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eprints@whiterose.ac.uk https://eprints.whiterose.ac.uk/ Validation of a stochastic digital packing algorithm for porosity prediction in fluvial gravel deposits

- 3 Rui Liang <sup>a,\*</sup>, Tobias Schruff <sup>a</sup>, Xiaodong Jia <sup>b</sup>, Holger Schüttrumpf <sup>a</sup>, Roy M. Frings <sup>a</sup>
- <sup>a</sup> Institute of Hydraulic Engineering and Water Resources Management, RWTH
- 5 Aachen University, Aachen D 52056, Germany
- <sup>b</sup> Institute of Particle Science and Engineering, School of Process, Environmental
   and Materials Engineering, University of Leeds, Leeds LS2 9JT, UK
- 8 \* Corresponding author. Tel.: +49 (0)241 80 25757; Fax: +49 (0)241 80 25750.
- 9 E-Mail address: rui@iww.rwth-aachen.de (R. Liang).

#### 10 Abstract

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Porosity as one of the key properties of sediment mixtures is poorly 11 12 understood. Most of the existing porosity predictors based upon grain size characteristics have been unable to produce satisfying results for fluvial 13 sediment porosity, due to the lack of consideration of other porosity-14 15 controlling factors like grain shape and depositional condition. Considering 16 this, a stochastic digital packing algorithm was applied in this work, which provides an innovative way to pack particles of arbitrary shapes and sizes 17 18 based on digitization of both particles and packing space. The purpose was to test the applicability of this packing algorithm in predicting fluvial 19 20 sediment porosity by comparing its predictions with outcomes obtained from laboratory measurements. Laboratory samples examined were two 21 natural fluvial sediments from the Rhine River and Kall River (Germany), 22 23 and commercial glass beads (spheres). All samples were artificially combined into seven grain size distributions: four unimodal distributions 24 and three bimodal distributions. Our study demonstrates that apart from 25 26 grain size, grain shape also has a clear impact on porosity. The stochastic digital packing algorithm successfully reproduced the measured variations 27 in porosity for the three different particle sources. However, the packing 28 algorithm systematically overpredicted the porosity measured in random 29

dense packing conditions, mainly because the random motion of particles
during settling introduced unwanted kinematic sorting and shape effects.
The results suggest that the packing algorithm produces loose packing
structures, and is useful for trend analysis of packing porosity.

34 Keywords: Porosity; Sediment; Grain shape; Random packing; Rivers

# 35 **1. Introduction**

36 Porosity prediction of sedimentary deposits is of interest in a fluvial environment. Previous studies have shown that porosity, as a key structural 37 38 property, plays an important role in the morphological, ecological and geological characteristics of fluvial systems. Morphologically, porosity 39 governs the initiation of sediment motion and bank collapse (e.g., Wilcock, 40 41 1998; Vollmer and Kleinhans, 2007). Ecologically, porosity determines the interstitial space of the hyporheic zone for aquatic habitats (e.g., Boulton et 42 al., 1998). Geologically, porosity dominates the exploitable reserve of oil, 43 gas, and groundwater stored in the voids of fluvial deposits (e.g., Athy, 44 1930). To date, existing porosity predictors can generally be classified into 45 two types: (1) empirical predictors; and (2) theoretical predictors. Most 46 efforts to predict porosity have been empirically driven, to a large extent 47 based upon median grain size D<sub>50</sub> (e.g., Carling and Reader, 1982; Wu and 48 Wang, 2006), sorting coefficient  $\sigma$  (e.g., Wooster et al., 2008), or a 49 combination of different grain size characteristics (e.g., Frings et al., 2011; 50 51 Desmond and Weeks, 2014). Theoretical predictors such as geometrical models (e.g., Ouchiyama and Tanaka, 1984; Suzuki and Oshima, 1985) or 52 analytical models (e.g., Yu and Standish, 1991; Koltermann and Gorelick, 53 1995; Esselburn et al., 2011) relate porosity to the full grain size distribution 54

of perfect spheres. The performance of these predictors has been investigated by comparing porosity values measured in situ with those computed by the predictors (e.g., Frings et al., 2008, 2011). Unfortunately, these predictors produced unsatisfying results in predicting fluvial sediment porosity (Frings et al., 2011), probably because such predictors mainly focused on grain size characteristics, ignoring other porosity-controlling factors such as depositional environment and grain shape.

Effects of grain shape on porosity have received less attention, due to 62 63 the complexity of arbitrary shapes of natural particles. Over the past decade, the application of computer simulations for the study of granular particle 64 packings has become more popular, supported by developments in the 65 66 computer hardware industry. However, most of the computer simulations have been limited to simple analytical geometries such as cylinders (Zhang 67 et al., 2006), disks (Desmond and Weeks, 2009), ellipsoids (Donev et al., 68 69 2007; Zhou et al., 2011) and spherocylinders (Abreu et al., 2003; Williams and Philipse, 2003; Zhao et al., 2012). The major reason is the practical 70 71 difficulty of representing and handling irregular shapes using vector-based approaches. Traditional ways to construct an irregular particle require the 72 73 user to place spherical elements within a meshed polyhedral body (e.g., 74 Wang et al., 2007; Matsushima et al., 2009; Ferellec and McDowell, 2010; 75 Fukuoka et al., 2013), which consumes high computational costs with large 76 numbers of components (spheres) involved (Hubbard, 1996; Song et al., 77 2006). Although techniques using 3D polyhedral (Latham et al., 2001) or continuous superquadric functions (Williams and Pentland, 1992; Lu et al., 78 79 2012) provide a straightforward way to generate irregular particle shapes,

complex contact-detection algorithms are needed, leading to deterioration in
simulation speed as particle complexity increases (Johnson et al., 2004).

In order to overcome these difficulties, a stochastic digital packing 82 83 algorithm was developed (Jia and Williams, 2001). The packing algorithm is distinguished from the traditional vector-based packing models by 84 digitization of both particles and packing space, allowing for a much easier 85 86 and computationally efficient way to pack particles of irregular shapes with no more than an ordinary PC. These advantages make it attractive to create 87 88 packings of complex fluvial deposits, and to study the grain shape effects on porosity. Applications of this stochastic digital packing algorithm have 89 proven to provide relatively accurate porosity predictions for both fine 90 91 powders (Jia et al., 2007) and large spheres (Caulkin et al., 2006, 2007) in 92 the fields of material science and engineering chemistry. Nevertheless, the packing algorithm has not yet been used for generating packings of fluvial 93 94 deposits. Therefore, the primary purpose of this work was to test the applicability of the stochastic digital packing algorithm in predicting fluvial 95 sediment porosities. In this study, we focused on fluvial gravel mixtures and 96 did so by comparing the predicted porosities with those obtained from 97 laboratory measurements. 98

99 **2. Materials and methods** 

100 2.1. Particle acquisition and analysis

101 The particles employed for this study came from three different sources: 102 (1) fluvial gravels from the Rhine River (Germany), (2) fluvial gravels from 103 the Kall River (Germany), and (3) commercial glass beads. The Rhine 104 sediments were collected from the channel bed between the barrage of Iffezheim and the German-Dutch border between July 2008 and April 2011.
Quartz is the dominant lithology. The Kall sediments were collected from
the channel bed near the river mouth in June 2014. Slate is the dominant
lithology.

After acquisition, the fluvial sediments were carefully cleaned by 109 flushing with fresh water, dried in an oven at 105 °C for 48 h and sieved 110 into seven size fractions: 2.8-4 mm, 4-5.6 mm, 5.6-8 mm, 8-11.2 mm, 11.2-111 16 mm, 16-22.4 mm, 22.4-31.5 mm. Subsequently, these fractions were 112 113 combined into seven grain size distributions: four unimodal ones with logarithmic standard deviations ( $\sigma_{0}$ ) of 0.00, 0.32, 0.49 and 0.71, and three 114 115 bimodal ones, with the finer mode, making up either k=30, k=50 or k=70 116 percent of the distribution (Fig. 1). The glass beads with seven size fractions of 3, 4, 6, 8, 11, 16 and 22 mm were also combined into the same 117 distributions as above. 118

For the fluvial sediments, nine representative particles were chosen 119 based on visual judgments from each of the seven sieve fractions, and 120 digitized (Fig. 2) using a nonmedical X-ray computed tomography (CT) 121 scanner. Shape analysis was done according to the classic Zingg diagram 122 123 (Zingg, 1935), which categorizes particle shape into sphere, disc, blade and 124 rod categories on the basis of the elongation ratio (b/a) and flatness ratio (c/b), where a, b and c are the long, intermediate and short orthogonal axes 125 respectively of the smallest volume imaginary box that can contain the 126 127 particle (Blott and Pye, 2008). It can be seen in Figure 3 that most of the Rhine sediments locate within the sphere area while the Kall sediments are 128 dominated by disks and blades. According to Krumbein's (1941) equation 129

(1), the intercept sphericity (ψ) for each selected particle was calculated,
with an average intercept sphericity of 0.74 gained for the Rhine sediments
and 0.55 for the Kall sediments.

133 
$$\psi = \sqrt[3]{\frac{b*c}{a^2}} \qquad (1)$$

### 134 2.2. Laboratory porosity measurements

The water displacement method (Bear, 1972) was used for porosity 135 measurements. The experimental procedure was as follows: firstly, a plastic 136 cylinder with an inner diameter of 104 mm was partially filled with a known 137 volume of water  $V_w$  larger than the expected pore volume of the particles to 138 be added. Then, particles of 3 kg mass were added into the cylinder in small 139 140 well-mixed portions, together with gently tapping the side of the cylinder in 141 order to dislodge trapped air bubbles and obtain a stable, dense packing. The 142 final water level was visually read to obtain the whole accumulated volume  $V_a$  ( $V_a = V_w + V_s$ , where  $V_s$  is the volume of the solid fraction). The jagged 143 surface of the particle packing caused by the wide range of sizes and shapes 144 145 was then smoothed by hand and the total volume of the particle packing  $V_t$ 146 (including pores) was obtained through reading the height of the particle packing. Eventually, the porosity was computed as  $n = V_p/V_t$ , where  $V_p$ 147  $(= V_t - (V_a - V_w))$  is the pore volume of the particle packing. 148

In total, 42 laboratory porosity experiments were performed as a basis for the validation of the stochastic digital packing algorithm: 14 experiments with the sub-spherical Rhine sediments (7 distributions, each two times), 14 experiments with low-sphericity Kall sediments (again  $7\times 2$ ) and 14 experiments with the spherical glass beads (again  $7\times 2$ ).

# 154 2.3. Porosity simulation

The stochastic digital packing algorithm of Jia and Williams (2001) is 155 designed to pack particles of arbitrary sizes and shapes in a confined space 156 157 of arbitrary geometry. In this packing algorithm, every element is digitized: 158 each particle as a coherent collection of voxels, the packing space (in a container) as a lattice grid, and the movements take place in units of grid 159 160 cells. During the simulation, the movements of particles, both translational and rotational, are random. In 3D, there are 26 possible translational 161 162 directions: 6 orthogonal and 20 diagonal. The diagonal moves are treated as a combination of two orthogonal moves. To ensure particles settle while still 163 164 make use of every available space, a rebounding probability is used. An 165 upward movement (which may be an orthogonal move or part of a diagonal 166 move) is only realized with this probability. After translation, a trial rotation follows, and it is accepted if the rotation does not result in overlaps. 167 168 Compared with vector-based approaches and for complex shapes, this digital approach is advantageous in several respects. First, there is no 169 170 conversion or parameterization required, since objects digitized by modern imaging devices, such as X-ray tomography (e.g., Richard et al., 2003) or 171 172 nuclear magnetic resonance imaging (e.g., Kleinhans et al., 2008), are 173 already in the digital volumetric format required by the packing algorithm. 174 Secondly. collision and overlap detection (normally the most computationally expensive part of packing simulations) is much easier to 175 176 implement as computer code, and usually faster to execute for complex shapes. Thirdly, the number of voxels used to represent objects, and hence 177 to a large extent the simulation runtime, does not necessarily increase with 178

shape complexity. The reverse is also true: it does not necessarily reduce
with shape simplification either. Further details on the stochastic digital
packing algorithm can be found elsewhere (Jia and Williams, 2001; Caulkin
et al., 2006, 2007).

In order to produce porosity results comparable to those aforementioned 183 measurements, simulation conditions need to be set up to resemble the 184 185 laboratory experiments, with respect to the packing space, the particle mixtures and the packing process. The digital objects (i.e., packing space 186 187 and particles) were prepared with DigiUtility, which is a bundled tool for viewing, manipulating and preparing digital files for this packing algorithm. 188 In DigiUtility, a cylinder (packing space) with solid boundary was built with 189 190 the size of 104 mm in diameter, and 300 mm in height, which is slightly 191 higher than the largest real packing heights (about 250 mm) to ensure all the particles would drop into it. The particle mixtures (i.e., number of particles 192 193 in each of the fractions) employed in these simulations were derived on a weight-to-weight basis. For glass beads, the numbers of particles in each 194 195 fraction were determined as the ratio of the real mass of each fraction to the single particle mass (density of 2500 kg/m<sup>3</sup> used for glass beads). The 196 regular spherical shapes with different sizes were directly created in digital 197 198 formats using DigiUtility. In the case of the fluvial sediments, we used nine digitized typical particles to represent each fraction and repeated them as 199 many times as needed to make up the feedstock according to the required 200 grain size distributions. The density of fluvial gravels was set to  $2650 \text{ kg/m}^3$ . 201 Resolution of 0.5 mm/voxel for the digital objects was assigned as it offers 202

relatively precise representation of the real particles in both dimension andshape, and also limits the computational cost to a feasible amount.

Having the digital objects created, a range of options and parameters 205 206 was set to mimic the packing process. The source was set to "rain-dropping" mode to let the particles randomly drop from a circular area above the 207 208 cylinder. In addition to the translational movements, particles were also 209 allowed to rotate randomly during the simulation. Optimized values of the parameters (rebounding probability, addition rate and number of time steps) 210 211 that control the generated packing structures were chosen such as to create 212 the densest possible packings. By doing so, simulation conditions (Table 1) 213 matched the experimental setups as close as possible. Finally, the porosity 214 of the digital packings was determined as the ratio of the number of empty 215 voxels to the total number of voxels within the corresponding packing space. Porosity was calculated for the lower 90% of the mixture to exclude effects 216 217 of surface irregularities. Each simulation was also done twice and 42 simulations were achieved in total. 218

# 219 **3. Results**

## 220 3.1. Measured porosity

The porosity measured in the laboratory experiments is shown in Figure 4. For the unimodal particle mixtures, porosity decreases with increasing logarithmic standard deviations, while the bimodal particle mixtures generally have lower porosity than the unimodal mixtures. This variation in porosity reflects the mixing effect between small and large particles.

226 Porosity comparisons between the three different particle sources show227 the low-spherical Kall sediments and the spherical glass beads produced

higher porosity than the sub-spherical Rhine sediments, which confirms that
there is a decrease and then increase in porosity as particle shape varies
from spherical to platy (Tickell and Hiatt, 1938; Zou and Yu, 1996). On the
other hand, in the case of the bimodal particle mixtures, different tendencies
toward the porosity are appreciable (Fig. 4B), suggesting grain shape exerts
a quite complicated influence on porosity, not merely in variation of amount
but in variation of trend.

It should be noted that the dense sediment deposits packed by hand in 235 236 the laboratory experiments are not fully representative of natural situations where grain arrangement is determined by depositional conditions, such as 237 238 flow impact (with near-bed turbulence playing an important role) and burial 239 depth (compaction mechanism). This topic is beyond the current effort. 240 Nonetheless, based on the comparisons between field measurements of porosity in the River Rhine (28 measurements on the channel bed and 18 241 242 measurements on the river banks, focusing on subsurface sediments) and measurements in the laboratory (Frings et al., 2011), it was found that in 243 244 most cases (59%), the difference between is less than 0.03 (Fig. 5), with an average porosity of 0.24 obtained ex situ and 0.22 in situ. 245

246 3.2. Algorithm behavior

The behavior of the stochastic digital packing algorithm is presented in Figure 6. In order to validate the packing algorithm, comparisons were made between the measured and simulated porosity outcomes. Figure 7 clearly shows that the packing algorithm successfully captures the measured variation in porosity due to grain size distributions for each particle source. While the packing algorithm also seems to be able to mimic the measured variation due to grain shape for a given grain size distribution, providingthat the glass beads (spheres) are not taken into account (Fig. 8).

255 However, nearly all simulated porosities were systematically 256 overestimated compared to the experimental measurements. To easily recognize these discrepancies, relative errors between the measured and 257 simulated porosities were calculated (Table 2). The average relative error is 258 259 29.4% for the Rhine sediments, 21.7% for the Kall sediments and 6.6% for the glass beads, indicating that the packing algorithm predicted relatively 260 261 higher porosities when it comes to fluvial sediments with irregular shapes. Figure 9 shows the comparison between these discrepancies over the seven 262 grain size distributions. For the unimodal particle mixtures, the 263 264 discrepancies are growing as logarithmic standard deviation increases (Fig. 265 9A). For the bimodal particle mixtures, with the finer mode increasing from 30% to 70%, the discrepancies for fluvial sediments decrease while the 266 267 discrepancies for glass beads increase (Fig. 9B).

### 268 **4. Discussion**

269 The purpose of determining the porosities of the samples was twofold: first, to point out that apart from grain size, grain shape also has a clear 270 271 impact on porosity (shown in section 3.1), and second, to serve as a basis of 272 comparison for the porosities predicted from the stochastic digital packing algorithm. It is shown in section 3.2 that although the packing algorithm is 273 able to follow the experimental trend, systematic overestimation of the 274 275 porosity is noticeable, especially for the fluvial sediments. The remarkable discrepancies between can be caused by (1) measurement inaccuracies, 276 277 and/or (2) simulation inaccuracies.

#### 278 4.1. Measurement inaccuracies

For the laboratory measurements, the reading errors related to the water 279 280 levels and packing heights dominate the accuracy of outputs. The water levels were visually read to obtain the whole accumulated volumes V<sub>a</sub> with 281 282 a deviation of about 1 mm, and readings of the packing heights for gaining 283 the total volume of particle packing  $V_t$  (including pores) were achieved with an accuracy of ~3 mm. These inevitable reading errors can lead to the 284 absolute error of the porosity to be  $\sim 0.01$  for the measurements. However, 285 measured inaccuracies are small compared to the apparent differences 286 between the measured and simulated porosities, particularly for fluvial 287 288 sediments.

289 4.2. Simulation inaccuracies

# 4.2.1. Digitization inaccuracy

As can be seen in Figure 10, the arrangements of particles leave 291 292 unexpected pore spaces. One reason for this may be the digitization errors of digital objects represented at a resolution of 0.5 mm/voxel. The effect can 293 294 be supported by the fact that the porosity of 0.355 simulated for glass beads is less than the limit of 0.36 in a random dense packing of spheres (Scott, 295 296 1960; Allen, 1985; Yu and Standish, 1991; Weltje and Alberts, 2011). This is probably because the spherical shape of glass beads is not perfectly 297 described at such a resolution (0.9% digitization error), causing a reduction 298 299 of porosity. Korte and Brouwers (2013) also observed the same effects in the simulation of packing 3D digitalized spheres under different resolutions. 300 301 For this reason, a test for the ID 5 case (Table 2) was carried out with a 302 higher resolution of 0.25 mm/voxel to decrease the digitized errors,

especially for smaller particles. This gave a slightly lower porosity of 0.37
instead of 0.38 at 0.5 mm/voxel resolution, indicating that effects of
digitization errors are not too significant when compared to the
discrepancies between measured and simulated porosities.

Another error arises from the strict non-overlap requirement in the algorithm. Imagine two large objects side by side. If for any reason, there is a voxel protruded from either of the objects, this single voxel can stop the two objects from coming closer, thus leaving a large gap. In reality or in DigiDEM simulations, where forces instead of probabilities are used to determine in which direction and by how much each object moves in the next time step, this would not have happened.

314 4.2.2. Process control parameters

315 Another cause of simulation inaccuracy is the settings of process control parameters that affect the simulated packing structures, which are 316 317 rebounding probability, addition rate and number of time steps. We did a sensitivity analysis to define the effects of these parameters on porosity. 318 319 This was done by running a number of simulations in which one of the parameters was varied while keeping the others constant. To perform these 320 321 simulations, 750 spherical particles (6.4 mm in diameter) and a cylinder 322 (64mm in both diameter and height) were used. Resolution was set to 0.4 mm/voxel, giving a slight difference (<1% digitization error) between the 323 digital volumes and real volumes. 324

Rebounding probability, designed to allow particles to move upwards, provides a non-physical way to generate vertical vibrations. The original intention of having a rebounding probability is to make it possible for

328 particles to escape from their cramped places and continue to explore more suitable space to fit in, thereby simulating sediment compaction. The 329 rebounding probability can be set between 0 and 1. A value of 0 means no 330 331 rebounding and hence no vertical vibration applied. A value of 1 means particles having the same probability to move up or down, and hence kept 332 suspended. To investigate its effects on porosity, seven rebounding 333 334 probabilities varying from 0.1 to 0.7 were tested, while the addition rate and number of time steps remained the same (Table 3). The sensitivity analysis 335 336 shows that bulk porosities vary parabolically as a function of the rebounding probability (Fig. 11A). The lowest porosity values appear at rebounding 337 probabilities of 0.3-0.5, while lower and higher rebounding probabilities 338 339 give higher porosities.

340 Addition rate controls the speed of introduction of particles into the packing space. Simulations with seven fixed addition rates were performed 341 342 with the same sets of rebounding probability, and number of time steps (Table 4). Slower addition rates tend to generate denser packing structures, 343 344 with bulk porosities decreasing from 0.46 to 0.42 (Fig. 11B). This effect is because with slower addition rates, particles have more time to find a better 345 346 fitting position before being locked-in by new additions, resulting in denser 347 packing structures.

In the packing algorithm, three types of time steps are defined: normal time steps, extra time steps and wind up time steps. Normal time steps are those during which particles drop into the packing space. They are closely related to the addition rate. For example, if the addition rate is chosen such that one particle drops down every 10 time steps, 1000 normal time steps

353 are needed to introduce 100 particles into the packing space. In the case that 354 a previously introduced particle still remains on top of the container, the next particle might be prevented from being introduced. In this instance, the 355 next particle has to "wait" and extra time steps are needed to finish the 356 packing. Wind up time steps are time steps at the end of a simulation during 357 which no more particles are added and the rebounding probability is set to 358 359 zero. These time steps enable the whole packing structure to consolidate. During the sensitivity analysis, only the effect of wind up time steps on 360 361 porosity was assessed, since the effect of normal and extra time steps is directly related to the addition rate. The number of wind up time steps was 362 varied between 0 and 32000 (Table 5), and shows no systematic effect on 363 364 porosity (Fig. 11C).

The sensitivity analysis confirms that the settings we chose for the validation of the stochastic digital packing algorithm (Table 1) result in the densest possible packings. This shows that the overestimation of porosity by this packing algorithm cannot be solved by choosing different settings for the simulations.

# 370 4.2.3. Random walk-based algorithm

The reasons why the simulations failed to yield random dense packing structures can be explored in the random walk-based packing algorithm, by which the translational and rotational movements of particles during the simulation are completely random. Looking at the cross sections of the digital packings (Fig. 10) closely, the mixing of the particles is not uniform as smaller particles are more likely to concentrate at the bottom layer, particularly for the bimodal distributions with percentage of small particles 378 increasing from 30% up to 70%. The phenomenon can be interpreted by 379 kinematic sorting (i.e., segregation) effects. This is because particles kept moving randomly throughout the simulation, thus giving more chances for 380 381 smaller particles to move through the pore spaces between larger particles and reach the bottom layer. Observations from Figure 10 also suggest that 382 shape effects strongly affect the simulated packing structures of fluvial 383 sediments compared to the packings of glass beads. Because of random 384 rotational motions during the simulation, the arrangements of particles with 385 386 irregular shapes lead to create larger voids, especially between larger 387 particles. For the simulations of glass beads, shape effects are 388 inconsequential because the rotation of a sphere has no impact on particle 389 packing. Therefore, kinematic sorting can fully explain the growing 390 discrepancy trend for glass beads over the seven grain size distributions, while shape effects are the dominant reason that causes the porosity to be 391 392 significantly overestimated for fluvial sediments (Fig. 9).

### **393 5.** Conclusions

394 The applicability of a stochastic digital packing algorithm in predicting porosity of fluvial gravel deposits was validated. The conclusions are 395 396 summarized as follows: (1) Apart from grain size, grain shape has a clear 397 impact on porosity. (2) The packing algorithm provides an innovative way to simulate fluvial sediment mixtures with arbitrary shapes. (3) The packing 398 algorithm correctly reflects the mixing effect on porosity for unimodal 399 400 particle mixtures and also reproduces the differences in porosity for bimodal particle mixtures. However, in all cases, the packing algorithm 401 systematically overestimates porosity mainly due to the unwanted kinematic 402

sorting effects as well as shape effects introduced by the random motion of
particles. (4) The packing algorithm is useful for trend analysis of packing
porosity; but for a quantitative match a more rigorous model such as
Discrete Element Method (DEM) where particle motion is physics-based
may be needed.

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#### 547 **Figure captions**

- Fig. 1. Four unimodal (A, B, C, D) and three bimodal (E, F, G) grain sizedistributions used for the porosity measurements and simulations.
- 550 Fig. 2. Nine representative digitized particles in the 22.4-31.5 mm fraction
- of (A) Rhine sediments, and (B) Kall sediments, represented at a resolution
- 552 of 0.5 mm/voxel.
- **Fig. 3.** Shape properties of (A) Rhine sediments, and (B) Kall sediments in
- the Zingg classification.  $(n = 9 \times 7 = 63)$

**Fig. 4.** Measured porosity for the Rhine sediments, Kall sediments and glass beads over the four unimodal distributions represented by logarithmic standard deviation (A) and three bimodal distributions represented by percentage of fine mode (B).

- **Fig. 5.** Porosity difference between field measurements and laboratory measurements, based on the porosity data set provided by Frings et al. (2011). The study area was the 520 km long river reach between the barrage of Iffezheim (Rhine kilometer 334) and the German-Dutch border (Rhine kilometer 865).
- **Fig. 6.** Generated digital packings for (A) Rhine sediments, (B) Kall sediments, and (C) glass beads. From left to right, the packings represent the four unimodal distributions (1, 3, 5, 7 fractions), and three bimodal distributions (30%, 50%, 70% proportion of fine mode).
- Fig. 7. Comparison of model predictions with experimental data for each
  particle source over the four unimodal distributions (A, C, E) and three
  bimodal distributions (B, D, F).

Fig. 8. Comparison of model predictions with experimental data between the three different particle sources (i.e., the spherical glass beads, the subspherical Rhine sediments and the low-spherical Kall sediments) for a given grain size distribution. A to G represents the four unimodal distributions (1, 3, 5, 7 fractions), and three bimodal distributions (30%, 50%, 70% percentage of fine mode).

577 Fig. 9. Comparisons between relative errors over the four unimodal578 distributions (A), and three bimodal distributions (B).

579 Fig. 10. Cross section images of the generated digital packings for (A)

580 Rhine sediments, (B) Kall sediments, and (C) glass beads. From left to right,

the packings represent the four unimodal distributions (1, 3, 5, 7 fractions),

and three bimodal distributions (30%, 50%, 70% percentage of fine mode).

583 Fig. 11. Sensitivity analysis of process control parameters on porosity,

584 including (A) Rebounding probability, (B) Addition rate, and (C) Windup

585 timesteps. Each simulation was conducted three times and the error bar

shows 95% confidence interval for the simulated porosities.

Parameters	Values
Resolution	0.5 mm/voxel
Container diameter	104 mm
Dropping height	300 mm
Sediment density	$2650 \text{ kg/m}^3$
Glass density	$2500 \text{ kg/m}^3$
Adding source	Rain-dropping mode
Rotation	Complete random
Rebounding probability	0.35
Addition rate	1 particle/every 50 timesteps
Windup timestesps	2000

 Table 1. Set-up conditions applied in simulations

Table 2. Porosity outcomes attained from laboratory measurements and simulations (a, standard deviation; b

ID	Deserie		Laboratory Measurements					
ID	Descrip	otion of grain size distri	1#	2#	Mean	SD <sup>a</sup>		
1			1 Fraction	0.370	0.372	0.371	0.001	
2		Unimodal	3 Fractions	0.359	0.353	0.356	0.003	
3	Rhine	distributions	5 Fractions	0.346	0.342	0.344	0.002	
4	sediments		7 Fractions	0.317	0.313	0.315	0.002	
5	seaments	Dimedal	30% <sup>b</sup>	0.272	0.267	0.270	0.003	
6		Bimodal distributions	50% <sup>b</sup>	0.284	0.294	0.289	0.005	
7		distributions	70% <sup>b</sup>	0.300	0.297	0.299	0.002	
8			1 Fraction	0.383	0.380	0.382	0.002	
9		Unimodal	3 Fractions	0.385	0.380	0.383	0.003	
10	17 11	distributions	5 Fractions	0.368	0.364	0.366	0.002	
11	Kall		7 Fractions	0.331	0.324	0.328	0.004	
12	sediments		30% <sup>b</sup>	0.325	0.315	0.320	0.005	
13		Bimodal	50% <sup>b</sup>	0.316	0.317	0.317	0.001	
14		distributions	70% <sup>b</sup>	0.314	0.312	0.313	0.001	
15			1 Fraction	0.365	0.362	0.364	0.002	
16		Unimodal	3 Fractions	0.383	0.377	0.380	0.003	
17	Glass	distributions	5 Fractions	0.368	0.368	0.368	0.000	
18			7 Fractions	0.353	0.344	0.349	0.005	
19	beads	Dimodol	30% <sup>b</sup>	0.317	0.314	0.316	0.002	
20		Bimodal	50% <sup>b</sup>	0.314	0.310	0.312	0.002	
21		distributions		0.330	0.324	0.327	0.003	

ID Rebounding Probability	Addition Rate			E (m	W/' a 1	Simulated porosity					
	Amount	Every Timesteps	Normal Timesteps	Extra Timesteps	Windup Timesteps	1#	2#	3#	Mean		
1	0.1	1	10	7500	0	500	0.437	0.440	0.441	0.439	0
2	0.2	1	10	7500	0	500	0.433	0.436	0.436	0.435	0
3	0.3	1	10	7500	0	500	0.434	0.429	0.432	0.432	0
4	0.4	1	10	7500	0	500	0.434	0.429	0.430	0.431	0
5	0.5	1	10	7500	0	500	0.434	0.438	0.434	0.435	0
6	0.6	1	10	7500	0	500	0.433	0.433	0.438	0.435	0
7	0.7	1	10	7500	0	500	0.446	0.447	0.443	0.445	0

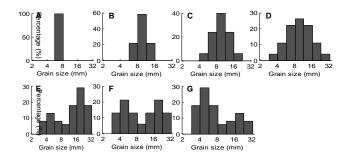
Table 3. Simulated porosity with varied rebounding probabilities (a, standard deviation)

# Table 4. Simulated porosity with varied addition rates (a, standard deviation)

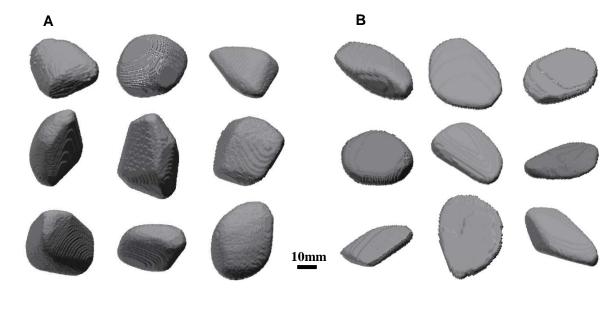
ID Rebounding Probability	Addition Rate			Eveteo	W/: d	Simulated porosity					
	Amount	Every Timesteps	Normal Timesteps	Extra Timesteps	Windup Timesteps	1#	2#	3#	Mean	4	
1	0.25	1	2	1500	0	500	0.460	0.463	0.457	0.460	0
2	0.25	1	5	3750	0	500	0.446	0.448	0.441	0.445	0
3	0.25	1	10	7500	0	500	0.434	0.432	0.434	0.433	0
4	0.25	1	20	15000	0	500	0.424	0.427	0.428	0.427	0
5	0.25	1	30	22500	0	500	0.423	0.421	0.422	0.422	0
6	0.25	1	40	30000	0	500	0.421	0.421	0.420	0.421	0
7	0.25	1	50	37500	0	500	0.420	0.420	0.421	0.420	0

**Table 5.** Simulated porosity with varied windup timesteps (a, standard deviation)

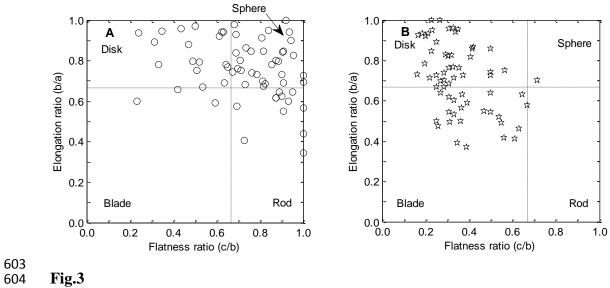
		<b>^</b>	<u>.</u>								
ID Rebounding Probability		Addition Ra	te	- Evites	Windun		Simu	ilated poi	rosity		
	Amount	Every Timesteps	Normal Timesteps	- Extra Timesteps	Windup Timesteps	1#	2#	3#	Mean		
1	0.25	1	10	7500	500	0	0.434	0.435	0.437	0.435	0
2	0.25	1	10	7500	500	1000	0.432	0.432	0.434	0.433	0
3	0.25	1	10	7500	500	2000	0.434	0.431	0.433	0.433	0
4	0.25	1	10	7500	500	4000	0.435	0.432	0.435	0.434	0
5	0.25	1	10	7500	500	8000	0.432	0.434	0.432	0.432	0
6	0.25	1	10	7500	500	16000	0.434	0.433	0.436	0.435	0
7	0.25	1	10	7500	500	32000	0.431	0.432	0.430	0.431	0
	·			<i>,</i>							-



596 Fig.1

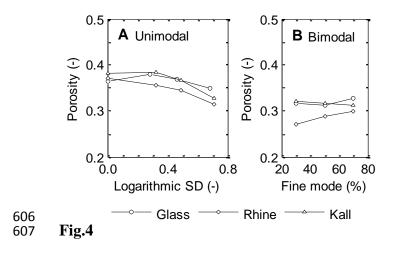


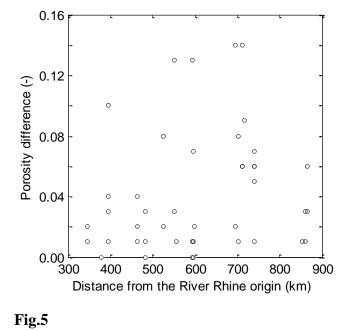
**Fig.2** 





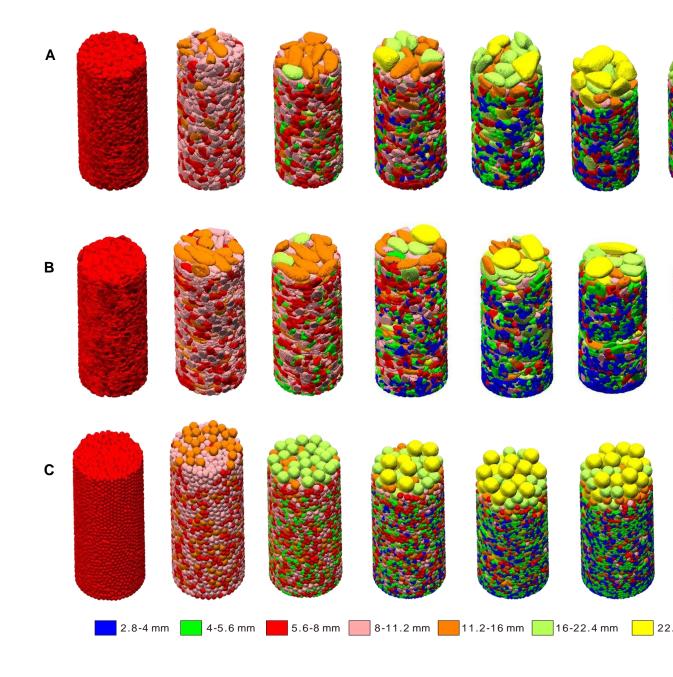




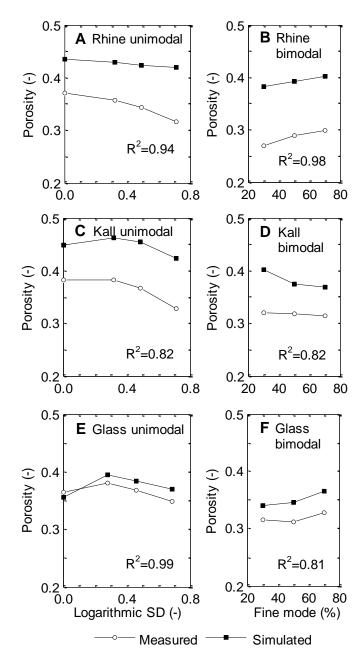






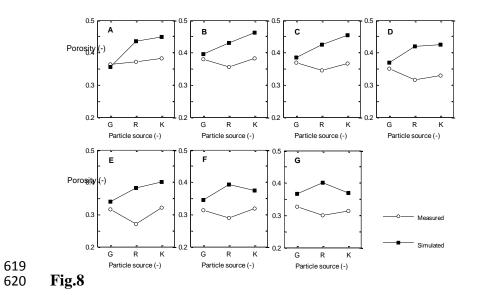


- **Fig.6**

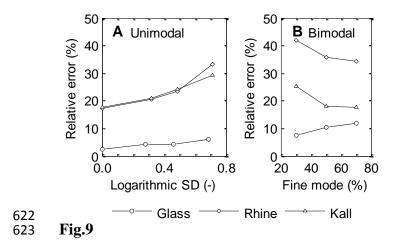




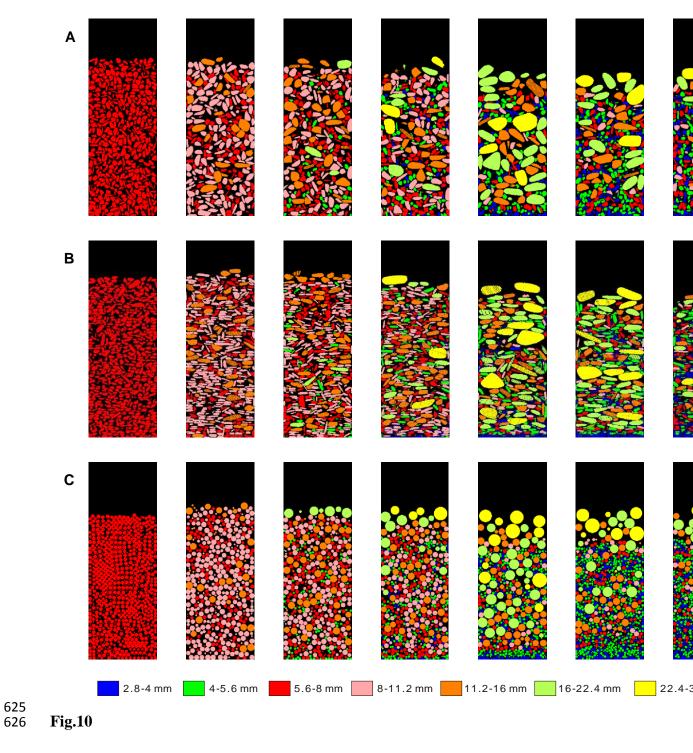


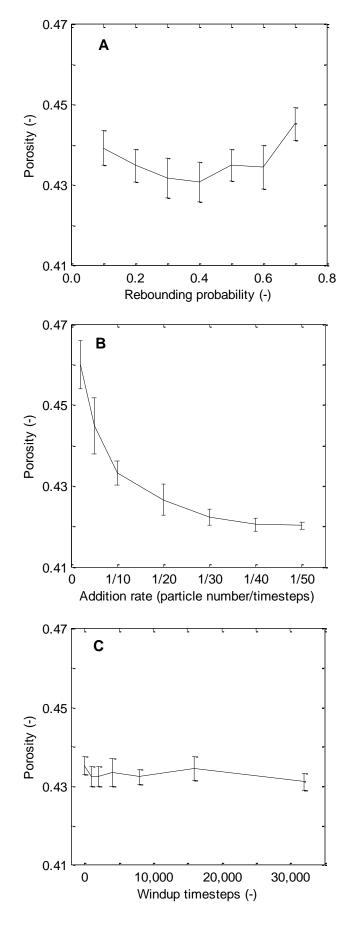












**Fig.11**