# A New Computational Method Based on Integral Transform for Solving Linear and Nonlinear Fractional Systems 

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#### Abstract

Abstrak. Pada artikel ini, Metode Elzaki homotopy perturbation (EHPM) diterapkan untuk menyelesaikan sistem fraksional stiff. Metode EHPM adalah kombinasi dari modifikasi transformasi integral Laplace yang disebut transformasi Elzaki dan metode gangguan homotopi. Metode yang diusulkan telah diterapkan pada beberapa contoh sistem fraksional stiff linier dan nonlinier. Hasil yang diperoleh dengan metode ini dibandingkan dengan hasil yang diperoleh dengan metode Kernel Hilbert space (KHSM) menunjukkan bahwa metode EHPM merupakan metode yang efektif dan akurat untuk menyelesaikan sistem persamaan diferensial orde fraksional.


#### Abstract

In this article, the Elzaki homotopy perturbation method is applied to solve fractional stiff systems. The Elzaki homotopy perturbation method (EHPM) is a combination of a modified Laplace integral transform called the Elzaki transform and the homotopy perturbation method. The proposed method is applied for some examples of linear and nonlinear fractional stiff systems. The results obtained by the current method were compared with the results obtained by the kernel Hilbert space KHSM method. The obtained result reveals that the Elzaki homotopy perturbation method is an effective and accurate technique to solve the systems of differential equations of fractional order.


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## 1. Introduction

Many scientific disciplines can be modeled by initial value problems of fractional order which lead to a better understanding in characterizing the natural phenomena in various fields of science such as engineering, biology, economy, computer science, and physics. The exact and analytical solutions of fractional order differential equations are difficult to be found, thus the numerical methods are used to investigate the solutions of differential equations of fractional orders. The proposed system at first was highlighted by Hirschfelder and Curtiss [1]. Then stiff systems were studied by many researchers in different branches [2]-[4]. Several numerical methods have been developed to obtain the analytical and approximate solutions of stiff systems of integer and fractional orders, such as Adomian decomposition method [5], homotopy perturbation method [6], homotopy analysis method [7], variation iteration method [8], rational homotopy perturbation method[9], Block method [10], Laplace adomian decomposition method and modified decomposition method [11], multistage Bernstein method[12], and fractional power series method [13].

Homotopy perturbation method (HPM) is an important semi-analytic technique for solving differential equations [14]-[16] and an efficient technique to study the different types of nonlinear functional equations. Volterra-Fredholm nonlinear systems were solved by HPM [17], hyperbolic PDEs [18], Zakharov-Kuznetsov[19], and system of nonlinear differential equations[20].

So far, many modified techniques have been described to solve linear and nonlinear differential equations of integer and fractional orders. Those of which are Laplace homotopy perturbation method, Laplace Adomian decomposition method, and recently, Elzaki homotopy transformation perturbation method is used to solve a range of problems such as a family of Fisher's equation [21], spatial diffusion of Biological population [22], nonlinear oscillators [23], system of linear and nonlinear PDEs of fractional orders [24], time-fractional Navier-Stokes equations [25], and An algorithm for solving the BurgersHuxley equation using the Elzaki transform [26]

Elzaki integral transform is a modification of the Laplace and Sumudu transforms which was invented by Tariq [27], Elzaki transformation is an efficient and powerful technique that has found the exact solutions to several differential equations which cannot be solved by Sumudu transform [28]. Elzaki integral equation is a powerful and efficient technique that has been used to solve many differential equations of integer and fractional orders [25][29]-[33].

The objective of this paper is to expand the applications of EHPM and illustrate the efficiency of the proposed method, thus we consider the stiff systems fractional ordinary differential equations:
$D^{\mu_{i}} z_{i}(t)+F_{i}\left(t, z_{1}(t), z_{2}(t), \ldots, z_{n}(t)\right)=f_{i}(t), z_{i}\left(t_{0}\right)=a_{i, 0}, m-1<\mu_{i}<m, m \in N$.

## 2. Preliminaries

In this section, we introduce some definitions and properties of fractional calculus and Elzaki transform which are used in this article.
Definition 2.1. [4] A real valued function $g(y), y>0$ is belongs to the space $\mathrm{C}_{\sigma}, \sigma \in R$ if there exists at least a real number $d>\sigma$, such that $g(y)=y^{d} g_{1}(y)$ where $g_{1}(y) \in$ $C(0, \infty)$, and it is said to be in the space $C_{\sigma}^{n}$ if $g^{n} \in R_{\sigma}, n \in N$.
Definition 2.2. [34] The function $f(u)$ is called R-L fractional integral of order $\propto \geq 0$ if it is defined as:
$J^{\alpha} f(u)=\frac{1}{\Gamma(\alpha)} \int_{0}^{u} \quad(u-t)^{\alpha-1} f(t) d t, \alpha>0, t>0$.
In particular $J^{0} f(u)=f(u)$.
For $\theta \geq 0$ and $\vartheta \geq-1$, some properties of the operator $J^{\alpha}$

1. $J^{\alpha} J^{\theta} f(u)=J^{\alpha+\theta} f(u)$,
2. $J^{\alpha} J^{\theta} f(u)=J^{\theta} J^{\alpha} f(u)$,
3. $J^{\alpha} x^{\vartheta}=\frac{\Gamma(\vartheta+1)}{\Gamma(\alpha+\vartheta+1)} x^{\alpha+\vartheta}$.

Definition 2.3. [34] The function $f \in C_{-1}^{n}, n \in N$ is called Caputo fractional derivative if it is defined as

$$
D^{\alpha} f(u)=\frac{1}{\Gamma(n-\alpha)} \int_{0}^{u}(u-t)^{n-\alpha-1} f^{n}(t) d t, n-1<\alpha \leq n
$$

Definition 2.4. [27] The Elzaki-transform of the function $f(u)$ is defined as:
$E[f(u)]=T(v)=v \int_{0}^{\infty} \quad f(u) e^{\frac{-u}{v}} d u \quad u>0$.
Suppose that $f$ is piecewise continuous, then we can calculate $E\left[\frac{\partial f}{\partial x}\right]$ as follows:
$E\left[\frac{\partial f(x, u)}{\partial x}\right]=\int_{0}^{\infty} \quad v e^{\frac{-u}{v}} \frac{\partial f(x, u)}{\partial x} d u=\frac{\partial}{\partial x} \int_{0}^{\infty} \quad v e^{\frac{-u}{v}} f(x, u) d u=\frac{\partial}{\partial x} T(x, v)$,
Similarly, we can have:

$$
E\left[\frac{\partial^{2} f(x, u)}{\partial x^{2}}\right]=\frac{d^{2} T(x, u)}{d x^{2}}
$$

Assume that $\frac{\partial f}{\partial u}=h$, then we have:

$$
\begin{gathered}
E\left[\frac{\partial^{2} f(x, u)}{\partial u^{2}}\right]=E\left[\frac{\partial h(x, u)}{\partial u}\right]=\frac{1}{v} E[h(x, u)]-v h(x, 0) \\
E\left[\frac{\partial^{2} f(x, u)}{\partial u^{2}}\right]=\frac{T(x, u)}{v^{2}}-f(x, 0)-v \frac{\partial f}{\partial u}(x, 0)
\end{gathered}
$$

By mathematical induction one can extend this result to the $n^{t h}$ partial derivative as follows:

$$
\begin{equation*}
E\left[\frac{\partial^{n} f(x, u)}{\partial u^{n}}\right]=\frac{T(x, t)}{v^{n}}-\sum_{i=0}^{n-1} \quad v^{2-n+i} \frac{\partial^{i} f(x, 0)}{\partial u^{i}} \tag{2}
\end{equation*}
$$

## 3. Elzaki homotopy perturbation method analysis

Assume the following system of nonlinear differential equations

$$
\begin{equation*}
M_{i}\left(z_{i}\right)+N_{i}\left(z_{i}\right)=g_{i}(t), \quad t \in \Lambda, i=1,2, \ldots, n \tag{3}
\end{equation*}
$$

where $z_{i}$ and $g_{i}(t)$ are sought and known functions respectively, $M_{i}$ and $N_{i}$ are linear and nonlinear operators respectively, where:

$$
N_{i}\left(z_{i}\right)=\sum_{k=0}^{\infty} A_{i k}
$$

where $A_{i k}$ represents the Adomian polynomial as follows:

$$
A_{i k}=\frac{1}{m!} \frac{d^{m}}{d \tau^{m}}\left(N_{i}\left(\sum_{k=0}^{\infty} \quad z_{1 k} \tau^{m}, \ldots, \quad \sum_{k=0}^{\infty} \quad z_{n k} \tau^{m}\right)\right.
$$

Now, we assume the following homotopy

$$
\begin{equation*}
И_{i}\left(\Upsilon_{i}, p\right)=(1-p)\left(M_{i}\left(\Upsilon_{i}\right)-z_{i, 0}\right)+p\left(M_{i}\left(\Upsilon_{i}\right)+N_{i}\left(\Upsilon_{i}\right)-g_{i}\right) \tag{4}
\end{equation*}
$$

or equivalently, $\quad U_{i}\left(\Upsilon_{i}, p\right)=M_{i}\left(\Upsilon_{i}\right)-z_{i, 0}+p z_{i, 0}+p\left(N_{i}\left(\Upsilon_{i}\right)-g_{i}\right)$,
Since $M_{i}\left(\Upsilon_{i}\right)-z_{i, 0}=0$ and $M_{i}\left(\Upsilon_{i}\right)+N_{i}\left(\Upsilon_{i}\right)-g_{i}=0$ which are equivalent to the operator equations $\Lambda_{i}(z, 0)=0$ and $U_{i}(z, 1)=0$ respectively, also $p \in[0,1]$ is a homotopy parameter, $z_{i, 0}$ are initial approximations and $\Upsilon_{i}:(t, p): \Lambda \times[0,1] \rightarrow R$.
Now, we apply the Elzaki transform on (4), we obtain

$$
E\left[M_{i}\left(\Upsilon_{i}\right)-z_{i, 0}+p z_{i, 0}+p\left(N_{i}\left(\Upsilon_{i}\right)-g_{i}\right)\right]=0
$$

By differential property of of Elzaki transform, we get

$$
\begin{gathered}
\frac{1}{v^{\mu}} E\left[\Upsilon_{i}\right]-v^{2-\mu} \Upsilon_{i, 0}-v^{3-\mu} \Upsilon_{i, 0}^{\prime}-\cdots-v^{n+1-\mu} \Upsilon_{i, 0}^{(n-1)} \\
=E\left[z_{i, 0}-p z_{i, 0}+p\left(N_{i}\left(\Upsilon_{i}\right)-g_{i}\right)\right]
\end{gathered}
$$

Applying the inverse Elzaki transform, we obtain

$$
\begin{align*}
\Upsilon_{i}= & E^{-1}\left[v^{2} \Upsilon_{i, 0}+v^{3} \Upsilon_{i, 0}^{\prime}+\cdots+v^{n+1} \Upsilon_{i, 0}^{(n-1)}\right]+E^{-1}\left[v ^ { \mu } E \left[z_{i, 0}-p z_{i, 0}+\right.\right. \\
& \left.\left.p\left(N_{i}\left(\Upsilon_{i}\right)-g_{i}\right)\right]\right] . \tag{5}
\end{align*}
$$

By HPM, the following series represent the solution of equation (5), where $p \in[0,1]$.

$$
\begin{equation*}
\Upsilon_{i}(t)=\sum_{j=0}^{\infty} \quad p^{j} \Upsilon_{i, j}, \quad i=1,2, \ldots, n \tag{6}
\end{equation*}
$$

Substituting (6) in (5), we obtain

$$
\begin{aligned}
\sum_{n=0}^{\infty} p^{n} \Upsilon_{i, n}= & E^{-1}\left[v^{2} \Upsilon_{i, 0}+v^{3} \Upsilon_{i, 0}^{\prime}+\cdots+v^{n+1} \Upsilon_{i, 0}^{(n-1)}\right] \\
& +E^{-1}\left[v^{\mu} E\left[z_{i, 0}-p z_{i, 0}+p\left(N_{i}\left(\sum_{n=0}^{\infty} p^{n} \Upsilon_{i, n}\right)-g_{i}\right)\right]\right]
\end{aligned}
$$

Now, by comparing the coefficients of the same powers of the embedded parameter $p$, leads to

$$
\begin{gathered}
p^{0}: \Upsilon_{i, 0}=E^{-1}\left[v^{2} \Upsilon_{i, 0}+v^{3} \Upsilon_{i, 0}^{\prime}+\cdots+v^{n+1} \Upsilon_{i, 0}^{(n-1)}+E\left[z_{0}\right]\right] \\
p^{1}: \Upsilon_{i, 1}=E^{-1}\left[v^{\mu} E\left[N_{i}\left(\Upsilon_{i, 0}\right)-z_{i, 0}-g_{i}\right]\right] \\
p^{2}: \Upsilon_{i, 2}=E^{-1}\left[v^{\mu} E\left[N_{i}\left(\Upsilon_{i, 0}, \Upsilon_{i, 1}\right)\right]\right] \\
p^{3}: \Upsilon_{i, 3}=E^{-1}\left[v^{\mu} E\left[N_{i}\left(\Upsilon_{i, 0}, \Upsilon_{i, 1}, \Upsilon_{i, 2}\right)\right]\right] \\
\vdots \\
p^{j}: \Upsilon_{i, j}=E^{-1}\left[v^{\mu} E\left[N_{i}\left(\Upsilon_{i, 0}, \Upsilon_{i, 1}, \Upsilon_{i, 2}, \ldots, \Upsilon_{i, j-1}\right)\right]\right]
\end{gathered}
$$

Assume that $\Upsilon_{i, 0}=\beta_{i, 0}, \Upsilon_{i, 0}^{\prime}=\beta_{i, 1}, \ldots, \Upsilon_{i, 0}^{(n-1)}=\beta_{i, n-1}$ are the initial approximations, where $i=1,2, \ldots, n$. Finally, the solution of (3) is obtained by the following series:

$$
\begin{equation*}
z_{i}(t)=\Upsilon_{i}=\Upsilon_{i, 0}+\Upsilon_{i, 1}+\Upsilon_{i, 2}+\cdots+ \tag{7}
\end{equation*}
$$

The convergence and uniqueness of the series (7) is discussed in [21].

## 4. Applications of Elzaki homotopy perturbation method

The numerical patterns are employed to confirm the efficiency and accuracy of the Elzaki homotopy transform perturbation technique for solving the fractional stiff systems. For this purpose, at the first example, the comparison of error analysis with the kernel Hilbert space technique [13] [35] is made, at the second example, the exact solution of the proposed problem is achieved where $\beta=1$, and at the third example the comparison between the proposed method and (ADM and MLDM [11]) is made.

Example 4.1 Consider the following linear stiff system of fractional order $0<\beta \leq 1$

$$
\begin{equation*}
\frac{D_{t}^{\beta} z(t)=-z(t)+95 w(t)}{D_{t}^{\beta} w(t)=-z(t)-97 w(t)} \tag{8}
\end{equation*}
$$

with the initial conditions $z(0)=1, w(0)=1$. The exact solution of the system when $\beta=$ 1 is

$$
z(t)=\frac{95}{47} e^{-2 t}-\frac{48}{47} e^{-96 t}, \quad w(t)=\frac{48}{47} e^{-96 t}-\frac{1}{47} e^{-2 t}
$$

To solve system (8) by the EHPM, we use the following homotopy

$$
\begin{gathered}
D_{t}^{\beta} Z(t)-z_{0}(t)+p\left(z_{0}(t)+Z(t)-95 W(t)\right)=0 \\
D_{t}^{\beta} W(t)-w_{0}(t)+p\left(w_{0}(t)+Z(t)+97 W(t)\right)=0
\end{gathered}
$$

Applying the Elzaki transform, we have

$$
\begin{gathered}
E\left[D_{t}^{\beta} Z(t)-z_{0}(t)+p\left(z_{0}(t)+Z(t)-95 W(t)\right)\right]=0 \\
E\left[D_{t}^{\beta} W(t)-w_{0}(t)+p\left(w_{0}(t)+Z(t)+97 W(t)\right)\right]=0 .
\end{gathered}
$$

Using the differential property of Elzaki transform, we obtain $\frac{1}{v^{\beta}} E[Z(t)]=v^{2-\beta} Z_{0}(t)+$

$$
\begin{gathered}
E\left[z_{0}(t)-p\left(z_{0}(t)+Z(t)-95 W(t)\right)\right] \\
\frac{1}{v^{\beta}} E[W(t)]=v^{2-\beta} W_{0}(t)+E\left[w_{0}(t)-p\left(z_{0}(t)+Z(t)+97 W(t)\right)\right] .
\end{gathered}
$$

Applying inverse Elzaki transform, we have

$$
\begin{equation*}
\frac{Z(t)=E^{-1}\left[v^{2} Z_{0}(t)+v^{\beta} E\left[z_{0}(t)-p\left(z_{0}(t)+Z(t)-95 W(t)\right)\right]\right]}{W(t)=E^{-1}\left[v^{2} W_{0}(t)+v^{\beta} E\left[w_{0}(t)-p\left(z_{0}(t)+Z(t)+97 W(t)\right)\right]\right]} \tag{9}
\end{equation*}
$$

Here, the following form is the solution of equation (9):

$$
\begin{equation*}
\frac{Z(t)=Z_{0}(t)+p Z_{1}(t)+p^{2} Z_{2}(t)+\cdots,}{W(t)=W_{0}(t)+p W_{1}(t)+p^{2} W_{2}(t)+\cdots,} \tag{10}
\end{equation*}
$$

Substituting (10) in to (9) and collecting the coefficients of equivalent powers of embedded parameter $p$, we obtain the following:

$$
\begin{aligned}
& p^{0}:\left\{Z_{0}(t)=E^{-1}\left[v^{2} Z_{0}(0)+v^{\beta} E\left[z_{0}(t)\right]\right], W_{0}(t)=E^{-1}\left[v^{2} W_{0}(0)+v^{\beta} E\left[w_{0}(t)\right]\right],\right. \\
& p^{1}:\left\{Z_{1}(t)=E^{-1}\left[-v^{\beta} E\left[\left(z_{0}(t)+Z_{0}(t)-95 W_{0}(t)\right)\right]\right], W_{1}(t)\right. \\
& \quad=E^{-1}\left[-v^{\beta} E\left[\left(z_{0}(t)+Z_{0}(t)+97 W_{0}(t)\right)\right]\right], \\
& p^{2}:\left\{Z_{2}(t)=\right. E^{-1}\left[-v^{\beta} E\left[\left(Z_{1}(t)-95 W_{1}(t)\right)\right]\right], W_{2}(t)=E^{-1}\left[-v^{\beta} E\left[\left(Z_{1}(t)\right.\right.\right. \\
&\left.\left.\left.+97 W_{1}(t)\right)\right]\right], \\
& p^{3}:\left\{Z_{3}(t)=\right. E^{-1}\left[-v^{\beta} E\left[\left(Z_{2}(t)-95 W_{2}(t)\right)\right], W_{3}(t)=E^{-1}\left[-v^{\beta} E\left[\left(Z_{2}(t)\right.\right.\right.\right. \\
&\left.\left.\left.+97 W_{2}(t)\right)\right]\right],
\end{aligned}
$$

Here, we set $Z_{0}(0)=z_{0}(t)=1$ and $W_{0}(0)=w_{0}(t)=1$. Thus, the above equations lead to the following results:

$$
\begin{array}{ll}
Z_{0}(t)=1+\frac{t^{\beta}}{\Gamma(\beta+1)^{\prime}} & W_{0}(t)=1+\frac{t^{\beta}}{\Gamma(\beta+1)^{\prime}} \\
Z_{1}(t)=\frac{93}{\Gamma(\beta+1)} t^{\beta}+\frac{94}{\Gamma(2 \beta+1)} t^{2 \beta}, & W_{1}(t)=-\frac{99}{\Gamma(\beta+1)} t^{\beta}-\frac{98}{\Gamma(2 \beta+1)} t^{2 \beta}, \\
Z_{2}(t)=-\frac{9498}{\Gamma(\beta+1)} t^{2 \beta}-\frac{9404}{\Gamma(3 \beta+1)} t^{3 \beta}, & W_{2}(t)=\frac{9510)}{\Gamma(2 \beta+1)} t^{2 \beta}+\frac{9412}{\Gamma(3 \beta+1)} t^{3 \beta} \\
Z_{3}(t)=\frac{9129486}{\Gamma(3 \beta+1)} t^{3 \beta}+\frac{9035446}{\Gamma(4 \beta+1)} t^{4 \beta}, & W_{3}(t)=-\frac{912972}{\Gamma(3 \beta+1)} t^{3 \beta}-\frac{903560}{\Gamma(4 \beta+1)} t^{4 \beta},
\end{array}
$$

Thus, the solution of system (8) using the series (7) is

$$
\begin{aligned}
& z(t)=1+\frac{94}{\Gamma(\beta+1)} t^{\beta}-\frac{9404}{\Gamma(2 \beta+1)} t^{2 \beta}+\frac{903544}{\Gamma(3 \beta+1)} t^{3 \beta} \ldots, \\
& w(t)=1-\frac{98}{\Gamma(\beta+1)} t^{\beta}+\frac{9412}{\Gamma(2 \beta+1)} t^{2 \beta}-\frac{903560}{\Gamma(3 \beta+1)} t^{3 \beta} \ldots
\end{aligned}
$$

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Figure 1. The behavior of EHPM solution of system (8) for different values of $\beta$ and exact solution.
Figure 1: illustrates the behavior of the EHPM solution of system (8) for different values of $\beta$ and exact solution. It can be seen that there is a good agreement between the exact solution and the approximate solution using EHPM, especially when $\beta=1$.

Table 1. The Absolute Error of kernel Hilbert space method KHSM and EHPM of (8)

| $\boldsymbol{t}$ | $\boldsymbol{z}(\boldsymbol{t}) \mathbf{E H P M}$ | $\boldsymbol{z}(\boldsymbol{t}) \mathbf{K H S M}$ | $\boldsymbol{w}(\boldsymbol{t}) \mathbf{E H P M}$ | $\boldsymbol{w}(\boldsymbol{t}) \mathbf{K H S M}$ |
| :---: | :---: | :---: | :---: | :---: |
| 0.00 | 0 | 0 | 0 | 0 |
| 0.5 | $3.0163443 \times 10^{-15}$ | $4.76706570 \times 10^{-4}$ | $3.0163483 \times 10^{-15}$ | $4.44439990 \times 10^{-4}$ |
| 0.1 | $5.0178941 \times 10^{-14}$ | $4.11985309 \times 10^{-5}$ | $3.4785609 \times 10^{-13}$ | $1.37132721 \times 10^{-5}$ |
| 0.15 | $6.4189477 \times 10^{-11}$ | $2.37714435 \times 10^{-5}$ | $1.1162905 \times 10^{-10}$ | $4.49704615 \times 10^{-7}$ |
| 0.2 | $7.3684661 \times 10^{-9}$ | $1.99138385 \times 10^{-5}$ | $1.2169510 \times 10^{-9}$ | $2.1210257 \times 10^{-7}$ |
| 0.25 | $3.4096780 \times 10^{-7}$ | $1.67357026 \times 10^{-5}$ | $3.6738962 \times 10^{-7}$ | $1.76192995 \times 10^{-7}$ |

Table 1: shows the comparison between the absolute errors of the approximate solution of the proposed technique at different values of $t$ when $\beta=1$ for 20 order approximation and the approximate solution of KHSM. The results illustrate that the proposed technique is superior to KHSM and an efficient method for investigating the solution of proposed fractional systems.

Example 4.2: Consider the following non-linear stiff system of fractional order

$$
\begin{equation*}
\frac{D_{t}^{\beta} z(t)=-\left(\rho^{-1}+2\right) z(t)+\rho^{-1} w^{2}(t),}{D_{t}^{\beta} w(t)=z(t)-w(t)-w^{2}(t),} \tag{11}
\end{equation*}
$$

with the initial conditions $z(0)=1, w(0)=1$. The exact solution of the system when $\beta=$ 1 is

$$
z(t)=e^{-2 t}, \quad w(t)=e^{t} .
$$

To solve system (11) by the EHPM, we use the following homotopy

$$
\begin{aligned}
& D_{t}^{\beta} Z(t)-z_{0}(t)+p\left(z_{0}(t)+\left(\rho^{-1}+2\right) Z(t)-\rho^{-1} W^{2}(t)\right)=0 \\
& D_{t}^{\beta} W(t)-w_{0}(t)+p\left(w_{0}(t)-Z(t)+W(t)+W^{2}(t)\right)=0 .
\end{aligned}
$$

Applying the Elzaki transform, we have

$$
\begin{gathered}
E\left[D_{t}^{\beta} Z(t)-z_{0}(t)+p\left(z_{0}(t)+\left(\rho^{-1}+2\right) Z(t)-\rho^{-1} W^{2}(t)\right)\right]=0 \\
E\left[D_{t}^{\beta} W(t)-w_{0}(t)+p\left(w_{0}(t)-Z(t)+W(t)+W^{2}(t)\right)\right]=0 .
\end{gathered}
$$

Using the differential property of Elzaki transform, we obtain $\frac{1}{v^{\beta}} E[Z(t)]=v^{2-\beta} Z_{0}(t)+$

$$
\begin{gathered}
E\left[z_{0}(t)-p\left(z_{0}(t)+\left(\rho^{-1}+2\right) Z(t)-\rho^{-1} W^{2}(t)\right)\right] \\
\frac{1}{v^{\beta}} E[W(t)]=v^{2-\beta} W_{0}(t)+E\left[w_{0}(t)-p\left(w_{0}(t)-Z(t)+W(t)+W^{2}(t)\right)\right] .
\end{gathered}
$$

Applying inverse Elzaki transform, we have

$$
\begin{equation*}
\frac{Z(t)}{}=E^{-1}\left[v^{2} Z_{0}(t)+v^{\beta} E\left[z_{0}(t)-p\left(z_{0}(t)+\left(\rho^{-1}+2\right) Z(t)-\rho^{-1} W^{2}(t)\right)\right]\right] . \tag{12}
\end{equation*}
$$

Here, the following form is the solution of system (11):

$$
\begin{equation*}
\frac{Z(t)=Z_{0}(t)+p Z_{1}(t)+p^{2} Z_{2}(t)+\cdots,}{W(t)=W_{0}(t)+p W_{1}(t)+p^{2} W_{2}(t)+\cdots} \tag{13}
\end{equation*}
$$

Substituting (13) in to (12) and collecting the coefficients of equivalent powers of embedded parameter $p$, we obtain the following:

$$
\begin{aligned}
p^{0}: & \left\{Z_{0}(t)=E^{-1}\left[v^{2} Z_{0}(0)+v^{\beta} E\left[z_{0}(t)\right]\right], W_{0}(t)=E^{-1}\left[v^{2} W_{0}(0)+v^{\beta} E\left[w_{0}(t)\right]\right],\right. \\
& p^{1}:\left\{Z_{1}(t)=E^{-1}\left[-v^{\beta} E\left[z_{0}(t)+\left(\rho^{-1}+2\right) Z_{0}(t)-\rho^{-1} W_{0}^{2}(t)\right]\right], W_{1}(t)\right. \\
& =E^{-1}\left[-v^{\beta} E\left[w_{0}(t)-Z_{0}(t)+W_{0}(t)+W_{0}^{2}(t)\right]\right], \\
& p^{j}:\left\{Z_{2}(t)\right. \\
= & E^{-1}\left[-v^{\beta} E\left[\left(\rho^{-1}+2\right) Z_{j-1}(t)\right.\right. \\
- & \left.\left.\rho^{-1} \sum_{l=0}^{j-1} W_{j}(t) W_{j-l-1}(t)\right]\right], \\
+ & \left.\left.W_{j-1}(t)+\sum_{l=0}^{j-1} W_{j}(t) W_{j-l-1}(t)\right]\right], \quad j=2,3, \ldots,
\end{aligned}
$$

Here, we set $Z_{0}(0)=z_{0}(t)=1$ and $W_{0}(0)=w_{0}(t)=1$. Thus, the above equations lead the following results.

$$
Z_{0}(t)=1+\frac{t^{\beta}}{\Gamma(\beta+1)^{\prime}}
$$

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$$
\begin{gathered}
W_{0}(t)=1+\frac{t^{\beta}}{\Gamma(\beta+1)^{\prime}} \\
Z_{1}(t)=-3 \frac{t^{\beta}}{\Gamma(\beta+1)}+\left(\rho^{-1}-2\right) \frac{t^{2 \beta}}{\Gamma(2 \beta+1)}+\frac{\rho^{-1} \Gamma(2 \beta+1) t^{3 \beta}}{\Gamma^{2}(\beta+1) \Gamma(3 \beta+1)^{\prime}} \\
W_{1}(t)=-2 \frac{t^{\beta}}{\Gamma(\beta+1)}-2 \frac{t^{2 \beta}}{\Gamma(2 \beta+1)}-\frac{\Gamma(2 \beta+1) t^{3 \beta}}{\Gamma^{2}(\beta+1) \Gamma(3 \beta+1)^{\prime}} \\
Z_{2}(t)=\left(6-\rho^{-1}\right) \frac{t^{2 \beta}}{\Gamma(2 \beta+1)}-\left(\frac{1}{\rho^{2}}+4 \rho^{-1}+\frac{4 \rho^{-1} \Gamma(2 \beta+1)}{\Gamma^{2}(\beta+1)}-4\right) \frac{t^{3 \beta}}{\Gamma(3 \beta+1)} \\
-\left(\frac{1}{\rho^{2}}+4 \rho^{-1}+\frac{4 \rho^{-1} \Gamma(3 \beta+1) \Gamma(\beta+1)}{\Gamma^{2}(2 \beta+1)}\right) \frac{\Gamma(2 \beta+1) t^{4 \beta}}{\Gamma^{2}(\beta+1) \Gamma(4 \beta+1)} \\
-\frac{2 \rho^{-1} \Gamma(2 \beta+1) \Gamma(4 \beta+1) t^{5 \beta}}{\Gamma^{3}(\beta+1) \Gamma(3 \beta+1) \Gamma(5 \beta+1)^{\prime}} \\
W_{2}(t)= \\
3 \frac{t^{2 \beta}}{\Gamma(2 \beta+1)}+\left(\rho^{-1}+4+\frac{4 \Gamma(2 \beta+1)}{\Gamma^{2}(\beta+1)}\right) \frac{t^{3 \beta}}{\Gamma(3 \beta+1)} \\
+\left(2 \rho^{-1}+6+\frac{4 \Gamma(3 \beta+1)}{\Gamma^{2}(\beta+1) \Gamma(2 \beta+1)}\right) \frac{t^{4 \beta}}{\Gamma(4 \beta+1)} \\
+\frac{2 \Gamma(2 \beta+1) \Gamma(4 \beta+1) t^{5 \beta}}{\Gamma^{3}(\beta+1) \Gamma(3 \beta+1) \Gamma(5 \beta+1)^{\prime}}
\end{gathered}
$$

Thus, the solution of system (11) using the series (7) is

$$
\begin{gathered}
z(t)=Z_{0}(t)+Z_{1}(t)+Z_{2}(t)+\cdots, \\
w(t)=W_{0}(t)+W_{1}(t)+W_{2}(t)+\cdots, \\
z(t)=1-\frac{2 t^{\beta}}{\Gamma(\beta+1)}+\frac{4 t^{2 \beta}}{\Gamma(2 \beta+1)}+\frac{\rho^{-1} \Gamma(2 \beta+1) t^{3 \beta}}{\Gamma^{2}(\beta+1) \Gamma(3 \beta+1)} \\
-\left(\frac{1}{\rho^{2}}+4 \rho^{-1}+\frac{4 \rho^{-1} \Gamma(2 \beta+1)}{\Gamma^{2}(\beta+1)}-4\right) \frac{t^{3 \beta}}{\Gamma(3 \beta+1)} \\
-\left(\frac{1}{\rho^{2}}+4 \rho^{-1}+\frac{4 \rho^{-1} \Gamma(3 \beta+1) \Gamma(\beta+1)}{\Gamma^{2}(2 \beta+1)}\right) \frac{\Gamma(2 \beta+1) t^{4 \beta}}{\Gamma^{2}(\beta+1) \Gamma(4 \beta+1)} \\
-\frac{2 \rho^{-1} \Gamma(2 \beta+1) \Gamma(4 \beta+1) t^{5 \beta}}{\Gamma^{3}(\beta+1) \Gamma(3 \beta+1) \Gamma(5 \beta+1)}+\cdots, \\
w(t)=1-\frac{t^{\beta}}{\Gamma(\beta+1)}+\frac{t^{2 \beta}}{\Gamma(2 \beta+1)}-\frac{\Gamma(2 \beta+1) t^{3 \beta}}{\Gamma^{2}(\beta+1) \Gamma(3 \beta+1)} \\
+\left(\rho^{-1}+4+\frac{4 \Gamma(2 \beta+1)}{\Gamma^{2}(\beta+1)}\right) \frac{t^{3 \beta}}{\Gamma(3 \beta+1)} \\
\quad+\left(2 \rho^{-1}+6+\frac{4 \Gamma(3 \beta+1)}{\Gamma^{2}(\beta+1) \Gamma(2 \beta+1)}\right) \frac{t^{4 \beta}}{\Gamma(4 \beta+1)} \\
\\
\quad+\frac{2 \Gamma(2 \beta+1) \Gamma(4 \beta+1) t^{5 \beta}}{\Gamma^{3}(\beta+1) \Gamma(3 \beta+1) \Gamma(5 \beta+1)}+\cdots
\end{gathered}
$$

Here, setting $\beta=1$ and follow the above solution, the following results are obtained:

$$
\begin{array}{ll}
Z_{0}(t)=1+t, & W_{0}(t)=1+t \\
Z_{1}(t)=-3 t+\left(\rho^{-1}-2\right) \frac{t^{2}}{2!}+\frac{2 \rho^{-1} t^{3}}{3!}, & W_{1}(t)=-2 t-2 t^{2}-\frac{2 t^{3}}{3!}, \\
Z_{2}(t)=\left(6-\rho^{-1}\right) \frac{t^{2}}{2}-\left(\frac{1}{\rho^{2}}+4 \rho^{-1}+8 \rho^{-1}-4\right) \frac{t^{3}}{3!}-\left(\frac{2}{\rho^{2}}+20 \rho^{-1}\right) \frac{t^{4}}{4!}-\frac{2 \rho^{-1} t^{5}}{15}, \\
W_{2}(t)=3 \frac{t^{2}}{2!}+\left(\rho^{-1}+12\right) \frac{t^{3}}{3!}+\left(2 \rho^{-1}+\frac{3}{4}\right) \frac{t^{4 \beta}}{4!}+\frac{2 t^{5}}{15},
\end{array}
$$

One can express the above values in a series after finding the other terms of the solution

$$
\begin{gathered}
z(t)=1-2\left(t-t^{2}\right)-\frac{8}{6} t^{3}-\frac{4}{6} t^{4}+\cdots=\sum_{k=0}^{\infty} \frac{(2 t)^{k}(-1)^{k}}{k!}=e^{-2 t} \\
w(t)=1-t+\frac{t^{2}}{2}-\frac{t^{3}}{6}+\frac{t^{4}}{24}+\cdots=\sum_{k=0}^{\infty} \frac{(t)^{k}(-1)^{k}}{k!}=e^{-t}
\end{gathered}
$$

which is the exact solution of system (11).
Table 2. The Absolute Error of kernel Hilbert space method KHSM and EHPM of (8)

| $t$ | $z(t) \mathbf{E H P M}$ | $z(t) \mathbf{K H S M}$ | $w(t) \mathbf{E H P M}$ | $w(t)$ KHSM |
| :---: | :---: | :---: | :---: | :---: |
| 0.00 | 0 | 0 | 0 | $1.2 \times 10^{-6}$ |
| 0.4 | $5.55 \times 10^{-17}$ | $1.20 \times 10^{-6}$ | 0 | $2.47 \times 10^{-2}$ |
| 0.8 | $5.55 \times 10^{-16}$ | $1.28 \times 10^{-6}$ | 0 | $1.19 \times 10^{-1}$ |
| 1.2 | $1.70 \times 10^{-12}$ | $9.10 \times 10^{-7}$ | $5.55 \times 10^{-17}$ | $2.62 \times 10^{-1}$ |
| 1.6 | $6.90 \times 10^{-10}$ | $5.59 \times 10^{-7}$ | 0 | $4.38 \times 10^{-1}$ |
| 2.0 | $7.30 \times 10^{-8}$ | $3.18 \times 10^{-7}$ | $5.56 \times 10^{-16}$ | $1.2 \times 10^{-6}$ |

Table 2: illustrates the comparison between the absolute errors of the numerical solution of the proposed technique at different values of $t$ where $\beta=1$ for 20 terms iterations and the solution obtained by KHSM. It is clear that the EHPM technique is more accurate and converges faster than the KHSM technique.

## 5. Conclusions

In this work, a new computational technique namely Elzaki homotopy perturbation method are constructed to solve linear and nonlinear systems of differential equation of fractional order. The construction scheme of the current method is discussed and implemented to show the performance of the computational method for this problem. The results obtained showed that the proposed method is a powerful and efficient method for solving linear and nonlinear stiff systems of fractional orders and compared the obtained results of the current method with the results obtained by the other methods. Table 1 and table 2 observed that the suggested technique is superior to kernel Hilbert space method KHSM.

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