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Ultrasonic sensing using Yb³⁺/Er³⁺-codoped distributed feedback fibre grating lasers

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Abstract

Yb³⁺/Er³⁺-codoped distributed feedback (YbEr-DFB) fibre grating lasers have been investigated for the detection of ultrasonic Lamb waves generated in an aluminium plate. Two sensor configurations have been investigated. A path-length-imbalanced readout Mach-Zehnder interferometer was used to measure the laser frequency shift induced by ultrasonic vibration. The resolution to out-of-plane displacement was 3×10^{-15} m Hz^{-1/2} for a continuous ultrasonic wave at a frequency of 470 kHz. Ultrasonic waves were also detected by directly monitoring the laser output power. This simplified configuration is useful for tone burst AE signals generated by a low energy impact.

(Some figures in this article are in colour only in the electronic version)

1. Introduction

Structural impact or fracture events are associated with the generation of acoustic stress waves, typically in a frequency range between 100 kHz and a few MHz. Acoustic-emission (AE) detection has been established as one of the main non-destructive testing techniques for monitoring structural health and integrity [1-4]. For airframe and large civil structures in particular, the monitoring of impacts and damage using the AE detection method is important for safety Fibre Bragg grating (FBG) and cost reduction [1, 4]. strain sensors have considerable advantages over conventional piezoelectric sensors, such as immunity to electromagnetic interference, ease of installation and multiplexing capability. Different techniques employing FBGs to detect AE have been demonstrated [2-7]. A minimum detectable dynamic strain in the order of 10 nanostrain has been demonstrated for sensors using passive FBGs [3, 4]. Higher strain sensitivity at a few hundred hertz was obtained using a pi-phase shifted FBG which was interrogated with a tuneable laser [7]. The

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We have observed that the output power of a DFB fibre laser is sensitive to ultrasonic waves that impinge on the laser

maximum measurable AE frequency using this technique is limited by the frequency scanning speed of the tuneable laser.

advantages associated with FBG sensors, but also the potential

to measure AE signals with much higher sensitivities than FBG

sensors [8-12]. The linewidth of fibre lasers is typically a

few tens of kilohertz, corresponding to a coherence length of

the order of tens of kilometres. In this sensor configuration,

the AE signals perturb the laser cavity, thus modulating the

lasing frequency which can be measured using a path-length-

imbalanced readout Mach-Zehnder interferometer (MZI).

Since the laser linewidth is narrow, a long path length can be

used and high AE sensitivity can be realized. For example, an

axial strain sensitivity of 10^{-14} Hz^{-1/2} at 7 kHz using a single-

mode distributed Bragg reflector (DBR) Er³⁺-doped fibre laser

has been reported [8]. Distributed feedback (DFB) fibre

grating lasers provide more robust single-frequency operation

without mode-hopping [13], and are expected to offer better

performance for AE sensing. The high acoustic sensitivity also

AE detection using fibre grating lasers offers not only the



Figure 1. Schematic diagram of the YbEr-DFB laser deployment: LD, 977 nm laser diode; WDM, 980/1550 nm wavelength division multiplexer; FI-1, polarization-insensitive optical fibre isolator.

cavity, in particular when the vibration frequencies are close to the relaxation oscillation frequency of the laser [14]. A weak AE signal either initiates a relaxation oscillation of the laser, or modulates the laser output power with the same frequency as the AE signal. Therefore, the ultrasonic waves can be detected by directly monitoring the laser output power.

In this paper we report the development of AE sensors based on Yb^{3+}/Er^{3+} -codoped DFB (YbEr-DFB) fibre grating lasers. The response of the fibre laser to ultrasonic Lamb waves generated in a 3 mm thick aluminium plate has been investigated. Two sensor configurations have been investigated.

- A path-length-imbalanced readout Mach–Zehnder interferometer was used to measure the laser frequency shift induced by ultrasonic vibration.
- (2) The detection of ultrasonic waves by directly monitoring the laser output power.

2. YbEr-DFB fibre grating laser

Figure 1 shows a schematic diagram of the YbEr-DFB fibre laser deployed for AE detection. The laser head was a phaseshifted FBG written into an Yb³⁺/Er³⁺-codoped fibre (Koheras, C15). The length of the grating was 5 cm, and it was recoated with polyimide to a diameter of 155 μ m. Standard telecommunication fibre was spliced to one or both sides of the doped fibre, approximately 5 mm from the grating ends. A 977 nm fibre pigtailed solid state pump laser was coupled into the Yb³⁺/Er³⁺-codoped fibre through a 980/1550 nm fibre wavelength division multiplexer (WDM). Two YbEr-DFB fibre lasers were investigated, operating at wavelengths of 1549.7 (YbEr-DFB-1) and 1548.0 (YbEr-DFB-2) nm. The laser output was separated from the pump light through the WDM and was measured after a polarization-insensitive optical fibre isolator FI-1.

Part of the fibre containing the phase-shifted FBG must be placed on a heat sink to dissipate the heat from the selfheating due to pump laser absorption. The laser output was monitored using a scanning Fabry–Perot interferometer with 80 MHz resolution. The laser was observed to lase on two orthogonally polarized modes, separated by approximately 1.49 GHz. When the fibre of the first laser, YbEr-DFB-1, was twisted by approximately 180° over the 10 cm long fibre section containing the phase-shifted FBG, the laser operated



Figure 2. (a) Laser relative intensity noise spectrum of YbEr-DFB-1 for different pump powers. (b) Frequency of relaxation oscillation (squares) and laser power (triangles) versus pump power for the two YbEr-DFB lasers.

in a single polarization mode [15]. Using the delayed selfheterodyne method [16] with a 2 km fibre delay line, the linewidth of the laser was measured to be 100 kHz, limited by the measurement resolution.

Single-polarization-mode operation was also achieved by separating part of laser YbEr-DFB-2 from the heat sink and not twisting the fibre. A possible reason for this is that the nonuniform distribution of temperature along the fibre made the lasing conditions more favourable for one of the orthogonally polarized modes than the other. It is beyond the scope of this paper to discuss the laser mechanisms in more detail.

The measured laser output powers and relative intensity noise spectra are shown in figure 2. The noise spectra recorded



Figure 3. Schematic diagram of the fibre laser AE sensor. FI-1, FI-2, polarization-insensitive optical fibre isolators; FG, synthesized function generator; PA-g, PA-s, piezoelectric pressure transducers (Physical Acoustics, model WD, diameter 18 mm); A1, A2, Physical Acoustics 1220A pre-amplifier and AE1A amplifier; F-P, scanning Fabry–Perot interferometer; C1, C2, C3, 1550 nm 50:50 fibre couplers; AOM, acousto-optic modulator; PC, polarization controller; WDM, 980/1550 nm wavelength division multiplexer.

using an RF spectrum analyser with a bandwidth of 3 kHz for different pump powers are shown in figure 2(a). The background noise spectrum of the photoreceiver $(P_{\rm p} = 0)$ was recorded when the pump laser was turned off. The laser relaxation oscillation frequency increases as the pump power increases. In figure 2(b), the peak frequency of relaxation oscillation and laser power output from the isolator FI-1 are plotted against pump power for the two lasers. The threshold pump power was less than 10 mW. The single-frequency laser power was greater than 0.5 and 1.0 mW respectively for the two lasers for 20 mW of pump power, demonstrating that YbEr-DFB fibre lasers are more efficient than DFB fibre lasers using Er³⁺-doped fibre. Such a high-power signal source is an advantage for constructing the sensor system for highfrequency measurement as the photoreceiver amplifier can be optimized for high frequency without the requirement for high gain.

3. AE sensing experiment

Figure 3 shows the configuration of the AE sensing experiment using an YbEr-DFB fibre laser. An YbEr-DFB laser was placed on an aluminium plate, that was 3 mm thick and with lateral dimensions of $50 \times 50 \text{ cm}^2$.

3.1. Generation of ultrasonic Lamb waves in an aluminium plate

Ultrasonic waves were generated by a piezoelectric pressure transducer (PA-g) (Physical Acoustics, model WD), which was driven directly by a synthesized function generator without using an RF amplifier. The vibration was coupled to the aluminium plate using a Perspex wedge with the wedge angle approximately equal to the critical angle for the interface between Perspex and aluminium ($\theta = 25^{\circ}$), so that shear wave propagation is dominant in the aluminium plate [17]. In general the ultrasonic waves generated interact in a complex manner with the laser sensor. To simplify interpretation of the signals, the vibration frequency was kept below 500 kHz, so that the frequency–thickness product is smaller than 1.5 MHz mm and the ultrasonic wave in the aluminium plate could be approximately treated as a zero-order mode Lamb wave, i.e. antisymmetric mode A_0 and symmetric mode S_0 [17].

3.2. Coupling of vibration from the aluminium plate to the laser

When the ultrasonic vibration was generated in the aluminium plate, the laser frequency shift and power change results from the vibration-induced perturbations in the laser cavity. In our preliminary experiments [14], the entire length of the YbEr-DFB laser was placed on an aluminium plate using two pieces of 5 mm wide tape. However, the laser also responds when only part of the laser is exposed to the ultrasonic vibration. This allows high spatial resolution sensing. As shown in figure 1, a layer of lens tissue was used as the AE isolation layer. The 3 mm long sensing element of the DFB laser was covered by a layer of lens tissue and held on the plate using a piece of tape. Vacuum grease was used to improve the coupling between the fibre and the aluminium plate surface.



Figure 4. Measured phase spectrum of the read-out MZI. The test signal was at 470 kHz, with out-of-plane displacement $\Delta X \approx 3 \times 10^{-11}$ m. Bandwidth of spectrum analyser: 3 kHz.

3.3. Measurement of Lamb-wave-induced laser frequency shift

The laser output was divided into two using a 50:50 fibre coupler after the fibre isolator FI-1. A path-length-imbalanced readout MZI was used to measure the laser frequency shift induced by ultrasonic vibration. The optical path difference (OPD) of the MZI was 125 m. An acousto-optic modulator (AOM) was connected in one of the MZI arms to introduce a frequency shift of 34.98 MHz and enable heterodyne signal processing. A polarization controller was used to ensure that the polarization states in the two arms were matched. The two outputs from the MZI were detected using a balanced photoreceiver, and then processed by either a 100 MHz eightbit digitizing card and computer, or a digital oscilloscope/RF spectrum analyzer. The use of a balanced photoreceiver minimizes the laser power modulation signal when the two outputs of the MZI are approximately equal. The laser output from the second port of fibre coupler C1 was simultaneously monitored using a scanning Fabry-Perot interferometer.

A second piezoelectric pressure transducer (PA-s) positioned 10 mm from the centre of the YbEr-DFB laser was used to detect the Lamb waves, providing an independent measurement for the Lamb waves at the approximate location of the DFB laser. The output signal from PA-s was amplified by a pre-amplifier (A1), then by the amplifier (A2). PA-s and PA-g were the same type of pressure transducers. Vacuum grease was used to improve the coupling between the piezoelectric transducers and the aluminium plate surface.

Figure 4 shows the phase spectrum of the fibre laser recorded from the output of the MZI using a spectrum analyser. The first peak at frequency $f_0 = 34.98$ MHz resulted from the AOM. The lower peak at ($f_0 + 320$) kHz corresponds to the laser noise. The peak at ($f_0 + 470$) kHz appeared when the vibration at 470 kHz was generated in the aluminium plate.

In order to estimate the laser sensitivity to the Lamb waves, it is appropriate to assume that the out-of-plane displacement was dominant on the surface for the test signal at 470 kHz. This approximation was made since the amplitude ratio of the out-of plane displacement to the in-plane displacement of the A_0 mode on the surface of a 3 mm thick aluminium plate was approximately 1.5, and the ratio of the out-of plane



Figure 5. Response to the ultrasonic wave at 250 kHz generated by PA-g. (a) Voltage applied to PA-g, (b) output from PA-s and (c) laser YbEr-DFB-2.

displacement of the A_0 mode to the in-plane displacement of the S_0 mode was approximately 6 [17]. In addition, since the fibre was not bonded to the plate, the coupling coefficient for in-plane-displacement between the fibre and the plate must be lower than that for out-of-plane displacement.

The out-of plane displacement on the plate surface associated with the test ultrasonic wave propagating in the aluminium plate was approximately 3×10^{-11} m, based on measurements using a single-point laser Doppler vibrometer (Polytec, OFV-3001/303/OVD-02). The minimum detectable velocity was 0.5 μ m s⁻¹, corresponding to a displacement of 2×10^{-13} m at 470 kHz. The minimum detectable out-of plane displacement measured using a fibre laser was 1.5×10^{-13} m, corresponding to a resolution of 3×10^{-15} m Hz^{-1/2}.

3.4. Measurement of Lamb-wave-induced laser power modulation

The laser output power was found to respond directly to the ultrasonic Lamb waves. Figure 5 shows the response of the YbEr-DFB laser and PA-s sensor to the ultrasonic waves generated by the piezoelectric transducer, PA-g. The voltage applied to PA-g was a sine wave at 250 kHz (figure 3(a)). The output from PA-s and from the YbEr-DFB-2 laser are shown in figures 5(b) and (c) respectively. The output power of the laser is modulated at the same frequency as that of the ultrasonic wave. The laser power variation possibly results from the vibration-induced perturbation of the value of the phase shift constant of the DFB laser. Quantifying this result with the underlying physics is the subject of further research.

3.4.1. Effects of pump power. The vibration-induced laser power modulation was observed to be strongly frequency dependent. Since the vibration amplitude of the transducer, the vibration coupling loss from the transducer to the aluminium plate and the properties of Lamb waves in the plate depend strongly on frequency, it is complex to measure the vibration amplitude coupled to the laser as a function of frequency. The operation state of the YbEr-DFB laser itself has a strong effect



Figure 6. Laser relative intensity noise spectrum measured when an ultrasonic wave at 340 kHz is generated in the plate. (YbEr-DFB-2, bandwidth of the spectrum analyzer 3 kHz).

on its response to ultrasonic waves. Since the laser relaxation oscillation frequency increases as the pump power increases, the laser sensitivity will be different for the same test vibration signal if the pump power is varied. Therefore, the laser output power was measured for different pump powers at the same ultrasonic frequency.

When an ultrasonic wave at 340 kHz was generated in the plate, the relative intensity noise spectrum of YbEr-DFB-2 was measured for different pump powers using an RF spectrum analyser. The voltage amplitude applied to the transducer PAg was 5 V. Figure 6 shows the measured spectrum for a pump power of 23 mW. The peak frequency of relaxation oscillation, fr, was 476 kHz. The peak at 340 kHz results from the laser response to the test ultrasonic wave. The signal to noise ratio (S/N) at the test frequency was 27 dB, which can be used as the characteristic sensitivity of the laser to the ultrasonic wave. In figure 7, the measured S/N at 340 kHz is plotted against peak relaxation oscillation frequency, fr, which increased with pump power. The S/N increased as fr increased. With such laser sensitivity dependence on pump power, the laser sensor can be tuned to optimize the measurement of vibrations in a specified frequency range.

3.4.2. Laser response to a tone burst. The DFB fibre laser can be used to monitor the transient AE signals on the surface of the aluminium plate. Figure 8 shows the response to a five-cycle sine tone burst at 250 kHz. PA-g was 18 cm away from the centre of the DFB sensing section.

The group velocities of A_0 and S_0 modes at 250 kHz in an aluminium plate are 3.0 and 5.3 km s⁻¹ respectively [17]. The approximate 51 μ s delay before either PA-s or the fibre laser responds to the tone burst corresponds to the time taken for the A_0 mode to propagate to the sensors. The size and shape of the sensing area for each sensor type are very different. The optical fibre has a cylindrical shape and therefore very little is in direct contact with the surface of the plate and has an interaction area at least two orders of magnitude less than the piezoelectric sensor (approximately 1 mm² compared to 250 mm²). It is therefore probable that although the response of the two sensor types will be similar in form they will differ in detail. From



Figure 7. The measured signal to noise ratio at 340 kHz is plotted against peak frequency of the relaxation oscillation, fr, for YbEr-DFB-2.

these results it appears that the laser sensor has comparable sensitivity to the piezoelectric device, but when normalized for sensing area the fibre laser has much higher sensitivity. This issue will be the subject of future investigation.

3.4.3. Low energy impact detection. The laser response to impacts generated by dropping a steel ball is shown in figure 9. The ball of 4.7 mm diameter and 0.43 g in weight was dropped from a height of 7 cm above the horizontal aluminium plate, corresponding to an initial impact energy of 0.3 mJ. The impact location on the plate was 15 cm from the centre of the YbEr-DFB-1 fibre laser on an aluminium plate of $13 \times 50 \times 0.3$ cm³. The subsequent bursts of laser power shown in figure 9 resulted from the bouncing of the ball on the plate. The impacts occurring after the ball had bounced were weaker than the initial one, indicated by the smaller amplitude and shorter duration of the bursts of the laser power. Such a high sensitivity to impact could have applications in composite structures where the AE attenuation is very high.

4. Conclusions

YbEr-DFB fibre lasers have been used to detect ultrasonic Lamb waves generated in an aluminium plate. The vibration frequency was kept below 500 kHz, so that the ultrasonic wave in the 3 mm thick aluminium plate could be approximately treated as a zero-order mode Lamb wave.

A path-length-imbalanced readout Mach–Zehnder interferometer was used to measure the laser frequency shift induced by ultrasonic vibration. The use of a balanced photoreceiver minimizes the effects of laser power modulation. The out-of-plane displacement resolution was 3×10^{-15} m Hz^{-1/2} for a continuous ultrasonic wave at a frequency of 470 kHz.

We have demonstrated that the output power of an YbEr-DFB fibre laser can be modulated by the ultrasonic waves generated in an aluminium plate, which has been exploited to develop an alternative method for AE sensing. By directly monitoring the laser power, the sensor system configuration has been considerably simplified. Since the laser sensitivity to



Figure 8. Response to a five-cycle sine tone burst at 250 kHz. (a) Voltage applied to PA-g, (b) output from PA-s and (c) laser YbEr-DFB-2.



Figure 9. The laser YbEr-DFB-1 response to the impact of a steel ball of 0.43 g dropped from a height of 7 cm.

the vibration depends on the pump power, the laser sensor can be tuned to suite the measurement of vibrations at a specified frequency range. This simplified configuration is useful for transient AE signals. The detection of low energy impacts has been demonstrated, indicating the potential for applications in composite structures where the AE attenuation is high.

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References

- Vahaviolos S J (ed) 1998 Acoustic Emission: Standards and Technology Update (Philadelphia, PA: American Society for Testing and Materials) STP 1353
- [2] Spillman W B and Claus R O 2002 Optical-fiber sensors for the detection of acoustic emission *Mater. Res. Soc. Bull.* 27 396–9
- [3] Perez I, Cui H L and Udd E 2001 Acoustic emission detection using fiber Bragg gratings Proc. SPIE 4328 209–15
- [4] Betz D C, Thursby G, Culshaw B and Staszewski W J 2003 Acousto-ultrasonic sensing using fiber Bragg gratings Smart Mater. Struct. 12 122–8
- [5] Takahashi N, Yoshimura K and Takahashi S 2000 Detection of ultrasonic mechanical vibration of a solid using fiber Bragg grating Japan. J. Appl. Phys. 39 3134–8
- [6] Fisher N E, Webb D J, Pannell C N, Jackson D A, Gavrilov L R, Hand J W, Zhang L and Bennion I 1998 Ultrasonic field and temperature sensor based on short in-fibre Bragg gratings *Electron. Lett.* 34 1139–40
- [7] Leblanc M, Kirkendall C and Dandridge A 2000 Acoustic sensing using free and transducer-mounted fiber Bragg gratings *Proc. OFS-14* pp 592–5
- [8] Koo K P and Kersey A D 1995 Bragg grating-based laser sensors systems with interferometric interrogation and wavelength-division multiplexing J. Lightwave Technol. 13 1243–9

- [9] Ronnekleiv E 2001 Frequency and intensity noise of single frequency fiber Bragg grating lasers Opt. Fiber Technol. 7 206–35
- [10] Lovseth S W, Kringlebotn J T, Ronnekleiv E and Blotekjaer K 1999 Fiber distributed-feedback lasers used as acoustic sensors in air Appl. Opt. 38 4821–30
- [11] Hill D J and Nash P J 2000 In-water acoustic response of a coated DFB fibre laser sensor Proc. OFS-14 pp 33–6
- [12] O'Neill S F, Webb D J, Jackson D A, Zhang L and Bennion I 2000 Acoustic detection using erbium fibre Bragg grating lasers *Proc. OFS-14* pp 264–7
- [13] Loh W H, Samson B N, Dong L, Cowle G J and Hsu K 1998 High performance single frequency fiber grating-based

erbium:ytterbium-codoped fiber lasers J. Lightwave Technol. 16 114–8

- [14] Ye C-C and Tatam R P 2003 Acoustic-emission detection using Yb³⁺/Er³⁺-codoped fibre grating lasers *Proc. OFS-16* pp 218–21
- [15] Haratjunian Z E, Loh W H, Laming R I and Payne D N 1996 Single polarisation twisted distributed feedback fibre laser *Electron. Lett.* 32 346–8
- [16] Okoshi T, Kikuchi K and Nakayama A 1980 Novel method for high resolution measurement of laser output spectrum *Electron. Lett.* 16 630–1
- [17] Rose J L 1996 Ultrasonic Waves in Solid Media (Cambridge: Cambridge University Press)