# On a Parabolic Strongly Nonlinear Problem on Manifolds 

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Dedicated to Professor Marcondes Clark on the occasion of his 54th birthday


#### Abstract

In this work we will prove the existence uniqueness and asymptotic behavior of weak solutions for the system $\left(^{*}\right)$ involving the pseudo Laplacian operator and the condition $\frac{\partial u}{\partial t}+\sum_{i=1}^{n}\left|\frac{\partial u}{\partial x_{i}}\right|^{p-2} \frac{\partial u}{\partial x_{i}} \nu_{i}+|u|^{\rho} u=f$ on $\Sigma_{1}$, where $\Sigma_{1}$ is part of the lateral boundary of the cylinder $Q=\Omega \times(0, T)$ and $f$ is a given function defined on $\Sigma_{1}$.


Key Words: Weak solutions, manifolds, monotonic operator, compactness.
Mathematics Subject Classications: 35L85, 35L05, 35L20, 35L70, 49 A29

[^0]
## 1 Introduction

Let $\Omega$ be a bounded open set of $\mathbb{R}^{n}(n \geq 1)$ with smooth boundary $\Gamma$. We consider $\left\{\Gamma_{0}, \Gamma_{1}\right\}$ a partition of $\Gamma$, that is, $\Gamma=\Gamma_{0} \cup \Gamma_{1}$, with $\Gamma_{0}$ and $\Gamma_{1}$ having positive Lebesgue measure and with $\overline{\Gamma_{0}} \cap \overline{\Gamma_{1}}=\emptyset$.

Let $\nu$ be the outward normal to $\Gamma$ and $T>0$ is a real number. We denote by $Q=\Omega \times(0, T)$ the cylinder of the $\mathbb{R}^{n+1}$.

The goal of this work is to solve following strongly nonlinear boundary problem:

$$
(*) \left\lvert\, \begin{array}{rrr}
\mathcal{A} u=0 & \text { in } & Q=\Omega \times(0, T) \\
u=0 & \text { on } & \Sigma_{0}=\Gamma_{0} \times(0, T) \\
\frac{\partial u}{\partial t}+\sum_{i=1}^{n}\left|\frac{\partial u}{\partial x_{i}}\right|^{p-2} \frac{\partial u}{\partial x_{i}} \nu_{i}+|u|^{\rho} u=f & \text { on } & \Sigma_{1}=\Gamma_{1} \times(0, T) \\
u(x, 0)=u_{0}(x) & & \text { on } \Gamma,
\end{array}\right.
$$

where $\mathcal{A}$ is the pseudo Laplacian operator defined by

$$
\begin{aligned}
\mathcal{A}: W_{0}^{1, p}(\Omega) & \rightarrow W^{-1, p^{\prime}}(\Omega), \quad \frac{1}{p}+\frac{1}{p^{\prime}}=1 \\
w & \mapsto \mathcal{A} w
\end{aligned}
$$

with $\mathcal{A} w=-\sum_{i=1}^{n} \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial w}{\partial x_{i}}\right|^{p-2} \frac{\partial w}{\partial x_{i}}\right), \quad 2<p<\infty ; \rho$ is a positive real constant satisfying the conditions (1) and $f$ is a known real value function.

As the solution of system depend of $x$ and $t$ and the equation $(*)_{1}$ does not have temporal derivative of the function $u$, this system is not CauchyKovalevsky type.

This problem associated with evolution equation on lateral boundary, with $p=2$, was study in Araruna-Antunes-Medeiros [1] and Domingos-Cavalcante [4], both motivated by the idea applied in Lions ([6], pp. 134), which consists to reduce the problem in a model of mathematical physics on the manifolds $\Sigma_{1}$. Also, Araruna-Araujo in [2] studied the system (*) in your form more simple, that is, $p=2$. Recently, O.A.Lima, at al has been researching in

Partial Differential Equations involving the pseudo Laplacian operator [10]. In this work we use a technique due to Lions [6], which transforms the system $(*)$ in a Cauchy-Kovalevsky type one by means of a suitable perturbations in the equation $(*)_{1}$. The solution of $(*)$ is obtained as limit of solutions of the perturbed problem.

For $p>2$, the operator $\mathcal{A}$ brings great difficulties, because it is non-linear, mainly to establish concepts of solutions, in passage to the limit, to work with the trace application and immersion in spaces $W^{s, p}(\Omega), s \in \mathbb{R}$ (for this we consult Něcas [5]) and to obtain a estimative for derived of the approximate function(here we use strongly the proprieties of the trace application). Finally all the difficulties will be overcome through careful handling of the proprieties of the operator $\mathcal{A}$.

This paper is organized as follows: In Section 1, we will give some notations, hypothesis and results. In Section 2, we will introduce the perturbed problem. In Section 3, we will prove the existence of the solutions for the perturbed problem. In Section 4, we will treat of the uniqueness for the solution of the perturbed problem. Finally in Section 5, we will prove the main result of this work.

## 2 Hypotheses and Notations

We denote by $W^{\frac{1}{p^{\prime}, p}}(\Gamma)$ the vectorial space of functions $\left.v\right|_{\Gamma}$ when $v \in W^{1, p}(\Omega)$, for $\frac{1}{p}+\frac{1}{p^{\prime}}=1$. By $W^{-\frac{1}{p^{\prime}, p^{\prime}}}(\Gamma)$ denotes the dual of the space $W^{\frac{1}{p^{\prime}, p}}(\Gamma)$. Let $p>2$ be and $V_{0}$ the Banach space given by $V_{0}=\left\{v \in W^{1, p}(\Omega) ;\left.v\right|_{\Gamma_{0}}=0\right\}$ equipped with the norm $\|v\|_{V_{0}}=\left(\sum_{i=1}^{n} \int_{\Omega}\left|\frac{\partial v}{\partial x_{i}}\right|^{p} d x\right)^{\frac{1}{p}}$. Note that the application $\gamma: V_{0} \longrightarrow W^{\frac{1}{p^{\prime}}, p}\left(\Gamma_{1}\right)$ is linear, continuous and surjective.

In (*) assume that $\rho$ is a real number such that

$$
\begin{equation*}
\rho>0, \quad \text { if } n=1 \quad \text { or } \quad 0<\rho \leq \frac{(n+2) p-2(n+1)}{n-p+1} \text { se } n \geq 2 \text {. } \tag{1}
\end{equation*}
$$

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With the choice we have $W^{\frac{1}{p}, p}\left(\Gamma_{1}\right) \subset L^{\rho+2}\left(\Gamma_{1}\right)$ with continuous and dense immersion, for $\frac{1}{p^{\prime}}+\frac{1}{p}=1$. Therefore $L^{\frac{\rho+2}{\rho+1}}\left(\Gamma_{1}\right)=\left(L^{\rho+2}\left(\Gamma_{1}\right)\right)^{\prime} \subset W^{-\frac{1}{p^{\prime}, p^{\prime}}}\left(\Gamma_{1}\right)$ with immersions are dense and continuous.

Let $a: V_{0} \times V_{0} \rightarrow \mathbb{R}$ defined by

$$
\begin{equation*}
a(u, v)=\sum_{i=1}^{n} \int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{p-2} \frac{\partial u}{\partial x_{i}} \frac{\partial v}{\partial x_{i}} d x . \tag{2}
\end{equation*}
$$

which is linear with respect the second variable. Note that, the application $v \rightarrow a(u, v)$ is continuous on $V_{0}$ for $u \in V_{0}$ fix.

Let $V=L^{p}\left(0, T ; V_{0}\right)$ and the operator $\mathcal{B}$ from $V$ in $V^{\prime}$ defined by

$$
\begin{equation*}
(\mathcal{B}(u), v)_{V^{\prime} \times V}=\int_{0}^{T} a(u(t), v(t)) d t, \quad \forall u, v \in V \tag{3}
\end{equation*}
$$

Thus $\mathcal{B}$ is a hemicontinuous, monotonic operator and $\|\mathcal{B} u\|_{V^{\prime}} \leq C\|u\|_{V}^{p-1}, \quad \forall u \in$ $V$.

To facilitate the understand of this work, introduce the followings notations:

$$
(f, g)_{\Omega}=\int_{\Omega} f g d x, \quad(\varphi, \psi)_{\Gamma}=\int_{\Gamma} \varphi \psi d \Gamma
$$

## 3 Perturbed Problem

The Problem (*) is not of the Cauchy-Kowaleska's type. Thus, consider the following perturbed problem:

For all $\varepsilon>0$, the family of functions $u_{\varepsilon}(x, t)$ is defined by:

$$
(* *) \left\lvert\, \begin{array}{ccc}
\varepsilon \frac{\partial u_{\varepsilon}}{\partial t}+\mathcal{A} u_{\varepsilon}=0 & \text { in } & Q=\Omega \times(0, T), \\
u_{\varepsilon}=0 & \text { on } & \Sigma_{0}=\Gamma_{0} \times(0, T), \\
\frac{\partial u_{\varepsilon}}{\partial t}+\sum_{i=1}^{n}\left|\frac{\partial u_{\varepsilon}}{\partial x_{i}}\right|^{p-2} \frac{\partial u_{\varepsilon}}{\partial x_{i}} \nu_{i}+\left|u_{\varepsilon}\right|^{\rho} u_{\varepsilon}=f & \text { on } & \Sigma_{1}=\Gamma_{1} \times(0, T), \\
u_{\varepsilon}(x, 0)=w_{0}(x) & x \in \Omega,
\end{array}\right.
$$

where $w_{0}=\gamma^{-1} u_{0} \in V_{0}$.
The solution concept for $(* *)$ is established by Gauss's Theorem as follows:

For $v \in V_{0} \cap C^{2}(\bar{\Omega})$ we have

$$
\int_{\Omega} \frac{\partial}{\partial x_{i}}\left\{\left(\left|\frac{\partial u_{\varepsilon}}{\partial x_{i}}\right|^{p-2} \frac{\partial u_{\varepsilon}}{\partial x_{i}}\right) \cdot v\right\} d x=\int_{\Gamma_{1}}\left|\frac{\partial u_{\varepsilon}}{\partial x_{i}}\right|^{p-2} \frac{\partial u_{\varepsilon}}{\partial x_{i}} \cdot \nu_{i} \cdot v d \Gamma,
$$

hence
$\int_{\Omega} \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u_{\varepsilon}}{\partial x_{i}}\right|^{p-2} \frac{\partial u_{\varepsilon}}{\partial x_{i}}\right) \cdot v d x+\int_{\Omega}\left|\frac{\partial u_{\varepsilon}}{\partial x_{i}}\right|^{p-2} \frac{\partial u_{\varepsilon}}{\partial x_{i}} \cdot \frac{\partial v}{\partial x_{i}} d x=\int_{\Gamma_{1}}\left|\frac{\partial u_{\varepsilon}}{\partial x_{i}}\right|^{p-2} \frac{\partial u_{\varepsilon}}{\partial x_{i}} \cdot \nu_{i} \cdot v d \Gamma$.
Summing up from $i=1$ to $n$ on both sides of the above equation yields:

$$
\left(\mathcal{A}\left(u_{\varepsilon}\right), v\right)_{\Omega}=a\left(u_{\varepsilon}, v\right)-\sum_{i=1}^{n} \int_{\Gamma_{1}}\left|\frac{\partial u_{\varepsilon}}{\partial x_{i}}\right|^{p-2} \frac{\partial u_{\varepsilon}}{\partial x_{i}} \cdot \nu_{i} \cdot v d \Gamma .
$$

From this and observing that (see $\left.(* *)_{3}\right)$
$-\sum_{i=1}^{n} \int_{\Gamma_{1}}\left|\frac{\partial u_{\varepsilon}}{\partial x_{i}}\right|^{p-2} \frac{\partial u_{\varepsilon}}{\partial x_{i}} \cdot \nu_{i} \cdot v d \Gamma=\int_{\Gamma_{1}} \frac{\partial u_{\varepsilon}}{\partial t} v d \Gamma_{1}+\int_{\Gamma_{1}}\left|u_{\varepsilon}\right|^{\rho} u_{\varepsilon} v d \Gamma_{1}-\int_{\Gamma_{1}} f v d \Gamma_{1}$, we obtain $\left(\mathcal{A}\left(u_{\varepsilon}\right), v\right)_{\Omega}=a\left(u_{\varepsilon}, v\right)+\left(\gamma u_{\varepsilon}^{\prime}, \gamma v\right)_{\Gamma_{1}}+\left(\left|\gamma u_{\varepsilon}\right|^{\rho} \gamma u_{\varepsilon}, \gamma v\right)_{\Gamma_{1}}-(f, \gamma v)_{\Gamma_{1}}$. Substituting this identity in $(* *)_{1}$ we get:

$$
\varepsilon\left(u_{\varepsilon}^{\prime}, v\right)_{\Omega}+\left(\gamma u_{\varepsilon}^{\prime}, \gamma v\right)_{\Gamma_{1}}+a\left(u_{\varepsilon}, v\right)+\left(\left|\gamma u_{\varepsilon}\right|^{\rho} \gamma u_{\varepsilon}, \gamma v\right)_{\Gamma_{1}}=(f, \gamma v)_{\Gamma_{1}},
$$

where $u_{\varepsilon}^{\prime}$ means $\frac{\partial u_{\varepsilon}}{\partial t}$. Therefore, a solution of the problem ( $* *$ ) is understood in the following sense.

Definition 3.1. A real value function $u_{\varepsilon}(x, t)$ is a solution of the problem (**) if, only if,

$$
\begin{gather*}
u_{\varepsilon} \in L^{p}\left(0, T ; V_{0}\right) \cap L^{\infty}\left(0, T ; L^{2}(\Omega)\right),  \tag{4}\\
\gamma u_{\varepsilon} \in L^{\rho+2}\left(0, T ; L^{\rho+2}\left(\Gamma_{1}\right)\right) \cap L^{\infty}\left(0, T ; L^{2}\left(\Gamma_{1}\right)\right),  \tag{5}\\
u_{\varepsilon}^{\prime} \in L^{p^{\prime}}\left(0, T ; W^{-1, p^{\prime}}(\Omega)\right),  \tag{6}\\
\gamma u_{\varepsilon}^{\prime} \in L^{p^{\prime}}\left(0, T ; W^{\frac{-1}{p^{\prime}, p^{\prime}}}\left(\Gamma_{1}\right)\right),  \tag{7}\\
\varepsilon\left(u_{\varepsilon}^{\prime}, v\right)_{\Omega}+\left(\gamma u_{\varepsilon}^{\prime}, \gamma v\right)_{\Gamma_{1}}+a\left(u_{\varepsilon}, v\right)+\left(\left|\gamma u_{\varepsilon}\right|^{\rho} \gamma u_{\varepsilon}, \gamma v\right)_{\Gamma_{1}}=(f, \gamma v)_{\Gamma_{1}}, \quad \forall v \in V_{0} . \tag{8}
\end{gather*}
$$

and satisfying the initial conditions

$$
\begin{array}{cc}
u_{\varepsilon}(0)=w_{0} & \text { in } \quad \Omega, \\
\gamma u_{\varepsilon}(0)=u_{0} & \text { in } \quad \Gamma_{1} \tag{10}
\end{array}
$$

with $u_{0}$ belongs to $W^{\frac{1}{p^{\prime}}, p}\left(\Gamma_{1}\right)$.

## 4 Existence Theorem

In this section we will establish a theorem of existence of solutions.
Theorem 1. Suppose $f \in L^{p^{\prime}}\left(0, T, W^{-\frac{1}{p^{\prime}, p^{\prime}}}\left(\Gamma_{1}\right)\right)$ and $w_{0} \in V_{0}$. Then, for each $\varepsilon>0$ the problem (**) has a unique solution $u_{\varepsilon}$ in the sense of Definition (3.1).

Remark 1. Note that, the date $w_{0}$ is taken such that $\gamma w_{0}=u_{0}$, since, given $u_{0} \in W^{\frac{1}{p^{\prime}, p}}\left(\Gamma_{1}\right)$ there exists $w_{0} \in V_{0}$ such that $\gamma w_{0}=u_{0}$ because the application $\gamma: V_{0} \rightarrow W^{\frac{1}{p^{\prime}, p}}\left(\Gamma_{1}\right)$ is surjective.

Proof: We will employ the Faedo-Galerkin's method. In fact, for $V_{0}$ we construct a special Hilbertian basis $\left(w_{\mu}\right)_{\mu \in \mathbb{N}}$ of $V_{0}$. By $V_{0 m}=\left[w_{1}, \ldots ., w_{m}\right]$ we will denote the subspace spanned by the $m$ first vectors of $V_{0}$. The approximated problem consist to find a function $u_{\varepsilon m}(t) \in V_{0 m}$ of the type $u_{\varepsilon m}(x, t)=\sum_{j=1}^{m} g_{j \varepsilon m}(t) w_{j}(x)$ solution of the initial value problem for the system of ordinary differential equations:

$$
\begin{align*}
& \varepsilon\left(u_{\varepsilon m}^{\prime}(t), v\right)_{\Omega}+\left(\gamma u_{\varepsilon m}^{\prime}(t), \gamma v\right)_{\Gamma_{1}}+a\left(u_{\varepsilon m}(t), v\right)+  \tag{11}\\
& +\left(\left|\gamma u_{\varepsilon m}(t)\right|^{\rho} \gamma u_{\varepsilon m}(t), \gamma v\right)_{\Gamma_{1}}=(f(t), \gamma v)_{\Gamma_{1}} \quad \forall v \in V_{0} \\
& u_{\varepsilon m}(0)=u_{\varepsilon 0 m} \longrightarrow w_{0} \quad \text { in } \quad V_{0} .
\end{align*}
$$

The system (11) has a local solution on the interval $\left[0, t_{m}\left[\right.\right.$, with $t_{m}<T$. This solution can be extended to the whole interval $[0, T]$ as consequence of the a priori estimates that shall be proved in the next step.

## Estimates

Taking $v=u_{\varepsilon m}(t)$ in (11) and integrating from 0 to $t<t_{m}$ we obtain

$$
\begin{aligned}
& \frac{\varepsilon}{2}\left|u_{\varepsilon m}(t)\right|_{L^{2}(\Omega)}^{2}+\frac{1}{2}\left\|\gamma u_{\varepsilon m}(t)\right\|_{L^{2}\left(\Gamma_{1}\right)}^{2}+\int_{0}^{t}\left\|u_{\varepsilon m}(s)\right\|_{V_{0}}^{p} d s+\int_{0}^{t}\left\|\gamma u_{\varepsilon m}(s)\right\|_{L^{\rho+2}\left(\Gamma_{1}\right)}^{\rho+2} d s \leq \\
& \int_{0}^{t}\|f(s)\|_{W^{-\frac{1}{p^{\prime}, p^{\prime}}\left(\Gamma_{1}\right)}}\left\|\gamma u_{\varepsilon m}(s)\right\|_{W^{\frac{1}{p^{\prime, p}}\left(\Gamma_{1}\right)}} d s \leq C \int_{0}^{t}\|f(s)\|_{W^{-\frac{1}{p^{\prime}, p^{\prime}}\left(\Gamma_{1}\right)}}\left\|u_{\varepsilon m}(s)\right\|_{V_{0}} d s+ \\
& \frac{\varepsilon}{2}\left|u_{\varepsilon 0 m}\right|_{L^{2}(\Omega)}^{2}+\frac{1}{2}\left\|\gamma u_{\varepsilon 0 m}\right\|_{L^{2}\left(\Gamma_{1}\right)}^{2} \leq \frac{C}{p^{\prime}} \int_{0}^{T}\|f(s)\|_{W^{-\frac{1}{p^{\prime}, p^{\prime}}\left(\Gamma_{1}\right)}}^{p^{\prime}}+\frac{1}{p} \int_{0}^{t}\left\|u_{\varepsilon m}(s)\right\|_{V_{0}}^{p} d s+ \\
& \frac{\varepsilon}{2}\left|u_{\varepsilon 0 m}\right|_{L^{2}(\Omega)}^{2}+\frac{1}{2}\left\|\gamma u_{\varepsilon 0 m}\right\|_{L^{2}\left(\Gamma_{1}\right)}^{2} .
\end{aligned}
$$

From hypotheses about the initial conditions and the continuity of the application $\gamma$, we obtain:

$$
\begin{align*}
& \frac{\varepsilon}{2}\left|u_{\varepsilon m}(t)\right|_{L^{2}(\Omega)}^{2}+\frac{1}{2}\left\|\gamma u_{\varepsilon m}(t)\right\|_{L^{2}\left(\Gamma_{1}\right)}^{2}+\frac{1}{p^{\prime}} \int_{0}^{t}\left\|u_{\varepsilon m}(s)\right\|_{V_{0}}^{p} d s+ \\
& \int_{0}^{t}\left\|\gamma u_{\varepsilon m}(s)\right\|_{L^{\rho+2}\left(\Gamma_{1}\right)}^{\rho+2} d s \leq C \tag{12}
\end{align*}
$$

where $C$ is constant which is independent of $t$ and $m$. This estimate implies that we can prolong the approximate solution $u_{\varepsilon m}(t)$ to interval $[0, T]$ and too we obtain:

$$
\begin{align*}
& \left(u_{\varepsilon m}\right) \text { is bounded in } L^{\infty}\left(0, T ; L^{2}(\Omega)\right) ; \\
& \left(\gamma u_{\varepsilon m}\right) \text { is bounded in } L^{\infty}\left(0, T ; L^{2}\left(\Gamma_{1}\right)\right) ; \\
& \left(u_{\varepsilon m}\right) \text { is bounded in } L^{p}\left(0, T ; V_{0}\right) ; \\
& \left(\gamma u_{\varepsilon m}\right) \text { is bounded in } L^{p}\left(0, T ; W^{\frac{1}{p}, p}\left(\Gamma_{1}\right)\right) \text {; }  \tag{13}\\
& \left(\gamma u_{\varepsilon m}\right) \text { is bounded in } L^{\rho+2}\left(0, T ; L^{\rho+2}\left(\Gamma_{1}\right)\right) ; \\
& \left(u_{\varepsilon m}(T)\right) \text { is bounded in } L^{2}(\Omega) ; \\
& \left(\gamma u_{\varepsilon m}(T)\right) \text { is bounded in } L^{2}\left(\Gamma_{1}\right) .
\end{align*}
$$

Note that $\left.\left.\int_{0}^{T} \int_{\Gamma_{1}}| | \gamma u_{\varepsilon m}\right|^{\rho} \gamma u_{\varepsilon m}\right|^{\frac{\rho+2}{\rho+1}} d t=\int_{0}^{T} \int_{\Gamma_{1}}\left|\gamma u_{\varepsilon m}\right|^{\rho+2} d t \leq C$. Thus

$$
\begin{equation*}
\left(\left|\gamma u_{\varepsilon m}\right|^{\rho} \gamma u_{\varepsilon m}\right) \quad \text { is bounded in } \quad L^{\frac{\rho+2}{\rho+1}}\left(0, T ; L^{\frac{\rho+2}{\rho+1}}\left(\Gamma_{1}\right)\right) . \tag{14}
\end{equation*}
$$

From $(11)_{1}$ we get

$$
\left\langle\left\{\varepsilon u_{\varepsilon m}^{\prime}, \gamma u_{\varepsilon m}^{\prime}\right\},\{v, \gamma v\}\right\rangle=\left(f(t)-\left|\gamma u_{\varepsilon m}(t)\right|^{\rho} \gamma u_{\varepsilon m}(t), \gamma v\right)_{\Gamma_{1}}-a\left(u_{\varepsilon m}(t), v\right),
$$

where $<., .>$ represent the duality paring between $V_{0}^{\prime} \times W^{-\frac{1}{p^{\prime}, p^{\prime}}}\left(\Gamma_{1}\right)$ and $\left\{V_{0} \times W^{\frac{1}{p^{\prime}, p}}\left(\Gamma_{1}\right)\right\}$.

Hence

$$
\begin{aligned}
& \left|\left\langle\left\{\varepsilon u_{\varepsilon m}^{\prime}, \gamma u_{\varepsilon m}^{\prime}\right\},\{v, \gamma v\}\right\rangle\right| \leq\left|a\left(u_{\varepsilon m}(t), v\right)\right|+ \\
& \left(\|f(t)\|_{W^{-\frac{1}{p^{\prime}, p^{\prime}}}\left(\Gamma_{1}\right)}+\left\|\left.\gamma u_{\varepsilon m}\right|^{\rho} \gamma u_{\varepsilon m}\right\|_{\left.W^{-\frac{1}{p^{\prime}, p^{\prime}}}{ }_{\left(\Gamma_{1}\right)}\right)}\right)\|\gamma v\|_{W^{\frac{1}{p^{\prime}, p}}\left(\Gamma_{1}\right)} \leq \\
& \left(\|f(t)\|_{W^{-\frac{1}{p^{\prime,}, p^{\prime}}}{ }_{\left(\Gamma_{1}\right)}}+C\left\|\left.\gamma u_{\varepsilon m}\right|^{\rho} \gamma u_{\varepsilon m}\right\|_{L^{\frac{\rho+2}{\rho+1}\left(\Gamma_{1}\right)}}\right)\|\gamma v\|_{W^{\frac{1}{p}, p}\left(\Gamma_{1}\right)}+ \\
& \left\|u_{\varepsilon m}(t)\right\|_{V_{0}}^{p-1}\|v\|_{V_{0}} \leq\left(\|\gamma v\|_{W^{\frac{1}{p^{, p}}\left(\Gamma_{1}\right)}}+\|v\|_{V_{0}}\right) \times \\
& {\left[\|f(t)\|_{W^{-\frac{1}{p^{\prime}, p^{\prime}}\left(\Gamma_{1}\right)}}+C\left\|\left.\gamma u_{\varepsilon m}\right|^{\rho} \gamma u_{\varepsilon m}\right\|_{L^{\frac{\rho+2}{\rho+1}\left(\Gamma_{1}\right)}}+\left\|u_{\varepsilon m}(t)\right\|_{V_{0}}^{p-1}\right] .}
\end{aligned}
$$

From estimates above we have

$$
\begin{gathered}
\left|\left\langle\left\{\varepsilon u_{\varepsilon m}^{\prime}, \gamma u_{\varepsilon m}^{\prime}\right\},\{v, \gamma v\}\right\rangle\right| \leq \\
{\left[\|f(t)\|_{\left.W^{-\frac{1}{p^{\prime}, p^{\prime}}}{ }_{\left(\Gamma_{1}\right)}+C\left\|\left.\gamma u_{\varepsilon m}\right|^{\rho} \gamma u_{\varepsilon m}\right\|_{L^{\frac{\rho+2}{\rho+1}}\left(\Gamma_{1}\right)}+\left\|u_{\varepsilon m}(t)\right\|_{V_{0}}^{p-1}\right] \times}^{\|\{v, \gamma v\}\|_{V_{0} \times W^{\frac{1}{p^{\prime}, p}\left(\Gamma_{1}\right)}} .}\right. \text {. }}
\end{gathered}
$$

Therefore, from bounded (14) and (13) ${ }_{3}$, we get

$$
\begin{equation*}
\left(\left\{\varepsilon u_{\varepsilon m}^{\prime}, \gamma u_{\varepsilon m}^{\prime}\right\}\right) \quad \text { is bounded in } \quad L^{p^{\prime}}\left(0, T ; V_{0}^{\prime} \times W^{-\frac{1}{p^{\prime}, p^{\prime}}}\left(\Gamma_{1}\right)\right) . \tag{15}
\end{equation*}
$$

## Passage to the Limit

From estimates (13), (14) and (15) we obtain

$$
\left\lvert\, \begin{align*}
& u_{\varepsilon m} \stackrel{*}{\rightharpoonup} u_{\varepsilon} \quad \text { weak-star in } \quad L^{\infty}\left(0, T ; L^{2}(\Omega)\right) ;  \tag{16}\\
& \gamma u_{\varepsilon m} \stackrel{*}{\rightharpoonup} \gamma u_{\varepsilon} \quad \text { weak-star in } \quad L^{\infty}\left(0, T ; L^{2}\left(\Gamma_{1}\right)\right) ; \\
& u_{\varepsilon m} \rightharpoonup u_{\varepsilon} \quad \text { weak in } \quad L^{p}\left(0, T ; V_{0}\right) ; \\
& \gamma u_{\varepsilon m} \rightharpoonup \gamma u_{\varepsilon} \quad \text { weak in } \quad L^{p}\left(0, T ; W^{\frac{1}{p^{\prime}, p}}\left(\Gamma_{1}\right)\right) ; \\
& \gamma u_{\varepsilon m} \rightharpoonup \gamma u_{\varepsilon} \quad \text { weak in } \quad L^{\rho+2}\left(0, T ; L^{\rho+2}\left(\Gamma_{1}\right)\right) \equiv L^{\rho+2}\left(\Sigma_{1}\right) ; \\
& u_{\varepsilon m}(T) \rightharpoonup \chi \quad \text { weak in } \quad L^{2}(\Omega) ; \\
& \gamma u_{\varepsilon m}(T) \rightharpoonup \varsigma \quad \text { weak in } \quad L^{2}\left(\Gamma_{1}\right) \\
& \left|\gamma u_{\varepsilon m}\right|^{\rho} \gamma u_{\varepsilon m} \rightharpoonup \eta \quad \text { weak in } \quad L^{\frac{\rho+2}{\rho+1}}\left(0, T ; L^{\frac{\rho+2}{\rho+1}}\left(\Gamma_{1}\right)\right) \equiv L^{\frac{\rho+2}{\rho+1}}\left(\Sigma_{1}\right) \\
& u_{\varepsilon m}^{\prime} \rightharpoonup u_{\varepsilon}^{\prime} \quad \text { weak in } \quad L^{p^{\prime}}\left(0, T ; V_{0}^{\prime}\right) ; \\
& \gamma u_{\varepsilon m}^{\prime} \rightharpoonup \gamma u_{\varepsilon}^{\prime} \quad \text { weak in } \quad L^{p^{\prime}}\left(0, T ; W^{-\frac{1}{p^{\prime}, p^{\prime}}}\left(\Gamma_{1}\right)\right)
\end{align*}\right.
$$

Remark 2. Note that by the convergence $(16)_{1},(16)_{9}$ and $(16)_{2},(16)_{10}$ it makes sense to calculate $u_{\varepsilon}(0), u_{\varepsilon}(T)$ and $\gamma u_{\varepsilon}(0), \gamma u_{\varepsilon}(T)$ respectively.

Let $V=L^{p}\left(0, T ; V_{0}\right)$ and $\mathcal{B}$ the operator from $V$ given by

$$
\begin{equation*}
(\mathcal{B}(u), v)_{V^{\prime} \times V}=\int_{0}^{T} a(u(t), v(t)) d t, \quad \forall u, v \in V, \tag{17}
\end{equation*}
$$

hence, $\mathcal{B}$ is hemicontinuous, monotonic operator and $\|\mathcal{B} u\|_{V^{\prime}} \leq C\|u\|_{V}^{p-1}, \quad \forall u \in$ $V$. Thus, from estimative $(16)_{3}$ we have $\left(\mathcal{B} u_{\varepsilon m}\right)_{m}$ is bounded in $V^{\prime}$, hence

$$
\begin{equation*}
\mathcal{B} u_{\varepsilon m} \rightharpoonup \zeta \quad \text { is } \quad V^{\prime} \tag{18}
\end{equation*}
$$

From convergence (16) $)_{9}$ we obtain $\left\langle u_{\varepsilon m}^{\prime}, \varphi>\rightarrow\left\langle u_{\varepsilon}^{\prime}, \varphi\right\rangle, \quad \forall \varphi \in L^{p}\left(0, T ; V_{0}\right)\right.$, that is,
$\int_{0}^{T}\left(u_{\varepsilon m}^{\prime}, v\right)_{\Omega} \theta d t \rightarrow \int_{0}^{T}\left(u_{\varepsilon}^{\prime}, v\right)_{\Omega} \theta d t, \quad v \in V_{0 m} \subset V_{0}, \quad \forall \theta \in \mathcal{D}(0, T) \subset L^{p}(0, T)$, or

$$
\begin{equation*}
\left(u_{\varepsilon m}^{\prime}, v\right)_{\Omega} \rightarrow\left(u_{\varepsilon}^{\prime}, v\right)_{\Omega} \quad v \in V_{0 m}, \quad \text { in } \quad \mathcal{D}^{\prime}(0, T) . \tag{19}
\end{equation*}
$$

From convergence (16) $)_{10}$ we have $<\gamma u_{\varepsilon m}^{\prime}, \gamma \varphi>\rightarrow<\gamma u_{\varepsilon}^{\prime}, \gamma \varphi>, \quad \forall \varphi \in$ $L^{p}\left(0, T ; V_{0}\right)$, that is,

$$
\begin{equation*}
\left(\gamma u_{\varepsilon m}^{\prime}, \gamma v\right)_{\Gamma_{1}} \rightarrow\left(\gamma u_{\varepsilon}^{\prime}, \gamma v\right)_{\Gamma_{1}} \quad v \in V_{0 m}, \quad \text { in } \quad \mathcal{D}^{\prime}(0, T) . \tag{20}
\end{equation*}
$$

Analogously, we have

$$
\begin{equation*}
\left(\mathcal{B} u_{\varepsilon m}, v\right) \rightarrow(\zeta, v) \quad v \in V_{0 m}, \quad \text { in } \quad \mathcal{D}^{\prime}(0, T), \tag{21}
\end{equation*}
$$

and

$$
\begin{equation*}
\left(\left|\gamma u_{\varepsilon m}\right|^{\rho} \gamma u_{\varepsilon m}, \gamma v\right)_{\Gamma_{1}} \rightarrow(\eta, \gamma v)_{\Gamma_{1}} \quad v \in V_{0 m}, \quad \text { in } \quad \mathcal{D}^{\prime}(0, T) . \tag{22}
\end{equation*}
$$

Thus, taking the limit as $m \rightarrow \infty$ in the approximated equation (11) $)_{1}$, using the convergence (19) - (22) and the density of $V_{0 m}$ in $V_{0}$, we obtain:

$$
\begin{gather*}
\varepsilon\left(u_{\varepsilon}^{\prime}, v\right)_{\Omega}+\left(\gamma u_{\varepsilon}^{\prime}, \gamma v\right)_{\Gamma_{1}}+(\zeta, v)_{\Omega}+(\eta, \gamma v)_{\Gamma_{1}}=(f(t), \gamma v)_{\Gamma_{1}}  \tag{23}\\
\forall v \in V_{0} \quad \text { in } \quad \mathcal{D}^{\prime}(0, T) .
\end{gather*}
$$

To follow we will proof that: $\left|\gamma u_{\varepsilon}\right|^{\rho} \gamma u_{\varepsilon}=\eta$ and $\mathcal{B} u_{\varepsilon}=\zeta$.

Proof. In fact, from the estimate above we have

$$
\begin{array}{ll}
\left(\gamma u_{\varepsilon m}\right) \text { is bounded in } & L^{p}\left(0, T ; W^{\frac{1}{p^{\prime}, p}}\left(\Gamma_{1}\right)\right) \text {; } \\
\left(\gamma u_{\varepsilon m}^{\prime}\right) \text { is bounded in } & L^{p^{\prime}}\left(0, T ; W^{-\frac{1}{p^{\prime}, p^{\prime}}}\left(\Gamma_{1}\right)\right) .
\end{array}
$$

As $W^{\frac{1}{p^{\prime}, p}}\left(\Gamma_{1}\right) \stackrel{c}{\hookrightarrow} L^{p}\left(\Gamma_{1}\right) \hookrightarrow L^{2}\left(\Gamma_{1}\right) \hookrightarrow L^{p^{\prime}}\left(\Gamma_{1}\right) \hookrightarrow W^{-\frac{1}{p^{\prime}, p^{\prime}}}\left(\Gamma_{1}\right)$, we have $W^{\frac{1}{p^{\prime}, p}}\left(\Gamma_{1}\right) \stackrel{c}{\hookrightarrow} L^{p}\left(\Gamma_{1}\right) \hookrightarrow W^{-\frac{1}{p^{\prime}, p^{\prime}}}\left(\Gamma_{1}\right)$. Thus, of the Aubin-Lions's Theorem we obtain a subsequence, still denoted by $\left(\gamma u_{\varepsilon m}\right)$, such that $\gamma u_{\varepsilon m} \rightarrow \gamma u_{\varepsilon}$ in $L^{p}\left(\Sigma_{1}\right)$, where still we can extract other subsequence which we insist in denote by $\left(\gamma u_{\varepsilon m}\right)$, such that: $\gamma u_{\varepsilon m} \rightarrow \gamma u_{\varepsilon}$ a.e $\Sigma_{1}$, thus,

$$
\begin{equation*}
\left|\gamma u_{\varepsilon m}\right|^{\rho} \gamma u_{\varepsilon m} \rightarrow\left|\gamma u_{\varepsilon}\right|^{\rho} \gamma u_{\varepsilon} \quad \text { a.e } \quad \Sigma_{1} . \tag{24}
\end{equation*}
$$

From estimative (14) we obtain:

$$
\begin{equation*}
\left\|\left|\gamma u_{\varepsilon m}\right|^{\rho} \gamma u_{\varepsilon m}\right\|_{L^{\frac{\rho+2}{\rho+1}\left(\Sigma_{1}\right)}} \leq C \tag{25}
\end{equation*}
$$

Hence, from (24), (25) and the Lions's Lema, we obtain

$$
\left|\gamma u_{\varepsilon m}\right|^{\rho} \gamma u_{\varepsilon m} \rightharpoonup\left|\gamma u_{\varepsilon}\right|^{\rho} \gamma u_{\varepsilon} \quad \text { weak } \quad \text { in } \quad L^{\frac{\rho+2}{\rho+1}}\left(\Sigma_{1}\right) .
$$

Thus $\left|\gamma u_{\varepsilon}\right|^{\rho} \gamma u_{\varepsilon}=\eta$.
Now we will prove that $\mathcal{B} u_{\varepsilon}=\zeta$.
Proof. In fact for this purpose we needed prove that: $(i) u_{\varepsilon}(0)=w_{0},(i i)$ $\chi=u_{\varepsilon}(T),(i i i) \gamma u_{\varepsilon}(0)=u_{0}$ and (iv) $\gamma u_{\varepsilon}(T)=\varsigma$.

In fact, to prove $(i)$ we use the convergence $(16)_{1}$ and $(16)_{9}$ which yield

$$
\begin{gathered}
\int_{0}^{T}\left(u_{\varepsilon m}, v\right) \varphi^{\prime} d t \rightarrow \int_{0}^{T}\left(u_{\varepsilon}, v\right) \varphi^{\prime} d t \quad \forall v \in L^{2}(\Omega), \\
\varphi \in C^{1}([0, T]), \varphi(0)=1, \varphi(T)=0
\end{gathered}
$$

and

$$
\int_{0}^{T} \frac{d}{d t}\left(u_{\varepsilon m}, v\right) \varphi d t \rightarrow \int_{0}^{T} \frac{d}{d t}\left(u_{\varepsilon}, v\right) \varphi d t \quad \forall v \in L^{2}(\Omega) \subset V_{0}^{\prime}
$$

$$
\varphi \in C^{1}([0, T]), \varphi(0)=1, \varphi(T)=0
$$

where

$$
\begin{gathered}
\int_{0}^{T} \frac{d}{d t}\left\{\left(u_{\varepsilon m}, v\right) \varphi\right\} d t \rightarrow \int_{0}^{T} \frac{d}{d t}\left\{\left(u_{\varepsilon}, v\right) \varphi\right\} d t \quad \forall v \in L^{2}(\Omega), \\
\varphi \in C^{1}([0, T]), \varphi(0)=1, \varphi(T)=0
\end{gathered}
$$

hence,

$$
\left(u_{\varepsilon m}(0), v\right) \rightarrow\left(u_{\varepsilon}(0), v\right), \quad \forall v \in L^{2}(\Omega)
$$

that is,

$$
u_{\varepsilon m}(0) \rightharpoonup u_{\varepsilon}(0) \quad \text { in } \quad L^{2}(\Omega)
$$

As $u_{\varepsilon m}(0) \rightarrow w_{0}$ in $V_{0} \hookrightarrow L^{2}(\Omega)$, we have $u_{\varepsilon m}(0) \rightarrow w_{0}$ in $L^{2}(\Omega)$, where $u_{\varepsilon m}(0) \rightharpoonup w_{0}$ in $L^{2}(\Omega)$. Hence

$$
u_{\varepsilon}(0)=w_{0} .
$$

Analogously, working as $\varphi(0)=0$ and $\varphi(T)=1$ we obtain

$$
u_{\varepsilon}(T)=\chi .
$$

In fact, to prove (iii) we use the convergence $(16)_{2}$ and $(16)_{10}$ which yield

$$
\left(\gamma u_{\varepsilon m}(0), v\right) \rightarrow\left(\gamma u_{\varepsilon}(0), v\right), \quad \forall v \in L^{2}\left(\Gamma_{1}\right),
$$

where

$$
\gamma u_{\varepsilon m}(0) \rightharpoonup \gamma u_{\varepsilon}(0) \quad \text { em } \quad L^{2}\left(\Gamma_{1}\right) .
$$

As $u_{\varepsilon m}(0) \rightarrow w_{0}$ in $V_{0}$ and $\gamma$ continuous from $V_{0}$ in $W^{\frac{1}{p}, p}\left(\Gamma_{1}\right)$, we have that $\gamma u_{\varepsilon m}(0) \rightarrow \gamma w_{0}$ in $W^{\frac{1}{p}, p}\left(\Gamma_{1}\right)$. Being $W^{\frac{1}{p^{\prime}, p}}\left(\Gamma_{1}\right) \hookrightarrow L^{2}\left(\Gamma_{1}\right)$, for $p>2>\frac{2 n}{2 n+1}$, by Fractionary Sobolev's Theorem, we have $\gamma u_{\varepsilon m}(0) \rightarrow \gamma w_{0}$ in $L^{2}\left(\Gamma_{1}\right)$. Therefore $\gamma u_{\varepsilon}(0)=\gamma w_{0}=u_{0}$, by remark 1. Analogously, we have $\gamma u_{\varepsilon}(T)=$ $\varsigma$

We will show that: $\mathcal{B} u_{\varepsilon}=\zeta$. In fact, being the operator $\mathcal{B}: V \rightarrow V^{\prime}$ mononotonic, we obtain:

$$
\begin{equation*}
\left(\mathcal{B} u_{\varepsilon m}, u_{\varepsilon m}\right)-\left(\mathcal{B} u_{\varepsilon m}, v\right)-\left(\mathcal{B} v, u_{\varepsilon m}-v\right) \geq 0 \tag{26}
\end{equation*}
$$

Taking $v=u_{\varepsilon m}$ and integrating of 0 the $T$ in the approximated equation (11) ${ }_{1}$ we obtain:

$$
\begin{aligned}
\left(\mathcal{B} u_{\varepsilon m}, u_{\varepsilon m}\right)= & \int_{0}^{T} a\left(u_{\varepsilon m}, u_{\varepsilon m}\right) d t=\int_{0}^{T}\left(f, u_{\varepsilon m}\right)_{\Gamma_{1}} d t-\varepsilon \int_{0}^{T}\left(u_{\varepsilon m}^{\prime}, u_{\varepsilon m}\right)_{\Omega} d t- \\
& \int_{0}^{T}\left(\gamma u_{\varepsilon m}^{\prime}, \gamma u_{\varepsilon m}\right)_{\Gamma_{1}} d t-\int_{0}^{T}\left(\left|\gamma u_{\varepsilon m}\right|^{\rho} \gamma u_{\varepsilon m}, \gamma u_{\varepsilon m}\right)_{\Gamma_{1}} d t .
\end{aligned}
$$

Thus, substituting in (26), we have

$$
\begin{aligned}
0 \leq & \int_{0}^{T}\left(f, u_{\varepsilon m}\right)_{\Gamma_{1}} d t-\frac{\varepsilon}{2}\left|u_{\varepsilon m}(T)\right|_{L^{2}(\Omega)}^{2}+\frac{\varepsilon}{2}\left|u_{\varepsilon m}(0)\right|_{L^{2}(\Omega)}^{2}-\frac{1}{2}\left|\gamma u_{\varepsilon m}(T)\right|_{L^{2}\left(\Gamma_{1}\right)}^{2}+ \\
& \frac{1}{2}\left|\gamma u_{\varepsilon m}(0)\right|_{L^{2}\left(\Gamma_{1}\right)}^{2}-\int_{0}^{T}\left\|\gamma u_{\varepsilon m}\right\|_{L^{\rho+2}\left(\Gamma_{1}\right)}^{\rho+2} d t-\left(\mathcal{B} u_{\varepsilon m}, v\right)-\left(\mathcal{B} v, u_{\varepsilon m}-v\right) .
\end{aligned}
$$

Using the convergence obtained and applying the $\liminf _{m \rightarrow \infty}$ in both sides of the inequality above we have:

$$
\begin{align*}
0 \leq & \int_{0}^{T}\left(f, u_{\varepsilon}\right)_{\Gamma_{1}} d t-\frac{\varepsilon}{2}\left|u_{\varepsilon}(T)\right|_{L^{2}(\Omega)}^{2}+\frac{\varepsilon}{2}\left|w_{0}\right|_{L^{2}(\Omega)}^{2}-\frac{1}{2}\left|\gamma u_{\varepsilon}(T)\right|_{L^{2}\left(\Gamma_{1}\right)}^{2}+ \\
& \frac{1}{2}\left|u_{0}\right|_{L^{2}\left(\Gamma_{1}\right)}^{2}-\int_{0}^{T}\left\|\gamma u_{\varepsilon}\right\|_{L^{\rho+2}\left(\Gamma_{1}\right)}^{\rho+2} d t-(\zeta, v)-\left(\mathcal{B} v, u_{\varepsilon}-v\right) . \tag{27}
\end{align*}
$$

Taking $v=u_{\varepsilon}$ and integrating of 0 the $T$ in the equation (23) we obtain:

$$
\begin{aligned}
& \int_{0}^{T}\left(f, u_{\varepsilon}\right)_{\Gamma_{1}} d t=\left(\zeta, u_{\varepsilon}\right)+\frac{\varepsilon}{2}\left|u_{\varepsilon}(T)\right|_{L^{2}(\Omega)}^{2}-\frac{\varepsilon}{2}\left|w_{0}\right|_{L^{2}(\Omega)}^{2}+ \\
& \frac{1}{2}\left|\gamma u_{\varepsilon}(T)\right|_{L^{2}\left(\Gamma_{1}\right)}^{2}-\frac{1}{2}\left|u_{0}\right|_{L^{2}\left(\Gamma_{1}\right)}^{2}+\int_{0}^{T}\left\|\gamma u_{\varepsilon}\right\|_{L^{\rho+2}\left(\Gamma_{1}\right)}^{\rho+2} d t .
\end{aligned}
$$

If we substitute this expression in (27), we obtain

$$
0 \leq\left(\zeta-\mathcal{B} v, u_{\varepsilon}-v\right), \quad \forall \quad v \in V
$$

Consider $u_{\varepsilon}-v=\lambda w, \quad \lambda>0$. Thus, using the hemicontinuity of the operator $\mathcal{B}$, we obtain $0 \leq\left(\zeta-\mathcal{B} u_{\varepsilon}, w\right), \quad \forall w \in V$. Working with $\lambda<0$ we have:
$\left(\zeta-\mathcal{B} u_{\varepsilon}, w\right) \leq 0, \quad \forall w \in V$.
Therefore $\left(\zeta-\mathcal{B} u_{\varepsilon}, w\right)=0, \quad \forall w \in V$, thus $\mathcal{B} u_{\varepsilon}=\zeta$.
Note that $\left(\mathcal{B} u_{\varepsilon m}, w\right) \rightarrow\left(\mathcal{B} u_{\varepsilon}, w\right), \quad \forall w \in V=L^{p}\left(0, T ; V_{0}\right)$, hence,

$$
\int_{0}^{T} a\left(u_{\varepsilon m}, v\right) \theta d t \rightarrow \int_{0}^{T} a\left(u_{\varepsilon}, v\right) \theta d t, \quad \forall v \in V_{0}, \quad \forall \theta \in \mathcal{D}(0, T) \subset L^{p}(0, T)
$$

Thus $a\left(u_{\varepsilon m}, v\right) \rightarrow a\left(u_{\varepsilon}, v\right) \quad \forall v \in V_{0}$ in $\mathcal{D}^{\prime}(0, T)$. Therefore,

$$
\left\lvert\, \begin{align*}
& \varepsilon\left(u_{\varepsilon}^{\prime}, v\right)_{\Omega}+\left(\gamma u_{\varepsilon}^{\prime}, \gamma v\right)_{\Gamma_{1}}+a\left(u_{\varepsilon}, v\right)+  \tag{28}\\
& \quad\left(\left|\gamma u_{\varepsilon}\right|^{\rho} \gamma u_{\varepsilon}, \gamma v\right)_{\Gamma_{1}}=(f(t), \gamma v)_{\Gamma_{1}} \quad \forall v \in V_{0} \quad \text { in } \quad \mathcal{D}^{\prime}(0, T),
\end{align*}\right.
$$

## Uniqueness of the Solution

To obtain the uniqueness of the solution, we suppose that there exists two solutions such that $u_{\varepsilon}, \widehat{u_{\varepsilon}}$ in the conditions of the Theorem 1. It following that $w_{\varepsilon}=u_{\varepsilon}-\widehat{u_{\varepsilon}}$ satisfy:

$$
\begin{gather*}
w_{\varepsilon} \in L^{p}\left(0, T ; V_{0}\right) \cap L^{\infty}\left(0, T ; L^{2}(\Omega)\right),  \tag{29}\\
\gamma w_{\varepsilon} \in L^{\rho+2}\left(0, T ; L^{\rho+2}\left(\Gamma_{1}\right)\right) \cap L^{\infty}\left(0, T ; L^{2}\left(\Gamma_{1}\right)\right),  \tag{30}\\
w_{\varepsilon}^{\prime} \in L^{p^{\prime}}\left(0, T ; W^{-1, p^{\prime}}(\Omega)\right),  \tag{31}\\
\gamma w_{\varepsilon}^{\prime} \in L^{p^{\prime}}\left(0, T ; W^{\left.-\frac{1}{p^{\prime}, p^{\prime}}\left(\Gamma_{1}\right)\right),}\right.  \tag{32}\\
\varepsilon\left(w_{\varepsilon}^{\prime}, v\right)_{\Omega}+\left(\gamma w_{\varepsilon}^{\prime}, \gamma v\right)_{\Gamma_{1}}+a\left(u_{\varepsilon}, v\right)-a\left(\widehat{u_{\varepsilon}}, v\right)+  \tag{33}\\
\left(\left|\gamma u_{\varepsilon}\right|^{\rho} \gamma u_{\varepsilon}-\left|\gamma \widehat{u_{\varepsilon}}\right|^{\rho} \gamma \widehat{u_{\varepsilon}}, \gamma v\right)_{\Gamma_{1}}=0, \quad \forall v \in V_{0} .
\end{gather*}
$$

Taking $v=w_{\varepsilon}$ in (33) and integrating from 0 the $t \leq T$ we obtain:

$$
\begin{gathered}
\frac{\varepsilon}{2}\left|w_{\varepsilon}(t)\right|_{L^{2}(\Omega)}^{2}+\frac{1}{2}\left|\gamma w_{\varepsilon}(t)\right|_{L^{2}\left(\Gamma_{1}\right)}^{2}+\int_{0}^{t}\left(a\left(u_{\varepsilon}, w\right)-a\left(\widehat{u_{\varepsilon}}, w\right)\right) d t+ \\
\int_{0}^{t}\left(\left|\gamma u_{\varepsilon}\right|^{\rho} \gamma u_{\varepsilon}-\left|\gamma \widehat{u_{\varepsilon}}\right|^{\rho} \gamma \widehat{u_{\varepsilon}}, \gamma u_{\varepsilon}-\gamma \widehat{u_{\varepsilon}}\right) d t=0 .
\end{gathered}
$$

Using the monotoneity of the function $h(s)=|s|^{\rho} s$ and $a\left(u_{\varepsilon}, w\right)-a\left(\widehat{u_{\varepsilon}}, w\right) \geq 0$, we have

$$
\frac{\varepsilon}{2}\left|w_{\varepsilon}(t)\right|_{L^{2}(\Omega)}^{2}+\frac{1}{2}\left|\gamma w_{\varepsilon}(t)\right|_{L^{2}\left(\Gamma_{1}\right)}^{2} \leq 0 .
$$

Therefore, we have that $w_{\varepsilon}(t)=0 \quad \forall t \in[0, T]$. Thus the Theorem is proved.

## 5 Main Result

In this Section we will prove the following result
Theorem 2. When $\varepsilon \rightarrow 0$ we have

$$
\begin{equation*}
u_{\varepsilon} \rightharpoonup u \quad \text { in } \quad L^{p}\left(0, T ; V_{0}\right), \tag{34}
\end{equation*}
$$

being $u$ the solution of the problem (*).

Proof: Making $v=u_{\varepsilon}(t)$ in (8) and proceeding as in the previous Theorem we obtain

$$
\begin{align*}
& \left(\sqrt{\varepsilon} u_{\varepsilon}\right) \text { is bounded in } L^{\infty}\left(0, T ; L^{2}(\Omega)\right) ; \\
& \left(\gamma u_{\varepsilon}\right) \text { is bounded in } L^{\infty}\left(0, T ; L^{2}\left(\Gamma_{1}\right)\right) ; \\
& \left(u_{\varepsilon}\right) \text { is bounded in } L^{p}\left(0, T ; V_{0}\right) ; \\
& \left(\gamma u_{\varepsilon}\right) \text { is bounded in } L^{p}\left(0, T ; W^{\frac{1}{p}, p}\left(\Gamma_{1}\right)\right) ;  \tag{35}\\
& \left(\gamma u_{\varepsilon}\right) \text { is bounded in } L^{\rho+2}\left(0, T ; L^{\rho+2}\left(\Gamma_{1}\right)\right) ; \\
& \left(u_{\varepsilon}(T)\right) \text { is bounded in } L^{2}(\Omega) ; \\
& \left(\gamma u_{\varepsilon}(T)\right) \text { is bounded in } L^{2}\left(\Gamma_{1}\right) .
\end{align*}
$$

and

$$
\left\lvert\, \begin{align*}
& \left(\varepsilon u_{\varepsilon}^{\prime}\right) \text { is bounded in } L^{p^{\prime}}\left(0, T ; V_{0}^{\prime}\right) ; \\
& \left(\left|\gamma u_{\varepsilon}\right|^{\rho} \gamma u_{\varepsilon}\right) \text { is bounded in } L^{\frac{\rho+2}{\rho+1}}\left(0, T ; L^{\frac{\rho+2}{\rho+1}}\left(\Gamma_{1}\right)\right)  \tag{36}\\
& \left(\gamma u_{\varepsilon}\right) \text { is bounded in } L^{p}\left(0, T ; W^{\frac{1}{p^{\prime}, p}}\left(\Gamma_{1}\right)\right) ; \\
& \left(\gamma u_{\varepsilon}^{\prime}\right) \text { is bounded in } L^{p^{\prime}}\left(0, T ; W^{-\frac{1}{p^{\prime}, p^{\prime}}}\left(\Gamma_{1}\right)\right) \\
& \left(\mathcal{B} u_{\varepsilon}\right) \text { is bounded in } V^{\prime}=L^{p^{\prime}}\left(0, T ; V_{0}^{\prime}\right) .
\end{align*}\right.
$$

Hence there exists an subsequence, still represented by $\left(u_{\varepsilon}\right)$, such that, when $\varepsilon \rightarrow 0$

$$
\left\lvert\, \begin{align*}
& \sqrt{\varepsilon} u_{\varepsilon} \stackrel{*}{\rightharpoonup} 0 \quad \text { weak-star in } \quad L^{\infty}\left(0, T ; L^{2}(\Omega)\right) ; \\
& \gamma u_{\varepsilon} \stackrel{*}{\rightharpoonup} \gamma u \quad \text { weak-star in } \quad L^{\infty}\left(0, T ; L^{2}\left(\Gamma_{1}\right)\right) ; \\
& u_{\varepsilon} \rightharpoonup u \quad \text { weak in } \quad V=L^{p}\left(0, T ; V_{0}\right) ; \\
& \gamma u_{\varepsilon} \rightharpoonup \gamma u \quad \text { weak in } \quad L^{p}\left(0, T ; W^{\frac{1}{p^{\prime}, p}}\left(\Gamma_{1}\right)\right) ; \\
& \gamma u_{\varepsilon} \rightharpoonup \gamma u \quad \text { weak in } L^{\rho+2}\left(0, T ; L^{\rho+2}\left(\Gamma_{1}\right)\right) \equiv L^{\rho+2}\left(\Sigma_{1}\right) ;  \tag{37}\\
& \gamma u_{\varepsilon}(T) \rightharpoonup \chi \quad \text { weak in } L^{2}\left(\Gamma_{1}\right) \\
& \varepsilon u_{\varepsilon}^{\prime} \rightharpoonup 0 \quad \text { weak in } V^{\prime}=L^{p^{\prime}}\left(0, T ; V_{0}^{\prime}\right) ; \\
& \left|\gamma u_{\varepsilon}\right|^{\rho} \gamma u_{\varepsilon} \rightharpoonup \eta \quad \text { weak in } \quad L^{\frac{\rho+2}{\rho+1}}\left(0, T ; L^{\frac{\rho+2}{\rho+1}}\left(\Gamma_{1}\right)\right) \\
& \gamma u_{\varepsilon}^{\prime} \rightharpoonup \gamma u^{\prime} \quad \text { weak in } L^{p^{\prime}}\left(0, T ; W^{-\frac{1}{p^{\prime}, p^{\prime}}}\left(\Gamma_{1}\right)\right) ; \\
& \mathcal{B} u_{\varepsilon} \rightharpoonup \zeta \text { weak in } V^{\prime}=L^{p^{\prime}}\left(0, T ; V_{0}^{\prime}\right) .
\end{align*}\right.
$$

Analogously to the Theorem 1 we can to show that $|\gamma u|^{\rho} \gamma u=\eta$ and $\mathcal{B} u=\zeta$. Using the convergence (37) in (28) we obtain the variational formulation of the problem (*), when $\varepsilon \rightarrow 0$
$\left(\gamma u^{\prime}, \gamma v\right)_{\Gamma_{1}}+a(u, v)+\left(|\gamma u|^{\rho} \gamma u, \gamma v\right)_{\Gamma_{1}}=(f(t), \gamma v)_{\Gamma_{1}} \quad \forall v \in V_{0} \quad$ in $\quad \mathcal{D}^{\prime}(0, T)$.

On the other hand, with a analogously analysis as in the Remark 2, we have as in the Theorem 1: $\gamma u_{\varepsilon}(0) \rightharpoonup \gamma u(0)$ in $L^{2}\left(\Gamma_{1}\right)$. Therefore

$$
\gamma u(0)=u_{0} \quad \text { on } \quad \Gamma_{1} .
$$

In this sense, we have the solution of the problem $(*)$ as limit of the perturbed problem ( $* *$ ).

## 6 Boundary Stabilization

The aim of section is study the algebric decay for the energy $E(t)$ associated to weak solution of the problem (*). To asymptotic behavior, we use the Nakao's
method [9]. This energy is given by

$$
\begin{equation*}
E(t)=\frac{1}{2}|\gamma u(t)|_{L^{2}\left(\Gamma_{1}\right)}^{2} \tag{38}
\end{equation*}
$$

Remark 3. Note that, the solution of the problem (*) we can be extend to $[0, \infty)$, when $f=0$.

Theorem 3. Let $E(t)$ a energy associated the weak solution of problem $(*)$.
Then, there exists a constant $\delta>0$ such that the energy satisfies

$$
E(t) \leq C \frac{1}{(1+t)^{\frac{1}{\delta}}}, \forall t \geq 0
$$

Considering (28) with $f=0$ and $v=u_{\epsilon}$ we get:

$$
\frac{1}{2} \frac{d}{d t}\left|\sqrt{\epsilon} u_{\epsilon}(t)\right|^{2}+\frac{1}{2} \frac{d}{d t}\left|\gamma u_{\epsilon}(t)\right|^{2}=-\left\|u_{\epsilon}(t)\right\|_{V_{0}}^{p}-\left\|\gamma u_{\epsilon}(t)\right\|_{L^{\rho+2}(\Omega)}^{\rho+2}
$$

Let $E_{\varepsilon}(t)=\frac{1}{2}\left|\sqrt{\epsilon} u_{\epsilon}(t)\right|^{2}+\frac{1}{2}\left|\gamma u_{\epsilon}(t)\right|^{2}$, then

$$
\begin{equation*}
\frac{d}{d t} E_{\varepsilon}(t) \leq 0, \forall t \geq 0 \tag{39}
\end{equation*}
$$

and

$$
E_{\varepsilon}(t) \leq E_{\varepsilon}(0)=\frac{1}{2}\left|\sqrt{\epsilon} u_{\epsilon}(0)\right|^{2}+\frac{1}{2}\left|\gamma u_{\epsilon}(0)\right|^{2} \leq C\left(\gamma u_{0}\right)
$$

because, $\frac{1}{2} \sqrt{\varepsilon} u_{\varepsilon}(0) \rightarrow 0$ in $L^{2}(\Omega)$ and $\gamma u_{\varepsilon}(0) \rightarrow \gamma u(0)$ in $L^{2}\left(\Gamma_{1}\right)$ when $\varepsilon \rightarrow 0$. Therefore, $E_{\varepsilon}(t)$ is increasing and bounded.

We have that

$$
\begin{equation*}
\frac{d}{d t} E_{\varepsilon}(t)=-\left\|u_{\varepsilon}(t)\right\|_{V_{0}}^{p}-\left\|\gamma u_{\varepsilon}(t)\right\|_{L^{\rho+2}(\Omega)}^{\rho+2} \tag{40}
\end{equation*}
$$

then

$$
E_{\varepsilon}(t+1)-E_{\varepsilon}(t)=-\int_{t}^{t+1}\left\|u_{\varepsilon}(t)\right\|_{V_{0}}^{p}-\int_{t}^{t+1}\left\|\gamma u_{\varepsilon}(t)\right\|_{L^{\rho+2}(\Omega)}^{\rho+2} .
$$

Therefore

$$
\begin{equation*}
\int_{t}^{t+1}\left\|u_{\varepsilon}(t)\right\|_{V_{0}}^{p}+\int_{t}^{t+1}\left\|\gamma u_{\varepsilon}(t)\right\|_{L^{\rho+2}(\Omega)}^{\rho+2}=E_{\varepsilon}(t)-E_{\varepsilon}(t+1) \tag{41}
\end{equation*}
$$

Since $V_{0} \hookrightarrow L^{2}(\Omega)$ and $L^{\rho+2}\left(\Gamma_{1}\right) \hookrightarrow L^{2}\left(\Gamma_{1}\right)$, we have

$$
\begin{equation*}
\int_{t}^{t+1}\left|u_{\varepsilon}(t)\right|_{L^{2}(\Omega)}^{p}+\int_{t}^{t+1}\left|\gamma u_{\varepsilon}(t)\right|_{L^{2}\left(\Gamma_{1}\right)}^{p} \leq C_{1}\left(E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right) \tag{42}
\end{equation*}
$$

Therefore exist $t^{*} \in[t, t+1]$, such that

$$
\left|u_{\varepsilon}\left(t^{*}\right)\right|_{L^{2}(\Omega)}^{p}+\left|\gamma u_{\varepsilon}\left(t^{*}\right)\right|_{L^{2}\left(\Gamma_{1}\right)}^{p} \leq C_{1}\left[E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right],
$$

where,

$$
\begin{aligned}
& \frac{1}{2}\left|u_{\varepsilon}\left(t^{*}\right)\right|^{2} \leq C_{2}\left[E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right]^{\frac{2}{p}} \\
& \frac{1}{2}\left|\gamma u_{\varepsilon}\left(t^{*}\right)\right|^{2} \leq C_{2}\left[E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right]^{\frac{2}{p}}
\end{aligned}
$$

Thus

$$
\begin{aligned}
& \frac{\varepsilon}{2}\left|u_{\varepsilon}\left(t^{*}\right)\right|^{2} \leq \varepsilon C_{2}\left[E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right]^{\frac{2}{p}} \\
& \frac{1}{2}\left|\gamma u_{\varepsilon}\left(t^{*}\right)\right|^{2} \leq C_{2}\left[E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right]^{\frac{2}{p}}
\end{aligned}
$$

we obtain

$$
\frac{1}{2}\left|\sqrt{\varepsilon} u_{\varepsilon}\left(t^{*}\right)\right|^{2}+\frac{1}{2}\left|\gamma u_{\varepsilon}\left(t^{*}\right)\right|^{2} \leq(1+\varepsilon) C_{2}\left[E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right]^{\frac{2}{p}}
$$

Therefore,

$$
\begin{equation*}
E_{\varepsilon}\left(t^{*}\right) \leq(1+\varepsilon) C_{2}\left[E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right]^{\frac{2}{p}} . \tag{43}
\end{equation*}
$$

Integrating from (40) of $t$ to $t^{*}$, we obtain

$$
\begin{aligned}
& E_{\varepsilon}(t)=E_{\varepsilon}\left(t^{*}\right)+\int_{t}^{t^{*}}\left\|u_{\varepsilon}(t)\right\|_{V_{0}}^{p}+\int_{t}^{t^{*}}\left\|\gamma u_{\varepsilon}(t)\right\|_{L^{p}\left(\Gamma_{1}\right)}^{p} \leq \\
& \left((1+\varepsilon) C_{2}\right)\left(E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right)^{\frac{2}{p}}+\left(E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right)
\end{aligned}
$$

Therefore,

$$
\begin{aligned}
& E_{\varepsilon}(t)^{\frac{p}{2}} \leq \max \left\{1,(1+\varepsilon) C_{2}\right\}^{\frac{p}{2}}\left(E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right)+\left(E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right)^{\frac{p}{2}} \leq \\
& \max \left\{1,(1+\varepsilon) C_{2}\right\}^{\frac{p}{2}}\left[E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right]\left[1+\left(E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right)^{\frac{p-2}{2}}\right] .
\end{aligned}
$$

Being, $E_{\varepsilon}(t)$ limited for all $\varepsilon>0$, follows that

$$
E_{\varepsilon}(t)^{\frac{p}{2}} \leq C_{3} \max \left\{1,(1+\varepsilon) C_{2}\right\}^{\frac{p}{2}}\left(E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right)
$$

or

$$
E_{\varepsilon}(t)^{\frac{p}{2}} \leq C_{1 \varepsilon}\left(E_{\varepsilon}(t)-E_{\varepsilon}(t+1)\right),
$$

where $C_{1 \varepsilon}=C_{3} \max \left\{1,(1+\varepsilon) C_{2}\right\}^{\frac{p}{2}}$.
Thus, by Nakao's Lemma, there exists $\delta>0$ such that

$$
E_{1 \varepsilon}(t) \leq C_{1 \varepsilon} \frac{1}{(1+t)^{\frac{1}{\delta}}}
$$

Note that $C_{1 \varepsilon} \rightarrow \max \left\{1, C_{2}\right\}^{\frac{p}{2}}, \gamma u_{\varepsilon}(t) \rightharpoonup \gamma u(t)$ weak in $L^{2}\left(\Gamma_{1}\right)$ and $\sqrt{\varepsilon} u_{\varepsilon}(t) \rightharpoonup$ 0 weak in $L^{2}(\Omega)$.

From

$$
\liminf _{\varepsilon \rightarrow 0} E_{1 \varepsilon}(t) \leq \liminf _{\varepsilon \rightarrow 0} C_{1 \varepsilon} \frac{1}{(1+t)^{\frac{1}{\delta}}}, \forall t \geq 0
$$

implies that

$$
\liminf _{\varepsilon \rightarrow 0}\left\{\frac{1}{2}\left|\sqrt{\varepsilon} u_{\varepsilon}(t)\right|_{L^{2}(\Omega)}^{2}+\frac{1}{2}\left|\gamma u_{\varepsilon}(t)\right|_{L^{2}\left(\Gamma_{1}\right)}^{2}\right\} \leq \max \left\{1, C_{2}\right\}^{\frac{p}{2}} \frac{1}{(1+t)^{\frac{1}{\delta}}}, \forall t \geq 0
$$

Thus,

$$
\frac{1}{2}|\gamma u(t)|_{L^{2}\left(\Gamma_{1}\right)}^{2} \leq C_{1} \frac{1}{(1+t)^{\frac{1}{\delta}}}, \forall t \geq 0
$$

where $C_{1}=\max \left\{1, C_{2}\right\}^{\frac{p}{2}}$.

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