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Ion Beam Irradiation of ABO₄ Compounds with the Fergusonite, Monazite, Scheelite, and Zircon Structures

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Abstract

The effects of irradiation on CaWO₄, SrWO₄, BaWO₄, YVO₄, LaVO₄, YNbO₄, and LaNbO₄ were investigated on thin crystals using 1.0 MeV Kr ions at 50-1000 K. All of the ABO₄ compounds can be amorphized with cross sections ($\sigma_a = 1/F_{c_a}$) in the range of ~ 0.30-1.09 × 10⁻¹⁴ cm² ion⁻¹ at zero Kelvin. Analysis of fluence-temperature data returned critical temperatures for amorphization (T_c) of 311 ± 1 , 358 ± 90 , 325 ± 19 , 415 ± 17 , 541 ± 6 , 636 ± 26 , and 1012 ± 1 K, respectively for the compounds listed above. Compared with previous in situ irradiation of ABO4 orthophosphate samples using 0.8 MeV Kr ions, the T_c values of LaVO₄ and YVO₄ are higher than those of LaPO₄ and YPO₄ by 82 K and 124 K, respectively. The T_c values of the three scheelite structures, CaWO₄, SrWO₄, and BaWO₄, indicate that they are the most radiation tolerant compounds under these conditions. The A-B cation anti-site energies, $E_{f_{AB}}$, determined by DFT range from 2.48 to 10.58 eV and are highly correlated with the A-B cation ionic radius ratio, r_A/r_B , but are not correlated with T_c across the different structure types, suggesting that Frenkel defects play a more important role in damage recovery in these compounds. We also discuss the role of cation and anion charge/iconicity as determined by DFT. ABO₄ compounds with the zircon structure and B = P or V have a distinct advantage over those with B = Si as the damaged regions do not appear to be significantly affected by polymerization of $(PO_4)^{3-}$ or $(VO_4)^{3-}$ groups which might stabilize the amorphous fraction and ultimately lead to phase separation as observed in zircon (ZrSiO₄).

Keywords: Amorphization, ion irradiation, fergusonite, monazite, scheelite, zircon

I. Introduction

The development of radiation resistant materials has been an area of increasing interest over the years, particularly with the development of fusion and Generation IV nuclear reactors as potential contributors to sustainable future energy.¹ In all reactor systems, the existence of both elevated

temperatures and harsh irradiation environments are known to have a profound effect on core structural materials, causing the degradation and alteration of their physical and chemical properties.^{2,3} Similar problems also extend to candidate ceramic and glass-ceramic materials for nuclear waste immobilization, experiencing radiation damage effects caused by the alpha decay of encapsulated plutonium and minor actinides.^{4–6}

A great deal of effort has centered upon the mechanisms that govern radiation tolerance in materials, especially complex oxides. With the aid of ion beam irradiation techniques and validation from computer models, experimental investigations have focused primarily on understanding the cumulative decay processes of encapsulated radionuclides over time.^{7–9} In most cases, the degree of damage differs significantly as a function of composition and dose, with structural defects occurring due to varied phenomena that rely heavily upon the intricate balance between damage production, structural stability, ability to accommodate lattice disorder, and the energetics of defect formation and migration. In this regard, tailoring the response of materials to irradiation constitutes a major challenge in materials research; therefore, a fundamental understanding of a material's behavior under such conditions is necessary to justify its practical applicability.

In amongst a number of promising materials for nuclear waste immobilization, the ABO₄ compounds have gained significant interest, especially in the area of radiation tolerant crystalline host phases for actinides and fission products from spent nuclear fuel. The materials based on this general formula exhibit extensive chemical flexibility, with A = Ca, Sr, Ba, Y, lanthanides, Th, and U and B = Si, P, As, Nb, Ta, and W, among other elements. The most common ABO₄ compounds are based on the structures of fergusonite, monazite, scheelite, and zircon. These are also represented in nature as potential ores or minor carriers of Nb, Ta, W, Y, lanthanides, Th, and U in geological systems, enabling researchers to study the system evolution as a function of temperature, pressure, melt-fluid composition, and time.^{10,11} Radiation damage effects in these materials have been studied extensively, especially on zircon and monazite structured

orthophosphates and silicates using ion irradiation techniques that cater specifically to the determination of damage models and mechanisms under strictly controlled conditions including temperature, ion mass, energy and flux.^{12–14} However, very little is known in regards to the radiation response of other ABO₄ structure types, particularly those with B = V, Nb, W, and other elements apart from Si. In order to assess these materials for nuclear waste immobilization, this study aims at comparing the effects of ion beam irradiation on CaWO₄, SrWO₄, BaWO₄, LaVO₄, YVO₄, LaNbO₄ and YNbO₄, spanning the various chemistries of ABO₄ compounds with either a monoclinic or tetragonal structure. This method aids in the identification of characteristic trends between different ABO₄ compounds due to the variation of structure/composition with respect to the critical amorphization dose (D_c) and the critical temperature (T_c), above which the material remains crystalline.

II. Crystal Chemistry of ABO₄ Structure Types

This section is based in part on the review articles by Chakoumakous et al.¹⁵ and Dacheux et al.,¹⁶ together with a number of specific crystal structure refinements and other studies relevant to this work.¹⁷⁻²⁴ As shown in Figure 1, all four of the structure types in the system described above are composed of edge sharing 8 or 9 coordinated cation polyhedra (A-sites) and smaller isolated cation tetrahedra (B-sites) connected to the A-sites via either corner sharing (scheelite and fergusonite) or edge sharing (monazite and zircon). The zircon structure (I4₁/*amd*) is related to the monazite structure (P2₁/*n*) by a shift in the cation positions, rotation of anion positions about the tetrahedra, and the change in coordination of the A-site from 8 to 9. In zircon, there are two sets of four equal A-O bond lengths and four equal B-O bond lengths. However, the BO₄ tetrahedral sites are significantly distorted with two shorter shared O-O edges and four longer unshared O-O edges. Similarly, there are two O-B-O tetrahedral angles associated with the shared edges that are significantly smaller or greater than the ideal value of 109.5° and four associated with the unshared edges that are larger. In the monazite structure there are nine individual A-O

bond lengths and four individual B-O distances. Like zircon, the two shared O-O edges are shorter and have the smaller O-B-O angles; however, due to the particular distortion of the BO₄ tetrahedron in monazite there are now only three angles significantly greater than 109.5° and one of which is close to this angle.

The crystal structure of scheelite (space group $I4_1/a$) consists of a framework of edge sharing AO₈ polyhedra connected to isolated BO₄ tetrahedra via corner sharing. There are two groups of four A-O distances in scheelite. All four B-O distances are equal; however, tetrahedron is somewhat distorted as shown by two sets of three O-B-O angles that are significantly smaller or larger than the ideal tetrahedral bond angle. Scheelite is related to fergusonite (standard space group C2/c, the non-standard space groups I2/a or I2/c are often used for comparison with scheelite) via small cation displacements and more significant changes in the anion positions. Phase transformations in fergusonite occur as a function of temperature and pressure and are generally of a second order, ferroelastic type. In YNbO₄ and LaNbO₄, the temperature driven phase transformations to scheelite occur at approximately 1073 and 790 K, respectively. Fergusonite, the low temperature or high pressure phase, is effectively "distorted and compressed" version of scheelite with four distinct pairs of A-O bond distances in the AO₈ polyhedron. Two of the oxygen atoms are displaced toward the B-site in fergusonite, therefore, the local coordination environment consists of two pairs of B-O distances (~ 1.8 and 1.9 Å) in a distorted tetrahedral configuration with another pair of distances at about 2.5 Å.

For these structure types, the cation-anion coordination environments can be expressed as $^{VIII}A^{IV}B^{III}O_4$ for zircon and scheelite, $^{IX}A^{IV}B^{III}O_3^{IIV}O$ for monazite, and either $^{VIII}A^{IV}B^{III}O_4$ or $^{VIII}A^{VI}B^{III}O_2^{IIV}O_2$ for fergusonite depending upon whether or not the B-site is considered to be IV or VI coordinated. Although in the latter case the choice is somewhat subjective, the configurations illustrated above point to important differences in the local coordination environments between these compounds. These local atomic configurations, coupled with differences in the chemistry of the B-site cation in particular, may have a major role in the atomic

scale processes that affect the production and recovery of defects during irradiation, e.g., the energetics of defect formation and migration on picosecond time scales.

III. Experimental Methods

(1) Materials Synthesis

Polycrystalline ceramic samples of CaWO₄, YNbO₄, LaVO₄ and YVO₄ were prepared via conventional solid-state oxide methods using stoichiometric amounts of La₂O₃ (99.99%), Nb₂O₅ (99.99%), Y₂O₃ (99.99%), WO₃ (99%), V₂O₅ (99%) and CaCO₃ (99%). Starting oxides were first calcined at 600 °C for 10 hours and CaCO₃ was dried in an oven at 120 °C before the powders were weighed then milled using Y₂O₃-stabilized zirconia balls in Teflon containers in a cyclohexane media for 16 h. The cyclohexane was evaporated from stainless steel pans at 110 °C. The homogeneously mixed powder was then uniaxially pressed into pellets at 5 MPa, cold isostatically pressed (400 MPa for the tungstates and vanadates and 300 MPa for the niobates) and sintered for 16 hours at 1400 °C for the niobates and tungstates and at 1100 °C for the vanadates.

(2) Materials Characterization

X-ray diffraction patterns of were measured using a Bruker D8 instrument, with weighted CuK(radiation and a Sol-X detector. Samples were analyzed in the angular range of 10 - 80 degrees two theta, with a step size of 0.03° and a counting time of 10 s per step. A Zeiss Ultra Plus scanning electron microscope operating at 15 kV and equipped with an Oxford Instruments X-Max 80mm² SDD X-ray microanalysis system was used for microstructural and phase-composition analysis. Samples were mounted in an epoxy resin and polished to 1 µm diamond finish. A carbon film (~ 5 nm) was deposited onto the polished surface for charge neutralization.

(3) Ion Irradiation

Before irradiation, selected area diffraction patterns (SAED) were obtained using a JEOL 2010F transmission electron microscope (TEM) operated at an accelerating potential of 200 kV and calibrated for selected area diffraction over a range of objective lens currents using a gold film standard. TEM specimens were prepared by lightly grinding a small amount of powder dispersed in ethanol and drop casted by a plastic pipette onto holey carbon coated copper TEM grids and dried in air. Ion irradiation experiments with 1.0 MeV Kr ions were carried out in situ at the IVEM-Tandem User Facility at Argonne National Laboratory using a Hitachi TEM interfaced to a NEC ion accelerator.²⁵ All TEM observations were carried out using an accelerating potential of 300 kV. Ion irradiations were performed at temperature range of 50-1000 K with the electron beam of the TEM turned off, using a counting rate of 50 ion counts s⁻¹ and a flux of 6.25×10^{11} ions cm⁻² s^{-1} within a beam of ~ 2 mm diameter. Each sample was irradiated using incremental irradiation steps and selected grains were observed using bright field imaging and selected area diffraction after each irradiation step. The critical amorphization fluence (F_c) was determined from the last fluence increment in which weak Bragg diffraction spots were observed and the next increment for which only diffuse rings occur in the diffraction pattern. For each sample, F_c was determined from an average of 3 to 6 grains. A zirconolite reference sample with previously established F_c was also irradiated in order to validate the experiments.

(4) Monte Carlo Simulations of Damage Production

The irradiation damage to a TEM specimen by 1 MeV Kr ions was calculated using SRIM-2008 (Stopping and Range of Ions in Matter) with full damage cascades.^{26,27} Here, displacement energies of 50 eV were assumed for all atoms, coupled with surface binding energies of 3 eV in order to assess the effects of sample thickness and composition on defect production, stopping power, and energy transfer to atoms in the target. Table 2 summarizes the results of these simulations. For the ABO₄ compounds used in this study, the nuclear stopping power ranges from 1113 to 1342 eV ion⁻

¹ nm⁻¹, the electronic stopping power ranges from 773 to 979 eV ion⁻¹ nm⁻¹, and the electronic-tonuclear stopping power (ENSP) ratio ranges from 0.643 to 0.783.

(5) DFT Simulations

Static calculations were carried out on five of the ABO₄ compounds studied here: CaWO₄ (scheelite), LaVO₄ (monazite), YVO₄ (zircon), and YNbO₄ and LaNbO₄ (fergusonite). The atomic bonding was predicted using density functional theory (DFT) within the VASP code.^{28,29} The DFT approach was chosen to account for the variation in bonding expected to occur as the structure and composition of the ABO₄ system is varied. An energy cutoff value of 450 eV was used together with a $2 \times 2 \times 2$ k-point grid. The GGA-PBE exchange correlation was employed with the PAW pseudopotential. For each of the five compounds, representing the four different structure types noted above, the energy involved in the exchange of Ca-W, La-V, Y-V, Y-Nb, and La-Nb next-nearest-neighbor (NNN) pairs was determined. We also investigated some of the fundamentals of bonding in these compounds by calculating the A-B cation charge ratio and the effective charge on the oxygen atom. Results of these calculations are listed in Table 1.

IV. Results

(1) Structure and Composition of the Ceramics

Structural details of the nominal ABO₄ compounds are shown in Table 1, based on high-quality data from the literature.¹⁶⁻²⁴ The table also includes the A-B cation anti-site defect formation energies, $E_{f_{AB}}$, calculated in this study. Figure 3 shows the XRD patterns which were in excellent agreement with the literature data. The samples were identified by laboratory XRD as essentially single-phase materials with either monoclinic or tetragonal crystal structure. Trace amounts of monoclinic ZrO₂ in the YNbO₄ sample may be indicated by very weak reflections appearing close to 28 and 31° 2 θ , probably incorporated during ball-milling. An SEM image collected for each sample with the backscatter detector confirmed the single-phase nature of the samples and EDS

analyses verified the phase composition. There was no evidence for ZrO_2 in the SEM images collected with the backscatter detector.

(2) Ion irradiation

The radiation-induced crystalline-to-amorphous transition in all ABO₄ compounds was monitored through a sequence of SAED patterns, mapping the gradual decrease in Bragg intensity and complementary increase in diffuse scattering from amorphous domains under continuous irradiation, until complete amorphization was observed as diffuse halos (no Bragg beams present) above the critical amorphization fluence, F_c . This critical threshold value generally increases as a function of irradiation temperature due to the defect recombination and epitaxial recrystallization processes at elevated temperatures which suppress the amorphization process. At this point, a certain critical temperature (T_c) is reached where the transition to complete amorphization does not occur. A plot of this relationship for the seven compounds investigated in this study can be seen in Figure 4, conferred by the expression:

$$F_{C} = \frac{F_{c0}}{1 - \exp\left[\left(\frac{E_{a}}{k_{b}}\right)\left(\frac{1}{T_{c}} - \frac{1}{T}\right)\right]}$$

In this equation, F_c is the fluence for complete amorphization at 0 K, T_c the critical temperature at which full amorphization ceases, E_a the activation energy for crystalline recovery, k_b the Boltzmann constant and T the measured experimental temperature.^{30,31} Data from this refinement, including the extrapolated critical fluence at 0 K (D_{c_0}), critical temperature (T_c) and defect annealing activation energy (E_a), are summarized in Table 3.

Under the stated experimental conditions, none of the compounds studied here are inherently radiation tolerant, with calculated intercept fluences of $F_{c0} = 0.9-3.3 \times 10^{-14}$ ions cm⁻². This is equivalent to radiation damage cross-sections for amorphization $\sigma_a = 1/F_{c_0} = 0.3-1.1 \times 10^{14}$ cm² ion⁻¹. The low temperature behavior observed here is similar to that found for a range of complex oxides and the intercept fluences are just below values observed for the prototypical zirconolite and titanate pyrochlore nuclear waste form compositions. Of the seven ABO₄ compounds, the three scheelite samples became amorphous below 400 K, but a systematic trend between the three compounds could not be determined due in part to the large error in the T_c determination for SrWO₄ (see Table 3). In contrast to this result, LaNbO₄ was found to be the least radiation tolerant compound with a T_c of 1012 K, followed by YVO₄ (636 K), YNbO₄ (541 K), and LaVO₄ (415 K). During ion irradiation experiments at elevated temperature, LaNbO₄ should go through the monoclinic-tetragonal phase transformation at ~ 763 to 798 K,³² but with no apparent effect on the fluence-temperature response. Note also that we do not find any direct evidence for complex, multi-stage recovery of damage in these materials, as was reported in a previous study.¹³

As shown in Figure 5a-c, there are broad correlations between density, the electronic to nuclear stopping power ratio (ENSP) for these irradiations, and the A-B cation ionic radius ratio (r_A/r_B) . Density and r_A/r_B (Fig. 5a,c) are negatively correlated with T_c for these structure types as a group, with the fergusonite data showing a general reversal in behavior from La to Y with respect to the trend shown by the two monazite (LaVO₄) and zircon (YVO₄) samples. Figure 5b, on the other hand, shows that ENSP is positively correlated with T_c from the three scheelite compounds to the group at low to medium T_c including monazite, zircon, and fergusonite (YNbO₄), and then up to fergusonite (LaNbO₄) at the highest T_c value.

(4) DFT Calculations

The cation anti-site defect formation energies ($E_{f_{AB}}$) obtained in this study are listed in Table 1 and plotted in Figure 5d and Figure 7a. In general, the energies, $E_{f_{AB}}$ required for an A-B cation exchange event to occur range from about 2.4 to 10.6 eV for the five ABO₄ compounds investigated here and are generally uncorrelated with T_c across the range of structure types and compositions. Interestingly, the lowest $E_{f_{AB}}$ values were found for the fergusonite structure type: 2.48 eV for YNbO₄ and 4.04 eV for LaNbO₄. This range points to a possible systematic difference in cation anti-site behavior for the fergusonite structure as a function of composition. An intermediate value of $E_{f_{AB}} = 6.58$ eV was determined for scheelite CaWO₄. The highest cation antisite energies were found for the zircon and monazite structure types, YVO₄ and LaVO₄, respectively, with $E_{f_{AB}} = 8.25$ and 10.58 eV, respectively. Figure 5d shows that the $E_{f_{AB}}$ values determined for the cation anti-site energy are highly correlated with the cation radius ratio, r_A/r_B , for all samples as a group, given by the linear relationship $r_A/r_B = 0.134(E_{f_{AB}}) + 1.8132$ with $R^2 =$ 0.988.

We also calculated the A/B cation charge ratio ($q_{A/B}$) and the average charge on the oxygen atom (q_0) for each of the five different compounds (Table 1, Figure 7 b and c). Overall, the A/B charge ratios determined from DFT calculations decrease systematically from the samples with B = V (zircon or monazite) to Nb (fergusonite) to W (scheelite) for which the $q_{A/B}$ values are 1.07, 1.06, 0.84, 0.82, and 0.56, respectively. The ideal ionic values of $q_{A/B}$ are 0.60 based on the formal charges, A³⁺ and B⁵⁺, for the zircon, monazite, and fergusonite compounds and 0.33 based on the formal charges, A²⁺ and B⁶⁺, for the scheelite compound investigated here. In terms of the corresponding value of the oxygen charge, q_0 , across this series in the same order, we find that q_0 decreases from -1.09 and -1.07 for the samples with B = V (zircon or monazite) to -1.23 and -1.20 for B = Nb (fergusonite) and then increases again to -1.14 for B = W (scheelite). The ideal ionic charge for O is, of course, -2.00 in these compounds.

V. Discussion

Under the conditions used for these *in situ* experiments all of the irradiated materials investigated in this study are inherently susceptible to amorphization induced by the 1.0 MeV Kr ions. The intercept fluences given in Table 3, as derived from the fluence-temperature response curves, are somewhat higher than the range of $1.0-1.6 \times 10^{14}$ ions cm⁻² reported for the monazite and xenotime (zircon) orthorphosphates by Meldrum et al.¹² using similar *in situ* irradiations but with 800 keV Kr ions. In another study using similar experimental conditions, Meldrum et al.¹³ determined F_{c_0} values of 1.1-1.5 × 10¹⁴ ions cm⁻² for the orthosilicates with the zircon (ZrSiO₄, HfSiO₄, ThSiO₄) and monazite (ThSiO₄ = huttonite) structure. Furthermore, heavy ion irradiation of ZrSiO₄ single crystals with 540 keV Pb ions at 77 K shows that complete amorphization occurs at a fluence of approximately 10¹⁴ ions cm⁻².³³

The factors that influence damage recovery in this complex system of structures and compositions are poorly understood; however, the temperature response based on the T_c values follows the general pattern of II-VI scheelite < III-V monazite < III-V zircon < III-V fergusonite. Furthermore, in relation to monazite and zircon, the trend in radiation response of fergusonite is reversed, i.e., the compound with the largest A-site cation (La) is the least radiation tolerant. As a result of this reversal, the T_c value of YNbO₄ is significantly lower than that of YVO₄ but slightly higher than the value of 512 K reported for YPO₄ (zircon). Results of this study demonstrate that the response to external irradiation is complicated in the broader ABO₄ system. Our results demonstrate that the three scheelite samples are the most radiation tolerant compounds investigated in this study, with a T_c values below that of LaPO₄ monazite.¹² The low T_c value obtained for CaWO₄ in this study is supported by the work of Mendoza et al.³⁴ who conducted ion irradiations of ceramic samples of powellite (CaMoO₄ with the scheelite structure) doped with lanthanides and Na. Following ion irradiation with 8 MeV Ar ions to 1.2×10^{16} ions cm⁻² or 108 MeV Pb ions to 10¹⁴ ions cm⁻² (presumably at room temperature), the authors suggest that the crystalline structure of the powellite is retained. Based on evidence obtained from $Eu^{3+5}D_0 \square {}^7F_0$ luminescence spectra, it appears that the Ar irradiations only have a minor effect on the local environment of Eu. Although the Eu local environments were randomized by the Pb irradiations, the structure of the CaMoO₄ remained crystalline.

In a study designed to simulate the effects of alpha decay damage, Picot et al.³⁵ conducted irradiation experiments on LaPO₄ and La_{0.73}Ce_{0.27}PO₄ monazite ceramics using 1.7 MeV He ions

to 1.2×10^{16} ions cm⁻² and 1.35 and 7 MeV Au ions to 2.3×10^{15} ions cm⁻². Results of the study demonstrate that (La,Ce)PO₄ monazite can be rendered amorphous by irradiation with Au ions under these conditions and also show that it undergoes a volume expansion of ~ 8%. As expected, irradiation with He ions does not lead to amorphization of monazite, but the volume does increase by ~ 0.8%, probably due to the accumulation of point defects in the structure. Natural monazites, however, are always found in the crystalline state even with high Th-U contents (over 20 wt.% in some cases) and very old geological ages. However, it is also known that synthetic samples doped with 8.1 wt% ²³⁸Pu also remain crystalline on the time scale of the experiments (several years), but show evidence of damage accumulation and retention in the form of decreased X-ray diffraction peaks relative to the starting material (Burakov et al.³⁶). Therefore, the barriers to damage recovery are accessible and may be overcome in natural systems where monazite samples are stored for geological time periods (lower dose rate) at somewhat elevated temperatures, averaging ~ 100-200

C in numerous cases.^{37,38}

The recovery of damage as a function of temperature in the ABO₄ orthovanadates (A = La, Y) with the monazite and zircon structure can be compared directly with similar experiments reported previously for phosphate and silicate systems. Firstly, it is instructive to compare our results for LaVO₄ (monazite) and YVO₄ (xenotime) with the orthophosphates studied by Meldrum and coworkers.¹² Our temperature-fluence data give $T_c = 410$ K for LaVO₄ and 644 K for YVO₄ whereas LaPO₄ and YPO₄ gave $T_c = 333$ K for LaPO₄ and 512 K for YPO₄ using a slightly lower ion energy of 800 keV.¹² These results demonstrate that the vanadates studied here are probably less radiation tolerant (with respect to T_c) than the corresponding phosphates based on the critical temperature values. This is consistent with the trends observed for the thermochemistry of orthophosphate and orthovanadate compounds with the zircon and monazite structures, e.g., the enthalpies of formation reported by Ushakov et al.³⁹ and Dorogova et al.⁴⁰ show that the vanadates

are consistently less stable than the phosphates with respect to the oxides by $\sim 190 \text{ kJ mol}^{-1}$ with the light lanthanide compounds being more stable in both systems.

In the experiments conducted via *in situ* ion irradiation, the orthosilicates with the zircon (ZrSiO₄, HfSiO₄, ThSiO₄) or monazite structure type (ThSiO₄) have high T_c values on the order of 900-1200 K.¹³ This is a clear illustration of the major effect of chemistry in the IV-IV versus III-V monazite and zircon structure types. In the orthosilicates, Si probably plays a major role in the stabilization of defects via polymerization effects, forming local -Si-O-Si- units, Q^n , with $n \le 4$. A major outcome of the polymerization of Si is local segregation of Si and A-site cations in collision cascades (e.g., Farnan and Salje⁴¹). Little work has been published on radiation effects in IV-IV ABO₄ compounds with other B-site cations; however, recent results have been reported for 93 MeV Ni irradiations conducted on ceramic samples of ThGeO₄ in both the zircon and scheelite structure types (Patel et al.⁴²). At room temperature, these "swift heavy ion" irradiation experiments caused amorphization of the zircon and scheelite phases at fluences of 6×10^{13} and 0.75×10^{13} ions cm⁻², respectively. These data indicate that the radiation tolerance of the zircon structure type is greater than that of the scheelite type for the ThGeO₄ polymorphs.

Fergusonite compounds have not received much attention with regard to experimental ion irradiation damage effects, but most natural samples are found in the metamict (amorphous) state, indicating that this structure type is generally susceptible to amorphization (see Lumpkin and Geisler⁴³). The results of this study reveal that fergusonite is the least resistant structure type to amorphization induced by 1.0 MeV Kr ions. Furthermore, the trend in radiation tolerance with respect to A-site cation radius is reversed relative to the monazite and zircon compounds with B = P and V. Based on a simple calculation using the differences in A-site cation radii and measured T_c values, we estimate that the T_c of LuNbO₄ is ~ 398 K. Thus, the range of critical temperatures for the fergusonite type niobates may span a temperature range of ~ 650 K, considerably greater than the range of ~ 250 K for the monazite and zircon orthophosphates combined.¹²

From the above discussion, the effects of 1.0 MeV Kr ion irradiation on thin crystals of the ABO₄ compounds investigated in this study are very complicated. In spite of the complexity of the results, we observe a distinct positive correlation between ENSP and T_c for CaWO₄, YNbO₄, and LaNbO₄, the three compounds with corner sharing tetrahedra. Together with the T_c values, this suggests that the electronic component of the stopping power assists in the production of defects in YNbO₄ and LaNbO₄ in particular, and that energy barriers for migration may be generally higher in the fergusonite structure type. However, embedded within the data set are the values for the monazite and zircon type orthovanadates, LaVO₄ and YVO₄, which have calculated ENSP values of 0.77 and 0.71. Although this trend is reversed in comparison to the other three samples of this study, it is consistent with the data presented in previous work for the orthophosphates¹² wherein ENSP is negatively correlated with T_c in these compounds, an indication that electronic interactions between the Kr ions and the atoms of the crystal may assist in defect recombination during the irradiations.

Provided that we separate the data for the two fergusonite samples from the others, results of this study show that density and r_A/r_B are negatively correlated with the critical temperature for amorphization (see Figure 5). In general, these observations are consistent with previous work on the TiO₂ polymorphs, for example, where the trend in radiation tolerance is highly correlated with density, molar volume, and packing index.⁴⁴ However, the situation is much more complex than this, e.g., for the monazite and zircon type orthophospates, ENSP is actually negatively correlated with T_c within each structure type. This is illustrated in detail in Figure 6, where we have plotted literature data for the phosphates¹² showing the variation of ENSP and r_A/r_B against T_c , clearly emphasizing the trends and, in particular, the steeper negative slopes for the zircon type orthophosphates relative to the monazites.

Finally, we have provided a view of the role of cation anti-site defect formation energies, cation charge ratio, and oxygen charge in the radiation damage of these structure types at the atomic scale using DFT calculations. Based on the values of $E_{f_{AB}}$, it is unlikely that A-B cation anti-site

disorder occurs to the extent that the materials transform to a disordered state, even in the case of YNbO₄ fergusonite with $E_{f_{AB}} = 2.48$ eV. However, anti-site disorder is not the only important defect mechanism and the energetics of cation and anion Frenkel defects may play a role radiation damage recovery. In a study of CaWO₄, Shao et al.⁴⁵ used empirical potentials to determine Frenkel defect energies of 1.23 eV for O and 4.04 eV for Ca. The Frenkel pair defect energy for a W vacancy and interstitial were not reported; however, data were presented suggesting that O atoms and vacancies have low migration energies of ~ 0.2 eV along 4 of the six WO₄ tetrahedral edges in the structure. This is generally consistent with experimental data and other atomistic studies of scheelite and related structure types^{46,47} and may play a role in the determining the radiation damage recovery as a function of temperature. In comparison with the above results, Yi et al.⁴⁸ used DFT methods to calculate Frenkel defect formation energies of 12.4 eV and 11.0 eV for Ce and O, respectively, in CePO₄ with the monazite structure. These results indicate that there are significant differences in the energetics of defect formation and potential migration mechanisms between the scheelite and monazite structure types relevant to recovery processes during irradiation. Additionally, it was recently⁴⁹ demonstrated that the threshold displacement energy of O in LaPO₄ monazite is only 8 eV, much lower than the values obtained for La (56 eV) and P (75 eV). This raises interesting questions about the behavior of the cation and anion sublattices across the different ABO₄ structure types during irradiation and requires further experimental and atomistic simulation work.

VI. Summary and Conclusions

None of the ABO₄ compounds studied here are inherently radiation tolerant with respect to *in situ* 1 MeV Kr irradiation, returning damage cross-sections for amorphization, $\sigma_a = 1/F_{c_0} = 0.3-1.1 \times 14 \text{ cm}^2$ ion⁻¹, on extrapolation to zero Kelvin. However, there are significant differences in the temperature response of the compounds, following the order II-VI scheelite < III-V monazite < III-V zircon < III-V fergusonite based on the critical temperature for amorphization (*T_c*).

Previous work on the zircon structure silicates ($ZrSiO_4$, $ThSiO_4$, and $USiO_4$), all with T_c values up around 1000 K under similar irradiation conditions, demonstrates that radiation damage recovery and post-recovery processes, e.g., due to annealing and/or long-term storage, are strongly influenced by Si-Si polymerization. This does not appear to be an important factor in this case, where the B-site tetrahedral are occupied by V, Nb, or W.

The wide range of observed T_c values in these ABO₄ compounds and structure types display a broad positive correlation with the electronic to nuclear stopping power ratio (ENSP) and negative correlations with density and r_A/r_B . In general, this indicates that a combination of thermal and irradiation assisted recovery processes occur during the irradiation of the thin TEM specimens.

The A-B cation anti-site energies, $E_{f_{AB}}$, determined by DFT range from 2.48 to 10.58 eV and are positively correlated with the A-B cation ionic radius ratio, r_A/r_B , but are not correlated with T_c across the different structure types, consistent with the observation that none of the materials disorder during irradiation; nor do they have stable disordered structures as a function of temperature, pressure, and composition within the range of experimental conditions use here. These observations suggest that Frenkel defects play a more important role in damage recovery in these compounds, consistent with a limited set of published data. However, the lowest $E_{f_{AB}}$ values are found for the fergusonite structure and have a positive correlation with T_c over a range of 471 K, in which case a significant role in damage recovery cannot be ruled out.

Tables

Table 1. Literature data for the crystal structures, mean bond lengths for A-O and B-O

coordination polyhedra, packing index (PI), and A-B cation ionic radius ratio (r_A/r_B), for the compounds examined in this study. Included here are our DFT results for the A-B cation anti-site defect formation energy $E_{f_{AB}}$, cation charge ratio $q_{A/B}$, and oxygen charge q_0 , as determined for a subset of five compounds.

	YVO ₄	LaVO ₄	YNbO4	LaNbO ₄	CaWO ₄	SrWO ₄	BaWO ₄
Structure	Ζ	М	F	F	S	S	S
Space group	I4 ₁ /amd	$P2_1/n$	I2/a	I2/a	$I4_1/a$	$I4_1/a$	$I4_1/a$
<i>a</i> (Å)	7.1183	7.047	5.317	5.565	5.243	5.4268	5.6065
b (Å)	7.1183	7.286	10.999	11.519	5.243	5.4268	5.6065
<i>c</i> (Å)	6.2893	6.725	5.090	5.202	11.376	11.9688	12.7084
β (°)	90.00	104.85	94.53	94.10	90.00	90.00	90.00
$V(\text{\AA})^3$	318.68	333.76	296.74	332.61	312.72	352.48	399.46
(Å-O) (Å)	2.3654	2.5965	2.363	2.505	2.4585	2.5405	2.724
⟨ B-O ⟩ (Å)	1.7088	1.7090	1.894	1.874	1.7854	1.869	1.822
$ ho_{ m calc}~({ m g~cm^{-3}})$	4.25	5.05	5.58	5.91	6.12	6.32	6.40
PI	58.70	60.33	63.41	59.11	61.84	57.69	54.52
r _A /r _B	2.870	3.268	2.123	2.417	2.667	3.000	3.714
$E_{f_{AB}}$ (eV)	8.25	10.58	2.48	4.04	6.58		
q _{A/B}	1.07	1.06	0.84	0.82	0.56		
qo	-1.09	-1.07	-1.23	-1.20	-1.14		

Table 2. Results of SRIM simulations of damage production, including the electronic $(dE/dx)_e$ and nuclear $(dE/dx)_n$ components of stopping power, and electronic to nuclearstopping power ratio (ENSP).

Sample	(dE/dx) _e eV nm ⁻¹	(dE/dx) _n eV nm ⁻¹	ENSP Ratio
YVO ₄	795	1113	0.714
LaVO ₄	872	1126	0.774
YNbO ₄	979	1342	0.730
LaNbO ₄	976	1246	0.783
CaWO ₄	801	1245	0.643
SrWO ₄	786	1201	0.654
BaWO ₄	773	1116	0.693

Table 3. Critical amorphization fluence (F_{c_0}) , displacements per atom (dpa), and radiation damage cross sections (σ_a) , all at zero Kelvin; together with the calculated defect annealing activation energies (E_a) and critical temperatures for amorphization (T_c) of seven ABO₄ compounds irradiated with 1.0 MeV Kr ions as a function of temperature.

Sample	$F_{c_{\theta}}(10^{14} \mathrm{cm}^{-2})$	dpa	$\sigma_a(10^{-14}\mathrm{cm}^2)$	<i>E</i> _a (eV)	$T_c(\mathbf{K})$
YVO ₄	1.63	0.13	0.61	1.95	636 (26)
LaVO ₄	1.30	0.22	0.77	1.26	415 (17)
YNbO4	3.30	0.52	0.30	1.69	541 (6)
LaNbO ₄	1.57	0.29	0.64	3.09	1012 (1)
CaWO ₄	2.04	0.36	0.49	0.96	311 (1)
SrWO ₄	1.00	0.19	1.00	1.10	358 (90)
BaWO ₄	0.92	0.20	1.09	0.98	325 (19)

Figures



Figure 1. Crystal structures of scheelite, monazite, zircon, and fergusonite.



Figure 2. Schematic drawing showing the relationships between scheelite, monazite, zircon, and fergusonite in terms of temperature, pressure, and ionic radii based on the nominal compositions of the four structure types. The typical valence states of ions on the A- and B-sites are also shown. Laboratory synthesis of fergusonite typically involves A^{3+} and B^{5+} cations; light blue numbers in parentheses illustrate the other cation valence states found in natural samples.



Figure 3. X-ray diffraction patterns of (a) CaWO₄, (b) YVO₄, (c) LaVO₄ and (d) YNbO₄. Trace monoclinic ZrO₂ is indicated by the star in (d). Otherwise, the ceramics are essentially phase pure at the level of detection typical for laboratory X-ray diffraction.



Figure 4. A plot of the critical dose for amorphization (D_c) as a function of temperature for the ABO₄ compounds using *in-situ* experimental techniques and thin TEM samples. Extrapolation of these curves to infinity gives the critical temperature for amorphization. a) A schematic plot summarizing the results obtained by Meldrum et al.^{12,13} for phosphates with monazite and zircon structures irradiated with 0.8 MeV Kr ions. Also shown for comparison (dashed line) is the general trend for zircon, thorite, coffinite, and huttonite silicate structures. b) New data obtained in this study for seven vanadate, niobate, and tungstate compounds (see legend) irradiated using 1.0 MeV Kr ions.



Figure 5. Plots of the critical temperature for amorphization (T_c) against the density in (a) and the electronic to nuclear stopping power (ENSP) in (b). The purple arrows show the possible trends of fergusonite compounds from YNbO₄ to LaNbO₄, the red arrows show the possible trends from LaVO₄ (monazite) to YVO₄ (zircon), and the circle and oval enclose the data for the scheelite compounds with A = Ca, Sr, and Ba. (c) A plot of the A-B cation radius ratio (r_A/r_B)) versus the critical temperature for amorphization (T_c). Arrows and oval as per figures (a) and (b) above. In (d) the relationship between the A-B cation anti-site defect formation energy ($E_{f_{AB}}$) and the A-B cation ionic radius ratio (r_A/r_B) is plotted for five of the seven compounds, revealing a close relationship between these parameters in these compounds.



Figure 6. Plots of ENSP (a,c) and r_A/r_B (b,d) versus the critical temperature for amorphization (T_c) for the monazite structure (a,b) and zircon (c,d) structure phosphates reported by Meldrum et al.¹² using 0.8 MeV Kr ions using the IVEM Tandem Facility. These figures reveal the relationships between T_c and these experimental and structural parameters in greater detail than previously reported, indicating that there are differences in the behavior of the two structure types. The monazite phosphates in (a,b) range from A = La to Gd (highest T_c) and the zircon phosphates in (c,d) range from A = Sc, Y, and Tb to Lu (highest T_c).



Figure 7. Plots obtained from experimental and DFT data showing the trends of the cation antisite energy $E_{f_{AB}}$ (a), A-B cation charge ratio $q_{A/B}$ (b), charge on the oxygen atoms q_0 (c), and for reference, the unit mass (d) versus the critical temperature for amorphization for five of the compounds investigated in this study.

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