Title:

Flow modeling in Pelton turbines by an accurate Eulerian and a fast Lagrangian evaluation method.

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ABSTRACT

The recent development of Computational Fluids Dynamics (CFD) has allowed the flow modeling in impulse hydro turbines that includes complex phenomena like free surface flow, multi fluid interaction, and unsteady, time dependent flow. Some commercial and open-source CFD codes, which implement Eulerian solving methods, have been validated against experimental results showing satisfactory accuracy. Nevertheless, further improvement of the flow analysis accuracy is still a challenge, while the computational cost is very high and unaffordable for multi-parametric design optimization of the turbine's runner. In the present work, a CFD Eulerian approach is applied at first, in order to simulate the flow in the runner of a Pelton turbine model installed at the laboratory. Then, a particulate method, the Fast Lagrangian Simulation (FLS), is used for the same case, which is much faster than the Eulerian approach, and hence potentially suitable for numerical design optimization, providing that it can achieve adequate accuracy. The results of both methods for various operation conditions of the turbine, as also for modified runner and bucket designs, are presented and discussed in the paper. In all examined cases the FLS method shows very good accuracy in predicting the hydraulic efficiency of the runner, although the computed flow evolution and torque curve during the jet-runner interaction exhibit some systematic differences from the Eulerian results.

Keywords: Pelton hydro turbines; Jet-runner interaction; CFD; Eulerian modeling; Fast Lagrangian Simulation; Design optimization

1. Introduction

The design of Pelton and other impulse type turbines was based on existing know-how, and any design improvements were mainly conducted after extensive experimental testing by the trial-and-error method. In recent years significant effort has been directed towards a better understating of the details of the complex unsteady flow in the runner, with the aid of Computational Fluid Dynamics (CFD) and modern numerical modeling [1]. Commercial software have been developed and used based on Eulerian

mesh-type codes like Ansys-CFX [2, 3] and Ansys-Fluent [4, 5], that are designed to simulate the flow in various physical problems. Also, a number of non-commercial software have been developed like the Open FOAM [6], as also various Lagrangian approaches, like the Smoothed Particle Hydrodynamics method [7] or other in-house codes like the Fast Lagrangian Simulation (FLS) [8, 9] that were developed to solve specific flow problems.

At first, the Eulerian mesh-type methods were used in order to simulate the flow in Pelton turbines. Several papers have been published investigating the flow in the injector [10-12], in the stationary buckets [13], and in the rotating runner [2, 5], as well as to analyzing specific flow mechanisms like cavitation [14].

In most cases, commercial CFD software were used as it can simulate the flow with adequate accuracy, in order to calculate the hydraulic efficiency and evaluate the geometrical characteristics of the turbine. Nevertheless, the computational cost is too high to be used for multi-parametric design optimization, where thousands of flow evaluations are required to obtain an optimum design solution. Other meshless simulation methods based on the Lagrangian approach have also been developed. The most popular is the Smooth Particle Hydrodynamics (SPH) methodology [15, 16], the computational cost of which is, however, comparable to the Eulerian methods, while further development and validation is needed to improve its accuracy.

In order to reduce the computational cost, a particulate method based on the Lagrangian approach has been developed by the Laboratory of Hydraulic Turbomachines, NTUA [8, 9]. This Fast Lagrangian Simulation (FLS) method, introduces appropriate adjustable terms in the flow particle equations to approximate the various viscous and pressure effects of their trajectories.

The aim of the present paper is to evaluate the performance and accuracy of the FLS method by comparing the results against the corresponding of the Ansys-Fluent commercial software [17, 18] that uses the Volume-of-Fluid (VOF) technique to simulate the jet and free surface flow in the runner. The accuracy of the latter has been validated against experimental results in Pelton runners [19-22].

In the present study, the appropriate computational domain, mesh density and settings of the VOF technique are at first investigated to achieve the best compromise of accuracy and computational cost. Then, an initial test case is defined including the exact geometrical characteristics of a Pelton model runner installed in the laboratory. The FLS and Fluent software are applied to various operating condition cases of this turbine, as well as for runners of different design, in order to compare their numerical performance based on the time history of torque development and on the hydraulic efficiency of the runner.

2. Numerical modeling of the flow

Two computer codes were used for the simulation of the complex flow in a rotating Pelton runner, the commercial software Fluent and the in-house software FLS. The Fluent code [17] is capable of solving multiphase flow problems with free surfaces, which are highly relevant to impulse turbines. It is based on the spatial discretization of the Reynolds Averaged Navier-Stokes equations on a computational mesh using cell-centered numerics (finite volumes), and offers flexibility in choosing between segregated based and coupled based solver.

The two fluid flow (air and water) problem with free surface, like the Pelton runner case, is being solved using the established Volume of Fluid (VOF) method. VOF is an Eulerian-Eulerian method based on tracking and locating the fluid-fluid interface. An additional factor, the volume fraction, is introduced which represents the percentage of each fluid volume in every cell. The method can model two or more immiscible fluids by solving a single set of momentum equations and tracking the volume fraction of each of the fluids throughout the domain.

For the two-phase (water-air) flow in Pelton runners the air is defined as the primary phase and the water as secondary, so as the tracking of the interface between them is accomplished by the solution of a continuity equation for the volume fraction of the secondary phase that has the following form:

$$\frac{1}{\rho_q}\left[\frac{\partial}{\partial t}\left(a_q\rho_q\right) + \nabla\cdot\left(a_q\rho_q\vec{v}_q\right)\right] = \sum_{p=1}^{n}\left(\dot{m}_{pq} - \dot{m}_{qp}\right) \quad (1)$$

where $\dot{m}_{pq}$ is the mass transfer from phase p to phase q and $\dot{m}_{qp}$ the opposite, a_q is the volume fraction and ρ the density. The present application of Fluent-VOF software for flow simulation in Pelton runners is analyzed below, in section 4.

2.1. Fast Lagrangian Simulation method

The Fast Lagrangian Simulation (FLS) [8, 9] is a single-phase flow simulation method, and it is based on the tracking of a representative number of fluid particles in order to model and calculate the flow pattern and the energy exchange in the rotating runner of impulse hydro turbines. The exact water-air interface pattern is not required for such computations, though the flow width on the bucket surface could be estimated from the properties of the tracked particles.

The water jet is being separated into discrete particles, as shown in Fig. 1, and their equations of motion are numerically integrated until they exit from the inner bucket surface. The jet is considered ideal, i.e. with uniform initial velocity, and the flow frictionless. Also, due to the periodic symmetry conditions only two consecutive buckets were modeled.

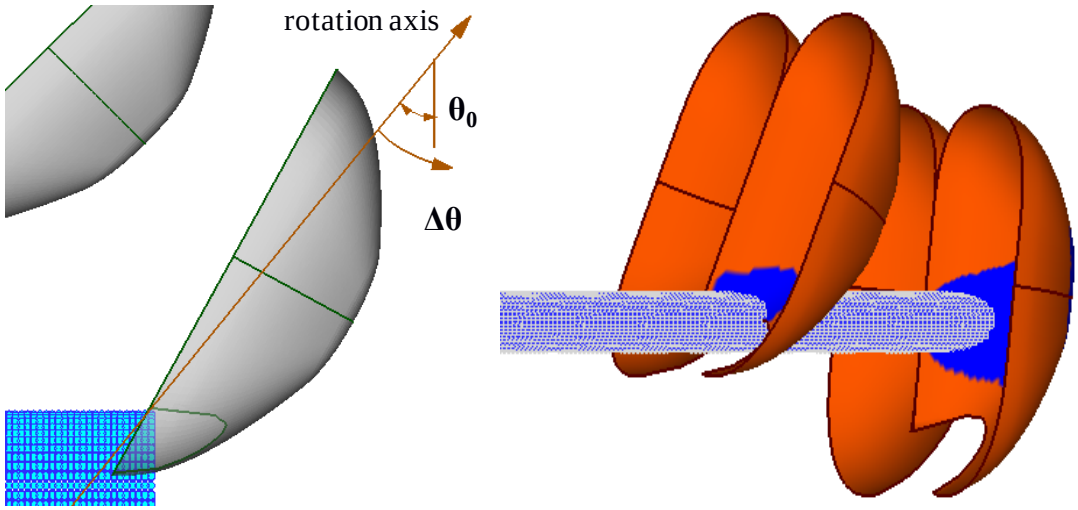


Fig. 1. FLS modeling: jet discretization and jet-bucket interaction starting angle (left); surface flow simulation and representation (right).

The fluid particle equations are solved in a rotating orthogonal system of reference, and are expressed in Cartesian coordinates as follows (bucket rim is on the xy level):

$$\begin{aligned}\frac{d^2x}{dt^2} &= f_x(x,y)\\ \frac{d^2y}{dt^2} &= f_y(x,y)+\omega^2\, y+2\,\omega\frac{dz}{dt}\\ \frac{d^2z}{dt^2} &= f_z(x,y)+\omega^2\, z-2\,\omega\frac{dy}{dt}\end{aligned} \qquad (2)$$

where ω is the angular rotation speed of the runner and f_x, f_y, and f_z are the additional terms, functions of the local surface geometrical characteristics.

The particle motion equations do not contain particle interaction or mechanical losses terms and hence they cannot reproduce the real flow picture in the bucket. For this reason, the FLS model introduces a number of additional terms in order to account for the various hydraulic losses (impact, friction, change direction), as well as for the pressure effects that control the spreading of the surface flow in the bucket, which are activated after the impact of a particle on the bucket surface. More specifically:

The friction losses on the bucket surface are modeled as reduction of particle kinetic energy by a factor analogous to the square of particle velocity and to sliding distance, hence the particle velocity magnitude after a time step Δt becomes:

$$V_p' \approx V_p \cdot \left(1 - C_f \cdot V_p \cdot \Delta t\right) \qquad (3)$$

The energy losses at the jet impact on the bucket are taken analogous to the square of the normal to the surface particle velocity component, and this gives:

$$V_p' \approx V_p \cdot \left(1 - C_i \cdot \cos^2 \varphi_i\right) \qquad (4)$$

where φ_i is the particle impingement angle.

The progressive change of particle's path direction (and momentum) as it slides along the curved bucket surface also causes minor energy losses, which are modeled using a similar to the impact losses term:

$$V_p' \approx V_p \cdot \left(1 - C_p \cdot \cos^2 \Delta\varphi\right) \qquad (5)$$

where $\Delta\varphi$ is the angular change in direction of the sliding particle during the time step Δt.

The adjustable coefficients C_f, C_i, and C_p are introduced in the above modelling equations (3-5) so as they can be tuned in order to match the results with the corresponding ones obtained by more accurate CFD methods or by experimental studies. This constitutes an important feature of the FLS model that can significantly improve the reliability and accuracy of its results, at almost no additional computer cost.

Finally, in order to model the pressure effects on the surface flow spreading and evolution, a particle acquires at the impact point an artificial "spreading" velocity component perpendicular to its impacting plane ($\vec{V}_S$, Fig. 2), while its main velocity component is correspondingly reduced to preserve kinetic

energy. The magnitude of this spreading velocity depends on the radial and circumferential position of a particle in the jet cross section, and it is determined by two corresponding additional adjustable coefficients [8].

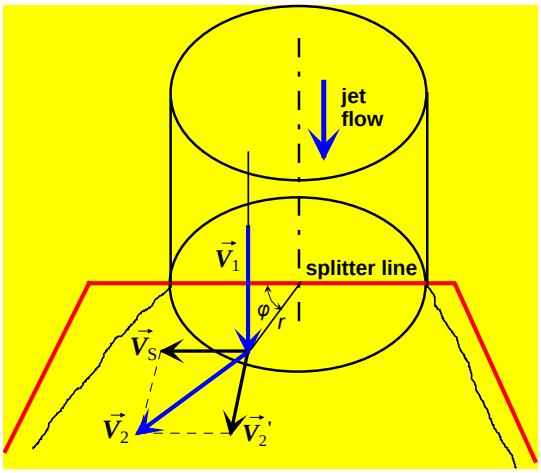


Fig. 2. Sketch of surface flow spreading modeling.

The first tuning of the above coefficients has been carried out with the aid of available experimental data in a Pelton turbine model, so as to minimize the squared differences between the measured and the computed characteristic curves of turbine efficiency [8]. This problem was solved with an optimization software based on evolutionary algorithms, considering these coefficients as free variables, the values of which are iteratively optimized from generation to generation.

It must be noted that some secondary flow mechanisms developed during the jet-runner interaction, like the jet cut or the attached flow at the back face of the bucket are not modeled by the FLS method. However, their effect on the energy transfer and hydraulic efficiency of the runner are taken into account in an implicit manner through the above adjustable terms. More details on the model formulation and its regulation procedure, as well as on the parametric design of the runner can be found in [8].

The application of the FLS model for the simulation of the jet-runner interaction in Pelton turbines requires only a mathematical/numerical description of the inner bucket surface, and since it does not use a computational mesh, it can be very easily adapted to any design or operation data of the runner. Also, a specific post-processing algorithm is developed and used for the presentation and comparison of FLS results (e.g. Fig. 1).

3. Test case

The geometrical characteristics of the Pelton model turbine used as a reference case corresponds to a Pelton turbine installed in the LHT, at the National Technical University of Athens (Fig. 3) [23]. The pitch diameter of the runner is 400 mm, and the axis is horizontal with two injectors of 36 mm nozzle

diameter. The net head, the rotational speed and the nozzle stroke can be adjusted for experimental reasons according to the IEC standards [24]. The runner contains 22 buckets which were designed and constructed in the lab, based on old literature guidelines [25], as shown in Fig. 4. Therefore the achieved efficiency of the turbine is smaller than the suggested in more recent bibliography [26], which refers to the state of the art turbines. Nevertheless, all geometrical characteristics of modern runners and buckets are included, and hence the flow mechanisms that take place during the interaction of the free jet with the runner are the same.

Fig. 3. Pelton model turbine and its runner installed at the LHT

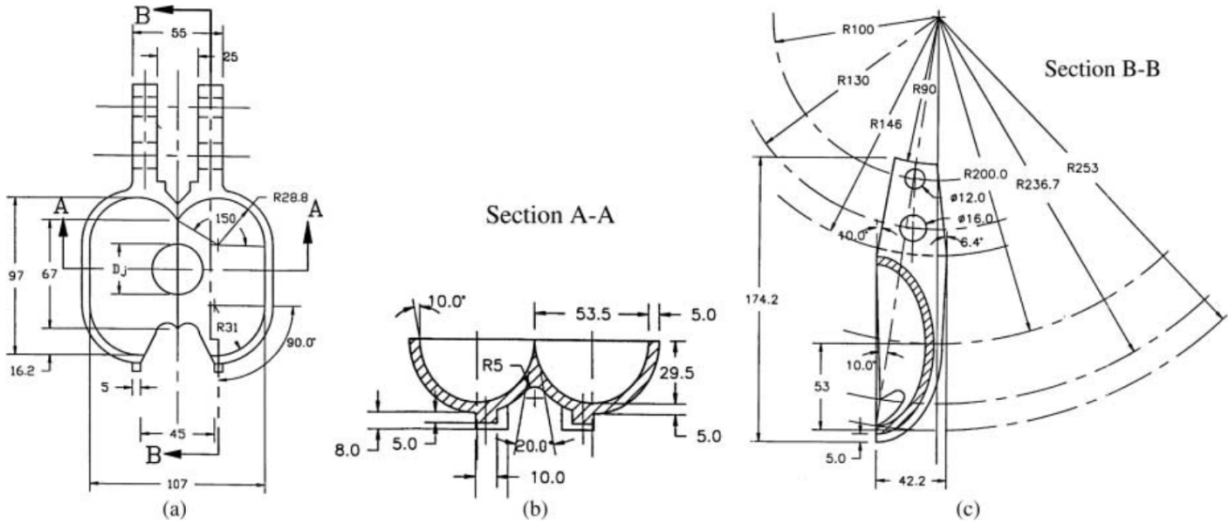


Fig. 4. Details of the runner design [23]

Although the Pelton turbine is installed at the laboratory test rig and the total efficiency of the turbine has been measured experimentally, a direct comparison between numerical and experimental results is not easy. The measured efficiency represents the total efficiency of the turbine including all losses according to the IEC standards [24], while only the hydraulic efficiency of the runner can be calculated numerically. So, the comparison would require the estimation of minor losses like the losses in the nozzle, the mechanical losses, and the windage losses caused by the movement of the runner and its interaction with the misty environment into the casing. These losses are highly depended on the quality of the turbine construction, the geometrical characteristics of the casing and the type of the bearings. So, their estimation would introduce considerable uncertainty to an experimentally derived efficiency of the runner. For this reason, the present work focuses on the comparison and evaluation of the performance and accuracy of the two software tools, when used for Pelton turbines analysis and design. The provided detailed geometric data of the runner (Fig. 4), can be used as benchmark for the validation and evaluation of other numerical modelling tools and methods.

The comparison of FLS and Fluent results was initially performed for the reference operating conditions of the Pelton turbine: One injector is operating with nozzle stroke 12 mm, which corresponds to the best efficient point (BEP) of the runner with one injector, according to the experimental results. The rotation speed was 1000 rpm, the diameter of the jet was 12 mm and its axial velocity was uniform and equal to 44.45 m/s. For modeling purposes, the velocity profile of the jet is taken uniform.

4. Accurate Eulerian flow evaluation

The basic steps for the numerical simulation with a mesh-type Eulerian method of Fluent software are the design of the appropriate 3D geometry, the construction of the grid covering the whole domain, and the solution of the system of flow equations. The simulated geometry is designed using the SolidWorks 12 commercial software, and it was essential to avoid very small edges and narrow faces or volumes. The geometry was separated into two domains, the rotating domain, which contains the buckets, and the stationary domain from where the free jet is injected. Due to the heavy computing requirements and thanks to the periodic symmetry of the runner, only two consecutive buckets were modeled (Fig. 5). Also, only the half symmetric part of the domain was considered, a common practice in modeling Pelton turbines.

4.1. Torque and efficiency calculation

The hydrodynamic torque and the hydraulic efficiency of the runner are computed after completing the evaluation of a jet-bucket interaction flow, starting from the moment of impingement until the evacuation of the bucket.

During the unsteady flow simulation the total torque on a bucket is calculated at every time step (or its tangential position φ) by adding the torque on the inner surface of the first bucket to the torque on the backside of the second bucket, as shown in Fig. 5. The latter is developed due to the adherence of a jet portion after its cut by the bucket (Coanda effect), as also due to possible interference between the back surface of the bucket and the outflow water from the first bucket.

$$T_r(\varphi) = T_{in}(\varphi) + T_b\left(\varphi + \frac{360}{N_b}\right) \quad (6)$$

where T_r, the torque in a single bucket, T_{in} and T_b the torque at the inner and backside surfaces, respectively, and N_b the number of buckets on the runner (22 for this case).

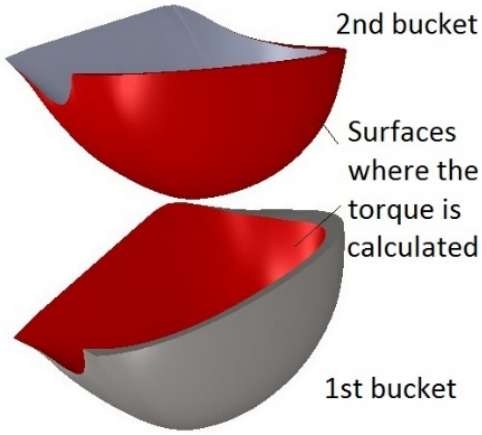


Fig. 5. Two consecutive half buckets used to calculate the total runner torque.

The torque curve of one bucket against the rotating angle is repeated periodically every 360°/N_b, where N_b is the total number of buckets (22 in this case). The sum of the torque values of all curves for every angular position represents the torque curve of the runner. So the total mechanical energy transferred to the shaft during one rotation can be calculated from the equation:

$$W = \int_0^{360} T(\varphi) d\varphi \quad (7)$$

where $T(\varphi)$ is the torque on a single bucket and φ the angular position of the runner ($\varphi = \omega \cdot \Delta t$, Δt the time step and ω the angular speed).

The runner power is then calculated as:

$$P = \frac{W \cdot \omega}{360} \quad (8)$$

while the power of the water jet is:

$$P_w = \frac{\pi \rho}{8} D_{jet}^2 u_{jet}^3 \quad (9)$$

Where ρ is the water density and u_{jet}, D_{jet}, the velocity and diameter of the jet.

Finally, the hydraulic efficiency of the runner is defined as:

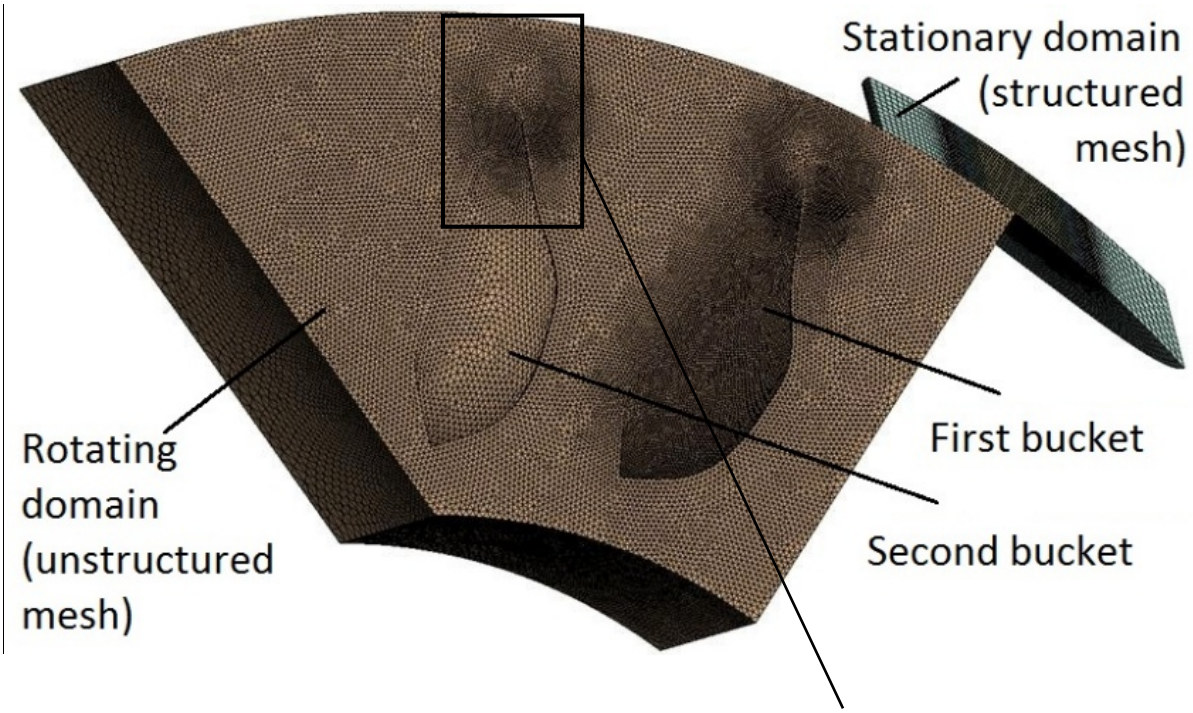
$$\eta = \frac{P}{P_w} \quad (10)$$

4.2. Computational grid and accuracy of results

The computational grid was made using the Ansys Workbench commercial software package after an extensive investigation due to its strong influence on the final speed of the simulation and the accuracy of the results. The elements around the rotating buckets and the cylindrical jet were tetrahedral and hexahedral respectively, as shown in Fig. 6. In addition, an inflation layer was adjusted at the inner and back side faces of the buckets in order to reduce the y+ value below 50 when the k-ω turbulence model is used. The inflation consists of 5 layers of hexahedral cells with very small thickness placed at the region where the boundary layer is being developed.

The quality of the mesh is acceptable as the minimum orthogonal quality is 0.15 while the average is 0.87. Since the position of the jet changes transiently it is not possible to refine the shear layer therefore the entire mesh had to be fine enough to have minimum impact on the simulated efficiency. In addition, as the unstructured tetrahedral mesh type can decrease the accuracy of the results, a very dense mesh was used at the area of high importance. The minimum and maximum cell volume is 1.1×10^{-12} m^3 and 1.5×10^{-8} m^3, respectively.

Moreover, a study about the appropriate mesh size was carried out in order to achieve independency between the size of the mesh and the accuracy of the results. The same case with the same settings was simulated for four different sizes of tetrahedral mesh. The densest mesh consists of about 5.16 million elements and it was found that the maximum possible accuracy in terms of the calculated efficiency was achieved. The efficiency dependence on the mesh density is represented in Fig. 7.



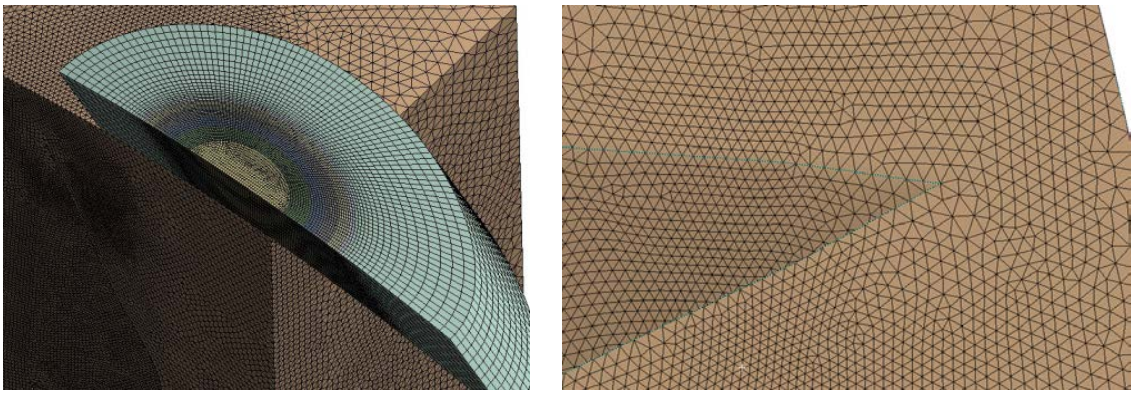


Fig. 6. Computational domain and structure of the 3D mesh

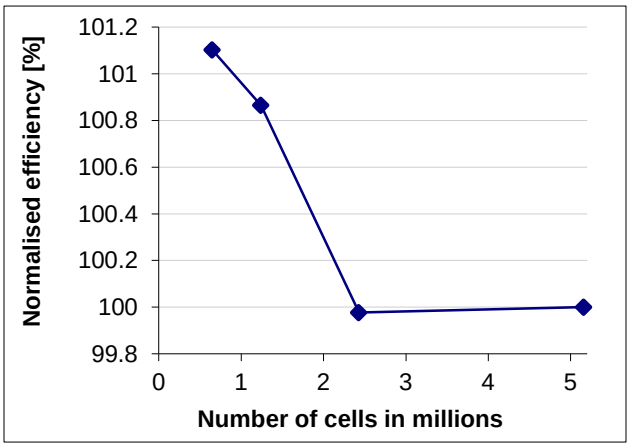


Fig. 7. Runner efficiency response to mesh refinement

From the above results it can be deduced that the independency was achieved for meshes with more than 2.4M cells. For the present simulations, a mesh with approximately 2.8M cells was adopted to ensure the accuracy of the results. The simulations were carried out using a four core Intel Xeon, 3.4GHz and 16GB memory RAM. The CPU-time was almost proportional to the size of the mesh and for the 2.4M cells is about 3.2 days.

The torque variation curves on a bucket obtained using different density meshes are presented in Fig. 8, where the angular position of the bucket is taken zero at its vertical or normal to the jet position (Fig. 1). At the start of jet-runner interaction (-40^o) a negative torque can be observed caused by the interaction of the jet with the back surface of the bucket. At this stage the jet is starting entering the bucket moving almost in parallel with the back surface, thus causing there a high pressure region, as shown in Fig. 9, responsible for the negative torque. Just afterwards, the torque is increasing as more water interacts with the inner surface. At an angle of about -22 degrees the second bucket begins to interact with the jet, leaving less water to move towards the first bucket. This reduced amount of water that moves towards the first bucket is not smooth due to divergence of the jet caused by Coanda effect (Fig. 10). This diverged portion of water hits the first bucket at around -10 degrees, causing an irregular increment

of the torque (Fig. 8). After this point, the developed torque decreases smoothly as the water is leaving the bucket, until evacuation completes at +30 degrees. The total jet-runner interaction period lasts about 70^{o} rotational angle of the runner.

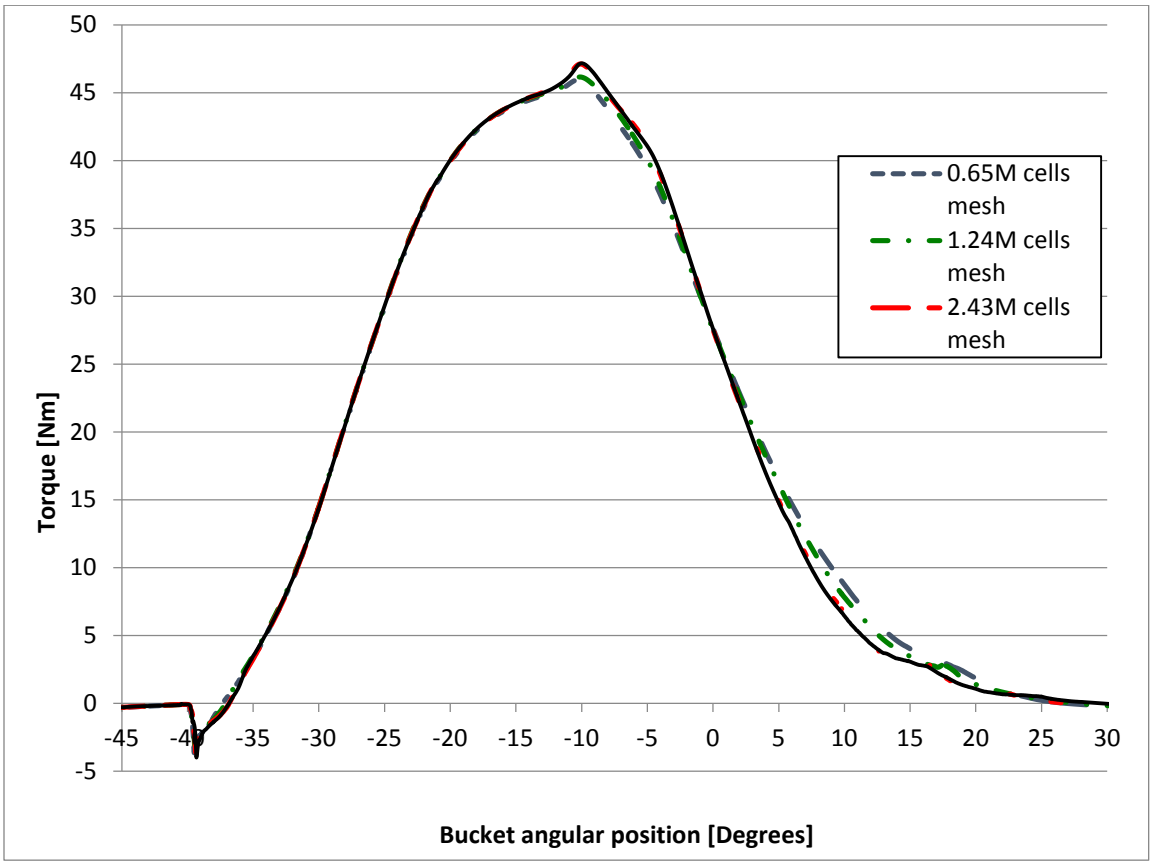


Fig. 8. Individual torque curves of mesh refinement simulations

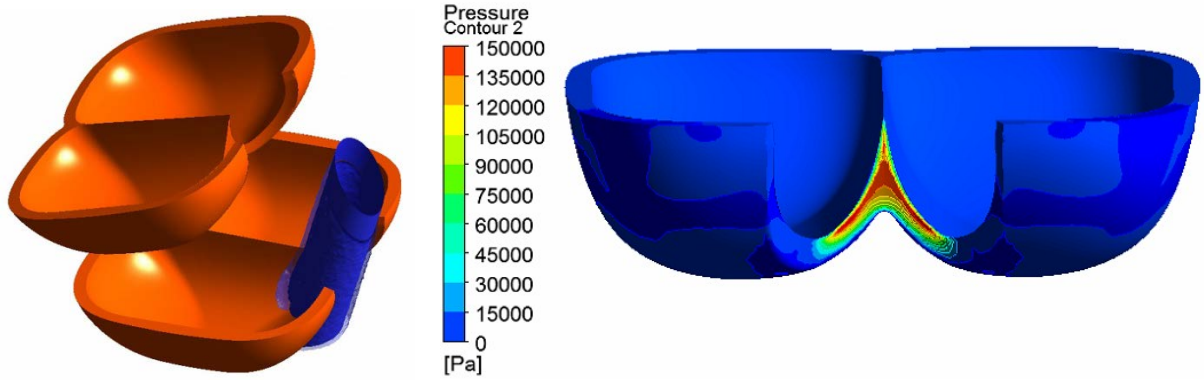


Fig. 9. High pressure development at the backside while the jet is entering the bucket

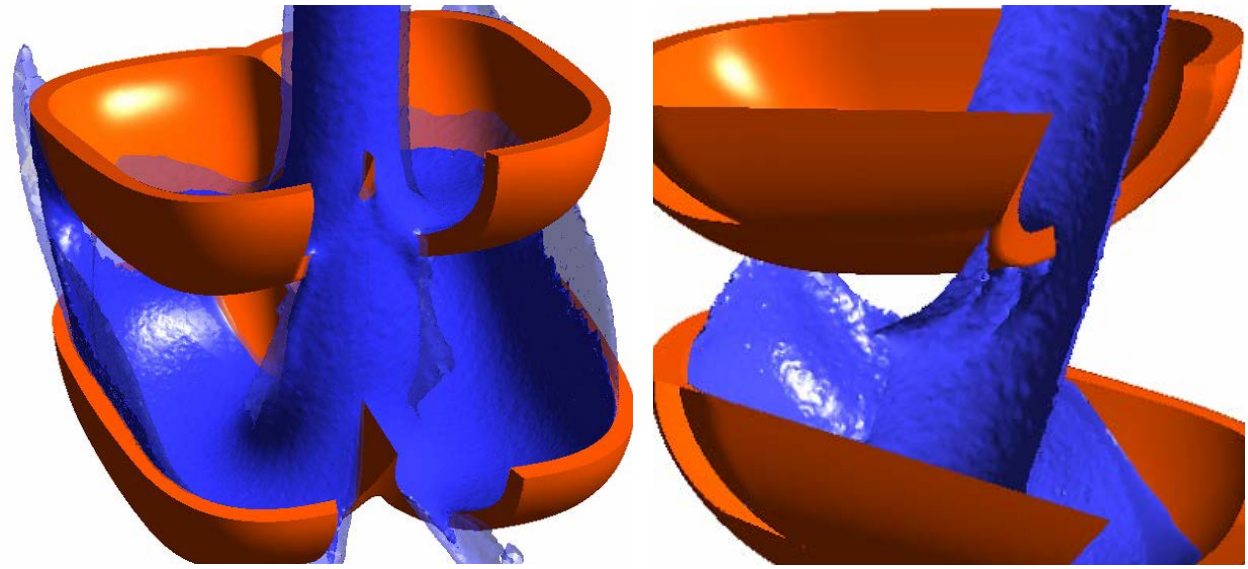


Fig. 10. Diverged flow moving towards the first bucket

4.3. Turbulence modeling and computational details

During the simulations using Fluent many different settings were tested, some of which had strong influence on the results. The 3D single precision solver was selected in order to reduce the computational cost as no important difference was observed compared to double precision simulation. The appropriate model for the simulation is the *Volume of fluid* (VOF), with pressure-velocity coupling solution method. The Pressure-based *coupled* scheme increases the computation time but provides a much more stable and reliable solution compared to the *simple* or *piso* schemes. The transient flow is simulated both as inviscid and as turbulent using the k-ω SST model, in order to assess the impact of the friction and turbulent losses on the efficiency. The time step was constant and equal to 5 μs, which corresponds to 0.03 degrees of runner rotation angle, while the maximum number of iterations was set to 10, as suggested in the literature [18]. The solution was not influenced by the time-step as soon as the continuity residuals were under 10^{-3} even if most of the time during simulation they were close to 10^{-4}, where the residuals target was set. The continuity residuals target for convergence was set to 10^{-4}, while the rest residuals were always below 10^{-5}.

The highest flow velocities in the computed field correspond to the free jet velocity, which is about 45 m/s. In addition, the coupled implicit scheme was used, and therefore the Courant number had a minor effect since the point Gauss-Seidel scheme used is unconditionally stable, according to the linear stability theory. Consequently, the default value of 200 was used while different values have negligible influence in the solution.

Finally, the possible effect of surface tension was considered by using the "continuum surface force" model, which is available in Fluent platform. However, and in agreement with the literature [2], the implementation of this model has negligible effect on the simulated flow in a Pelton runner. Moreover, the Weber number of the relatively high velocity flow in the bucket is large (> 50), and only at the last evacuation stage, where the remaining water forms very thin and separated films on the bucket surface,

the Weber number may be considerably reduced. But the contribution of this stage to the energy exchange and the developed torque is minor. Consequently, the surface tension effects are ignored in the simulations.

A comparison between the torque curves calculated with inviscid and k-ω SST turbulence model is shown in Fig. 11. As can be observed in this figure, the developed torque is slightly lower for the turbulent flow simulation, while the corresponding efficiency is about 3.5% reduced (84.8% compared to 88.3% of the inviscid flow). This is mainly due to the friction losses at the flow boundary layer along the inner bucket surface, and the impact losses of the entering jet flow.

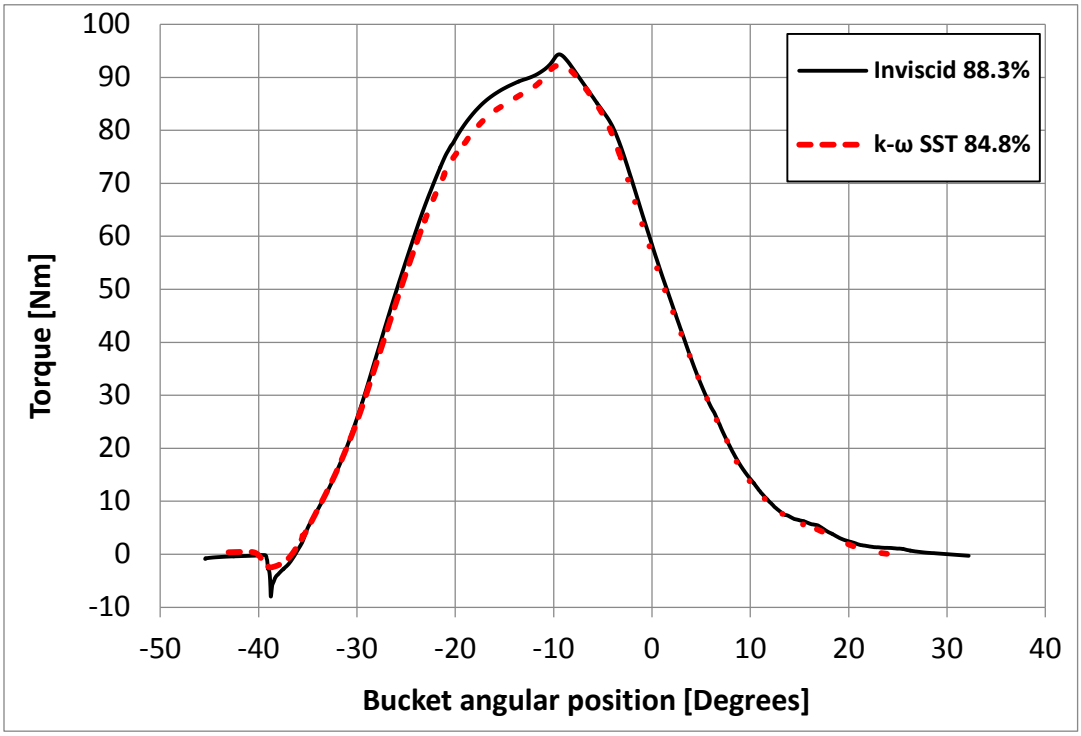


Fig. 11. Comparison of inviscid and turbulent flow numerical results

5. FLS method application and comparison with VOF

Tracking of about 5000 fluid particle trajectories using a time step for integration of 2×10^{-5} sec was found to produce statistically accurate results for all cases examined in the present study. A complete flow evaluation requires about 10 CPU sec in a modern PC, which is almost 4 orders of magnitude less than the corresponding evaluation time using Fluent. This significant advantage of the FLS method allows its application for multi-parametric design optimization studies of impulse turbine runners, in conjunction with modern optimization software.

After completion of the fluid particles tracking, the FLS post-processing algorithm computes the runner performance in terms of the developed torque on the blades and the hydraulic efficiency of the runner, while it can also calculate the local forces exerted on the blade during the energy conversion procedure. The mechanical energy transferred to each bucket is obtained from the equation of conservation of angular momentum [8]:

$$W = \rho Q_u \left(R_{run}\, u_{jet} - \frac{1}{N} \sum_i r_i w_i \right) \tag{11}$$

where Q_u is the cumulative flow that enters each bucket, R_{run} is the runner pitch radius, and w_i the tangential velocity component of particle *i* at the moment it exits the bucket at radial distance r_i. *N* is the total number of fluid particles that interact with a single bucket. The hydraulic efficiency of the runner can then be obtained as the ratio of the developed mechanical power divided by the corresponding net hydraulic power at the inlet, as in Eq. (10).

The performance and accuracy of the FLS method is evaluated by comparing its results with the corresponding ones of Fluent software, for a number of different cases. The comparison is based on the computed values of the runner hydraulic efficiency, the pattern of the torque curve, and the distribution of the surface flow on the inner bucket face. At first, the comparison is carried out for the design point of the laboratory model turbine, for which the coefficients of the FLS model are calibrated.

Next, two cases with different flow rates, one lower and one higher, were simulated maintaining the remaining parameters constant. An additional simulation took place for different head and rotating speed of the turbine that corresponds to a smaller specific speed of the runner. Also, in another test case the geometry of the runner was changed by reducing the number of buckets from 22 to 18. Finally, a comparison is made for another Pelton runner with different bucket design.

5.1. Comparison at the reference test point

Firstly, the computational results of the reference test point (BEP) using FLS and Fluent are being compared with each other. In Fig. 12 the resulted total torque in one bucket is compared. It can be observed that the agreement between the curves is good, as the general shape and the maximum torque are the same. The small differences can be explained due to the inability of the FLS method to model secondary mechanisms, like the previously mentioned negative torque at the beginning of the interaction with the jet, and the Coanda effect, thus resulting in a smoother torque curve. Another minor phenomenon is a small radial deviation of the jet caused when it is cut at the bucket tip, as shown in Fig. 13. The amount of deviated water hits the next bucket at a slightly different time (earlier), and this may explain the slightly higher torque values of the VOF in the increasing part of curve of Fig. 12.

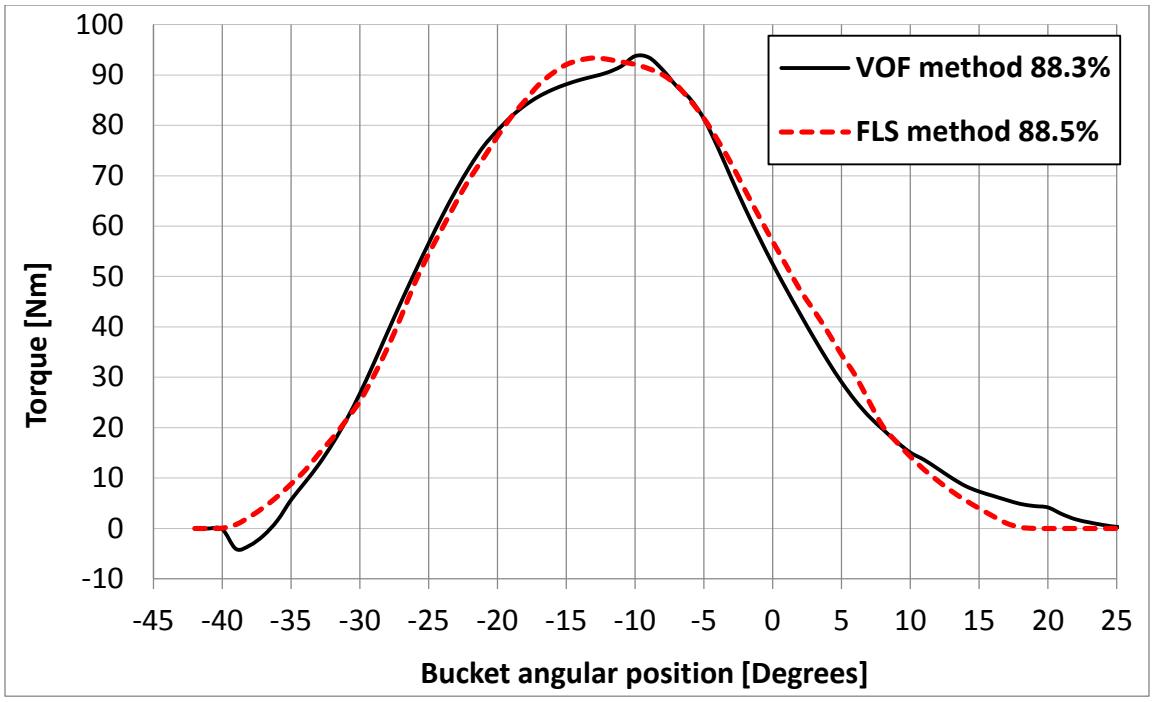


Fig. 12. Comparison of the computational results for the reference point

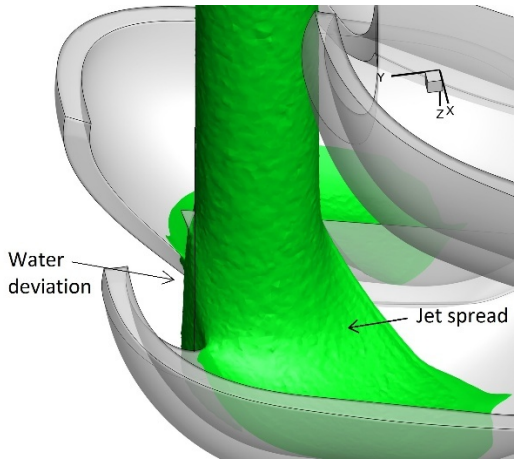


Fig. 13. Water deviation close to the solid surface calculated with Fluent

In addition to the torque results, the spreading rate of the surface flow on the inner bucket surface is also considered in order to adjust the spreading model coefficients of the FLS method. In Fig. 14 an iso-surface is being represented with green color as it is calculated by Fluent, whereas the black lines are the orbits of the particles obtained by FLS. It can be observed that the surface flow pattern on the solid surface calculated from Fluent and FLS is quite similar, but the spreading rate shows certain differences. After an analysis of the particles' orbits and the velocity fields obtained by the two methods it was ascertained that two main reasons cause this difference.

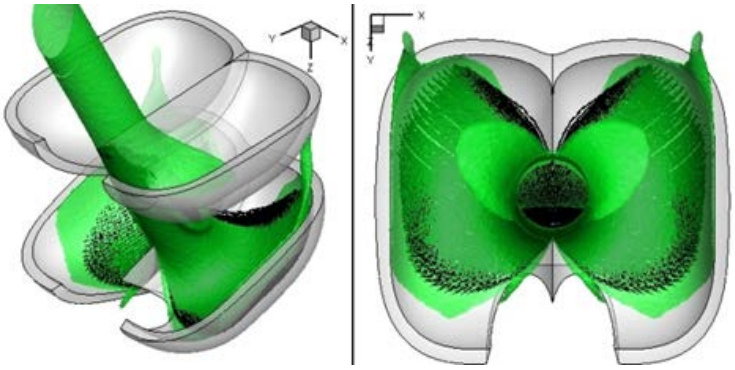


Fig. 14. Water phase comparison of FLS (black lines) and Fluent (green areas) simulations

Firstly, the jet flow as computed by VOF, exhibits a significant spread just before it reaches the bucket (Fig. 13), due to the high pressure field developed in the impact area below. Consequently, the flow path lines follow a smooth curved orbit before reach at the bucket surface, with the curvature radius being analogous to the radial position in respect to the jet axis. On the other hand, all flow particles of the FLS method move straight until they hit the inner bucket surface, and then change direction instantly to flow parallel to the bucket surface. Therefore, during the first interaction stages a flow particle modeled by FLS travels a longer distance to reach to a particular position on the surface than the corresponding one simulated by Fluent. Moreover, the small impact losses of the flow are accounted for at the impact point thus reducing the particles' velocity when they hit the bucket, whereas this is taking place more progressively with Fluent simulation.

For the above reasons, the spreading rate of the surface flow is computed lower of the FLS at the first jet-runner interaction stages, while the outflow as it is being calculated from Fluent starts earlier (Fig. 14). To compensate for this difference, the variation limits of the coefficients of the spreading model are properly regulated, by using proper constraints during the optimization procedure of the set of FLS coefficients for the reference case, in order to match also as best as possible both the torque variation curve and the hydraulic efficiency. As a result, the subsequent interaction stage becomes faster with the FLS simulation, and the evacuation of the bucket happens a little earlier (about 7^{o} in Fig. 12).

In spite of the inability of FLS method to simulate accurately all flow details and mechanisms, the resulting hydraulic efficiency of the runner is almost equal to that of Fluent (88.3% compared to 88.5%, Fig. 12). This confirms the capability of the regulated FLS model to reproduce satisfactory the main flow characteristics (e.g. torque curve), and also to implicitly take into account the effects of all secondary flow mechanisms on the efficiency.

5.2. FLS performance for different turbine loading

In order to evaluate further the FLS method performance, two test cases at different operating points of the turbine are simulated, by changing the flow rate (loading) of the runner. All other design and operation characteristics remain the same except of the jet diameter, which depends on the nozzle

opening (or spear valve stroke) of the injector. The first case is for smaller nozzle stroke which corresponds to 50% of the design point flow, while the second for higher nozzle stroke and 135% of the design flow. The torque curves and the runner efficiency values computed by VOF and FLS methods are presented in Fig. 15.

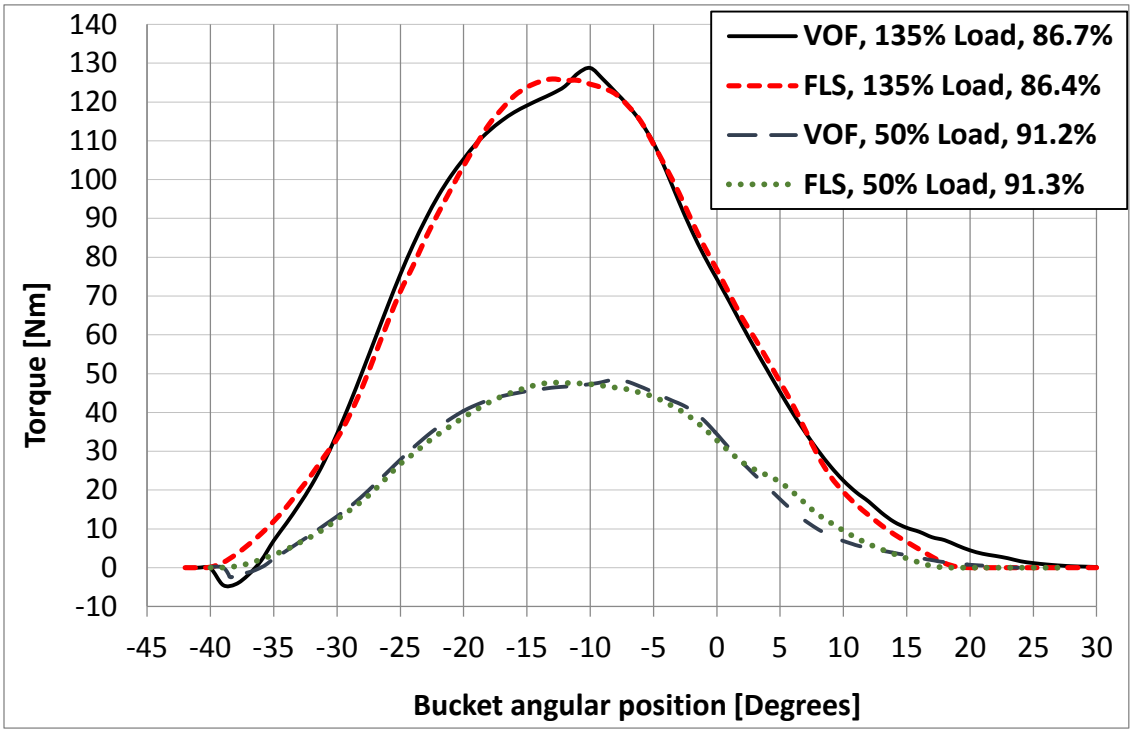


Fig. 15. Comparison of VOF and FLS method for off design turbine operation conditions

The shape of the torque curves shown in this figure is similar to those at the reference point (Fig. 12), but the maximum torque is, as expected, analogous to the loading of the runner. The agreement between the FLS and Fluent results is also good in both cases, with the differences of the FLS curves being again at the same regions (non-negative torque at the beginning, smoother curve, slightly displaced to the right, and faster bucket evacuation).

Moreover, in both cases the calculated efficiency by the FLS method are again very similar to that of the VOF method, with the differences being 0.1% and 0.3% (Fig, 15), namely within the order of numerical solution accuracy of Fluent software. It is noted that the increased efficiency of the runner at smaller flow rates is due to the inviscid simulation of the flow, in which the increased friction losses that have a thinner free surface flow along the bucket are not accounted. On the other hand, the obtained reduction of the efficiency at higher than the design flow rate is due to a different mechanism: During the evolution of the free surface flow on the bucket surface a portion of water may leave the bucket through the cut edge area, as shown in Fig. 14, but not with optimum (minimum) outflow velocity, thus causing a small drop in overall efficiency of the runner. This phenomenon becomes more intense for higher flow rates, due to the wider spread of the surface flow on the bucket.

5.3. FLS performance for different hydraulic head

In this case, the inlet water pressure was increased by 56% (from 100 m of the reference point), which corresponds to exit jet velocity of 55.7 m/s. The higher jet velocity requires higher rotating speed of the runner, which is calculated equal to 1250 rpm in order to maintain the same jet/runner speed ratio as in the reference point. The torque curves calculated by FLS and VOF methods are shown in Fig. 16. The shape of the curves and the calculated efficiencies are similar with that corresponding at the reference point, while the only differences are the higher values of torque caused due to the increased kinetic energy of the jet. The agreement between FLS and Fluent results is again very good, with the same small systematic differences in the pattern of torque curve, and runner efficiency that differs only 0.3 percentage units (Fig. 16).

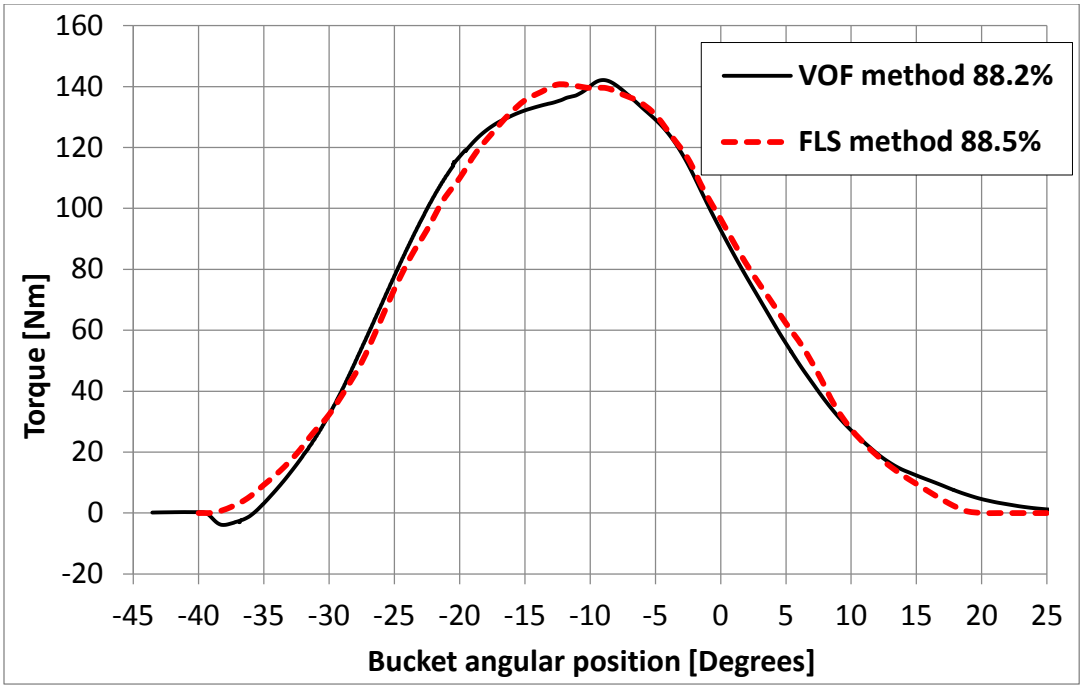


Fig. 16. Comparison of VOF and FLS methods for higher hydraulic head

5.4. FLS performance for modified runner design

Finally, a modified runner design was obtained by reducing the number of buckets from 22 to 18. In this case the angular distance between two consecutive buckets becomes 22.2% larger, while the amount of water jet that interacts with each bucket is also equally increased. This affects mainly the latest stage of jet-runner interaction, which becomes longer. As a result, the last part of outflow leaves the bucket lips at less optimum velocity, and in addition, more water exits from the cutout area. For this reason, the hydraulic efficiency of the runner becomes substantially lower, almost 2.5 percentage units as computed by Fluent (from 88.3% in Fig. 12 to 85.9%).

This efficiency reduction is again reproduced well by the FLS method, as shown in Fig. 17. As can be observed in this graph, the shape of the torque curves is also similar to the previous cases, but the deviation of FLS and Fluent values in the higher torque area (-20 to +5 degrees, Fig. 17) becomes more pronounced.

These results support the previous discussion, according to which the irregularity of the VOF curve in this area was attributed to the Coanda effect and the impact of the back flow from the previous bucket. The peak of the torque curve is now displaced to higher angular positions of the bucket, at about -3^{o}, compared to about -10^{o} in the reference case (Fig. 12). On the other hand, the FLS does not simulate this mechanism, hence its torque curve remains smooth and reaches maximum at about -13^{o}, as in the reference case.

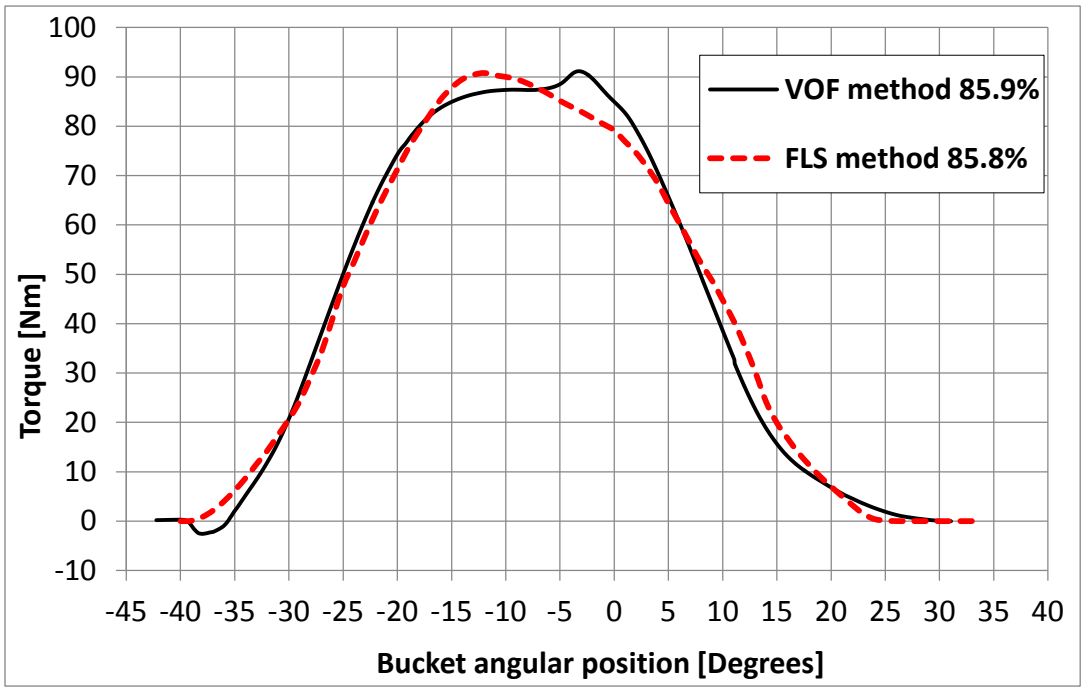


Fig. 17. Comparison of VOF and FLS method for different runner design of 18 buckets (from 22)

5.5. Models performance and comparison for different bucket design

The Fluent-VOF and FLS models were finally applied to simulate the flow in a modified runner design obtained in [5]. The new runner is drastically different in terms of the bucket shape, including changes of the main dimensions, the scheme of the cut, and the exit angle, as shown in Fig. 18. In addition, the number of buckets was 20 and their position was changed in terms of the radial distance from the center and their inclination.

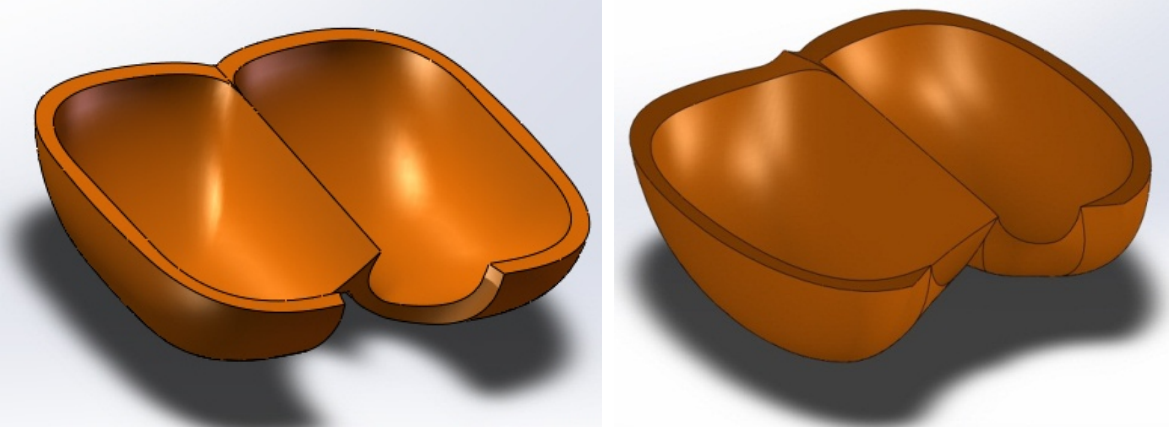


Fig. 18. Comparison of the reference bucket (left) and the modified bucket design (right).

The higher attainable efficiency of this new runner is confirmed by both Fluent and FLS computations, while its value obtained by the FLS exhibits again close agreement with the Fluent value, being only 0.5% higher (95% compared to 94.5%). However, the differences on the torque curves compared in Fig. 19 are more pronounced than in the reference runner, especially in the beginning of the jet impact, (at about -35 degrees), in the area of the maximum torque (-15 to -5 degrees), and to the end of the interaction (after +10 degrees).

These discrepancies where investigated by examining the characteristics of the flow in the new runner as computed by the VOF method, and it was found that the attachment of the flow at the backside of the new bucket is much more pronounced. A much larger portion of the jet flow remains attached there causing substantial torque due the Coanda effect, as shown in Fig. 19 (backside torque curve at -35 degrees). Afterwards, this diverged portion of water hits the leading bucket at an average angle of about -15 degrees, in agreement to available experimental results [20], but at reduced impact velocity, causing reduced peak torque values and delayed evacuation of the bucket. Although a different tuning of current FLS coefficients could mitigate the above deviations, the present results indicate that further development of the FLS model may be required in order to take into account the effects of this important flow mechanism. The structure and philosophy of the FLS method facilitates the introduction of such new relations.

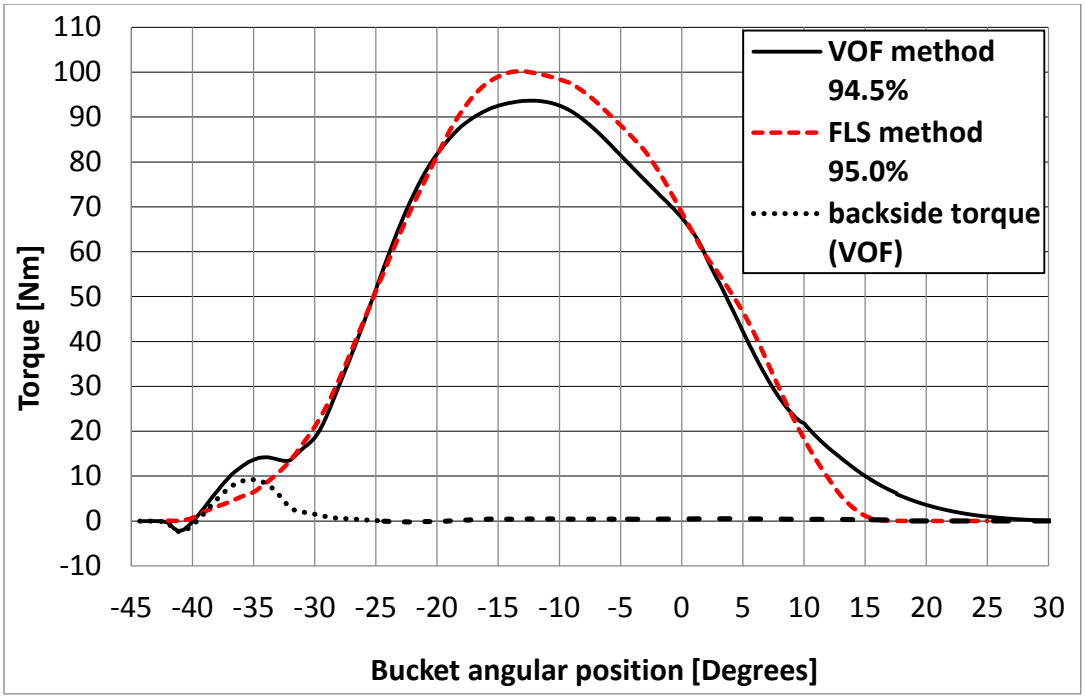


Fig. 19. Comparison of VOF and FLS methods for a new modified runner design

5.6. Summary of runner efficiency results

The results for the hydraulic efficiency of the runner obtained from the various test cases carried out in this study are concentrated in the following Table 1. As it can be seen, the agreement is in all cases very satisfactory, with maximum deviation of only 0.5 percentage units. Considering the much faster flow evaluation by the FLS algorithm, these results support its important advantage for being used for multiple evaluations, as required in parametric performance studies and numerical design optimization of Pelton turbines.

Table 1. Comparison of FLS and Fluent computational results

Case description	Discharge (% of BEP)	Net head (m)	Runner speed (rpm)	No of buckets	**Eff. VOF (%)**	**Eff. FLS (%)**
Reference point	100	100	1000	22	**88.3**	**88.5**
Low flow rate	50	100	1000	22	**91.2**	**91.3**
High flow rate	135	100	1000	22	**86.7**	**86.4**
High hydraulic head	100	156	1250	22	**88.2**	**88.5**
Modified runner	100	100	1000	18	**85.8**	**85.9**
Different bucket design	100	100	1000	20	**94.5**	**95.0**

6. Conclusions

This work aims to validate the capability of a particulate numerical method, the Fast Lagrangian Simulation (FLS) algorithm, to reproduce in a reliable and accurate way the very complex flow created during the jet-runner interaction in Pelton turbines.

The method was compared with a more accurate Eulerian mesh-type VOF method, which has been validated in various previous works and proved to provide satisfactory results. The FLS model is tuned based on the numerical results at a reference operation point of a laboratory model turbine, and then it is applied, together with the Eulerian method, to a number of different test cases, by modifying the flow rate, the hydraulic head and the number of buckets, and the bucket design of the runner.

In all these cases the hydraulic efficiency values obtained by the FLS method showed very good agreement with the Eulerian method results, while the predicted evolution of the free surface flow in the bucket, in terms of time variation of the developing torque, was also satisfactory.

The present results are very encouraging towards the implementation of the FLS tool to perform multi-parametric and multi-objective design optimization studies in modern Pelton runners at very low computational cost. Further validation tests are needed using different runner and bucket shapes for a better and more general adjustment of the FLS coefficients, while further development in order to take into account the effect of other important mechanisms like the backside flow could enhance the accuracy of its results.

Finally, the presented results, along with the provided detailed geometric dimensions of the Pelton model runner, can constitute a benchmark set of data for the validation and evaluation of other numerical modelling tools and methods.

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