

# Optimisation of the Key SOA Parameters for Amplification and Switching

A. Abd El Aziz, W. P. Ng, *Member, IEEE*, Z. Ghassemlooy, *Senior Member, IEEE*, Moustafa H. Aly, *Member, OSA* and M. F. Chiang,

**Abstract**—This paper presents the effects of the input signal on the gain and carrier density response of a semiconductor optical amplifier (SOA). The SOA is modelled using segmentation method. The optimum bias current required by the SOA for amplification and switching functions are investigated. The operation principle is simulated and the results show the input boundary conditions and requirements in which the SOA can be used as an amplifier and a switch.

**Index Terms**—semiconductor optical amplifier (SOA), carrier density, stimulated emission, gain response.

## I. INTRODUCTION

NOWADAYS, ultrafast photonic networks rely on photonic signal processing to overcome the speed bottleneck imposed by optoelectronic conversions. Optical amplifier is not limited only to amplify optical signals; it is also used as an important element in all-optical switching, regeneration and wavelength conversion. Among all-optical switches, ultrafast all-optical switches based on SOA, such as Mach-Zehnder Interferometers (MZIs) [1] are the most promising candidates for the realization of all-optical switching and processing applications due to their feature small size, high stability, low switching energy, high integration potential and their fast and strong nonlinearity characteristics [2-4]. There are three major nonlinear effects based on SOA; (i) cross-gain modulation (XGM), (ii) cross-phase modulation (XPM) and (iii) four-wave mixing (FWM).

SOA is also a key component for cascaded optical fiber systems and optical gating [5, 6]. Optical gates are needed in most all-optical functions such as wavelength conversion, add-drop multiplexing (wavelength and time), clock recovery, regeneration and simple bit-pattern recognition [7].

Moreover, the use of SOAs as in-line amplifiers is very suitable for bi-directional transmission in local and metropolitan systems and networks because of the lower cost of SOAs and they do not need any optical isolators as often used in different types of amplifiers such as erbium doped fiber amplifiers (EDFAs) [6].

In this paper, optimum parameters required for the SOA are simulated to perform amplification and switching functions. The key optimisations are achieved by controlling the bias current based on the input power to the SOA. The effect of the optimization on the carrier density and the gain responses are investigated. The principle of operation of the SOA is shown in the following section while section III presents the

mathematical analysis of the total gain and the change of the carrier density in terms of the rate equations. In section IV the boundary conditions and requirements for the SOA to perform amplification and switching are presented. The final section concludes the findings of the investigation.

## II. SOA PRINCIPLE OF OPERATION

The basic structure of the SOA is a semiconductor laser composed of an optical waveguide between a P-N junction with an input and output coated facets as displayed in Fig. 1 [8]. When a direct current (DC) is biased to the SOA, electrons get higher energy. Hence, the conduction and valence bands (energy levels) containing electrons and holes, respectively are formed [9].

The formation of these energy levels results in three radiative mechanisms within the SOA:

i. Spontaneous emission; which is the process where an electron from the conduction band drops to the valence band releasing a photon or generating heat. This process is considered as loss or noise because the generated photon is radiated with different phase and direction.

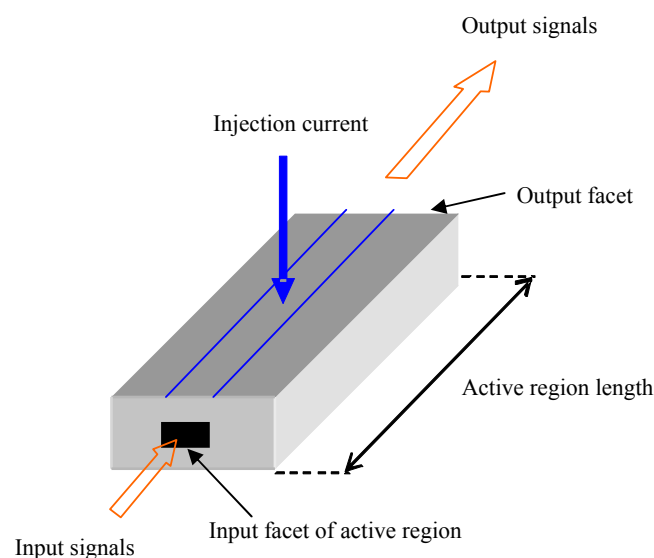


Fig. 1. Schematic diagram of the SOA.

ii. Stimulated absorption; this is the process by which an electron in the valence band absorbs enough energy from an incident photon to pass the energy gap to the conduction band and hence causes loss.

iii. Stimulated emission; this process occurs when an incoming optical beam is launched into the active waveguide of the SOA via the input facet of the amplifier, an incident photon collides with an excited electron from the conduction band releasing a stimulated photon with the same phase, frequency and direction (i.e. amplification takes place). More identical photons are released by the collision of the incident beam of photons with more excited electrons in the conduction band thus amplifying the input signal [9].

The reduction of excited electrons in the conduction band (i.e. the carrier density) will result in a decrease in the SOA gain because the gain is proportional to the carrier population. Moreover, it will increase the active refractive index due to the nonlinear refractive index being dependent on the carrier density [10, 11].

When a short input optical pulse is launched into the SOA, stimulation emission will take place resulting in signal amplification. Therefore, the carrier density will be reduced and will cause a drop in the SOA gain. The carrier non-equilibrium is governed mainly by the spectral hole burning effect [12, 13]. The distribution recovers to equilibrium by carrier-carrier scattering. Instantaneous mechanisms such as two-photon absorption [14, 15] and the optical Kerr effects [16, 17] will then influence on the SOA response. After few picoseconds, a quasi-equilibrium distribution will occur due to the carrier temperature relaxation process and then the carrier density will be recovered [9].

### III. MATHEMATICAL MODEL

The rate equations in small segments in an SOA are iteratively calculated while taking the carrier density change and the SOA length in account [18].

#### Rate Equations

When light is injected into the SOA, changes occur in the carrier and photon densities within the active region of the SOA. These changes can be described using the rate equations. The gain medium of the amplifier is described by the material gain coefficient,  $g$  (per unit length) which is dependent on the carrier density  $N$  and is given by [9]:

$$g = \alpha_g (N - N_0), \quad (1)$$

where  $N_0$  is the carrier density at transparency point and  $\alpha_g$  is the differential gain parameter. The net gain coefficient  $g_T$  is defined by:

$$g_T = \Gamma \cdot g - \alpha_s, \quad (2)$$

where  $\alpha_s$  is the internal waveguide scattering loss, and  $\Gamma$  is the confinement factor which is the ratio between the cross-

sectional area of the active medium and the transverse area of the optical waveguide.

The total gain  $G$  of an optical wave experienced at the location  $z$  of an SOA can be calculated according to

$$G = e^{g_T \cdot z}, \quad (3)$$

assuming a constant carrier density at any given location  $z$  within the active region of the SOA.

Therefore, the average output power  $P_{av}$  over the length of the SOA becomes:

$$P_{av} = \frac{1}{L} \int_0^L P_{in} G dz, \quad (4)$$

where  $L$  is the length of the SOA and  $P_{in}$  is the input signal power. The average output power can be rewritten as [19]:

$$P_{av} = P_{in} \frac{e^{g_T \cdot L} - 1}{g_T \cdot L}. \quad (5)$$

The dynamic equation for the change in the carrier density within the active region of the device is given by:

$$\frac{dN}{dt} = \frac{I_{DC}}{q \cdot V} - R(N) - \frac{\Gamma \cdot g \cdot P_{av} \cdot L}{V \cdot h \cdot f}, \quad (6)$$

where  $I_{DC}$  is the DC current injected to the SOA,  $q$  is the electron charge and  $V$  is the active volume of the SOA,

$$V = L \cdot W \cdot H, \quad (7)$$

where  $W$  and  $H$  denote the width and the thickness of the active region, respectively.  $R(N)$  is the recombination rate,  $h$  is the Planck constant and  $f$  is the light frequency.

There are two mathematical definitions for the recombination rate. The simple definition is given by:

$$R(N) = \frac{N}{T_{sp}}, \quad (8)$$

where  $T_{sp}$  is the spontaneous emission time and is given by:

$$T_{sp} = \frac{q \cdot V \cdot N_0}{I_{DC}}, \quad (9)$$

while the more complex form is called the Auger recombination and is given by:

$$R(N) = A \cdot N + B \cdot N^2 + C \cdot N^3, \quad (10)$$

where  $A$  is the surface and defect recombination coefficient,  $B$  is the radiative recombination coefficient and  $C$  is the Auger recombination coefficient. In this paper, the Auger recombination equation shown in (10) is used in developing the theoretical SOA model.

#### IV. RESULTS AND DISCUSSION

The rate equations shown in section III are carried out via Matlab™ to investigate the gain response of the SOA model while employing the segmentation method. The standard SOA parameters used are given in Table I.

The segmentation method involves dividing the SOA into five equally segments of length  $l=L/5$  each. The carrier density is assumed to be constant within a segment. However, the carrier density changes from one segment to another depending on its input power and the carrier density of the previous segment using equation (6). In all the equations in section III, the segment length  $l$  will replace the SOA length  $L$  for the segment total gain and carrier density calculations.

Figure 2 shows the normalized gain response of the SOA using the physical parameters in Table I. As it can be seen, the gain of the SOA increases rapidly until a steady state value where it becomes constant. This increase is due to is the injection of the bias current to the SOA hence resulting in a large number of electrons to overcome the energy gap to reach the conduction band. Therefore, the carrier density of the SOA will increase and hence the total gain of the SOA.

TABLE I  
Physical parameters of the SOA

<i>parameter</i>	<i>symbol</i>	<i>value/unit</i>
Carrier density at transparency	$N_0$	$1.4 \cdot 10^{24}/\text{m}^3$
Initial carrier density	$N_i$	$3 \cdot 10^{24}/\text{m}^3$
Differential gain	$a_g$	$2.78 \cdot 10^{-20}/\text{m}^2$
Internal waveguide scattering loss	$a_s$	$40 \cdot 10^2/\text{m}$
SOA length	$L$	$500 \mu\text{m}$
SOA width	$W$	$3 \mu\text{m}$
SOA height	$H$	$80 \text{nm}$
Confinement factor [20]	$\Gamma$	$0.3$
Light frequency	$f$	$193.1 \text{ THz}$
Plank constant	$h$	$6.62606896 \cdot 10^{-34}$
Electron charge	$q$	$1.602 \cdot 10^{-19} \text{ C}$
Surface and defect recombination coefficient	$A$	$1.43 \cdot 10^8 \text{ 1/s}$
radiative recombination coefficient	$B$	$1 \cdot 10^{-16} \text{ m}^3/\text{s}$
Auger recombination coefficient	$C$	$3 \cdot 10^{-41} \text{ m}^6/\text{s}$

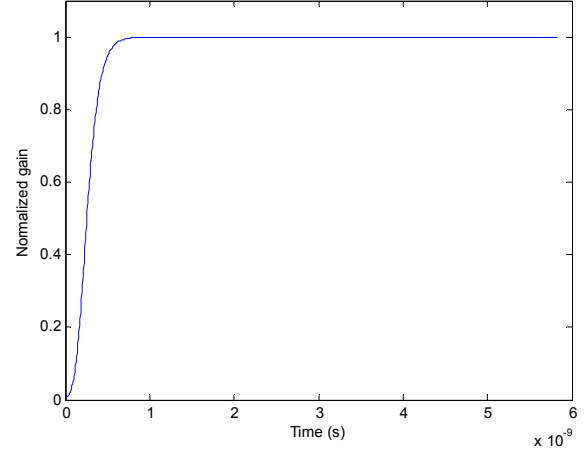


Fig. 2. Normalized gain response of the SOA.

In this simulation, two types of input signals are applied to the SOA separately for investigation. The first input is a short pulse input signal with a pulse width of  $l/v_g$  (i.e. 1.1667 ps), where  $v_g$  is the group velocity. As explained previously in section II, when the input pulse is injected to the active region of the waveguide, the total gain of the SOA will drop instantly and then will recover back to its steady state value as shown in Fig. 3. The drop of the total gain depends on the power and pulse width of the input pulse as will be discussed at the end of this section.

The second input signal is a continuous wave (CW) signal which will be launched into the SOA after the total gain reaches its steady state value. When a CW is injected into the SOA, the SOA gain will reduce until it reaches a saturation value where the gain becomes constant as shown in Fig. 4. The reason for such response is that a depletion of the carrier density happens due to the continuous stimulation emission process. The carrier density continues to decrease and hence the gain until excited electrons in the conduction band are no longer available. By using the segmentation method, the change of the carrier density along the SOA length after applying the CW can be investigated, as depicted in Fig. 5. Carriers are depleted and therefore gain is suppressed.

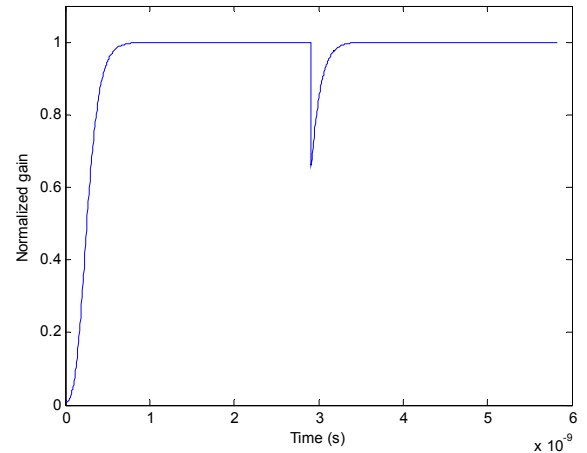


Fig. 3. Normalized gain response of the SOA due to the injection of a short input pulse.

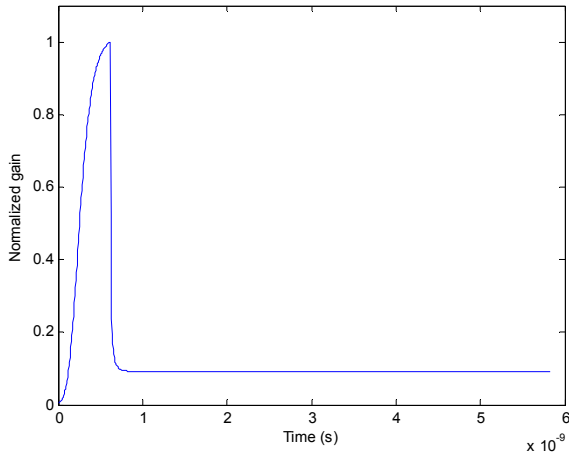


Fig. 4. Normalized gain response of the SOA due to the injection of a continuous input signal.

In order to use the SOA as an amplifier, we should ensure that the signal will not be affected by the SOA nonlinear response which occurs when the gain is saturated. This amplification can only be achieved if the SOA gain does not reach its saturation value that is shown in Fig. 4 when applying an input signal pulse. In this paper, the saturation value for a 1mW continuous input signal is used as the reference saturation gain value in which the SOA will achieve a linear response for amplification function.

From Fig. 6 for a given input signal power, one can choose the optimum bias current in order to achieve the required output gain. As expected, the increase of the input signal power results in the decrease of the output gain to a constant minimum gain. For the different values of the bias current, the minimum gain is different. Input signal of 1mW with  $I = 250\text{mA}$  achieves a high gain compared to signals with higher input signal power and to signals with lower bias current. This is because higher bias current gives more electrons enough energy to reach the conduction band and hence higher carrier density and higher output gain.

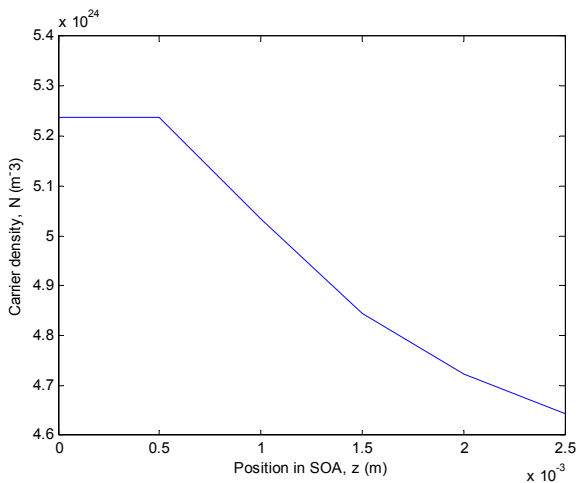


Fig. 5. Carrier density along the SOA length after applying a CW.

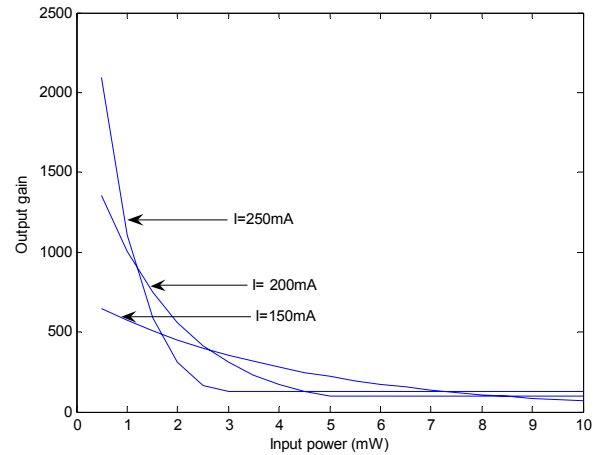


Fig. 6. The output gain corresponding to the input power at different bias currents.

Figure 6 also shows that the rate of decrease of the output gain is the highest in the case of  $I = 250\text{mA}$  while it is the lowest in the case of  $I = 150\text{mA}$ . The reference saturation gain of the SOA at 1mW continuous input signal is 66, 96 and 127 for bias currents of 150mA, 200mA and 250mA, respectively.

On the contrary, in order to use the XPM introduced by SOA for switching function [21], ideally the signal should be affected by the nonlinearity of the SOA and achieve a  $180^\circ$  phase shift for the destructive interference [21]. This can be obtained by the aid of a control pulse (CP) which is injected to the active region of the SOA in order to achieve a gain depletion that reaches the gain saturation value. When the total gain reduction reaches this value, a phase shift of  $180^\circ$  will occur [22] and then the input signal should be launched into the SOA in order to achieve this phase shift.

For any input signal to achieve the nonlinear response of the SOA, a corresponding CP is required to drive the total gain of the SOA into the SOA gain saturation. Figure 7 illustrates the corresponding CP needed by a range of input signal powers to obtain the desired phase shift for different injection currents.

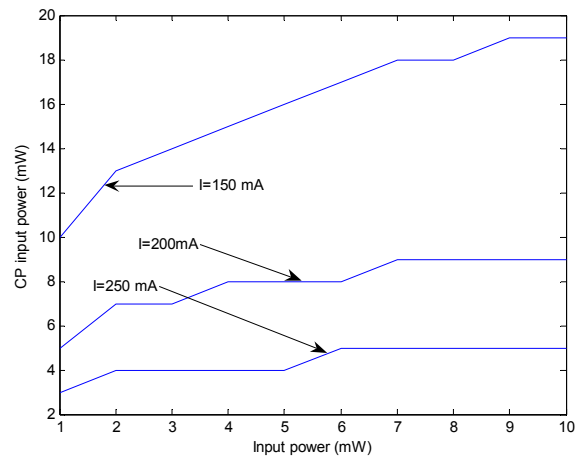


Fig. 7. The saturation CP for corresponding input power at different bias currents.

In Fig. 7, some input signal power would require the same CP power to drive the SOA into saturation, i.e. input power of 6 to 10 mW required CP of 5 mW at  $I = 250$  mA. In addition, the CP power required is higher for  $I = 150$  mA compared to higher bias currents while for  $I = 250$  mA a lower control pulse power is required to drive the SOA into saturation and hence perform switching function [22]. The highest control pulse power needed is of 19mW for an input signal power of 10mW at a bias current of 150mA. On the other hand, the lowest control pulse power needed is 3mW for an input of 3mW at a bias current of 250mA.

## V. CONCLUSION

This paper has simulated the total gain response of an SOA model using segmentation method. Continuous and short input signals were applied to the segmented SOA model to investigate the corresponding gain response in order to achieve amplification and all-optical switching operation respectively. The optimum performance conditions are investigated in this paper for the SOA to function as an amplifier and a switch.

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