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Specific heat capacity and thermal conductivity of PEEK/Ag nanoparticles composites determined by Modulated-Temperature Differential Scanning Calorimetry

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A B S T R A C T

The thermal conductivity accurate measurement of polymer based composites is a challenge: it would allow us to understand the mechanisms of thermal transport in such materials. Silver nanoparticles were introduced in Polyetheretherketone matrix and their influence on thermal properties was studied. Thermal conductivity and specific heat capacity of composites were determined by Modulated-Temperature Differential Scanning Calorimetry and analysed as a function of particles volume content and temperature. The specific heat capacity of the composites decreases with increasing silver particles content below the electrical percolation threshold. Above the electrical percolation threshold the specific heat capacity decreases more slowly and converge toward the specific heat capacity of compressed silver nanoparticles. The evolution of the thermal conductivity with filler content exhibits a non-linear profile. Experimental data are coherent with the Maxwell model suggesting continuity of the polymer matrix and a contribution of the silver particles to the effective thermal conductivity greater than volume effect. The temperature dependence of the composites thermal conductivity is characteristic of amorphous phase, while a transition from vitreous-like to crystalline-like behaviour of the specific heat capacity is observed with the introduction of metallic particles.

Keywords:

Thermal conductivity
Specific heat capacity
Composite
PEEK
Silver
Electrical percolation

1. Introduction

The development of polymer based composites has heightened the need for the multi-scale understanding of their properties. The increasing use of these materials for weight saving purpose, has led to the optimisation of filler content in order to enhance thermal and electrical properties while maintaining good mechanical properties. Mechanical and electrical properties have been extensively studied whereas thermal properties like thermal conductivity and specific heat capacity are less investigated.

The electrical conductivity mechanisms and their dependence on particle content in polymer based composites is well understood. We observe a sharp transition from insulating to conductive material through a percolation phenomenon as the filler content

increases. Above the percolation threshold, the electrical conductivity can be modelled by a power law that accounts for the particles aspect ratio and system dimensionality [1–3].

In this context, it appears interesting to compare thermal properties with electrical conductivity in order to identify possible characteristic behaviour of these properties as a function of filler content. The challenge relies on the identification of the mechanisms involved, and their contribution to the composite thermal properties. Filler chemical nature, filler content, aspect ratio, polymer physical structure, polymer/particles interactions and interfaces are parameters to be taken into account.

It has been shown that the introduction of conductive filler in the polymer matrix improves thermal conductivity [4,5]. In particular, carbon nanotubes [6,7] have been extensively studied for their high thermal conductivity, in various polymer matrices like polymethylmethacrylate [8], polyethylene [9] or polyvinylchloride [10]. Few studies have been devoted to polymer/metal composites for thermal properties improvement [4,11,12]. Nevertheless, high level of electrical conductivity can be reached for low filler content with the use of high aspect ratio metallic particles in thermoplastic polymers [13–15].

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Some models have been proposed to describe thermal conductivity dependence on filler content [16–18] for various particles shapes, including percolation theory [8,19–21], but experimental studies show some discrepancies between measured data and predictions [22–24]. In addition, the variability of measurement techniques and experimental conditions prevents a real consensus. The composite specific heat capacity can be expressed as the constituents specific heat capacities weighted average. Some studies confirm experimentally this model [25,26] and others report a change of behaviour for electrically conductive composites [27,28]. The correlation between the formation of an electrically conductive pathway and the specific heat capacity variation with filler content has not been demonstrated.

Silver nanoparticles were introduced in Polyetheretherketone matrix and their influence on thermal properties was studied. Thermal conductivity and specific heat capacity of composites were determined by Modulated-Temperature Differential Scanning Calorimetry. The thermal conductivity and specific heat capacity dependences on particle volume fraction have been studied and compared with the electrical conductivity evolution in the electrical percolation threshold region.

2. Materials and methods

2.1. Rough materials

Polyetheretherketone (PEEK) is a high-performance thermoplastic of the polyaryletherketones (PAEK) family. It is a semi-crystalline linear polymer with aromatic structure that exhibits high thermal stability, chemical resistance and mechanical properties. Its high glass transition temperature ($T_g = 140$ °C) and melting point ($T_m = 340$ °C) make this polymer suitable for aeronautical applications [29,30]. Its crystallinity is about 35%. The PEEK matrix selected for this study was the Vestakeep 2000PF from Evonik. This commercially available grade is in powder shape of 50 μm particle size.

Silver nanoparticles (AgNP) from Sigma Aldrich with particle size <100 nm and 99.5% purity, were used in this study to elaborate the PEEK/AgNP composites.

PEEK/AgNP composites were elaborated by melt blending. Silver nanoparticles were first dispersed in ethanol by means of ultrasonic stirring then added to PEEK powder and stirred again. The ethanol was removed from the mixture with a rotary vacuum evaporator and the powdered composite was hot pressed at 380 °C for 15 min. The obtained bulk material was then repeatedly cut in tiny pellets and pressed again (three to four times) until homogenization. Samples of compressed silver nanoparticles were elaborated with a hydraulic press under 0.4 GPa.

2.2. Volumic mass and Ag content determination

Volumic mass of composites ρ_{comp} was determined using a pycnometer and distilled water as filling liquid. The method [31] relies on the measurement of sample mass and volume, by means of successive weighting of empty pycnometer m_1 , sample in pycnometer m_2 , sample in pycnometer filled with water m_3 , and pycnometer filled with water m_4 . The composites volumic mass expression is given by Eq (1).

$$\rho_{comp} = \rho_{water} \times \frac{(m_2 - m_1)}{(m_4 - m_1) - (m_3 - m_2)} \quad (1)$$

The quantity $(m_2 - m_1)$ corresponds to the sample mass and $(m_4 - m_1) - (m_3 - m_2)$ to the mass of water displaced by the sample. The sample volume is equal to the mass of displaced water divided

by the water volumic mass ρ_{water} . The apparent volumic mass of the compressed AgNP sample was found at 8.3 g cm⁻³ which correspond to about 80% of the silver bulk density (10.50 g cm⁻³ at 25 °C [32]). Volume fraction of silver nanoparticle, x , was deduced from Eq (2) with $\rho_{PEEK} = 1.32$ g cm⁻³ and $\rho_{Ag} = 10.50$ g cm⁻³, volumic masses of PEEK and bulk silver respectively.

$$\rho_{comp} = (1 - x)\rho_{PEEK} + x\rho_{Ag} \quad (2)$$

11 composites were elaborated with particles content varying from 1.3 vol% (9.5 wt%) to 27.3 vol% (74.9 wt%).

2.3. Electrical conductivity

A Novocontrol broadband dielectric spectrometer (with SI 1260 gain/phase analyzer) has been used to determine DC conductivity of insulating composites ($Z > 10$ Ω). Dielectric experiments were carried out at 20 °C in the frequency range 10⁻² Hz–10⁶ Hz. In disordered materials, the real part $\sigma'(\omega)$ of the measured electrical conductivity is frequency dependant:

$$\sigma'(\omega) = \sigma_{DC} + A\omega^s \quad (3)$$

At 10⁻² Hz, the frequency dependant term $A\omega^s$ can be neglected so the reported conductivity is considered equivalent to the DC conductivity.

$$\sigma'(10^{-2}\text{Hz}) \approx \sigma_{DC} \quad (4)$$

The resistivity of conductive composites ($Z < 10$ Ω) was measured with a Keithley 2420 source meter using the four wire configuration. The resistivity of compressed silver nanoparticles was also measured with this protocol.

The samples were disks with a thickness of 500 μm and a diameter of 20 mm, placed between two gold plated electrodes with a 10 mm diameter. The samples were previously coated with a thin layer of silver ink to reduce contact resistivity. Data are averaged electrical conductivity of three to five measurements for each samples.

2.4. Modulated temperature differential scanning calorimetry

Modulated temperature differential scanning calorimetry (MT-DSC) relies on the analysis of the modulated heat flow response of a sample submitted to a modulated temperature program. This technique was introduced by Reading in the early 1990s [33], and enables to separate the in-phase and out-of-phase responses from the total modulated heat flow. The MT-DSC 2920 from TA Instrument has been used to measure specific heat capacity and thermal conductivity of PEEK/AgNP composites. Specific heat capacity C_p measurement is possible in quasi-isothermal conditions from the in-phase heat flow. In the absence of any thermal event [34], it can be expressed as:

$$C_p = K \cdot \frac{A_{HF}}{A_T} \quad (5)$$

With A_T the temperature modulation amplitude, A_{HF} the heat flow amplitude, and K the calibration constant that takes into account the thermal responses of the reference, pans and furnace [35].

We measured specific heat capacity of PEEK/AgNP composites from the five last of 20 min isotherms at 17 °C, 32 °C and 47 °C. Samples were flat disks encapsulated in aluminium pan of the same diameter. The modulation parameters were chosen to ensure uniform and stationary heat transfer across the sample [36]. A 100 s

modulation period and 1 °C modulation amplitude were used for samples with a diameter of 6 mm, a thickness inferior to 500 μm and masses between 10 mg and 30 mg. An empty sealed aluminium pan with mass matching that of sample pan was used as reference material. A sapphire disk of 26.4 mg was used as calibrant. Calibration runs were carried out at 27 °C before and after each measurement to determine the calibration constant K . With the same protocol, we also measured the specific heat capacity of compressed silver nanoparticles and bulk silver at 20 °C, 30 °C, 40 °C, 50 °C and 60 °C. The presented results are averaged C_p of three to five measurements for each sample.

The MT-DSC thermal conductivity can be obtained by the determination of a thick sample apparent heat capacity, in intentionally unfavourable experimental conditions. When applying a temperature modulation to a thick, non-encapsulated sample, we cannot consider that the whole sample follows the maximum variation, as it is supposed for a thin encapsulated sample. There is a temperature gradient through the sample which is the expression of finite thermal conductivity of the sample [37]. The apparent thermal conductivity λ_0 is given by Eq (6).

$$\lambda_0 = \frac{8eC^2}{C_p m d^2 T} \quad (6)$$

With C the apparent heat capacity of the thick sample, e its thickness, d its diameter and m its mass. C_p is the specific heat capacity of the sample and T the modulation period [38].

This value has to be corrected from heat losses through the thick sample lateral surface with the use of a Polystyrene calibrant of exactly the same size. The corrected thermal conductivity is then given by Eq (7) and Eq (8).

$$\lambda = \frac{1}{2} \left(\lambda_0 - 2D + \sqrt{\lambda_0^2 - 4D\lambda_0} \right) \quad (7)$$

With

$$D = \sqrt{\lambda_{0,PS}\lambda_{PS}} - \lambda_{PS} \quad (8)$$

where $\lambda_{0,PS}$ is the measured thermal conductivity of the Polystyrene calibrant and λ_{PS} the thermal conductivity of the Polystyrene from the literature [39].

Thermal conductivity of PEEK/AgNP composites at 17 °C, 32 °C and 47 °C were determined using the same modulation parameters as for specific heat capacity measurements. The apparent heat capacity, C , measurements were performed on right circular cylinders. The thicknesses varied from 2.9 mm to 3.8 mm and the diameter was 6.0 mm. The samples were placed into the furnace on aluminium disks (GoodFellow) with a 50 μm thickness and a 6.0 mm diameter. One of the disks was also used on the reference position. Each thermal conductivity measurement was corrected with a Polystyrene calibrant and the presented results are averages of three to five measurements. The thermal conductivity of Polystyrene calibrant was measured in the same conditions than composites.

3. Results and discussions

3.1. Electrical charge transport in composites

DC electrical conductivity of the composites as a function of AgNP volume content is displayed in Fig. 1. We observe a sharp transition from insulating to conductive material due to percolation phenomenon.

Below a critical filler volume fraction p_c , called the percolation

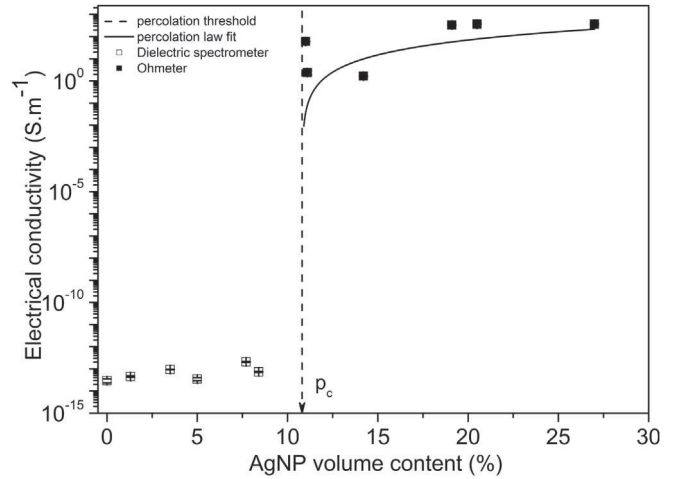


Fig. 1. Evolution of electrical conductivity with AgNP volume content at room temperature. Measurements with dielectric spectrometer (\square), measurements with ohmmeter (\blacksquare), percolation law fit (—), percolation threshold p_c (---).

threshold, the metallic particles are isolated in a continuous insulating polymer matrix. The matrix contribution to the electrical conductivity of the composite is predominant and there is no contribution of the metallic particles. The measured electrical conductivity of PEEK is $\sigma_{PEEK} = 2.9 \cdot 10^{-14} \text{ S m}^{-1}$ which is consistent with literature that gives a volume resistivity $\rho_{PEEK} \sim 10^{14} \Omega \text{ m}^{-1}$ [40].

In this study, the experimental percolation threshold p_c is around 10.8 vol%. For non-overlapping spherical particles dispersed in a continuous matrix, the theoretical percolation threshold is predicted between 14% and 16% [3]. The experimental value is lower than the theoretical prediction and can be attributed to the formation of silver nanoparticles aggregates, called grapes [41], of higher apparent aspect ratio. It then can form conductive paths for a lower volume content. We can still consider a homogeneous dispersion of these aggregates. This reflects however, the difficulty to obtain an optimized dispersion of particles for filler volume content above 10%.

Above the percolation threshold, the electrical conductivity follows a power law.

$$\sigma = \sigma_0 (p - p_c)^t \quad (9)$$

In Eq (9), σ_0 is the conductivity of the particles network, p the volume fraction of particles, p_c their volume fraction at the percolation threshold, and t the critical exponent which describes the system dimensionality. The parameters σ_0 and t strongly depend on percolation pathway morphology. For 3D networks, t has to lie between 1.60 and 2.00. The best fit of our experimental data gives $t = 1.98$. σ_0 represents the electrical conductivity of the silver nanoparticles network. In our case $\sigma_0 = 6.7 \cdot 10^4 \text{ S m}^{-1}$, this value is in good agreement with the measured electrical conductivity of compressed silver nanoparticles $\sigma_{AgNP} = 7.5 \cdot 10^4 \text{ S m}^{-1}$ and seems to accurately represent the silver particles network conductivity above the percolation threshold. The maximum electrical conductivity reached for composites beyond the percolation threshold is around $\sigma_{max} = 3.6 \cdot 10^2 \text{ S m}^{-1}$, this value is two decades lower than the electrical conductivity of compressed silver nanoparticles. This indicates that the connectivity of the system has not changed, i.e., particle aggregates are not in direct contact even for highly filled composites.

3.2. Specific heat capacity of composites

The specific heat capacity of composites, PEEK, compressed silver nanoparticles and bulk silver was measured and are displayed in Table 1.

The measured data are coherent with the literature and standard deviations are low. It shows that polymers have higher specific heat capacity than crystals, because the glass has lower density than the crystal [43].

Specific heat capacities of PEEK/AgNP composites at 17 °C are represented on Fig. 2 as a function of silver nanoparticles weight content. We observe a decrease of the composite specific heat capacity with AgNP content. This is associated to the low value of metallic particles heat capacity.

Below 49.5 wt%, the decrease of composite specific heat capacity can be fitted by Eq (10). The specific heat capacity is actually equal to the weighted average $C_p(w)$ of each constituent heat capacities in the case of an isotropic composite with constant pressure and volume (negligible thermal expansion) with no local strain or stress [44].

$$C_p(w) = wC_{p,Ag\ bulk} + (1 - w)C_{p,PEEK} \quad (10)$$

w is the weight fraction of silver particles, $C_{p,Ag\ bulk}$ the specific heat capacity of bulk silver, and $C_{p,PEEK}$ the specific heat capacity of PEEK. Boudenne et al. have also reported a good correlation between measured specific heat capacities and values predicted by weighted average, for polymer filled with metallic micro particles up to 45 wt% [26,45].

Above 49.5 wt%, Eq (10) no longer fits experimental data. The specific heat capacity decreases more slowly approaching the specific heat capacity of compressed silver nanoparticles (Table 1). This evolution of the specific heat capacity can be attributed to composite heterogeneity at high filler content. In highly filled composites, part of silver nanoparticles is well dispersed (up to 49.5 wt%), the rest is poorly dispersed and is comparable to compressed AgNP (regions with less polymer matrix). For spherical filler, the percolation threshold occurs at high filler content, above of which the heterogeneity may occur. These two phenomena are coincident but not related in terms of thermal transport.

To adjust the fitting for composites above 49.5 wt%, we modelled their structure by a continuous phase with an effective heat capacity $C_p(w_c)$ in which silver nanoparticles with a specific heat capacity $C_{p,AgNP}$ are dispersed. The conductive composites specific heat capacity $C_p'(w)$ was expressed as the weighted average of $C_p(w_c)$ and $C_{p,AgNP}$ as in Eq (11):

$$C_p'(w) = \frac{(w - w_c)}{(1 - w_c)} C_{p,AgNP} + \frac{(1 - w)}{(1 - w_c)} C_p(w_c) \quad (11)$$

In our case, $C_p(w_c)$ is the measured specific heat capacity for $w_c = 49.5$ wt% and equals $0.58 \text{ J g}^{-1} \text{ K}^{-1}$. This mixture rule seems to accurately describe the specific heat capacity of conductive composites.

In Fig. 3(a), is plotted the evolution of the composites specific heat capacity with temperature from 17 °C to 48 °C. Open symbols represent electrically conductive composites and filled symbols

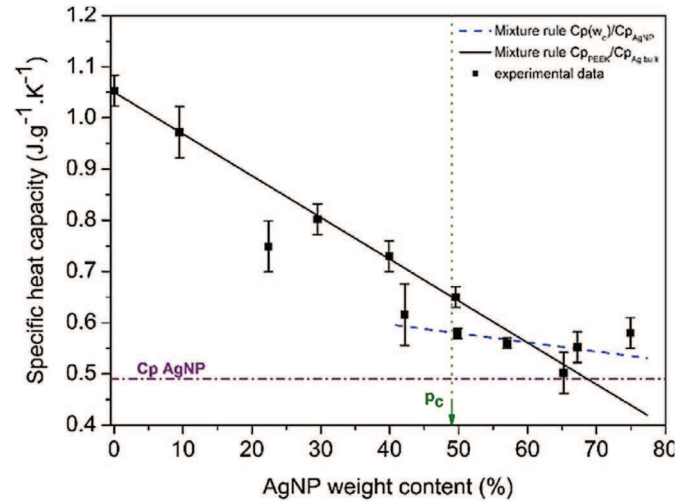


Fig. 2. Evolution of the specific heat capacity of the composites with silver nanoparticles weight content at 17 °C. Measured specific heat capacity (■), $C_{p,PEEK}$ (· · · ·), $C_{p,AgNP}$ (— · — ·), mixture rule $C_{p,PEEK}/C_{p,Ag bulk}$ (—), mixture rule $C_p(w_c)/C_{p,AgNP}$ (---). p_c is the electrical percolation threshold.

insulating ones. Fig. 3(b) shows the derivative of specific heat capacity with respect to temperature as a function of AgNP content. Fig. 3(a) and (b) underline that the heat capacity increases with temperature for every composites. But, the temperature dependence of the heat capacity tends to become less significant as the AgNP content increases. The specific heat capacity of PEEK increases significantly with temperature and the measured specific heat capacity of compressed AgNP barely increases over this temperature range, as represented by the solid line in Fig. 3(a). In crystalline materials, the heat capacity reaches a plateau above ambient temperature following the Dulong and Petit empirical behaviour [46]. In glasses, the temperature dependence of the heat capacity is more pronounced because of the greater molecular mobility [47]. In our case, as the AgNP content increases, the heat capacity increases more slowly with temperature. This behaviour can be attributed to the increasing proportion of metallic phase providing the overall material an ordered-like behaviour for the highly loaded composites. The temperature influence on the specific heat capacity highlights the contribution of the silver particles network to the composites specific heat capacity.

3.3. Thermal conductivity of the composites

We plotted on Fig. 4 the thermal conductivity of the composites as a function of silver nanoparticles volume content. We observe an increase with the introduction of silver nanoparticles up to 19.1 vol%. The measured thermal conductivity of PEEK is $0.26 \text{ W m}^{-1} \text{ K}^{-1}$. This value is consistent with the literature: Choy et al. found the PEEK thermal conductivity at $0.26 \text{ W m}^{-1} \text{ K}^{-1}$ measured by Laser Flash technique at 25 °C [48] and Diez-Pascual et al. at $0.22 \text{ W m}^{-1} \text{ K}^{-1}$ [29]. The maximum thermal conductivity is reached for 19.1 vol% AgNP and equals $0.43 \text{ W m}^{-1} \text{ K}^{-1}$ which represents a raise of 63%. Above 19.1 vol% AgNP, the thermal

Table 1
Specific heat capacity of PEEK at 17 °C, compressed AgNP and bulk silver at 20 °C measured by MT-DSC.

	Measured C_p ($\text{J g}^{-1} \text{ K}^{-1}$)	Standard deviation ($\text{J g}^{-1} \text{ K}^{-1}$)	C_p from literature ($\text{J g}^{-1} \text{ K}^{-1}$)
PEEK	1.05	0.03	1.09 [42]
Compressed AgNP	0.49	0.06	—
Bulk Ag	0.24	0.01	0.235 [32]

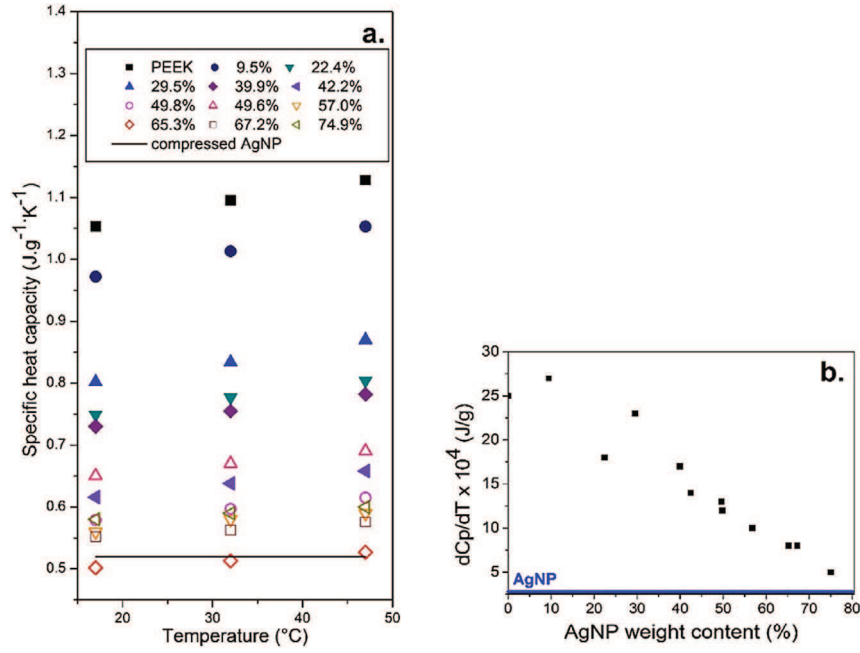


Fig. 3. (a) Evolution of the composites specific heat capacity with temperature. AgNP weight content: 0% (■), 9.5% (●), 22.4% (▼), 29.5% (▲), 39.9% (◆), 42.2% (◄), 49.6% (△), 49.8% (○), 57.0% (▽), 65.3% (◇), 67.2% (◻), 74.9% (◂); compressed AgNP is also represented (—). (b) Derivative of specific heat capacity with respect to temperature as a function of AgNP weight content. Experimental data for composites (■) and for compressed silver nanoparticles (—).

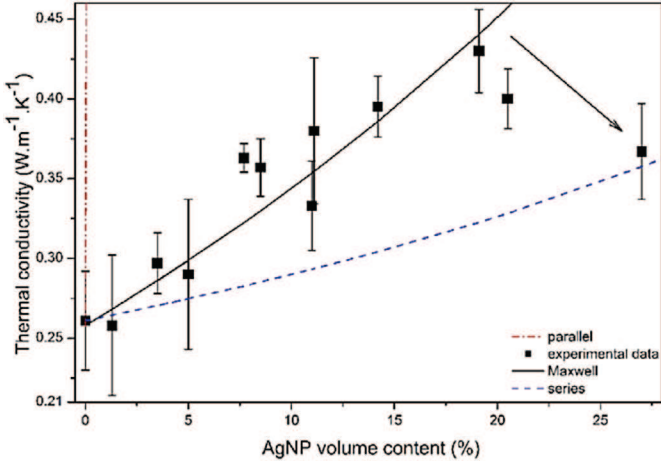


Fig. 4. Evolution of the thermal conductivity of the composites with silver nanoparticles volume content. Measured thermal conductivity at 17 °C (■), Maxwell model (—), Parallel model (---), Series model (···).

conductivity decreases. This fall is associated with an evolution of the dispersion quality. For high volume content, the dispersion process is less efficient and aggregates are observed. The polymer/particles interfaces are modified resulting in higher thermal resistances.

Most of the models used to predict thermal conductivity of two-phases composites are combinations of mixture rules taking into account phases relative content and morphologies. The series model describes thermal resistances in series configuration with respect to heat flow direction, and is represented by Eq (12):

$$\frac{1}{\lambda_{comp}} = \frac{x}{\lambda_{Ag}} + \frac{(1-x)}{\lambda_{PEEK}} \quad (12)$$

where λ_{comp} is the thermal conductivity of the composite, λ_{PEEK} is the thermal conductivity of PEEK, λ_{Ag} that of bulk silver and x is the volume fraction of AgNP [17].

The parallel model suggests co-continuous phases modelled by thermal resistances in parallel configuration with respect to heat flow direction [17], and can be expressed by Eq (13):

$$\lambda_{comp} = \lambda_{Ag}x + \lambda_{PEEK}(1-x) \quad (13)$$

These two fundamental models are generally in good agreement with experimental data for layered or anisotropic composites depending if the heat flows perpendicular or parallel to the interfaces, but tend to either minimize or maximize the contribution of the dispersed phase in the case of an isotropic composite [17]. The series and parallel models are the lower and upper bounds of mixture rules respectively. Our experimental data lie between these bounds. The deviation from the series model suggests a contribution of the silver nanoparticles to the thermal conductivity greater than a volume effect.

The Maxwell model is tailored for composites composed of a dispersed and a continuous phase, and gives the following expression for the thermal conductivity in the case of dispersed silver particles in PEEK matrix [49]:

$$\lambda_{comp} = \lambda_{PEEK} \frac{\lambda_{Ag} + 2\lambda_{PEEK} + 2x(\lambda_{Ag} - \lambda_{PEEK})}{\lambda_{Ag} + 2\lambda_{PEEK} + x(\lambda_{PEEK} - \lambda_{Ag})} \quad (14)$$

Our experimental data are in good correlation with the Maxwell model. It indicates that the polymer matrix is always continuous (except above 19.1 vol%), thus that particles distribution into the polymer matrix is satisfying. Nevertheless, there is no thermal manifestation of the electrical percolation, observed near 10.8 vol%. Similar results are observed in the literature. Mamunya et al. reported a monotonous increase of the thermal conductivity and no discontinuity in the electrical percolation threshold region in epoxy and PMMA matrixes filled with copper and nickel particles [4]. Huang et al. also observed a large increase of the thermal

conductivity of Polyvinylidene fluoride/AgNP composites, underestimated by the Effective Medium Approximation, while no electrical percolation was observed up to 20 vol% of AgNP [50]. The evolution of thermal conductivity as a function of AgNP content is not influenced by the electrical percolation pathway formation.

Fig. 5, represents the evolution of the composites thermal conductivity with temperature from 17 °C to 48 °C. Open symbols represent electrically conductive composites and filled symbols the insulating ones. We observe the same behaviour for both: the thermal conductivity increases slightly with the temperature. Classically, near room temperature, the temperature dependence of the thermal conductivity of a material is well known. For disordered solids and contrary to crystalline ones, the thermal conductivity increases with temperature [46,51,52]. This phenomenological behaviour is strongly associated with phonon mean free path. Our results suggest the predominant influence of the polymer matrix on the composite thermal conductivity dependence with temperature. In this temperature range, the thermal conductivity behaviour is similar to the heat capacity one.

4. Conclusion

The introduction of metallic nanoparticles into an insulating matrix enables an interesting optimization of electrical and thermal conductivities. The electrical conductivity drastically increases above a critical particle content $p_c = 10.8$ vol%, well described by a percolation law. The electrical characterization of the composites gives insights about connectivity and phases morphology. The specific heat capacity of the composites decreases with increasing

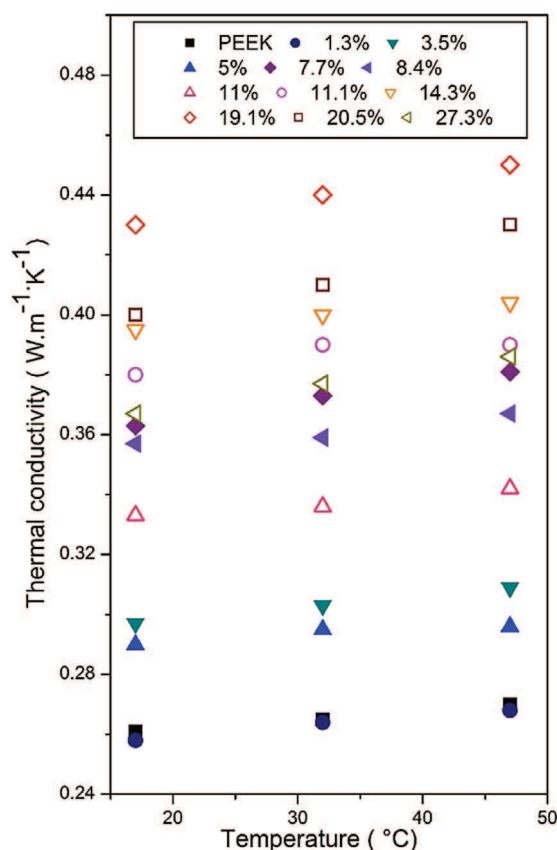


Fig. 5. Evolution of the composite thermal conductivity with temperature. AgNP volume content: 0% (■), 1.3% (●), 3.5% (▼), 5% (▲), 7.7% (◆), 8.4% (◀), 11% (△), 11.1% (○), 14.3% (▽), 19.1% (◇), 20.5% (□), 27.3% (◁).

particles content. This decrease is slower for conductive composites and their specific heat capacities converge toward the specific heat capacity of compressed silver nanoparticles. This behaviour has been attributed to the formation of heterogeneities above the electrical percolation threshold. The temperature dependence of the heat capacity decreases with the particles content. It indicates an evolution from glasses to an ordered-like behaviour for the highly loaded composites. Concerning the thermal conductivity, the discrepancy of our experimental data from the values predicted by the series model suggests a contribution of the silver particles to the effective thermal conductivity greater than a volume effect. A good agreement with the Maxwell model validates an isotropic distribution of silver particles into the polymer matrix up to 19.1 vol%. No manifestation of the electrical percolation phenomenon on the thermal conductivity is observed. The composites thermal conductivity increases with temperature as for glasses. The introduction of metallic particles does not induce any modification.

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References

- [1] S. Kirkpatrick, Percolation and conduction, *Rev. Mod. Phys.* 45 (1973) 574–588, <http://dx.doi.org/10.1103/RevModPhys.45.574>.
- [2] F. Lux, Models proposed to explain the electrical conductivity of mixtures made of conductive and insulating materials, *J. Mater. Sci.* 28 (1993) 285–301, <http://dx.doi.org/10.1007/BF00357799>.
- [3] C.W. Nan, Physics of inhomogeneous inorganic materials, *Prog. Mater. Sci.* 37 (1993) 1–116, [http://dx.doi.org/10.1016/0079-6425\(93\)90004-5](http://dx.doi.org/10.1016/0079-6425(93)90004-5).
- [4] Y.P. Mamunya, V.V. Davydenko, P. Pissis, E.V. Lebedev, Electrical and thermal conductivity of polymers filled with metal powders, *Eur. Polym. J.* 38 (2002) 1887–1897, [http://dx.doi.org/10.1016/S0014-3057\(02\)00064-2](http://dx.doi.org/10.1016/S0014-3057(02)00064-2).
- [5] G.W. Lee, M. Park, J. Kim, J.I. Lee, H.G. Yoon, Enhanced thermal conductivity of polymer composites filled with hybrid filler, *Compos Part A Appl. Sci. Manuf.* 37 (2006) 727–734, <http://dx.doi.org/10.1016/j.compositesa.2005.07.006>.
- [6] F.H. Gojny, M.H.G. Wichmann, B. Fiedler, I.A. Kinloch, W. Bauhofer, A.H. Windle, et al., Evaluation and identification of electrical and thermal conduction mechanisms in carbon nanotube/epoxy composites, *Polym. Guildf.* 47 (2006) 2036–2045, <http://dx.doi.org/10.1016/j.polymer.2006.01.029>.
- [7] Z. Han, A. Fina, Thermal conductivity of carbon nanotubes and their polymer nanocomposites: a review, *Prog. Polym. Sci.* 36 (2011) 914–944, <http://dx.doi.org/10.1016/j.progpolymsci.2010.11.004>.
- [8] P. Bonnet, D. Sireude, B. Garnier, O. Chauvet, Thermal properties and percolation in carbon nanotube-polymer composites, *Appl. Phys. Lett.* 91 (2007) 201910, <http://dx.doi.org/10.1063/1.2813625>.
- [9] R. Haggenueller, C. Guthy, J.R. Lukes, J.E. Fischer, K.I. Winey, Single wall carbon nanotube/polyethylene nanocomposites: thermal and electrical conductivity, *Macromolecules* 40 (2007) 2417–2421, <http://dx.doi.org/10.1021/ma0615046>.
- [10] Y.P. Mamunya, A. Boudenne, N. Lebovka, L. Ibos, Y. Candau, M. Lisunova, Electrical and thermophysical behaviour of PVC-MWCNT nanocomposites, *Compos. Sci. Technol.* 68 (2008) 1981–1988, <http://dx.doi.org/10.1016/j.compscitech.2007.11.014>.
- [11] K. Razeeb, E. Dalton, Nanowire-polymer nanocomposites as thermal interface material, in: Dr Reddy Borredy (Ed.), *Adv. Nanocomposites - Synth. Charact. Ind. Appl.*, Vol. 1, Rijeka: in Tech, 2011, pp. 685–706.
- [12] B. Weidenfeller, M. Höfer, F.R. Schilling, Thermal conductivity, thermal diffusivity, and specific heat capacity of particle filled polypropylene, *Compos Part A Appl. Sci. Manuf.* 35 (2004) 423–429, <http://dx.doi.org/10.1016/j.compositesa.2003.11.005>.
- [13] A. Lonjon, I. Caffrey, D. Carponcin, E. Dantras, C. Lacabanne, High electrically conductive composites of Polyamide 11 filled with silver nanowires: nanocomposites processing, mechanical and electrical analysis, *J. Non Cryst. Solids* 375 (2013) 199–204, <http://dx.doi.org/10.1016/j.jnoncrysol.2013.05.020>.
- [14] A. Lonjon, P. Demont, E. Dantras, C. Lacabanne, Low filled conductive P(VDF-TrFE) composites: influence of silver particles aspect ratio on percolation threshold from spheres to nanowires, *J. Non Cryst. Solids* 358 (2012) 3074–3078, <http://dx.doi.org/10.1016/j.jnoncrysol.2012.09.006>.
- [15] L. Quiroga Cortes, A. Lonjon, E. Dantras, C. Lacabanne, High-performance thermoplastic composites poly(ether ketone ketone)/silver nanowires: morphological, mechanical and electrical properties, *J. Non Cryst. Solids* 391 (2014) 106–111, <http://dx.doi.org/10.1016/j.jnoncrysol.2014.03.016>.
- [16] S. Torquato, Thermal conductivity of disordered heterogeneous media from

- the microstructure, *Rev. Chem. Eng.* 1987 (4) (1987) 151, <http://dx.doi.org/10.1515/REVCE.4.3-4.151>.
- [17] D.M. Bigg, Thermal conductivity of heterophase polymer compositions, *Adv. Polym. Sci.* 119 (1995) 1–30, <http://dx.doi.org/10.1007/BFb0021279>.
- [18] C.-W. Nan, R. Birringer, D.R. Clarke, H. Gleiter, Effective thermal conductivity of particulate composites with interfacial thermal resistance, *J. Appl. Phys.* 81 (1997) 6692, <http://dx.doi.org/10.1063/1.365209>.
- [19] M. Foygel, R. Morris, D. Anez, S. French, V. Sobolev, Theoretical and computational studies of carbon nanotube composites and suspensions: electrical and thermal conductivity, *Phys. Rev. B* 71 (2005) 104201, <http://dx.doi.org/10.1103/PhysRevB.71.104201>.
- [20] S. Kumar, M.A. Alam, J.Y. Murthy, Effect of percolation on thermal transport in nanotube composites, *Appl. Phys. Lett.* 90 (2007) 104105, <http://dx.doi.org/10.1063/1.2712428>.
- [21] S.Y. Kwon, I.M. Kwon, Y.G. Kim, S. Lee, Y.S. Seo, A large increase in the thermal conductivity of carbon nanotube/polymer composites produced by percolation phenomena, *Carbon N. Y.* 55 (2013) 285–290, <http://dx.doi.org/10.1016/j.carbon.2012.12.063>.
- [22] J. Koo, Y. Kang, C. Kleinstreuer, A nonlinear effective thermal conductivity model for carbon nanotube and nanofiber suspensions, *Nanotechnology* 19 (2008) 375705, <http://dx.doi.org/10.1088/0957-4484/19/37/375705>.
- [23] N. Shenogina, S. Shenogin, L. Xue, P. Keblinski, On the lack of thermal percolation in carbon nanotube composites, *Appl. Phys. Lett.* 87 (2005) 1–3, <http://dx.doi.org/10.1063/1.2056591>.
- [24] S.U.S. Choi, Z.G. Zhang, W. Yu, F.E. Lockwood, E.A. Grulke, Anomalous thermal conductivity enhancement in nanotube suspensions, *Appl. Phys. Lett.* 79 (2001) 2252–2254, <http://dx.doi.org/10.1063/1.1408272>.
- [25] A. Boudenne, *Etude expérimentale et théorique des propriétés thermophysiques de matériaux composites à matrice polymère*, PhD. Paris XII Val de Marne, 2003.
- [26] A. Boudenne, L. Ibos, M. Fois, J.C. Majesté, E. Géhin, Electrical and thermal behavior of polypropylene filled with copper particles, *Compos Part A Appl. Sci. Manuf.* 36 (2005) 1545–1554, <http://dx.doi.org/10.1016/j.compositesa.2005.02.005>.
- [27] M.Y. Razaq, M. Anhalt, L. Frommann, B. Weidenfeller, Thermal, electrical and magnetic studies of magnetite filled polyurethane shape memory polymers, *Mater. Sci. Eng. A* 444 (2007) 227–235, <http://dx.doi.org/10.1016/j.msea.2006.08.083>.
- [28] B. Weidenfeller, M. Höfer, F. Schilling, Thermal and electrical properties of magnetite filled polymers, *Compos Part A Appl. Sci. Manuf.* 33 (2002) 1041–1053, [http://dx.doi.org/10.1016/S1359-835X\(02\)00085-4](http://dx.doi.org/10.1016/S1359-835X(02)00085-4).
- [29] A.M. Díez-Pascual, M. Naffakh, C. Marco, G. Ellis, M.A. Gómez-Fatou, High-performance nanocomposites based on polyetherketones, *Prog. Mater. Sci.* 57 (2012) 1106–1190, <http://dx.doi.org/10.1016/j.pmatsci.2012.03.003>.
- [30] J.C. Seferis, Polyetheretherketone (PEEK): processing-structure and properties studies for a matrix in high performance composites, *Polym. Compos* 7 (1986) 158–169, <http://dx.doi.org/10.1002/pc.750070305>.
- [31] N.A. Pratten, Review the precise measurement of the density of small samples, *J. Mater. Sci.* 16 (1981) 1737–1747, <http://dx.doi.org/10.1007/BF00540619>.
- [32] D.R. Lide, *CRC Handbook of Chemistry and Physics*, CRC Press, 2003.
- [33] M. Reading, A. Luget, R. Wilson, Modulated differential scanning calorimetry, *Thermochim. Acta* 238 (1994) 295–307, [http://dx.doi.org/10.1016/S0040-6031\(94\)85215-4](http://dx.doi.org/10.1016/S0040-6031(94)85215-4).
- [34] A. Boller, Y. Jin, B. Wunderlich, Heat capacity measurement by modulated DSC at constant temperature, *J. Therm. Anal.* 42 (1994) 307–330, <http://dx.doi.org/10.1007/BF02548519>.
- [35] B. Wunderlich, *Thermal analysis of Polymeric Materials*, 2005, <http://dx.doi.org/10.1007/b137476>.
- [36] L.C. Thomas, *Measurement of Accurate Heat Capacity Values*, TA Instruments, Modul DSC Pap #9, 2005, pp. 1–11.
- [37] S.M. Marcus, R.L. Blaine, Thermal conductivity of polymers, glasses and ceramics by modulated DSC, *Thermochim. Acta* 243 (1994) 231–239, [http://dx.doi.org/10.1016/0040-6031\(94\)85058-5](http://dx.doi.org/10.1016/0040-6031(94)85058-5).
- [38] R.L. Blaine, S.M. Marcus, Derivation of temperature-modulated DSC thermal conductivity equations, *J. Therm. Anal. Calorim.* 54 (1998) 467–476, <http://dx.doi.org/10.1023/A:1010178022559>.
- [39] U. Gaur, B. Wunderlich, Heat capacity and other thermodynamic properties of linear macromolecules. V. Polystyrene, *J. Phys. Chem. Ref. Data* 11 (1982) 313, <http://dx.doi.org/10.1063/1.555663>.
- [40] H.F. Mark, *Encyclopedia of Polymer Science and Technology*, third ed., Wiley Interscience, 2003.
- [41] A.I. Medalia, Electrical conduction in carbon black composites, *Rubber Chem. Technol.* 59 (1986) 432–454, <http://dx.doi.org/10.5254/1.3538209>.
- [42] S.Z. Cheng, S. Lim, L.H. Judovits, B. Wunderlich, Heat capacities of high melting polymers containing phenylene groups, *Polym. Guildf.* (1987), [http://dx.doi.org/10.1016/0032-3861\(87\)90313-2](http://dx.doi.org/10.1016/0032-3861(87)90313-2).
- [43] A.I. Chumakov, G. Monaco, A. Fontana, A. Bosak, R.P. Hermann, D. Bessas, et al., Role of disorder in the thermodynamics and atomic dynamics of glasses, *Phys. Rev. Lett.* (2014), <http://dx.doi.org/10.1103/PhysRevLett.112.025502>, 112: 025502.
- [44] B. Budiansky, Thermal and thermoelastic properties of isotropic composites, *J. Compos. Mater.* 4 (1970) 286–295, <http://dx.doi.org/10.1177/002199837000400301>.
- [45] B. Agoudjil, A. Boudenne, I. Krupa, L. Ibos, J. Majeste, *Propriétés thermophysiques de composites polymère/charges isolantes métallisées*. Congrès Français Therm. SFT 3 (2007) 1–6. Les Embiez.
- [46] C. Kittel, *Introduction to Solid State Physics*, third ed., Wiley, John, New York, 1967.
- [47] B. Wunderlich, The heat capacity of polymers, *Thermochim. Acta* 300 (1997) 43–65, [http://dx.doi.org/10.1016/S0040-6031\(96\)03126-7](http://dx.doi.org/10.1016/S0040-6031(96)03126-7).
- [48] C.L. Choy, K.W. Kwok, W.P. Leung, F.P. Lau, Thermal conductivity of poly(ether ether ketone) and its short-fiber composites, *J. Polym. Sci. Part B Polym. Phys.* 32 (1994) 1389–1397, <http://dx.doi.org/10.1002/polb.1994.090320810>.
- [49] J. Wang, J.K. Carson, M.F. North, D.J. Cleland, A new structural model of effective thermal conductivity for heterogeneous materials with co-continuous phases, *Int. J. Heat. Mass Transf.* 51 (2008) 2389–2397, <http://dx.doi.org/10.1016/j.ijheatmasstransfer.2007.08.028>.
- [50] X. Huang, P. Jiang, L. Xie, Ferroelectric polymer/silver nanocomposites with high dielectric constant and high thermal conductivity, *Appl. Phys. Lett.* 95 (2009) 242901, <http://dx.doi.org/10.1063/1.3273368>.
- [51] C.L. Choy, Thermal conductivity of polymers, *Polym. Guildf.* 18 (1977) 984–1004, [http://dx.doi.org/10.1016/0032-3861\(77\)90002-7](http://dx.doi.org/10.1016/0032-3861(77)90002-7).
- [52] R.C. Zeller, R.O. Pohl, Thermal conductivity and specific heat of noncrystalline solids, *Phys. Rev. B* 4 (1971) 2029–2041, <http://dx.doi.org/10.1103/PhysRevB.4.2029>.