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Co-catalytic absorption layers for controlled laser-induced chemical vapor deposition of carbon nanotubes

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ABSTRACT

The concept of co-catalytic layer structures for controlled laser-induced chemical vapor deposition of carbon nanotubes is established, in which a thin Ta support layer chemically aids the initial Fe catalyst reduction. This enables a significant reduction in laser power, preventing detrimental positive optical feedback and allowing improved growth control. Systematic study of experimental parameters combined with simple thermostatic modelling establishes general guidelines for the effective design of such catalyst/absorption layer combinations. Local growth of vertically-aligned carbon nanotube forests directly on flexible polyimide substrates is demonstrated, opening up new routes for nanodevice design and fabrication.

KEYWORDS: laser-induced, chemical vapor deposition, carbon nanotubes, optical feedback, co-catalytic, absorption layer, LICVD

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3 INTRODUCTION
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5 The application potential of carbon nanotubes (CNTs) critically depends on the
6 development of adequate growth and integration methods. Chemical vapor deposition (CVD),
7 in which CNTs nucleate and grow from catalyst nanoparticles exposed to hydrocarbon
8 precursors at elevated temperatures, has become the dominant growth technique due to its
9 versatility and scalability. To enable CMOS back-end compatibility and direct CNT
10 integration on temperature-sensitive flexible substrates, a range of local heating strategies
11 have been explored, in particular laser-induced CVD (LiCVD)^{1–25}. An opaque absorption
12 layer is widely used for effective local heating in combination with high laser powers^{2–4,8,10–}
13 ¹⁴. More controlled CNT LiCVD requires a detailed understanding of the catalytic growth
14 process and the related laser interactions. Crucial to this, albeit widely neglected in current
15 literature, are thereby the process-specific increases in optical absorption arising from initial
16 reduction of the transition-metal nanoparticle catalysts typically used²⁶, such as Fe, and the
17 ever increasing amount of carbon deposited. This can lead to detrimental positive optical
18 feedback, whereby the sample quickly overheats at the centre of the laser spot, leading to
19 uncontrolled inhomogeneous growth^{14,16}.

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38 Here, we introduce the concept of a co-catalytic absorption layer structure for controlled
39 LiCVD, in which a thin CNT catalyst support layer chemically aids the initial catalyst
40 reduction and hence enables a significant reduction of the applied laser power, preventing
41 detrimental positive optical feedback and allowing improved growth control. We focus on Ta
42 mediated solid-state reduction of Fe catalysts²⁷ and systematically investigate the interplay of
43 light absorption and heat conduction for varying Ta support layer thickness to establish
44 general guidelines for the effective design of catalyst/absorption layer combinations in
45 LiCVD. We highlight the potential of this approach by demonstrating the local growth of
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3 vertically aligned CNT forests directly on flexible polyimide substrates with potential
4 applications ranging from flexible, integrated devices to proteomics^{28–35}.
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8 9 10 EXPERIMENTAL METHODS

11 **Sample Preparation.** Transparent fused quartz plates (FQP, 1mm thick), thermally
12 oxidized Si wafers (200 nm SiO₂) and flexible polyimide substrates³⁶ (Kapton HN500,
13 Dupont, 125μm thick) are used as substrates. The polyimide is coated with a ~20 nm Al₂O₃
14 diffusion barrier layer^{37,38} by atomic layer deposition (ALD) using a Cambridge Nanotech
15 Savannah system and a 200°C process with tri[methyl]aluminium and water both carried in a
16 N₂(20 sccm) flow for 200 cycles^{39,40}. Ta layers are sputter deposited (100W, 35sccm Ar,
17 3.5×10⁻³ mbar) in a custom-built DC sputter coater, using a slit-shaped shadow mask⁴¹ to
18 create a gradual thickness profile across the sample (~0 to 130nm, measured by
19 profilometry). The samples are air exposed and subsequently a ~1 nm thick Fe catalyst layer
20 is sputter deposited (20W).
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34 **LiCVD Growth.** Samples are transferred in ambient air. CNT growth is carried out in a
35 custom-built LiCVD chamber (Figure 1; base pressure ~10⁻⁵ mbar). As carbon feedstock, a
36 constant flow of undiluted acetylene is used at 0.1-1mbar pressure. A laser (cw, 532nm, 1-
37 30mW) is focused on the catalyst side of the sample (front-illumination) to a 1μm spot size
38 (FWHM, measured using a knife-edge) with a 50x long-working-distance microscope
39 objective through a viewport (for which it is optically compensated).
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47 **Thermal CVD.** CNT growth is also carried out in a custom-built cold-wall CVD chamber
48 with a resistive graphite heater element. Samples are heated to ~670°C and annealed for 5
49 min in a non-reducing (~10⁻³ mbar vacuum) or reducing atmosphere (100sccm NH₃ at
50 1mbar). CNT growth is initiated by introducing a 10sccm C₂H₂ flow (0.1mbar partial
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3 pressure). After a 45min growth time, the gas flow is stopped and the chamber is pumped
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5 allowing cool down in vacuum.
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7 **Sample Characterization.** During LiCVD growth, real-time in-situ Raman spectroscopy
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9 (Stokes) is carried out with reflected laser light (see Fig. 1) using a Princeton Instruments
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11 Pixis/Acton spectrometer. For post-growth characterization of samples a combination of
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13 scanning electron microscopy (SEM, Zeiss Sigma), transmission electron microscopy (TEM,
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15 Philips Tecnai 20) and Raman spectroscopy (same spectrometer at reduced laser power of 1-
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17 5mW) is used.
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20 21 22 RESULTS

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24 **Solid State Reduction of Fe Catalyst.** We use thermal CVD, i.e. conventional global
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26 substrate back heating, to investigate the influence of the Ta support layer thickness on the
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28 solid-state reduction²⁷ of the initially oxidized Fe catalyst (from transport/storage in ambient
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30 air), and hence the effectiveness of the layer stack for CNT CVD. Figure 2 compares thermal
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32 CNT CVD results for Fe(1 nm)/Ta(~0-130nm) layer stacks for a non-reducing ($\sim 10^{-3}$ mbar
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34 vacuum; Figure 2a) and, as reference, for a reducing pre-treatment atmosphere (100sccm NH₃
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36 at 1mbar; Figure 2b). In both cases CNT growth was then initiated by addition of 10sccm
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38 C₂H₂ (0.1mbar partial pressure) to the pre-treatment atmosphere. Figure 2a shows that for the
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40 non-reducing pre-treatment atmosphere the CNT yield significantly increases at medium and
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42 high Ta thicknesses. SEM imaging (
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47 Figure 2c,d) shows a vertically-aligned CNT forest in a narrow sample region for $\sim 12 \pm 1$ nm
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49 Ta thickness. This CNT forest region broadens for shorter annealing times (5min to 30s),
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51 whereas no or less dense CNT growth was seen for prolonged annealing (not shown). As-
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53 grown CNTs are multi-walled, similar to previous reports^{16,27,42}. Reference samples without a
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55 Ta support layer did not give CNT growth under non-reducing conditions. This clearly
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3 confirms that Ta is required for catalyst activation and acts as the primary reducing agent for
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5 the Fe²⁷, with no significant catalyst activation from the C₂H₂ gas itself under these
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7 conditions.
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10 In contrast, under reducing NH₃ pre-treatment and growth conditions Figure 2b shows that
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12 the narrow region of CNT forest growth is shifted towards lower Ta thicknesses and also Ta-
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14 free samples show growth. This is consistent with NH₃ acting as reducing gas, activating the
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16 Fe without the need for Ta. For such NH₃ pre-treatment conditions, no CNT growth occurs at
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18 high Ta thicknesses, which can be understood by assuming that once reduced the catalyst can
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20 diffuse into the metal support and hence can be lost from nucleating CNT growth depending
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22 on how quickly the hydrocarbon is introduced^{42,43}.
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26 **LiCVD Optimization.** With the above thermal CVD experiments in non-reducing pre-
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28 treatment atmosphere showing how the Ta support layer chemically aids the initial Fe
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30 catalyst reduction, we now focus on the laser induced CVD process where the Ta based layer
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32 stack additionally acts as an absorption layer, and thus fulfills an enhanced functionality. We
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34 employ a simple LiCVD growth procedure in 1mbar C₂H₂ with a growth time of 60s (see
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36 Experimental and Fig. 1). Figure 3 shows that a significant variation in CNT growth is
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38 observed with laser power and Ta support layer thickness. The laser power threshold, below
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40 which no CNT growth is observed, is 8mW (see Supporting Information). The SEM analysis
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42 shows that irradiation with 20mW laser power at low (3.5nm) and medium (15nm) Ta
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44 thickness respectively (Figs. 3a,b) leads to a 'donut'-shaped CNT growth profile, which is
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46 caused by overheating at the centre of the laser spot. For thicker Ta (70 nm, Fig. 3c) the
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48 overheating in the central CNT growth spot can be avoided, but the CNT growth yield
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50 drastically decreases and for even thicker Ta no CNT growth is observed. In comparison, at
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52 lower laser powers (8.5 mW, Fig. 3d) for 15 nm Ta thickness a circular growth spot with
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54 comparatively high CNT yield is observed. It is noteworthy that the growth spot diameter is
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3 400-500nm, which is significantly smaller than the actual laser spot size ($\sim 1\mu\text{m}$). To the best
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5 of our knowledge, this is the smallest continuous LiCVD-grown CNT spot reported to date.
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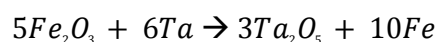
7
8 Figure 4 shows the time evolution of a growth spot within the first 60s of laser illumination
9 based on 15nm of Ta thickness and low laser power conditions, as optimized above. The
10 SEM analysis shows visible CNT growth starting at $\sim 15\text{s}$ (Fig. 4a), after which the CNT spot
11 diameter steadily increases. This potentially enables tailoring of the desired CNT growth spot
12 size. We find the results of time-resolved in-situ Raman spectroscopy during LiCVD ($\sim 1\text{s}$
13 resolution, see Supporting Information) to agree well with post-growth Raman spectra (Fig.
14 5b). Starting at $t=20\text{s}$, 5s after CNT nucleation, the in-situ measured time-resolved integrated
15 G-band intensity of the central growth spot increases linearly with time and hence correlates
16 well with the SEM analysis of Figure 4a. Furthermore, the D/G-ratio remains at a constant
17 level throughout CNT growth, indicating no rise in defect density. This confirms that the
18 laser induced growth process on such optimized Fe/Ta layer stacks is not compromised by
19 detrimental positive optical feedback.
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34 **LiCVD on Flexible Substrate.** We use the optimized Fe/Ta layer stack and LiCVD
35 conditions to demonstrate direct CNT growth on a flexible thin polyimide foil. These
36 polymer substrates show fundamentally different properties than the fused quartz plates and
37 oxidized Si wafer substrates used above. As polyimide is non-transparent at the given laser
38 wavelength, light absorption of the substrate contributes to the overall heat generation.
39 Further, polyimide is stable only up to 400°C and its thermal diffusivity is ten times lower
40 than that of quartz. We consistently find that the laser power threshold for CNT LiCVD
41 reduces to $\sim 3.5\text{mW}$ for polyimide substrates. Figure 5a shows that for optimized conditions
42 (2nm Fe on 8nm Ta, 4mW, 60s) pillars of vertically aligned CNT forests can be grown with a
43 spot size of $\sim 3\mu\text{m}$ and without any visible damage to the polyimide substrate. High-
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3 resolution TEM analysis (Figure 5b) further confirms that the as-grown CNTs are multi-
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5 walled, consistent with the Raman data.
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8 9 DISCUSSION

10 **Co-catalytic layer structure.** Figure 6 schematically depicts the LiCVD growth process and
11 highlights the multiple functions of the Ta/Fe layer stack. One function of the additional Ta
12 layer is that it absorbs laser light and transforms it into heat (I), which it subsequently
13 conducts radially away from the laser spot (II). The thermostatics of this will be considered
14 below. The Ta also aids the initial reduction of the Fe catalyst enabling the formation of
15 catalytically active metal nanoparticles (III)^{27,42}. We note that Ta itself is not catalytically
16 active for CNT growth⁴⁴ under the conditions used. The Fe/Ta solid state reduction
17 mechanism is driven by the higher oxygen affinity of Ta⁴⁵. At temperatures as low as 300°C,
18 Fe oxide can be fully reduced to elemental Fe, based on diffusion of oxygen ions from Fe-
19 oxide (Fe³⁺) to form stable Tantalum pentoxide (Ta⁵⁺). This co-catalytic reaction can be
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41 We have however previously found that Ta pre-oxidation from ambient air exposure or
42 residual oxygen can change the balance of the equation and that the desirable ratio of Fe:Ta
43 thicknesses for solid state reduction and co-catalyzed CNT growth is hence strongly affected
44 by the annealing conditions²⁷. Residual oxygen in the CVD growth chamber can cause re-
45 oxidation of the Fe catalyst (IV, Fig. 6) and a direct gas phase oxidation of the Ta. Hence
46 under oxygen-residual-rich low vacuum conditions (as used in this study, base pressure $\sim 10^{-5}$
47 mbar) a comparably thicker Ta film is needed for CNT growth. A thicker Ta support can
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3 however lead to increased depletion of the active Fe catalyst by its diffusion into the Ta bulk
4 through grain boundaries (V, Fig. 6)⁴³. Since only catalytically active Fe nanoparticles at the
5 Ta surface exposed to the hydrocarbon precursor can nucleate CNTs (VI, Fig. 6), this
6 indicates that an optimized Fe/Ta-thickness ratio is based on a compromise between these
7 two competing effects. The narrow region of CNT forest growth seen in Figure 2 can be seen
8 to reflect the balance between the ongoing catalyst diffusion and (re-)oxidation, from which
9 only nanoparticles with nucleated CNTs are exempt. Pre-annealing in a reducing NH₃
10 atmosphere can increase the time for which Fe is reduced before CNT growth and hence
11 increases the likelihood of catalyst deactivation by diffusion into the Ta bulk, which explains
12 the lack of CNT growth at high Ta for these conditions. For LiCVD, a decrease in CNT
13 nucleation density was observed only at comparatively high Ta thicknesses, further
14 highlighting the strong link between catalyst depletion and annealing time.

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30 **Thermostatics of light absorbing layers.** To gain insight into the thermostatics of laser-
31 heated absorption layers we analyze the balance of heat gain and loss mechanisms at the
32 center of the laser spot. Whilst a detailed quantitative analysis of the temperature profile
33 resulting from laser heating is beyond the scope of this study, we can apply simple models to
34 rationalize the experimentally observed absorption layer thickness dependencies. The
35 following heat flux equation can be applied:
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$$\frac{dQ}{dt} = \frac{dT}{dt} C = \dot{q}_{in}(z, I_0) - \dot{q}_{out}(T, z)$$

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48 where Q is the heat energy, T the temperature and C the specific heat capacity of the heated
49 material of thickness z , while \dot{q} is the power density, meaning that for a given area σ ,
50 $\dot{q} = \frac{1}{\sigma} \frac{dQ}{dt}$. The incoming power density \dot{q}_{in} based on light-heat conversion equals the intensity
51 of absorbed laser light I_{abs} via the Beer-Lambert law:
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$$\dot{q}_{in}(z, I_0) = I_{abs} = (1 - R)(1 - e^{-\alpha z})I_0 = (1 - R)A(z)I_0$$

Here R is the reflectivity of the sample surface, α is the attenuation coefficient, I_0 is the incoming laser intensity and $A(z)$ the optical absorptivity. The peak intensity, I_0 , at the center of the laser spot is simply given by twice the incoming laser power, P_0 , divided by the modal area (w_0 : modal radius):

$$\dot{q}_{in}(z, P_0) = \frac{2}{\pi w_0^2} (1 - R)A(z)P_0 := \beta A(z)P_0$$

In order to further simplify the discussion, we combine all the factors which are taken as constant into β . Given the thin (<130nm) absorption layers used in this study in comparison to the characteristic heat penetration depths, the temperature is considered constant throughout the thickness of the absorption layer. In this case, the only term dependent on the layer thickness is the heat loss via lateral heat conduction radially away from the spot centre, $\dot{q}_{lat}(T, dz)$. Heat losses by radiation and those by conduction to the gas phase and into the substrate $\dot{q}_{rgs}(T)$ only depend on the temperature. This leads to the following simplified expression for heat losses:

$$\dot{q}_{out}(T, z) = \dot{q}_{lat}(T, z) + \dot{q}_{rgs}(T)$$

For the case of radial heat flow in an infinite slab with a continuous, axially-symmetric heat source, T increases linearly with the input power density per unit thickness⁴⁶, $\frac{\dot{q}_{in}}{z}$. Therefore to maintain a constant T for increasing absorption layer thickness, \dot{q}_{in} must increase linearly with z . Hence for a given T , $\dot{q}_{out}(z)$ has the following linear form:

$$\dot{q}_{out}(z) = az + b$$

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5 Figure 7a depicts the laser power threshold required for CNT growth as a function of the Ta
6 layer thickness. The data points represent the lowest (arrow down) and the highest (arrow up)
7 Ta thickness for which CNT growth was clearly observable with in-situ Raman spectroscopy
8 (integrated G-Band, see Supporting Information) using laser powers of 8-20mW. Growth
9 spots were set along an entire Ta gradient sample with a separation of 50 μ m, corresponding
10 to a maximum relative change of Ta thickness between individual growth spots of ~10%.
11 These data points therefore define a phase boundary between CVD conditions that yield CNT
12 growth and conditions which do not lead to growth.
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22 The threshold laser power shows a minimum for 15nm Ta, which is confirmed by marginal
23 CNT growth using 8mW laser power (as seen in SEM, Supporting Information). Therefore
24 the peak temperature, induced by irradiation with a given laser power, is maximized at this
25 Ta thickness. We note that this ought not to be confused with most efficient light-to-heat
26 conversion, but simply means that this Ta thickness marks the best compromise between the
27 competing effects of simultaneously increasing heat gain and heat loss with rising thickness
28 of the absorption layer. Towards lower and higher Ta thicknesses the required laser power
29 increases due to the lower optical absorption of thinner Ta layers or the higher conductive
30 heat losses for thicker Ta layers, respectively.
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43 Under the assumption that the temperature needed to trigger CNT CVD, T_{gr} , is not
44 significantly affected by changing the Ta thickness, we identify the phase boundary in Figure
45 7a as a temperature contour line. Due to the rapid heating in LiCVD, we can further assume
46 that the incoming and outgoing heat fluxes are in equilibrium, which gives the defining
47 equation for the temperature contour line (where ' indicates $T = T_{gr}$):
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$$P_0(z, T_{gr}) = \frac{1}{\beta A(z)} [\dot{q}_{lat}(z) + \dot{q}_{rgr}]_{T=T_{gr}}$$

To obtain a measure for the heat losses, $A(z)P_0(z, T_{gr})$ is plotted in Figure 7a (blue circles). As can be seen, heat losses increase linearly with the Ta thickness above $\sim 8-10\text{nm}$, and approach a constant value below this thickness. This relates to the morphology of the thin Ta layers on these substrates: sputter deposition initially leads to island formation, before forming continuous layers for increasing deposition time⁴⁷. Hence lateral heat conduction is negligible in the initial island formation stage and heat losses equate to \dot{q}_{rgr} . To reconstruct the phase boundary, we fit the acquired data using tabulated material properties of Ta⁴⁸ ($\alpha = 4.42 \times 10^5 \text{ cm}^{-1}$, $R = 0.37$, from Fresnel equations). Thereby $\beta^{-1}\dot{q}_{lat}(z)$ is determined from a linear fit to $A(z)P'_0(z)$ (blue dotted line in Fig. 7a) at high Ta thickness, where the film is continuous. The constant $\beta^{-1}\dot{q}_{rgr}(z) \sim 2.8\text{mW}$ is obtained by fitting $P_0(z, T_{gr})$ at low Ta thickness where lateral heat conduction is neglected, $\dot{q}_{lat}(z) = 0$. In contrast to the experimental data, the resulting fit curve (red dashed line in Fig. 7a) converges to a minimum at high Ta thickness, which further illustrates the important effect of lateral heat conduction in the absorption layer at the threshold laser power.

The experimental LiCVD growth results can be understood on the basis of this model. The conditions of growth spots (A-D) from Figure 3 are discussed in the context of the contour plots of Figure 7a and their temperature profiles are schematically indicated in Figure 7b. Points (A,C,D) are close to the phase boundary and therefore all have comparable peak temperatures. The linear relation between the Ta thickness and heat conduction suggests that no broadening of the temperature profile via lateral heat conduction occurs, which is consistent with the observed growth spots having similar diameter. In Figure 3, spots (A) and (B) both show donut-shaped CNT growth profiles, resulting from sample overheating.

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3 Graphitization in the spot centre of (A) suggests that for this spot the light-induced
4 temperature profile was initially well adjusted for CNT growth, but after nucleation its
5 temperature increased (profile A') by optical feedback¹⁶ due to increased uncontrolled light
6 absorption within the carbon material. In contrast, the absence of CNTs in the centre of spot
7 (B) is not due to a positive optical feedback effect, but indicates extensive overheating ab
8 initio, caused by the far-higher-than-needed laser power (20mW with 8mW threshold). This
9 can be rationalized by considering the combined effect of faster Fe catalyst diffusion into the
10 Ta film and decreasing growth rates at the increasingly excessive temperatures towards the
11 spot centre, which prevents CNT growth as highlighted for thermal CVD in previous
12 literature^{42,49,50}. CNT growth temperatures were attained in outer regions only where the
13 incident laser intensity was lower, which is in good agreement with the observed increased
14 spot size (Figure 3b). Similarly, the thick Ta film at spot (C) formed a large diffusion sink for
15 active Fe nanoparticles, which led to a drastically reduced nucleation density (Figure 3c)⁴².
16 Furthermore, spots (B-D) do not show detrimental graphitization due to optical feedback,
17 suggesting that the higher optical absorption of thicker Ta films indeed gives a more stable
18 temperature evolution thus indicating an effective strategy for optical feedback control.
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41 CONCLUSIONS

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43 Controllable, co-catalytic LiCVD growth of carbon nanotubes was demonstrated, whereby
44 a light absorbing tantalum layer also aids in chemical activation of the initially oxidized iron
45 catalyst. Using transparent substrates and laser irradiation directly onto the catalyst, the
46 influence of absorption layer thickness on CNT growth morphologies and laser threshold
47 powers were identified. A simple thermostatic model was used to rationalize the obtained
48 optimum experimental conditions and to establish general guidelines for the effective design
49 of catalyst/absorption layer combinations in LiCVD, enabling controlled CNT growth
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3 without detrimental effects from sample overheating by optical feedback. The potential of
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5 this approach was highlighted by achieving localized growth of vertically aligned CNT
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7 forests directly on flexible polyimide substrates.
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ASSOCIATED CONTENT

Supporting Information. SEM image of LiCVD grown spot using 8mW and 15nm Ta; Waterfall plot of in-situ Raman spectra during first 60s of growth (10mW, 15nm Ta); Plotted integrated G-peak intensity of growth along the Ta gradient under various laser powers (Raw data of Figure 7a). This material is available free of charge via the Internet at <http://pubs.acs.org>.

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Notes

The authors declare no competing financial interest.

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FIGURES

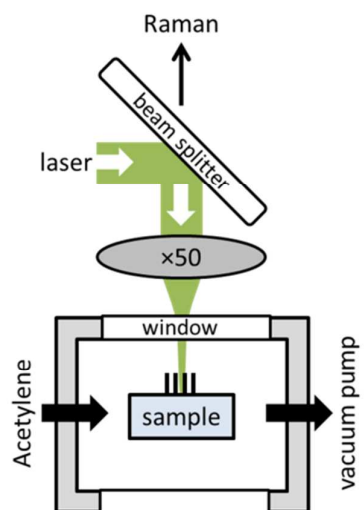


Figure 1: Schematic of LiCVD set-up with in-situ Raman spectroscopy.

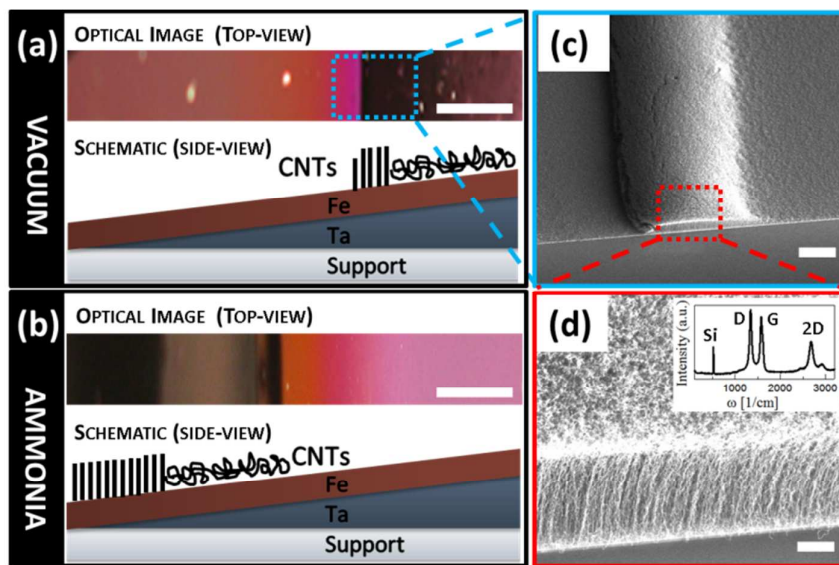


Figure 2: Thermal CNT CVD results for Fe(1nm)/Ta/SiO₂(200nm)/Si wafer samples with exponential Ta thickness profile (~0-130 nm). (a,b) Top-view optical micrographs (upper panels) showing CNT growth areas when annealing in (a) vacuum and (b) in ammonia. Black region maps the presence of CNTs with density (as observed in SEM) illustrated in the corresponding side-view schematics below (straight lines denote high-density vertically aligned CNT ‘forest’ growth, squiggly lines represent unaligned low density growth; layer sketches not to scale). (c) SEM image taken from boxed region in (a) showing severely-confined zone of CNT forest growth at ~12nm Ta thickness. (d) High magnification SEM image of the CNT forest shown in (c). Inset shows the corresponding Raman spectrum for this region of the sample, with peaks typical of multi-wall CNTs. Scale bars: (a,b) 1mm; (c) 20 μ m; (d) 2 μ m.

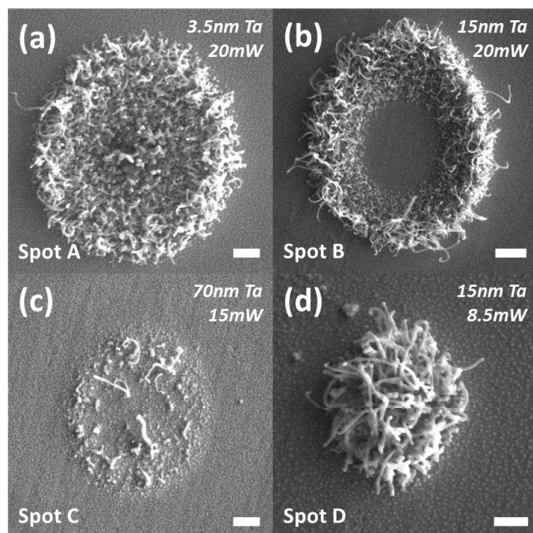


Figure 3: SEM images representative of different LiCVD CNT growth regimes for different laser power and Ta thickness: (a) 20mW, 3.5nm Ta (b) 20mW, 15nm Ta, (c) 15mW, 70nm Ta, (d) 8.5mW, 15nm Ta. Scale bars: (a) 300nm, (b) 1 μ m, (c) 300nm, (d) 200nm.

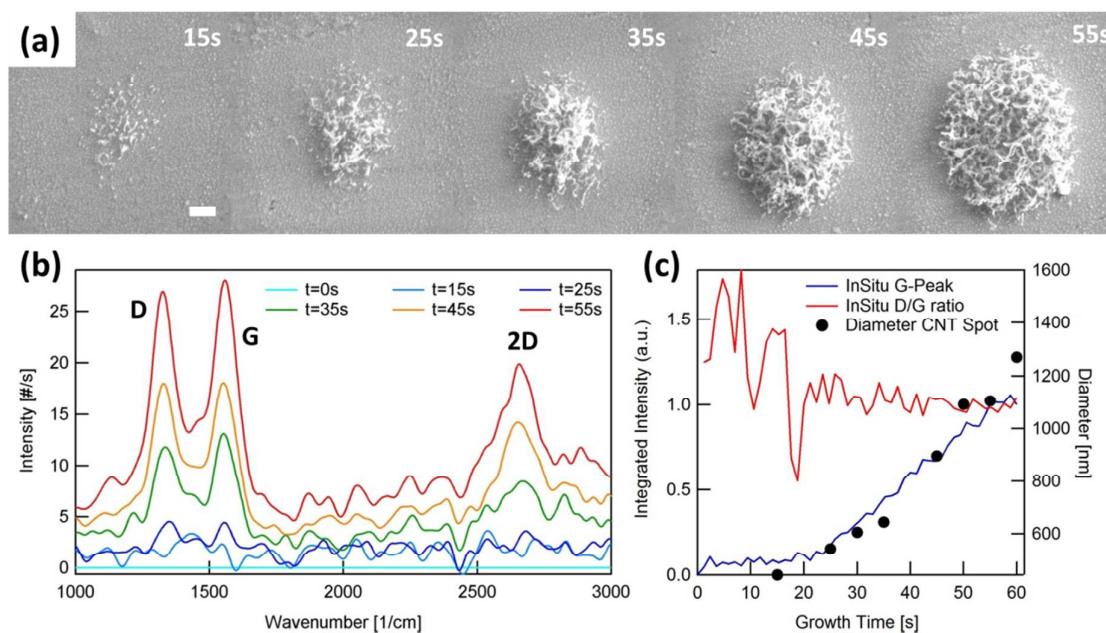


Figure 4: Time series of CNT growth under optimized growth conditions (10mW, 15nm Ta). (a) SEM images of CNT growth spots for various exposure times (Scale bar: 200nm). (b) Post-growth Raman spectra of the CNT growth spots in (a). (c) G-peak intensity and D/G ratio during the first 60s of CNT growth (~ 1 s resolution) showing constant defect rates and linear growth rate after nucleation of CNTs ($t > 20$ s) and correlation between G-Peak intensity and the diameter of a growth spot.

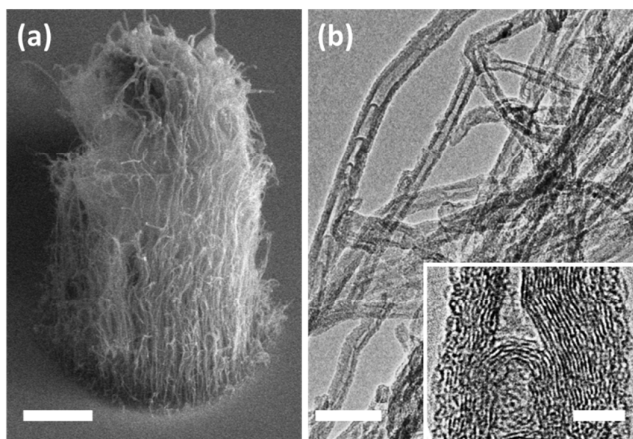


Figure 5: LiCVD-grown CNTs on flexible polyimide substrate. (a) SEM image of vertically aligned CNT forest grown for 60s at 4mW laser power (45° tilt); (b) HR-TEM images representative of as-grown CNTs. Scale bars: (a) 1 μ m, (b) 50nm, inset: 5nm.

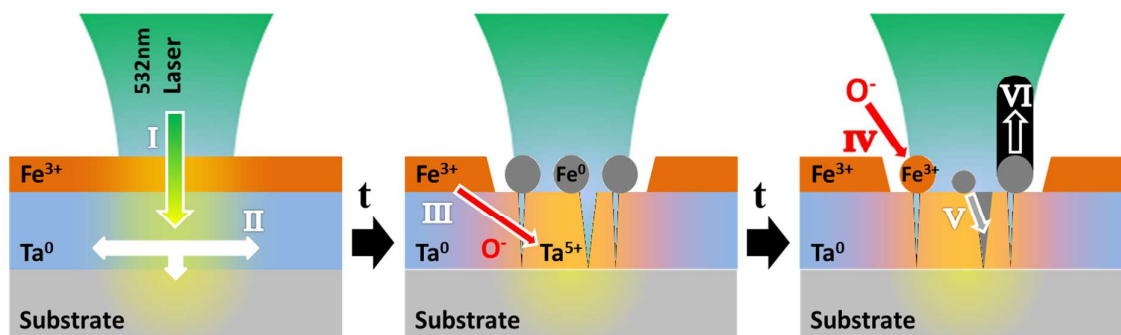


Figure 6: Schematic of the processes involved in LiCVD growth of CNTs on FeTa co-catalyst. (I) Light absorption and conversion to heat. (II) Heat conduction within the Ta film away from the centre of the Gaussian laser beam. The resulting temperature profile is illustrated in yellow (intense colour denotes high temperature). (III) Solid state reduction of Fe^{3+} to catalytically active Fe^0 , which de-wets into catalyst nanoparticles. In this process Ta⁰ is oxidized to Ta⁵⁺. (IV) Re-oxidation of Fe^0 by residual oxygen in the growth chamber. (V) Fe^0 diffusion into bulk Ta through grain boundaries. (VI) Nucleation and growth of CNTs.

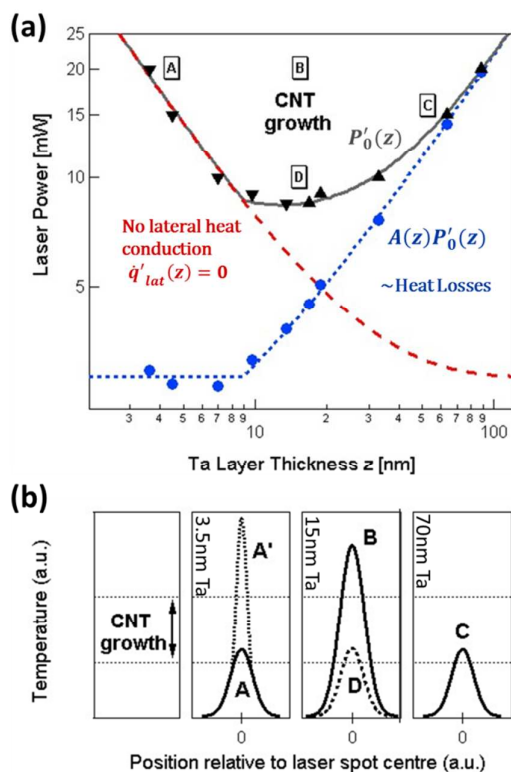
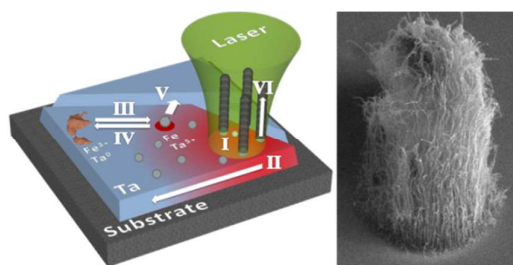


Figure 7: (a) LiCVD phase diagram: Measured lower (down-pointing triangles) and upper threshold (up-pointing triangles) of Ta layer thickness which leads to CNT growth (Raman G-band) under sample illumination using various laser powers. These data points mark a contour line for the lowest temperature which triggers the CVD reaction. The predicted laser power turning off heat conduction through the Ta film (red dashed) and for constant absorptivity (blue dotted) are also shown. The latter reveals the onset of non-negligible lateral heat conduction at ~ 8 -9nm thickness where the Ta film becomes continuous. (b) Schematic spatial temperature profiles during CNT growth of the spots shown in Figure 3, under conditions marked in (a). Curve A' marks the heat profile for A after being altered during growth by optical feedback.

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