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厦门大学

博士 学位 论文

分数阶热流方程, 积分方程以及向列型液晶模型解的存在性及其相关问题

**Existences to Fractional Heat Flow, Integral Equation
and Liquid Crystal Flow and Related Topics**

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目 录

中文摘要	V
英文摘要	VII
第一 章 绪论	1
1.1 分数阶热流方程	1
1.2 奇异积分方程组	4
1.3 向列型液晶模型的整体有限能量弱解性存在性	7
第二 章 基础知识	9
2.1 符号约定	9
2.2 预备引理	11
第三 章 分数阶热流方程	21
3.1 引言	21
3.2 \wedge^α 算子的性质与对应的线性系统	24
3.3 惩罚逼近系统	32
3.4 定理 3.1.1 的证明	41
第四 章 积分方程组	47
4.1 引言	47
4.2 一些记号和引理	51
4.3 定理 4.1.1 的证明	53
4.4 定理 4.1.2 的证明	53
第五 章 液晶模型	59
5.1 介绍与主要结果	59
5.2 预备性引理	61
5.3 逼近解	62
5.4 定理 5.1.1 的证明	68
参考文献	75
致谢	81

厦门大学博硕士论文摘要库

Contents

Chinese Abstract	V
English Abstract	VII
Chapter I Preface	1
1.1 Heat Flow Equation Involving Fractional Laplacian	1
1.2 Singular Integral Equations	4
1.3 Global Existence of the Finite Energy Weak Solutions to a Nematic Liquid Crystals Model	7
Chapter II Preliminary	9
2.1 Basic Notation	9
2.2 Preliminary Lemma	11
Chapter III Generalized Heat Flow Equation Involving Fractional Laplacian	21
3.1 Introduction and Main Results	21
3.2 Properties of Operators \wedge^α and Linear System	24
3.3 Penalized Approximate System	32
3.4 Proof of Theorem 3.1.1	41
Chapter IV Integral Equations	47
4.1 Introduction	47
4.2 Some Notations and Lemmas	51
4.3 Proof of Theorem 4.1.1	53
4.4 Proof of Theorem 4.1.2	53
Chapter V Global Existence of the Finite Energy Weak Solutions to a Nematic Liquid Crystals Model	59
5.1 Introduction and Main results	59
5.2 Some Preliminary Lemma	61
5.3 Approximate Solution	62
5.4 Proof of Theorem 5.1.1	68
Chapter VI References	75

Acknowledgements	81
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厦门大学博硕士论文摘要库

中文摘要

本文讨论了三个偏微分方程组解的存在性以及解的相关性质。全文共分为三个部分。第一部分为本文的第三章，主要研究一个广义的热流方程，即：

$$\begin{cases} \frac{\partial \mathbf{u}}{\partial t} = -\wedge^{2\alpha} \mathbf{u} + (\wedge^{2\alpha} \mathbf{u}, \mathbf{u}) \mathbf{u}, \\ \mathbf{u}(x, 0) = \mathbf{u}_0(x), \\ |\mathbf{u}_0(x)| = 1. \end{cases}$$

这里 $\mathbf{u} = (u_1, u_2, u_3)$ 为定义在 $\Omega \times [0, +\infty)$ 上的一个有限旋转向量，即 $|\mathbf{u}| = 1$ 。 $\Omega \subset \mathbb{R}^3$ 是周期域或全空间。我们也用 \wedge 表示 $(-\Delta)$ 的平方根算子， $0 < \alpha \leq 1$ 。 (\cdot, \cdot) 表示通常的向量内积。

当 $\alpha = 1$ 时，这系统变为了一个从 $\Omega \subset R^3$ 到 S^2 上的传输调和热流映射。实际上上面的传输调和热流方程可以推广到更一般的形式，即从紧致 Riemannian 流型 M (m 维) 到紧致 Riemannian 流型 $N = S^n$ 上的一个传输热流映射^[32,33]。为了区别这种经典的传输调和映射，我们把上面这个偏微分方程称为广义热流方程或分数阶热流方程。我们证明了：如果初值 $u_0 \in H^\alpha(\mathbb{I}^n) \cap L^p(\mathbb{I}^n)$, $p > \frac{n}{\alpha}$ 则分数阶热流方程至少存在一个整体弱解。

本章的主要贡献在于我们推广了文献 [47] 中阶为 $\frac{1}{2}$ 的传输热流模型。下面我们简单介绍一下我们处理这个系统所遇到的困难和所用的技巧。

对于 $\alpha = \frac{1}{2}$ 和 $\alpha = 1$ 而言，此时 \wedge 和 $-\Delta$ 分别是 Cauchy 半群和 Gaussian 半群的生成元，其半群的 Fourier 变化有具体的显示，因而我们容易得到其对应的惩罚方程解的 L^∞ 估计。但是对于一般的热流方程而言， $-(-\Delta)^\alpha$ 是对称 α -稳定过程的生成元^[7]，其半群的 Fourier 变化不再具有显示，这里我们利用 Fourier 变化的在周期域上的性质得到了对应逼近方程的 L^∞ 估计，并证明了惩罚逼近方程的整体存在性。

第二个困难来源于非线性项的收敛问题。由于逼近解的正则性不高以及分数阶算子不满足莱布尼兹法则即 $\wedge^\alpha(f \cdot g) \neq g \cdot \wedge^\alpha f + f \cdot \wedge^\alpha g$ ，所以我们不能像经典调和热流方程那样直接利用逼近解的某种紧性去证明分数阶热流方程解的存在性。同时我们也注意到对于 $\frac{1}{2} < \alpha < 1$ ，文献 [47] 中的方法也是失效的。这里我们利用著名的交换子引理和方程自身的结构来克服这个困难。

第二部分是本文的第四章，在这一章中，我们主要的目标是研究下列奇异积分方程

组正解的存在性和不存在性。

$$\begin{cases} u(x) = \int_{\mathbb{R}^n} \frac{v^q(y)}{|x-y|^{n-\alpha}} dy, \\ v(x) = \int_{\mathbb{R}^n} \frac{u^p(y)}{|x-y|^{n-\mu}} dy. \end{cases}$$

我们证明了：当 $p \leq \frac{\mu}{n-\alpha}$ ($q \leq \frac{\alpha}{n-\mu}$)，或 $q = 1, 0 < p < \frac{n}{n-(\alpha+\mu)}$ ($p = 1, 0 < q < \frac{n}{n-(\alpha+\mu)}$) 时，上面的积分系统不存在正解。相似的，当 $q = 1, \frac{n}{n-(\alpha+\mu)} \leq p < \frac{n+(\alpha+\mu)}{n-(\alpha+\mu)}$ ($p = 1, \frac{n}{n-(\alpha+\mu)} \leq q < \frac{n+(\alpha+\mu)}{n-(\alpha+\mu)}$) 上面的积分系统在 $u \in L_{loc}^{\frac{n(p-1)}{(\alpha+\mu)}}(R^n)$ ($v \in L_{loc}^{\frac{n(q-1)}{(\alpha+\mu)}}(R^n)$) 中不存在正解。并且，利用 W.X.Chen, C.M.Li 和 B.Ou^[79,80,82,84] 所建立的积分型的移动平面法，我们也证明了如果 $1 < p, q < \infty$, $p_1 \geq 1$, $q_1 \geq 1$ 满足下列不等式

$$\frac{p-1}{p_1} + \frac{q-1}{q_1} \geq \frac{\alpha+\mu}{n},$$

$$1 \geq \frac{p}{p_1} > \frac{\mu}{n}, 1 \geq \frac{q}{q_1} > \frac{\alpha}{n}.$$

且 $u \in L^{p_1}(\mathbb{R}^n) \cap L^\infty(\mathbb{R}^n)$, $v \in L^{q_1}(\mathbb{R}^n) \cap L^\infty(\mathbb{R}^n)$ 是系统 (1.1) 一个正解，则 (u, v) 一定是关于 \mathbb{R}^n 中的某一点是径向对称且单调递减的。

第三部分为本文的第五章，在这章中我们主要讨论了下列向列型液晶模型。

$$\begin{cases} \rho_t + \operatorname{div}(\rho \mathbf{u}) = 0, \\ (\rho \mathbf{u})_t + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u}) + \nabla P = \mu \Delta \mathbf{u} - \lambda \nabla \cdot (\nabla \mathbf{d} \odot \nabla \mathbf{d}), \\ \mathbf{d}_t + \mathbf{u} \cdot \nabla \mathbf{d} = \gamma(\Delta \mathbf{d} - \mathbf{f}(\mathbf{d})), \\ \nabla \cdot \mathbf{u} = 0. \end{cases}$$

上面的方程是定义在区间 $Q_T := \Omega \times (0, T)$ 上，其中 $\Omega \subset \mathbb{R}^3$ 是有界正则区域， $\mathbf{u}(x, t), \mathbf{d}(x, t)$ 分别表示流体的速度与液晶分子的方向， $\rho(x, t), P(x, t)$ 分别表示流体密度和压强。 \mathbf{f} 是一个光滑函数。

在初始密度属于 $L^r(\Omega)$ $r > \frac{3}{2}$ 且 \mathbf{f} 为惩罚函数的条件下，我们证明了向列型液晶模型存在整体能量弱解。对比于 Jiang-Tan 在 [2] 中文献的结论，我们利用逼近解的能量泛函以及标准的椭圆正则性估计去掉了系统对于 \mathbf{f} 光滑有界的要求。

关键词：惩罚函数；弱解；向列型液晶模型；积分方程组；对称 α - 稳态过程。

Abstract

This dissertation which consists of three parts, is devoted to studying the existence of three partial differential equations and related properties. Part 1, which is the third chapter in this paper, studies the existence of generalized heat flow equation involving fractional Laplacian. i.e.,

$$\begin{cases} \frac{\partial \mathbf{u}}{\partial t} = -\wedge^{2\alpha} \mathbf{u} + (\wedge^{2\alpha} \mathbf{u}, \mathbf{u}) \mathbf{u} \\ \mathbf{u}(x, 0) = \mathbf{u}_0(x) \\ |\mathbf{u}_0(x)| = 1. \end{cases}$$

in which $\mathbf{u} = (u_1, u_2, \dots, u_n) : \Omega \times [0, +\infty) \rightarrow \mathbb{R}^n$ is a finite spin vector, i.e., $|\mathbf{u}| = 1$. $\Omega \subset \mathbb{R}^n$ is a period region or the whole space, \wedge denotes the square root of $(-\Delta)$ and $0 < \alpha \leq 1$. (\cdot, \cdot) represents the normal inner product.

As $\alpha = 1$, this system turns into the classical evolution problems of Harmonic maps from $\Omega \subset R^3$ into S^2 , which is extended to more general harmonical map between different compact Riemannian manifolds and has been studied extensively in [32,33]. To distinguish the classical evolution heat flow, the above system is called a generalized evolution heat flow or evolution heat flow equation involving fractional Laplacian. We show that there exist at least a weak solution to the generalized evolution heat flow, provided $u_0 \in H^\alpha(\mathbb{I}^n) \cap L^p(\mathbb{I}^n)$ with $p > \frac{n}{\alpha}$ and $\mathbb{I}^n = \mathbb{R}^n / \mathbb{Z}^n$.

The main contribution of this chapter is that we extend the evolution heat flow with $\alpha = \frac{1}{2}$ studied by Z.A.Yao in [47], Main difficulties and techniques involved in studying our problem are as follows:

Firstly, when $\alpha = \frac{1}{2}$ or $\alpha = 1$, \wedge and $-\Delta$ are respectively the infinitesimal generator of the Cauchy semigroup and Gaussian semigroup in R^n , the Fourier transform of which has concrete expression, therefore we easily obtain the L^∞ -estimation of the corresponding approximated system of fractional heat flow equation. However, for general symmetric α -stable processes^[7], we can't obtain the concrete expression of the Fourier transform of corresponding semigroup. Here we make use of the properties of Fourier transform in period region to obtain the L^∞ -estimate of penalized equation.

Secondly, as the regularity of approximate solutions is not strong enough and the fractional operator \wedge^α doesn't meet the law of Leibniz ($\wedge^\alpha(f \cdot g) \neq g \cdot \wedge^\alpha f + f \cdot \wedge^\alpha g$),

this brings difficulties to the nonlinear term of the convergence. Therefore, we can't use directly the compactness of approximate solutions to show the existence of fractional Laplacian heat flow equation which is done by the classical evolution equation of harmonic map. On the other hand, we observe that the method used by Z.A.Yao in[47] is invalid to the case $\frac{1}{2} < \alpha < 1$. Here, with the help of the well-known commutator theorem and the feature of system, we overcome this difficulty.

The second part is the fourth chapter in this paper. Its main aim is to study the existence and non-existence of positive solution to a class of singular integral system

$$\begin{cases} u(x) = \int_{\mathbb{R}^n} \frac{v^q(y)}{|x-y|^{n-\alpha}} dy, \\ v(x) = \int_{\mathbb{R}^n} \frac{u^p(y)}{|x-y|^{n-\mu}} dy. \end{cases}$$

we prove that as $p \leq \frac{\mu}{n-\alpha}$ ($q \leq \frac{\alpha}{n-\mu}$), or $q = 1, 0 < p < \frac{n}{n-(\alpha+\mu)}$ ($p = 1, 0 < q < \frac{n}{n-(\alpha+\mu)}$) there doesn't exist positive solutions to the above integral system and as $q = 1, \frac{n}{n-(\alpha+\mu)} \leq p < \frac{n+(\alpha+\mu)}{n-(\alpha+\mu)}$ ($p = 1, \frac{n}{n-(\alpha+\mu)} \leq q < \frac{n+(\alpha+\mu)}{n-(\alpha+\mu)}$) the integral system does not have any positive solution in $u \in L_{loc}^{\frac{n(p-1)}{(\alpha+\mu)}}(R^n)$ ($v \in L_{loc}^{\frac{n(q-1)}{(\alpha+\mu)}}(R^n)$). Furthermore, With the help of moving plane increased by W.X.Chen, C.M.Li and B.Ou in [79,80,82,84], we also show that if $p_1, q_1 \geq 1$ and $1 < p, q < \infty$ satisfy

$$\frac{p-1}{p_1} + \frac{q-1}{q_1} \geq \frac{\alpha+\mu}{n},$$

$$1 \geq \frac{p}{p_1} > \frac{\mu}{n}, 1 \geq \frac{q}{q_1} > \frac{\alpha}{n}.$$

and $u \in L^{p_1}(\mathbb{R}^n) \cap L^\infty(\mathbb{R}^n)$, $v \in L^{q_1}(\mathbb{R}^n) \cap L^\infty(\mathbb{R}^n)$ are a pair of positive solutions of the singular integral system. Then (u, v) must be radially symmetric and monotone decreasing about some point in \mathbb{R}^n

The final part is the fifth chapter in this paper. we consider the following simplified system modeling the flows of liquid crystal materials

$$\begin{cases} \rho_t + \operatorname{div}(\rho \mathbf{u}) = 0, \\ (\rho \mathbf{u})_t + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u}) + \nabla P = \mu \Delta \mathbf{u} - \lambda \nabla \cdot (\nabla \mathbf{d} \odot \nabla \mathbf{d}), \\ \mathbf{d}_t + \mathbf{u} \cdot \nabla \mathbf{d} = \gamma(\Delta \mathbf{d} - \mathbf{f}(\mathbf{d})), \\ \nabla \cdot \mathbf{u} = 0. \end{cases}$$

in $Q_T := \Omega \times (0, T)$ with the bounded regular domain $\Omega \subset \mathbb{R}^3$. Here $\mathbf{u}(x, t), \mathbf{d}(x, t)$ denote the velocity of the fluid and the orientation of the liquid crystal molecules, respectively. $\rho(x, t)$ is the density of the fluid and $P(x, t)$ is the pressure. $\mathbf{f}(d)$ is a smooth vector function.

With $u_0 \in L^r(\Omega)$ $r > \frac{3}{2}$, we show the global existence of finite energy weak solutions to the liquid crystal system, as \mathbf{f} is the penalized function. In comparison with the result obtained by Jiang and Tan in [2], we remove the constrained condition that \mathbf{f} is a bounded and smooth function with the help of the energy function to approximation solution and standards elliptic regularity estimate.

Key words: penalized function; weak solution ; integral equations; symmetric α -stable process.

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