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Asymmetric impacts on Mars' polar vortices from an equinoctial Global Dust Storm

Paul M. Streeter¹, Stephen R. Lewis¹, Manish R. Patel^{1,2}, James A. Holmes¹,
Anna A. Fedorova³, David. M Kass⁴, and Armin Kleinböhl⁴.

¹School of Physical Sciences, The Open University, Walton Hall, Milton Keynes, UK

²Space Science and Technology Department, Science and Technology Facilities Council, Rutherford

Appleton Laboratory, Oxfordshire, UK

³Space Research Institute of the Russian Academy of Sciences (IKI RAS), Moscow, Russia

⁴Jet Propulsion Laboratory, California Institute of Technology, Pasadena, California, USA

Key Points:

- The 2018 Global Dust Storm had significant and asymmetrical impacts on the morphology and intensity of Mars' two polar vortices
- The southern vortex was substantially diminished while the northern vortex remained robust, but both were reduced in ellipticity
- These vortex changes enhance transport into Mars' southern (but not northern) polar region and modify longitudinal transport patterns

Corresponding author: Paul Streeter, paul.streeter@open.ac.uk

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Abstract

Mars possesses dynamical features called polar vortices: regions of cold, isolated air over the poles circumscribed by powerful westerly jets which can act as barriers to transport to dust, water, and chemical species. The 2018 Global Dust Storm was observed by multiple orbiters and offers a valuable opportunity to study the effects of such a storm on polar dynamics. To this end, we assimilate data from the Mars Climate Sounder and Atmospheric Chemistry Suite into a Mars Global Climate Model. We find that the storm had asymmetrical hemispherical impacts, with the northern vortex remaining relatively robust while the southern vortex was substantially diminished in its intensity. We propose that this asymmetry was due both to the storm's latitudinal extent, as it extended further south than north, and to its equinoctial timing, occurring as the southern vortex was already decaying. We show that both polar vortices, in particular the northern, were reduced in ellipticity by the storm. There was a well-correlated reduction in stationary topographic wave activity at high latitudes in both hemispheres. We demonstrate that the characteristic elliptical martian polar vortex shape is the pattern of the stationary waves, which was suppressed by the shifting of the polar jet away from regions of high mechanical forcing (north) or reduction of polar jet intensity by a reduced meridional temperature gradient (south). These asymmetric effects suggest increased transport into the southern (but not northern) polar region during Global Dust Storms at northern autumn equinox, and more longitudinally symmetric transport around both poles.

Plain Language Summary

Like Earth, Mars has regions of cold air around its winter poles. The temperature contrast creates powerful polar jets, polar vortices, which can block transport of atmospheric aerosols and chemicals. Unlike Earth, Mars regularly experiences Global Dust Storms which have enormous effects on atmospheric temperatures and winds. The most recent storm occurred at Mars' equinox and was observed by multiple spacecraft. We combined these observations with a numerical model. We find that the northern vortex remained relatively strong and coherent, while the southern was greatly disrupted, showing atmospheric warming and a diminished polar jet. This was because of the greater southward extent of the storm and its seasonal timing. Both vortices normally show a distinct elliptical shape; the storm made both vortices more longitudinally symmetrical. We link this to a corresponding decrease in the amplitude of topographic planetary-scale waves, showing that the elliptical shape is that of the planetary wave structure. These results suggest that equinoctial storms may enhance transport into the southern pole due to the diminished vortex, while the more robust northern vortex continues to act as an effective barrier. The reduced ellipticity of both vortices may also lead to more longitudinally symmetric transport into the polar regions.

1 Introduction

Mars' winter atmosphere is characterized by a polar vortex of low temperatures around the winter pole, circumscribed by a strong westerly jet (Waugh et al., 2016; Mitchell et al., 2015). These vortices are a key part of the atmospheric circulation and are heavily involved with dust and volatile transport (e.g. J. A. Holmes et al., 2017). Planetary polar vortices are a common feature of atmospheres in the Solar System, but Mars' differ from Earth's in several important respects. Among the most notable and visible is their peculiar annular structure. On Earth, the potential vorticity (a way of diagnosing the presence and strength of the polar vortex; see discussion below) of the polar vortices increases monotonically towards the pole; on Mars, there is a distinctive ring of higher potential vorticity around the pole, then a minimum over the pole itself (Waugh et al., 2016; Mitchell et al., 2015). This annular potential vorticity structure should be barotropically unstable but appears to persist over seasonal timescales; modelling indicates that Mars'

67 low radiative relaxation timescales can help maintain this equilibrium (Seviour et al., 2017).
68 The current best explanation for the annular structure itself appears to be diabatic heat-
69 ing from CO₂ condensation over the winter pole; as the CO₂ condenses, it releases la-
70 tent heat energy, warming the lower atmosphere and causing a local reduction in poten-
71 tial vorticity (Toigo et al., 2017; Rostami et al., 2018; Scott et al., 2020). Mars’ polar
72 vortices also show a hemispheric asymmetry, with the northern vortex being stronger in
73 reanalyses than the southern (Mitchell et al., 2015).

74 Another feature of the martian polar vortices is their elliptical shape, particularly
75 in the northern hemisphere (Waugh et al., 2016). It has been speculated by Mitchell et
76 al. (2015) and Rostami et al. (2018) that this ellipticity could be linked to topography,
77 something not incorporated into their simplified model of the martian circulation. This
78 elliptical shape is only visible when averaged over time periods of 10s of sols; over smaller
79 timescales, the polar vortex structure is less coherent and composed of smaller regions
80 of high potential vorticity (Waugh et al., 2016). Rostami et al. (2018) attributed this to
81 inhomogeneous deposition of condensing CO₂ ice. Meanwhile, Mitchell et al. (2015) found
82 that the martian polar vortices are consistently centred over the pole itself at the sol-
83 stices, suggesting a relatively lesser (compared to Earth) role for wave-mean flow inter-
84 actions in controlling the shape of the polar vortices.

85 The martian polar vortices appear to have a complex relationship with atmospheric
86 dust loading. Mitchell et al. (2015) found that in the MACDA reanalysis, there was less
87 seasonal variability in the polar vortex structure due to planetary Rossby wave activ-
88 ity (and resulting sudden stratospheric warming), as there is on Earth; rather, any vari-
89 ability was linked to dust-induced changes to the Hadley circulation (and resulting in-
90 tensified polar warming). Specifically, they investigated the effect of a regional dust storm
91 at $L_S=320^\circ$ in MY 26 (a “C”-type storm; see Kass et al. (2016)), and found that it acted
92 to shift the northern vortex towards the equator by $\sim 10^\circ$ in latitude and weaken the vor-
93 tex circulation overall. Guzewich et al. (2016) used an MGCM with an analytically pre-
94 scribed dust scenario to investigate the effects of high southern hemisphere dust load-
95 ing on the northern polar vortex, and vice-versa. They found that regional and Global
96 Dust Storm (GDS) events could produce sudden transient vortex warming, disrupting
97 the northern polar vortex for periods of up to 10s of sols, by shifting the downwelling
98 branch of the cross-equatorial Hadley cell poleward. By contrast, the southern polar vor-
99 tex was significantly more robust to high northern hemisphere dust loading. The exact
100 relationship between the polar vortices and atmospheric dust content is an important
101 one to understand, as it has implications for the transport of both volatiles and dust it-
102 self through the vortices (e.g. I. B. Smith et al., 2017; McCleese et al., 2017).

103 Previous studies have used reanalyses to investigate Mars’ polar vortices, but to
104 date have only used those which assimilate Thermal Emission Spectrometer (TES) data,
105 namely the MACDA (Mitchell et al., 2015; Waugh et al., 2016) and EMARS (Waugh
106 et al., 2016) reanalyses. This study assimilates Mars Climate Sounder (MCS) column
107 dust optical depth (CDOD) products and MCS and Atmospheric Chemistry Suite (ACS)
108 temperature retrievals. MCS CDOD products are not limited, as TES CDOD nadir mea-
109 surements were, to areas with relatively warm surface temperatures ($>220\text{ K}$) (M. D. Smith,
110 2004), allowing greater coverage over the seasonal CO₂ caps and therefore more CDOD
111 data for assimilation over these regions. This study also uses the newer 2D MCS retrievals,
112 one goal of which are improved retrievals over the polar regions specifically (Kleinböhl
113 et al., 2017). ACS temperature retrievals have the advantage of observing over differ-
114 ent martian local times.

115 Crucially for the understanding of how large-scale dust loading affects Mars’ po-
116 lar dynamics, this study assimilates CDOD and temperature data from the real GDS that
117 occurred in MY 34. The MY 34 GDS was first detected at $L_S=186^\circ$ and matured into
118 a global, planet-encircling event by $L_S=200^\circ$, before beginning its long decay phase around
119 $L_S=213^\circ$ and returning to climatological levels around $L_S=270^\circ$ (Kass et al., 2019a). At

120 the height of the GDS, the planet-encircling dust cloud extended to approximately 45° N
121 and 70° S, and the greatest atmospheric temperature response (at 50 Pa) was seen in the
122 southern hemisphere, from the equator to the southern pole itself (Kass et al., 2019a).
123 GDS in general have been found to have significant impacts on martian atmospheric dy-
124 namics (e.g. Fedorova et al., 2020; McDunn et al., 2013; Guzewich et al., 2014). The MY
125 34 event was observed by an unprecedented number of spacecraft, including the Mars
126 Reconnaissance Orbiter (MRO) and the ExoMars Trace Gas Orbiter (TGO), making it
127 an ideal candidate for using data assimilation to investigate its impacts on polar dynam-
128 ics. Previously, Guzewich et al. (2016) used a prescribed dust scenario to investigate the
129 effects of a GDS-like event at solstice. The use of a reanalysis allows the study of a re-
130 alistic GDS-level dust loading closer to equinox, affording a chance to study GDS effects
131 at a different season. This also allows comparison with a recent study on the MY 34 GDS
132 using MCS data directly, and its findings of significant diurnal variation in southern po-
133 lar vortex structure (Kleinböhl et al., 2020). Finally, this study devotes time to inves-
134 tigate wave-mean flow interactions, and specifically study how the horizontal structure
135 of the polar vortices at this season might be affected by such interactions.

136 2 Methods

137 2.1 Model

138 The model used for this study is a Mars Global Climate Model (MGCM), a four-
139 dimensional numerical model which exists as a collaborative effort between the Labo-
140 ratoire de Météorologie Dynamique, the University of Oxford, the Open University, and
141 the Instituto de Astrofísica de Andalucía (Forget et al., 1999). This version of the MGCM
142 uses a spectral dynamical core to solve the equations of fluid motion, with a finite-difference
143 scheme in the vertical dimension and a semi-Lagrangian scheme for tracer advection (Lewis
144 et al., 2007). The MGCM advects dust by use of a two-moment scheme with a log-normal
145 size distribution, and total CDOD is scaled at each column to match assimilated obser-
146 vations (Madeleine et al., 2011; Streeter, Lewis, Patel, Holmes, & Kass, 2020). The dust
147 distribution in the vertical is allowed to evolve without constraint. Dust in the MGCM
148 is radiatively active, using radiative properties derived from observations (Wolff et al.,
149 2006, 2009). The MGCM radiative transfer scheme is reliable to within ~10% error even
150 at the very high dust loadings observed during the 2018 GDS (Toon et al., 1989; Streeter,
151 Lewis, Patel, Holmes, & Kass, 2020).

152 2.2 Retrievals and data assimilation

153 The MGCM was run with a modified version of the Analysis Correction data as-
154 similation scheme (Lorenz et al., 1991), tuned for use on the martian atmosphere (Lewis
155 et al., 1997, 2007). Orbitally retrieved temperature profiles were assimilated using the
156 method previously used in this scheme for TES (Lewis et al., 2007; J. A. Holmes et al.,
157 2018) and MCS (J. A. Holmes et al., 2019; Steele et al., 2014) data, while CDOD de-
158 rived from MCS limb dust profiles was assimilated to constrain MGCM dust columns
159 (Lewis et al., 2007). The assimilation scheme and MGCM were the same as those used
160 for the OpenMARS reanalysis dataset, which currently extends to MY 32 (J. Holmes et
161 al., 2019).

162 The retrieved temperature profiles used were from MCS (McCleese et al., 2007) aboard
163 MRO (Zurek & Smrekar, 2007), and from ACS (O. Korablev et al., 2018) aboard TGO.
164 MCS temperature profiles extend to ~85 km, with an intrinsic vertical resolution of ~5 km
165 (Kleinböhl et al., 2009). The sun-synchronous orbit of the Mars Reconnaissance Orbiter
166 results in two approximately fixed local times for MCS observations, namely 0300 and
167 1500 at nonpolar latitudes (Zurek & Smrekar, 2007). The retrieval version used was v5.2,
168 the latest version which incorporates two-dimensional radiative transfer to correct for
169 lateral gradients in temperature and aerosol, resulting in improved retrievals over the

170 poles (Kleinböhl et al., 2017). The exception was during the period of the 2018 GDS,
171 for which a reprocessed version (v5.3.2) was used; this reprocessed version incorporated
172 information from extra channels on MCS (Kleinböhl et al., 2020). ACS temperatures for
173 MY 34 were also assimilated, in the form of temperature profile retrievals from solar oc-
174 culations by the NIR (near-infrared) channel, with an intrinsic vertical resolution of 1-
175 3 km and altitude range of 0-90 km depending on the TGO orbit (Fedorova et al., 2020).
176 NIR performed on average nine occultations per martian Sol throughout the relevant pe-
177 riod of MY 34 (Fedorova et al., 2020). TGO’s non-sun-synchronous orbit means that ACS
178 temperatures were available at sunrise/sunset local times, near the terminator.

179 Assimilated CDOD data was from MCS, which does not directly measure dust columns
180 but provides a derived column product based on extrapolation of retrieved dust profiles.
181 As stated above, however, while the CDOD in the MGCM was given by assimilated MCS
182 CDOD, the vertical distribution of dust in the MGCM was allowed to evolve freely with-
183 out prescription. As MCS CDOD is reported at infrared wavelengths while the MGCM
184 uses visible wavelengths for dust radiative transfer calculations, CDOD values were first
185 converted from $21.6\ \mu\text{m}$ to 670 nm via a conversion factor of 7.3 (Kleinböhl et al., 2011).
186 Dayside equatorial CDOD values were filtered out before assimilation to avoid spuriously
187 high values (Montabone et al., 2015), except during the period of the GDS itself (Montabone
188 et al., 2020). As with MCS temperatures, v5.2 retrievals were used except during the GDS
189 period, when v5.3.2 retrievals were used.

190 2.3 Simulations performed

191 A single MGCM simulation utilising data assimilation (“reanalysis”) was performed
192 covering multiple Mars Years, which included the periods $L_S=200-220^\circ$ for both MY 33
193 and MY 34. These periods were focussed on for the purposes of this article. The reanal-
194 ysis assimilated MCS and ACS temperature profiles and MCS CDOD products. The MGCM
195 was run at a spatial spectral resolution T42, corresponding to a spatial resolution of $\sim 3.75^\circ$
196 (~ 215 km at the equator), with 50 topography-following vertical levels with midpoints
197 from at shallowest ~ 5 m and at deepest ~ 105 km above the surface. The MGCM was
198 run without water cycle parametrizations in order to isolate the effects of dust.

199 For the purposes of this article, MY 33 was chosen as a non-GDS year to compare
200 against MY 34. MY 33 was a very typical MCS year in terms of dust loading, includ-
201 ing in the timing and magnitude of its regional “A”-, “B”-, and “C”-storms. The very
202 average dust loading of MY 33 allows comparison between a situation of “normal” dust
203 loading and a situation of GDS-scale dust loading. The $L_S=200-220^\circ$ period was focussed
204 on as it represents the height of global, homogeneous dust loading during the GDS pe-
205 riod in MY 34 (Kass et al., 2019b).

206 2.4 Potential vorticity diagnostic

207 A diagnostic used frequently throughout this article is Ertel potential vorticity (PV).
208 PV is a measure of air circulation derived from the vorticity and stratification of the at-
209 mosphere, and is valuable for being conserved like a material tracer under adiabatic pro-
210 cesses (Haynes & McIntyre, 1987), making it especially useful for the study of polar dy-
211 namics: the polar vortices can be defined as regions of high PV around the poles. PV
212 has both dynamical components, in the vorticity of both the air mass itself and the planet,
213 and thermodynamic elements, in the form of the potential temperature structure and
214 static stability of the atmosphere. The PV of an air mass on an isentropic surface is con-
215 served, and cannot be created, destroyed, or transported across isentropic surfaces (Haynes
216 & McIntyre, 1987); therefore a large-scale local reduction in PV implies significant mix-
217 ing along the isentropic surface, associated with diabatic and/or frictional processes. PV
218 can be defined as

$$PV = -g(\xi + f) \frac{\partial \theta}{\partial p} \quad (1)$$

where g is the gravitational acceleration (3.72 m/s^2 on Mars), ξ is the relative isentropic vorticity (the relative vorticity of the air mass on that particular isentropic surface, an isentropic surface being a surface of constant potential temperature), f is the Coriolis parameter (the vorticity associated with the planetary rotation at a particular latitude), θ is the potential temperature, and p is the pressure.

PV is given as a value on a particular isentropic surface; this study uses the 300 K isentropic surface for consistency with previous studies of the martian polar atmosphere (e.g. Mitchell et al., 2015; Waugh et al., 2016). This corresponds to an approximate altitude range of 20-30 km, and when winds are presented in this article they are integrated between 20-30 km.

PV is typically positive/negative in the northern/southern hemisphere, and increases in magnitude near the poles due to the value of f . The term “magnitude” is used throughout this article for PV values, to make it clear that a larger negative PV value means a greater absolute value of PV. For simplicity, 1 “MPVU” (Mars potential vorticity unit) is defined throughout this article as $1 \times 10^{-4} \text{ K m}^2 \text{ kg}^{-1} \text{ s}^{-1}$, or 100 PVU (a standard unit used for terrestrial studies).

3 Results

3.1 Diurnally averaged changes

This section explores the effects of the MY 34 GDS over both the north and south poles (“NP” and “SP”) as averaged over all local times.

Fig. 1 shows the average structure of the NP and SP polar vortices between $L_S=200\text{-}220^\circ$ for MYs 33 and 34, and the difference between them. NP PV saw an overall reduction (Fig. 1.e) up to 15 MPVU, with the greatest reduction around latitudes $60\text{-}70^\circ \text{ N}$, though PV actually increased slightly over the pole itself. The reduction was highly longitudinally asymmetric, with maxima in PV reduction in the eastern hemisphere at $180\text{-}90^\circ \text{ E}$ and $-30\text{-}30^\circ \text{ E}$. Despite this, the gross morphology of the NP vortex (Fig. 1.a,c) showed minor changes compared to the SP. The MY 34 vortex showed reduced ellipticity compared to MY 33, and the disappearance of a local PV minimum over the pole itself.

NP zonal winds also changed, tending to increase in the GDS case up to 20 m/s (Fig. 1.e) north of 60° N , particularly around $90\text{-}180^\circ \text{ E}$, but decreasing southward of 60° N . The wind speed increases align with the PV decreases, occurring primarily in the east. In the western hemisphere there was little change except between latitudes $70\text{-}80^\circ \text{ N}$, where zonal winds decreased by around 8 m/s between $-120\text{-}30^\circ \text{ E}$. At the lower latitudes $50\text{-}60^\circ \text{ N}$, in (primarily) the eastern hemisphere, zonal wind speeds decreased up to 16 m/s , showing a poleward jet shift. The GDS-induced changes were asymmetric and made the MY 34 wind structure less elliptical than in MY 33. A strong local wind maximum up to 120 m/s in the eastern hemisphere (55° N , 60° E) in MY 33 was substantially reduced in MY 34, creating a more longitudinally symmetric jet. Despite the generally more symmetrical MY 34 wind structure, the jet appears shifted off-pole towards -30° E . Finally, the closer clustering of contours in MY 34 shows a latitudinal narrowing of the westerly jet. There was a visible anticorrelation between PV and zonal wind speeds.

SP PV also showed overall absolute reduction, up to 5 MPVU. The pattern of decrease (Fig. 1.f) correlates exactly with the MY 33 polar vortex structure (Fig. 1.b). The annular PV structure in MY 33 is both longitudinally asymmetric, weaker in the west/stronger in the east, and centred off-pole, with the central PV minimum around 80° S , -30° E . The MY 34 PV structure is far more uniform; there is still a (weaker) annular pattern, cen-

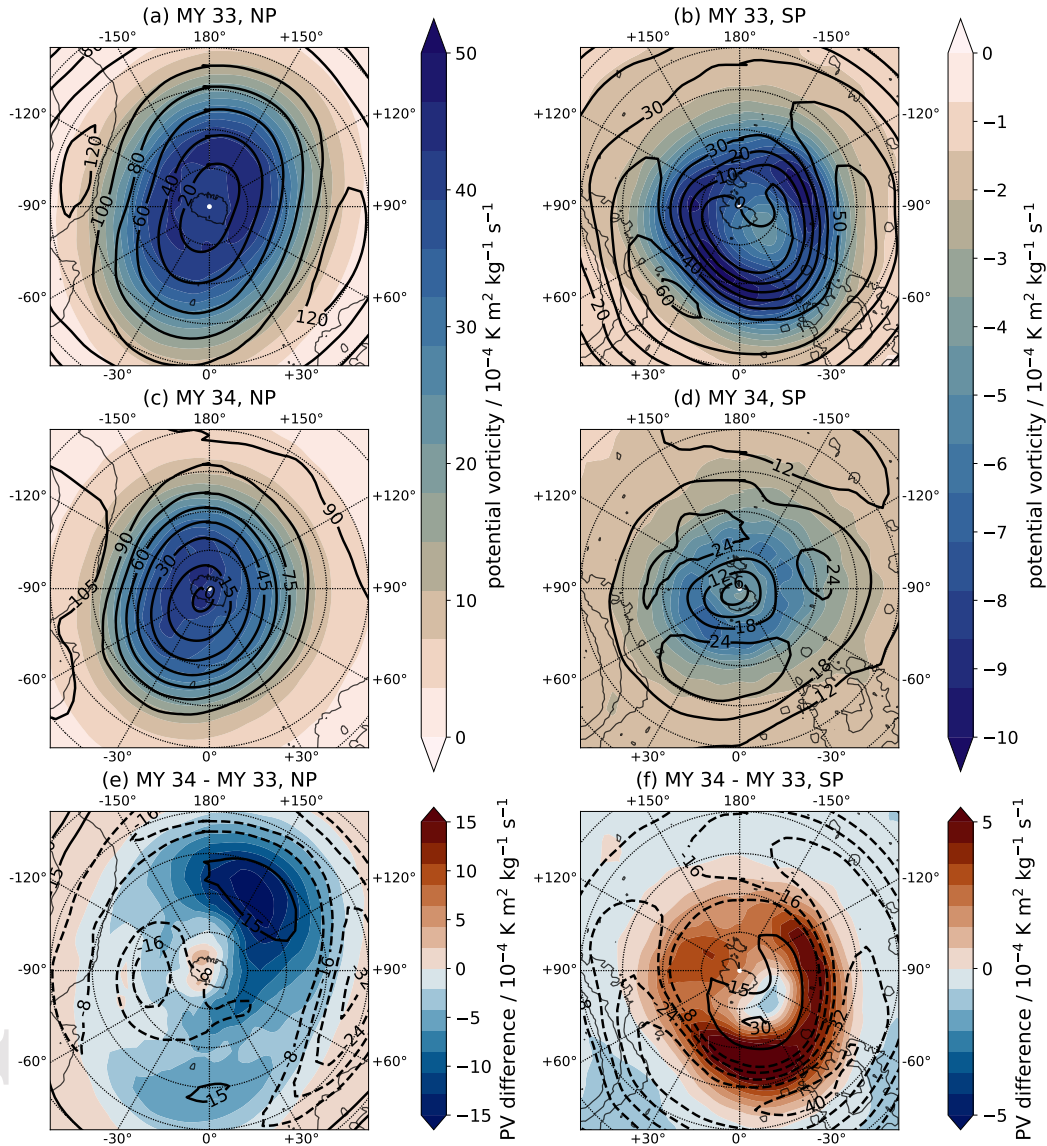


Figure 1. Potential vorticity (colours) as averaged over the $L_S=200-220^\circ$ period on the 300 K isentropic level (located approximately between 20-30 km altitude) and zonal wind speeds (contours) as averaged between 20-30 km for (top) MY 33, (middle) MY 34, and (bottom) the difference between MY 34 - MY 33, over the (left) north pole and (right) south pole. Plots are stereographic projections where each latitude circle is 10° separate from its neighbours and the innermost circle represents the 80° latitude band. Lower absolute values indicate lower magnitude of potential vorticity, and vice-versa. Note the difference in scales and colours. Thinner black contour lines represent surface topography.

265 tred over the pole itself, and reduced longitudinal asymmetry. The proportional decrease
 266 of PV at the SP (up to 50%) was much greater than at the NP, resulting in a drastically
 267 altered morphology and substantially reduced PV across the entire vortex, suggesting
 268 large-scale PV mixing from dynamical and/or diabatic changes. PV mixing is also im-
 269 plied by the fact that absolute PV increased slightly equatorward of 60° S.

270 SP zonal winds generally increased where PV decreased, up to 30 m/s, and decreased
 271 elsewhere (outside the vortex) by similar amounts (Fig. 1.f). The greatest wind increases
 272 were at the MY 33 local PV minimum (Fig. 1.b); this calm “eye” had very low wind speeds
 273 in MY 33. This signifies a shift of the remnant vortex from ~5° off-pole to over the pole
 274 itself. Zonal wind gradients were weaker in MY 34, showing a less coherent jet core. The
 275 decrease in zonal wind speeds equatorward of 60° S was likely due to increased dust load-
 276 ing reducing the meridional temperature gradient. The MY 34 wind structure was more
 277 longitudinally symmetric than in MY 33, and centred around the pole itself. Again, there
 278 was a visible anticorrelation between PV and zonal wind speeds.

279 The MY 34 GDS significantly boosted the mean meridional circulation (MMC),
 280 strengthening both the dominant cross-equatorial clockwise and the southern anticlock-
 281 wise Hadley cells (Fig. 2.I,II,III). A stronger Hadley cell signifies greater transport of
 282 (warmer) air from lower towards higher latitudes, causing adiabatic heating. The MY
 283 34 cross-equatorial Hadley cell also extended further poleward, indicating transport to-
 284 wards higher northern latitudes than under non-GDS conditions; from ~60° N in MY 33
 285 to ~75° N in MY 34. This was not seen in the southern anticlockwise Hadley cell.

286 In addition to thermally direct circulatory cells, there is frequently a thermally in-
 287 direct cell at mid-high northern latitudes. Such features at mid-latitudes on Earth are
 288 called “Ferrel cells”, but are not technically real circulatory cells; rather, they indicate
 289 the presence of mechanical forcing from planetary wave activity due to thermal contrasts
 290 at mid-latitudes (Salby, 2011), and their presence in the Eulerian mean meridional cir-
 291 culation is an artefact of the averaging process (see Andrews et al., 1987, Chapter 3).
 292 Given the thermal contrast at northern mid-latitudes at this time of year ($L_S=200-220^\circ$,
 293 early northern winter) the presence of these cells implies baroclinic planetary wave ac-
 294 tivity. The weakening of the northern thermally indirect cell in MY 34 (Fig. 2) indicates
 295 that the GDS suppressed baroclinic wave activity over Mars’ northern lowlands, where
 296 such activity is generally strongest (e.g. Barnes et al., 1993).

297 The meridional flow was altered at different longitudes, reflecting the increased lon-
 298 gitudinal symmetry in the MY 34 polar vortex structure. Fig. 2 shows meridional winds
 299 as calculated for six different longitude ranges, at increments of 60°, which serve as a proxy
 300 for the meridional circulation at these longitude ranges. There is significant variation in
 301 the meridional flow at different longitudes.

302 In MY 33, the seasonally typical large-scale flow of the MMC is replicated at some
 303 longitudes: -180-120° E and 0-60° E, where the northward meridional flow extends to the
 304 NP. However, northern hemisphere southward (away from pole) flow occurs at longitudes
 305 -120-0° E and 60-180° E. These correspond to northern topographic depressions (Fig. S1)
 306 where baroclinic wave activity is strongest (e.g. Barnes et al., 1993). This pattern sug-
 307 gests a wavenumber 2 stationary wave. In the south, at 0-180° E (eastern hemisphere),
 308 the southward flow extends to the SP while in the western hemisphere there is north-
 309 ward flow, indicating wavenumber 1 stationary wave activity. These flow patterns cor-
 310 relate with the MY 33 polar vortex structure (Fig. 1.a,b). In both hemispheres, pole-
 311 ward flow correlates with where the vortex is compressed poleward, while equatorward
 312 flow correlates with where the vortex extends further equatorward.

313 The equivalent longitude ranges in MY 34 show complex changes, but some broad
 314 patterns are identifiable. In general, the flows show greater zonal symmetry than in MY
 315 33, resembling the MMC more closely, especially at southern mid-high latitudes. In the

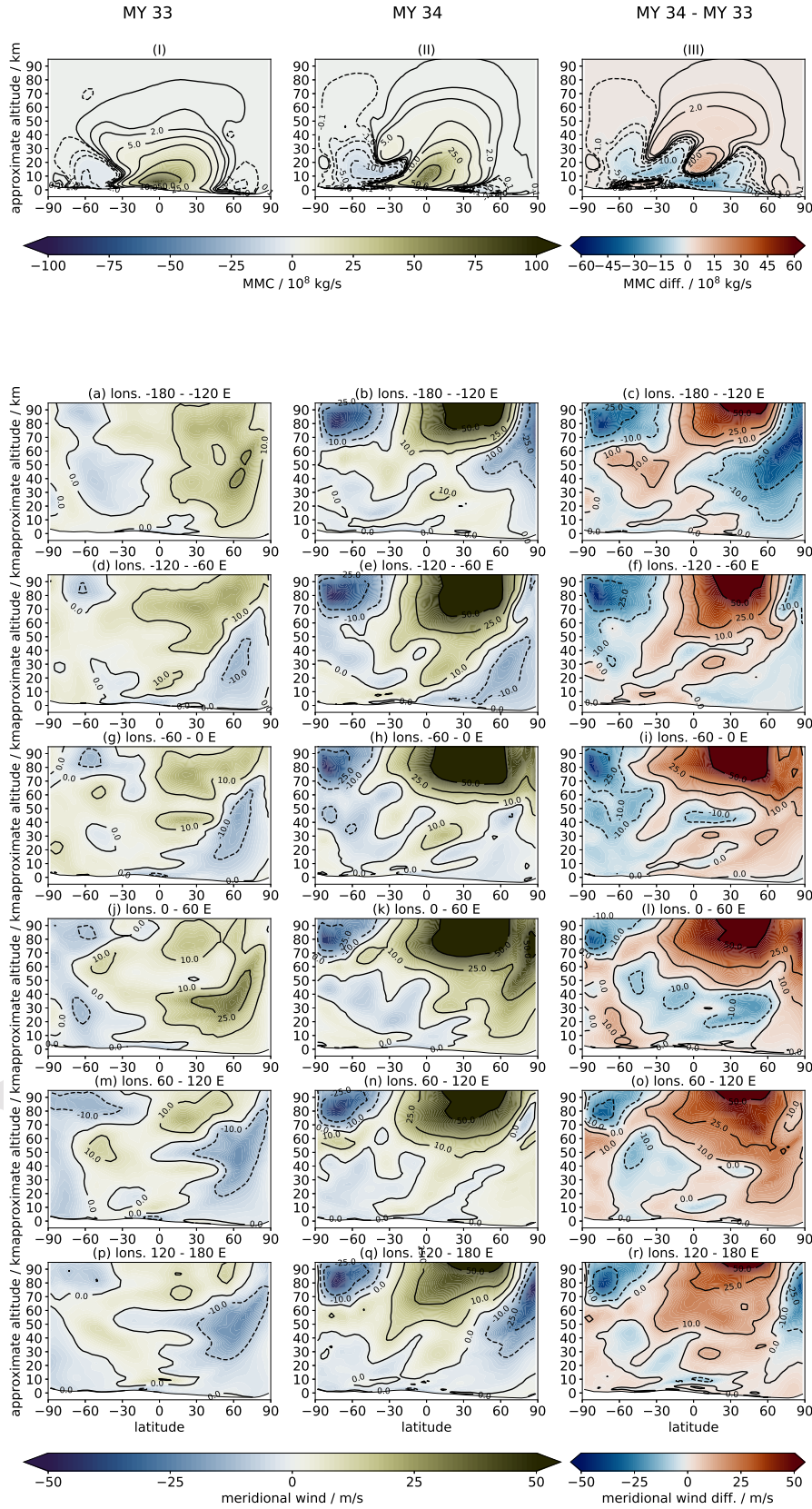


Figure 2. Top row: Mean meridional circulation calculated for the $L_S=200-220^\circ$ period. Positive/negative values indicate clockwise/anticlockwise flow. All other rows: meridional wind averaged over different longitude ranges for the same period. Positive/negative values indicate northward/southward flow.

316 north, the wavenumber 2-like pattern in MY 33 changes to a more dominant wavenum-
 317 ber 1-like pattern, with roughly half the planet showing northward meridional flow and
 318 half showing southward meridional flow. This generally more longitudinally symmetri-
 319 cal flow matches the more longitudinally symmetrical NP and SP polar vortices (Fig.
 320 1.c,d). In the eastern hemisphere where there is southward flow extending to the SP in
 321 MY 33, in MY 34 this is significantly weakened. Likewise, in the west where there is flow
 322 away from the SP in MY 33, in MY 34 this is diminished and/or reversed. This altered
 323 flow matches the highly symmetric MY 34 SP vortex (Fig. 1.d). In the north, the MY
 324 34 meridional flow pattern generally shows northward flow towards the NP between lon-
 325 gitudes $-60-120^\circ\text{E}$, and southward flow away from the NP between longitudes $120-60^\circ\text{E}$.
 326 This GDS flow pattern therefore implies an MY 34 NP vortex which is compressed pole-
 327 ward between $-60-120^\circ\text{E}$, but extends further equatorward between $120-60^\circ\text{E}$. Exam-
 328 ination of Fig. 1.c shows that the highest PV values can be found on the $120-60^\circ\text{E}$ side
 329 of the planet; for example, the same ~ 35 MPVU contour which extends to 70°N at -150°E
 330 extends to only 80°N at 30°E .

331 It has been shown that the longitudinal asymmetries in the polar vortices are relat-
 332 ed to longitudinally asymmetric meridional wind patterns; but what is the cause of
 333 this longitudinal asymmetry? Fig. 3 shows the meridional wind deviation integrated be-
 334 tween altitudes 20-30 km. The meridional wind deviation is defined as the difference be-
 335 tween the time mean (in this case, between $L_S=200-220^\circ$) of meridional wind and the
 336 time and zonal mean of meridional wind ($[v]-[\bar{v}]$, where the brackets and overbar rep-
 337 resent time and zonal means respectively), and indicates the presence of stationary plan-
 338 etary waves. Fig. 3.a shows the stationary wave pattern in MY 33. A clear spatial wavenum-
 339 ber 2 signal is visible at northern mid-high latitudes due to the zonal topographic dif-
 340 ferences present at northern mid-high latitudes, with two notable depressions at Acidalia
 341 and Utopia plains and higher topography around Alba Patera and northern Arabia Terra
 342 (see Fig. S1), which induce differential heating and thus drive adjacent clockwise and
 343 anticlockwise circulations (Nayvelt et al., 1997; Hollingsworth & Barnes, 1996; Haberle
 344 et al., 2019). The presence of this wavenumber 2 feature and associated wind directions
 345 has been seen in observations of NP hood clouds (Haberle et al., 2019). There is a north-
 346 ward flow between approximately $-10-70^\circ\text{E}$ and $160-120^\circ\text{E}$, and a southward flow be-
 347 tween approximately $-110-10^\circ\text{E}$ and $70-160^\circ\text{E}$. There is also stationary wave activity
 348 at southern high latitudes indicating a wavenumber 2 feature at mid latitudes which tran-
 349 sitions into a wavenumber 1 feature poleward of $75-80^\circ\text{S}$; other modelling work suggests
 350 that this becomes a more unambiguous wavenumber 1 at southern winter (Hollingsworth
 351 & Barnes, 1996). There is a northward flow between approximately $-70-10^\circ\text{E}$, and a south-
 352 ward flow between approximately $90-180^\circ\text{E}$, both extending all the way to the pole. There
 353 are also flows at southern mid-latitudes: a north-to-south flow between approximately
 354 $-20-60^\circ\text{E}$, and a less well-defined south-to-north flow between approximately $60-120^\circ\text{E}$.

355 The MY 34 GDS drastically changed stationary wave structures (Fig. 3.b,c.); north-
 356 ern amplitudes were decreased up to 20 m/s. There remained two significant remnants:
 357 southward flow around $-110-10^\circ\text{E}$, and adjacent northward flow around $-10-70^\circ\text{E}$. These
 358 remnant flows were also constrained to poleward of $\sim 45^\circ\text{N}$. The other two flows were
 359 almost completely destroyed, creating a wavenumber 1-like pattern. In the south, the
 360 mid-high latitude wavenumber 2 structure was completely destroyed, leaving only a weak-
 361 ened wavenumber 1 feature consisting of weak northward flow between $-180-60^\circ\text{E}$ and
 362 weak southward flow between $60-180^\circ\text{E}$.

363 These stationary wave structures and changes match the polar vortex structures
 364 and changes. In the north, the MY 33 wavenumber 2 pattern matches the longitudinal
 365 asymmetry of the elliptical polar vortex (Fig. 1.a). The elongated parts of the ellipse are
 366 where the meridional wind deviation consists of southward flow, roughly between $-110-$
 367 -10°E and $70-160^\circ\text{E}$. The narrow parts of the ellipse are where the meridional wind de-
 368 viation consists of northward flow, roughly between $-10-70^\circ\text{E}$ and $170-110^\circ\text{E}$. The MY

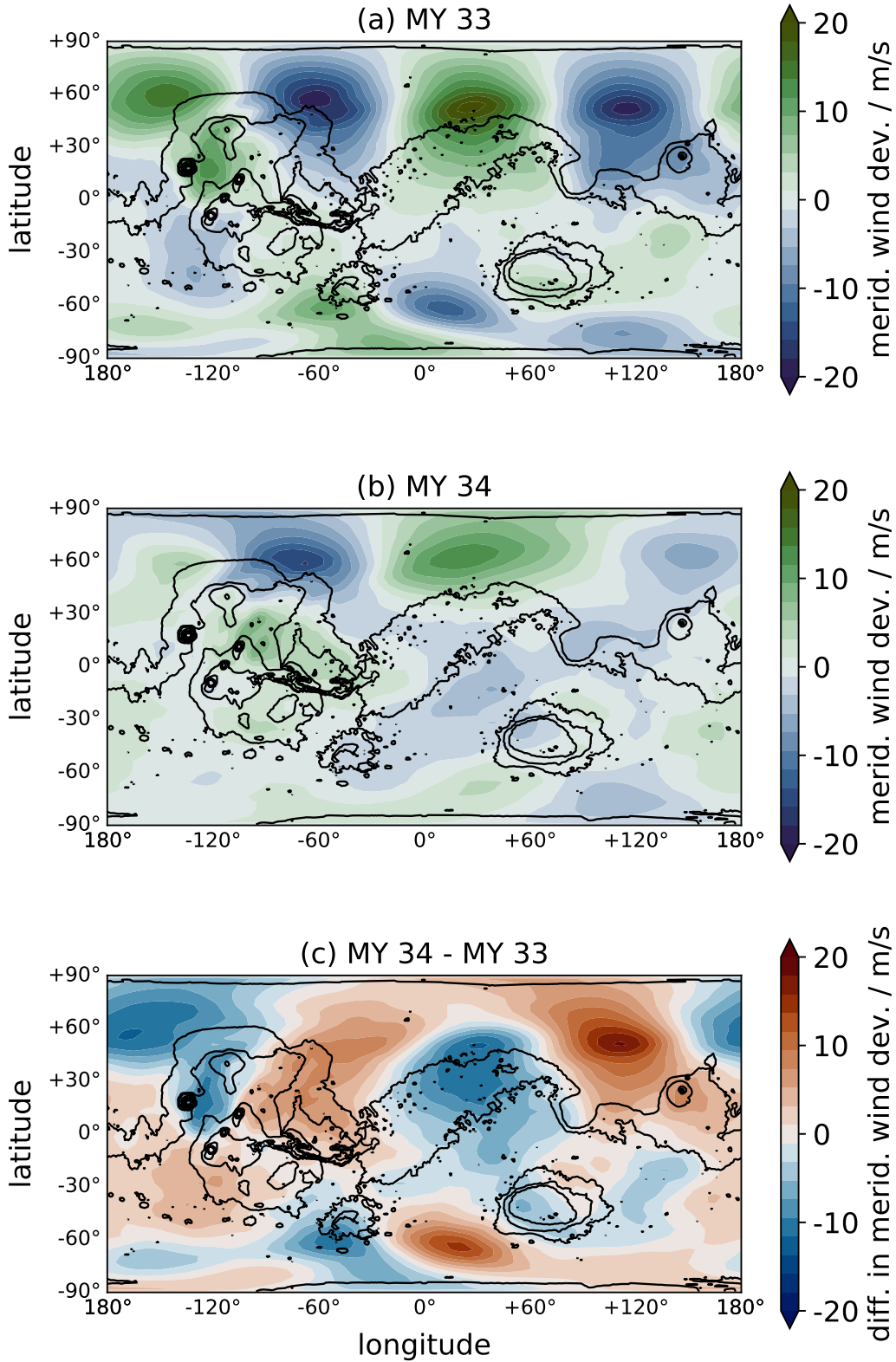


Figure 3. Meridional wind deviation ($[v] - [\bar{v}]$) at 20-30 km for $L_S=200-220^\circ$ for (top) MY 33, (middle) MY 34, and (bottom) the difference between MY 34 - MY 33. Red (blue) in the bottom plot corresponds to either increased (decreased) northward meridional wind deviation or decreased (increased) southward meridional wind deviation.

369 34 GDS significantly diminished northern stationary wave amplitudes, particularly be-
 370 tween -180 - -120° E and 70 - 180° E. This correlates well with the longitudes of greatest PV
 371 reduction (Fig. 1.e). Where the pre-existing stationary wave structure remained rela-
 372 tively intact, such as -120 - -60° E, PV changes were minimal.

373 Likewise in the south, equatorward vortex extension matches where the meridional
 374 wind deviation consists of northward flow, between -70 - 10° E. Where the vortex is more
 375 constrained towards the pole is where the meridional wind deviation consists of south-
 376 ward flow, between 10 - 180° E. The MY 34 GDS almost completely destroyed southern
 377 stationary wave activity (Fig. 3.b), matching the highly symmetrical MY 34 vortex (Fig.
 378 1.d). The greatest change occurred where the MY 33 stationary wave amplitudes were
 379 greatest, between -60 - 60° E; this is where the MY 33 vortex was most latitudinally ex-
 380 tended.

381 While the NP PV structure was largely unaltered during the GDS except becom-
 382 ing more longitudinally symmetric and less annular, the SP polar vortex saw a propor-
 383 tionally much greater and more extensive PV reduction, even at longitudes where the
 384 stationary planetary wave pattern was not significantly altered; this suggests that in ad-
 385 dition to dynamics, other factors contributed to the SP vortex morphological changes.
 386 Fig. 4 shows zonally averaged temperatures and dust opacities for MYs 33 and 34. It
 387 is apparent that GDS-induced atmospheric heating at southern high latitudes was much
 388 greater (up to 33 K at 20-30 km) than at northern high latitudes (up to 15 K at 20-30 km,
 389 and significantly less within the vortex). The southern diurnally-averaged vortex saw sig-
 390 nificant diminishment, while the northern diurnally-averaged polar vortex narrowed in
 391 latitudinal extent but remained coherent and clearly defined. Crucially for the SP ther-
 392 mal winds, the southern hemisphere meridional temperature gradient was substantially
 393 reduced except nearest the pole, visible in the increased spacing between vertical con-
 394 tour lines. This matches with the GDS-induced increase in dust opacity, which occurred
 395 primarily between latitudes 75° S- 50° N (Fig. 4.f), though there was also a large increase
 396 (up to 0.02) between 60 - 90° S up to 20 km. There was little increase in dust opacity north
 397 of 60° N (<0.01), and any increase occurred below ~ 10 km. This implies that diabatic
 398 heating from the increased dust presence may have had a role in affecting the SP but
 399 not NP vortex.

400 The plots of zonally-averaged dust opacity (Fig. 4.b,d,f) also indicate differences
 401 in tracer transport into the polar vortices from the non-GDS and GDS cases. There was
 402 a slight increase in dust opacity below 10 km around 65° N, but a slight decrease further
 403 north around 75 - 80° N. In the south, there was a greater dust opacity increase, up to 0.02
 404 between 5-15 km, with smaller increases up to 30 km. This indicates minimal (if any) dust
 405 transport into the more coherent northern vortex, and possibly greater dust exclusion
 406 from the vortex itself (which could lead to increased opacities on the vortex boundary).
 407 In the south, by contrast, dust opacity increased up to 30% at around 10 km and up to
 408 65% around 30 km (where absolute opacities are lower). The results suggest enhanced
 409 MY 34 tracer transport into the southern vortex, while the northern vortex remained
 410 a coherent transport barrier.

411 To better ascertain the roles of diabatic and adiabatic/dynamical heating, short-
 412 wave (SW), longwave (LW), and net (SW+LW) atmospheric radiative heating rates were
 413 calculated, and their difference between MYs 33 and 34 (Fig. 5). SW heating relies solely
 414 on the atmospheric dust presence, while LW heating is dependent on dust and atmospheric
 415 emission. Given the martian atmosphere's short radiative timescales, SW and LW heat-
 416 ing rates should be in approximate balance in a radiative-convective model (except in
 417 the boundary layer); in a model containing dynamical processes, deviations from radi-
 418 ative balance indicate the presence of dynamical heating/cooling (Wolff et al., 2017). For
 419 example, in MY 33 dynamical cooling is visible between 0 - 60° S at 10-30 km, showing the
 420 upwelling Hadley cell branches, while dynamical heating is visible at high altitudes be-

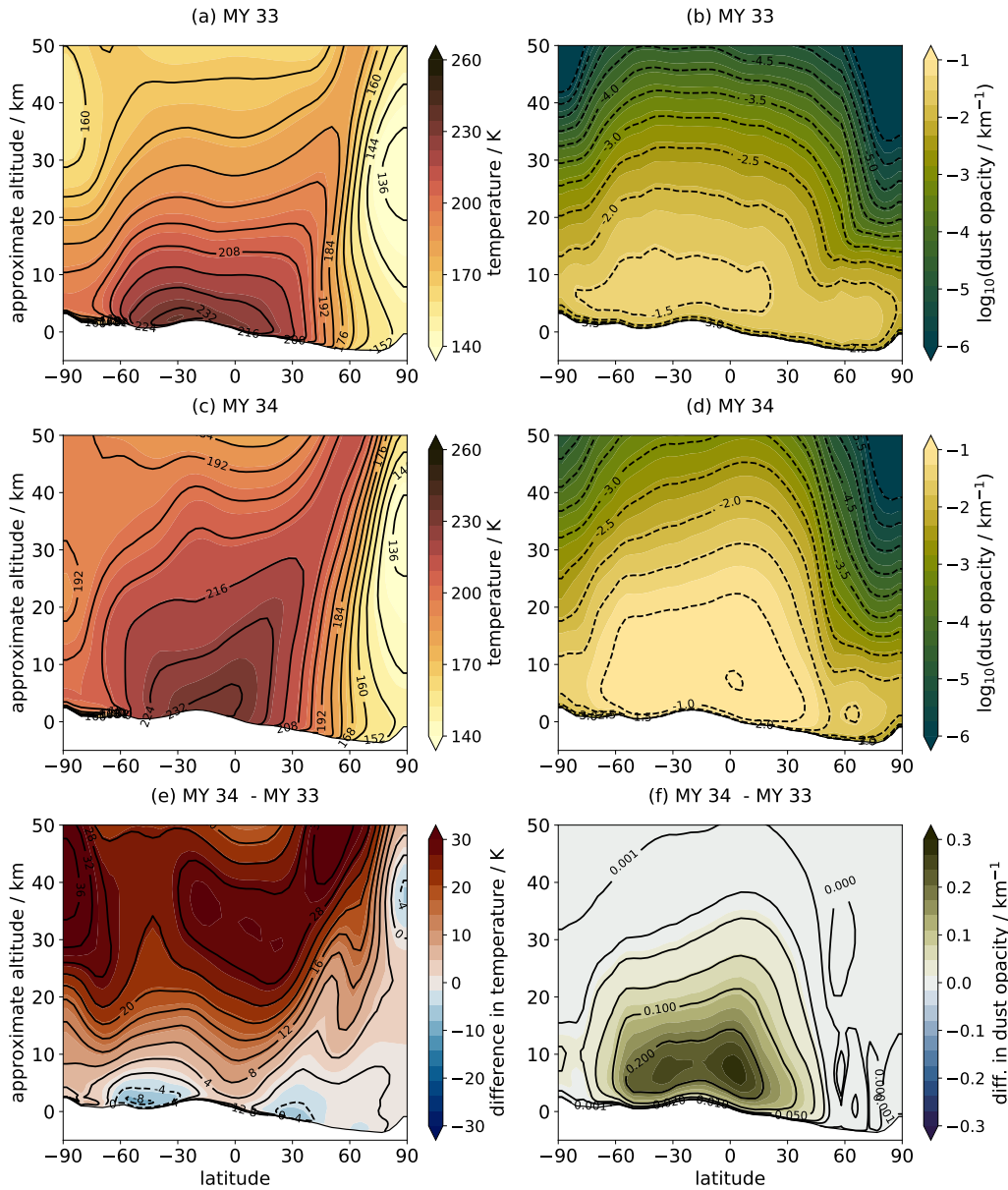


Figure 4. Zonally-averaged (left) temperatures and (right) dust opacities averaged between $L_S=200\text{-}220^\circ$ for (top) MY 33, (middle) MY 34, and (bottom) the difference between MY 34 - MY 33. The plotted dust opacities in subplots b. and d. are \log_{10} values; those in subplot f. are actual dust opacities. All dust opacities are at 670 nm, as used in the MGCM. Altitude was cropped to 50 km for the plots.

421

tween 30-60° N and 60-90° S, showing the downwelling Hadley cells and Mars' characteristic polar warming (e.g. Wilson, 1997; McDunn et al., 2013).

422

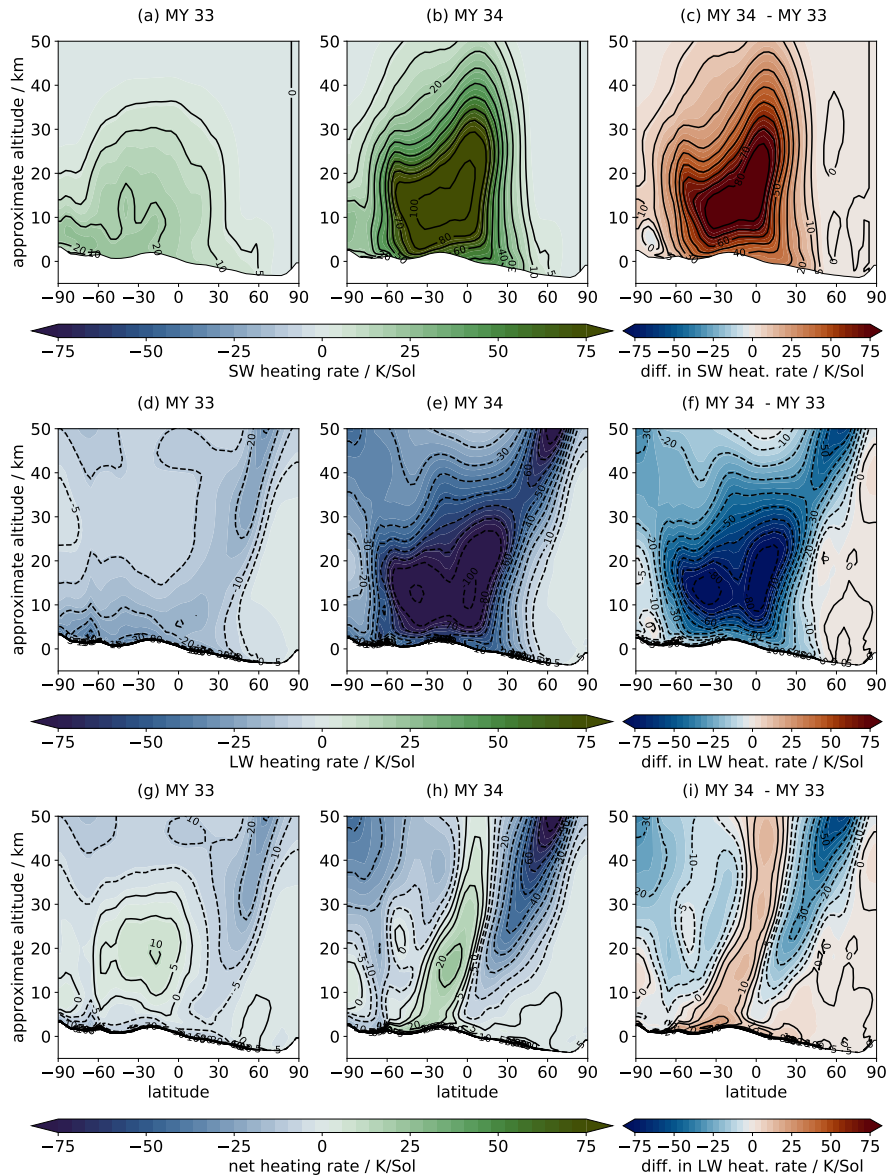


Figure 5. Zonally averaged (top) shortwave (SW), (middle) longwave (LW), and (bottom) net (SW+LW) radiative heating rates averaged between $L_S=200-220^\circ$ for (left) MY 33, (middle) MY 34, and (right) the difference between MY 34 - MY 33. Altitude was cropped to 50 km for the plots.

423

The increased dust loading in MY 34 caused increases in both SW heating and LW cooling rates, particularly between latitudes 75° S to 45° N. Southern SW heating rates

424

425 increased south to 75°S and up to 30 km altitude by at a minimum 6 K/sol (Fig. 5.c),
 426 indicating the important role of direct dust-related atmospheric heating at southern high
 427 latitudes. This heating helped reduce the meridional temperature gradient at southern
 428 high latitudes, weakening the thermal polar jet and thus reducing relative vorticity, caus-
 429 ing local PV reduction via equation 1. This local PV reduction implies equatorward PV
 430 mixing, visible in the small absolute PV increase equatorward of 60°S (Fig. 1.f) as well
 431 as, in the presence of diabatic processes like diabatic descent, destruction of isentropic
 432 PV (Hoskins et al., 1985). PV destruction at higher latitudes has been estimated to oc-
 433 cur on timescales of ~ 5 sols around solstice (Barnes & Haberle, 1996), though this is likely
 434 shorter under significant diabatic heating (e.g. sunlight, dust diabatic heating, CO_2 con-
 435 densation) and could be much longer during polar night given estimated air ages there
 436 of up to 300 sols (Waugh et al., 2019). In the north, by contrast, there was minimal al-
 437 teration to SW heating rates poleward of 45°N , due to the absence of dust. Dynamical
 438 heating was also altered by the GDS: there was an increase in polar warming at both
 439 the NP and SP due to the enhanced Hadley circulation, indicated by the increase in LW
 440 cooling rates (Fig. 5.f,i). The increased dynamical heating over the SP occurred primar-
 441 ily above 30 km, though there was also a ~ 10 K/sol increase between 20-30 km at 65°S .
 442 In the north, the increased dynamical heating narrowed the size of the polar vortex, but
 443 did not cause large-scale local PV reduction as in the south.

444 The results from analysis of atmospheric temperatures, dust opacities, and radia-
 445 tive heating rates suggest that the GDS induced significant dynamical changes at both
 446 poles, but also impacted the SP alone through increased diabatic heating. The effect in
 447 the north was to narrow the latitudinal extent of the polar vortex and make the polar
 448 vortex structure less elliptical and more longitudinally symmetric, but local PV reduc-
 449 tion in the vortex was relatively minimal. The effect in the south, by contrast, was both
 450 to create a more symmetrical polar vortex and also to significantly reduce PV across the
 451 entire pole by reduction of the polar jet through atmospheric heating, leading to a much
 452 diminished diurnally-averaged SP vortex.

453 3.2 Diurnal behaviour

454 This section investigates how each polar vortex was affected by the GDS at differ-
 455 ent local times. As well as the NP/SP notation described above, this section also em-
 456 ploys the concept of Mars Universal Time (MUT). MUT is the local time at longitude
 457 0° , e.g. MUT 00:00 is when it is midnight at longitude 0° , midday at longitude 180° , etc.
 458 Note that the MGCM uses Mars hours and minutes, of which there are the same num-
 459 ber in a martian sol as there are SI hours and minutes in a terrestrial day; seconds are
 460 SI seconds.

461 The SP vortex exhibited a high degree of diurnal variation. This behaviour was first
 462 noted in MCS temperature and dust extinction profile observations from MCS by Kleinböhl
 463 et al. (2020): a mass of colder, more isolated air (indicated by higher PV) and depletion
 464 of dust (indicated by lower dust amounts, in particular at higher altitudes) following the
 465 planet's nightside, centred around MUT 06:00. This is reproduced in the reanalysis (Fig.
 466 6.g-1): the higher absolute PV mass is centred at MUT 06:00, and followed around by
 467 a corresponding CDOD minimum (< 0.6). Interestingly, there does appear to also be di-
 468 urnal variation in the magnitude of the high PV mass, with its absolute PV being greater
 469 at MUT 02:00 and 14:00 than at 06:00 and 18:00. This asymmetry suggests that there
 470 may still be longitudinally asymmetric processes at work, despite almost complete de-
 471 struction of the southern high-latitude wavenumber 1 feature (Section 3.1).

472 Kleinböhl et al. (2020) attribute the diurnal behaviour of the SP vortex during the
 473 GDS to the variation of the MMC throughout the day, with the GDS-enhanced circula-
 474 tion amplifying a pre-existing pattern. To argue for this point, they present the MY
 475 34 MMC as calculated for different local time ranges in an MGCM, with the spatial dust

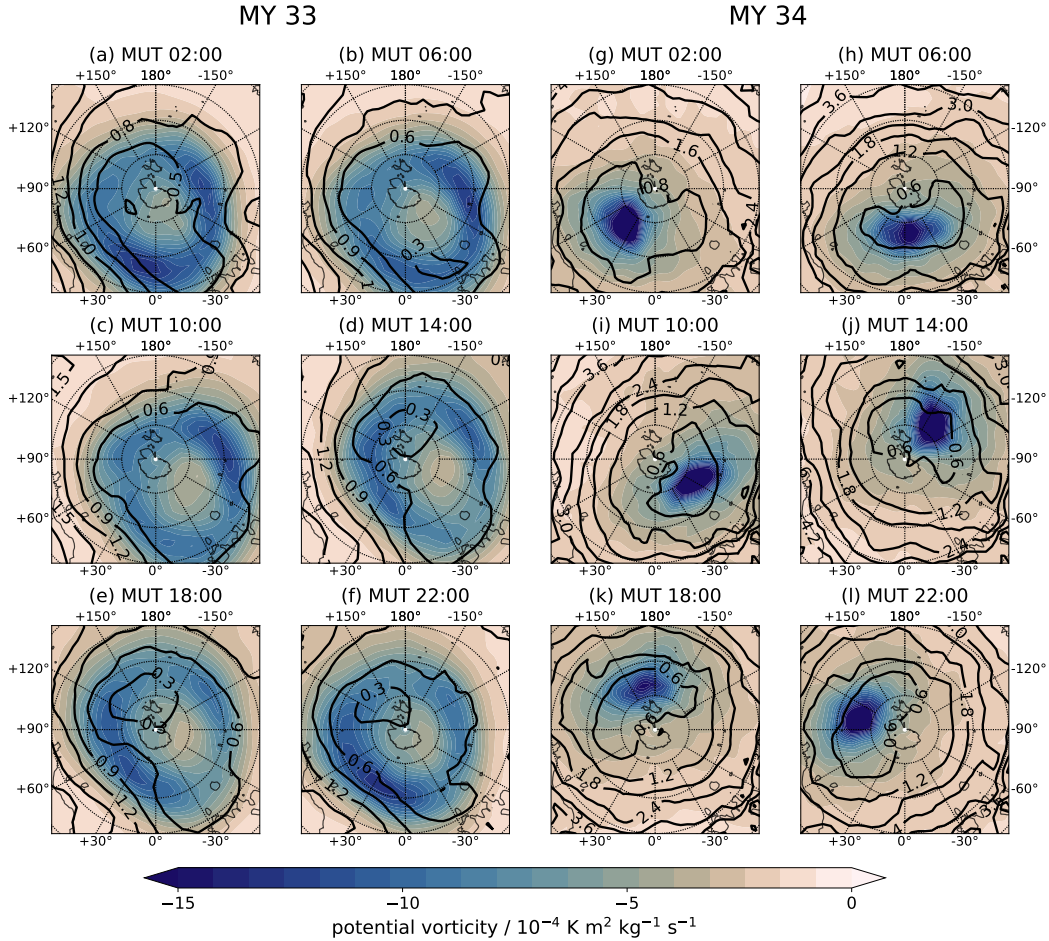


Figure 6. Potential vorticity (colours) as averaged over the $L_S=200-220^\circ$ period on the 300 K isentropic level and CDOD at 610 Pa (contours) over the SP for MY 33 and 34 at six different MUTs. Each MUT is averaged over the 2 hours before and after, or 4 hours in total. Plots are stereographic projections where each latitude circle is 10° separate from its neighbours and the innermost circle represents the 80° latitude band. Lower absolute values indicate lower potential vorticity, and vice-versa.

476 distribution set using a diurnally-averaged MY 34 CDOD map (Montabone et al., 2020).
477 They show a daytime circulation that transports air to the SP, and a nighttime circula-
478 tion that transports air away from it, exhibiting diurnal tidal variations and explain-
479 ing the presence of dust and warmer air (expressed as lower absolute PV) on the SP day-
480 side.

481 The MMC at different local times was calculated for MYs 33 and 34 (Fig. S2), to
482 see if the described pattern holds under non-GDS conditions. The MY 34 results show
483 very good agreement with Kleinböhl et al. (2020): a strong nighttime circulation away
484 from the SP (MUTs 02:00, 22:00), a strong daytime circulation towards the SP (MUTs
485 10:00, 14:00), and a weaker dawn/dusk transitional circulation towards the SP (MUTs
486 06:00, 18:00). This day-night pattern also agrees well with the MY 34 EMARS reanal-
487 ysis, which shows a boosted GDS poleward circulation and nighttime equatorward
488 circulation early during the GDS (Gillespie et al., 2020). The presence of a (weak) cir-
489 culation towards the SP at dawn/dusk explains why the PV minimum covers a small lon-
490 gitudinal extent. The MY 33 circulation closely resembles MY 34, but weaker. One would
491 therefore expect a similar effect in MY 33: a higher absolute PV air mass, and a CDOD
492 minimum, following Mars' nightside. The distinctive annular shape of the vortex is main-
493 tained throughout, including the central eye location around $80^{\circ}\text{S}, -60^{\circ}\text{E}$ (Fig. 6). How-
494 ever, there is also definite diurnal variation in PV and CDOD, with a localised increase
495 in absolute PV following the nightside. The magnitude of this increase is highly longi-
496 tudinally asymmetric: the PV minimum is over 3 MPVU greater at midnight at longi-
497 tude 0° (Fig. 6.a) than longitude 180° (Fig. 6.d), likely due to stationary wave influence.
498 There also appears to be a CDOD minimum following MUT $\sim 06:00$ (Fig. 6.b,d,e,f). These
499 results corroborate the Kleinböhl et al. (2020) finding that the observed MY 34 SP vor-
500 tex behaviour is due to an enhanced MMC boosting an already present mechanism by
501 which warmer, dustier air is transported to the dayside of the pole and colder, clearer
502 air remains on the nightside. As well as the local time effect, there is also a complex in-
503 terplay with southern stationary waves.

504 Finally, given the weaker northern circulation at this season one would expect a
505 similar but weaker diurnal cycle at the NP vortex. Fig. S3 shows some diurnal varia-
506 tion in the PV maximum's longitude, roughly following the nightside, but it is less clear
507 than in the south, and more dominated by apparent intrinsic longitudinal asymmetries.
508 PV is consistently high, up to 50 MPVU, around longitudes -60° - 0°E for half the diur-
509 nal cycle, and around longitudes 150° - -150°E for the other half; elsewhere, it is notably
510 weaker even during local nighttime (e.g. MUT 02:00). This pattern is consistent with
511 the observed residually elliptical PV structure visible in the diurnally averaged plot (Fig.
512 1.c). There is also an apparent minimum in CDOD diurnal variation following the PV
513 maximum, as at the SP. Again, the pattern is less obvious than in the south. These re-
514 sults suggest local time variations at the NP similar to those in the SP, but largely ob-
515 scured by planetary wave effects.

516 4 Discussion

517 The MY 34 GDS had a significant effect on the morphology of both the southern
518 and northern polar vortices, through both altered dynamics (in the north) and a com-
519 bination of altered dynamics and altered radiative heating (in the south). For both hemi-
520 spheres, the change in the dynamics came in the form of changes to the meridional flow,
521 linked to changes to the high latitude stationary wave structure. This stationary wave
522 structure is an alternative interpretation of the morphology of both polar vortices at this
523 time of year, with the northern wavenumber 2 feature following the NP vortex's char-
524 acteristic elliptical shape. These are two possible perspectives, wave-based or PV-based,
525 on the same phenomenon.

526 The GDS-induced changes to the stationary waves were reflected in changes in the
527 shape of polar vortices, specifically reducing their longitudinal and, in the southern case,
528 latitudinal asymmetry. Indeed, the amplitude changes in the northern hemisphere wavenum-
529 ber 1 feature exactly correlate to modelled changes in PV. In the south, the GDS-induced
530 changes to the heating rates in the atmosphere also had a crucial effect on the diurnally-
531 averaged polar vortex structure. Increased atmospheric heating from the high southern
532 dust loading in conjunction with dynamical heating reduced the meridional temperature
533 gradient, substantially reducing the strength of the polar jet and causing large-scale PV
534 mixing and dilution along the isentropic surface. The remnant westerly jet was shifted
535 to a tighter area circumscribing the pole. This was a dramatic acceleration of the already
536 ongoing decay of the SP vortex. In the north, by contrast, there were minimal changes
537 in radiative heating rates at high latitudes, preserving the strength of the polar vortex
538 even as its area and ellipticity decreased by the (longitudinally asymmetric) expansion
539 of the Hadley cell to higher latitudes.

540 The effect of the GDS on northern stationary wave activity was twofold. Most rel-
541 evant for the polar vortex, one effect was to shift the latitude of the NP jet northwards,
542 away from the high amplitude zonal topography of the northern mid-latitudes which me-
543 chanically forces the wavenumber 2 stationary wave (Nayvelt et al., 1997), thereby re-
544 ducing high latitude stationary wave activity associated with the polar jet. There was
545 residual ellipticity in the MY 34 NP vortex, visible in the zonal wind and PV structure
546 in Fig. 1.c and in remnant stationary wave activity between -90° E and 60° E, which can
547 be likely attributed to the presence of high-latitude topographic gradients north of Alba
548 Patera and in Acidalia Planitia. Secondly at lower, sub- 45° N latitudes, where differen-
549 tial heating due to zonal topographic gradients plays a greater role in driving station-
550 ary wave activity (Nayvelt et al., 1997), the high dust loading of the GDS dramatically
551 reduced zonal surface temperature gradients (Streeter, Lewis, Patel, Holmes, & Kass,
552 2020), thereby dramatically reducing stationary wave activity across all longitudes. This
553 latter effect was less relevant for the high-latitude polar vortex. At southern high-latitudes,
554 the already weak (relative to the north) stationary wave activity was substantially re-
555 duced as well. South of 60° S, where the dust cloud extended, the reduced meridional ther-
556 mal gradient reduced zonal wind speeds, decreasing mechanical forcing and the wavenum-
557 ber 1 stationary wave.

558 The stationary wave interpretation explains not only the elliptical polar vortex shape,
559 but why the elliptical shape is more prevalent in the northern hemisphere, as noted by
560 Waugh et al. (2016). The non-GDS northern hemisphere has, at the $L_S=200-220^\circ$ pe-
561 riod, a strong and latitudinally extended wavenumber 2 feature. The southern hemisphere,
562 by contrast, has a more ambiguous stationary wave structure at this time, with a mid-
563 high latitude wavenumber 2 feature transitioning to a wavenumber 1 feature near the
564 pole. These wave features most likely have a topographically-induced origin, both through
565 mechanical forcing and differential heating (e.g. Nayvelt et al., 1997; Hollingsworth &
566 Barnes, 1996; Haberle et al., 2019). Mitchell et al. (2015) and Rostami et al. (2018) hy-
567 pothesised that the elliptical polar vortex shape might be linked to large-scale topogra-
568 phy; here we provide compelling evidence that the elliptical shape of the vortices is the
569 shape of the stationary planetary waves, by showing that the suppression of one entails
570 the suppression of the other.

571 There was some apparent disagreement with previous literature on the effects of
572 high dust loading on polar vortex structure, though this may be due to the time of year
573 under consideration, among other factors. Mitchell et al. (2015), using the MACDA re-
574 analysis, found that a regional-scale dust storm at $L_S=320^\circ$ caused a $\sim 10^\circ$ latitude shift
575 in the NP vortex, as well as an overall weakening in PV. This study shows little change
576 in the broad morphology of the NP vortex even from the very high dust loading of a GDS,
577 and only a localised weakening in PV. A couple of factors could be responsible for this
578 disagreement. First, and likely most importantly, the times of year are different. At $L_S=320^\circ$,

579 the NP vortex is already weakening as the planet approaches equinox, while in the $L_S=200-$
580 220° period the NP vortex is strengthening as the planet approaches northern winter sol-
581 stice. While a proper intercomparison between this reanalysis and MACDA is called for,
582 this suggests that seasonal differences could be crucial in determining polar vortex re-
583 sponse to sudden high dust loadings. The second factor is the nature of the regional storm
584 in question itself, and specifically its spatial location. This could have an impact on its
585 resultant dynamical and radiative effects.

586 Guzewich et al. (2014) investigated the impact of the MY 25 GDS on stationary
587 waves, primarily the wavenumber 1 mode, finding that this mode was enhanced at both
588 northern and southern high latitudes. The results presented here are consistent with Guzewich
589 et al. (2014) for the northern hemisphere, as they show a marked decrease in amplitude
590 of the wavenumber 2 mode and a transition towards a more apparently dominant wavenum-
591 ber 1 mode; this is even more marked at higher altitudes than those shown here. The
592 picture in the southern hemisphere is more complex however, as these results show de-
593 creased wavenumber 1 activity, even higher than 20-30 km (not shown). This is puzzling
594 due to the close similarity in season and general structure between the MY 25 and MY
595 34 GDS, and further study is required to explain this difference.

596 Guzewich et al. (2016) found that the NP vortex was significantly disrupted in an
597 MGCM with a prescribed high-dust peak at $L_S=270^\circ$, but not with a dust peak of stan-
598 dard magnitude at $L_S=200^\circ$. Reversing the seasonal dust loading did not impact the SP
599 vortex. In the study presented here, the SP vortex was disproportionately affected by
600 an equinoctial/early perihelion season GDS, when the NP/SP vortices are strengthen-
601 ing/weakening, respectively. This is not inconsistent with Guzewich et al. (2016), as they
602 investigated the effects of high dust loading at southern summer solstice, $L_S=270^\circ$, and
603 standard dust loading at other times of year. Comparing the two sets of results suggests
604 a key role for seasonality of high dust loading in resultant polar vortex behaviour. Given
605 that GDS have been observed to occur at various times within the dusty season $L_S=180-$
606 360° (Shirley, 2015), with the three most recent events occurring at equinox (MY 25), sol-
607 stice (MY 28), and equinox (MY 34), it is worth investigating both equinoctial and sol-
608 sticial events and the likely significant effect of GDS seasonality.

609 At this point it should be noted that despite the evidence of wave-related effects
610 on the SP vortex structure, wave structure alone is not the whole story in explaining its
611 broad morphology. Notably, planetary wave structure does not explain the off-pole pres-
612 ence of the eye of lower absolute PV. Following the hypothesis that the central absolute
613 PV minimum is the result of localised diabatic heating from latent heat release as CO_2
614 condensation occurs over the pole (Toigo et al., 2017; Rostami et al., 2018), one would
615 expect that this location in the reanalysis - approximately 80°S , -60°E - undergoes greater
616 CO_2 condensation than the surrounding seasonal cap. The reanalysis and the MCS sur-
617 face CO_2 ice observations from Piqueux et al. (2015) offer an opportunity to try to fur-
618 ther validate this hypothesis. MCS observations from this period show a clear ice rem-
619 nant present between $70-90^\circ\text{S}$ and $-120-60^\circ\text{E}$, even as the ice in the opposite hemisphere
620 has almost completely disappeared (Piqueux et al., 2015). While the reanalysis shows
621 a greater latitudinal extent of CO_2 ice coverage than the observations, it does agree with
622 the observations regarding where the most CO_2 ice mass is located, at the same latitude
623 range noted earlier. The location of this remnant surface ice feature agrees very well with
624 the location of the low absolute PV eye seen in for example Fig. 1.b.

625 MCS observations suggest that this location is indeed an area of greater CO_2 con-
626 densation, resulting in a thicker and therefore longer lasting ice cap; alternatively, it could
627 be that this region is colder than the surrounding areas, thereby promoting greater CO_2
628 condensation and less rapid sublimation. These “cold spots” could also be linked to the
629 presence of CO_2 ice clouds (Hayne et al., 2012), and are an ongoing area of study. In any
630 case, this would seem to be further supporting evidence for the Toigo et al. (2017) hy-
631 pothesis for the cause of the annular PV structure. However, Figure 6 appears to show

632 that the local PV minimum or “eye” has lower absolute PV during local daytime. This
 633 is contrary to what might be expected if CO₂ condensation was the driving mechanism
 634 behind the eye, as condensation would be expected to occur preferentially at local night
 635 due to lower atmospheric temperatures. Indeed despite low local radiative flux, the in-
 636 terconnectedness of the global thermal tide can still force large diurnal atmospheric vari-
 637 ations even at high latitudes (Lee et al., 2009). One way to account for this would be
 638 to consider the diurnal pattern of mass transport (Figure S2); air mass is transported
 639 towards the south pole during the daytime, thereby supplying more mass for condensa-
 640 tion onto the seasonal cap. This may outweigh diurnal changes in radiative flux and at-
 641 mospheric temperature causing sublimation/condensation, which will be low at such high
 642 latitudes. This remains to be investigated in future work. Finally, it is interesting to note
 643 that the GDS appears to shift the local PV minimum towards the pole itself, at least for
 644 the $L_S=200-220^\circ$ period; this shows that the GDS affects whatever mechanisms are re-
 645 sponsible for location of the off-pole eye in non-GDS conditions.

646 Lastly, analysis of local time variations in the SP vortex during the GDS reproduce
 647 the Kleinböhl et al. (2020) results of an isolated (high absolute PV), dust-clear air mass
 648 trapped on the nightside. This study provides further evidence that the proposed mech-
 649 anism, a boosting of the meridional circulation in MY 34 intensifying an existing local
 650 time pattern of nightside transport away from the SP and dayside transport towards,
 651 is indeed correct. This pattern is shown to exist at the same time period in MY 33, with
 652 a weaker but still visible impact on the diurnal variation of PV and CDOD. The weaker
 653 meridional circulation and stronger stationary wave activity means that the diurnally-
 654 averaged structure, consisting of an off-pole eye of lower absolute PV with a surround-
 655 ing annulus of higher PV, remains visible at all local times. There is some diurnal vari-
 656 ation in PV at the NP in MY 34, but this is much weaker than in the south.

657 5 Conclusions

658 The GDS had a significant impact on both the northern and southern polar vortices.
 659 In the north, the polar vortex structure (reflected via PV and zonal wind speeds)
 660 became less elliptical and more longitudinally symmetric. PV increased slightly at the
 661 pole itself and decreased at lower latitudes, indicating a narrowing of the area of the po-
 662 lar vortex as the Hadley cell descending branch extended further poleward. In the south,
 663 there was significant mixing and therefore local reduction of PV across the entire polar
 664 vortex, and a corresponding increase in near-pole wind speeds and decrease in wind speeds
 665 at lower latitudes. The remaining diurnally-averaged polar vortex was significantly more
 666 symmetric than in the non-GDS case, and centred at the pole itself rather than off-pole.
 667 The non-GDS asymmetries in polar vortex shape at both poles were found to be linked
 668 to a longitudinally asymmetric meridional flow, reflected in the stationary planetary wave
 669 structure in each hemisphere. The GDS significantly reduced the amplitude of the north-
 670 ern wavenumber 1 feature and almost entirely destroyed the southern stationary wave
 671 feature, which was expressed in more symmetrical polar vortices.

672 The MY 34 GDS also caused a peculiar feature at the southern pole: a mass of colder,
 673 isolated air coupled with a minimum in dust which followed the nightside of the planet,
 674 first seen in direct MCS observations (Kleinböhl et al., 2020). This feature was clearly
 675 visible in the reanalysis and its suggested causal mechanism, a boosted mean meridional
 676 circulation enhancing a pre-existing (non-GDS) pattern of nighttime air transport away
 677 from the southern pole and daytime air transport towards it, further corroborated. This
 678 pattern was shown to be present at the northern pole in a weaker form, and at the south-
 679 ern pole, albeit in a weaker form, in the non-GDS year MY 33.

680 The intensity and shape of the polar vortices and their corresponding westerly zonal
 681 jets control what can be transported above and onto the poles themselves. The non-GDS
 682 shape and GDS-induced alteration of the vortices therefore has potentially important

consequences for both seasonal and long-term transport of dust, water, and chemical species into the polar regions and onto the seasonal CO₂ caps. The substantially diminished diurnally-averaged southern vortex during an equinoctial GDS should provide a weaker barrier to transport into the southern polar region and onto the surface, while the relatively robust northern vortex should remain an effective barrier to transport. This is especially relevant considering the boosted meridional circulation during GDS. Additionally, if the poleward shift of the westerly jets is a consistent effect of equinoctial GDS, this implies that quantities such as dust, water, and chemical species can be transported to higher latitudes during these intense dust events. If this kind of pattern holds over the course of the thousands of years that Mars maintains its particular axial obliquity, this has further implications for the record of deposited dust at the southern and northern polar layered deposits, for example as measured by the SHARAD subsurface radar (Seu et al., 2018). Over long timescales, the spatially anchored morphology of the polar vortices may influence the longitudinal pattern of dust deposition over the polar regions.

Validation is an important consideration for results from a meteorological reanalysis, particularly when based on fields (PV, wind speeds) not directly measured from orbit. The most valuable possible future measurements would be of atmospheric winds from an orbiter. These would allow better constraints on the structure of the polar vortices, and (combined with temperature measurements) direct calculation of PV. In the near future, ongoing retrievals of temperature, dust, and long-lived trace gases from TGO and its NOMAD (Vandaele et al., 2015; Patel et al., 2017) and ACS (O. I. Korabev et al., 2015; O. Korabev et al., 2018) spectrometer suites will enable further investigation of tracer transport and an opportunity to cross-validate and jointly assimilate multiple orbiter observations. NOMAD/ACS provide the crucial feature of observing over a range of martian local times, enabling exploration of the diurnal cycles of tracer transport and polar dynamics.

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