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# **Reaction zone visualisation in swirling spray n-heptane flames**

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#### **Abstract**

Joint PLIF measurements of CH2O and OH were performed in a swirl-stabilized nheptane spray flame at conditions close to extinction. Simulations of laminar counterflow heptane flames at different strain rates showed that the heat release could be approximately represented by the product  $CH<sub>2</sub>O<sub>x</sub>OH$ , but also that this product cannot visualise the heat release at the lean side of stoichiometry. The simulations suggest that the outline of  $CH<sub>2</sub>O$ regions in PLIF images could be an approximate indicator of the stoichiometric mixture fraction iso-line. Due to the intense turbulence and local extinction, individual PLIF images show a very variable behaviour. They indicate rich zones, reaction sheet breaks, lift-off, and they suggest that this flame is mostly of non-premixed character. The mean heat release rate as represented here is consistent with inverse Abel-transformed OH\* chemiluminescence imaging. The usefulness of this technique for spray flames is discussed.

**Keywords:** Extinction, swirl, formaldehyde fluorescence, heptane spray

#### **1. Introduction**

Local extinction is an important limiting factor for lean combustion. Knowledge of the structure of the flame close to blow-off is important from a practical perspective, but also from the viewpoint of validating advanced turbulent combustion models. In gas turbines and industrial furnaces, the flame is virtually always stabilised by swirl. It can be argued that we know little about the fundamental processes of extinction of spray flames in recirculation zones, as we do not have the same degree of information as we do for gaseous flames [1-4]. Cavaliere *et al.* studied n-heptane spray flames and saw local extinction holes present along the flame sheet and randomly-occurring lift-off [5]. These observations were based on 5kHz OH-PLIF, which cannot reveal with certainty whether a flame is extinguished or not. In particular, the OH signal alone may not reveal whether the flame is of premixed character (and therefore there may be OH in the post-flame gases) or of non-premixed character (and therefore the OH is found roughly along the stoichiometric contour only).

As a second scalar for flame structure studies, formaldehyde  $(CH_2O)$  has been used extensively. CH2O is produced in the low temperature oxidation process and consumed in the subsequent high temperature oxidation. Formaldehyde LIF was used for autoignition of methane jets [6], methanol, ethanol and acetone spray jet flames [7,8], and diesel fuel [9-11] and n-heptane [11,12] in HCCI engines. Najm et al. [13] conducted detailed chemical kinetic computations of methane premixed flame and found that the concentration of formyl radical (HCO) is very well correlated with flame heat release rate. They also point that the production of HCO is directly dependent on the availability of its precursor  $CH_2O$  ( $CH_2O$ )  $\overrightarrow{OH,H,O,M}$  HCO). However, as the PLIF of HCO is difficult for single shot imaging diagnostics when applied to fully turbulent flows due to its short fluorescence lifetime and low concentrations [13], the product of simultaneous OH and CH2O PLIF was considered as an alternative. Simultaneous CH<sub>2</sub>O and OH imaging and then taking the product  $X_{CH2O} \times X_{OH}$  (X being the mole fraction, taken as proportional to the PLIF signal intensity) has been used successfully as a marker of heat release in non-premixed and premixed flames [14-17]. Even when quantitative imaging has not been possible, the qualitative comparison of  $CH<sub>2</sub>O$  and OH regions has been very fruitful. Medwell and Dally [18] discussed formaldehyde in MILD combustion conditions, while Masri and co-workers looked at simultaneous CH2O and OH in jet flames of ethanol, methanol and acetone [7,8] and very interesting comments on the nature of the reaction zones in these spray flames could be made.

In this paper, we use simultaneous  $CH<sub>2</sub>O$  and OH imaging in the stabilisation region of an n-heptane spray flame in a swirl-induced recirculation zone, a configuration of great relevance to gas turbine combustion, and in particular for flames close to the stability (i.e. extinction) limits. Our specific objectives are: (i) to examine the performance of the  $CH<sub>2</sub>O-$ OH technique for n-heptane; (ii) to visualise reaction zones in swirl spray flames; (iii) to examine how these zones are altered close to blow-off. The paper continues with a presentation of the experimental method, some results from laminar flame simulations, and typical images from the application of the technique.

#### **2. Methods**

An enclosed bluff-body swirling n-heptane spray flame was stabilized at two flow conditions: far from and close to blow-off. The burner (Fig. 1) was identical to the one examined previously by Cavaliere *et al.* [5], with the exception of a new atomiser giving a slightly narrower spray. Four quartz plates formed the combustion chamber of 97 mm width square and 150 mm length. The air was supplied through a 350 mm long circular duct of *D*=37 mm inner diameter fitted with a conical bluff body of diameter  $d=25$  mm. A 60° swirler provided strong swirl at the exit, which, together with the bluff body, contributed to the establishment of a recirculation zone. For information, see Ref. [19] for velocity measurements in this burner for the n-heptane case, which should not be too far from the present flame due to the dominance of the air flow. A pressurized atomizer (Lechler,  $\text{\#212.054.17.AC}$ ; a flow rate up to 0.45 g/s at 5 bar) was fitted inside the bluff-body holder, providing a hollow cone spray profile with a nominal spray angle 60°.

The air flow was dried and particle-filtered and supplied by an air compressor. The heptane was supplied from a nitrogen-compressed container. Air and fuel flow rates were metered separately using an Alicat mass flow controller (MFC 5, 1000 SLPM) and a Bronkhorst liquid flow controller (LIQUI-flow, L30, 0-2 g/s). The two flow conditions studied as well as the blow-off condition are listed in Table 1.

Two 10Hz Nd:YAG lasers and one dye laser were used for the joint CH2O and OH PLIF system [20]. The OH-PLIF was excited at 283 nm, while the CH<sub>2</sub>O-PLIF was excited at 355 nm. The fluorescence of OH was captured in the range of 309-375 nm with WG305 and UG11 Schott glass filters applied. The fluorescence of CH<sub>2</sub>O was captured in the range 375-610 nm using an intensified CCD camera through a filter set (laser line notch filter: OD>4, 17.8 nm FWHM, GG395 and BG 40) to eliminate scattering, flame luminosity, and soot emission interference. A 355 nm laser line notch filter was added to the CH2O camera to further cut-off any scattering of laser light by droplets. Because heptane (liquid and/or vapor) and PAH may also fluoresce in the range 350-600 nm [21,22], possible interference from fluorescence of vapor and liquid and of PAH could not be subtracted from the  $CH<sub>2</sub>O$  signals. Some comments on the importance of this interference are made later based on laminar flame simulations and on some preliminary experiments. The nominal resolution of the imaging system was 0.05 mm/pixel. The gain of the image intensifier was 30,000 FL/FC and 80,000 FL/FC for CH<sub>2</sub>O and OH respectively. The laser powers were 12 mJ/pulse for OH PLIF and 300 mJ/pulse for CH2O PLIF. The intensifiers were triggered off the Q switch of each laser with a gate of 220 ns. The CH<sub>2</sub>O PLIF laser was triggered at 300 ns delay from the triggering of the OH laser. The laser pulse width was 10-12 ns.

Both OH and CH2O PLIF images were transformed to be spatially aligned with each other on a pixel-by-pixel basis. The alignment needed to be precise and followed several stages. Firstly, a target image (a grid pattern with several reference points) was taken using both cameras. Secondly, the CH2O image was mirrored and the corresponding coordinates of the same reference point were identified. Thirdly, a transformation matrix relating the two sets of coordinates was computed and then used to map the OH images onto the coordinates system of the CH2O images. The images were further processed via Matlab with background subtraction, laser profile correction,  $3\times3$  median filtering, and a threshold below which the signal was put to zero. The product  $CH<sub>2</sub>O<sub>×</sub>OH$  was then calculated and plotted as indicator of local reaction zone. 400 images were recorded for each flow condition.

Finally, a laminar counterflow diffusion flame was simulated using Cosilab [23] with a detailed n-heptane mechanism [24]. The main species as well as the heat release were obtained for a range of strain rates and used as a reference for the joint  $CH<sub>2</sub>O-OH$  PLIF measurement.

#### **3. Results and discussion**

#### 3.1 *Laminar flame simulations*

Figures 2 to 4 show results from the counterflow non-premixed flame calculations for two strain rates in terms of mole fractions  $(X)$  *vs.* mixture fraction  $(\xi)$ . C<sub>2</sub>H<sub>2</sub> is plotted here as an indicator of soot precursors. Formaldehyde is present in the fuel-rich region, with  $X<sub>CH2O</sub>$ increasing from  $\xi=1$  toward the stoichiometric  $\xi_{st}$  ( $\xi_{st}=0.062$  for heptane), decreases sharply close to  $\xi_{st}$ , and is not present at all in the lean side. In contrast,  $X_{OH}$  peaks near  $\xi_{st}$  just before the temperature and the formaldehyde peaks. At the higher strain rate shown, X<sub>CH2O</sub> reaches

its maximum around  $\zeta=0.4$ , and the values are higher compared to the low strain rate. C<sub>7</sub>H<sub>16</sub> and  $C_2H_2$  are present in the fuel rich region, but the overlap between these species and OH is minor. This suggests that any fuel or PAH fluorescence will not coincide with the OHcontaining region.

Figure 3 shows the OH and CH2O normalised by their peak values and includes the similarly normalised heat release rate (HRR). It can be seen that the HRR curve contains the curve showing the product  $X<sub>CH2O</sub> \times X<sub>OH</sub>$  (Fig. 3a). The reaction pathway analysis shows that high percentage of CH<sub>2</sub>O originates from the decomposition of the methoxy redical  $(CH_3O+M=\geq CH_2O+H+M)$  and from  $CH_3$   $(CH_3+O=\geq CH_2O+H)$ . Also, formaldehyde is consumed by H-atom abstraction reaction mainly through the reaction  $CH_2O+H=\geq HCO+H_2$ and  $CH_2O+OH = > HCO + H_2O$ . At low strain rates, the reaction  $CH_3 + O = > CH_2O + H$  and CH<sub>3</sub>O+M=>CH<sub>2</sub>O+H+M is seen approximately in the region  $\zeta=0.066$  to  $\zeta=0.097$ , and the reaction CH<sub>2</sub>O+OH=>HCO+H<sub>2</sub>O is mostly observed in mixture fraction space from  $\zeta=0.070$ to  $\zeta$ =0.099. This region correlates well with the product  $X_{CH2O} \times X_{OH}$  region. However, due to the consumption of  $CH<sub>2</sub>O$  in the lower mixture fraction region, the discrepancy between HRR and the product  $X_{CH2O} \times X_{OH}$  is high and the heat release rate comes mainly from other reactions (such as  $CO+OH = >CO<sub>2</sub>+H$ ,  $HO<sub>2</sub>+OH = >H<sub>2</sub>O+O<sub>2</sub>$ ). Thus, at low strain rate, the HRR is not quantitatively represented by the product  $CH<sub>2</sub>O<sub>×</sub>OH$ : the latter could only contribute around 50% of the true HRR. In contrast, at high strain rate, the two quantities are quite close. A further observation is that the OH peaks at the lean side of the HRR peak, while for both low and high strain rates the peak  $CH<sub>2</sub>O<sub>×</sub>OH$  coincides with the HRR peak. Unlike the strong variation of CH2O with strain rate, the OH curve and levels do not vary too much, although a decrease in the peak OH is evident. Finally, Fig. 4 shows the normalised HRR plotted *vs.* the normalised CH<sub>2</sub>O×OH and *vs.* the normalised  $(CH_2O+C_7H_{16})\times OH$  for various strain rates. It is evident that for high strain rate the correlation is high and that there is likely no effect from any fuel fluorescence. As the strain rate decreases, it is evident that the product underestimates the heat release, however it can still be used as a spatial marker of the location of a large fraction of the heat release.

#### 3.2 *Instantaneous* CH2O×OH *images*

Figure 5 shows an image obtained in un-ignited flow with the fuel running and an example of reacting flow image, both captured by the CH<sub>2</sub>O camera. The inert flow image suggests that the CH<sub>2</sub>O images contain some information from the droplets. This is not likely to be Mie scattering, as the laser wavelength has been filtered out. The droplet-like shape of the high intensity regions suggests fluorescence of heptane from the liquid phase, but the intensity is not very high. Considering the density ratio between vapour and liquid suggests that the fuel vapour is not likely to be picked up by the system. The flame images are usually smooth, but occasionally contain similar bright circular spots: see, for instance, the circled spot in Fig. 5b. These are thought to originate from fuel droplets. The percentage of images showing droplet-like shapes in the dataset is small (around 10%) and such images are removed before the calculation of averages. In addition, such droplet-like regions virtually always do not overlap with the OH image.

Figure 6 presents a few typical OH,  $CH<sub>2</sub>O$ , and  $CH<sub>2</sub>O<sub>×</sub>OH$  (HR) images from the two flames. For both flow conditions, the OH signal shows discontinuities, similar to what has been observed previously with 5kHz OH-PLIF imaging [5]. The OH-containing regions are quite thin and the recirculation zone is void of OH, contrary to premixed flames [5,25] that have significant amounts of OH inside the recirculation zone. In both H1S1 and H1S2, the OH regions are overall thin, indicating that the present spray flames have non-premixed flame characteristics. The OH line follows a relatively straight path along the spray in the inner part of the flame, and a very convoluted part roughly aligned with the shear layer outer flame, which is occasionally anchored at the corner of the bluff body. The location of OH seems more variable in H1S2 than in H1S1, which indicates that at a condition far from blow-off (H1S1), the flame sheet is less often extinguished along the spray path.

The region A (inside the recirculation zone) and region C (in the annular air jet), Row I, Fig 6a, have very low CH2O and OH. Occasionally, some formaldehyde appears in the outer recirculation zone (further outwards from the annular air stream), likely to be due to the capture of unburnt or partially-burnt fuel from the wall region where the spray impinges on the wall. Between the spray and the air flow in the annular jet (i.e. region B), intense  $CH<sub>2</sub>O$  is visible. Inside the hollow cone spray (region A, i.e. close to the axis), where hot gases recirculate, the CH2O signal is virtually zero. There is a transition region that overlaps with the OH signal and this marks the inner flame branch. This inner branch is likely to be a diffusion flame between evaporated fuel and the air that recirculates that is mixed with hot products from the flame itself. Such a flame is less prone to extinction [26] and this might explain why the outer branch along the shear layer detaches and seems extinguished more often than the inner part. This is also consistent with the temporal analysis of visualisation of blow-off events in this spray flame by 5 kHz OH\* imaging [5] that showed that the inner part of the flame was the last to extinguish.

Individual images vary significantly. Some images show an attached flame (circled in Fig. 6), some present a flame lifted from the bluff-body edge (square), and some show an absence of flame sheet either in the outer shear layer (IIIa, IIb) or the inner recirculation zone (IVb). In the region close to the bluff body and radially outwards from the spray, a significant amount of  $CH<sub>2</sub>O$  is present and, when the OH shows a continuous sheet, the  $CH<sub>2</sub>O$  region is bounded by the OH. In conjunction with the laminar flame simulation results, that show that the CH<sub>2</sub>O-containing regions are fuel-rich ( $\zeta \zeta_{st}$ ), we may conclude that the region between the spray and the annular air is fuel-rich, while the region inside the spray cone is lean. Note also that the region between the spray and the annular air occasionally contain no CH2O signal. This can be either unburnt fuel or air penetrated through the sheet, both possibilities alluding to an extinction event. Row IV, Fig 6b, shows a region void of both OH and  $CH<sub>2</sub>O$ ; such regions tend to exist when the OH has breaks.

The image showing the product  $CH<sub>2</sub>O<sub>×</sub>OH$  representing heat release is very strongly correlated to the corresponding OH image, contrary to premixed flames [5,17,25] that show OH even in regions very far from the region of coincidence with the CH2O. This provides support to the use of OH-PLIF for marking reaction zones in such spray flames [5].

#### 3.3 *Averaged* CH2O×OH *images*

Figure 7 presents the average images of OH, CH2O, and HR from the two cases. Again, the HR regions are similar to the OH regions. The region of the recirculation zone between the spray and the shear layer has significant CH<sub>2</sub>O. Close to the blow off condition, the flame seems closer to the bluff body, as concluded from the mean OH and HR images. In addition, higher values of CH2O are seen, consistent with the laminar flame simulations that show higher CH2O mole fraction at rich mixture fractions as the strain rate increases (Fig. 2).

Figure 8 shows the Abel-transformed mean OH\* chemiluminescence images taken at the same conditions. These two images show very similar results to the mean HR images (Fig. 7). A relatively high intensity region at the inner reaction zone, aligned with the spray, and an outer branch, along the shear layer, are evident by both techniques. The inner branch starts a few mm from the fuel nozzle and the outer branch seems lifted, but by a distance that is decreasing at the high velocity condition (H1S2). The inner branch seems to have a higher mean heat release rate than the outer branch. These details are seen in both OH\* and HR data. As the OH\* images do not have interference from Mie scattering or fuel fluorescence, the agreement supports the view that the flame shape and location of heat release zones obtained with CH<sub>2</sub>O×OH images is trustworthy.

#### **4. Conclusions**

Joint CH2O and OH PLIF was applied to swirl n-heptane spray flames close to extinction. To assist the image interpretation, laminar counterflow non-premixed flame calculations were also performed. The simulations suggest that the product  $X_{CH2O} \times X_{OH}$ marks relatively well the spatial location of the true heat release, especially at high strain rates, but neglects the heat release in lean mixture fractions, which could amount to about 50% of the total at low strain rates. The CH2O-PLIF measurement has limitations due to the interference by fuel and PAH fluorescence. Nevertheless, because the laminar flame simulations show a sharp destruction of CH<sub>2</sub>O, soot precursors, and fuel at stoichiometry, we can assume that  $CH_2O-PLIF$  marks fuel-rich regions, and that the outline of the  $CH_2O$  region can be thought of as close to the stoichiometric mixture fraction iso-line.

The PLIF showed that the HR regions followed very well the OH regions, which were thin, characteristic of non-premixed combustion at high strain. With an increase in bulk air velocity, the flame was fragmented more and extinguished more often. The instantaneous images showed great variability and significant parts of the flame sheet could at times be extinguished. The good correlation of the HR image with the OH signal indicates that highspeed OH-PLIF measurements could be a reliable marker for HR regions in swirling spray flames. Compared with line-of-sight inverse Abel-transformed OH\* chemiluminescence measurements, the  $CH<sub>2</sub>O<sub>×</sub>OH$  heat release rate estimation reveals quite similar flame shape and locations.

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## **Tables**

Name	$U_a$ [m/s]	$U_a/U_{BO}$	Φ
H1S1	17.1	0.75	0.32
H <sub>1</sub> S <sub>2</sub>	20.0	0.88	0.27
H1B	22.8		0.24

Table 1. Flow conditions. For all, the fuel flow rate was 0.27 g/s.

## **Figures**



Figure 1: The burner (left) and laser system (right) schematic.



**Figure 2:** Temperature and various species molar fractions *vs.* mixture fraction for strain rate (a)  $100 s^{-1}$  and (b)  $500 s^{-1}$ .



**Figure 3:** Normalised mole fractions of OH and CH2O, normalised HRR, and normalised  $X_{\text{CH2O}} \times X_{\text{OH}}$  at strain rate (a) 100 s<sup>-1</sup> and (b) 500 s<sup>-1</sup>.



**Figure 4:** (a) HRR vs.  $X_{CH2O} \times X_{OH}$  (b) HRR vs.  $(X_{CH2O} + X_{C7H16}) \times X_{OH}$  for the indicated strain rate.



**Figure 5:** (a) Un-ignited spray, CH<sub>2</sub>O image; (b) example CH<sub>2</sub>O image of H1S1 flame. The colour map is logarithmic.



**Figure 6:** Instantaneous CH2O, OH, and images for H1S1 (a) and H1S2 (b). I-IV are discussed in the text. The CH2O image colour map is logarithmic.



**Figure 7:** Mean OH, CH<sub>2</sub>O, and CH<sub>2</sub>O×OH for (a) H1S1 and (b) H1S2.



**Figure 8:** Inverse Abel-transformed mean OH\* chemiluminescence images: (a) H1S1 and (b) H1S2.The black square marks the PLIF region.

#### **Figure captions**

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