NOTCHED CELL



CELL WITHOUT NOTCHES

Fig. 3.

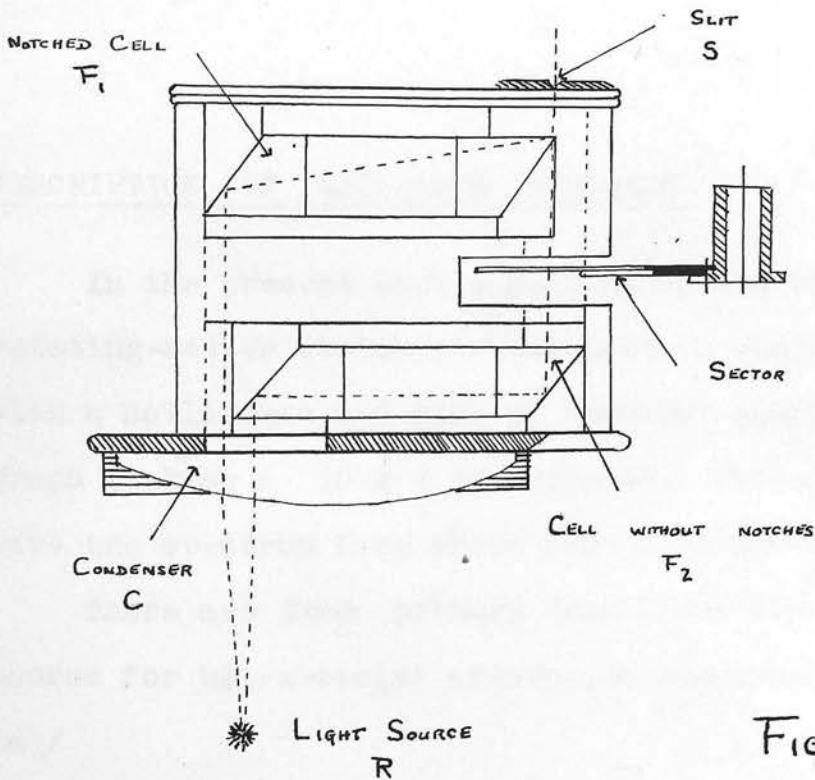


Fig. 4.

the edge of the cell against the slit, from each of which is reflected onto the slit of the spectrograph two beams of light which have passed respectively through two successive layers of liquid whose thicknesses are determined by the echelon. Through the adjacent space between the notches passes a beam of light which has passed through a control cell without notches and whose intensity is reduced by known amount by a rotating sector. The whole arrangement is shown in Fig. 4. The echelon being only 2 cm. in total height the whole cell stands in front of the slit, and the whole twenty bands (i.e. ten through the echelon and ten through the control cell.) are photographed in one exposure. Its special value is twofold: (1) in studying liquids subject to rapid change by ultra-violet radiation, and (2) as a speedy and convenient means of experiment.

## 2. DESCRIPTION OF APPARATUS EMPLOYED.

In the present work a Bellingham and Stanley rotating-sector photometer was used in conjunction with a Bellingham and Stanley "medium" quartz spectrograph (taking a 10 x 4 inch plate). This combination gave the spectrum from about 180  $\mu$  to the red.

There are four primary conditions for a light source for ultra-violet absorption spectrophotography.

(a)/

(a) The source should be as free from fluctuations as possible, (b) it should have a high intrinsic brightness, (c) it should yield a spectrum as nearly continuous as possible, (d) it should be a point source. In the present case an iron spark 2-3 mm. in length was used in series with a second spark gap and in parallel with a condenser of about 0.025 microfarad. capacity. The spark was worked by a transformer on the 230 volt A.C. supply, giving about 6000 volts across the secondary. As the cells of the photometer were mounted on convergent axes, the spark had to be at the intersection of these; a latitude in its position of a few mm. each way was allowable. When the spark was in the correct position, the fact that the two cells had identical "entrance pupils" determined that two exactly equal beams should be produced. The electrodes had one end sharpened to a wedge and this edge was placed so as to be colinear with the optical axis of the spectrophotometer. Then any wandering of the discharge will only take place along this axis and will have no ill effect upon the results.

The photometer had two sectors, one with fixed and one with variable opening; the latter read directly in terms of  $\log \frac{I_0}{I}$ , giving values of this quantity from 0 to 2. A range of cell sizes of 0.1 to 4.0 cm. was available - cells with screw-on quartz ends for 1.0, 2.0 and 4.0 cm., and glass distance-pieces with/

with quartz end-plates for 0.1, 0.2 and 0.4 cm. These thicknesses in the latter case were not quite accurate, but measurement with a micrometer enabled a corrected value of  $\log d$  to be fixed. It was found not advisable to use values of  $\log \frac{I_0}{I}$  less than about 0.30 as then the two spectra of the pair were of such nearly equal intensity in the region of no absorption, that it was very difficult to fix the point of equal blackening. The exposures did not, of course, need to be timed accurately. Before each set of exposures, an approximate calculation showed the correct value of  $d$  to use with the concentration in question. Each sector-opening corresponds to a certain value of  $\log Kcd$  according to the relation  $\log \frac{I_0}{I} = Kcd \therefore \log K = \log(\log \frac{I_0}{I}) - \log c - \log d$ . Hence a table was drawn up, giving the values of  $\log Kcd$  for each sector-opening and also  $\log Kc$  for each combination of opening and size of cell. Hence  $\log K$  by the above equation. The sector-openings used were chosen so as to give values of  $\log K$  differing by about 0.1.

Ilford "Thin-film Half-tone" plates were used for wave-lengths less than  $400 \text{ m}\mu$  and were very satisfactory on account of their high contrast; but as they were most insensitive to longer wave-lengths, "Zenith Super-Sensitive" (650 H & D) plates were used for this region. The finished plates were examined/

examined with a magnifying glass in order to determine the points of equal blackening for each pair of spectra. The wave-lengths were fixed by comparison with an enlarged photograph of the iron spark spectrum with the wave-lengths of the lines marked. Each plate was examined twice at least, with an interval of time in between, and it was also found advantageous to reverse the plate between examinations; when this led to two different points being chosen, a mean position was taken.

### 3. PREPARATION OF MATERIALS.

Sn I<sub>4</sub> . This was prepared by the method of McDermott. 12 grammes tin ("British Chemical Standard" 99.98% pure granulated) and 50 grammes A.R. iodine were refluxed in 75 grammes carbon tetrachloride (free of chlorine, HCl and sulphur) until the violet colour of the iodine vapour was gone. The stannic iodide was filtered off when cool and twice recrystallised from carbon tetrachloride. Melting point 142.5<sup>o</sup> C. (uncorrected).

Sn Br<sub>4</sub> . Prepared by the method of Lorenz (25). Tin was placed in a distilling flask, closed at the top with a cork fitted with a separating funnel drawn out/

out to a capillary point; the funnel was filled with bromine, which was added to the tin so slowly that the temperature remained at  $35 - 59^{\circ}$  and the side-tube did not become filled with bromine vapour. When the greater portion of the tin had disappeared the contents of the flask were distilled; the first few drops contained bromine but that which subsequently passed over was pure stannic bromide. Boiling Point =  $200^{\circ}\text{C}$ . (uncorrected).

Sn Br<sub>2</sub> . By the method of Butler and Hamilton for the preparation of zinc chloride. (26)

Pure dry tin (30 gms.) was placed in a flask with sodium-dried redistilled ether (100 cc.).

Pure dry H Br (Erdmann "Anleitung zur Darstellung chemische Präparate" prepared by dropping redistilled bromine onto pure sodium-dried benzene) was passed into the flask until the tin had completely dissolved. The clear liquid separated into two layers. The flask was warmed on a water bath and ether and excess H Br distilled in vacuo, further heating in vacuo yielded the pure product.

Sn Cl<sub>4</sub> . The product supplied by Hopkin and Williams was purified by redistillation over tin in an all glass distilling apparatus.

C<sub>2</sub>I<sub>2</sub> . By the method of Dehn (27) (J.A.C.S. 33. 1598) Acetylene was passed into a solution of KI in water, into which was dropped a solution of Na OCl (prepared/

(prepared by passing  $\text{Cl}_2$  into an aqueous solution of  $\text{NaOH}$ ) until all the  $\text{KI}$  was used up, i.e. until the  $\text{KI}$  lost its power of giving a brown colouration with  $\text{NaOCl}$ . The  $\text{C}_2\text{I}_2$  formed as a white flocculent precipitate, this was filtered off, washed thoroughly with water, dissolved in low boiling ligroin, dried over  $\text{CaCl}_2$  and decanted off and allowed to crystallise. Very fine white crystals with a melting point of  $80^\circ\text{C}$ . were obtained.

$\text{C}_2\text{I}_4$ . By the method of Datta and Prosad (28) Nitrogen iodide was first obtained as a black precipitate by adding strong ammonia to a  $\text{KI}$  solution of iodine. It was washed thoroughly with water by decantation, placed under water, and acetylene slowly passed in with shaking from time to time. At first the greenish yellow supernatant liquid began to ~~dis-~~colourise, on shaking the colour reappeared due to the dissolution of some nitrogen iodide which was again decolourised in a similar way. The reaction became more rapid as the reaction proceeded. Crystals began to appear which gradually increased until the whole of the nitrogen iodide was converted into a mass of dirty white fine crystals which on the further action of the gas turned orange-yellow. The substance was recrystallised from ether, when small light yellow needle-shaped crystals with a melting point of  $185^\circ\text{C}$  were obtained.

CHI/

CHI = CHI. By the method of Biltz (29)

Acetylene was passed slowly into a flask containing 100 gms. finely ground iodine, covered with 200 gms. EtOH for five days. The acetylene was slowly absorbed, and white needle-like crystals formed. The contents of the flask were poured into water, the excess iodine removed with NaOH, the crystals filtered off, washed and recrystallised twice from EtOH. White needles with a melting point of  $73^{\circ}\text{C}$  were obtained.

CH<sub>2</sub>I - CH<sub>2</sub>I. By the method of Semenow (30)

A paste of iodine in EtOH was saturated with ethylene with constant shaking and with repeated addition of iodine. The iodine solution was decanted from the crystalline mass formed; the latter was washed with cold EtOH and purified (31) by two recrystallisations from ether, free iodine being removed by shaking up the ethereal solution with a solution of NaOH.

Melting point  $81 - 82^{\circ}\text{C}$ . It was found that if left exposed to the atmosphere the crystals slowly gave off iodine, but this decomposition was prevented if the substance was kept in an atmosphere of ethylene.

CH<sub>2</sub>I<sub>2</sub>. By the method of Auger (32) 1 gm.-mol.

weight of CHI<sub>3</sub> was mixed with an excess of trisodium-arsenite (obtained by dissolving 1 gm. mol. weight of As<sub>2</sub>O<sub>3</sub> in 6 gm.-mol. weights of a 33% caustic soda solution/



solution) and the mixture steam-distilled. The  $\text{CH}_2\text{I}_2$  formed as a straw-coloured oily liquid and was purified by a vacuum distillation. The  $\text{CH}_2\text{I}_2$  decomposed readily with the formation of free iodine. Such decompositions are known to be catalysed by free iodine and it was found that in the presence of a little mercury the  $\text{CH}_2\text{I}_2$  was quite stable.

Hexane. The "specially purified for spectroscopy" product of the B.D.H. was used.

Carbon Tetrachloride. The ordinary sulphur-free A.R. tetrachloride was found to be sufficiently transparent not to require further purification.

Ethyl Acetate. The A.R. product in this case also needed no further purification.

Ethyl Alcohol. A very large number of different methods for purifying this substance have been given. As it was desired here to obtain a product with the greatest transparency in the ultra-violet, the method of Castille and Henri (33) which was specially designed to give a product suitable in this respect, seemed to be the most promising. After some experiment the following modification of the above method was adopted. Ordinary absolute alcohol was refluxed for 3 - 4 hours with 2 gms. iodine per litre. 30 - 35 gms. calcium (to remove the iodine and water simultaneously) were then added, and the refluxing continued for a day. The alcohol was then distilled off/

off, while a current of air dried by passage through a 5 foot tube filled with  $\text{CaCl}_2$  was passed through it in order to remove any traces of ammonia. The first product was rejected, in case it contained ammonia, and the distillation was stopped when about two-thirds had come over. The alcohol so obtained showed the following transmission:-

In a 4 cm. layer to 230  $\text{m}\mu$

In a 1 mm. " " 215  $\text{m}\mu$

(Exposure: just sufficient to record to about 208  $\text{m}\mu$  in air). The spectrum down to the limit of transmission of the apparatus could easily be photographed through the alcohol using increased exposures

#### Methyl Alcohol.

"Absolute"  $\text{MeOH}$  ( $\frac{1}{2}$  - 1%  $\text{H}_2\text{O}$ ) reacts actively with magnesium (10 gms. / litre) and is quite dehydrated. The reaction proceeds of itself and the  $\text{MeOH}$  boils. When the  $\text{Mg}$  has disappeared, the  $\text{MeOH}$  is refluxed gently for several hours and then distilled in a stream of dry air, the first portion of the distillate being rejected. (34)

Acetonitrile. A.R. acetonitrile was dehydrated over phosphorus pentoxide for several days and redistilled (35).

The solutions were prepared by dissolving a weighed amount of the substance in the solvent, with subsequent/

subsequent dilution of necessary. The concentra-  
tions used in calculating log K were simply based on  
the weight of substance dissolved.

Previous workers (1) have found that a plot  
of  $\log K$  versus  $\log C$  for a series of solutions  
of a given substance in a given solvent at a  
constant temperature is a straight line. This  
is in agreement with the view of Scherer  
that the activity of a solid is determined  
only by the surface area. It was the original  
aim of this research to see if this similarity  
held for other solids.

For solids such as sugar, which are  
small particles, the determination of the  
weight of the dissolved substance is not  
difficult. For larger particles, the  
determination of the weight of the dissolved  
substance is more difficult. In this  
research, the weight of the dissolved  
substance was determined by the weight of  
the residue left after the solvent had  
evaporated.

A series of solutions of sugar in  
water were prepared in which the  
concentration of the dissolved substance  
was varied. The weight of the dissolved  
substance was determined by the weight of  
the residue left after the solvent had  
evaporated. The results are shown in  
Figure 1. The results show that the  
plot of  $\log K$  versus  $\log C$  is a  
straight line. This is in agreement  
with the view of Scherer that the  
activity of a solid is determined  
only by the surface area.

### III. EXPERIMENTAL.

#### 1. STANNIC HALIDES.

Previous workers (60) had found that a solution of stannic iodide in ethyl alcohol gave an absorption curve very similar to that of a solution of iodine in ethyl alcohol. This seemed to support Scheibe's view that absorptions of this type are determined only by the halogen atom present (36). It was the original aim of this research to find out if this similarity held for other solvents.

For solvents such as hexane, which have very small dipoles, the deformation of the electron system of the dissolved molecule is very slight, so that the absorption of a substance dissolved in hexane may be taken as closely approximating to that of the non-deformed unsolvated molecule. For this reason hexane was the first solvent used in this research.

A solution of iodine in hexane is purple and has an absorption curve similar to that of iodine vapour (Graph 1, curve <sup>1</sup>~~2~~). Stannic iodide in hexane gives a yellow solution with an absorption curve similar to that of iodine in ethyl alcohol (Graph 1, curve <sup>2</sup>~~1~~), but it was found that this solution did not obey Beer's Law, and on standing the colour slowly changed to brown and later to purple, and a white deposit formed. This solid was filtered off, washed thoroughly with/

with hexane, and on examination was found to contain tin but no iodine. This pointed to the stannic iodide having decomposed with the formation of free iodine; to test this the purple solution was shaken up with finely divided metallic silver, the solution became colourless and a yellow deposit (which proved to be silver iodide) formed on the surface of the silver.

Means were then sought by which to prevent the decomposition of the stannic iodide, and it was found that if allowed to stand over finely-divided metallic tin the solution did not darken on standing although a white deposit gradually formed. This solution gave an absorption curve (Graph 2, curve 4) which obeyed Beer's Law within the limits used.

A solution of iodine in hexane was then shaken up with finely-divided tin, and it was observed that the colour gradually changed from violet through brown to pale yellow; the final solution gave the same curve as that of stannic iodide in hexane over tin (Graph 2, curve 4) if in calculating the concentration of iodine the molecule is taken as  $I_4$ , and obeyed Beer's Law. Tin and iodine combine directly with the formation of stannic iodide, so this curve may be taken as the true curve for stannic iodide in hexane.

It seemed probable that the observed similarity between the curves for stannic iodide and iodine in ethyl/

ethyl alcohol was due to the slow decomposition of the stannic iodide with the liberation of free iodine, so a solution of stannic iodide in ethyl alcohol was made up and allowed to stand over metallic tin. This solution was pale yellow and gave a curve (Graph 2, curve 3) quite different from that of iodine in ethyl alcohol (Graph 1, curve 1) and obeyed Beer's Law. Stannic iodide in ethyl alcohol with no metallic tin present had an absorption curve which changed on standing (Graph 2, curves 1 and 2). Ethyl alcohol was allowed to stand over tin for a month but no absorption was observed. A solution of iodine in ethyl alcohol was shaken up with metallic tin and the colour gradually changed from dark brown to pale yellow; the resulting solution gave the same curve as that of stannic iodide in ethyl alcohol if in calculating the concentration of iodine the molecule is taken as  $I_4$  (Graph 2, curve 3); this solution also obeyed Beer's Law. This curve may be taken as the true curve for stannic iodide in ethyl alcohol.

Similarly stannic iodide in methyl alcohol in the presence of metallic tin gave a curve (Graph 2, curve 5) quite different from that of iodine in methyl alcohol without tin (Graph 1, curve 3), but the same as that for iodine in methyl alcohol over tin.

A solution of stannic iodide in acetonitrile over tin/

tin decomposed too rapidly under the action of the ultra-violet light for any curve to be obtained.

As the curves for stannic iodide in ethyl alcohol and in hexane are so different a series of solutions of stannic iodide in hexane containing varying amounts of ethyl alcohol over metallic tin was examined in an attempt to trace the transition from one curve to the other. With small amounts of ethyl alcohol (.0318 and .09 molar) the curve differed very little from that of stannic iodide in hexane over tin (Graph 3, curves 1 and 2). But with increasing amount of alcohol (Graph 3, curves 3-6) the short-wave maximum rose and shifted towards the red, while the long-wave maximum fell and shifted towards the ultra-violet, the curve gradually flattening out until at length both maxima practically disappeared. All the curves passed through the same point (within the limits of experimental error) at about  $\lambda = 343 \text{ m}\mu$  and  $\log K = 3.66$ .

A solution of stannic bromide in hexane (which is colourless) was examined and it was found that, even in the presence of metallic tin, the absorption altered with time (Graph 4, curves 1-3). The maximum wavelength did not change but  $\log K_{\text{max}}$  gradually decreased.

The absorption of stannic bromide in methyl alcohol over tin also changed with time (Graph 5, curves 1, 2, 4). Both in the case of hexane and methyl alcohol/

alcohol a white deposit formed when the solution was allowed to stand, but the solution itself remained colourless. Stannous bromide in methyl alcohol has a curve (~~Fig.~~ <sup>GRAPH</sup> 5, curve 3.) similar to that of stannic bromide in methyl alcohol but further towards the ultra-violet. Stannous bromide is insoluble in hexane.

In order to complete the series a solution of stannic chloride in hexane was examined. As with even a trace of moisture stannic chloride reacts rapidly to form stannic hydroxide, the following method was used:- Some stannic chloride was distilled into some hexane in an all glass distilling apparatus, the cloudy solution so obtained was filtered and the filtrate analysed for stannic tin (37); in the analysis a known volume of the hexane solution was added to 250 cc. of water containing the equivalent of 50 cc. of concentrated hydrochloric acid, and the solution boiled for some time to drive off the hexane. The solution was then reduced by heating with lead shot (38) in an atmosphere of carbon dioxide, when the solution has cooled it was titrated with standard iodine using starch as indicator. Two concordant results were obtained. This standard solution of stannic chloride was then diluted to the correct concentrations as required. As the absorption was well over into the ultra-violet a copper spark was used as the source of ultra-violet light, and the plates used were/



were "Zenith Super-Sensitive" (650 H & D ) sensitised with oil, and long exposures were given. Hexane was the only solvent possible because stannic chloride was found to react rapidly with the others (Graph 4. curve 4.).

It was intended to have measured the absorption spectrum of stannic iodide in acetonitrile (this was later found to be impossible owing to the instability of the solution) and as a preliminary, a solution of iodine in acetonitrile was examined. The absorption of this solution changed with time, and a series of curves were obtained (Graph 6). Similar results were obtained by Ley and Engelhardt (39) for iodine in ethyl alcohol.

All the tin solutions were kept under a blackened bell-jar as they tended to decompose if exposed to bright sunlight; for the same reason, as short exposures as possible were used and the solution in the cell was changed before each exposure.

## 2. CARBON - IODINE COMPOUNDS.

### Previous Work.

Crymble, Stewart and Wright (31) measured the absorption spectra of  $C_2H_2I_2$  and  $CH_2I_2$  in ethyl alcohol, and from these results in conjunction with/

with results from similar work with  $\text{CHI}_3$  ,  $\text{KI}_3$  ,  $\text{C}_2\text{H}_5\text{I}$  and  $\text{I}_2$  put forward a theory to account for the absorption (This theory is discussed later in this essay).

Macbeth and Stewart (40) examined the absorption of  $\text{C}_2\text{I}_2$  ,  $\text{C}_2\text{H}_2\text{I}_2$  and  $\text{C}_2\text{H}_4\text{I}_2$  and put forward an explanation of their results on the basis of the change produced in the absorptive power of the iodine atoms by the introduction of double and triple bonds into the  $\text{C}_2\text{H}_4\text{I}_2$  molecule.

Massol and Faucon (41) studied the absorption of  $\text{CH}_2\text{I}_2$  in the course of a work undertaken with a view to correlating the absorption of the iodine derivatives of methane with that of iodine.

None of these authors measured the absorption coefficients of the substances and their work was not of a strictly quantitative nature. In view of the small amount of work which has been done on these compounds it was thought that a systematic study of the absorption curves of the series  $\text{C}_2\text{I}_2$   $\text{C}_2\text{I}_4$   $\text{C}_2\text{H}_2\text{I}_2$   $\text{C}_2\text{H}_4\text{I}_2$  and  $\text{CH}_2\text{I}_2$  in various solvents might yield interesting results.

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PRELIMINARY EXPERIMENTS.

The white crystals of  $C_2I_2$  were found to be easily soluble in hexane, ethyl alcohol, methyl alcohol and acetonitrile, but it was observed that the solution became coloured (hexane solution became violet, and the others brown) if they were left exposed to bright sunlight or ultra-violet light. It was thought that perhaps it was only the far ultra-violet light that was causing the decomposition, so various filters were tried. It was finally found that a thin sheet of vita-glass (transmitting to  $\lambda = 275 \mu$  for a 5 secs. exposure) placed between the source of ultra-violet light and the solution prevented the decomposition of the latter; the same was found to hold good for all the other iodine compounds of carbon studied, so it was decided to use this filter during the measurement of their absorptions; this was found to be quite satisfactory and gave reproducible results.

TABLE./

## SOLUTIONS MEASURED.

No.	Solute	Solvent	Conc. in		Graph.	Curve.
			Mols. I.	per litre II.		
1.	$C_2I_2$	Hexane	.000642	.00642	8	1.
2.		$CCl_4$	.00641	.0320	8	2.
3.		Ethyl acetate	.00642	.0321	8	3.
4.		EtOH	.00642	.0321	8	4.
5.		MeOH	.00642	.0321	8	4.
6.		$CH_3CN$	.00642	.0321	8	4.
7.	$C_2I_4$	Hexane	.000642	.000321	7	1.
8.		$CCl_4$	.001002	.000334	7	2.
9.		Ethyl acetate	.001062	.000334	7	3.
10.		EtOH	.000668	.000334	7	4.
11.		MeOH	.001068	.000356	7	4.
12.		$CH_3CN$	.000975	.000670	7	5.
13.	$C_2H_2I_2$	Hexane	.001670	.00835	9	1.
14.		$CCl_4$	.001674	.00837	9	2.
15.		Ethyl acetate	.001697	.008486	9	3.
16.		EtOH	.001670	.00835	9	3.
17.		MeOH	.001677	.008387	9	3.
18.		CH CN	.001683	.008414	9	4.

SOLUTIONS MEASURED. (Contd.)

No.	Solute	Solvent	Conc. in		Graph.	Curve.
			Mols. per I.	Litre II.		
19.	$C_2H_4I_2$	Hexane	.008301	.001660	10	1.
20.		$CCl_4$	.008171	.001634	10	2.
21.		Ethyl acetate	.008171	.001634	10	1.
22.		EtOH	.008130	.001626	10	1.
23.		MeOH	.008215	.001643	10	3.
24.		$CH_3CN$	.008228	.001645	10	4.
25.	$CH_2I_2$	Hexane	.005546	.001109	11	1.
26.		$CCl_4$	.005224	.001045	11	1.
27.		Ethyl acetate	.004286	.001200	11	2.
28.		EtOH	.005952	.001190	11	2.
29.		MeOH	.005662	.001132	11	2.
30.		$CH_3CN$	.004349	.001230	11	3.

In the case of  $CH_2I_2$  a small amount of mercury was placed in the cells during the exposures as this was found to act as a negative catalyst in the decomposition of  $CH_2I_2$ .

DESCRIPTION OF CURVES.C<sub>2</sub>I<sub>4</sub> .

C<sub>2</sub>I<sub>4</sub> in hexane has an absorption curve with two maxima (1) at  $\lambda = 332 \text{ m}\mu$  and (2) at  $\lambda =$  approximately  $290 \text{ m}\mu$ . There is a shallow minimum between these two maxima. The curve then rises steeply into the ultra-violet.

The values for  $\log K$  at the two maxima are (1) 2.94 and (2) 3.08.

The curve changes with change in solvent, and with increasing polarity of solvent the two maxima appear to move together, until with ethyl alcohol only one appears, and with acetonitrile there is no definite maximum at all.

C<sub>2</sub>I<sub>2</sub> .

The solution in hexane has a curve rising fairly steeply to about  $\lambda = 334 \text{ m}\mu$  ( $\log K = 1.75$ ) where the curve flattens out until about  $\lambda = 304 \text{ m}\mu$  ( $\log K = 2.0$ ) and then rises steeply to  $\log K = 2.64$  where the first part of a maximum is obtained at

$\lambda = 273 \text{ m}\mu$ . The whole of the maximum could not be measured because it was beyond the region of transmission of the vita-glass.

The solution in CCl<sub>4</sub> gave the same curve as that in hexane except at the flattened-out part of the curve where a shift towards the red is observed.

The/



The curve for ethyl acetate is further to the ultra-violet than that for hexane and the region of inflexion is less apparent; Ethyl alcohol, methyl alcohol and acetonitrile all give the same curve which is further to the ultra-violet than the ethyl acetate curve and the region of inflexion is still less pronounced.

C<sub>2</sub>H<sub>2</sub>I<sub>2</sub> .

All the curves (except that for CH<sub>3</sub>CN) just consist of a branch rising fairly steeply towards the ultra-violet. In the case of hexane there is a change in direction at  $\lambda = 307 \text{ m}\mu$  ( $\log K = 2.1$ ) this is not apparent in the polar solvents, but appears in CCl<sub>4</sub> at  $\lambda = 305 \text{ m}\mu$  ( $\log K = 2.15$ ).

The curve for CH<sub>3</sub>CN flattens out from  $\lambda = 330 \text{ m}\mu$  ( $\log K = 1.27$ ) to  $\lambda = 370 \text{ m}\mu$  ( $\log K = 1.1$ ) and then falls off into the visible.

The curves for EtOH MeOH and ethyl acetate are coincident and further to the ultra-violet than the curve for hexane.

The curve for CCl<sub>4</sub> shows a shift towards the red compared with hexane.

C<sub>2</sub>H<sub>4</sub>I<sub>2</sub> .

There are only four distinct curves for this the substance, the curves for ethyl acetate, EtOH and hexane being coincident. They each consist of a/

a single branch rising steeply towards the ultra-violet.

$\text{CCl}_4$  shows a shift towards the red compared with the position for hexane, while the curves for MeOH and  $\text{CH}_3\text{CN}$  are further towards the ultra-violet (again compared with hexane) the shift for  $\text{CH}_3\text{CN}$  being greater than that for MeOH.

### $\text{CH}_2\text{I}_2$ .

The curve in hexane rises steeply to a maximum at  $\lambda = 297 \text{ m}\mu$  ( $\log K = 3.14$ ) and then falls away again towards the ultra-violet. The curve for  $\text{CCl}_4$  is coincident with that for hexane.

The curves for ethyl acetate, EtOH and MeOH are coincident (MeOH diverges on the lower part of the long-wave branch) and show a maximum at  $\lambda = 293 \text{ m}\mu$  ( $\log K = 3.11$ )

The Acetonitrile curve is still further to the ultra-violet  $\lambda_{\text{max}} = 291 \text{ m}\mu$  ( $\log K = 3.04$ )

Beer's Law was found to hold for all the solutions of all the curves measured.



IV. DISCUSSION.1. INTRODUCTION.

Lowry and Sass (42) hold that the absorption due to the bond C - X where X is a halogen atom might be shifted towards the red as a result of the weakening of the C - X bond under the influence of contiguous negative atoms. In order to obtain experimental proof of this theory they examined the series of compounds  $\text{CH}_3\text{I}$   $\text{CH}_2\text{I}_2$   $\text{CHI}_3$   $\text{CI}_4$ . The last of the series was found to be too unstable to examine.

They showed that the absorption shifts towards the red in passing from  $\text{CH}_3\text{I}$  to  $\text{CH}_2\text{I}_2$  to  $\text{CHI}_3$  at which point two maxima appear. They state from the general form and position of the curves that they are dealing with the same band throughout, and that although the second and third iodine atoms may reinforce and develop the band they are not needed to initiate it.

Massol and Faucon (41) had studied the same series and compared the absorption spectra in ethyl alcohol with that of iodine in the same solvent and in the case of  $\text{CHI}_3$  observed two bands which they said corresponded to two of the iodine bands. They had hoped to find a correspondence between the bands of all the compounds/

compounds and one or other of the iodine bands but this was found to be impossible with  $\text{CH}_2\text{I}_2$  and  $\text{CH}_3\text{I}$ .  $\text{CI}_4$  gave an absorption very similar to that of iodine, but we now see that this was probably due to the decomposition of  $\text{CI}_4$  with the formation of free iodine.

Crymble Stewart and Wright (31) attribute the methylene iodide band to the juxtaposition of two atoms of iodine in the same molecule, i.e. they suggest that  $\text{I} - \overset{\text{I}}{\underset{|}{\text{C}}} - \text{I}$  gives the same band as  $\text{I} - \text{I}$  but that  $\text{I} - \overset{\text{I}}{\underset{|}{\text{C}}} - \overset{\text{I}}{\underset{|}{\text{C}}} - \text{I}$  behaves like  $-\overset{\text{I}}{\underset{|}{\text{C}}} - \text{I}$ .

Lowry and Sass (42) show that there is a great similarity between the curves for  $\text{CHI}_3$  and for  $\text{KI}_3$  (i.e. iodine in a solution of potassium iodide (Fig. 5.)), but find that this is not accounted for by the formulae commonly assigned to these compounds, and state that it appears to be beyond the scope of our present theories of valency. They tentatively put forward the suggestion that it may be due to the presence of three contiguous atoms of iodine in the space models of the two compounds; they picture both molecules as dominated by a cluster or ring of three contiguous iodine atoms, just as (according to W.L. Bragg's determination) the crystal structures of most of the oxysalts and of many of the metallic oxides are dominated by clusters of contiguous oxygen atoms, leaving the co-ordinating atom to occupy a mere cavity between/

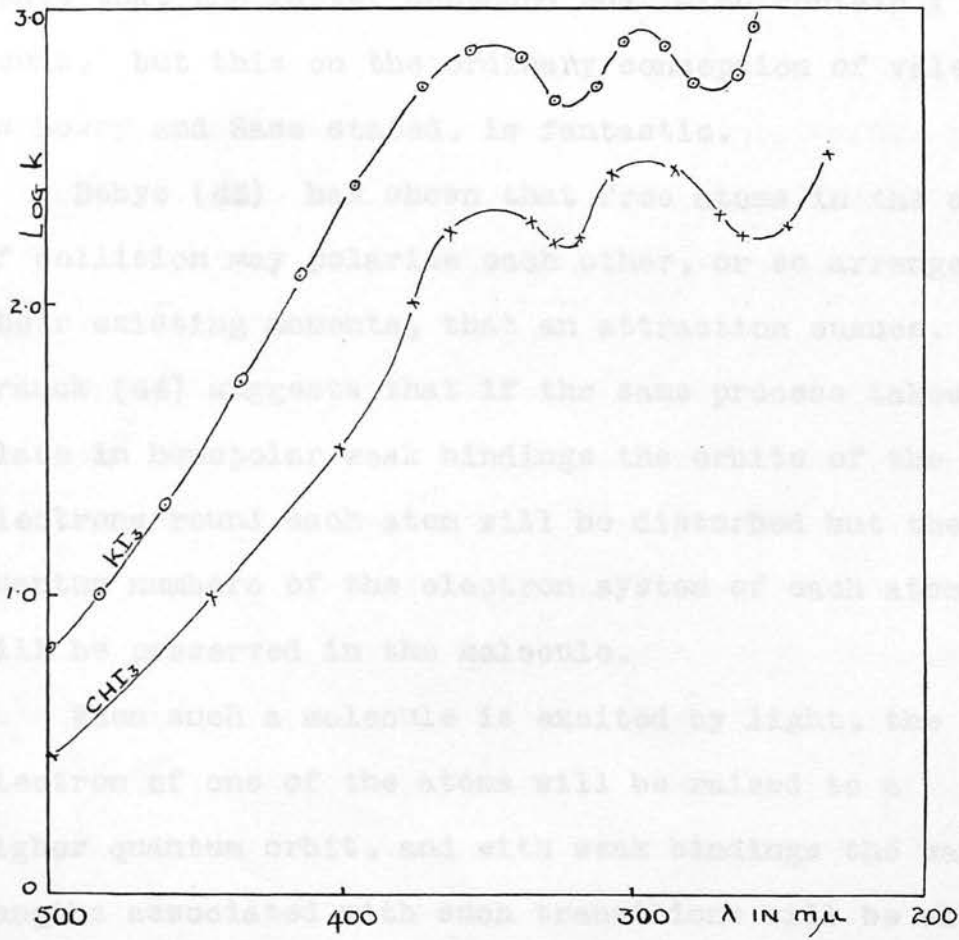


Fig. 5.

between the oxygens. The colour of  $KI_3$  is always attributed to the negative ion  $I_3^-$  which (like iodine itself) can only contain I - I bonds, so that the similarity in the spectra of  $KI_3$  and  $CHI_3$  suggests that the latter compound must also contain I - I bonds, but this on the ordinary conception of valency, as Lowry and Sass stated, is fantastic.

Debye (43) has shown that free atoms in the act of collision may polarise each other, or so arrange their existing moments, that an attraction ensues. Franck (44) suggests that if the same process takes place in homopolar weak bindings the orbits of the electrons round each atom will be disturbed but the quantum numbers of the electron system of each atom will be preserved in the molecule.

When such a molecule is excited by light, the electron of one of the atoms will be raised to a higher quantum orbit, and with weak bindings the wavelengths associated with such transitions will be almost identical with those absorbed by the free atoms.

The work of Mecke (45) on the absorption spectrum of gaseous iodine has shown the I - I binding to be of this type. So that the absorption spectrum of iodine is that of two iodine atoms held together by attractive forces but each retaining their own electronic orbits although in a deformed state.

In/

In the case of solutions of iodine Getman (46) (J.A.C.S. 50 2883) who measured the absorption coefficients of iodine through the visible spectrum in a series of carefully purified solvents showed that the solutions may be divided into two distinct groups as follows:-

(a) solns. in n-hexane, benzene, toluene, carbon disulphide, carbon tetrachloride, and chloroform which are violet with absorption maxima ranging from 520  $m\mu$  to 540  $m\mu$  .

(b) Solutions in ethyl alcohol, isopropyl alcohol, acetone, acetophenone, and pyridine which are brown with absorption maxima ranging from 460  $m\mu$  to 480  $m\mu$  .

The first group of solvents he calls "saturated compounds" while the second group "unsaturated compounds".

He finds that the maximum intensity of absorption is found with the "saturated" solvents and is practically constant throughout the group, while in the second group of "unsaturated" solvents the intensity of

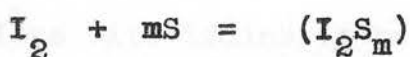
absorption varies appreciably from solvent to solvent and is in every case less than with "saturated" solvents. Also the solutions tend to undergo a change, with corresponding change in absorption, on standing.

Getman concludes that "the colour and stability of iodine solutions on exposure to ultra-violet radiation appear to be more or less dependent upon the degree

of/

of "saturation" of the solvent." It is not very clear what Getman implies by an "unsaturated" solvent.

Ley and Engelhardt (47) after studying the change in the absorption of solutions of iodine with time, adopt the view of Beckmann that in the brown solution there is an equilibrium between the iodine and a solvation product.



This view is supported by Waentig whose work proves a connection between the shift in the absorption and the heat of solution.

de Boer and Custers (48) measured the absorption of iodine adsorbed on calcium fluoride at different degrees of saturation of the surface. They found that at small degrees of saturation the absorption curve shows two decided maxima at  $283 \text{ m}\mu$  and  $342 \text{ m}\mu$ ; if more molecules are adsorbed the curve becomes less definite and the maxima (especially the long-wave one) are shifted towards the visible. They observe that there is a marked similarity between the spectrum for  $KI_3$  and the spectrum of adsorbed iodine on calcium fluoride.

In discussing their results de Boer and Custers employ the formula

$$h\nu = \Delta E + Q_A - Q'_A$$

where  $\nu$  is the frequency of the light absorbed,  $\Delta E$  the energy difference of the excited and unexcited molecule/

molecule  $A'$  in the free gaseous state,  $Q_A$  is the energy of binding in the unexcited state, and  $Q'A'$  is the energy of binding in the excited state.

So they state that the position of the absorption will depend on  $(Q_A - Q'A')$  and this quantity will vary with the state of the iodine.

But in the case of the solvents which give brown solutions with iodine it has been shown that compound formation occurs between the solute and solvent, (the compound with pyridine has been actually isolated) so that the energy of binding of the solvate should also be taken into account i.e.

$$h\nu = \Delta E + Q_A - Q'A' + S_A - S'A'$$

where  $S_A$  and  $S'A'$  are the energies of solvation of the excited and unexcited molecules. The fact that the pyridine compound has been isolated shows that, in some cases at any rate the  $S$  terms are of the same order as the  $Q$  terms and therefore cannot be neglected. The equation now contains so many variables that it is impossible to predict from it the shift in absorption with change in solvent.

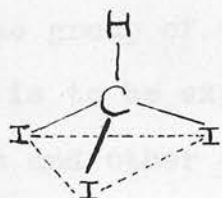
Chatelet (49) states that the absorption curves for brown solutions of iodine only seem compatible with the hypothesis of the combination of iodine with solvent.

Getman's two groups of solvents may therefore be defined/

defined as follows:-(1) "saturated" solvents are those which have little effect on the strength of binding of the dissolved iodine molecule, i.e. in which  $Q_A$   $Q'_A$  in the above equation and  $S_A$  and  $S'_A$  are negligible so that the absorption of the solution approximates to that of free gaseous iodine; and (2) "unsaturated" solvents which enter into a type of combination with the iodine molecule and thus affect the strength of the binding of the latter.

The iodine derivatives of methane are all fairly unstable and decompose under the action of light with the formation of free iodine. This indicates that the homopolar  $\overset{\cdot}{\text{C}} - \text{I}$  binding is a weak one. Thus it may be anticipated that it comes into the same category as the  $\text{I} - \text{I}$  linkage in the iodine molecule, i.e. consists of a carbon atom and an iodine atom held together by attractive forces, each retaining its own electronic orbits though in a deformed state.

If we consider the space formula of  $\text{CHI}_3$  we see that it consists of three iodine atoms held by attractive forces to a  $-\text{CH}$  radical. The attraction between



the carbon and iodine atoms is represented by the ordinary valency bonds directed out from the carbon at the centre of the tetrahedron to

the three iodine atoms at the vertices. But as these iodine atoms are only held to the carbon by weak forces and/



and each retains its own electronic orbits although in a deformed state, it is probable that the iodine atoms will exert similar attractive forces on each other, i.e. the molecule will approximate to a ring of three I - I bonds held by attractive forces to a CH group.  $KI_3$  is always considered to be made up of a  $K^+$  ion and an  $I_3^-$  ion, so this compound also contains a group of three <sup>iodine</sup> atoms held by attractive forces, so that this may help to explain the similarity in the absorption spectra of the two compounds.

When we consider the space formula of carbon tetraiodide we see that the molecule consists of a carbon atom at the centre of a regular tetrahedron with an iodine atom at each of the vertices, so that, by reasoning similar to that employed in the case of  $CHI_3$ , we can imagine the C atom surrounded by a shell of deformed I - I linkages; thus the spectrum of carbon tetraiodide might approximate to the spectrum of iodine.

Carbon tetraiodide is found to be too unstable to be convenient to work with. However, tin is in the same group of the Periodic System as carbon, so that it is to be expected that the type of linkage between tin and other elements is similar to the type of linkages between carbon and the same elements. On this basis the Sn - I linkage may be considered as a weak homopolar/

homopolar binding, this is supported by the fact that stannic iodide decomposes fairly easily under the action of light with the formation of free iodine.

## 2. STANNIC HALIDES.

### (a) Stannic Iodide in Hexane.

Hexane is a practically non-polar solvent, so that its molecules have little or no modifying effect on the absorption of the molecules of substances dissolved in it. Hence the molecules of stannic iodide are merely dissolved in the hexane and not solvated (i.e. not associated with solvent molecules) to any appreciable extent, and the absorption curve in hexane is that for stannic iodide in the normal unsolvated state. This curve shows two maxima (1) at  $\lambda = 365 \text{ m}\mu$  and (2) at  $\lambda = 285 \text{ m}\mu$ . Thus the spectrum of stannic iodide very closely resembles the spectrum of iodine in solvents with which it gives a brown solution, i.e. with solvents in which the iodine molecule is combined with solvent molecules. This is an interesting analogy when we consider that on the above theory the stannic iodide molecule is pictured as a shell of deformed I - I linkages held together by a tetravalent tin atom; while in iodine in brown solutions, the I - I linkages will be deformed by the solute - solvent forces.

(b)/

(b) Stannic Iodide in Ethyl Alcohol.

In the case of the solution of stannic iodide in ethyl alcohol, the condition of the stannic iodide molecule is greatly altered by strong association with the solvent, and now the association complex is unable to absorb radiation at  $\lambda = 365 \text{ m}\mu$  as no maximum appears at that wave-length on the curve for stannic iodide in ethyl alcohol. The ability to absorb radiation is a question of resonance, and the natural frequencies of the association complex do not respond to  $\lambda = 365 \text{ m}\mu$ .

Stannic Iodide in Methyl Alcohol.

This curve shows neither of the two maxima of the curve of stannic iodide in hexane. Methyl alcohol is a more highly polar solvent than ethyl alcohol, so that the degree of solvation of the dissolved stannic iodide molecules is even greater and the distribution of the radiation absorbed by them differs from that when ethyl alcohol is the solvent so that a different absorption curve would be expected, and the results obtained fulfil this expectation.

(c) Stannic iodide in hexane and ethyl alcohol.

When increasing amounts of ethyl alcohol are added to the hexane used as solvent, the dissolved stannic iodide molecules gradually become solvated. The curves all pass through one point so that an equilibrium/

equilibrium exists between the two states of stannic iodide (solvated and unsolvated) in the solution (51). Only the unsolvated molecule absorbs to any great extent at  $\lambda = 365 \text{ m}\mu$ ; on the addition of ethyl alcohol some of the stannic iodide molecules become solvated, so that there are fewer in the state capable of absorbing the radiation at  $\lambda = 365 \text{ m}\mu$  and the value of  $\log K$  at that point falls. The degree of solvation can be calculated as follows:-

Let  $\alpha$  = degree of solvation

"  $c$  = total molar concentration of stannic iodide.

$\therefore$  conc. of solvated molecules =  $\alpha c$

" " unsolvated " =  $(1 - \alpha) c$

Let  $K_m$  = measured absorption coefficient at  $\lambda = 365 \text{ m}\mu$

"  $K_1$  = absorption coefficient for unsolvated  $\text{SnI}_4$   
at  $\lambda = 365 \text{ m}\mu$

"  $K_2$  = absorption coefficient for solvated  $\text{SnI}_4$   
at  $\lambda = 365 \text{ m}\mu$

( $K_1$  and  $K_2$  are obtained from the curves for stannic iodide in hexane and in ethyl alcohol respectively).

$$K_m c = K_2 \alpha c + K_1 c (1 - \alpha)$$

$$K_m = K_2 \alpha + K_1 (1 - \alpha)$$

This equation gives  $\alpha$  for each concentration of ethyl alcohol. If  $\alpha$  is plotted against the concentration of ethyl alcohol a smooth curve is obtained (Fig. 5A).

The curve for stannic iodide in ethyl alcohol  
and/

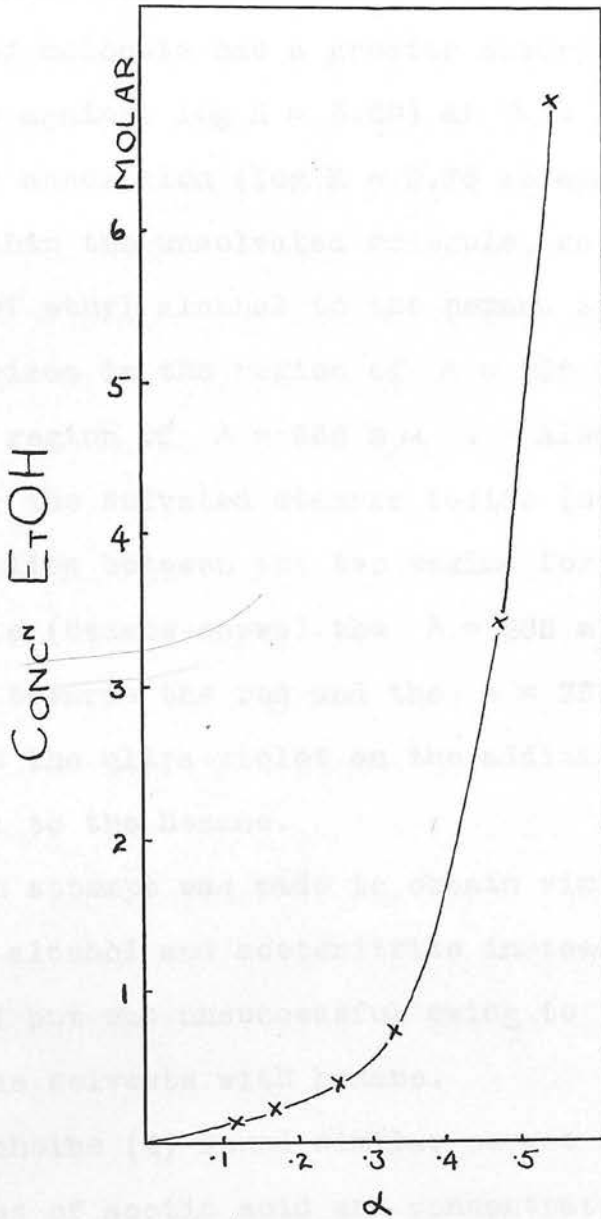


FIG. 5A.

and hexane is the resultant of the curves for stannic iodide in ethyl alcohol and in hexane; on increasing the amount of ethyl alcohol, more and more stannic iodide molecules become solvated and the ethyl alcohol curve gradually predominates. At  $\lambda = 285 \text{ m}\mu$  the solvated molecule has a greater absorption ( $\log K = 3.67$  as against  $\log K = 3.60$ ) at  $\lambda = 365 \text{ m}\mu$  a much smaller absorption ( $\log K = 2.76$  as against  $\log K = 3.90$ ) than the unsolvated molecule, so that on addition of ethyl alcohol to the hexane solution the curve rises in the region of  $\lambda = 285 \text{ m}\mu$  and falls in the region of  $\lambda = 365 \text{ m}\mu$ . Also as the maximum for the solvated stannic iodide (ethyl alcohol curve) lies between the two maxima for the unsolvated molecule (hexane curve) the  $\lambda = 285 \text{ m}\mu$  maximum shifts towards the red and the  $\lambda = 365 \text{ m}\mu$  maximum towards the ultra-violet on the addition of ethyl alcohol to the hexane.

An attempt was made to obtain similar series with methyl alcohol and acetonitrile instead of ethyl alcohol but was unsuccessful owing to the immiscibility of these solvents with hexane.

Scheibe (4) found similar curves for phoron in mixtures of acetic acid and concentrated sulphuric acid and explains his results by assuming that there are two states of phoron, one stable in concentrated sulphuric acid and the other in acetic acid.

(d)/

(d) Stannic Bromide.

Stannic bromide can be pictured as having a molecule similarly constituted to that of  $\text{SnI}_4$  but with the iodine atoms replaced by bromine atoms, so that the tin atom is surrounded by a shell of deformed Br - Br linkages, and the spectrum of stannic bromide might, on these grounds, be similar to that of bromine in one of Getman's "unsaturated" solvents. Owing to the fact that the bromine atom is smaller than the iodine atom it will have a greater deforming effect on the electronic orbits of the tin orbit, so that the bond between the tin and bromine in  $\text{SnBr}_4$  will be stronger than the bond between tin and iodine in  $\text{SnI}_4$ ; hence the absorption of Sn - Br bond will approximate less closely to that of a bromine atom in a deformed bromine molecule; therefore there will probably be less similarity between the curves for  $\text{Sn Br}_4$  and  $\text{Br}_2$  than between the curves for  $\text{SnI}_4$  and  $\text{I}_2$ . In practice it was found that  $\text{Sn Br}_4$  in hexane gave a single maximum at  $\lambda = 239 \text{ m}\mu$ , while Bovis (50) for a solution of bromine in ethyl alcohol found a single maximum at  $\lambda = 264 \text{ m}\mu$ . Thus there is found to be a certain similarity between the two curves, but it is less than in the case of iodine.

(e)/

(e) Comparison of Behaviour of the Curves for Stannic Iodide and Stannic Bromide.

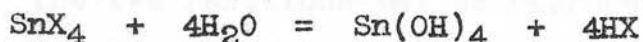
Stannic iodide and stannic bromide being very similar compounds it might have been expected that their absorption curves would behave in a similar manner under the same influence, but we see from the above results that this is not the case. Solutions of stannic iodide standing over metallic tin are stable and obey Beer's Law, while if no tin is present the solution gradually turns brown. On the other hand, solutions of stannic bromide remain colourless but their absorption changes with time whether metallic tin is present or not. In the case of both halides a white deposit forms slowly.

The stannic halides are fairly readily hydrolysed by water (52) and the following explanation of the above observations is based on this fact.

The solvent, however carefully dried, absorbs moisture from the air whenever the stopper of the flask is removed, and from the surfaces of the various vessels used. These traces of moisture gradually hydrolyse the stannic halide with the eventual formation of stannic hydroxide which appears as the white deposit observed. Probably various oxyhalides are formed as intermediate products during the hydrolysis, but they in their turn are hydrolysed until finally stannic/



stannic hydroxide is reached. The complete reaction may be represented by the equation

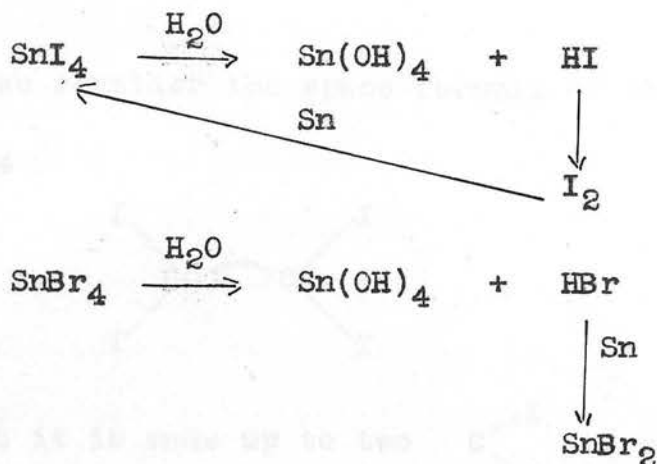


In the case of stannic iodide the hydrogen iodide liberated is unstable (especially in sunlight) (53) and readily forms free iodine which produces the brown colouration of the solution; iodine, however, reacts with metallic tin with the formation of stannic iodide, so that, in the presence of tin, the iodine is taken up as soon as it is formed, and the concentration of stannic iodide in the solution remains constant, and the equivalent of the tin that goes into solution as stannic iodide is deposited as stannic hydroxide.

Hydrogen bromide, on the other hand, does not readily decompose with the formation of free bromine so the solution remains colourless, but hydrogen bromide reacts with metallic tin to form stannous bromide. Thus in solutions of stannic bromide the concentration of stannic bromide gradually decreases, and that of stannous bromide increases (in solvents in which it is soluble) as the hydrolysis proceeds, so that the maximum on the curve for stannic bromide gradually falls. Stannous bromide is insoluble in hexane but soluble in methyl alcohol so that with the former solvent there is no shift of the maximum wavelength, but with the latter solvent the curve gradually moves/

moves towards that for stannous bromide, i.e. towards the ultra-violet.

The two reactions may be represented diagrammatically thus:-



(f) Stannic Chloride in Hexane.

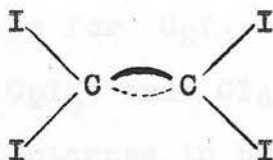
Such a small proportion this curve was measurable with the apparatus employed, that it is almost impossible to discuss its position except in relation to the curves for stannic iodide and stannic bromide, for on the theory that the absorption of the stannic halides is mainly due to the absorption of deformed halogen atoms it is interesting to note that the absorptions of the series  $\text{Sn I}_4$   $\text{SnBr}_4$  and  $\text{Sn Cl}_4$  changes in a manner similar to the change in absorption in the series  $\text{I}_2$   $\text{Br}_2$  and  $\text{Cl}_2$ .

### 3. SOME IODINE DERIVATIVES OF CARBON.

#### (a) The effect of change in molecular structure.

##### C<sub>2</sub>I<sub>4</sub> .

When we consider the space formula of the compound C<sub>2</sub>I<sub>4</sub>



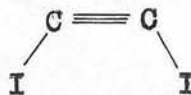
we see that it is made up to two  $\begin{matrix} \text{I} \\ \text{C} \\ \text{I} \end{matrix}$  groups held together by a double bond; so that it consists of two regular tetrahedra with one edge (and two vertices) in common. The other two vertices in each case are occupied by iodine atoms. On the Baeyer strain theory there will be a change in the angle between the valencies which join the carbon atoms, i.e. they will be brought closer together than they are in such compounds as CH<sub>4</sub>. But the angles between the valencies joining the iodine atoms to the carbons will not be appreciably altered from that in CI<sub>4</sub>, so that the distance between the iodine atoms approximates to that between the iodine atoms in CI<sub>4</sub> or CHI<sub>3</sub> .

The resemblance of the CHI<sub>3</sub> and KI<sub>3</sub> spectra, discussed earlier, led to the suggestion that the C - I bond is a weak one, sufficiently so as to permit/

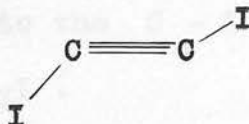
permit us to consider the  $\text{CHI}_3$  and  $\text{KI}_3$  spectra as arising from the iodine atoms in the  $\text{I}_3$  group. In  $\text{C}_2\text{I}_4$  we have again weak C - I bonds, and two pairs of two adjacent iodine atoms. In  $\text{CI}_4$  the carbon is completely surrounded by I - I linkages, but this is not the case in either  $\text{CHI}_3$  or  $\text{C}_2\text{I}_4$  so it would be more reasonable to look for a similarity between the curves for  $\text{C}_2\text{I}_4$  and  $\text{CHI}_3$  than between the curves for  $\text{C}_2\text{I}_4$  and  $\text{CI}_4$ . (In all cases the curves of the substances in hexane, i.e. the curves for the substances in the normal unsolvated states are under consideration; solvation effects will be examined later.) In fulfilment of this expectation it was found in practice that the curve for  $\text{C}_2\text{I}_4$  is very similar to that for  $\text{CHI}_3$ .

### $\text{C}_2\text{I}_2$

There are two possibilities in the space model of  $\text{C}_2\text{I}_2$ . (1) in which the iodine atoms are on the same side of the carbon linkage



and (2) in which the iodine atoms are on different sides of the carbon linkages



But  $\text{C}_2\text{I}_2$  has been shown (54) xxxxxxxxxxxxxx to have a dipole moment of 0.33 so that formula (1) is the more probable. In this molecule the iodine atoms are further/

further apart than they are in  $C_2I_4$  so that the attractive forces between the iodine atoms will be much less but there will still be some attraction. Owing to the decrease in the attraction between the iodine atoms the attraction between the carbon and iodine atoms will be more apparent. Also owing to the fact that there is only one I - I "bond" in the  $C_2I_2$  molecule (compared with four in  $C_2I_4$ ) it would be expected that the intensity of absorption would be decreased.

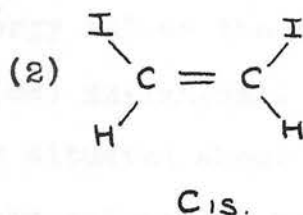
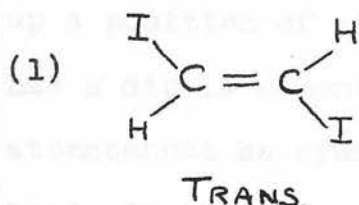
This is supported by the curve obtained experimentally for  $C_2I_2$ . The inflexion stretching from approximately  $\lambda = 335 \text{ m}\mu$  to  $\lambda = 300 \text{ m}\mu$  corresponds to the region of absorption in  $C_2I_4$  which was attributed to the I - I bonds. The value of  $\log K$  (which on the quantum theory is a measure of the frequency of the occurrence of the absorption process) is only 1.9 at the inflexion on the  $C_2I_2$  curve compared with an average value of  $\log K = 3.0$  for the same on the  $C_2I_4$  curve.

The maximum to which the  $C_2I_2$  curve rises at  $\lambda = 270 \text{ m}\mu$  is probably that due to the C - I linkage which appears at  $\lambda = 250 \text{ m}\mu$  in  $CH_3I$ .

### $C_2H_2I_2$ .

$C_2H_2I_2$  can exist in two isomeric forms, one solid and one liquid. The two possible formulae are/

are



The trans-form has been shown as would be expected, to have zero dipole moment (55) xxxxxxxxxxxx and a dipole moment of .75 has been found for the cis-form (56) xxxxxxxxxxxx. By the preparation of fumaric and maleic acids from the two forms Keiser (57) (Am: Chem: Jour. 48 518) decided that the solid variety had the trans-form and the liquid variety the cis-form. When either form is heated to 146°C for several hours an equilibrium between the forms is established. In the present work the solid form (i.e. the non-polar form) was prepared.

In the trans-form the iodine atoms are on opposite sides of the molecule so that there will be only a slight absorption due to an I - I "bond" but instead it will be almost all due to the C - I bonds. In accordance with this the curve for C<sub>2</sub>H<sub>2</sub>I<sub>2</sub> consists of a branch rising towards the ultra-violet; this is probably the absorption due to the C - I bonds. Just about  $\lambda = 307 \text{ m}\mu$  there is an inflexion in the curve which may be caused by the I - I "bond".

#### C<sub>2</sub>H<sub>4</sub>I<sub>2</sub>

There is only one possible formula for this compound because of the free rotation of the atoms/

atoms round the bonds. The atoms will tend to take up a position of minimum energy but as the molecule has a dipole moment of 1.3 (58) xxxxxxxxxxxx the iodine atoms cannot be symmetrically situated about the C - C bond; in that case some slight attraction between the iodine atoms will exist; in agreement with this it is found that although the main absorption consists of a branch rising towards the ultra-violet (presumably due to the C - I linkages) there is a kink in the curve at  $\lambda = 307 \text{ m}\mu$  which is in the region of absorption attributed to the I - I "bond".

### CH<sub>2</sub>I<sub>2</sub>

In the CH<sub>2</sub>I<sub>2</sub> molecule where the iodine and hydrogen atoms are at the vertices of a tetrahedron at the centre of which is the carbon atom, the iodine atoms are bound to be adjacent to each other in space, so that there must be some interaction between them, and one would expect it to be of the same order of magnitude as in CHI<sub>3</sub>. It is rather difficult to deduce any agreement with this from the experimental curve owing to the fact that the maximum attributed to the C - I linkage has moved so far towards the visible compared with its position in CH<sub>3</sub>I (to  $\lambda = 297 \text{ m}\mu$  in fact) that it overlaps into the region attributed to the I - I linkage. However it may be noted that a change in the direction of the curve occurs at  $\lambda = 354 \text{ m}\mu$  while the first maximum of the I - I "bond"/

"bond" in  $C_2I_4$  occurs at  $\lambda = 330 \text{ m}\mu$  .

The shift of the C - I maximum from its position at  $\lambda = 250 \text{ m}\mu$  in  $CH_3I$  to  $\lambda = 297 \text{ m}\mu$  in  $CH_2I_2$  is due according to Lowry and Sass to the weakening of the C - I bond by the other iodine atom in the molecule.

(b) The Solvent Effect.

When an absorbing substance is dissolved in a solvent the electrons of the absorbing group will be under the influence of the dipole (if any) of the solvent. But as the field of a dipole falls off very rapidly with distance the effect of the solvent dipole is negligible compared with that of the dipole of the chromophore group; and any deformation of the electron system of the absorbing group caused by the polar solvent will be small compared with the potential energy changes in the solute molecule due to the formation of solvation complexes and may, therefore, be neglected in a consideration of the effect of change of solvent on the absorption of the solute.

Let  $\eta$  = the excitation energy of the molecule when not influenced by solvent.

There are several cases to be considered:-

1. A polar molecule dissolved in a non-polar unpolarisable solvent e.g. hexane.

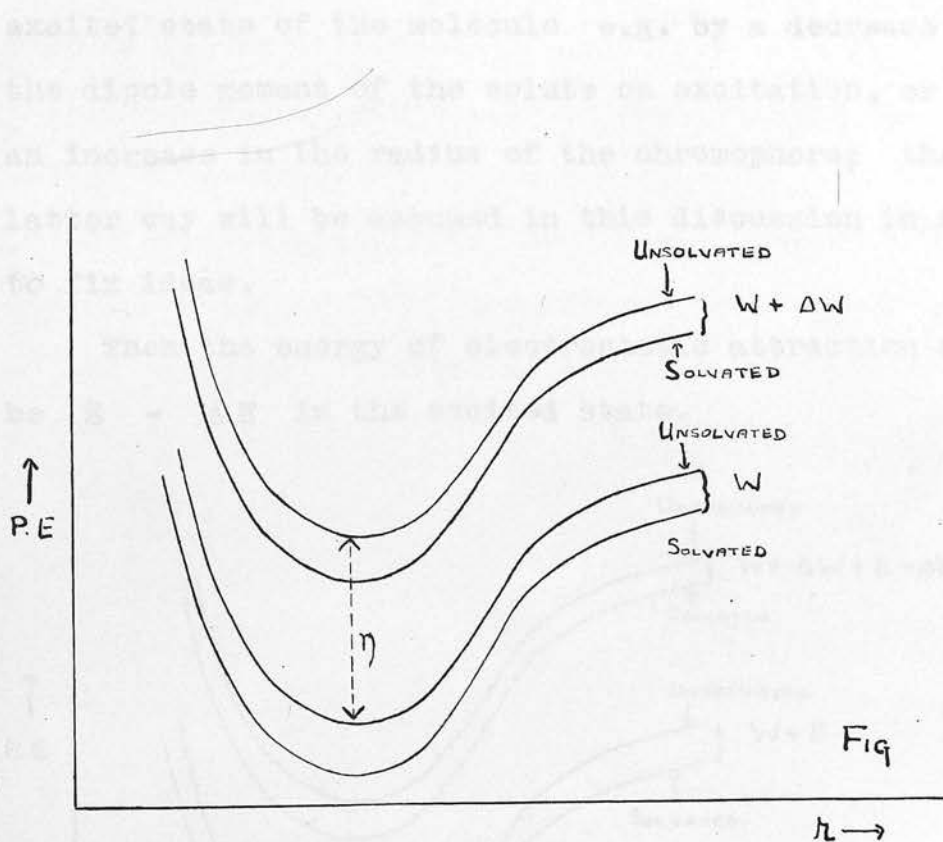
In/



In this case there will be no electrostatic interaction with the solvent.

Let the solvation energy in the unexcited state be  $W$ , this energy is due to van der Waal's type of forces, and represents the energy required to remove the molecule from its solvation shell.

When the molecule is excited the solvation energy will change to  $W + \Delta W$ . It will be assumed that the solvation energy is altered to approximately the same extent  $\Delta W$  in all the cases discussed, this assumption is permissible because in each case discussed here we are dealing with the same chromophore.



Energy of light absorbed by solvated molecule  
on being excited =  $\eta - \Delta W$

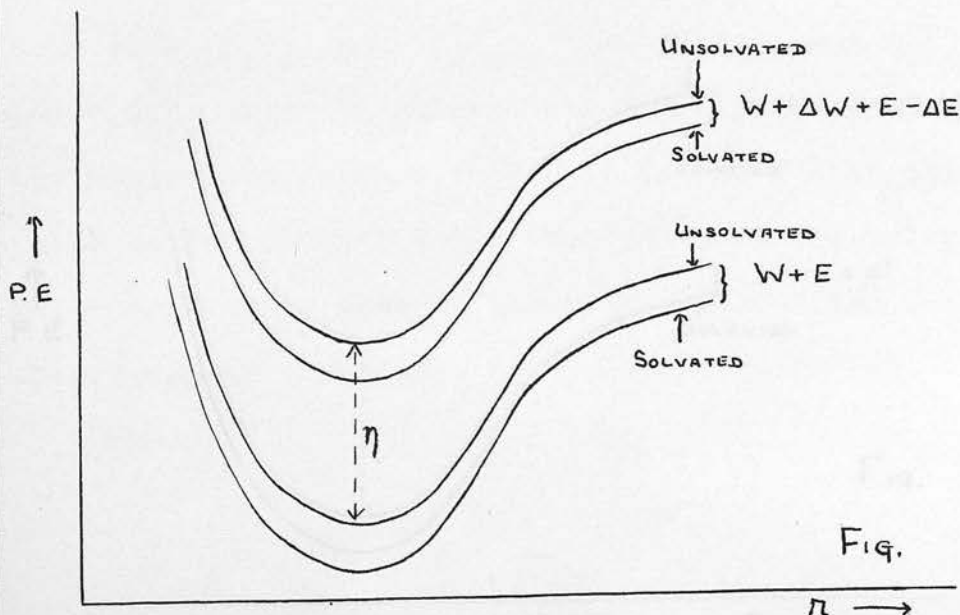
$$\text{Wave-length of absorbed light} = \frac{hc}{\eta - \Delta W} \text{----- (1)}$$

(2) A polar molecule dissolved in a polar but unpolarisable solvent such as EtOH .

In the unexcited state the energy of solvation is  $W$  as before; added to this is the energy  $E$  of electrostatic attraction between the dipoles of the solute and of the solvent.

In the excited state the solvation energy becomes  $W + \Delta W$ . The electrostatic energy may be decreased in several ways in passing from the unexcited to the excited state of the molecule e.g. by a decrease in the dipole moment of the solute on excitation, or by an increase in the radius of the chromophore; the latter way will be assumed in this discussion in order to fix ideas.

Then the energy of electrostatic attraction will be  $E - \Delta E$  in the excited state.



Energy of light required to excite solvated mole-

$$\text{cule} = \eta - \Delta W + \Delta E$$

$$\text{Wave-length of absorbed light} = \frac{hc}{\eta - \Delta W + \Delta E} \quad (2).$$

This is less than (1) so that a shift in absorption towards the ultra-violet would be expected.

(3) A polar molecule dissolved in a non-polar but polarisable solvent e.g.  $\text{CCl}_4$ .

In the unexcited state the solvation energy is  $W$ . The solute molecule being polar will induce a dipole in the solvent molecules. To withdraw the solute molecule from the solvent, energy  $W - H$  must be exerted where  $H$  is the energy required to create the dipoles in the solvent.

Putting in the electrostatic attraction term  $E'$  the energy of solution is  $W - H + E'$ .

In the excited state both the  $H$  and  $E'$  terms will be diminished due to the increased radius of the molecule, so the energy of solution is now  $W + \Delta W - (H - \Delta H) + (E' - \Delta E')$ .

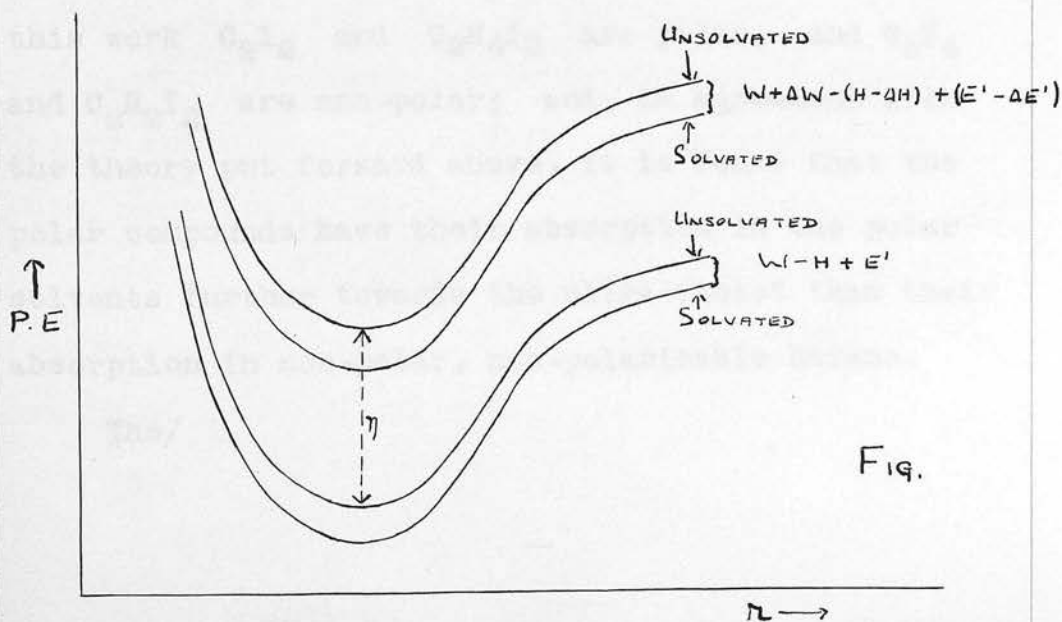


Fig.

Energy of light required to excite molecule in solvated state =  $\eta - \Delta W - \Delta H + \Delta E'$

$$\text{Wave-length of absorbed light} = \frac{hc}{\eta - \Delta W - \Delta H + \Delta E'} \quad (3).$$

where  $\Delta E'$  must be less than  $\Delta H$ .

The value of the wave-length of the absorbed light in this case is greater than in (1) so that the absorption should shift towards the red.

It may be mentioned that for this case, the present theory is in agreement with Kundt's Rule. This rule is no longer regarded as having general validity, but Scheibe has found also that polar solutes in non-polar polarisable solvents have their absorption curves shifted towards the red.

(4) A non-polar molecule in a non-polar solvent

whether the latter is polarisable or not. This is really the same as case 1. because there will be no electrostatic interaction between the solute and solvent.

In the series of carbon compounds examined in this work  $C_2I_2$  and  $C_2H_4I_2$  are polar, and  $C_2I_4$  and  $C_2H_2I_2$  are non-polar; and, in agreement with the theory put forward above, it is found that the polar compounds have their absorption in the polar solvents further towards the ultra-violet than their absorption in non-polar, non-polarisable hexane.

The/

The non-polar compounds here investigated also show a shift towards the ultra-violet in their absorption spectra, when dissolved in polar solvents. These non-polar compounds consist, however, in each case of a set of polar bonds arranged symmetrically so as to have no resultant moment. In a polar solvent, each bond will probably interact with the electrostatic field of the solvent molecule nearest to it, so that the case will approximate to that of a polar solute and a polar solvent.

The effect in  $\text{CCl}_4$  (non-polar but polarisable) is not however the same as for a polar solute, for even if each bond interacts with its adjacent solvent molecule, the effect will be symmetrical throughout the whole solvation complex; i.e. the solvent will on the whole be unaffected. The observed coincidence of the curves of compounds of this type in hexane and  $\text{CCl}_4$  is in agreement with this.

$\text{C}_2\text{I}_2$   $\text{C}_2\text{H}_2\text{I}_2$  and  $\text{C}_2\text{H}_4\text{I}_2$  all show a shift in absorption towards the red in  $\text{CCl}_4$  as compared with hexane.  $\text{C}_2\text{I}_2$  and  $\text{C}_2\text{H}_4\text{I}_2$  are both polar and so the results for these compounds support the above theory which predicts a shift towards the red under such conditions.

The  $\text{C}_2\text{H}_2\text{I}_2$  used was the trans-form, which is non-polar. The reason for the observed shift to the red in  $\text{CCl}_4$  is not evident.

The/

The polarity of  $\text{CH}_2\text{I}_2$  does not seem to have been investigated.

$\text{C}_2\text{I}_4$  is non-polar but it seems impossible to draw any conclusions from the results obtained as the variation in the curves is not regular and other effects seem to enter in.

In general there does not appear to be any strict parallelism between the polarity of the solvent and the amount of shift of the absorption. The reason for this is probably the occurrence of specific solute - solvent effects as distinct from general effects.

SUMMARY.

The absorption spectra of  $\text{SnI}_4$  ,  $\text{SnBr}_4$  ,  $\text{SnCl}_4$  ,  $\text{SnBr}_2$  ,  $\text{C}_2\text{I}_4$  ,  $\text{C}_2\text{I}_2$  ,  $\text{C}_2\text{H}_2\text{I}_2$  ,  $\text{C}_2\text{H}_4\text{I}_2$  , and  $\text{CH}_2\text{I}_2$  have been measured in various solvents. It is shown that results of previous workers for  $\text{SnI}_4$  have been subject to errors due to decomposition of the compound. This decomposition was prevented in the curves obtained in this work. The decomposition could not be stopped in the case of  $\text{SnBr}_4$  and a hypothesis is put to explain this difference in the behaviour of  $\text{SnI}_4$  and  $\text{SnBr}_4$  .

The spectra of the iodine compounds are discussed on the basis of the theory that in compounds containing several iodine atoms and weak X - I bonds part of the absorption is due to interaction between the iodine atoms.

The effect of change of solvent on the absorption of the carbon compounds was studied and a tentative explanation put forward.

In conclusion I wish to express my thanks to Dr E.B. Ludlam and Dr H.J. Walls of this University for their valuable encouragement and advice during the work.

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