

A Discontinuous Galerkin Formulation for Solid Dynamics

J. Peraire^{*}, J. Bonet^{**}, A. Huerta⁺, P.-O. Persson^{*} and Y. Vidal^{+*}

^{*} Massachusetts Institute of Technology

^{**} University College Swansea

⁺ Polytechnic University of Catalonia



WCCM - Los Angeles, July 20, 2006

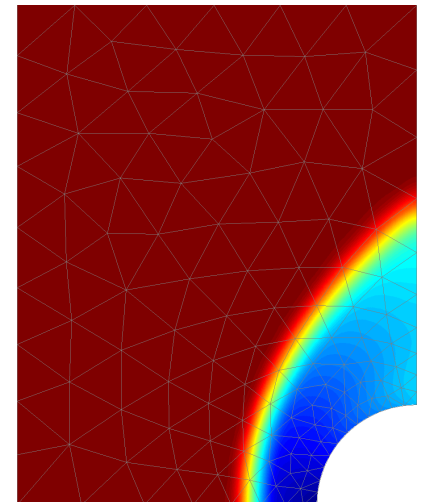
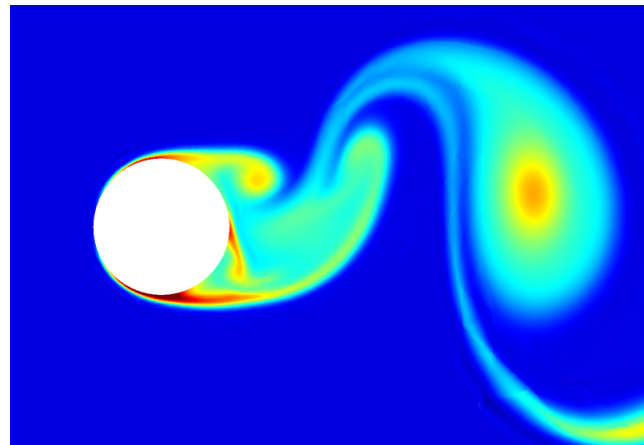
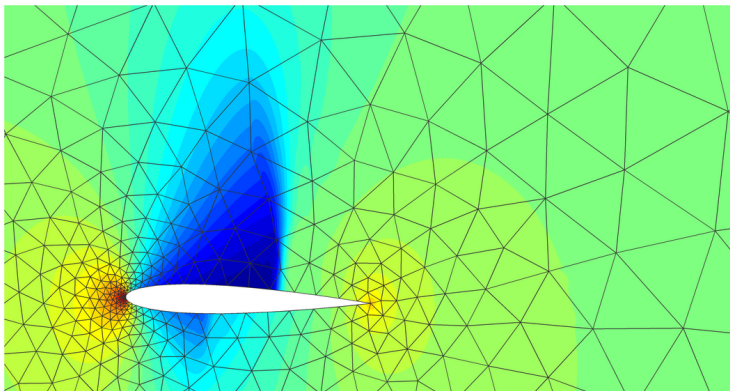
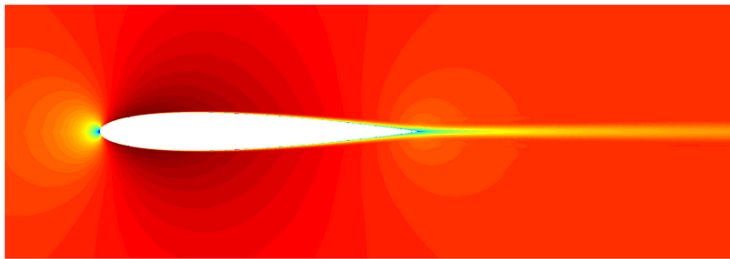
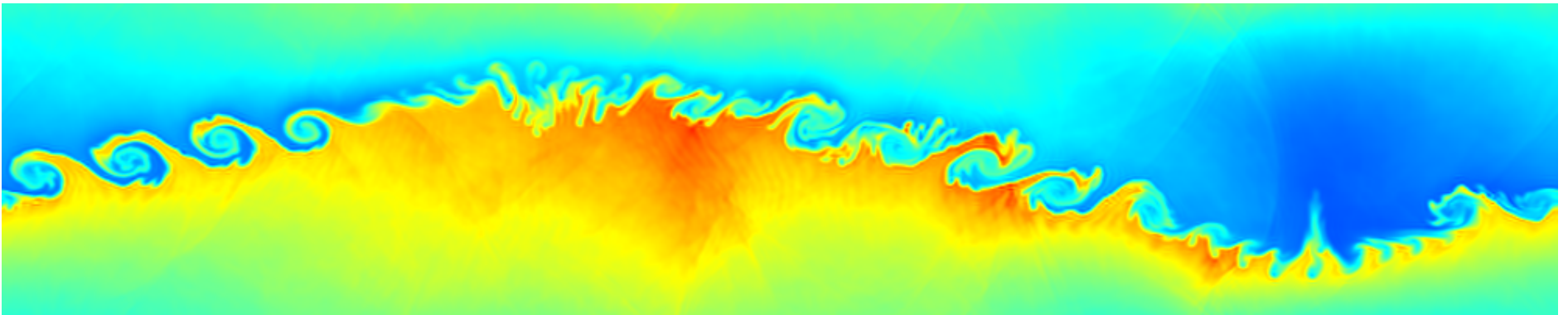


Acknowledgments

- Sandia National Labs
- Singapore MIT Alliance
- Department of Research - Catalan Government

DG for Eulerian CFD

High Resolution, Complex Geometries/Physics, $0.001 \leq M \leq 10 - 20$



Objectives

- General Materials, e.g. Neo-Hookean, Visco-elasto-plastic
- Arbitrary Description (Lagrangian/Eulerian)
- High Energy/Large Deformation
- Shock waves incorporating correct entropy production
- Accurate resolution of isentropic shear/volumetric waves
- Low Numerical Dissipation/Dispersion
- Multi-material Interactions
- Accuracy Guarantees

Approach

- Cast Governing Equations as System of First Order Conservation Laws
- Account for involutions by appropriate selection of approximating spaces
- Use High-Order DG
- For Dissipative/Non-local Fluxes introduce local variables (LDG/CDG)
- Use “standard” shock capturing techniques

Here, we will present basic formulation for Lagrangian non-dissipative Neo-Hookean material

The Discontinuous Galerkin Method

- (Reed/Hill 1973, Cockburn/Shu 1998-2001, etc)
- Consider non-linear hyperbolic system in conservative form:

$$\mathbf{u}_t + \nabla \cdot \mathcal{F}(\mathbf{u}) = \mathcal{S}(\mathbf{u})$$

- Triangulate domain Ω into elements $\kappa \in T_h$
- Seek approximate solution \mathbf{u}_h in space of element-wise polynomials:

$$V_h^p = \{ \mathbf{v} \in L^2(\Omega) : \mathbf{v}|_{\kappa} \in P^p(\kappa) \quad \forall \kappa \in T_h \}$$

- Multiply by test function $\mathbf{v}_h \in V_h^p$ and integrate over element κ :

$$\int_{\kappa} [(\mathbf{u}_h)_t + \nabla \cdot \mathcal{F}(\mathbf{u}_h) - \mathcal{S}(\mathbf{u}_h)] \mathbf{v}_h \, d\mathbf{x} = 0$$

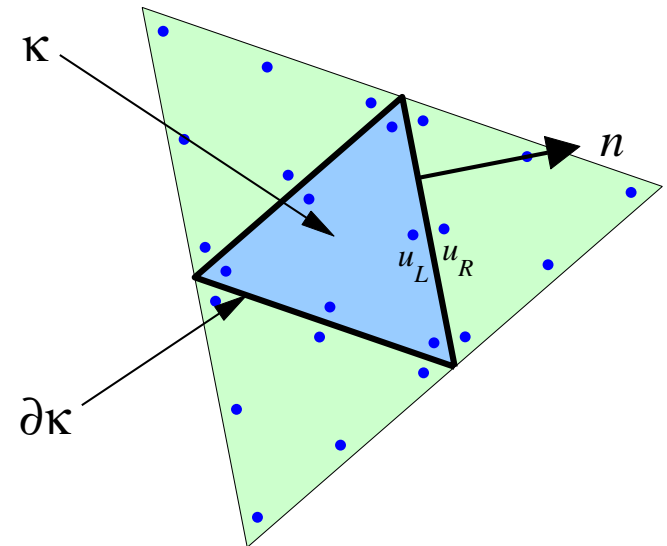
The Discontinuous Galerkin Method

- Integrate by parts:

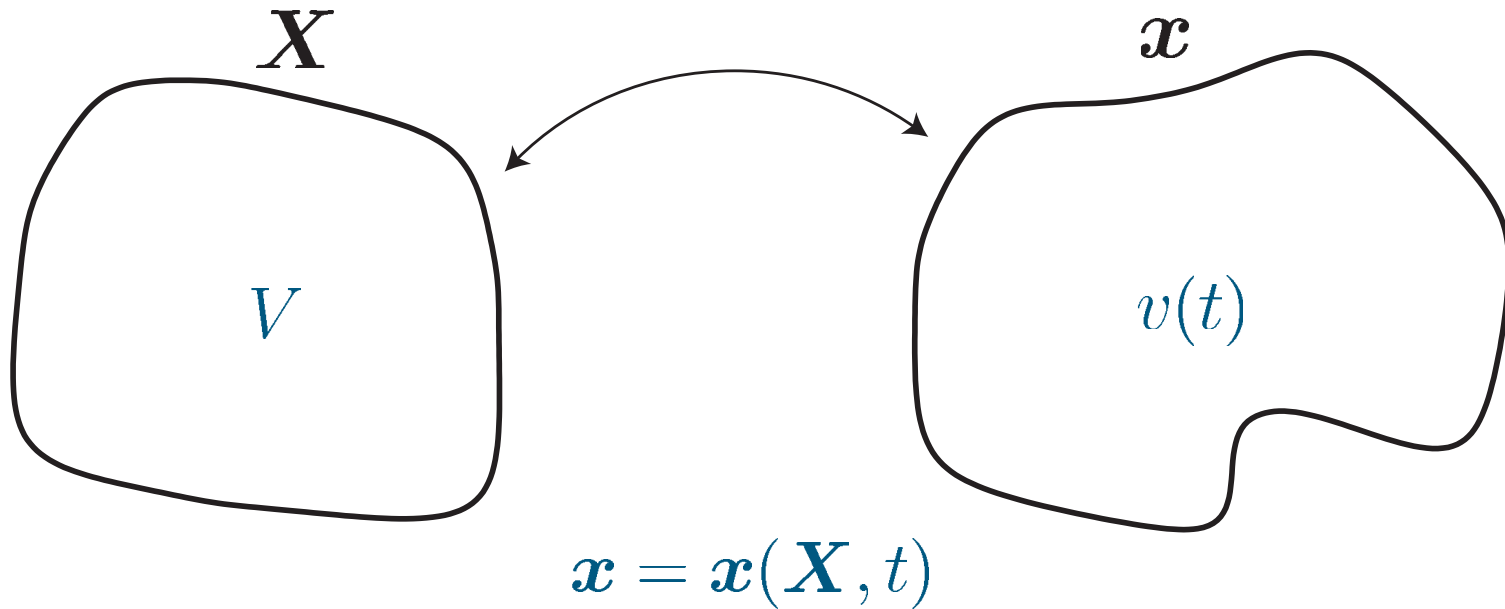
$$\int_{\kappa} [(\mathbf{u}_h)_t] \mathbf{v}_h d\mathbf{x} - \int_{\kappa} \mathcal{F}(\mathbf{u}_h) \nabla \mathbf{v}_h d\mathbf{x} + \int_{\partial\kappa} \mathcal{H}(\mathbf{u}_h^+, \mathbf{u}_h^-, \hat{\mathbf{n}}) \mathbf{v}_h^+ ds = \int_{\kappa} \mathbf{s}(\mathbf{u}_h) \mathbf{v}_h d\mathbf{x}$$

with numerical flux function $\mathcal{H}(\mathbf{u}_L, \mathbf{u}_R, \hat{\mathbf{n}})$ for left/right states $\mathbf{u}_L, \mathbf{u}_R$ in direction $\hat{\mathbf{n}}$ (Godunov, Roe, Osher, Van Leer, **Lax-Friedrichs**, etc)

- Global view: Find $\mathbf{u}_h \in \mathcal{V}_h^p$ such that this weighted residual is zero for all $\mathbf{v}_h \in \mathcal{V}_h^p$
- Spatial Error $\mathcal{O}(h^{p+1})$ for smooth problems
- Explicit Time Integration Using RK Method



Lagrangian Formulation



$$\mathbf{F} = \frac{\partial \mathbf{x}}{\partial \mathbf{X}}, \quad J = \det \mathbf{F}, \quad \mathbf{v}(\mathbf{X}, t) = \frac{\partial \mathbf{x}}{\partial t}, \quad \mathbf{p} = \rho_0 \mathbf{v}$$

Conservation Laws

- Momentum

$$\frac{d}{dt} \int_V \mathbf{p} dV = \int_V \rho_0 \mathbf{b} dV + \int_{\partial V} \mathbf{t} dA, \quad \mathbf{t} = \mathbf{P} \mathbf{N}$$

$$\frac{\partial \mathbf{p}}{\partial t} = \nabla \cdot \mathbf{P} + \rho_0 \mathbf{b}$$

- Deformation Gradient

$$\frac{\partial \mathbf{F}}{\partial t} = \nabla (\mathbf{p} / \rho_0) \quad \text{or} \quad \frac{\partial \mathbf{F}}{\partial t} = \nabla \cdot \left(\frac{1}{\rho_0} \mathbf{p} \otimes \mathbf{I} \right)$$

$$\frac{d}{dt} \int_V \mathbf{F} dV = \int_{\partial V} \frac{1}{\rho_0} \mathbf{p} \otimes \mathbf{N} dA$$

“generalized” form of continuity

Constitutive Model

- Strain Energy Potential

$$\psi(\mathbf{F}) = \psi_{iso}(J^{-1/3}\mathbf{F}) + \psi_{vol}(J)$$

$$\psi_{iso} = \frac{\mu}{2}[\text{tr}(J^{-2/3}\mathbf{F} : \mathbf{F}) - 3], \quad \psi_{vol} = \frac{1}{2}\kappa(J - 1)^2$$

- First Piola-Kirchhoff Stress Tensor

$$\mathbf{P} = \mathbf{P}_{dev} + \mathbf{P}_{vol}; \quad \mathbf{P}_{dev} = \frac{\partial\psi_{iso}}{\partial\mathbf{F}}, \quad \mathbf{P}_{vol} = \frac{\partial\psi_{vol}}{\partial\mathbf{F}}$$

$$\mathbf{P}_{vol} = \frac{d\psi_{vol}(J)}{dJ} \frac{\partial J}{\partial\mathbf{F}} = pJ\mathbf{F}^{-T}; \quad p = \frac{d\psi_{vol}(J)}{dJ} = \kappa(J - 1)$$

$$\mathbf{P}_{dev} = \mu J^{-2/3} \left[\mathbf{F} - \frac{1}{3}(\mathbf{F} : \mathbf{F})\mathbf{F}^{-T} \right]$$

Conservative Form of the Governing Equations

$$\frac{\partial \mathbf{p}}{\partial t} = \nabla \cdot \mathbf{P}(\mathbf{F}) + \rho_0 \mathbf{b}$$

$$\frac{\partial \mathbf{F}}{\partial t} = \nabla \cdot \left(\frac{1}{\rho_0} \mathbf{p} \otimes \mathbf{I} \right)$$

$$\Rightarrow \quad \mathcal{U}_t + \nabla \cdot \mathcal{F}(\mathcal{U}) = \mathcal{S}(\mathcal{U})$$

$$\mathcal{U} = \begin{pmatrix} \mathbf{p} \\ \mathbf{F} \end{pmatrix}, \quad \mathcal{F} = \begin{pmatrix} \mathbf{P} \\ \frac{1}{\rho_0} \mathbf{p} \otimes \mathbf{I} \end{pmatrix}, \quad \mathcal{S} = \begin{pmatrix} \rho_0 \mathbf{b} \\ \mathbf{0} \end{pmatrix}$$

If \mathbf{x} is required we can solve (independently)

$$\frac{\partial \mathbf{x}}{\partial t} = \mathbf{v} = \frac{1}{\rho_0} \mathbf{p}$$

Conservative Form of the Governing Equations

$$\mathcal{U}_t + \nabla \cdot \mathcal{F}(\mathcal{U}) = \mathcal{S}(\mathcal{U})$$

$$\mathcal{U} = \begin{pmatrix} p_1 \\ p_2 \\ p_3 \\ F_{11} \\ F_{12} \\ F_{13} \\ F_{21} \\ F_{22} \\ F_{23} \\ F_{31} \\ F_{32} \\ F_{33} \end{pmatrix}; \quad \mathcal{F}_I = \begin{pmatrix} -P_{1I}(\mathbf{F}) \\ -P_{2I}(\mathbf{F}) \\ -P_{3I}(\mathbf{F}) \\ -\delta_{I1}p_1/\rho_0 \\ -\delta_{I2}p_1/\rho_0 \\ -\delta_{I3}p_1/\rho_0 \\ -\delta_{I1}p_2/\rho_0 \\ -\delta_{I2}p_2/\rho_0 \\ -\delta_{I3}p_2/\rho_0 \\ -\delta_{I1}p_3/\rho_0 \\ -\delta_{I2}p_3/\rho_0 \\ -\delta_{I3}p_3/\rho_0 \end{pmatrix} \quad I = 1, 2, 3; \quad \mathcal{S} = \begin{pmatrix} \rho_0 b_1 \\ \rho_0 b_2 \\ \rho_0 b_3 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \end{pmatrix}$$

Eigenvalue Structure

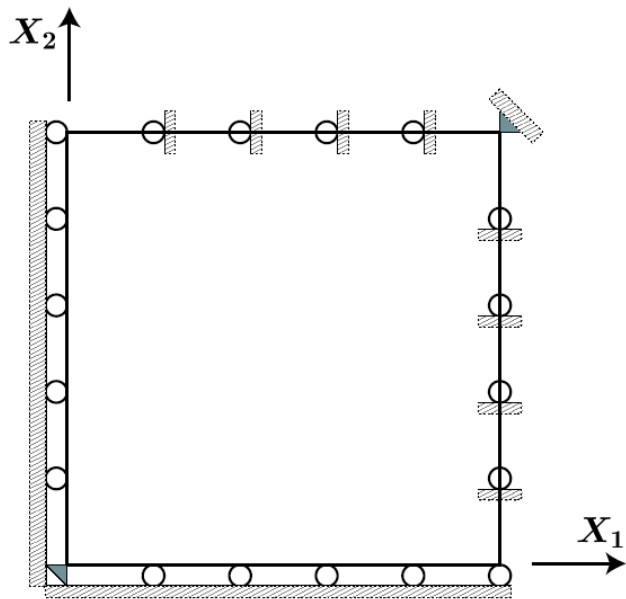
$$\mathcal{A}_N = \frac{\partial \mathcal{F}_N}{\partial \mathcal{U}} = N_I \mathcal{A}_I, \quad \mathcal{A}_I = \frac{\partial \mathcal{F}_I}{\partial \mathcal{U}} \quad I = 1, 2, 3$$

- Two Acoustic Waves
- Four Shear Waves
- Six zero eigenmodes ($\delta \mathbf{F}^T = -\delta \mathbf{F} +$ “Plane Stress States”)

$$\lambda_{max} = \sqrt{\frac{\beta + \left(\frac{\alpha}{\varepsilon^2} + 2\gamma\right)}{\rho_0}},$$

$$\alpha = \kappa J^2 + \frac{5}{9} \mu J^{-2/3} (\mathbf{F} : \mathbf{F}), \quad \beta = \mu J^{-2/3}, \quad \gamma = -\frac{2\mu}{3} J^{-2/3}$$

Example - Plane Strain



Unit Square

Linear Solution

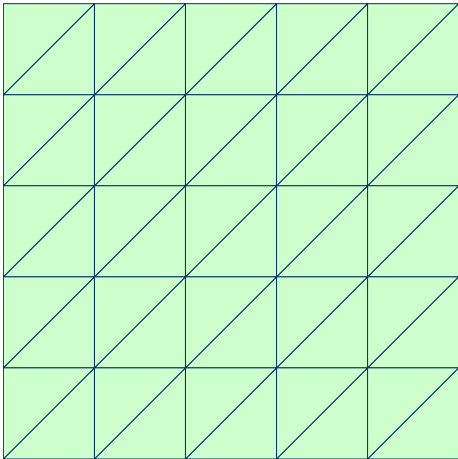
$$\mathbf{u}(t) = U_0 \cos\left(\frac{c_d \pi t}{\sqrt{2}}\right) \begin{bmatrix} \sin\left(\frac{\pi X_1}{2}\right) \cos\left(\frac{\pi X_2}{2}\right) \\ -\cos\left(\frac{\pi X_1}{2}\right) \sin\left(\frac{\pi X_2}{2}\right) \end{bmatrix} \cdot$$

$$c_d = \sqrt{\frac{\mu}{\rho}}$$

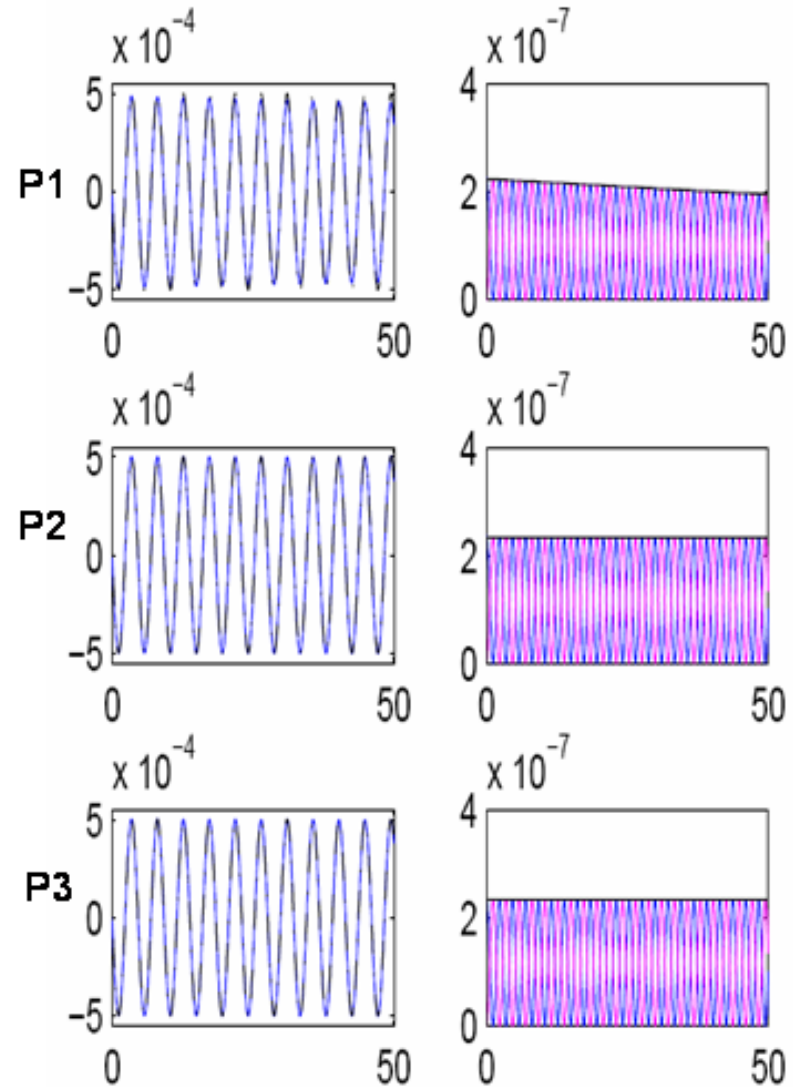
Example - Plane Strain Solution

$$U_0 = 5e - 4 \text{ (linear)}$$

$(5 \times 5 \times 2)$ mesh

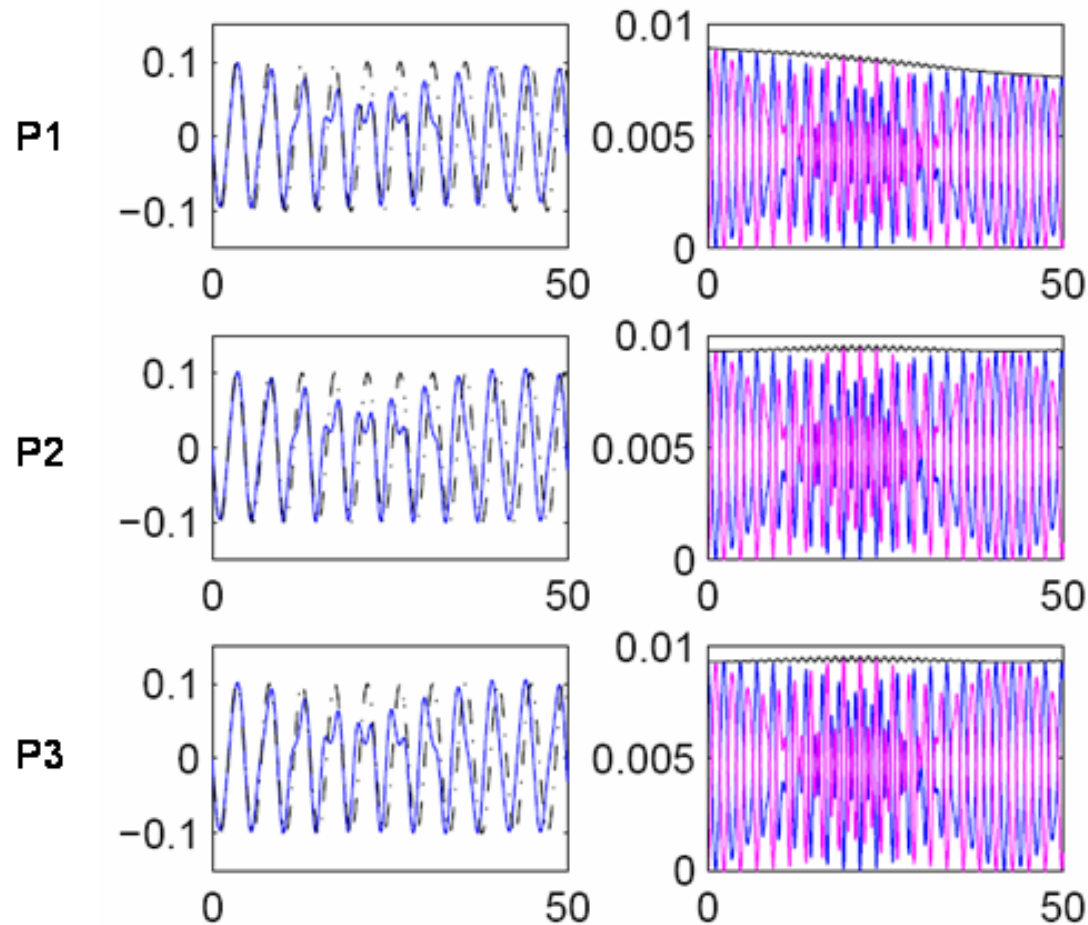


Displacement
(Bottom Right Corner)
and
Energy



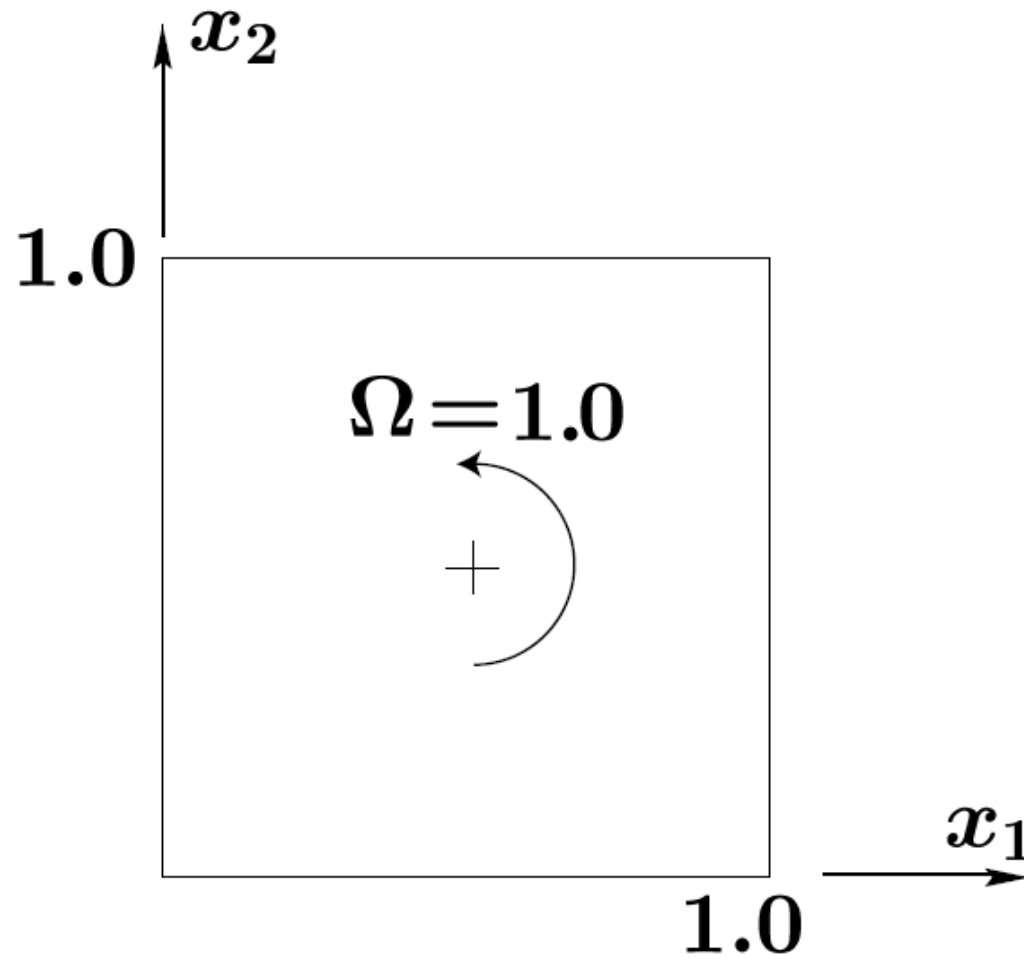
Example - Plane Strain Solution

$U_0 = 0.1$ (non-linear), $(5 \times 5 \times 2)$ mesh



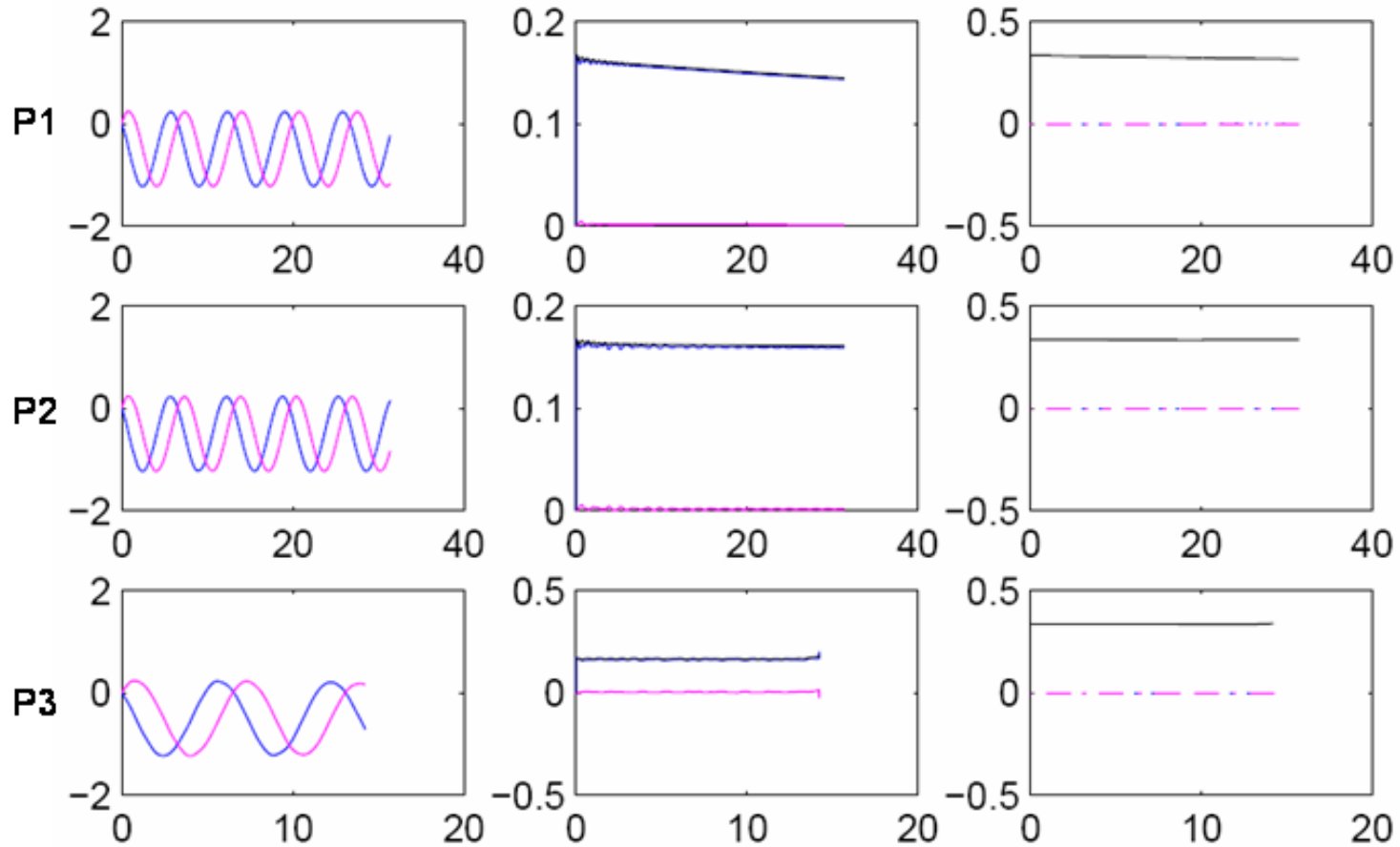
Displacement
(Bottom Right Corner)
and
Energy

Example - Rotating Plate



“Plane Strain” Rotating Plate

Example - Rotating Plate



Displacement of point (1,1), Energy, Linear and Angular Momentum

The Problem

We solve for the evolution of \mathbf{F}

$$\frac{\partial \mathbf{F}}{\partial t} = \nabla \cdot \left(\frac{1}{\rho_0} \mathbf{p} \otimes \mathbf{I} \right)$$

but \mathbf{F} is a gradient

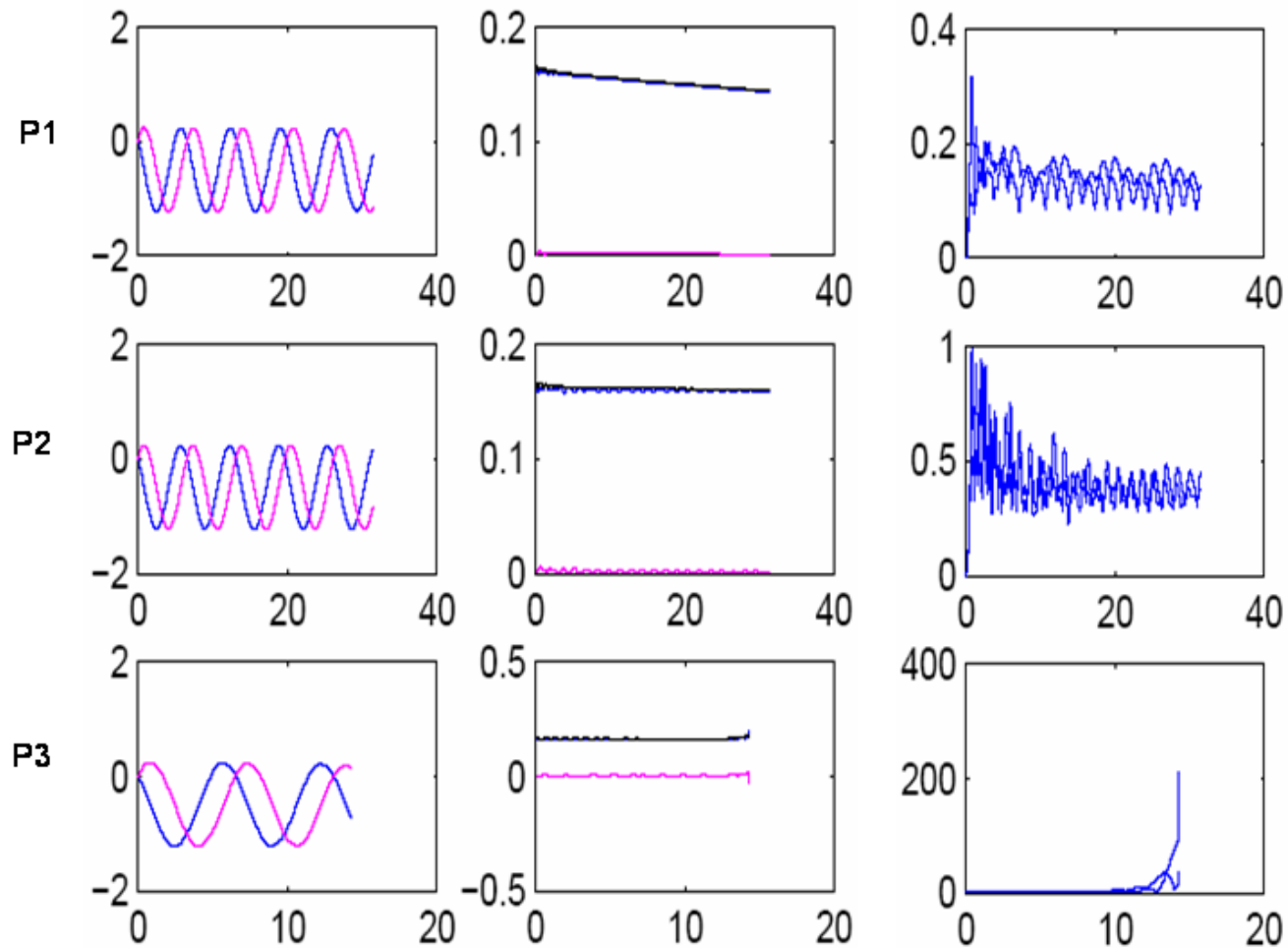
$$\mathbf{F} = \frac{\partial \mathbf{x}}{\partial \mathbf{X}} = \begin{pmatrix} x_{1,1} & x_{1,2} & x_{1,3} \\ x_{2,1} & x_{2,2} & x_{2,3} \\ x_{3,1} & x_{3,2} & x_{3,3} \end{pmatrix} = \begin{pmatrix} \mathbf{F}^1 \\ \mathbf{F}^2 \\ \mathbf{F}^3 \end{pmatrix}$$

Therefore, we have

$$\nabla \times \mathbf{F}^i = \mathbf{0}, \quad i = 1, 2, 3 \quad \forall t$$

... this is called an **INVOLUTION**

Example - Rotating Plate



Displacement of point (1,1), Energy, $F_{11,2} - F_{12,1}$ and $F_{21,2} - F_{22,1}$

Dealing with the Involutions ($\nabla \times \mathbf{F}^i = \mathbf{0}$)

- Calculate a curl-free vector basis $\{\phi_k\}$ within each element (SVD)

$$\nabla \times \phi_k = \mathbf{0}, \quad \forall k$$

- Approximate each row of \mathbf{F} using the curl-free basis

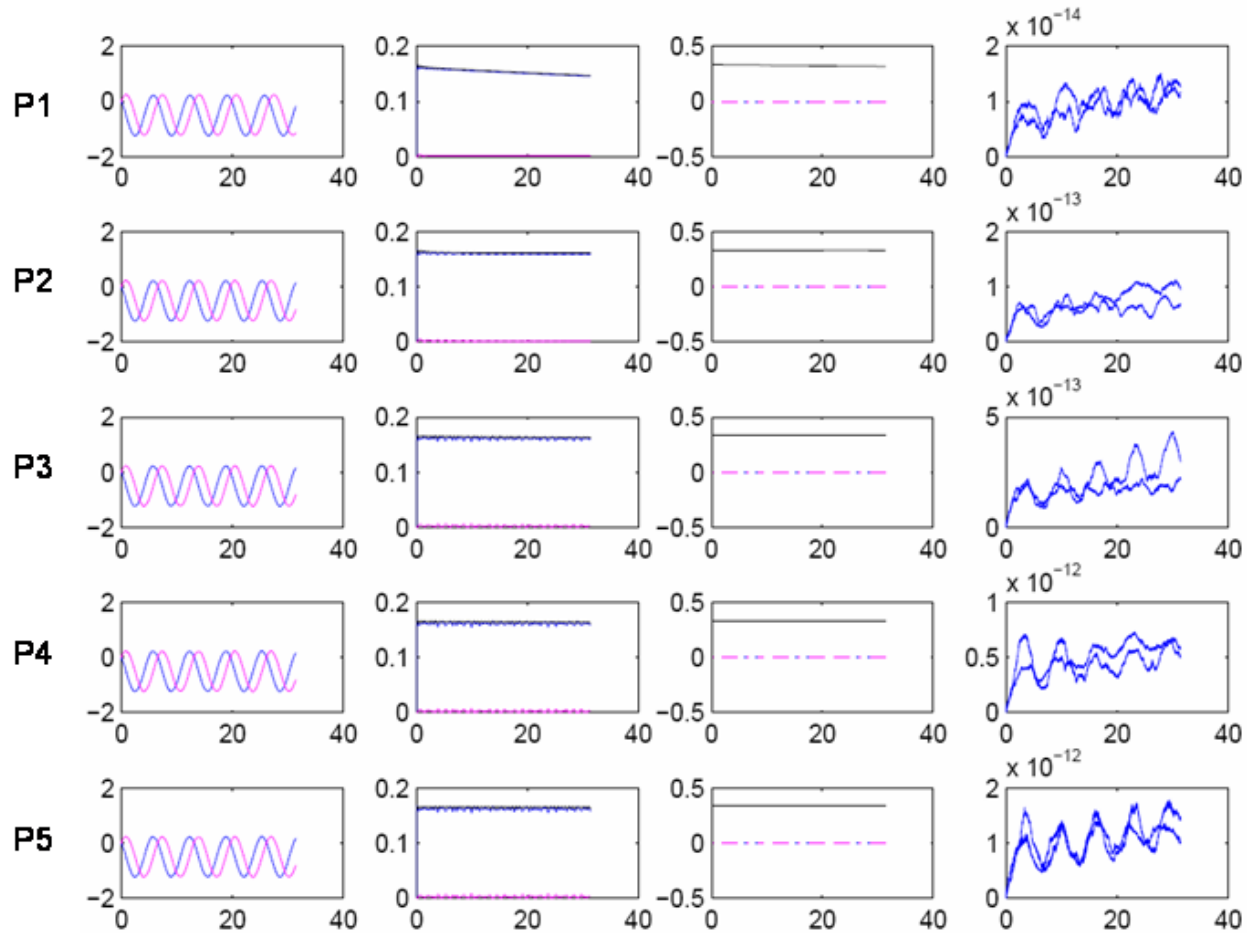
$$\mathbf{F}^i = \sum_k \mathbf{F}_k^i \phi_k$$

- Use standard DG approach with modified basis functions
- In practice, it can be implemented very efficiently by pre-calculating a modified inverse mass matrix \mathbf{M}_*^{-1} for each element

$$\mathbf{M}^{-1} \Rightarrow \mathbf{M}_*^{-1} = \mathbf{W}(\mathbf{W}^T \mathbf{M} \mathbf{W})^{-1} \mathbf{W}^T$$

\mathbf{W} is a rectangular matrix which expresses the curl-free basis in terms of the standard basis

Example - Rotating Plate

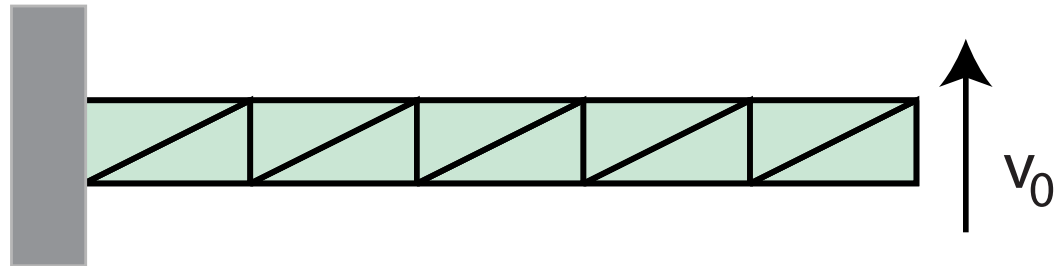


Displacement, Energy, Momentum, $\nabla \times \mathbf{F}^i$

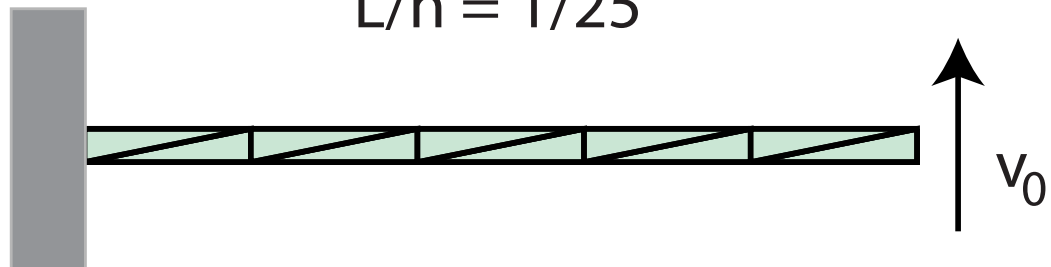
Bending Beam

Initial velocity given by first eigenmode

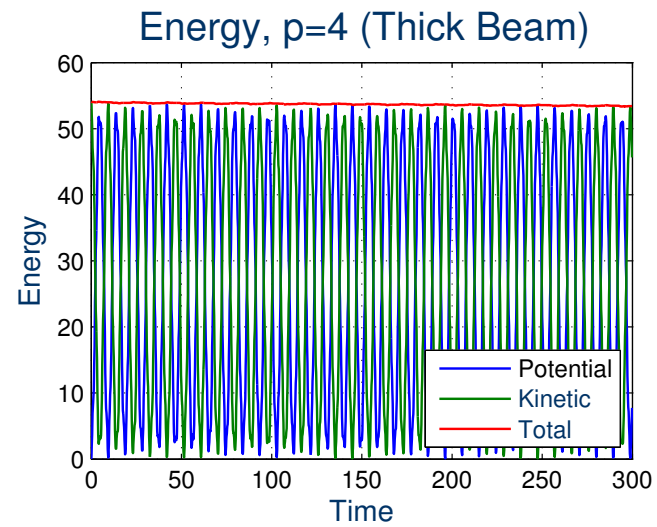
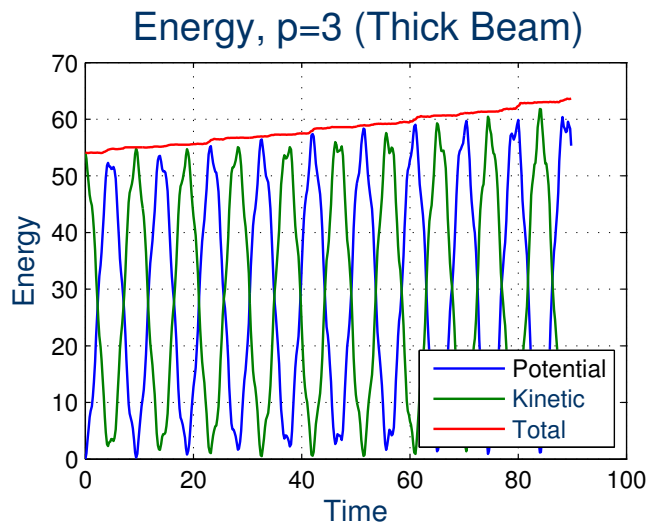
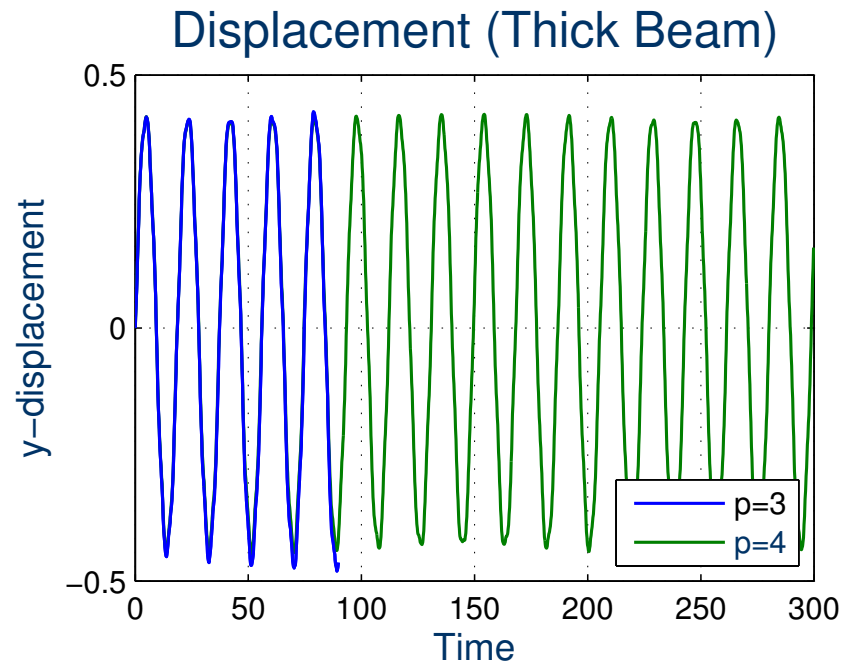
$$L/h = 1/10$$



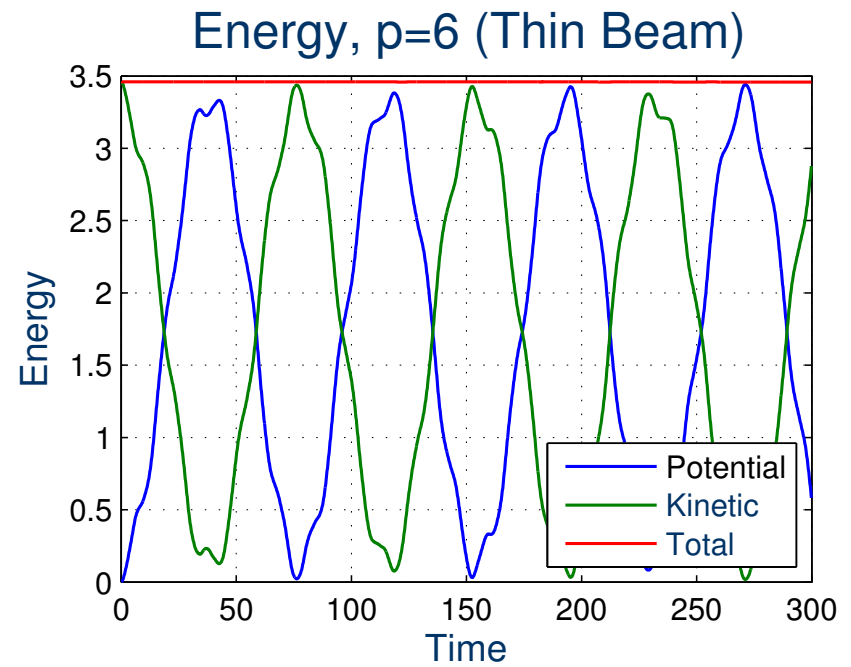
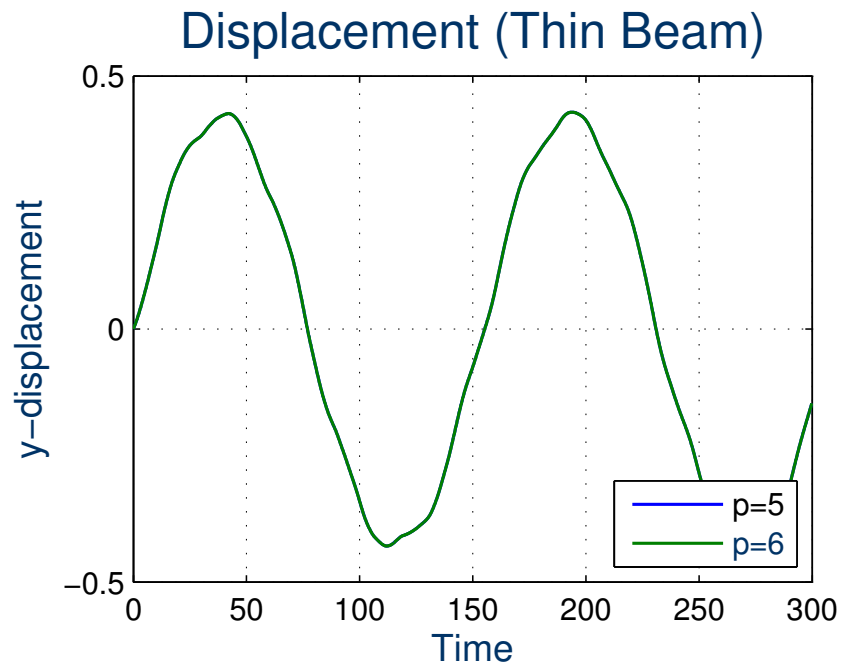
$$L/h = 1/25$$



Thick Beam (1/10)



Thin Beam (1/25)



Current/Future Work

- More General Constitutive Relations $\psi(\mathbf{F}, \theta)$ e.g. Mie-Gruneisen
- Solve for Energy equation

$$\frac{\partial E}{\partial t} + \nabla \cdot \left(\mathbf{q} - \frac{1}{\rho_0} \mathbf{P}^T \mathbf{p} \right) = 0$$

...plus additional constitutive internal variables

- Dissipation models (LDG, CDG)

$$\begin{aligned} \mathcal{U}_t + \nabla \cdot \mathcal{F}(\mathcal{U}, \mathcal{Q}) &= \mathcal{S}(\mathcal{U}) \\ \mathcal{Q} - \nabla \mathcal{U} &= \mathbf{0} \end{aligned}$$

- Shock Capturing
- Arbitrary Lagrangian/Eulerian Description
- Fluid Structure Interaction