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Electrospun black titania nanofibres: Influence of hydrogen plasma induced disorder on the electronic structure and photoelectrochemical performance

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3 KEYWORDS: Black Titania, Plasma Treatment, Photoelectrochemical Properties,
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5 Electrospinning.
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10 11 **Abstract**

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16 This work encompasses a facile method for tailoring surface defects in electrospun TiO₂
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18 nanofibres by employing hydrogen plasma treatments. This amiable processing method was
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20 proven with SQUID, EPR, and XPS to be highly effective in generating oxygen vacancies,
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22 accompanied by the reduction of Ti⁴⁺ centres to Ti³⁺, resulting in the formation of black titania.
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24 The treatment temperature was found to affect the Ti³⁺/Ti⁴⁺ ratios and surface valence, while
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26 preserving the original 1D morphology of the titania fibres. *Ab initio* DFT calculations showed
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28 that a high concentration of oxygen vacancies is highly efficient in producing midgap states that
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30 enhance the system absorption over the whole visible range, as observed with UV/Vis/NIR
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32 diffuse reflectance spectroscopy. Pristine TiO₂ nanofibres produced a photocurrent density of
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34 ~0.02 mA/cm² at 1.23 V vs. RHE, whereas the hydrogen plasma treatment resulted in up to a
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36 ten-fold increase in the photoelectrochemical performance.
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42 **1) Introduction**

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44 During the last decade, solar-assisted water splitting has received considerable attention for the
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46 sustainable production of hydrogen, an attractive eco-friendly energy vector, from largely
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48 abundant and renewable resources.¹⁻⁹

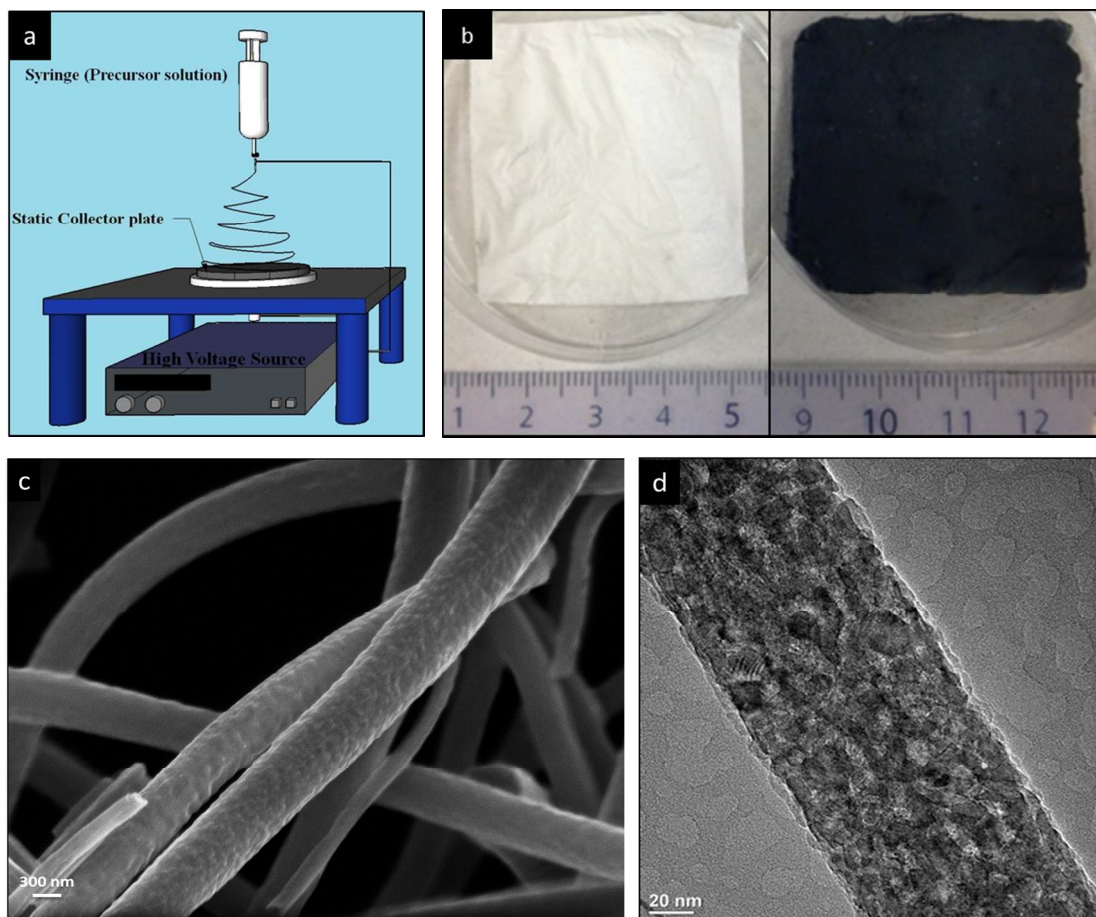
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51 Among the various water oxidation catalysts, TiO₂ has been one of the most extensively studied
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53 semiconductors, since the pioneering work of Fujishima and Honda in 1972.¹⁰ Nevertheless, the
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55 large band gap (3.2 eV) of TiO₂ limits its absorption especially in the visible part (400-800 nm)
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3 of the solar spectrum, which represents a bottleneck for its large-scale utilization.¹¹ This intrinsic
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5 limitation has been addressed in a large body of data related to structural (core-shell, multi-layer,
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7 composite) and electronic (doped, co-doped, composite) modifications of titania. Tailoring and
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9 nano-engineering of TiO₂ systems has been investigated by modification of both fabrication
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11 routes and *ex-situ* treatments of titania-based compositions.¹² To this aim, various efforts have
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13 been devoted to chemically functionalize titania by incorporation of metal¹³ or non-metal¹⁴
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15 entities. Approaches based on the surface modification of TiO₂ pigments by high pressure
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17 hydrogen treatment resulted in black titania powder.¹⁵⁻¹⁶ Although the exact nature of the
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19 occurring chemical modification is not completely understood, the formation of midgap states
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21 above the valence band maximum is likely responsible for the creation of Ti³⁺ centres within the
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23 TiO₂ lattice, thus darkening the color of TiO₂ by modifying the band gap, which is of key
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25 importance for photocatalytic processes.^{15, 17-18} So far black TiO₂ containing Ti³⁺ centres have
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27 been often accomplished by creating surface defects on the TiO₂ nano-particles/-rods under high
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29 pressure and extended reaction periods.¹⁵⁻¹⁶ Particular attention was devoted to the use of milder
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31 conditions and shorter processing time with respect to previous reports showing that hydrogen
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33 plasma treatment is effective in improving functional performance of hematite photoanodes.¹⁹

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36 Among the possible nanomaterial synthetic strategies, electrospinning is a very versatile
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38 approach for fabricating continuous and interwoven fiber networks with high active surface
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40 areas. One dimensional microstructure and reduced grain boundary regions ensure an efficient
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42 charge carrier transfer during photo-activated processes.²⁰⁻²² In this study, we have fabricated
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44 free-standing electrospun 1-D TiO₂ nanofibres and adopted hydrogen plasma treatments for the
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46 partial conversion of pristine TiO₂ to black titania.
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2) Experimental Section

TiO₂ nanofibres were produced *via* the electrospinning of Ti sols (Figure 1a).²³ Typically, a solution of 1.4 g of Ti(IV) tetra-*iso*-propoxide (Ti(OⁱPr)₄) in 5.0 mL of *iso*-propanol was mixed under vigorous stirring, followed by 0.318 mg acetic acid and 300 mg of polyvinylpyrrolidone (PVP, 1300000 g/mol) as the spinning-aid. As-obtained electrospun nanofibers were subjected to *ex-situ* annealing at 500 °C for 5 h to obtain a crystalline phase. The surface reduction was carried out using a Radio Frequency (RF)-plasma system (Domino CVD/PE-CVD, Plasma electronic), with a gas flow rate of 500 sccm, power 15 W and purity of 99,99 % in the temperature range 300-500° C for 3 h. The treatment of the target materials under hydrogen plasma resulted in a darkening of the initially white fibres, giving rise to black TiO₂, as shown in Figure 1b.



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5 **Figure 1** (a) Representative diagram showing an electrospinning process. (b) Optical difference
6 between a pristine TiO₂ nanofibers (white) and hydrogen plasma treated fibres (black).
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10 (c) FE-SEM- and (d) HR-TEM- of the untreated TiO₂ nanofibres.

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18 X-ray photoelectron spectroscopy (XPS) characterizations were conducted on a custom
19 designed UHV system using Mg K_α excitation at a photon energy of 1253.6 eV and a Specs 100
20 hemispherical analyzer at 10 eV pass energy. The base pressure of the system is 3·10⁻¹⁰ bar. The
21 fitting of the XPS spectra was done using the program XPSPEAK Version 4.1. The
22 concentration of the detected Ti³⁺ and Ti²⁺ species in relation to Ti⁴⁺ were calculated using the
23 relative areas of their 2p_{3/2} peaks. UV

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32 All PEC measurements were performed in a three-electrode electrochemical cell to obtain
33 current-voltage (I-V) characteristics of the hydrogen plasma treated nanofibres and of the
34 untreated pristine TiO₂ fibres. The fibrous networks were fabricated as electrodes by attaching
35 the electrospun and annealed nanofibres to titanium gel covered FTO glass substrates (FTO-
36 TEC, Sigma-Aldrich, ~8 Ω/sq) and heating to 500 °C for two hours, after which the electrodes
37 were hydrogen plasma treated. Exposed FTO on the electrodes was connected to a silver wire
38 with silver paste and covered with an epoxy resin. The electrodes had a well defined active
39 surface area of 0.5 cm² and were employed as the working electrode under dark and illuminated
40 conditions with a 150 W Xenon lamp (Newport, Model: 67005) and an AM 1.5G filter. 1M
41 NaOH electrolyte (pH = 13.6) was used as the electrolyte to record I-V characteristics at varying
42 applied potentials from +0.8 V/SCE (anodic bias) to -0.8 V/SCE (cathodic bias) with a scan rate
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3 of 10 mV/s for all samples. The applied potentials versus SCE were calculated versus the
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5 reversible hydrogen electrode (RHE) using the following equation:¹⁹
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$$E_{RHE} = E_{SCE} + E_{SCE}^0 + 0.059pH$$

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15 The photocurrent density was calculated from the linear-sweep voltammograms under dark
16 and illumination conditions and plotted as a function of applied potential.
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27 **3) Results and Discussion**

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29 Field Emission-Scanning Electron Microscopy (FE-SEM) and High Resolution-Transmission
30 Electron Microscopy (HR-TEM) micrographs (Figures 1c and 1d) revealed the spatial
31 organization of highly interwoven nanofibres with diameter ~400 nm. The performed plasma
32 treatment did not induce any appreciable variation in the morphology of the fibres such as
33 increased roughness, cleavage, peeling or branching phenomena [Figures 3–6 in the Electronic
34 Supplementary Information (ESI)]. Irrespective of the adopted synthesis and plasma-treatment
35 conditions, X-ray diffraction (XRD) indicated the formation of the *anatase* phase (Figure 2a;
36 JCPDS Card: 78-2486), free from other possible titania polymorphs. A careful inspection of the
37 XRD results show a decrease in the peak intensity upon increasing the hydrogen plasma
38 treatment temperature. This phenomenon, particularly evident for the (101) reflection, could be
39 related to the formation of atomistic surface defects that do not cause any extensive surface
40 roughening, since the plasma power used in our treatment is lower than that required for
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etching.²⁴⁻²⁵ Nevertheless, the bombardment of the target nanomaterials by hydrogen ions with high kinetic energy produces surface vacancies by removing oxygen species, thus increasing the structural disorder of the nanofiber surfaces.²⁶ In addition, the chemical composition of the nanofibres can change due to the incorporation of hydrogen into the TiO₂ lattice to form gradient (TiO₂@TiO_{2-x}H_x) structures with reduced Ti centers [Ti(IV) → Ti(III)].

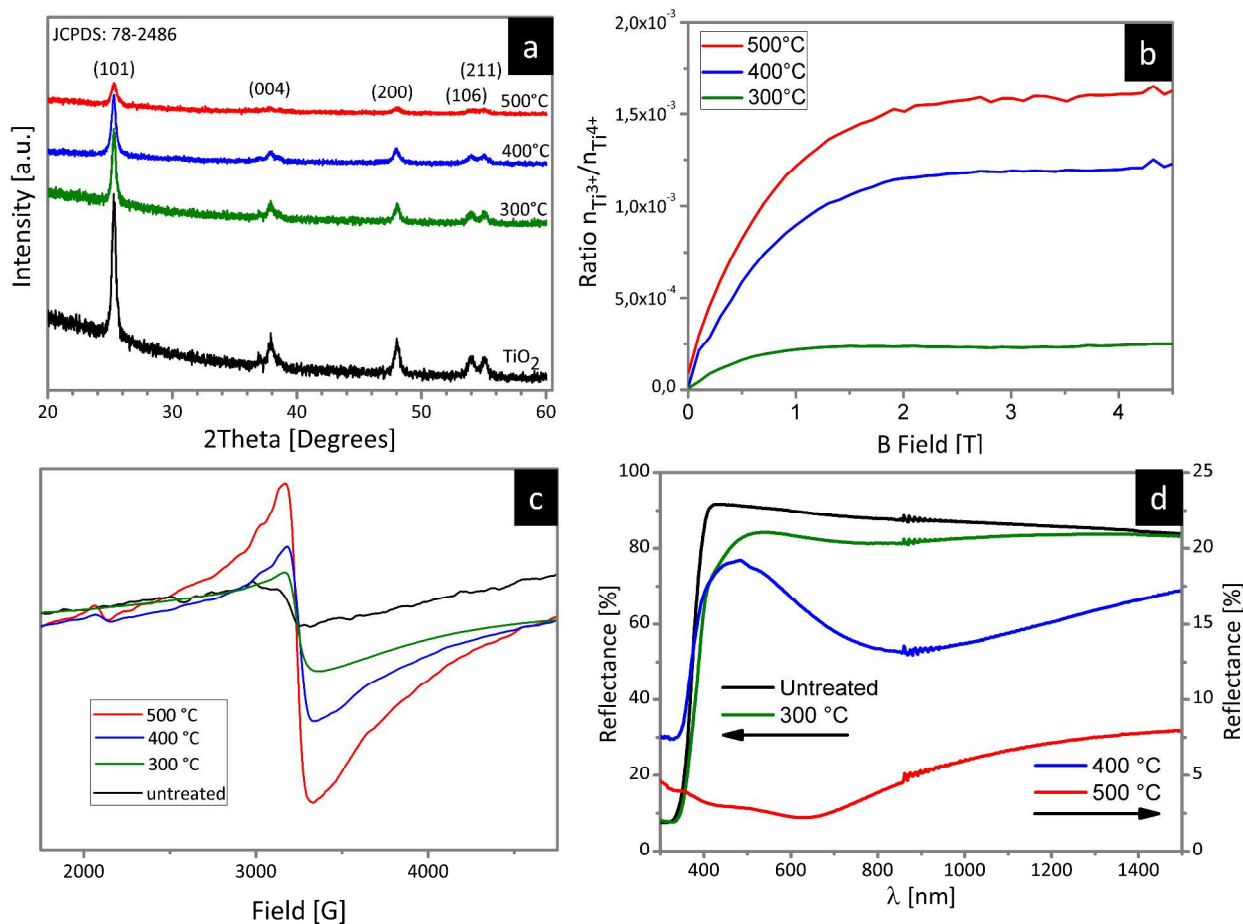


Figure 2 (a) XRD patterns of the untreated TiO₂ and H₂ plasma treated TiO₂ nanofibres different plasma treatment temperatures (b) SQUID measurements of H₂ plasma treated TiO₂ samples performed at 2 K. (c) EPR measurements of the pure TiO₂ and H₂ plasma treated TiO₂. (d) UV/Vis/NIR reflectance spectra of the untreated and H₂ plasma treated TiO₂ samples.

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3 In order to attain a deeper insight into the mentioned phenomena, particularly regarding the
4 formation of oxygen vacancies upon hydrogen plasma treatment, magnetic measurements were
5 carried out. Whereas pure anatase TiO_2 is diamagnetic, the partial reduction of Ti^{4+} centers to
6 Ti^{3+} upon formation of oxygen defects leads to paramagnetic behavior. SQUID measurements of
7 the plasma treated samples clearly showed that an increase in the processing temperature
8 corresponded with an overall increment in the content of paramagnetic species in the fiber
9 surfaces (Figure 2b). This evidences that H_2 plasma treatment induces the formation of oxygen
10 vacancies and the partial reduction of the neighbouring Ti^{4+} centers to Ti^{3+} , a phenomenon
11 further promoted by an increased processing temperature.²⁶

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13 To corroborate the presence of paramagnetic Ti^{3+} centers in the target systems and to
14 substantiate the magnetic measurements, Electron Paramagnetic Resonance (EPR) spectroscopy
15 was performed. In particular, X-Band EPR-spectra recorded before and after H_2 plasma
16 treatment of TiO_2 fibres exhibited EPR signals for all samples. The low intensity signals with
17 g -values of 1.982 G are assigned to Ti^{3+} species.²⁷ Figure 2c shows the EPR-spectrum of the
18 TiO_2 fibres H_2 plasma treated at 500 °C, 400°C and 300°C compared to that of untreated TiO_2 .
19 The signal intensity significantly increased upon hydrogen plasma treatment when compared to
20 the untreated sample. The EPR results suggest that the electrons are located in the d-orbitals.
21 Thus, $J = S = \frac{1}{2}$ is a fair approximation for the electrons, and the saturation moment for each Ti^{3+}
22 equals 1 μB . Hence, the saturation moment for each curve represents the relative amount of Ti^{3+}
23 in a diamagnetic matrix.²⁸

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25 Optical reflectance spectra versus a Spectralon reference were measured in the UV (300 nm) to
26 near IR (1500 nm) region as seen in Figure 2d. The untreated TiO_2 fibres have no absorption in
27 the visible and near infrared regions with the band gap absorption beginning at 405 nm (optical
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band gap ~ 3.1 eV). The H_2 plasma treatment has a drastic effect on the sample absorption as indicated by the observed gradual color change to black. The plasma treatment at $300^\circ C$ results in a broad but low intensity absorbance around 800 nm and a redshift of the band gap absorption to 450 nm. Plasma treatment at higher temperatures results in a noticeably lowered reflectance in the whole measurement range. The broad absorption band has a maximum at 860 nm on the $400^\circ C$ plasma treated sample, while the band gap absorption remains practically unchanged from the $300^\circ C$ plasma treated sample. On the other hand, no TiO_2 specific band gap absorption is observed anymore for the $500^\circ C$ plasma treated sample. The increased absorption of this sample has a maximum at 630 nm due to localised empty midgap states lying ~ 2.0 eV above the valence band level that coincides with the previously reported band gap of black TiO_2 .¹⁵ Therefore, from the diffuse reflectance data we deduce that the lack of a TiO_2 specific band gap absorption in the UV range suggests a saturation of the defects on the surface of the nanofibres, which is further supported by the drop in the XRD peak intensity (Figure 2a).

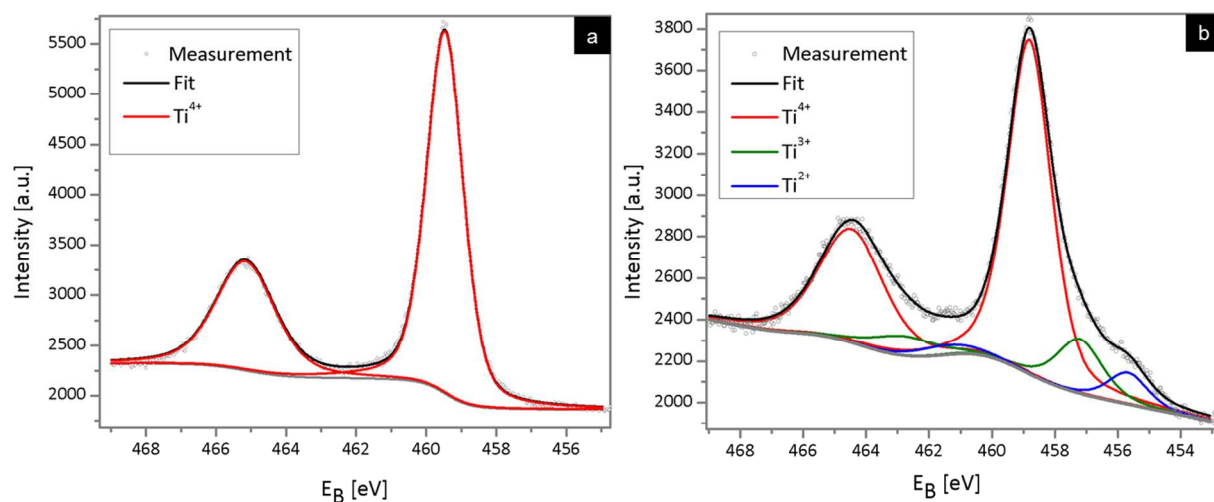


Figure 3 XPS spectra Ti $2p$ core level peaks for: (a) untreated TiO_2 and (b) hydrogen plasma treated TiO_2 samples at $500^\circ C$.

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5 Figure 3a displays the XPS spectra for the Ti $2p$ core level of nanofibres before hydrogen
6 plasma treatment. The spin orbit splitting between the two p peaks is 5.7 eV and the intensity
7 ratio is 2:1, as expected. The $2p_{3/2}$ component can be well fitted with a single peak, centered at a
8 binding energy (BE) of 459.5 eV, having a full width at half maximum (FWHM) of 1.24 eV.
9 These values are characteristic for octahedral Ti^{4+} in TiO_2 .²⁹⁻³⁰ In contrast, hydrogen plasma
10 treated fibres at 500°C (Figure 3b) show a change in the spectrum shape, with the appearance of
11 low-BE shoulders that suggest the co-presence of Ti oxidation states lower than 4. Spectral
12 deconvolution showed the presence of two additional bands with BEs of 457.20 eV and 455.68
13 eV (FWHM = 1.71 and 1.49 eV, respectively) which can be attributed to $2p_{3/2}$ peaks of Ti^{3+} and
14 Ti^{2+} .³¹⁻³² The creation of Ti^{2+} species could be a direct result of the high degree of reduction
15 during the hydrogen plasma treatment.³³ The relative amounts of Ti^{3+} and Ti^{2+} formed as a result
16 of the valence electron transition compared with Ti^{4+} were 13.5% and 6.8%, respectively. XPS
17 spectra for samples treated with hydrogen plasma at 300°C and 400°C do not show additional
18 oxidation states other than Ti^{4+} , despite a continuous peak broadening with the increase in the
19 hydrogen plasma treatment temperatures (Figure 8 in the ESI).
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40 *Ab-initio* calculations within the GGA+U framework of density functional theory (DFT) have
41 been used in order to investigate the influence of oxygen vacancies on the electronic structure of
42 anatase TiO_2 .³⁴⁻³⁶ A direct comparison to the effect of lattice disorder as proposed by Chen *et. al*
43 is made.¹⁰ A special focus is put on the effect of disorder in the crystal structure. Our results
44 indicate that a mild hydrogen treatment process that leaves the crystal structure intact is
45 favourable compared to hydrogen treatment processes that generate highly disordered TiO_2
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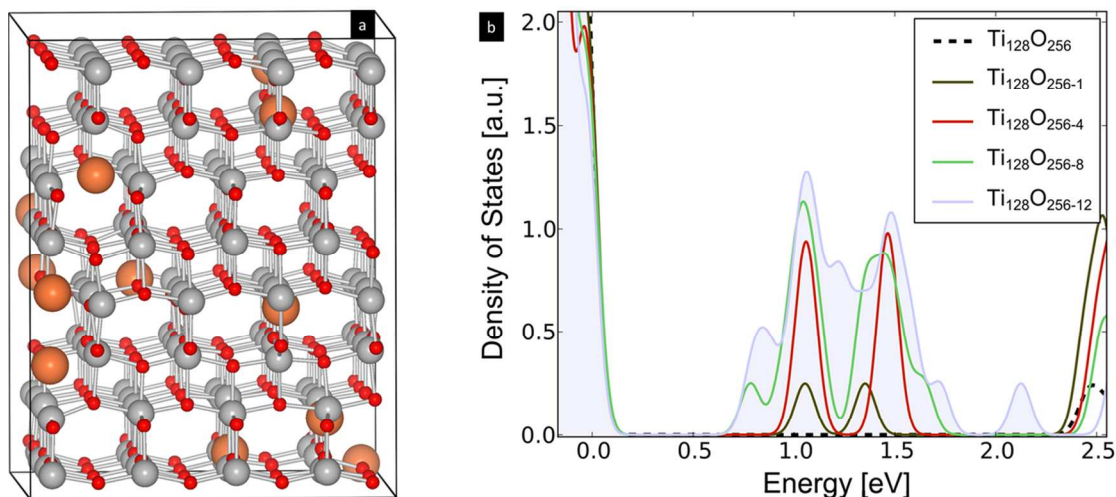


Figure 4 a) Anatase TiO_2 supercell with the composition $\text{Ti}_{128}\text{O}_{256-12}$. The 12 oxygen vacancies on random positions in the anatase cell are shown as large brownish spheres. Ti ions are shown in grey and oxygen ions in red. b) Density of electronic states (DoS) of TiO_2 for varying oxygen vacancy concentrations. Occupied energy levels are indicated by the shaded areas. With increasing oxygen vacancy concentrations the spectrum of midgap states formed by localized Ti^{3+} centers broadens.

In the first step, oxygen vacancies were generated with different concentrations in random sites in the TiO_2 lattice, shown in Figure 4a. The density of states (DoS) for the different configurations was calculated after a full relaxation of the employed supercells with the compositions $\text{Ti}_{128}\text{O}_{256-x}$, $x=0, 1, 4, 8, 12$ and is shown in Figure 4b. Below oxygen vacancy concentrations of 0.9 nm^{-3} , *i.e.* for the supercell $\text{Ti}_{128}\text{O}_{256-4}$, the DoS reveals two distinct peaks of occupied electronic states, which are localized at different Ti^{3+} centers near the defect sites. This is in agreement with previous work where single oxygen vacancies were investigated.³⁶ Upon increasing the defect concentration, the two peaks broaden considerably and finally merge, yielding a broad spectrum of midgap states. These results indicate that a rather high

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3 concentration of oxygen vacancies ($\sim 1\text{--}2\text{ nm}^{-3}$) is needed in order to enhance the absorption of
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6 TiO_2 over a broad photon spectrum, supporting the conclusion from DRS data that the surface
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8 defect density is possibly saturated for the sample treated at $500\text{ }^\circ\text{C}$.
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10 The SQUID measurements show that the concentrations of Ti^{3+} centers in the overall
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12 hydrogenated samples are rather small (see Figure 2b). This suggests that in our experiments
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14 oxygen vacancies are mainly formed within a small region near the surface of the nanofibres, due
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16 to the limited plasma penetration into the TiO_2 nanofibres.³⁷ This is consistent with the enhanced
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18 values of the $\text{Ti}^{3+}/\text{Ti}^{4+}$ ratio for our hydrogen treated samples in the surface sensitive XPS
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20 measurements.
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24 The TiO_2 (001) surface is usually considered to be catalytically more active for oxygen
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26 evolution reactions.³⁸ Subsequently, the electronic structure of the (001) TiO_2 surface in the
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28 presence of oxygen vacancies was simulated, investigating also the effect of lattice disorders,
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30 including hydrogen interstitials.¹⁶
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34 To this aim, we employed TiO_2 *anatase* slabs (thickness 2.8 nm) terminated by (001) surfaces
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36 and periodically repeated in the transverse directions. Whereas the lower half of the slabs were
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38 unmodified, 6 oxygen vacancies and 12 interstitial hydrogens were incorporated in the upper half
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40 of the unit cell with composition $\text{Ti}_{108}\text{O}_{216}$ (Figure 5a), corresponding to impurity concentrations
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42 of 3.1 nm^{-3} and 6.2 nm^{-3} , respectively. Since an interstitial hydrogen impurity donates one
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44 electron and a single oxygen vacancy creates two Ti^{3+} states, the same amount of Ti^{3+} states are
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46 generated in the two cases. The hydrogenated slab was additionally annealed to 3000 K for 2 ps
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48 and subsequently fully relaxed. The atoms in the lower half of the slab were kept fixed. In this
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50 way the model for a disordered surface structure shown in Figure 5a was obtained.
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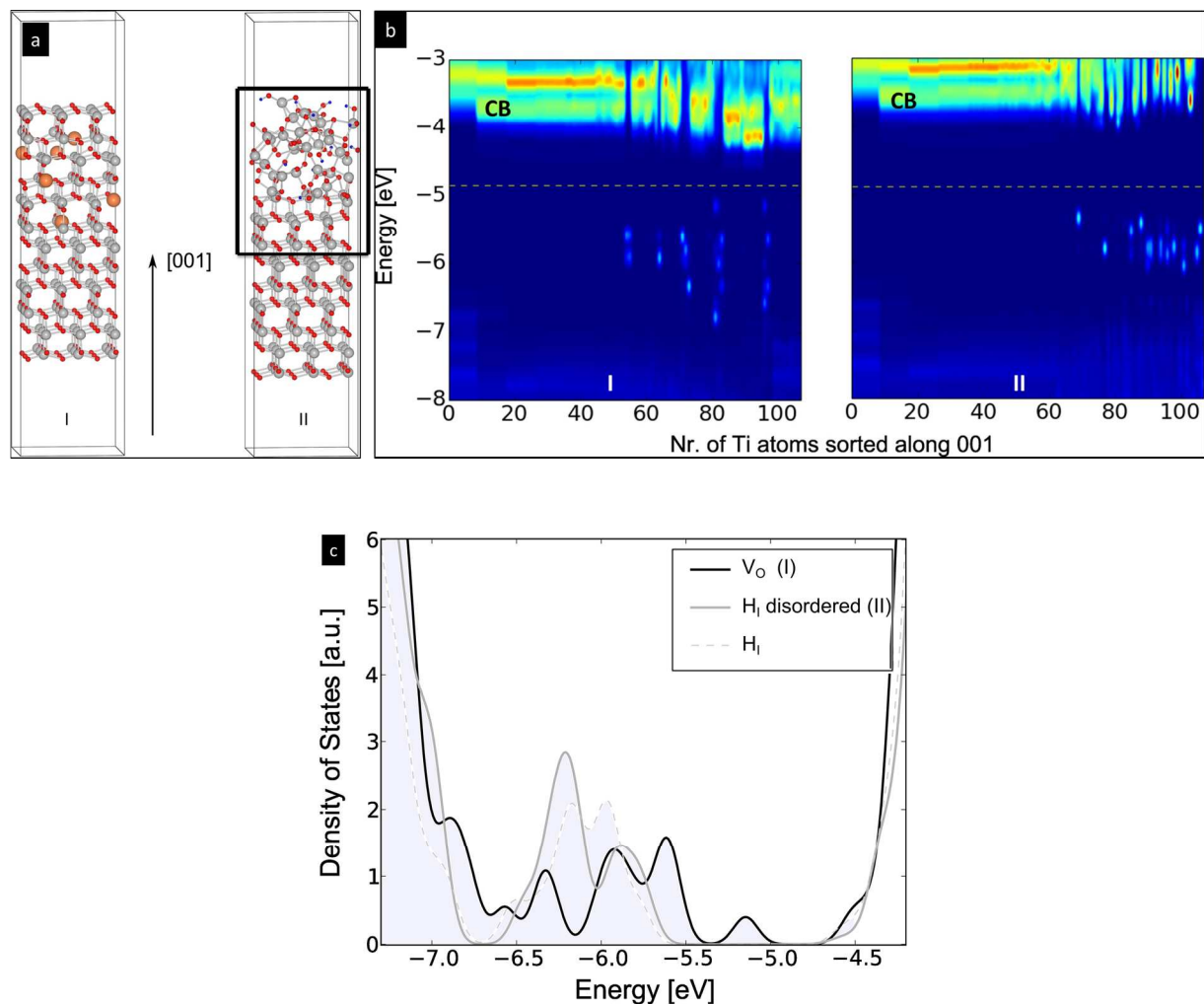


Figure 5 (a) TiO₂ slab with an oxygen vacancy concentration of approximately 3.1 nm⁻³ (I) and a disordered crystal structure with hydrogen interstitial concentration of 6.2 nm⁻³ (II) in the upper half within the black frame. (b) DoS projected on the Ti atoms for the simulated slab configurations I and II with oxygen vacancies and interstitial hydrogen atoms in a disordered lattice. The Fermi level is indicated by the yellow dashed lines. The Ti orbitals mainly form the midgap states as well as the conduction band (CB). (c) Densities of states (DoS) of the surface configurations I with oxygen vacancies (V_O) and II with a disordered crystal structure and interstitial H atoms (H_I disordered). Additionally, the DoS of interstitial H atoms in an ordered crystal structure (H_I) is depicted. The shown DoS is obtained by a summation of the atom

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3 projected DoS of the atoms in the upper half of the slabs. Occupied energy levels are indicated
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5 by the shaded areas. The vacuum potential is taken as the zero energy reference.
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9 In both cases, the atom projected density of states (PDoS) was calculated. The PDoS confirms
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11 that midgap states are introduced by both oxygen vacancies and hydrogen interstitials, with the
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13 states being occupied and located at the Ti sites, see Figure 5b. Hence, the formation of Ti^{3+}
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15 centres, as observed by XPS analyses, is substantiated. However, in the case of oxygen
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17 vacancies, the energy spectrum of the midgap states is considerably broader for identical Ti^{3+}
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19 concentrations. From the PDoS we also find that the disordered crystal structure of the annealed
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21 sample lead to a highly fluctuating energy spectrum at the conduction band edges as shown in
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23 Figure 5b. The same behavior is also found at the valence band edges (not shown). This indicates
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25 the formation of localized states at the band edges, which likely act as traps for electrons and
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27 holes, respectively. On the other hand, the incorporation of oxygen vacancies into the ordered
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29 crystal structure leaves the energy spectrum at the band edges relatively smooth, and therefore
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31 indicates superior transport properties compared to the disordered hydrogenated system.
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38 The finding that the midgap states show higher peaks for the hydrogenated disordered sample,
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40 while the configuration with the oxygen vacancies shows a broader level distribution, is also
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42 verified by analysing the sum of the densities of states projected on the atoms of the upper half of
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44 the slabs (Figure 5c).
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47 The density of states of a configuration with interstitial hydrogen defects in an ordered TiO_2
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49 crystal structure was also investigated. Here we again observe a rather narrow distribution of
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51 midgap states. However, at the hydrogen interstitial concentration of 6.2 nm^{-3} , we find that the
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53 occupation of states at the conduction band minimum indicates that at such high hydrogen
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55 doping levels delocalized charge carriers are generated, leading to enhanced conductivity of the
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3 TiO₂. Since the improvement of the electrical transport properties is considered as one of the
4 major issues that have to be addressed for highly efficient metal oxide photoelectrodes,³⁹ the
5 incorporation of a high concentration of hydrogen interstitials into the ordered *anatase* lattice
6 may lead to enhanced photoelectrochemical performance.
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12 The performed simulations indicate that compared with hydrogen interstitials the oxygen
13 vacancies are highly efficient in generating a broad spectrum of midgap states, enhancing light
14 absorption over a broad photon spectrum. This finding explains why a relatively mild plasma
15 treatment that generates oxygen vacancies is sufficient for the fabrication of black TiO₂, which
16 maintains its crystal structure intact, a very attractive aspect concerning the mobility of charge
17 carriers. We also observe a strong localization of states in the disordered slab configuration with
18 interstitial hydrogen atoms, which can explain a significant drop of charge carrier mobility in
19 disordered TiO₂ lattices.⁴⁰
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32 In order to elucidate the photoactivity of the nanofibres, photoelectrochemical measurements
33 (PEC) were carried out, shown in Figure 6. We observe a gradual increase in the photocurrent
34 density (j) at 1.23 V *vs.* RHE from ~0.02 mA/cm² for the pristine TiO₂ to a maximum of
35 ~0.2 mA/cm² for the TiO₂ nanofibres H₂ plasma treated at 500°C. This indicates a tenfold
36 increase in the photoelectrochemical performance due to the increased absorption of the fibres.
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39 The fact that photocurrent densities increase gradually with higher treatment temperatures
40 indicates clearly that the process temperature assists the reduction of the TiO₂ nanofibres. In
41 addition, an increase in the treatment temperature resulted in a cathodic shift of the photocurrent
42 onset voltage. This shift illustrates a reduction of the applied bias that is required to produce
43 photocurrent, enhancing the photoelectrochemical performance of the working electrode. The
44 growth of the photocurrent is linear for the pristine TiO₂ after ~0.6 V *vs.* RHE, whereas this
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value is lowered to 0.3 V (300°C), 0.25 V (400°C) and 0.2 V (500°C) vs. RHE for the H₂ plasma treated fibres. The rapid increase in the photocurrent densities at lower potentials demonstrates improved charge separation and transport of charge carriers. These enhancements due to hydrogen plasma treatment under mild conditions, combined with the outcomes of theoretical simulations clearly demonstrate that an improvement in the overall PEC performance occurs due to the creation of oxygen vacancies. When compared with existing data on 0–2 dimensional TiO₂ nanostructures, photocurrent densities achieved in this study are relatively modest.⁴¹ This discrepancy can be explained by the sparse contact at the interface between the conducting glass and the nanofibres, hindering efficient transfer of photogenerated carriers and producing at least a partial recombination. This is clearly observed for the untreated fibres as well. Further work to circumvent this problem is already in progress.

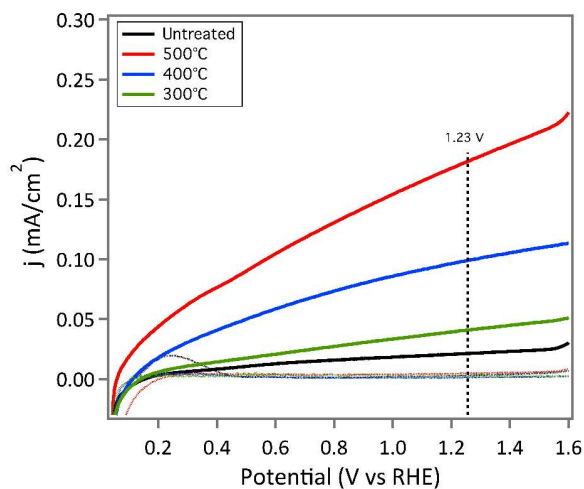


Figure 6 Photocurrent density vs. applied potential curves for *titania* nanofibres treated under different conditions. Dotted and solid lines represent the measurements in the dark and under illumination, respectively. For clarity, the j values at 1.23 V vs. RHE *i.e.* at the theoretical water oxidation potential, are marked by a dashed line.

Conclusions

In summary, this work presented a facile route for tailoring the surface defects in TiO₂ nanostructures fabricated by electrospinning by employing mild hydrogen plasma treatments. The adopted modification process was proven to be effective in generating oxygen vacancies in TiO₂, accompanied by the reduction of Ti⁴⁺ centres to Ti³⁺. Variations in the treatment temperatures were observed to be able to finely modulate the Ti³⁺/Ti⁴⁺ ratios and valence dynamics on the surface, while preserving the original one dimensional morphology of the titania nanofibres. Computational studies demonstrated that a high concentration of oxygen vacancies resulted in a broad spectrum of midgap states that enhanced the system absorption over the visible range. The increased absorption resulted in greatly enhanced PEC performances. However, the introduction of disorder in the TiO₂ lattice was identified to be detrimental for the electronic transport properties. Interestingly, this could be tailored as a function of the temperature adopted during the plasma processing. These results disclose attractive perspectives for the engineering of photoelectrochemical systems that could be adapted to a broad range of usages.

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5 ASSOCIATED CONTENT:
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9 **Supporting Information**

10 Detailed procedures for the fabrication of TiO₂ nanofibres, hydrogen plasma treatment and
11 characterization such as EDXS Mapping and HR-TEM are included in the Supporting
12 Information.
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