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APPLIED PHYSICS

p-type transparent superconductivity in a layered oxide

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Development of p-type transparent conducting materials has been a challenging issue. The known p-type transparent conductors unsatisfy both of high transparency and high conductivity nor exhibit superconductivity. Here, we report on epitaxial synthesis, excellent p-type transparent conductivity, and two-dimensional superconductivity of $Li_{1-x}NbO_2$. The LiNbO₂ epitaxial films with NbO₂ sheets parallel to (111) plane of cubic MgAl₂O₄ substrates were stabilized by heating amorphous films. The hole doping associated with Li⁺ ion deintercalation triggered superconductivity below 4.2 kelvin. Optical measurements revealed that the averaged transmittance to the visible light of ~100-nanometer-thick Li_{1-x}NbO₂ was ~77%, despite the large number of hole carriers exceeding 10²² per cubic centimeter. These results indicate that Li_{1-x}NbO₂ is a previously unknown p-type transparent superconductor, in which strongly correlated electrons at the largely isolated Nb 4d₂₂ band play an important role for the high transparency.

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INTRODUCTION

Search for new transparent conductors (TCs) is one of the most important subjects not only for practical use but also for advances in materials science (1). In case of n-type TCs, a variety of materials have been developed and used in industry. For example, Sn-doped In₂O₃ (ITO) exhibits high conductivity (~6000 S cm⁻¹) and high transparency (~80%) (2). Furthermore, even a transparent superconductor (TSC), which is regarded as an ultimate TC, was reported for LiTi₂O₄ epitaxial films (3). In contrast to n-type TCs, however, p-type TCs (p-TCs) are subjects of fundamental research since their performances are still low (4).

Zhang *et al.* (5) recently proposed new strategy for designing TCs based on transition-metal oxides (TMOs). In contrast to a traditional concept (taking ITO as an example, high-mobility semiconductor In_2O_3 is doped with a small amount of Sn), they shed light on correlated metals, where large number of carriers resides in low-mobility TMOs. On the basis of this strategy, perovskite-type AMO₃ (A = Ca, Sr, and La; M = V, Mo, and Cr) have been found to be good TCs, both n- and p-type (5–8). However, the studies have been so far limited to the perovskite-type TMOs. Thus, further investigation on other TMOs will pave a way for better p-TC performances.

To realize an excellent p-TC and p-type TSC (p-TSC), one may investigate two-dimensional materials, which have been enthusiastically studied in these days (9). Among various TMOs, we focus on a layered niobate, which is the only isostructural oxide to 2*H*-type transition-metal dichalcogenide (TMD) such as MoS₂ (10). The crystal structure of LiNbO₂ consists of NbO₆ triangular prisms (Fig. 1A), which are thought to create characteristic d-band splitting for realizing p-type conduction (11, 12). Similar to TMDs, exotic properties can be expected (13, 14). In addition, the wide bandgap nature of TMOs provides high transparency in the visible range (15). However, electronic properties of LiNbO₂ are still unclear. Although superconductivity was observed below 5 K in a Li-deficient phase Li_{1-x}NbO₂ ($x \sim 0.5$) (16), its electronic properties have not been investigated in detail (17). Moreover, there is no report about single-crystalline superconducting films, and their transport and optical properties are not yet fully understood.

Here, we fabricated superconducting $\text{Li}_{1-x}\text{NbO}_2$ epitaxial films using three-step synthesis method as illustrated in Fig. 1B and investigated their electronic properties. p-type doping associated with Li^+ ion deintercalation to LiNbO_2 triggered superconductivity and enhanced transparency in the visible range, and we created a p-TSC with strongly correlated electrons in two-dimensional NbO₂ conduction layers and isolated Nb $4d_{z2}$ band derived from the NbO₆ triangular prism coordination.

RESULTS

Three-step synthesis and structural properties

Figure 1C shows out-of-plane x-ray diffraction (XRD) profiles of Li-Nb-O films. The films prepared by the high-temperature direct growth (direct) showed 002 reflection of layered LiNbO₂ at $2\theta = 17.0^{\circ}$ and a number of reflections from secondary phases, Li₃NbO₄ 222, NbO₂ 440, and LiNbO₃ 006 at $2\theta = 36.8^{\circ}$, 37.3°, and 38.9°, respectively (*18*, *19*). The pulsed laser deposition (PLD) growth of Li-containing materials such as LiCoO₂ and LiTi₂O₄ often tends to form Li-deficient phases because Li atoms are easily scattered by gaseous species and/or reevaporated from the growing surface (*20*, *21*). In our case, however, the Li-rich phase (Li₃NbO₄) indicated much stronger peak than the Li-deficient phase (NbO₂) and stoichiometric phases (LiNbO₂ and LiNbO₃). The inclusion and unusual stability of Li₃NbO₄ were further verified by its excellent crystallinity irrespective of growth conditions (see section S1).

To avoid the formation of the secondary phases and to create layered LiNbO₂, we developed a three-step synthesis, which was initiated from PLD growth of Li-Nb-O films on the unheated substrates (step 1). Out-of-plane XRD profiles of the films at each step are also shown in Fig. 1C. In step 1, no film peak was observed, suggesting the amorphous phase. After in situ reductive annealing (step 2), clear 00*l* reflections of LiNbO₂ phase appeared without Li₃NbO₄ peaks despite the condition identical to the direct growth. A small peak of the oxidized phase (LiNbO₃) is thought to originate from partial oxidation of the LiNbO₂ phase. LiNbO₃ has wide bandgap and insulating nature and does not affect the transport and optical properties described below. Last, the LiNbO₂ films were immersed into an I₂/CH₃CN solution to deintercalate Li⁺ ions

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Fig. 1. Three-step synthesis method and structural analysis for Li_{1-x}**NbO**₂ **epitaxial films.** (**A**) Schematic crystal structure of LiNbO₂. The green, blue, and red spheres indicate lithium, niobium, and oxygen atoms, respectively. The magnified unit of NbO₆ takes a triangular prism structure. (**B**) Schematic illustrations of three-step synthesis method. The experimental conditions and produced compounds at each step are shown in each box. (**C**) Out-of-plane x-ray diffraction (XRD) profiles for the direct film and the films at each step. Filled circles indicate reflections coming from the MgAl₂O₄ substrates. Filled and open triangles indicate reflections coming from a sample stage and secondary phases, respectively. arb. units., arbitrary units. (**D**) XRD ϕ scans of reflections for MgAl₂O₄ 440 in the substrate, LiNbO₂ 103 in the step 2 film, and Li_{1-x}NbO₂ 103 in the step 3 film.

(step 3). This solution process is known to take place with the following redox reaction (22)

Anodic reaction: $LiNbO_2 \rightarrow Li_{1-x}NbO_2 + xLi^+ + xe^-$ (1)

Cathodic reaction:
$$I_2 + 2e^- \rightarrow 2I^-$$
 (2)

After the Li⁺ ion deintercalation, only small shift of LiNbO₂ 00*l* peaks was observed, suggesting that NbO₂ layers in a host structure were preserved. Note that the crystallinity of the films (full width at half maximum of ω -scan profiles for LiNbO₂ 002 reflections) was unchanged.

We also measured asymmetric reflections (Fig. 1D). For both step 2 and step 3 films, the LiNbO₂ 103 reflection was clearly observed with a sixfold rotational symmetry, demonstrating solid-phase epitaxy at hexagonal planes (*10*). The angles of 440 reflections of the substrate coincided with those of the film, indicating an epitaxial relationship of LiNbO₂ [100] || MgAl₂O₄ <1T0>. The lattice parameters for step 2 (step 3) films were calculated to be *a* = 2.906 (2.914) Å and *c* = 10.446 (10.468) Å. Bulk values of Li_{0.93}NbO₂ (Li_{0.45}NbO₂) were reported to be *a* = 2.906 (2.923) Å and *c* = 10.447 (10.455) Å (*16*), suggesting that the step 2 films had nearly stoichiometric composi-

tion, while the step 3 films had Li⁺ ion deficiency. In addition, the in-plane lattice parameters at each step are considerably larger than that of MgAl₂O₄ ($d_{2\overline{2}0} = 2.858$ Å), indicating unstrained film lattices.

Transport properties and strongly correlated electrons

The direct and step 1 films indicated resistivity (ρ) exceeding a measurement limit. Figure 2A shows the temperature dependence of p for the films at the following steps. The step 2 films showed ρ less than 10 milliohm·cm at 300 K and metallic conduction $(d\rho/dT > 0)$. The step 2 films showed positive Hall coefficient, and hole concentration (*n*) was estimated to be 2.4×10^{21} cm⁻³ at 300 K (see section S2). On the other hand, decrease in p and change of the film color were found over the course of days, suggesting that oxidation upon exposure to air led to hole doping (23). To avoid any reaction in air, we prepared a capping layer (several-nanometer-thick amorphous alumina) before removing from the chamber (capped step 2). The capped step 2 films showed insulating behavior $(d\rho/dT < 0)$ with higher ρ (50 milliohm·cm at 300 K) and lower *n* (6.4 × 10²⁰ cm⁻³ at 300 K). Therefore, it is concluded that the metallic conduction in the step 2 films arises from easily oxidized (hole dopable) nature of LiNbO₂. It should be emphasized that the electronic ground state of LiNbO₂ has been controversial (17). Our results strongly suggest that $LiNbO_2$ is an insulator and supports a theoretical study (24).



Fig. 2. Transport properties and Fermi-liquid analysis. (A) Temperature dependence of resistivity for the films at each step. The labeled values are hole concentration at 300 K (*n*) obtained by Hall measurements. Inset: Superconducting properties of the step 3 films. Magnification of that is indicated in the main panel, and its magnetic field (along the *c* axis) dependence was taken up to 0.9 T. (B) T and (C) T^2 dependence of ρ normalized by the values at 300 K for the films at each step. Dashed lines are linear fits to the plots.

The step 3 films showed much lower ρ (3 milliohm·cm at 300 K), higher $n (2.5 \times 10^{22} \text{ cm}^{-3} \text{ at } 300 \text{ K})$, and a superconductivity below 4.2 K, demonstrating superconducting transition in Li_{1-x}NbO₂ triggered by hole doping associated with the Li⁺ ion deintercalation. Inset of Fig. 2A shows the magnification in the low-temperature range and modulation of superconductivity under magnetic field applied parallel to the *c* axis of $Li_{1-x}NbO_2$. The superconductivity was completely suppressed under a magnetic field as low as 0.9 T, which supports two-dimensional hole superconductivity in the NbO₂ plane (25). The critical magnetic field is much lower for a polycrystal (26). Moreover, superconductivity remained up to 4.2 T, when magnetic field applied perpendicular to the c axis, details of which will be published elsewhere. Further investigation on the two-dimensional hole superconductivity in our ideal samples is important, especially in connection with two-dimensional superconductors such as TMD and high-critical temperature $(T_{\rm C})$ cuprate superconductors.

To elucidate how electron correlation works in $\text{Li}_{1-x}\text{NbO}_2$, we replot Fig. 2A for Fig. 2 (B and C), where the normalized resistivity $(\rho/\rho_{300\text{K}})$ is plotted as a function of *T* (Fig. 2B) and T^2 (Fig. 2C). A comparison between these plots allows us to elucidate a Fermi-liquid picture in the transport properties. We clearly demonstrate enhanced strength of electron correlation and non-Fermi liquid ground states in the superconducting $\text{Li}_{1-x}\text{NbO}_2$ films. We observed the linear T^2 dependence of ρ for the metallic step 2 and step 3 films, indicating Fermi-liquid ground states. However, their temperature ranges are different from one to another. The step 2 films showed good linearity lasting up to 300 K. In contrast, the step 3 films showed deviation from ~100 K and, instead, followed *T* linear dependence. This characteristic temperature, referred to as *T**, is regarded as an important

phase point and/or a boundary in electronic phase diagrams of wellknown cuprate, iron-pnictide superconductors, and heavy-fermion compounds (27, 28). We find that, in our $\text{Li}_{1-x}\text{NbO}_2$, T^* showed a monotonic decrease by increasing the hole-doping level similar to well-established quantum critical systems, which we wish to publish elsewhere.

We performed fitting of T^2 plots to Fermi-liquid equation, i.e., $\rho = \rho_0 + AT^2$. We obtained $A = 3.3 \times 10^{-2}$ and 5.4×10^{-2} microhm·cm K⁻² for step 2 and step 3 films, respectively. The *A* parameters are directly associated with effective mass of carrier m^* ($A \propto m^{*2}$). Conventional metals and metal-oxide conductors exhibit $A \sim 10^{-6}$ microhm·cm K⁻² (29). On the other hand, strongly correlated metal oxides in the vicinity of Mott-Hubbard transitions (30) and heavy-fermion systems (29) show *A* in a range of 10^{-3} to 10 microhm·cm K⁻². The values obtained for our Li_{1-x}NbO₂ films are comparable to or larger than those of the latter, reflecting strong electron correlations in Li_{1-x}NbO₂ are comparable to that in superconducting Cu_{0.07}TiSe₂, an intercalated TMDs (1.1×10^{-2} microhm·cm K⁻²) (31). Furthermore, the roles of electron correlation in electronic structures of Li_{1-x}NbO₂ are also expected theoretically (32).

Optical properties and strongly isolated band structure

We have observed marked color changes throughout the three-step synthesis as shown in Fig. 3A. The direct and step 1 films indicated dark color typical of doped TMOs. In contrast, the step 2 films indicated characteristic bloody red, which is consistent with the previous



Fig. 3. Sample photographs and optical properties. (A) Optical images of the films at each step. Substrate size is about 5 mm by 5 mm. The transparent regions at the diagonal corners are trace of clamps that were used to hold the substrate to a metal backplate. (B) Optical transmittance spectra for the films at each step and $MgAl_2O_4$ substrate.

results (27). Furthermore, the step 3 films indicated transparent olive. Figure 3B shows optical transmittance in infrared to visible range for the films at each step and a pristine substrate. The first characteristic coloration in the step 2 films originated from drop of transmittance below 600 nm, which was consistent with calculated bandgap of layered LiNbO₂ (~2 eV) (24, 33, 34). The second color bleaching in the step 3 films resulted in increasing transmittance in the visible range (1.5 to 3.0 eV) reached ~50% despite a heavily hole-doped medium. Note that this value is the external transmittance, where substrate and light reflection at the surface and interface cause intensity loss. The T_{ave} after subtracting these contributions was as high as ~77%.

Taking difference between internal absorption spectra of $LiNbO_2$ and $Li_{1-x}NbO_2$ films into account, notable effects of the Li^+ ion deintercalation (hole doping) can be clearly captured (see section S3). The near–band edge absorption (2 to 2.5 eV) is greatly suppressed upon hole doping. In addition, the near-infrared absorption is hardly enhanced, although *n* increases by an order magnitude. The origin of this behavior will be explained later.

Performance as a p-TC

The p-TC performance of $\text{Li}_{1-x}\text{NbO}_2$ is summarized and compared with other p-TCs in Fig. 4, where their conductivities (σ) at room temperature (RT) are plotted as a function of T_{ave} . In general, σ and T_{ave} tend to indicate trade-off relationship as recently demonstrated in high-performance p-TCs of $\text{La}_{1-x}\text{Sr}_x\text{CrO}_3$ (LSCO) films. In this system, increase in *x* (number of holes) leads to increase in σ and decrease in T_{ave} (8). The lightly doped metallic LiNbO₂ (step 2) marks almost on the LSCO line. In contrast, the heavily doped superconducting $\text{Li}_{1-x}\text{NbO}_2$ marks beyond the border, highlighted by gradient background, and even violates the general tendency. Given



Fig. 4. Performances as a p-TC. Graphical representation of conductivity and transmittance for the LiNbO₂ (step 2) and Li_{1-x}NbO₂ (step 3) epitaxial films and reported p-TCs. Some of p-TCs are merely labeled, but all details including other p-TCs are listed in section S5. n-type ITO was also plotted as the ultimate performance for the purpose of exploring p-TCs. A series of LSCO, which is regarded as the guidance of TMO-based p-TC, was highlighted by a dashed arrow (8).

the record large figure of merit as a p-TC and superconductivity, $Li_{1-x}NbO_2$ can be regarded the first p-TSC.

DISCUSSION

Let us explain the origin of the anomalous transparency of $\text{Li}_{1-x}\text{NbO}_2$. The degree of transparency in near-ultraviolet to near-infrared range can generally be divided into lower (near-infrared) and higher (nearultraviolet) photon energy regime, which must be maintained against (i) free-carrier absorption, also known as a plasma reflection, and (ii) interband optical absorption, respectively (5). In conventional TCs, the higher σ becomes, the lower T_{ave} becomes due to (i). As for TMO-based TCs, T_{ave} can relatively be maintained against (i) due to large effective mass m^* of quasiparticle carriers. In contrast, $\text{Li}_{1-x}\text{NbO}_2$ rather harvests T_{ave} in terms of (ii) as will be explained as follows (schematically illustrated in Fig. 5).

Regarding (i), strongly correlated electrons decreased plasma reflection. Plasma energy ($h\omega_p$) is associated with ratio of *n* to m^* (5). Our results allowed us to estimate m^* numerically, taking previously reported parameters into account (24). Suppose $h\omega_p \sim 1$ eV for the step 3 films from decrease of transmittance in near-infrared region (Fig. 5), we obtain $m^* \sim 100 m_0$ (see section S4). On the other hand, $h\omega_p \sim 0.5$ eV was estimated for the step 2 films. The heavy mass is



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consistent with large A $(3.3 \times 10^{-2} \text{ and } 5.4 \times 10^{-2} \text{ microhm} \cdot \text{cm K}^{-2}$ for step 2 and step 3 films, respectively) described above. Li_{1-x}NbO₂ turned out to form strongly correlated electronic states, resulting in the lower $h\omega_p$ even when *n* exceeds 10^{22} cm^{-3} . Upon hole doping into LiNbO₂, $h\omega_p$ did not shift into the visible range. Thus, further carrier doping is acceptable to enhance the performance of p-TSC.

Regarding (ii), strongly isolated Nb 4d_{z2} band decreased interband optical absorption. According to our results and the reported band calculations, the Fermi level $(E_{\rm F})$ was thought to locate in the band on Nb 4d_{z2} orbital, which was strongly isolated from the upper (degenerated Nb d_{xy} and d_{x2-y2} orbitals) and the lower (O 2p orbitals) bands. Note that highly isolated single band has a potential to create high density of states at $E_{\rm F}$ and strong electron correlation, both of which are favorable to superconductivity (34). The characteristic layered NbO₆ triangular prism structure derives this prominent nature (24) and leads to the anomalous transparency. According to Fig. 5, bandgap becomes wider as $E_{\rm F}$ drops with increasing *n*. This effect known as Burstein-Moss effect extends a short wavelength cutoff to blue range (35). Combining these two effects in $Li_{1-x}NbO_2$ realizes appropriate transparent window ranging from 1.0 to 2.4 eV. One expects to expand this window to violet upon further deintercalation of Li⁺ ions. Synergetic increase in σ allows Li_{1-x}NbO₂ to approach to ITO (Fig. 4). The transmittance can also be improved by reducing film thickness.

In summary, we have first grown two-dimensional $Li_{1-x}NbO_2$ epitaxial films using specialized three-step synthesis method. The hole doping associating the Li⁺ ion deintercalation substantially enhances visible-light transparency and leads to superconductivity. We conclude that strongly correlated electrons in the isolated Nb $4d_{z2}$ band realize the p-type transparent superconductivity in $Li_{1-x}NbO_2$. Our study demonstrates a new approach to creating high-performance TCs.

MATERIALS AND METHODS

Sample fabrication

Li-Nb-O films were grown on MgAl₂O₄ (111) substrates using PLD method with KrF excimer laser (1.0 J cm⁻²). A Li excess ceramic $(Li_{1,2}NbO_{2+\delta})$ was prepared by conventional solid-state reaction steps, starting from mixing Li₂CO₃ and Nb₂O₅ powders with a molar ratio of 1.2:1. The direct growth was conducted at substrate temperature $(T_{\rm s})$ of 800°C under a chamber pressure $(P_{\rm Ar/H2})$ of 0.1 mtorr set by continues flow of Ar/H_2 (1 volume percent H_2) gas. The three-step synthesis method is illustrated in Fig. 1B. In step 1, using $Li_{1,2}NbO_{2+\delta}$ target, amorphous Li-Nb-O films were deposited in vacuum (background pressure, 1.0×10^{-7} torr) at RT. In step 2, Ar/H₂ gas was fed into the chamber to set $P_{Ar/H2} = 0.1$ mtorr, and the step 1 films were annealed in situ at $T_s = 800^{\circ}$ C for 1 hour. In step 3, the step 2 films were exposed to air and immersed in 0.02 M I₂/CH₃CN solution for 15 min. Some films were capped by several-nanometer-thick alumina films using PLD at RT in vacuum for avoiding reactions with air (capped step 2).

Characterizations

Using a stylus profiler, film thickness was regulated to be ~100 nm. The crystal structures and epitaxial relationship were investigated by a laboratory XRD apparatus with Cu K α_1 radiation. The temperature dependence of resistivity and the Hall voltage were measured by a standard four-probe method using a physical property

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measurement system (Quantum Design). The optical properties were investigated by ultraviolet-visible near-infrared spectroscopy at RT.

SUPPLEMENTARY MATERIALS

Supplementary material for this article is available at http://advances.sciencemag.org/cgi/ content/full/6/29/eabb8570/DC1

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