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修士学位論文要約 (令和 2 年 9 月)

## ノンコリニア反強磁性 Mn-Sn 薄膜の作製と結晶構造・輸送特性の評価

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## Fabrication of noncollinear antiferromagnetic Mn-Sn thin films and evaluation of their crystal structures and transport properties

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I study the growth technique of non-collinear antiferromagnetic  $Mn_3Sn$  thin films deposited by sputtering with various substrates and underlayers. Relation between their crystal structure and magnetic/transport properties is also investigated. I achieve a formation of epitaxial films with both C-plane and M-plane orientations, whose Kagome lattices are parallel and perpendicular to the plane, respectively. Transverse resistivity originating from the anomalous Hall effect shows different trends reflecting the Kagome lattice orientations of each stack. The established technique and findings offer a platform to study functional devices utilizing unconventional physical properties of non-collinear antiferromagnets with controlled Kagome lattice orientation.

## 1. Introduction

Non-collinear antiferromagnetic materials with kagome lattice such as  $D0_{19}$ - $Mn_3Sn$  and  $Mn_3Ge$  have attracted increasing attention owing to their large anomalous Hall effect (AHE) originating from non-trivial Berry curvature mediated by the non-collinear spin texture<sup>1,2</sup>. In order to explore the functionalities of these phenomena for device applications, it is of importance to establish technique to grow thin films with a single crystalline phase.

Many physical properties of  $Mn_3Sn$  have been revealed in bulk single crystal<sup>1,2</sup>. Meanwhile, few studies on  $Mn_3Sn$  thin film have been performed and magneto-transport property of  $Mn_3Sn$  film has been investigated only for poly-crystalline samples<sup>3,4</sup>. In this work, I study the growth of epitaxial  $D0_{19}$ - $Mn_3Sn$  thin films with various crystal orientations. I also measure magneto-transport properties and discuss the relation between crystal orientation and transport properties.

## 2. Methods

I deposit  $Mn_3Sn$  thin films by magnetron sputtering at 400°C on various substrates and underlayers (ULs). The samples are annealed at 500°C for an hour. Crystal structure and transport properties are characterized by X-ray diffraction (XRD) and physical property measurement system, respectively. The transport measurement is performed for Hall-bar

devices at room temperature.

## 3. Analysis of crystal structures

Figure 1(a) shows  $2\theta$ - $\theta$  XRD spectra for Si sub./Ta(5 nm)/ $Mn_3Sn$ (50 nm) and Si sub./Ta(2 nm)/Ru(5 nm)/ $Mn_3Sn$ (30 nm) (thickness in nm). For UL = Ta, several weak peaks of  $Mn_3Sn$  indicating no preferential orientation are observed. For UL = Ta/Ru, on the other hand, prominent peaks of C-plane oriented  $Mn_3Sn$  are observed. However, as the  $Mn_3Sn$  thickness increases beyond 50 nm,  $(2\bar{2}01)$  peak indicating different orientation appears at  $2\theta \approx 42^\circ$  [Fig. 1(b)]. To overcome this issue and form epitaxial films, I use single crystalline MgO substrates. Figure 1(c) shows  $2\theta$ - $\theta$  spectra for MgO(111) sub./Ru(5 nm)/ $Mn_3Sn$  (30 or 80 nm). Only (0002) and (0004) peaks are observed regardless of the thickness, indicating C-plane orientation even for thicker films. I then turn to an M-plane oriented epitaxial  $Mn_3Sn$  film. I first use W as an underlayer because of a small lattice mismatch between W and  $Mn_3Sn$ . Figure 1(d) shows the  $2\theta$ - $\theta$  XRD spectra for MgO(110) sub./W (10 nm)/ $Mn_3Sn$  (50 nm). In addition to the expected  $(1\bar{1}00)$ ,  $(2\bar{2}00)$ ,  $(4\bar{4}00)$  peaks, unexpected peaks are observed at  $2\theta \approx 24^\circ$  and  $79^\circ$ , which can be attributed to  $WMn_2Sn$  formed through intermixing of W and  $Mn_3Sn$  layers. To avoid the formation of this intermixing layer, I insert a thin Ta layer between W and  $Mn_3Sn$ . As shown in Fig. 1(e),  $WMn_2Sn$  peaks are

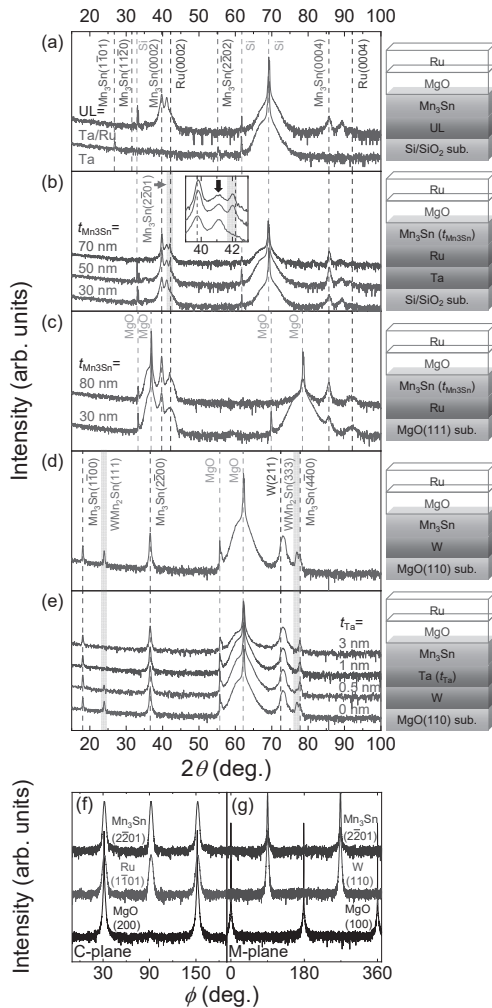


Fig. 1. (a)-(e) XRD spectra of  $2\theta$ - $\theta$  scans for the stack structures shown on the right. (f), (g) XRD spectra of  $\phi$ -scan for stack structures in Fig. 1 (c) and (e) respectively.

eliminated by the Ta insertion. Both the C-plane and M-plane oriented  $Mn_3Sn$  is epitaxial in in-plane orientation too as shown in Fig.1(f), (g).

#### 4. Magneto-transport properties

Figures 2(a)-(d) show Hall resistivity( $\rho_{xy}$ ) as a function of out-of-plane magnetic field ( $H$ ) in poly-crystalline, C-plane dominant poly-crystalline, epitaxial C-plane, and epitaxial M-plane  $Mn_3Sn$  thin films, respectively. Clear hysteresis with respect to the out-of-plane field is observed in M-plane-oriented stack [Fig. 2(d)], but not in the C-plane-oriented stack [Fig. 2(c)]. This result can be accounted for by AHE

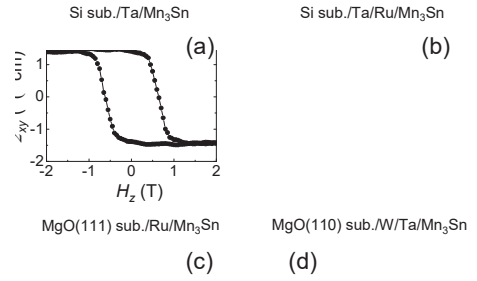


Fig. 2. (a)-(d)  $\rho_{xy}$  vs.  $H$  for  $Mn_3Sn$  thin films deposited with the stack structures above the figures.

induced by Berry curvature arising in the kagome plane.  $\rho_{xy}$  of poly-crystalline  $Mn_3Sn$  ( $1.4 \mu\Omega cm$ ) [Fig. 2(a)] agrees well with previous studies<sup>3,4</sup>. For stacks of Si sub./Ta/Ru/ $Mn_3Sn$ ,  $\rho_{xy}$  increases with  $Mn_3Sn$  thickness [Fig. 2(b)], in consistent with the formation of  $Mn_3Sn(2201)$  revealed by XRD.

#### 5. Summary

In summary, I study crystalline structure and transport properties of non-collinear antiferromagnetic  $Mn_3Sn$  thin films deposited by sputtering. Crystal orientation of  $Mn_3Sn$  is controlled by choosing appropriate substrates and ULs. AHE of  $Mn_3Sn$  depending on crystalline orientation is observed, proving anisotropic Berry curvature in non-collinear antiferromagnet. The obtained technique to prepare  $Mn_3Sn$  thin film with controlled crystal orientation provide the basis to study the functionalities of non-collinear antiferromagnets<sup>5</sup>.

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