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AN INVESTIGATION INTO PRACTICAL ASPECTS  
OF TRANSPIRATION FLOWS

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A dissertation presented in fulfilment of the  
requirements for the Degree of Master of Science  
in Engineering

MARCH 1978

TO MY MOTHER AND FATHER .....  
AND SUMMERTIME

DECLARATION

I, MARK HANS ROBERT FRIEDRICH KRIEG, hereby declare that the dissertation is my own work and has not been submitted by me for a degree at any other University.



M.H.R.F. KRIEG

## ABSTRACT

It was felt that by attaching a porous matrix to a rigid base with metering holes or slots, these materials would become more viable in transpiration cooling. Such an alteration was expected to affect the mean velocity profile of the injected turbulent boundary layer. This, and the more basic study of turbulent transport, required attention. An investigation of the literature showed that although a great deal of work had been done, there was little to compare with the present problem.

A wind tunnel was designed and constructed which produced a 2-D incompressible turbulent shear layer on a porous plate, through which a secondary stream was injected. A flow situation with a step change in injection was thus modelled. A main stream velocity of 21 m/s (corresponding to  $Re_x = 1,178 \times 10^6$ ), and an average injection ratio ( $v_w/U_\infty$ ) of 0,017, was attained. Results indicated that blockage below the porous surface resulted in an apparently higher injection ratio for the same total flow rate.

To further the understanding of turbulent transport, a programme was written which solved the momentum equation. The eddy-viscosity formulation suggested by Cebeci was tested in this solution. It could be used to calculate the velocity profiles, with varying injection and pressure gradients, of either laminar or turbulent external boundary layers.

A numerically calculated profile and the experimental results of another researcher compared very well. The trend for the 'transition' region from flat plate turbulent flow to the injected turbulent boundary layer was established experimentally, and reflected in the computer

generated profiles. Exact comparison was not possible, owing to lack of time to rewrite part of the programme.

Experimental results were cross-plotted, and this indicated the possibility of a short-cut method for deriving velocity profiles with injection. Skin friction was not measured in the tests, and was found to be very inadequately documented in previous work. This must be a very important aspect of the injected turbulent shear layer, and requires further work.

A preliminary investigation of the flow in the porous matrix, especially the spread of an axisymmetric jet, was done. Initially experiment was expected to precede analysis, and a rig was built for future research.

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## NOTATION

### ROMAN

$a$	Pipe diameter
$A$	Van Driest damping length
$A^+$	Van Driest damping constant
$[A]$	Coefficient matrix in $[A]\{X\} = \{B\}$
$\{B\}$	Vector of constants
$c_f$	Skin friction coefficient
$d$	Orifice diameter
$D$	Pipe diameter
$D_o$	Particle diameter
$E$	Voltage drop across hot-wire
$E_o$	Voltage drop across hot-wire for zero velocity
$E_n$	Voltage drop, corrected for temperature
$E_1(n)$	Contribution to energy of fluctuation for frequency $n + dn$
$\underline{f}$	Force vector (per unit volume)
$f$	Dimensionless stream function, general function
$F$	Local injection ratio, $v_w/U_\infty$
$F_{av}$	Average injection ratio, $v_w/U_{Cl}$
$g$	General function
$G$	Main stream mass flux, $G = \rho_\infty U_\infty$
$h$	First step in $n$ grid
$H$	Shape factor
$k$	Roughness height
$k^+$	Non-dimensionalised roughness height, $k^+ = ku_w^*/\nu$

$K$	Ratio of adjacent $\eta$ intervals
$l$	Mixing length
$L$	Length scale, porous matrix thickness in direction of flow
$m''$	Injection mass flux
$n$	Frequency
$N-4$	Number of equations in system $[A]\{X\} = \{B\}$
$p$	Pressure
$p^+$	Pressure gradient term, $-dp/dx \cdot \nu / \rho (u_w^*)^3$
$Re_L$	Reynolds number, $Re_L = U_\infty L / \nu$
$t$	Temperature
$T$	Turbulence intensity
$u$	Axial velocity, $u^+ = u / u_w^*$
$u_w^*, u_\tau$	Wall shear velocity
$\underline{v}$	Velocity vector
$v$	Velocity normal to wall
$v_w$	Injection velocity, $v_w^+ = v_w / u_w^*$
$x$	Co-ordinate parallel to flow
$\{X\}$	Vector of unknowns
$XTR$	Start of porous plate, measured from origin
$y$	Co-ordinate normal to flow, $y^+ = y u_w^* / \nu$

GREEK

$\alpha$	Constant in outer eddy-viscosity formulation
$\beta$	Pressure gradient parameter, $\beta = (2\epsilon) \cdot (du_e/d\epsilon) / u_e$
$\gamma$	Intermittency term
$\delta$	Boundary layer thickness

$\delta^*$	Displacement thickness
$\Delta$	Differential, e.g. $\Delta p = p_2 - p_1$
$\epsilon$	Eddy viscosity, $\epsilon^+ = \epsilon/\nu$
$\eta$	Transformed y-co-ordinate
$\theta$	Momentum thickness
$\kappa$	Universal constant in logarithmic portion of turbulent profile
$\mu$	Dynamic viscosity
$\nu$	Kinematic viscosity
$\xi$	Transformed x-co-ordinate
$\rho$	Density
$\tau$	Shear stress
$\phi$	Translated stream function
$\psi$	Stream function

#### SUBSCRIPTS

$CL$	Centre line at Working Section inlet
$e$	Edge of boundary layer
$i$	Inner region
$o$	Outer region, unblown condition
$s$	Static, laminar sub-layer
$w$	Wall condition
$\underline{\quad}$	Vector
$\underline{\underline{\quad}}$	Second order tensor

SUPERSCRIPTS

- + Non-dimensional
- ' Derivative with respect to  $\eta$ ,  
fluctuating quantity
- Integral averaged quantity (time based,  
unless otherwise stated)

MATHEMATICAL OPERATORS

grad,  $\nabla$   $(\delta/\delta x \ \delta/\delta y \ \delta/\delta z)$

div  $\underline{v}$ ,  $\nabla \cdot \underline{v}$   $\delta u/\delta x + \delta v/\delta y + \delta w/\delta z$

$\nabla^2$   $\delta^2/\delta x^2 + \delta^2/\delta y^2 + \delta^2/\delta z^2$

## CHAPTER 1

### INTRODUCTION

The need for cooling engineering components becomes necessary when the components are in a high temperature environment. This problem frequently arises in flow situations. Important examples are :- turbine blades and combustion chambers in jet engines, exposed surfaces of supersonic vehicles such as aircraft, missiles and space craft. A further example is the lining of nuclear reactors. At present, turbine vanes and blades are film cooled by blowing cool air through discrete holes near the leading and trailing edge of the blades. This is a very important application, because the efficiency of the engine increases with cycle temperature.

An improvement in efficiency can have a dramatic effect in increasing the range of aircraft. Suciu (1970) showed that an aircraft with constant gross weight can transport a payload some 40% further if the cycle turbine inlet temperature is increased from 950°C to 1260°C. Metallurgical developments have resulted in much higher allowable metal temperatures, and film cooling has also made a significant contribution.

The coolant which is injected through the porous surface will certainly affect the boundary layer, skin friction, aerodynamic characteristics of an aerofoil (lift, drag) as well as the heat transfer coefficient. It was therefore decided that the mass transfer problem had to be investigated first. A great deal of work, both experimental and semi-analytic, has been done in this field. Unfortunately, because of the magnitude of the problem, and the large number of variables, tests have not been very systematic. Considerable thought was devoted to film cooling, and the

possibility of using porous materials for the manufacture of turbine blades. As a result, an experimental research programme was embarked on.

In particular, a turbulent boundary layer was established on a smooth flat plate before it reached a porous plate, where the flow encountered injection. A computer programme was written to simulate this situation. A deeper understanding of turbulent flow thus became possible, which is necessary not only for the abovementioned applications. In the design of air intake ducts for aircraft, turbulence and intermittency must be accounted for to ensure the required flow rate. Further applications are boundary layer control on aircraft wings, using suction or injection, and hence drag reduction.

The work of various authors is discussed. It was found that there were fundamental differences between this work and that of other researchers. This made comparison very difficult. A new approach to the use of porous materials in the manufacture of turbine blades is shown. A stable computational technique was developed and the existence of a new two-parameter family of curves for the injected turbulent boundary layer is presented.

## CHAPTER 2

### OBJECT

Investigation of existing turbine blades in aircraft gas turbines, e.g. Pratt and Whitney JT9, General Electric CF6 and Rolls Royce RB211, showed that film cooling has become an integral part in the design of first stage stator vanes and rotor blades. Impingement cooling is used in the second stage. Clearly, the technology for manufacturing these blades with internal ducting is fully developed. Welding a porous envelope onto a blade was suggested by Grootenhuis (1959).

Because the pressure distribution on an aerofoil would result in minimum transpiration at the leading edge, a complete envelope would be unsatisfactory. This could be overcome by having porous inserts at the leading edge and mid-chord positions. A configuration of solid wall, followed by porous wall thus arose. Furthermore, the porous wall is notoriously weak, and would require to be bonded to a solid wall at the inner surface to maintain the aerodynamic loads.

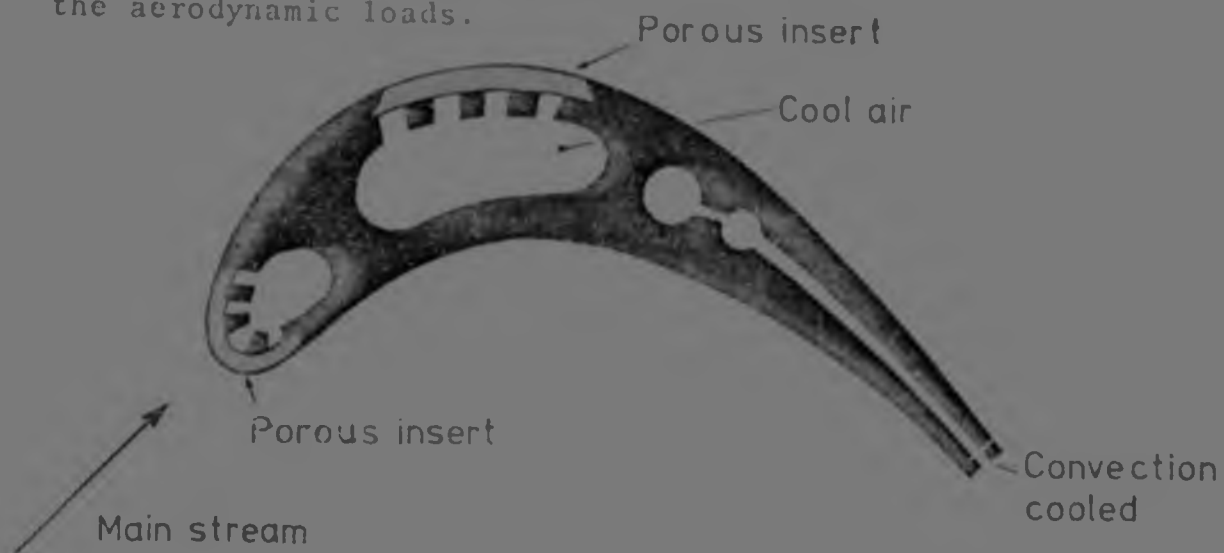


FIG 2.1 : CROSS SECTION OF BLADE WITH POROUS INSERTS



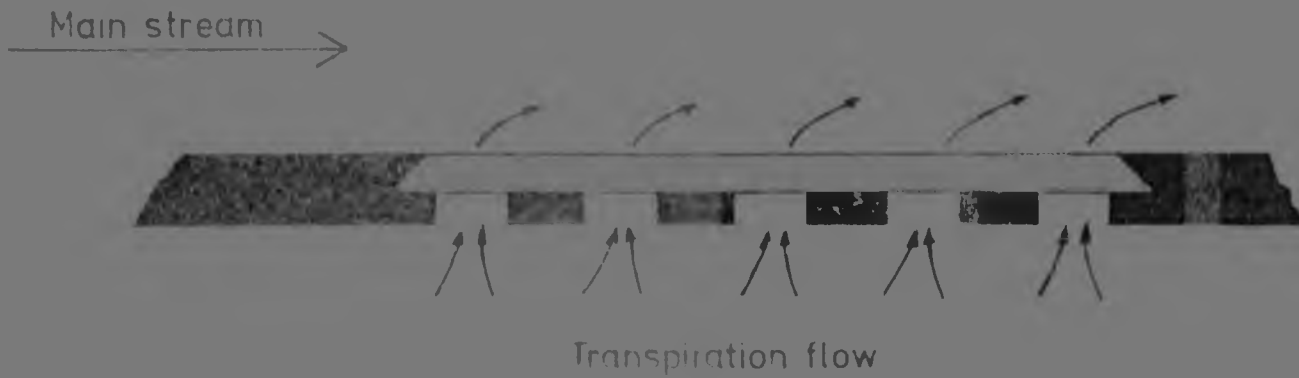


FIG 2.2 : BOUNDARY LAYER WITH STEP CHANGE IN ROUGHNESS AND INJECTION

It may be expected that the heat transfer from the metal to the coolant would be improved because the surface area is greatly increased in the porous medium. Furthermore, the webs between discrete holes in film cooling are very thin and hence prone to melting or combustion.

A great deal of research has been done in the field of transpiration cooling. However, the effect of a short porous insert with blockage below the porous matrix still remained to be explored.

To achieve this, an experimental rig, shown in Fig 2.3, was designed. This enabled tests on a flat plate in 2-D flow to be carried out. A set of data for uniform blowing was to be established, with which results for non-uniform injection could be compared.

The first stage was to consider only the effect of sudden mass transfer into an established turbulent boundary layer, without heat transfer.

The spread of the jet of coolant from the restricted inner porous surface to the open exposed surface clearly is an important consideration.

## CHAPTER 3

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### REVIEW OF LITERATURE ON THE TURBULENT BOUNDARY LAYER WITH INJECTION

---

#### 3.1 Introduction

One of the earliest papers to appear on transpiration cooling was published by Duwez and Wheeler (1948). The authors concentrated on a porous tube with high temperature axial flow and cold air injection through the walls. Their results compared well with that of Krieg (1975). In 1956 Mickley and Davis followed with an extensive experimental investigation which is still used for qualifying rigs, and for comparison with analytical or computer profiles. Chemical engineers have shown great interest in fluidized beds, ie. the flow inside a porous matrix. This latter line of investigation was considered relevant to the present work (see Appendix A3). The research may be categorised as follows :

- (i) Fluidized beds;
- (ii) Turbulent boundary layers with suction or injection at the wall :
  - a) in channels
  - b) in tubes
  - c) on a flat plate.

Each of these may be further subdivided into experimental, analytic and numerical results, with or without heat transfer, and into compressible or incompressible flow.

Note that injection through discrete holes or slots has enjoyed much attention (e.g. Hartnett *et al* 1961), but is not reported on here.

### 3.2 Experimental Work

As mentioned above, the first thorough experimental investigation was that of Mickley and Davis (1956). Free stream velocities ranged from 5,2 m/s to 18,3 m/s,  $Re_x$  from  $4 \times 10^4$  to  $3 \times 10^6$  and  $v_w/U_\infty$  values of 0; 0,001; 0,002; 0,003; 0,005 and 0,010. They found that the boundary layer decreased in thickness near the outlet of the tunnel. This latter problem was overcome by adding a flow divider to the outlet of the tunnel. The estimated values of  $c_f/2$  were shown plotted against  $Re_x$ . Where  $c_f$  tended to 0 or to negative values, the results were not reported. This occurred at values of  $F = 5 \times 10^{-3}$  and above. Negative  $c_f$  would indicate reverse flow, however, as will be discussed later, this need not be the case. Up to  $F = 3 \times 10^{-3}$ , the values of  $c_f$  compare with those predicted by McQuaid (1967), if these curves are extrapolated. Their plot of  $u^+$  vs  $\log y^+$  fitted the universal turbulent boundary layer profile acceptably. Graphs of  $u/U_\infty$  vs  $y/\delta$  were shown for  $F = 0,003$  and  $F = 0,005$ . To predict these profiles, it was necessary to allow  $\kappa$ , the mixing length constant, to vary.  $\kappa$  increased with  $F$ . The trend today is to assume  $\kappa$  constant, and to introduce a function in  $v_w$ .

Grootenhuis (1959) concentrated on heat transfer, as did most early investigators. Apart from the experimental results and correlation, this paper points out numerous applications of effusion cooling, discusses types of porous materials, shows thermal conductivity vs porosity for sintered bronze and stainless steel, and the internal heat transfer correlation for porous bronze. A list is given of manufacturers of porous materials, both sintered and woven, the various grades, their porosity and U.T.S. Blade designs were suggested, but required a complete porous envelope, spot welded to a hollow stem (see fig 3.1).

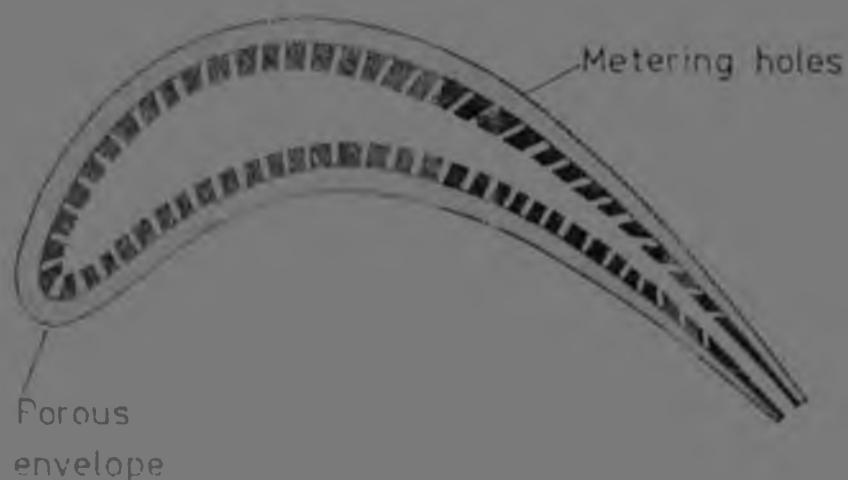


FIG 3.1 : PROPOSED DESIGN FOR A POROUS TURBINE BLADE OR VANE

A more detailed account was described by Moffat and Keys (1969). Their research was to establish a basis of experimental data for skin friction and heat transfer. Reynolds number, blowing fraction and Stanton number were varied systematically on a flat, porous plate. This floor was supported on 24 plenum chambers which allowed  $F$  to be controlled as a function of  $x$ . Electric heaters were installed below the porous plates, so that the transpiration flow could be heated. This unfortunately results, for an injected boundary layer, in an incorrect temperature profile. Results were presented for Stanton number varying with blowing fraction; these values showed good correlation.

The second part of this work was reported by Simpson *et al* (1969). Skin friction results were presented for  $v_w$  constant,  $v_w \propto x$ ,  $v_w \propto x^{-0.5}$  and  $v_w \propto x^{-0.5}$ .  $\sigma$  was presented as a function of  $h''/T(\sigma_{f,c}/T)$ . Various techniques for calculating  $\sigma_c$  from profile data were suggested. These included a momentum integral equation method, viscous sublayer model method and the heat transfer analogue method, which was the only method to yield results for

blowing fractions above 0,0078. Indications were that for  $b[-\dot{m}''/G(c_{f0}/2)]$  greater than 4,  $c_f$  tended to zero. Surface roughness was not included in the skin friction calculations.

Schetz and Nerney (1977) performed tests on an axisymmetric model. A skin friction balance was built into the porous wall : the floating element had its own transpiration supply. A crystal strain gauge was attached to the floating element support. The velocity profiles were obtained with a pitot rake. It was mentioned that the effect of roughness should first be ascertained for the porous wall without injection or suction. Note that the floating element is one of the more attractive techniques for measuring skin friction. Once true values are established, semi-analytic methods may be tested against these. It was found that blowing through the porous floating element made no difference to the results. The porous wall without injection showed significantly higher values of  $c_f$  than the equivalent smooth wall. The roughness increased the turbulence intensity in the boundary layer. With the new skin friction results, the laws of the wall, one by Simpson, one by Stevenson, could be tested. It was shown clearly that these laws are inadequate. The "deductive" approach of Van Driest (1956) could not be extended as had been suggested, but a new formulation using the Reichardt model  $\{\rho e = \kappa \rho \nu [y^+ - y_a^+ \tanh (y^+/y_a^+)]\}$  was presented. This new approach should yield significantly better results.

### 3.3 Analytical Developments

Before commencing this subsection, it should be noted that turbulent flow is of such a random nature that any analysis must contain empirical or experimentally determined constants. When considering quantities such as velocity, a time based average value is implied.

Mickley and Davis (1956) concentrated on experimental work, and the determination of  $c_f$ . This was done by evaluating  $d\theta/dx$ , and solving for  $c_f$  from :

$$\frac{d\theta}{dx} - \frac{v_w}{U_\infty} + (2 + \frac{\delta^*}{\delta}) \frac{\theta}{U_\infty} \frac{dU_\infty}{dx} = \frac{\tau_w}{\rho_\infty U_\infty} = \frac{c_f}{2} \quad \dots [3.3.1(a)]$$

For negligible pressure gradient, this reduces to :-

$$\frac{d\theta}{dx} - \frac{v_w}{U_\infty} = \frac{\tau_w}{\rho_\infty U_\infty^2} = \frac{c_f}{2} \quad \dots [3.3.1(b)]$$

The difficulty unfortunately lies in determining  $\frac{d\theta}{dx}$ .

The remainder of the analysis attempted to extend the mixing length theory for flow with  $v_w \neq 0$ . Reasonable correlation was obtained up to  $F = 0,005$ , but  $\kappa$  was no longer constant, but increased with  $F$ . To plot experimental data non-dimensionalized with respect to  $u_\tau^*$ , when this quantity is calculated from semi-analytic considerations, seems very unsatisfactory. Only if  $\tau_w$  is measured directly can an estimate of the reliability of the analysis be made.

It was assumed that there was only a laminar sublayer and a turbulent region (ie. no buffer layer), with a transition point at  $y_a$ . As do most researchers, they assumed the predominance of the  $v_w \delta u / \delta y$  term in the equation

$$u \frac{\delta u}{\delta x} + v \frac{\delta v}{\delta y} = \frac{1}{\rho} \left( \frac{\delta \tau}{y} - \frac{d\tau}{dx} \right) \quad \dots [3.3.2]$$

For an incompressible layer, the following expressions were obtained :

$$\text{Sublayer} \quad \frac{u \tau_w}{\nu} = 2 \frac{u \tau_w}{v_w} \ln \left( 1 + \frac{v_w u_a}{u \tau_w} \right) \quad \dots [3.3.3]$$

$$\text{Turbulent region} \quad \ln \frac{u \tau_w}{\nu} - \ln \frac{u \tau_w a}{\nu} = 2 \kappa \frac{u \tau_w}{v_w} \left[ \left( 1 + \frac{v_w u_a}{u \tau_w} \right)^{\frac{1}{2}} - \left( 1 + \frac{v_w u_a}{u \tau_w^2} \right)^{\frac{1}{2}} \right] \quad \dots [3.3.4]$$

Black and Sarnecki (1958) approached the problem in a very similar manner. The mixing length was considered proportional to distance from the wall,

$$l = \kappa y \quad \dots\dots [3.3.5]$$

Assuming incompressible, two dimensional flow, Prandtl's boundary layer equation [3.3.2] above, again forms the foundation of the analysis, with  $v_w \delta u / \delta y$  a dominant term.

The continuity equation is  $\frac{\delta u}{\delta x} + \frac{\delta v}{\delta y} = 0 \quad \dots\dots [3.3.6]$

and the shear is  $\tau / \rho = v \frac{\delta u}{\delta y} - u'v' \quad \dots\dots [3.3.7]$

with boundary conditions  $u = 0, v = v_w$  at  $y = 0 \quad [3.3.8]$

Close to the wall, with  $dp/dx > 0$ , equation [3.3.2] reduces to

$$v_w \frac{\delta u}{\delta y} = \frac{1}{\rho} \frac{\delta \tau}{\delta y} \quad \dots\dots [3.3.9]$$

[3.3.9] was integrated, to yield :

$$v_w u = \frac{\tau - \tau_w}{\rho} \quad \dots\dots [3.3.10]$$

and using the definition  $\frac{\tau_w}{\rho} = u_\tau^2 \quad \dots\dots [3.3.11]$

[3.3.12] was obtained :-

$$u_\tau^2 + v_w u = \tau / \rho \quad \dots\dots [3.3.12]$$

The momentum integral equation for the transpired boundary layer is thus :

$$\frac{d}{dx} (U_\infty^2 \theta) + \frac{d}{dx} u_w \delta^* = \frac{\tau_w}{\rho} + v_w u_w \quad \dots\dots [3.3.13]$$

Using the vorticity transfer theory of G.I. Taylor with  $l = \kappa y$  resulted in an expression which, for  $v_w = 0$ , did not reduce to the generally accepted law for flat plate flow. Prandtl's momentum transfer theory in conjunction with a linear mixing length yielded what Black and Sarnecki (1958) called the bilogarithmic law. This referred to the squared logarithmic term in their expression, derived as follows :

Substitute the momentum transfer formulation,

$$\frac{\tau}{\rho} = \kappa^2 y^2 \left(\frac{\delta u}{\delta y}\right)^2 \quad \dots\dots [3.3.14]$$

into [3.3.12] :

$$u_\tau^2 + v_w u = \kappa^2 y^2 \left(\frac{\delta u}{\delta y}\right)^2 \quad \dots\dots [3.3.15]$$

From [3.3.9],

$$\frac{\delta u}{\delta y} = \frac{1}{\rho v_w} \frac{\delta \tau}{\delta y} = \frac{1}{v_w} \frac{\delta(\tau/\rho)}{\delta y}$$

$$u_\tau^2 + v_w u = \kappa^2 \left(\frac{y}{v_w} \cdot \frac{\delta(u_\tau^2 + v_w u)}{\delta y}\right)^2 \quad \dots\dots [3.3.16]$$

$$(u_\tau^2 + v_w u) = \left(\frac{\kappa}{v_w}\right)^2 \left[\frac{\delta(u_\tau^2 + v_w u)}{\delta(\ln y)}\right] \quad \dots\dots [3.3.17]$$

This integrates to :

$$u_\tau^2 + v_w u = \left(\frac{v_w}{2\kappa} \ln y/d\right)^2 \quad \dots\dots [3.3.18]$$

where  $d$  is a constant of integration.

Equation [3.3.18] is expected to hold for 2-D incompressible flow, with  $\delta p/\delta y$  negligible (Prandtl's assumption for [3.3.2]), and a thin laminar sublayer.



[3.3.18] was expressed in a manner which allows both  $u_\tau$  and  $d$  to be determined from the slope and the intercept on the ordinate of a straight line drawn through the experimental points in the bilog region.  $\kappa$  was treated as a universal constant.

Thus

$$\frac{u}{U_\infty} - \frac{1}{4\kappa^2} \frac{v_w}{U_\infty} \left( \ln \frac{U_\infty d}{\nu} \right)^2 = \left[ \frac{1}{4\kappa^2} \frac{v_w}{U_\infty} \left( \ln \frac{U_\infty d}{\nu} \right)^2 - \frac{u_\tau^2}{v_w U_\infty} \right] - \left( \frac{1}{2\kappa^2} \frac{v_w}{U_\infty} \ln \frac{U_\infty d}{\nu} \right) \ln \frac{U_\infty y}{\nu} \dots [3.3.19]$$

$$\text{Let } \frac{1}{2\kappa} \sqrt{\frac{v_w}{U_\infty}} \ln \frac{U_\infty d}{\nu} = Y_i \dots [3.3.20]$$

$$n_i = - \frac{1}{2\kappa} \sqrt{\frac{v_w}{U_\infty}} \ln \frac{U_\infty d}{\nu} \dots [3.3.21]$$

$$p_i^2 = \frac{u_\tau^2}{v_w U_\infty} \dots [3.3.22]$$

$$\text{then } \frac{u}{U_\infty} - Y_i^2 = (n_i^2 - p_i^2) + 2n_i Y_i \dots [3.3.23]$$

Clearly, if  $\frac{u}{U_\infty} - Y_i^2$  is plotted against  $Y_i$ , and a straight line is fitted through the experimental points, the slope would be  $2n_i$ , and the ordinate intercept  $n_i^2 - p_i^2$ . This technique has numerous advantages over that of Mickley and Davis because the velocity scale is now the known, measured value of  $U_\infty$ , rather than  $u_\tau$ , obtained from the momentum integral equation.  $d\theta/dx$  could be calculated from :

$$U_\infty \frac{d\theta}{dx} = u_\tau^2 + v_w U_\infty \dots [3.3.24]$$

(3.3.13 for  $dU_\infty/dx = 0$ )

$$\frac{d\theta}{dx} = \frac{v_w}{U_\infty} (F_i^2 + 1) \dots [3.3.25]$$

For values of  $F > 0,005$ , it was found that the wall shear velocity became imaginary, ie.  $p_t^2 < 0$ . This gave  $\alpha_f < 0$ , which was what Mickley and Davis (1956) found, but ascribed to experimental error. It is very important to note that in the above technique, negative  $\alpha_f$  does not imply reversed flow, as was found in un-injected flows with adverse pressure gradient.

This lengthy discussion has been given because the bilog technique is used for comparing the present results with those of Mickley and Davis (1956). Values of  $u_\tau$  were not reported, as they were imaginary, and were only needed for the graphical presentation. Surface roughness was mentioned in the paper, but did not feature in the analysis.

McQuaid (1967) introduced intermittency into his formulation. A two parameter family of curves was mentioned, but not given in this paper. The velocity profiles generated from the family were compared with some experimental data. Results were good, also when fitted to the Mickley and Davis results.

The law of the wall for  $v_w = 0$  is

$$\frac{u}{u_\tau} = f\left(\frac{u_\tau y}{\nu}\right) \quad \dots [3.3.26]$$

This could be extended for injected flows by adding  $v_w$ , e.g.

$$\left. \begin{aligned} \frac{u}{u_\tau} &= f_1\left(\frac{u_\tau y}{\nu}, \frac{v_w}{u_\tau}\right) \\ \text{or } \frac{u}{v_w} &= f_2\left(\frac{u_\tau y}{\nu}, \frac{u_\tau}{v_w}\right) \end{aligned} \right\} \dots [3.3.27]$$

and to include surface roughness

$$\left. \begin{aligned} \frac{u}{u_\tau} &= f\left(\frac{u_\tau y}{\nu}, \frac{v_w}{u_\tau}, \frac{u_\tau k}{\nu}\right) \\ \text{ie. } u^+ &= f(y^+, v_w^+, k^+) \end{aligned} \right\} \dots [3.3.28]$$

McQuaid, however, suggested a function using  $U_m$  as a velocity scale, and retained wall shear in the term  $c_f$ , thus :

$$\frac{u}{U_\infty} = g(y/\theta, H, Re_\theta, \frac{v_w}{U_\infty}) \dots\dots\dots [3.3.29]$$

and

$$\left. \begin{aligned} c_f &= c_f(H, Re_\theta, \frac{v_w}{U_\infty}) \\ Re_{\delta_B} &= Re_{\delta_B}(H, Re_\theta, \frac{v_w}{U_\infty}) \end{aligned} \right\} \dots\dots\dots [3.3.30]$$

( $\delta_\theta$  was defined as "twice the distance from the wall to the position where  $\gamma = 0,5$ ").

McQuaid assumed the law presented by Black and Sarnecki for the sublayer [3.3.3], and modified Stevenson's 'law of the wall' for the turbulent region with suction or injection :

$$2 \frac{u}{v_w} \left[ \left( 1 + \frac{v_w}{U_\infty} \frac{U_\infty^2}{u_\tau^2} \right)^{1/2} - 1 \right] = A \log \frac{u_\tau \delta_B}{\nu} + B \dots [3.3.31]$$

A pressure gradient term was also derived, and was shown to be analogous to the pressure gradient term for unblown flows. This was given as :

$$\Delta_z = \frac{\nu}{(u_\tau^2 + v_w u_\tau)^{1/2}} \cdot \frac{1}{\theta} \cdot \frac{d\theta}{dz} \dots\dots\dots [3.3.32]$$

It was found that the model gave good agreement provided  $-0,004 < \Delta_z < 0,006$ . Furthermore, and this is of importance to the present work, the "fully developed state was not attained even at 33 boundary layer lengths downstream after a sudden change in injection rate".

Although this paper clearly indicates the importance of  $\theta$  in transpired boundary layers, the objections raised previously to  $\tau_w$  being calculated from the velocity profile, and not direct measurement, again apply.

Pletcher (1969) analysed an axisymmetric configuration, and developed a mixing length based on the Van Driest damping factor (Van Driest, 1956). An explicit finite difference technique is explained in the paper, and purports to give good results for flat plate, uninjected flows. The injected velocity profile shown agrees well with data from one of the Mickley and Davis tests. It was pointed out that the explicit method is more direct, and simpler to programme than an implicit method. Furthermore, it was not necessary to linearize the boundary layer equation before solving it, thus eliminating the need for an iterative process. However, it is well known that explicit methods are prone to instability, especially if the step length becomes too small.

Cebeci and Smith (1970a) also solved the Prandtl boundary layer equation for axisymmetric, compressible flow. This paper was preceded by numerous others that were very similar. This particular version showed the equations, boundary conditions, the finite difference grid and transformations in some detail. An implicit finite difference technique was employed to solve the linearized momentum equation with boundary conditions. A mixing length formulation based on Van Driest (1956) was used for the sublayer and the blending region, and a formulation using intermittency for the fully turbulent region. The energy equation was solved simultaneously, as this was a compressible flow situation. Grid spacing and CPU time required for solution were discussed.

Cebeci (1970b) showed how the Van Driest damping factor could be modified for pressure gradient and injection. Comparison with experimental results appeared excellent. The Van Driest velocity gradient relationship was extended to flows with injection :-

$$\frac{du^+}{dy^+} = \frac{2(v_w^+ u^+ + 1)}{1 + \{1 + 4(v_w^+ u^+ + 1)\kappa^2 (y^+)^2 [1 - \exp(-y^+/A^+)]^2\}^{1/2}} \dots [3.3.33]$$

where  $A^+$  is the modified Van Driest damping constant.

Equation [3.3.33] suggests an iterative solution for  $u^+$ .

This work by Cebeci will be discussed in greater detail in Chapter 4.

Cebeci and Smith (1974) and Cebeci and Bradshaw (1977) continued the above work, but suggest changes to the Reynolds stress formulation, and a very much more sophisticated finite difference technique, known as Keller's box method. The Falkner-Skan equation requires additional transformation before computer solution is possible. The second book is of more general interest showing how Runge-Kutta can be used, and gives examples of computer programmes.

Bale (1975) concentrated on axial flow through a tube with mass extraction at the walls. A potential flow solution was obtained, which gave a good representation of the flow field inside the porous pipe, but gave no information with respect to shear stress at the wall.

Schetz and Nerney (1977), as mentioned above, concentrated on experimental work, and the direct measurement of skin friction. It was shown that an extension to the Van Driest mixing length could not account for surface roughness, but the Reichardt model could. The momentum equation could thus be written :

$$\frac{du^+}{dy^+} = \frac{1 + v_w^+ u^+}{1 + \kappa (1 + v_w^+ u^+) (y^+ - y_a^+ \tanh (y^+/y_a^+))} \quad [3.3.34]$$

(cf. equation 3.3.33)

$y_a^+$  is a length scale, and is of the order of the laminar sublayer thickness. It was expected that  $y_a^+$  would decrease with roughness, and  $y_a^+ = y_a^+(k^+)$  was determined. The relationship  $y_a^+ = y_a^+(k^+, v_w^+)$  still needs to be found experimentally for various values of  $k$ ,  $v_w$  and types of roughness. For the unblown case, the Clauser shift for roughness was confirmed :

$$\frac{u}{u_\tau} = A \log \frac{u u_\tau}{\nu} + B - \Delta(u/u_\tau) \quad \dots\dots [3.3.35]$$

ie. roughness did not change the slope of the fully turbulent region, only the intercept on the ordinate.

Presented above are a few papers written on the subject of the transpired boundary layer. There are innumerable papers on this topic : studying each one would result in a great deal of repetition, and possibly be of little use until the skin friction for shear flows with transpiration has been truly established. From this section, it should be clear that the present situation bears close resemblance to the confusion which existed in the 1940's as regards the fully turbulent region of the solid-wall turbulent boundary layer. With the introduction of reliable skin friction transducers, that position was clarified.

## CHAPTER 4

### A THEORETICAL APPROACH TO THE TURBULENT BOUNDARY LAYER WITH TRANSPIRATION

#### 4.1 Introduction

The analysis follows the classical treatment by Prandtl for the boundary layer on a flat plate. For the turbulent shear layer, the Prandtl momentum transfer theory is considered and is modified by introducing the Van Driest damping factor. This technique is used to extend the analysis to include transpiration and pressure gradient. For the fully turbulent region, intermittency is introduced into the transport equation.

It was felt that more information regarding the boundary layer could be obtained following this type of analysis than using a potential flow solution. Knowledge with respect to skin friction, boundary layer growth and intermittency became available, and the analysis could be extended to include heat transfer and compressible flows.

#### 4.1.1 Definition of parameters featured in turbulent flow analysis

One of the most important quantities in any analysis of this type is  $\tau_w$ , the shear stress at the wall. In non-dimensional form,  $\tau_w$  appears in the local skin friction coefficient,  $c_f$ .

$$c_f = \frac{\tau_w}{\rho u_w^2} \quad \dots\dots [4.1.1]$$

When analysing the turbulent shear layer, the wall shear velocity is invariably the velocity scale. By definition,

$$u_{\tau}^2 = \frac{\tau_w}{\rho} \dots\dots\dots [4.1.2]$$

$$u_{\tau} = u_e \sqrt{\frac{\tau_w}{\rho u_e^2}} \dots\dots\dots [4.1.3]$$

The thickness of the boundary layer is denoted by  $\delta$ .  
 $y \rightarrow \delta$  as  $du/dy \rightarrow 0$ . A more practical definition is

$$y = \delta \quad \text{when } u = 0,99u_e \dots\dots\dots [4.1.4]$$

$$\text{Thus } u_e = 0,99u_{\infty} \dots\dots\dots [4.1.5]$$

The displacement thickness,  $\delta^*$ , is the mass flow deficit which occurs in the boundary layer owing to the reduced velocity near the surface :

$$\rho u_e \delta^* = \int_0^{\delta} \rho (u_e - u) dy \dots\dots\dots [4.1.6]$$

$$\frac{\delta^*}{\delta} = \int_0^1 \left(1 - \frac{u}{u_e}\right) \frac{dy}{\delta} \dots\dots\dots$$

A deficiency in momentum in the boundary layer may be expressed as a momentum thickness,  $\theta$ :

$$\rho u_e^2 \theta = \int_0^{\delta} \rho u (u_e - u) dy \dots\dots\dots [4.1.7]$$

$$\frac{\theta}{\delta} = \int_0^1 \frac{u}{u_e} \left(1 - \frac{u}{u_e}\right) \frac{dy}{\delta} \dots\dots\dots$$

Note that the upper integrated limit is given here as  $\delta$ , but is more correctly  $\delta^*$ .

The ratio of displacement to momentum thickness,  $H$ , is a very useful shape factor,

$$H = \delta^* / \theta \dots\dots\dots [4.1.8]$$

\* see Addendum for further explanation.



4.2 The Continuity and Momentum Equations

The derivation of the Prandtl boundary layer equations for an incompressible, newtonian fluid is well known (see Schlichting, 1968).

Continuity :  $\frac{\delta \bar{u}}{\delta x} + \frac{\delta \bar{v}}{\delta y} = 0$  ..... [4.2.1]

Momentum :  $\bar{u} \frac{\delta \bar{u}}{\delta x} + \bar{v} \frac{\delta \bar{u}}{\delta y} = u_e \frac{d u_e}{d x} + \frac{1}{\rho} \frac{\delta}{\delta y} (\mu \frac{\delta \bar{u}}{\delta y} - \rho \overline{u'v'})$  [4.2.2]

The bar denotes a time based average quantity, and is assumed in the following equations.

At the edge of the boundary layer :

$u_e \frac{d u_e}{d x} = - \frac{1}{\rho} \frac{d p}{d x}$  ..... [4.2.3]

4.3 Formulation for the Reynolds Stress

To solve equation [4.2.2], it is necessary to relate the Reynolds stress,  $\rho \overline{u'v'}$  to the mean velocity. Following Schlichting (1968), a formulation analogous to the laminar shear stress is assumed :

$\tau_x = \rho \epsilon \frac{\delta \bar{u}}{\delta y}$  ..... [4.3.1]

and then substitute Prandtl's momentum transfer theory for the eddy kinematic viscosity :

$\epsilon = l^2 \left( \frac{\delta \bar{u}}{\delta y} \right)$  ..... [4.3.2]

where  $l = \kappa y$  ..... [4.3.3]

This approach leads to a useful velocity distribution for the fully turbulent region. Near the wall the profile is linear. See Grimson (1971) for the derivation of [4.3.4] and [4.3.5].

Laminar sublayer :  $u^+ = y^+$  ..... [4.3.4]

Fully turbulent region :  $u^+ = \frac{1}{\kappa} \ln y^+ + B$  ..... [4.3.5]

The Buffer or Blending region is not predicted.

Van Driest (1956) derived a continuous velocity and shear distribution by modifying the mixing length. By considering Stokes flow (ie. an infinite plate oscillating sinusoidally parallel to itself).

The modified mixing length became :-

$l = \kappa y [1 - \exp(-y/A)]$  ..... [4.3.6]

Clearly, as  $y$  increases, the exponential term decreases, and  $l \rightarrow \kappa y$ , as before, for the fully turbulent region.

Using [4.3.6], [4.3.2] and [4.3.1],

$\tau_t = -\rho \overline{uv} = \rho \kappa^2 y^2 [1 - \exp(-y/A)]^2 \left| \frac{\delta u}{\delta y} \right| \left| \frac{\delta u}{\delta y} \right|$  ..... [4.3.7]

$A$  is the Van Driest damping length, dependent on the turbulence intensity and kinematic viscosity of the flow. It was determined experimentally, and in non-dimensional form,  $A^+ = 26$  for  $\kappa = 0,4$ .

$\tau = \tau_l + \tau_t$  ..... [4.3.8]

$\tau = \mu \frac{\delta u}{\delta y} + \rho \kappa^2 y^2 [1 - \exp(-y/A)]^2 \left| \frac{\delta u}{\delta y} \right| \left| \frac{\delta u}{\delta y} \right|$  ..... [4.3.9]

Cebeci (1970b) modified  $A$  to include injection and pressure gradient :

$A^+ = \frac{\Delta u \omega^+}{\nu}$  ..... [4.3.10]

Modified,  $A^+ = 26 \left\{ \frac{p^+}{v_w^+} [\exp(11.8 v_w^+) - 1] + \exp(11.8 v_w^+) \right\}^{-\frac{1}{2}}$   
 ..... [4.3.11]

where  $p^+ = \frac{dv}{dx} \frac{v}{\rho (u_w^+)^3}$ ;  $v_w^+ = \frac{v_w}{u_w^+}$  ..... [4.3.12]

4.4 Eddy Viscosity Formulation for the Fully Turbulent Region

In turbulent boundary layers, the potential core often extends well into the boundary layer. A pictorial representation of the instantaneous layer is depicted below :-

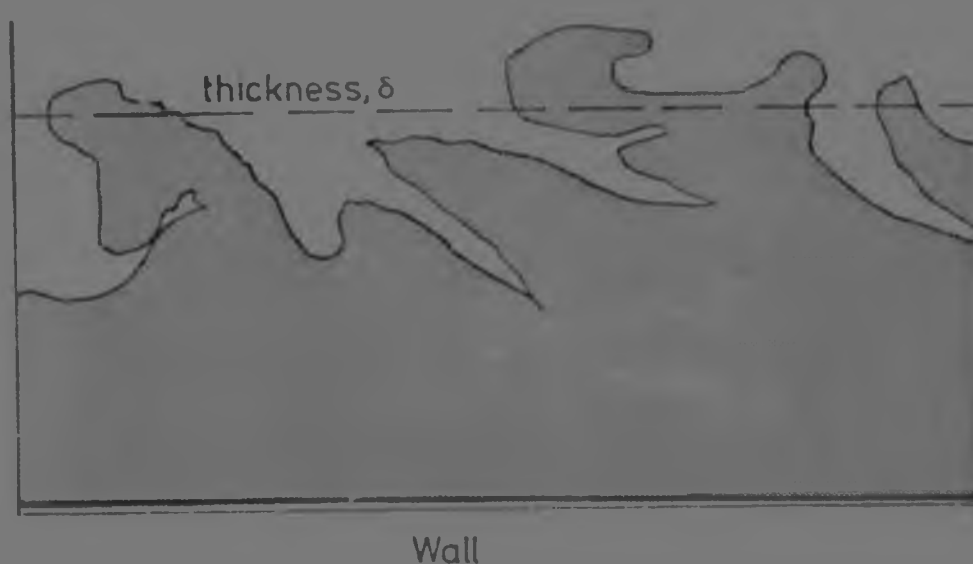


FIG 4.1 : SKETCH OF THE INSTANTANEOUS TURBULENT BOUNDARY LAYER, SHOWING THE EXTENSION OF THE POTENTIAL CORE BELOW THE AVERAGE THICKNESS

If the Reynolds stress in the potential core is zero, then the intermittent extension of this 'clean' flow into the boundary layer must influence the shear distribution across the outer turbulent region. Sarnecki postulated that  $\gamma$ , the intermittency, was the fraction of total time for which the flow was turbulent at a fixed distance from the plate.  $\gamma$  ranges from 0 to 1, where  $\gamma = 1$  inside the continuously turbulent boundary layer, and  $\gamma > 0$  at the edge of the layer.

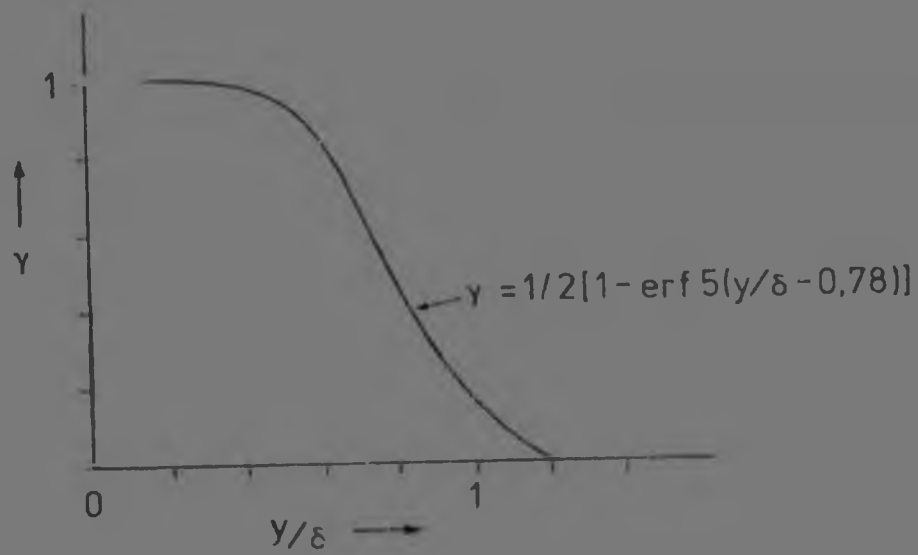


FIG 4.2 : INTERMITTENCY DISTRIBUTION ACROSS A  
TURBULENT BOUNDARY LAYER  
(CEBECI, 1974)

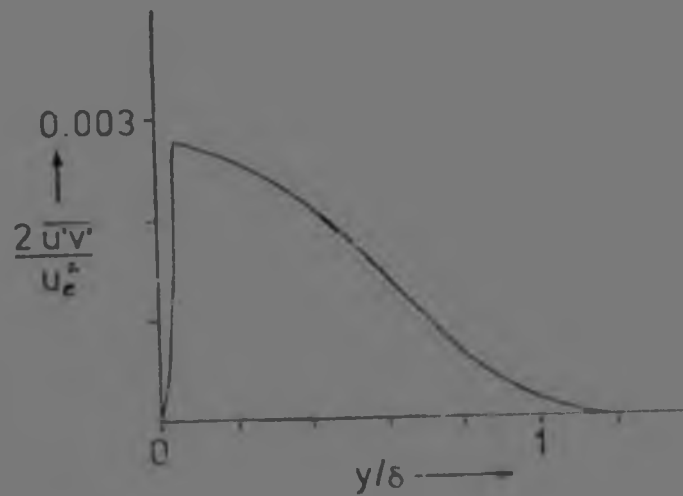


FIG 4.3 : REYNOLDS STRESS ACROSS THE BOUNDARY  
LAYER FOR A SMOOTH FLAT PLATE,  
ZERO PRESSURE GRADIENT

Intermittency and Reynolds stress were determined experimentally with hot wire anemometry by Klebanoff (Cebeci 1974, Schlichting 1968).

Clearly, even the modified mixing length theory cannot predict  $\epsilon$  across the entire boundary layer. For in the outer region,  $\epsilon$  must decrease, and experiments have shown that the decrease follows the intermittency curve, as could be expected. The eddy viscosity in the outer region was given as :

$$\epsilon_0 = 2 \left[ \int_0^y (u_e - u) dy \right] \times \frac{1}{2} [1 - \operatorname{erf} 5(y/\delta - 0,78)] \quad [4.4.1]$$

The intermittency curve may be approximated (Cebeci 1974) by :

$$\gamma = \frac{1}{[1 + 5,5(y/\delta)^6]} \quad \dots\dots [4.4.2]$$

$\gamma$  is a constant, with a value between 0,016 and 0,02. In the present work,  $\gamma = 0,0168$  is used. In the programme, the turbulent transport eddy viscosity was used until it was equal to the outer eddy viscosity. From this point outwards,  $\epsilon_0$  was used.

[4.4.1] may be written

$$\epsilon_0 = \alpha u_e \delta^* \gamma \quad \dots\dots [4.4.3]$$

#### 4.5 Boundary Conditions

The boundary conditions for the turbulent shear layer are :-

$$u(x,0) = 0 \quad \dots\dots [4.5.1]$$

$$v(x,0) = 0 \quad \text{or} \quad v_w \quad \text{with mass transfer} \quad \dots\dots [4.5.2]$$

$$\lim_{y \rightarrow \infty} u(x,y) = u_e(x) \quad \dots\dots [4.5.3]$$

4.6 The Transformed Momentum and Related Equations : A Summary

Only the transformed equations are given here. Details of the transformation appear in Appendix A1.  $\phi$  is a transformed, non-dimensionalized stream function.

Momentum :

$$[(1 + \epsilon^+) \phi'']' + (\phi + \eta) \phi'' - \beta(\phi' + 2) \phi' = 2\xi [(\phi' + 1) \frac{\delta \phi'}{\delta \xi} - \phi'' \frac{\delta \phi}{\delta \xi}] \quad \dots [4.6.1]$$

Boundary conditions :

$$\phi(\xi, 0) = 0 \quad \text{for no transpiration} \quad \dots [4.6.2a]$$

$$\phi(\xi, 0) = -\frac{1}{(2\xi)^{1/2}} \int_0^{\xi} \frac{v_w}{u_e} d\xi \quad \text{with } v_w \neq 0 \quad \dots [4.6.2b]$$

$$\phi'(\xi, 0) = -1 \quad \dots [4.6.3]$$

$$\lim_{\eta \rightarrow \eta_\infty} \phi'(\xi, \eta) = 0 \quad \dots [4.6.4]$$

Eddy viscosity equations :

$$\epsilon_i = \kappa^2 (2\xi)^{1/2} \eta^2 \frac{\rho''}{\rho} [1 - \exp\{-\eta((\frac{2\xi}{\mu})^{1/2} \cdot \phi_w'')^{1/2} / A^+\}] \quad \dots [4.6.5]$$

$$\epsilon_o = \alpha \cdot [ -(\frac{2\xi}{\rho})^{1/2} \int_0^{\eta_\infty} \phi' d\eta ] \times [1 + 5.5 (\frac{\eta}{\eta_\infty})^6]^{-1} \quad \dots [4.6.6]$$

Boundary layer parameters :

$$\sigma_f = \mu (\frac{2}{\xi})^{1/2} \phi''(\xi, 0) \quad \dots [4.6.7]$$

$$\delta^* = -\frac{(2\xi)^{1/2}}{\rho u_e} \int_0^{\eta_\infty} \phi'(\xi, \eta) d\eta \quad \dots [4.6.8]$$

$$\theta = -\frac{(2\xi)^{1/2}}{\rho u_e} \int_0^{\eta_\infty} \phi'(\xi, \eta) [1 + \phi'(\xi, \eta)] d\eta \quad \dots [4.6.9]$$

## CHAPTER 5

### COMPUTER SOLUTION OF THE MOMENTUM EQUATION

#### 5.1 Introduction

If a generalised computer programme is developed and proved with respect to reliable experimental results, it becomes possible to calculate such parameters as boundary layer thickness, skin friction and heat transfer. Such a programme could be simple, e.g. for the case of a flat plate, and may be made more versatile to include transition, pressure gradient, roughness, transpiration, heat transfer, compressible flow, axisymmetric flow and three dimensional flow.

A finite difference technique, rather than a finite element method, is used in this dissertation to solve equation [4.6.1]. Because of the well known stability problems experienced with explicit methods, an implicit scheme was selected, with its attendant requirement to linearise the equation. In this chapter the linearisation and matrix solution are discussed.

#### 5.2 Linearisation of the Momentum Equation

Linearisation is achieved by allowing those terms which make the equation non-linear to be set to values known from a previous station or iteration. Clearly an iterative process is necessary, in which the latest value approaches that of the previous iteration. When linearising, care must be taken to retain at least one term of each derivative. The subscript  $o$  here denotes a previously known quantity :-

$$[(1 + \epsilon^+) \phi'']' + (\phi_0 + \eta) \phi'' - \phi' B(\phi_0' + 2) = 2\epsilon[(\phi_0' + 1) \frac{\delta \phi'}{\delta \xi} - \phi_0'' \frac{\delta \phi}{\delta \xi}] \dots\dots [5.2.1]$$

### 5.3 The Finite Difference Grid

Cebeci (1974) suggested a grid in which the step lengths normal to the wall were very small, near the wall, ie. in the region of rapid change in velocity, and increasing in size towards the edge. He proposed a grid of the form :-

$$\Delta \eta_i = K \cdot \Delta \eta_{i-1} \dots\dots [5.3.1]$$

$\Delta \eta_i$  is the general step length, and for  $K > 1$ , this step length is larger than the previous step. If the first step length is  $h$ , the distance to the  $i^{\text{th}}$  grid point is :-

$$\eta_i = h \frac{K^{i-1} - 1}{K - 1} \dots\dots [5.3.2]$$

The exponential growth in grid spacing ties up with the turbulent velocity profile.

An alternate expression for the  $i^{\text{th}}$  grid spacing is :-

$$\Delta \eta_i = K^{i-1} h \dots\dots [5.3.3]$$

The programme was written in a manner to allow variable step length in the streamwise direction. This may be a completely random step length. Because the programme "marches" in the  $x$ -direction when solving, the grid points in this direction were called stations, and are denoted by, e.g. the  $n^{\text{th}}$  station.

The ordinate was divided into points or steps, denoted by  $i$ . Note that the first  $\eta$  value, ie. at the wall, is point  $i = 1$ , therefore  $\eta_1 = 0$ .



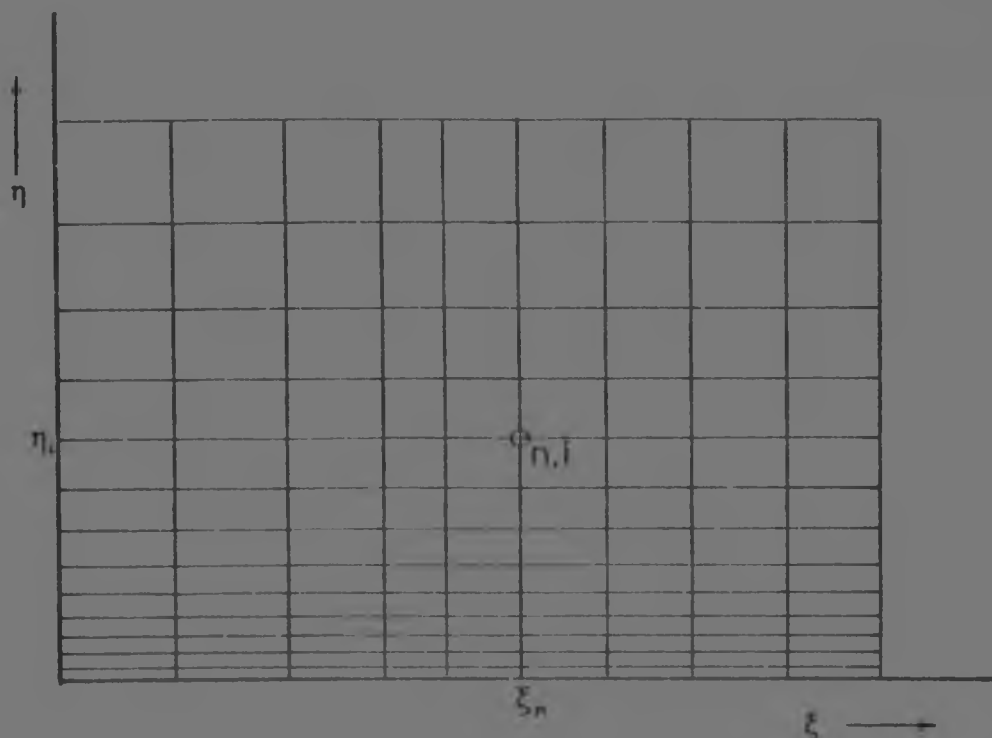


FIG 5.1 : THE FINITE DIFFERENCE GRID FOR  
 $k = 1,2$  AND  $h = 1,5$

5.4 The Momentum Equation in Finite Difference Form

Equation [5.2.1] is given with the subscripts as used in the programme :-

$$\begin{aligned}
 & [(1 + \epsilon_{n,i_0}^+) \phi_{n,i}'' ]' + (\phi_{n,i_0} + \eta_i) \phi_{n,i}'' - \phi_{n,i}' \cdot \epsilon_n (\phi_{n,i_0}' + 2) \\
 & = 2\xi_n [ (\phi_{n,i_0}' + 1) \frac{\delta \phi'}{\delta \xi} \Big|_{n,i} - \phi_{n,i_0}'' \frac{\delta \phi}{\delta \xi} \Big|_{n,i} ] \dots [5.4.1]
 \end{aligned}$$

By taking the derivative of the first term at this stage  $d\epsilon^+/dy$  would be introduced, which may be determined numerically, but with the attendant error inherent in any such calculation, with the added difficulty that  $\epsilon$  was not a continuous function.  $\phi_{n,i}'''$  would also appear. The derivative was solved implicitly in the final equation by treating the term as a separate function, say  $g$ . This is why the term  $(1 + \epsilon^+)_{n,i-1_0}$  and at  $(n,i + 1)$  feature in the following equation. Values of  $\phi$  are known at  $n-1$ ,  $n-2$  and  $0$ . All these terms were kept on the right hand side, while

those in  $n, i$  were retained on the left hand side. The capital letters denote constants based on step size, and are defined in Appendix A2.

$$\begin{aligned}
 & \phi_{n, i-2} \{A.X1. (1 + \epsilon_0^+)_{n, i-1}\} \\
 & + \phi_{n, i-1} \{A.X2. (1 + \epsilon_0^+)_{n, i-1} + B.D. (1 + \epsilon_0^+)_{n, i} + (\phi_{0, n, i} + \eta_i) D \\
 & \quad - B.A. (\phi'_{0, n, i} + 2) - 2\xi_n (\phi'_{0, n, i} + 1).A.N\} \\
 & + \phi_i \{A.X3. (1 + \epsilon_0^+)_{n, i-1} + B.E. (1 + \epsilon_0^+)_{n, i} + C.X4. (1 + \epsilon_0^+)_{n, i+1} \\
 & \quad + E. (\phi_{0, n, i} + \eta_i) - B.B. (\phi'_{0, n, i} + 2) - 2\xi_n (\phi'_{0, n, i} + 1)N.B \\
 & \quad + 2\xi_n \phi''_{0, n, i}.N\} \\
 & + \phi_{i+1} \{B.F. (1 + \epsilon_0^+)_{n, i} + C.X5. (1 + \epsilon_0^+)_{n, i+1} + F. (\phi_{0, n, i} + \eta_i) \\
 & \quad - B.C. (\phi'_{0, n, i} + 2) - 2\xi_n (\phi'_{0, n, i} + 1).N.C\} \\
 & + \phi_{i+2} \{C.X6. (1 + \epsilon_0^+)_{n, i+1}\} \\
 & = 2\xi_n \{ (\phi'_{0, n, i} + 1) [L.\phi'_{n-2, i} + M\phi'_{n-1, i}] - \phi''_{0, n, i} [L.\phi_{n-2, i} + M\phi_{n-1, i}] \} \\
 & \dots\dots\dots [5.4.2]
 \end{aligned}$$

Lumping the coefficients into single constants reduces [5.4.2] to :-

$$\phi_{n, i-2}.A_i + \phi_{n, i-1}.B_i + \phi_{n, i}.C_i + \phi_{n, i+1}.D_i + \phi_{n, i+2}.E_i = F_i$$

\dots\dots\dots [5.4.3]

Remembering that  $i = 1$  denoted  $\eta = 0$ , a system of linear simultaneous equations could be written starting with  $i = 3$ . If  $N$  is the number of  $\eta$  steps, a total of  $N-4$  equations can



From Fig 5.2 it is seen that as  $\eta \rightarrow \eta_\infty$ ,  $f'(\xi, \eta) \rightarrow 1$  and therefore  $\phi'(\xi, \eta) \rightarrow 0$ .

Equation [4.6.4] states that

$$\lim_{\eta \rightarrow \eta_\infty} \phi'(\xi, \eta) = 0 \quad \dots\dots\dots [5.4.6]$$

If  $N-2$ ,  $N-1$  and  $N$  are in this region, then using a simple finite difference approximation :-

$$\phi'_{N-1} = \frac{\phi_N - \phi_{N-1}}{\eta_N - \eta_{N-1}} = 0 \quad \dots\dots\dots [5.4.7]$$

$$\therefore \phi_N = \phi_{N-1} \quad \dots\dots\dots [5.4.8]$$

Similarly :-

$$\phi'_{N-1} = \frac{\phi_N - \phi_{N-2}}{\eta_N - \eta_{N-2}} = 0 \quad \dots\dots\dots [5.4.9]$$

$$\therefore \phi_N = \phi_{N-1} = \phi_{N-2} \quad \dots\dots\dots [5.4.10]$$

The  $N-3$  equation is :-

$$A_{N-3} \phi_{n, N-3} + B_{N-3} \phi_{n, N-4} + C_{N-3} \phi_{n, N-5} + \phi_{n, N-2} (D_{N-3} + E_{N-3}) = F_{N-3} \quad \dots\dots\dots [5.4.11]$$

and the  $N-2$  equation (last) is :-

$$A_{N-2} \phi_{n, N-2} + B_{N-2} \phi_{n, N-3} + \phi_{n, N-2} (C_{N-2} + D_{N-2} + E_{N-2}) = F_{N-2} \quad \dots\dots\dots [5.4.12]$$

The general equation may be stated :-

$$[A]\{\phi\} = \{B\} \quad (n=3 \text{ to } (n=N-2) \quad \dots\dots\dots [5.4.13]$$

### 5.5 The Finite Difference Molecule

The general finite difference molecule is shown in Fig 5.3

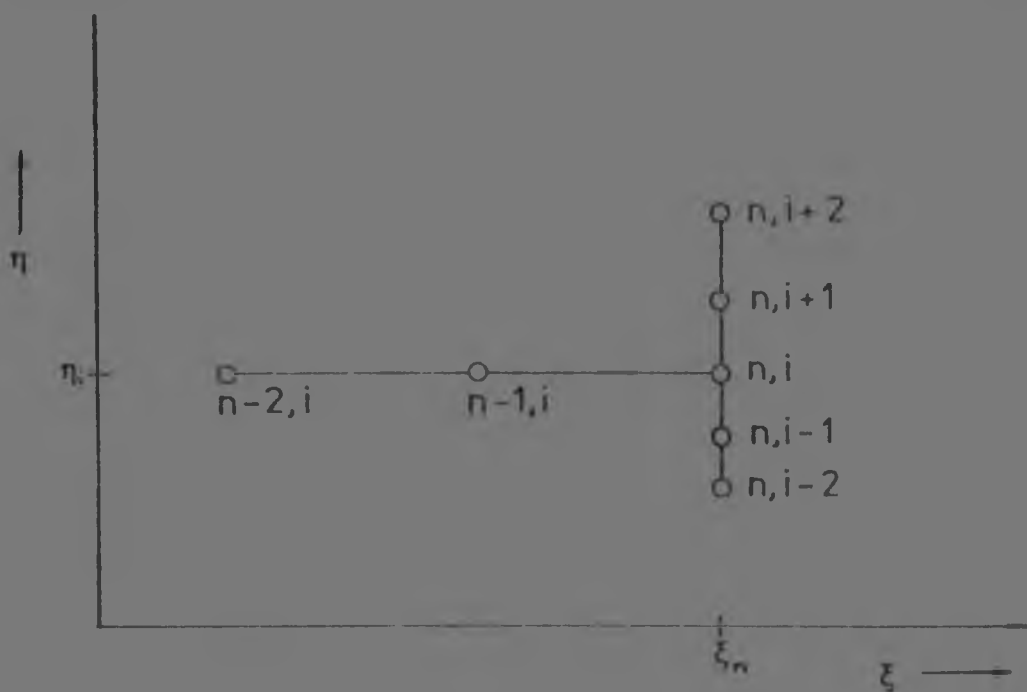


FIG 5.3 : THE FINITE DIFFERENCE MOLECULE

To solve for  $\phi$  at station  $n$ , it is necessary to know the values of  $\phi$ ,  $\phi'$  and  $\phi''$  at stations  $n-2$  and  $n-1$ . Two profiles were therefore required to start the solution at the third station. The programme then marched to  $n+1$  and used  $n-1$  and  $n$  to solve for  $\phi$  at  $n+1$ , and so on. For the first solution at say  $n$ , the values on the right hand side, denoted by  $o$ , were substituted by the values at  $n-1$ . Similarly in generating  $\epsilon$ . The iterative process was based on the convergence of  $\phi''_{n,1}$  and  $\phi''_{0,n,1}$  and  $\delta''_n$  and  $\delta''_0$ .

To economise on core storage on the computer, only four vectors for  $\phi$ ,  $\phi'$ ,  $\phi''$  and  $\epsilon'$  were used. Solutions at  $n-2$ ,  $n-1$  and  $n$  were stored in the first three vectors; the present iteration,  $o$ , in the fourth. The first three rotated cyclicly. If the vectors are denoted by  $V_1$ ,  $V_2$  and  $V_3$ , the solutions were stored :-

$n=1, V_1$

$n=2, V_2$

$0, V_3$

Once convergence was achieved,  $V_3 \rightarrow V_2$  and print/plot  $V_1$  could be overwritten and the solution march to  $n=4$

$n=2, V_2$

$n=3, V_3$

$0, V_4$

After convergence,  $V_4 \rightarrow V_3$ , print/plot, and delete  $V_2$

$n=3, V_3$

$n=4, V_4$

$0, V_5$

After convergence,  $V_5 \rightarrow V_4$ , print/plot, etc.

A tremendous saving was possible when storing matrix [A] because it was banded. It had two upper and two lower co-diagonals. Typically, the banded matrix required 500 storage locations, compared with 10 000 for full square matrix storage. In double precision, this is equivalent to 80 k bytes.

$n=1, V_1$   
 $n=2, V_2$   
 $0, V_4$

Once convergence was achieved,  $V_4 \rightarrow V_3$  and print/plot  $V_1$  could be overwritten and the solution march to  $n=4$

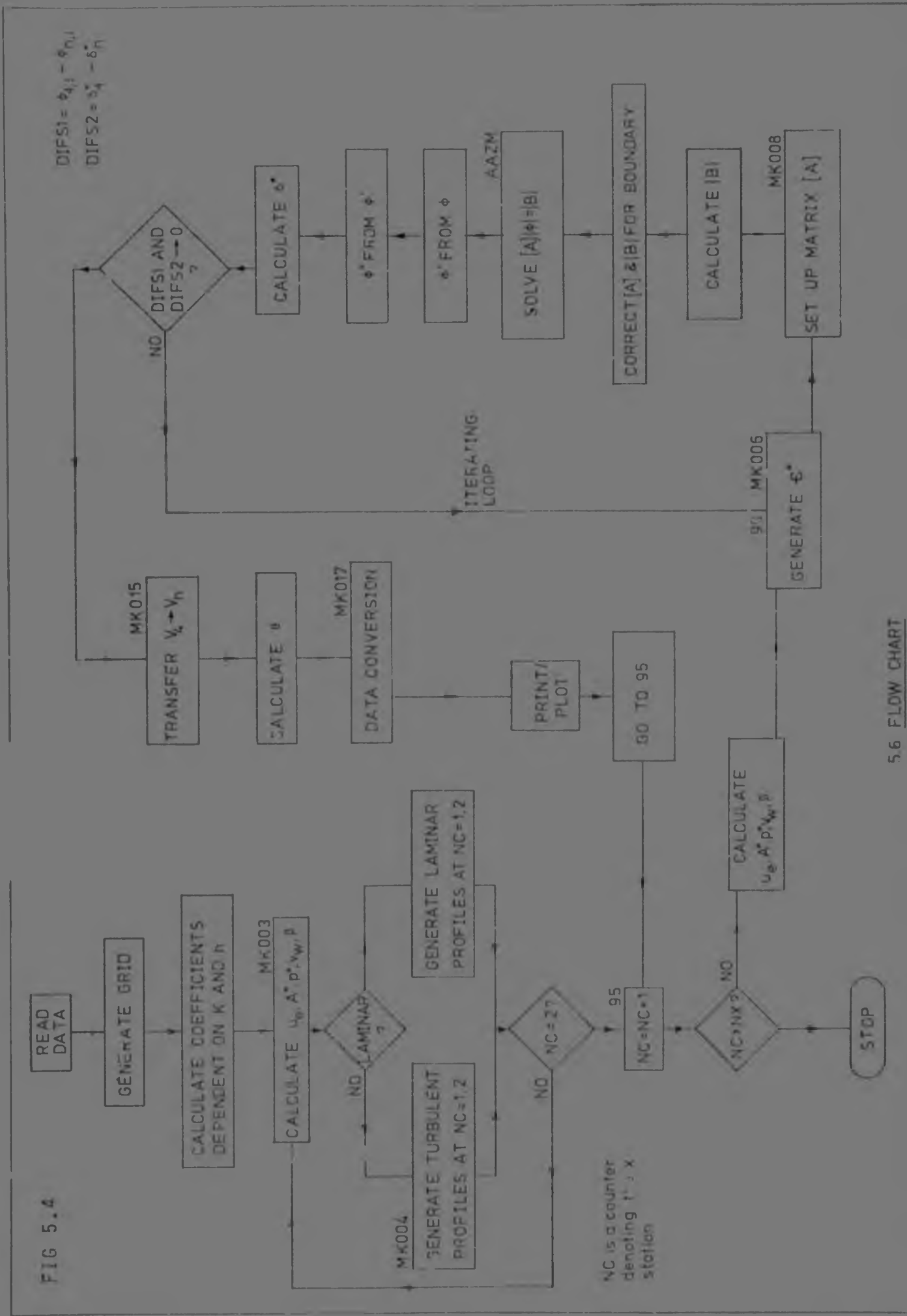
$n=2, V_2$   
 $n=3, V_3$   
 $0, V_4$

After convergence,  $V_4 \rightarrow V_1$ , print/plot, and delete  $V_2$

$n=3, V_3$   
 $n=4, V_1$   
 $0, V_4$

After convergence,  $V_4 \rightarrow V_2$ , print/plot, etc.

A tremendous saving was possible when storing matrix [A] because it was banded. It had two upper and two lower co-diagonals. Typically, the banded matrix required 500 storage locations, compared with 10 000 for full square matrix storage. In double precision, this is equivalent to 80 k bytes.





5.7 Discussion of Various Features of the Programme

5.7.1 Initial profiles

As stated above, profiles at  $n=1$  and  $n=2$  were required to start the solution. These were generated in MK001L (laminar) or MK004 (turbulent) - see Fig 5.5. Distance along the abscissa was used as a criterion for laminar or turbulent flow. Various stability techniques, critical Reynolds numbers and shape factors have been suggested (Cebeci 1974). However, these are often limited, and not sufficiently accurate, e.g. critical Reynolds number is for flat plate flow with zero pressure gradient, and spans a large range ( $3,5 \times 10^5$  to  $10^6$ ).

If  $\epsilon=0$ , the solution is laminar, and from a given  $\xi$  station,  $\epsilon>0$  and the programme entered the transition and hence turbulent solutions.

The laminar profile was generated from an equation given by Grimson (1971) :-

$$\frac{u}{U_\infty} = f(y/\delta) = f(\bar{y}) \dots\dots\dots [5.7.1]$$

If a fourth order polynomial is used, a pressure gradient may be included :-

$$\frac{u}{U_\infty} = A_0 + A_1\bar{y} + A_2\bar{y}^2 + A_3\bar{y}^3 + A_4\bar{y}^4 \dots\dots\dots [5.7.2]$$

A pressure gradient related shape factor was introduced :-

$$\Lambda = \frac{\delta^2}{\nu} \cdot \frac{du}{dx} \dots\dots\dots [5.7.3]$$

$$\frac{u}{U_\infty} = \bar{y}(2 - 2\bar{y}^2 + \bar{y}^3) + \frac{\Lambda}{6} \bar{y}(1 - 3\bar{y} + 3\bar{y}^2 - \bar{y}^3) \dots [5.7.4]$$

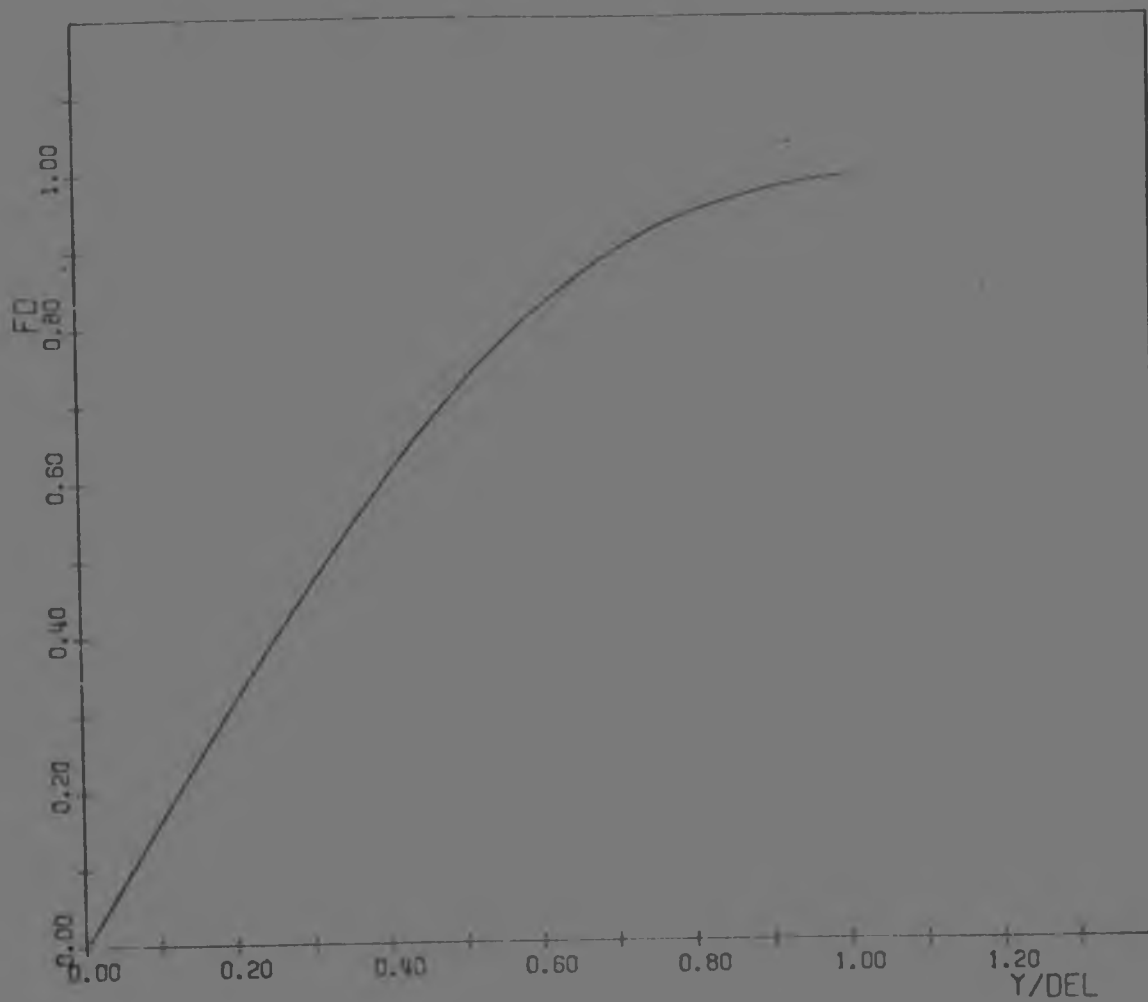


FIG 5.5 LAMINAR INPUT PROFILE

The turbulent profile was generated using Van Driest's mixing length theory. One difficulty here was that the edge of the boundary layer was over-estimated, and pressure gradient was not included. However, the programme would obtain the correct profile at  $n=6$  after some iteration. In developing this subroutine, it was found that the results were significantly improved if the unequal, exponential type  $\eta$  grid was used. See Fig 5.6 for equal step length, and Fig 5.7 for exponential grid, using the well known semi-log plot. Fig 5.7 is reproduced on simple non-dimensionalised axes in Fig 5.8. The dashed lines on Fig 5.7 are for :-

$$u^+ = y^+ \quad \dots\dots\dots [5.7.5]$$

$$\text{and } u^+ = 5,757 \log y^+ + 5,24 \quad \dots\dots\dots [5.7.6]$$

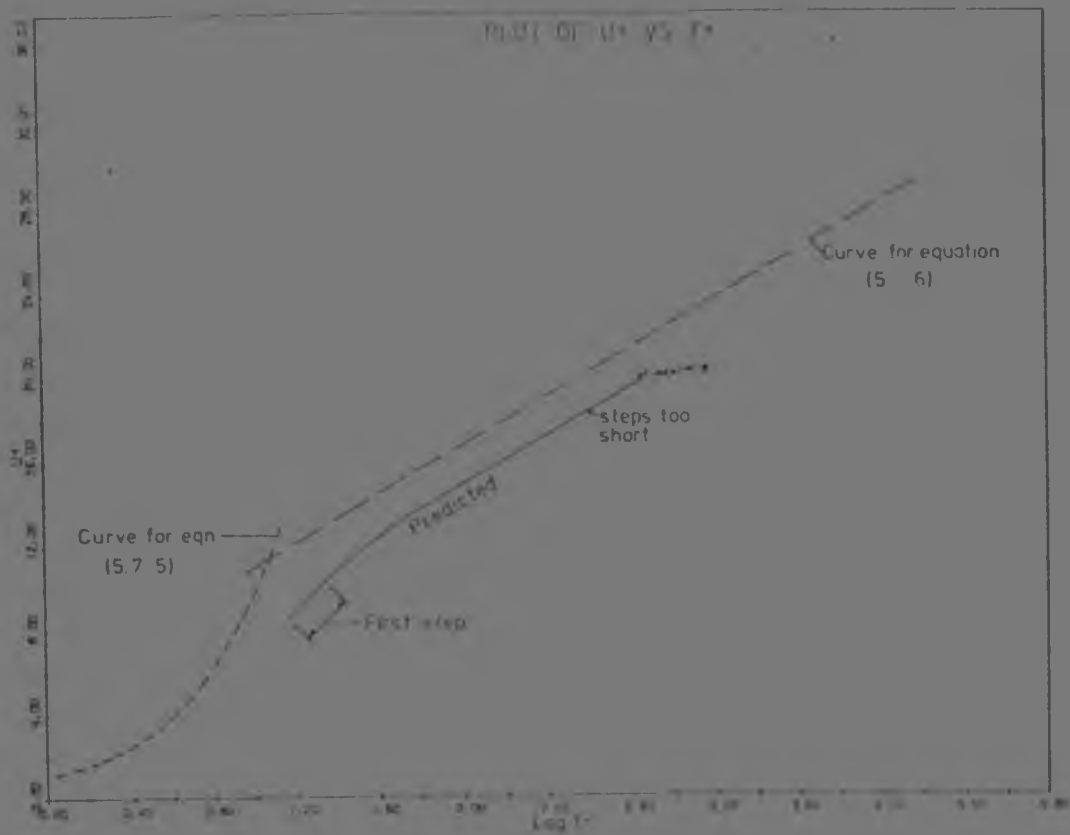


FIG 5.6 : TURBULENT PROFILE GENERATED WITH EQUAL STEP LENGTH

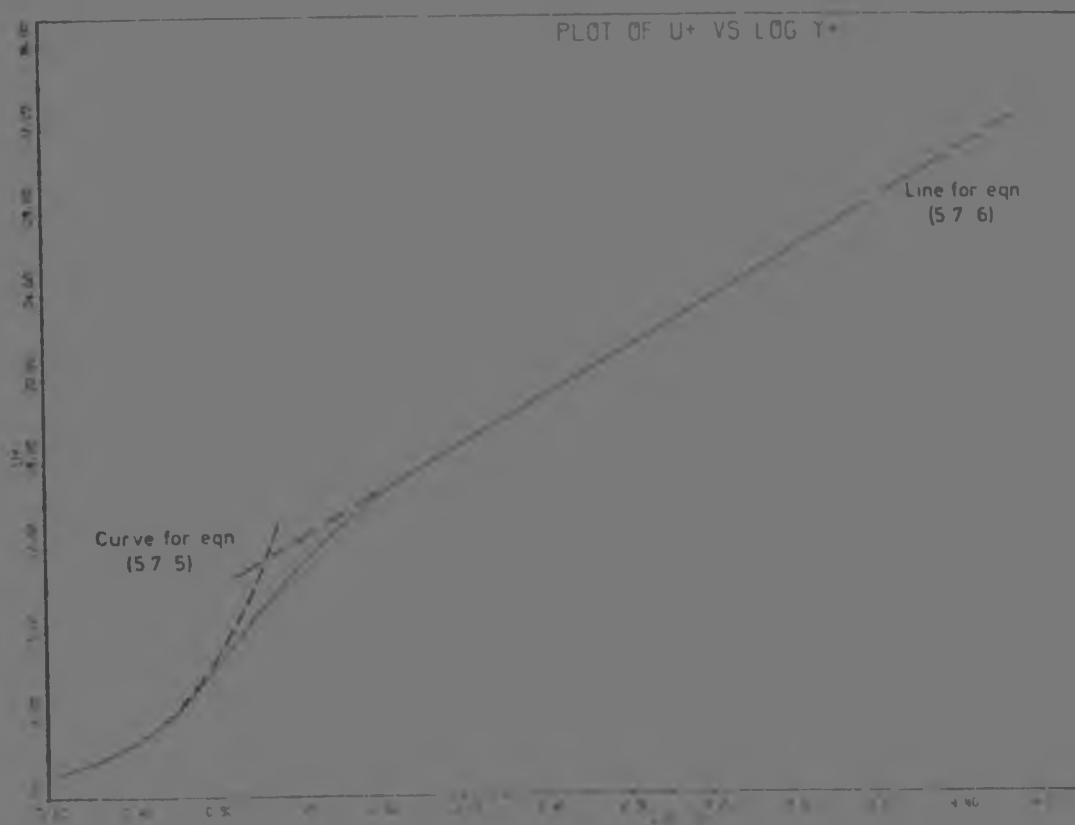


FIG 5.7 : TURBULENT PROFILE GENERATED WITH UNEQUAL STEP LENGTH

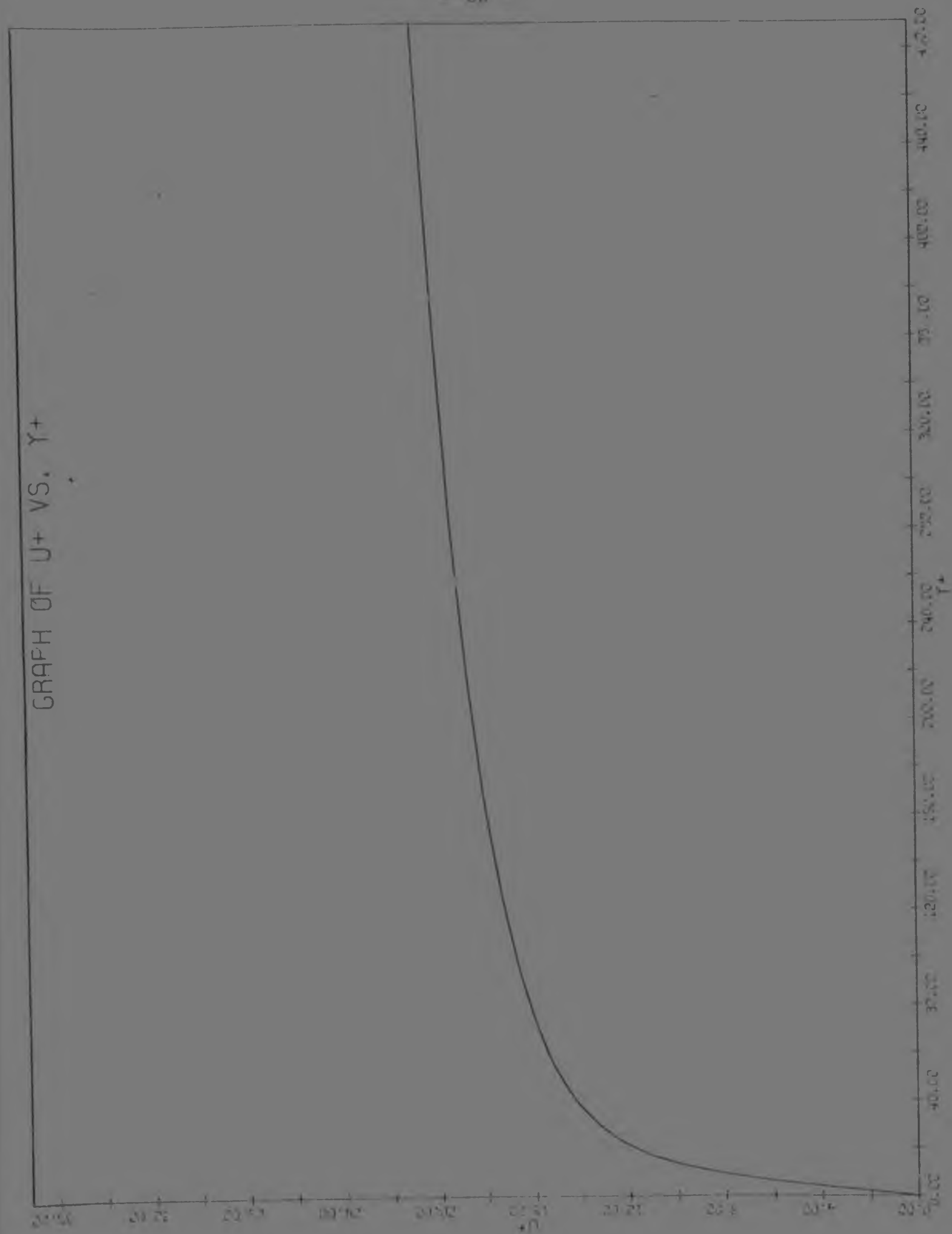


FIG 5.8 : TURBULENT PROFILE GENERATED WITH UNEQUAL STEP LENGTH

Equation [3.3.33], with  $v_w^+ = 0$ , simplified to that derived by Van Driest (1956) :-

$$\frac{du^+}{dy^+} = \frac{2dy^+}{1 + \sqrt{1 + 4\kappa^2 y^{+2} [1 - \exp(-y^+/A^+)]^2}} \quad \dots [5.7.7]$$

by integrating this expression,  $u_i^+ = \int_0^{y_i^+} f(y^+) dy^+$ . This was performed numerically in subroutine MK004.  $u_i^+$  had to be converted to  $u/u_e$ , ie.  $v_w^+$  was required.  $c_f$  was estimated in subroutine TAUF from the simplified equation

$$c_f = \frac{0,0595}{Re_x^{1/2}} \quad \dots [5.7.8]$$

In both laminar and turbulent profile generators, it was necessary to convert  $u/u_e (=f')$  to  $\phi'$ , and to calculate  $\phi$  and  $\phi''$ . The edge of the boundary layer was taken to be where  $\phi' < 0,0001$ .

To obtain  $\delta^*$  and  $\theta$ , the profile had to be integrated. This was done by using the trapezoidal rule, which was considered sufficiently accurate because the step length was so small. Note that the experimental profiles were integrated in a similar manner, and for a very coarse grid (20 points compared with 100 for the computer programme) the difference in  $\theta$  when using Simpson's rule was 1%. A typical equation appears in Appendix F1.

As the injection increased, the shear stress at the wall tends to zero, or even a negative value. However, reverse flow is not thought to occur. The programme was not designed to handle these possibilities. If  $\phi''$  was obtained from  $\phi$  using Newtons Divided Difference technique at the wall, negative shear was indicated at values of  $F \geq 0,005$ . This is in accordance with experiment, but resulted in difficulties when calculating  $u_w^+$ , and hence also  $v_w^+$  and  $p^+$ . The programme was written for a continuous velocity

profile and the condition of no-slip at the wall. Note too that the derivation of  $A'$  was for very small injection rates, and  $y_0^+ = 11,8$  was the value for unblown flat plate flow.

By using a reasonably large initial  $\eta$  step (0,06 compared with values as low as 0,001 suggested by Cebeci (1970a), and calculating  $\phi''$  from  $\phi'$ , the programme remained stable up to values of  $F = 0,017$ .

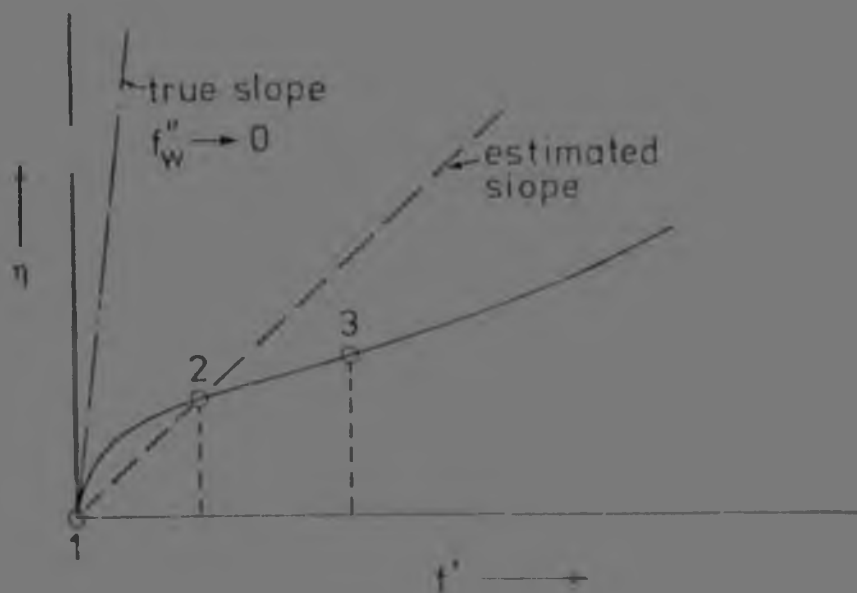


FIG 5.9 :  $u$ -VELOCITY PROFILE NEAR THE WALL

The numbers shown in Fig 5.9 denote  $\eta$  grid points. If 2 is outside the point of inflexion, the programme would effectively be unaware of this behaviour.  $\phi''$  would be predicted to be the dashed line (depending on the finite difference equation). This value is incorrect, but allows real values of  $u_w^+$ ,  $v_w^+$ ,  $p^+$  etc. to be calculated. Although these errors are large, the effect of this on the complete solution of [5.4.13] appeared to be very small.

Because this technique produced results which appeared very promising, the loss of accuracy in  $u_w^*$ , and hence meaningless  $c_f$  for injected flows, was considered unimportant in the present phase of research. As will be pointed out later in Chapter 7,  $u_w^*$  as characteristic velocity for injected boundary layers may well require reconsideration.

For zero pressure gradient flat plate flow,  $c_f$  predicted by the programme was very close to that experimentally measured and fitted by [5.7.8].

## CHAPTER 6

### EXPERIMENTAL APPARATUS AND PROCEDURE

#### 6.1 Introduction

The object of the experiments was to obtain a number of mean velocity profiles on a porous plate preceded by a solid plate, ie. to subject a fully developed turbulent profile to injection. The next step was to constrict the lower porous surface, to ascertain whether this made a measurable difference to the profiles. To this end an open circuit wind tunnel with rectangular working section was designed. Three different grades of porous plate were tested.

#### 6.2 The Wind Tunnel

##### 6.2.1 General

The open circuit wind tunnel consisted of a Sirrocco centrifugal fan (see Appendix B1 for the performance curve). Drive was provided by a Siemens A.C. induction motor rated at 22 kW, through three Fenner  $\beta$  type belts. By changing pulleys, the rotational speed of the fan could be varied, resulting in different flow rates and pressures. A flow damper at the fan inlet could be used to vary the flow rate, as could a by-pass at the beginning of the tunnel. This by-pass was only used at start-up (minimize starting load), because a power spectrum measured in the working section produced two additional spikes when the trap door was not fully closed (see 6.4).

The ducts leading to the Working Section were made up of a diffuser, settling chamber and contraction. Screens and honeycomb straightened the flow and decreased the turbulence intensity.





FIG 6.1 : GENERAL VIEW OF WIND TUNNEL SHOWING  
(1) THE SETTLING CHAMBER, (2) CONTRACTION,  
(3) ENTRY TO WORKING SECTION, (4) REAR OF  
INSTRUMENT PANEL (5) BETZ MANOMETER

#### 6 2.2 The working section

A rectangular working section 209 x 496 mm was built. The side walls were made of perspex, the roof and solid floor sections of aluminium. A boundary layer bleed was installed in the floor in the entry section. The plate preceding the porous section was 800 mm long, the leading edge milled to  $45^\circ$ . This configuration could be altered such that the initial solid plate was 150 mm long. Theoretically a laminar boundary layer would reach the porous plate, but tests were not done with the short plate installed.

A traversing rig (Disa type 55H01) was mounted on slides. The probe could be lowered through a slit in the roof of the working section. The distance of the probe from the wall was measured with a Mitutoyo Dial Gauge. This entire rig was mounted on a separate frame which was isolated from the tunnel to prevent vibration being transmitted to the probe. The traversing rig allowed 4 degrees of freedom.

Fig 6.2 shows the working section and the traversing rig. In the next photograph, Fig 6.3, a close-up of the working section interior is given. In this figure the hot wire probe can be seen mounted above the porous plate. Note that the short plate configuration is depicted.

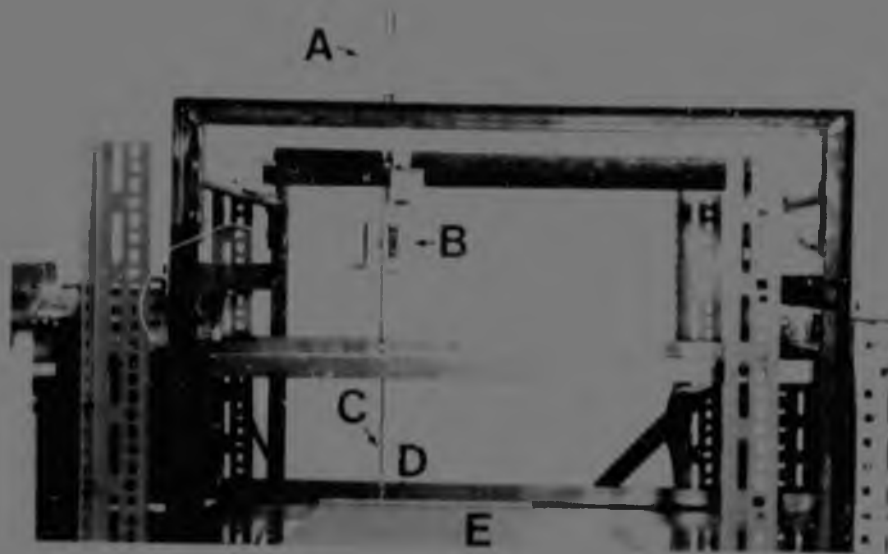


FIG 6.2 : THE WORKING SECTION AND THE TRAVERSING RIG :  
(A) DIAL GAUGE, (B) DISA TRAVERSING MECHANISM,  
(C) PROBE HOLDER, (D) HOT WIRE PROBE,  
(E) POROUS PLATE

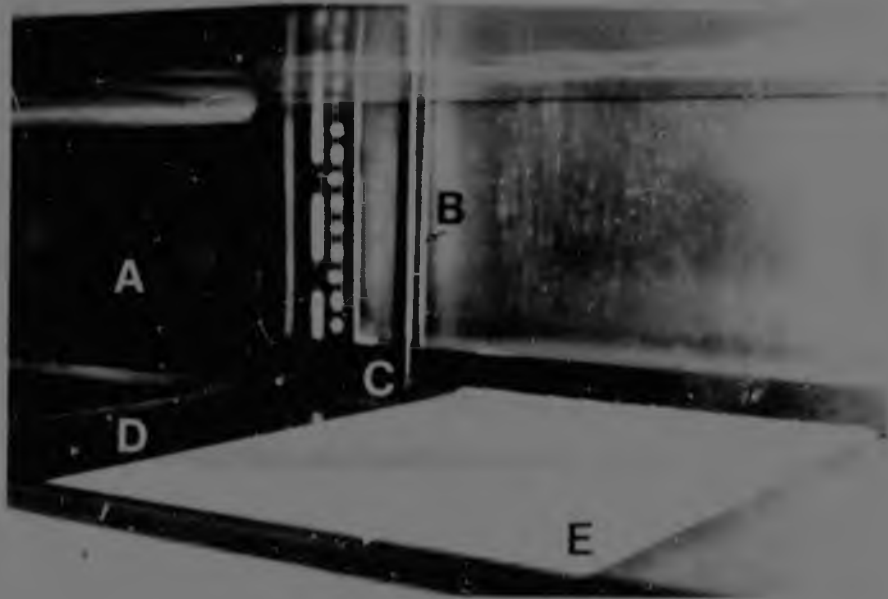


FIG 6.3 : INSIDE OF THE WORKING SECTION : (A) ENTRY TO  
WORKING SECTION, (B) PROBE HOLDER, (C) HOT WIRE  
PROBE, (D) SHORT LEADING PLATE, (E) POROUS PLATE

### 6.2.3 Transpiration flow supply

The transpiration or secondary flow was drawn from the laboratory compressed air supply. An Ingersoll-Rand compressor fed two receivers, from which a 50 mm i.d. pipe ducted the air to a pressure regulator (Norgren type AG20). The flow rate was measured with an orifice plate, preceded by an entry length of 30 diameters. The orifice plate was calibrated according to British Standards 1042. A fibre-type filter, supplied by Vokes, was placed in the plenum chamber below the porous plate. The filter removed particles down to 5 microns, as well as grease. The front cover of the plenum chamber had to be removed when changing porous plates. The height of the porous plate was adjusted with screws inside the chamber. The plates were made of sintered bronze by Sintered Products Ltd. of Nottingham, England. Each plate was 6,35 mm thick. The front and rear edges were milled to  $45^\circ$ . Grades E (highly porous) C and A were tested.



FIG 6.4 :  
THE WORKING SECTION  
AND THE PLENUM CHAMBER  
SHOWN WITH ITS FRONT  
COVER IN PLACE

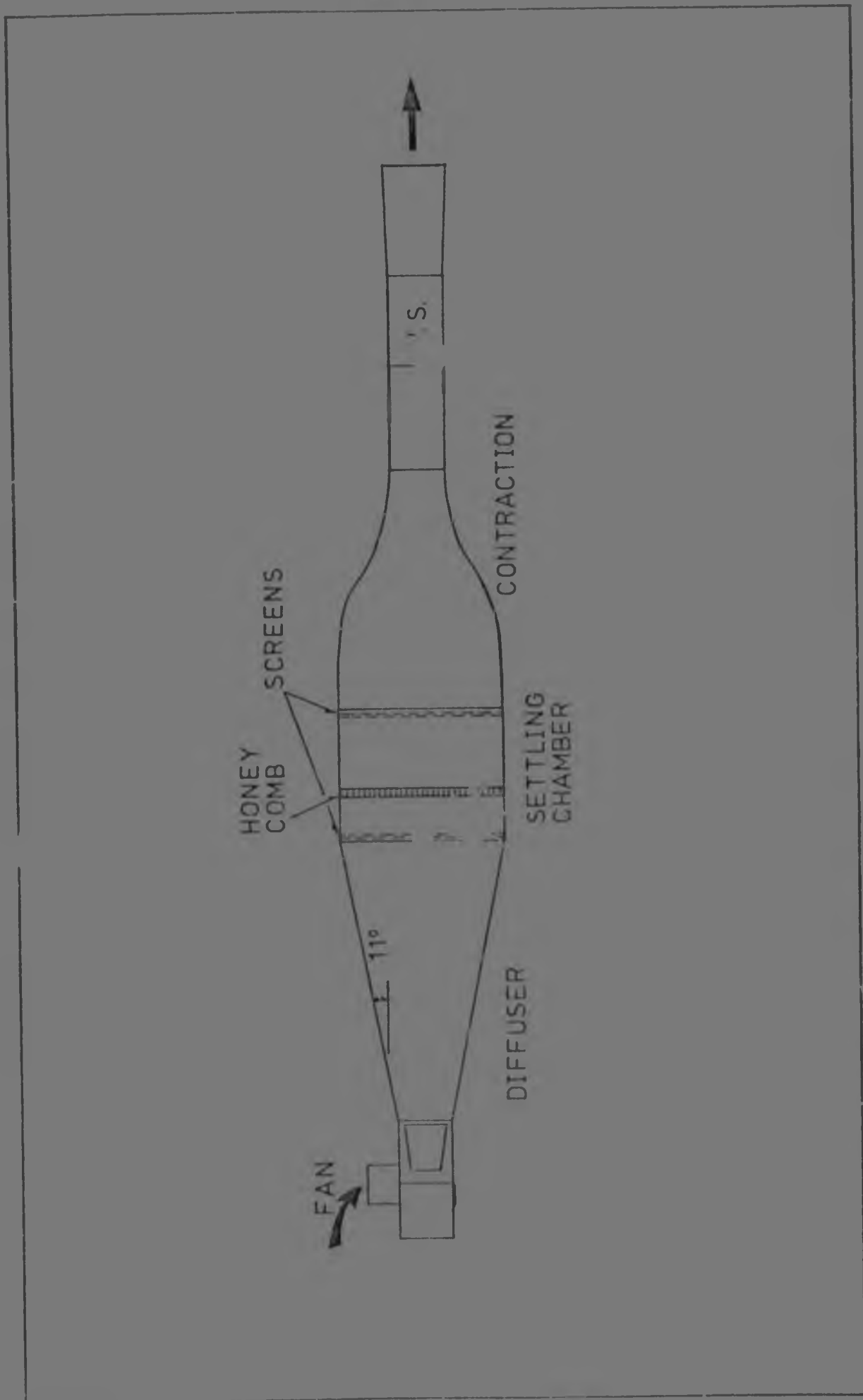


FIG 6.5 : PLAN VIEW OF THE WIND TUNNEL

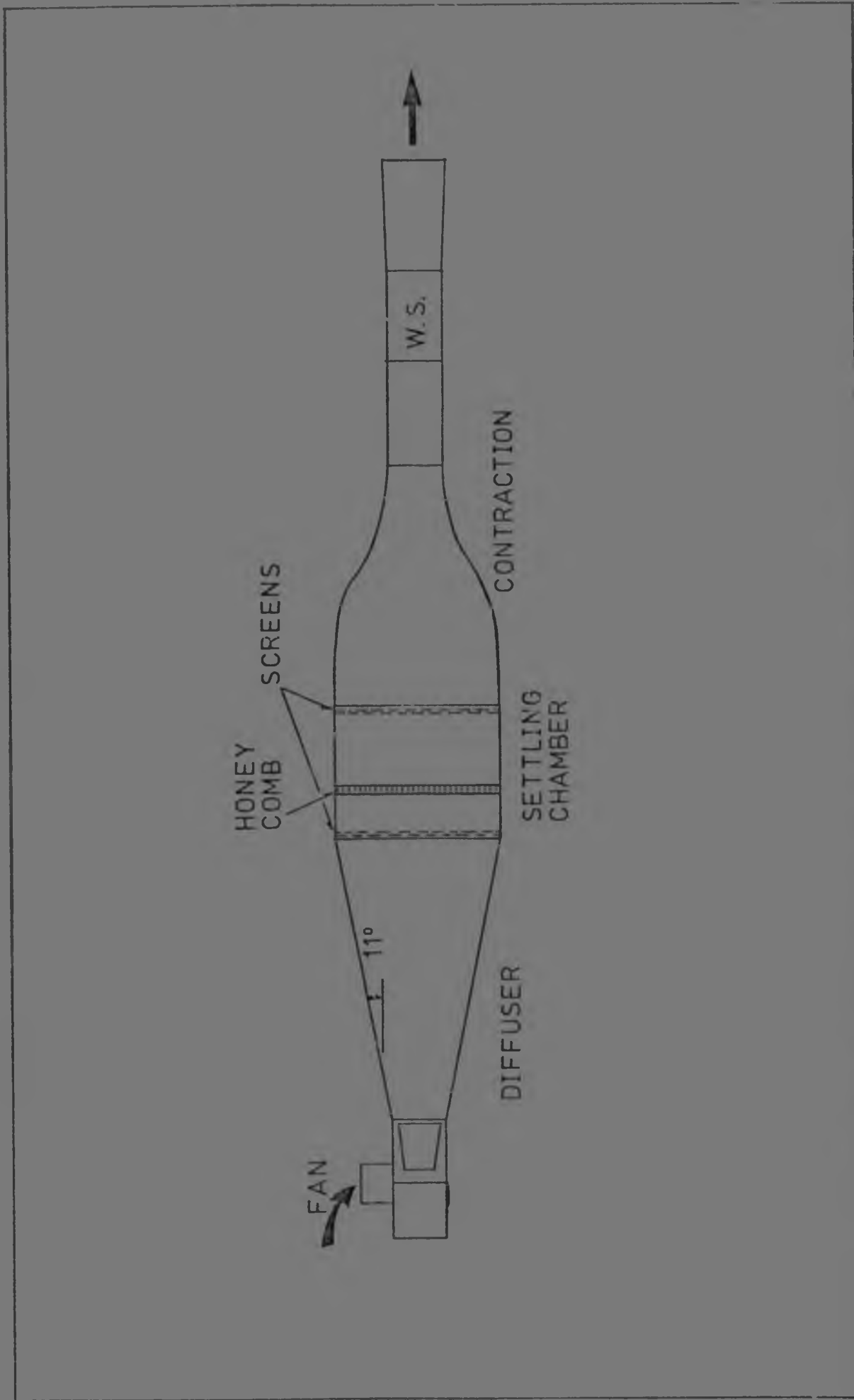


FIG 6.5 : PLAN VIEW OF THE WIND TUNNEL

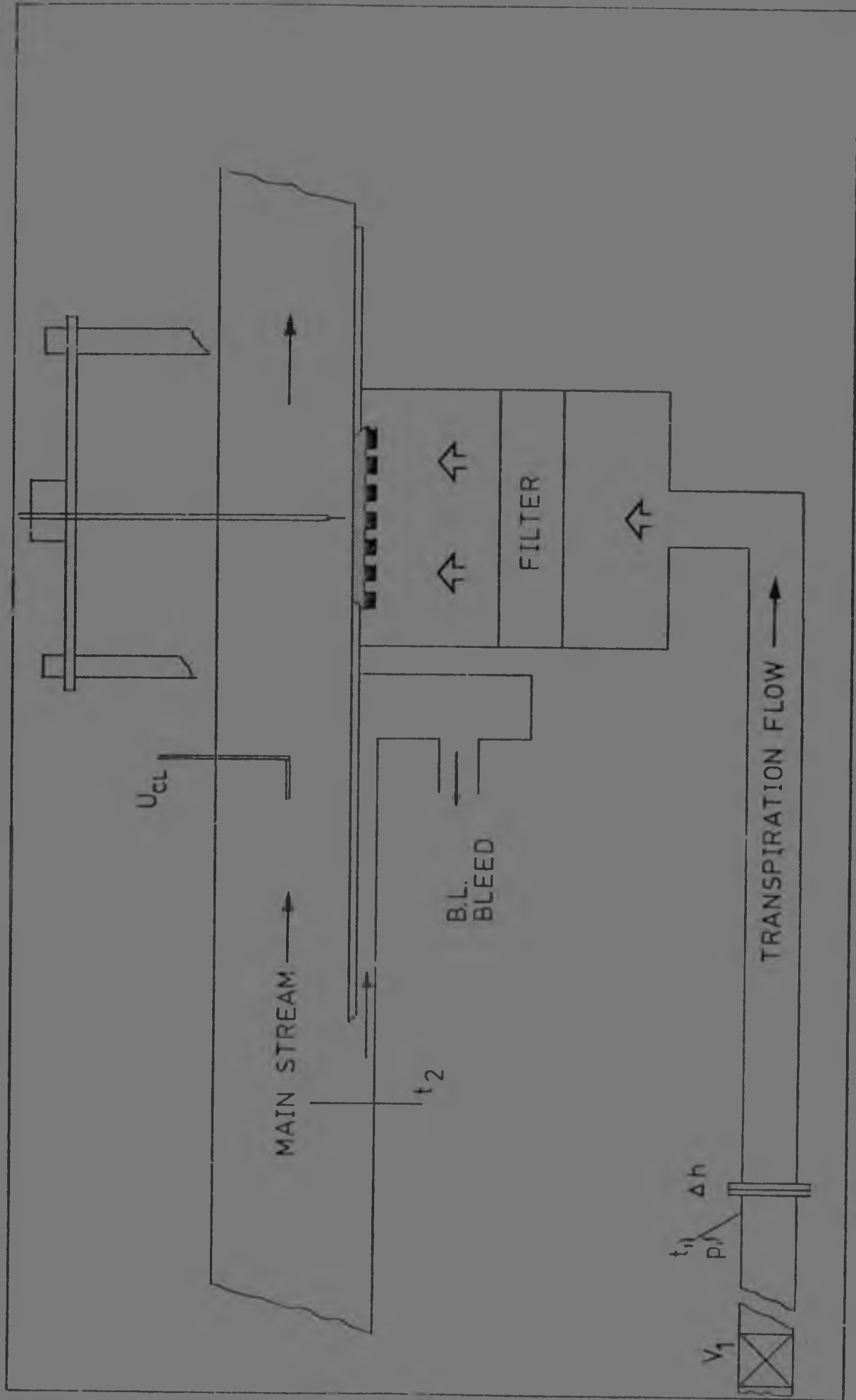


FIG 6.6 : SIDE VIEW OF THE WORKING SECTION SHOWING THE TRANSDUCERS

### 6.3 Hot-wire Anemometry

Velocity traverses were obtained using a Disa type 55D01 constant temperature anemometer. A single sensor, gold-plated wire probe (Disa type 55F01) was inserted in a probe holder with a protection pin attached. By adjusting this pin, it was possible to touch the pin on the working section floor, and treat that as the first point of traverse. By reading the dial gauge at that point, the height above the plate was determined, and subsequent increases in distance from the floor known to 0,01 mm.

The probe holder was held by the Disa traversing mechanism. This rig had a range of 100 mm, and was operated manually through a pinion drive, for slipless traversing. For increased accuracy, the dial gauge was mounted above the rig.

Apart from mean velocity, the turbulence intensity and power spectrum of the Working Section mainstream was measured. For this, the probe was mounted at a x-station, in the centre of the Working Section. Because the anemometer output is non-linear, it required processing. This was done with a Type 55D10 Lineariser. The linearised turbulent signal was integrated on a Type 55D35 R.M.S. Voltmeter, which gave  $\sqrt{v'^2}$ . This unit, also made by Disa, had a voltage range of 0 - 300V. The integrator time constant could be set to 0,3, 1, 3, 10 and 30 seconds.

The calibration of the probes is discussed in Appendix C3.

To obtain the mean velocity in the boundary layer, a simple R.C. circuit was included in the circuit.

A Doric type DS100 digital voltmeter was used to measure the output from the anemometer.

### 6.4 The Energy Spectrum

Turbulent flow is so irregular that early investigators gave up hope of capturing the flow in universal expressions. In its very randomness lies the opportunity of using statistical methods to characterise the motion. Thus the time based mean velocity and thickness could be defined. The hot wire anemometer with its high frequency response, and oscilloscopes to record the output, made more sophisticated investigations possible. Thus the correlation function and energy spectrum are suitable means of presenting data.

$E_1(n)$  is defined such that  $E_1(n)dn$  is the contribution to  $\overline{u'^2}$  of the frequencies in the band  $n$  to  $n + dn$ .

$$\overline{u'^2} = \int_0^{\infty} E_1(n) dn \quad \dots\dots\dots [6.4.1]$$

As the frequency increases,  $E_1(n)$  decreases rapidly because of the energy dissipation by viscous action.

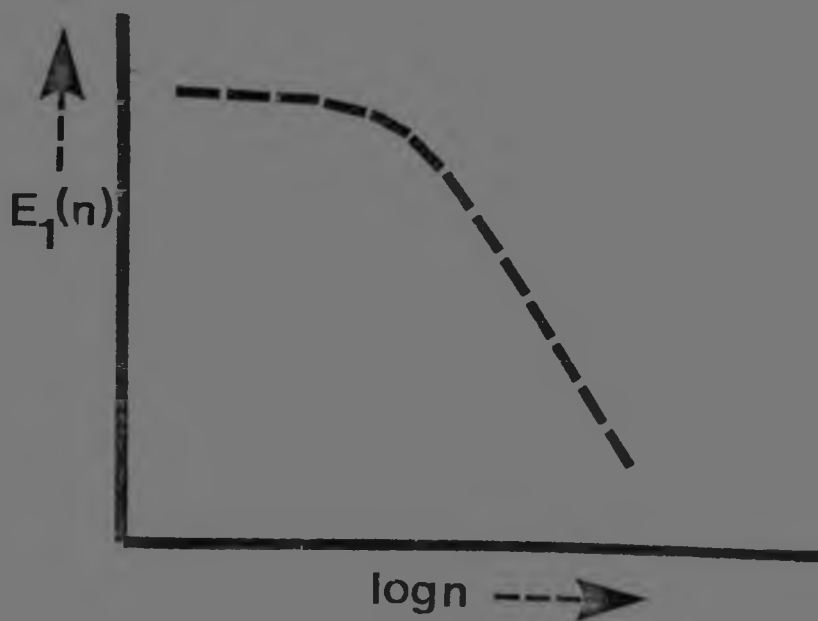


FIG 6.7 : SKETCH OF AN ENERGY SPECTRUM FOR TURBULENT FLOW



From Fig 6.7,  $E_1(n)$  is clearly greatest at low frequencies. The energy creating the vortices is drawn from the mainstream, and eventually dissipated through the action of viscosity.

Obtaining a plot of  $E_1(n)$  vs  $n$  required the turbulent oscillations to be represented by a Fourier series, e.g.

$$f(n) = A_0 + A_1 \cos \frac{\pi n}{L} + A_3 \cos \frac{2\pi n}{L} + \dots$$

and then plotting  $A_n$  vs  $n$ . Equipment employing a band pass filter could be used to generate the spectrum. For this work, a Hewlett-Packard Spectrum Analyser, model 3580A, was used in conjunction with the linearised output from the single sensor probe.

This unit produced a display on a screen, but a graphical output was obtained as well. The  $x$ -co-ordinate signal (frequency) was converted by a Hewlett-Packard Log Voltmeter/Converter, type 7562. Thus a log-log plot was produced directly on a Riken Denshi X-Y Recorder.

#### 6.5 Experimental Procedure

It was found that the anemometer (in particular, the lineariser) was prone to drift with temperature. Before an experiment could be started, the lineariser was allowed a warm-up period of two hours (suggested by manufacturer). A linearised calibration curve was obtained, and the hot wire transducer transferred to the working section. It was essential that the anemometer was not switched off between calibration and test.

At the start of a test, valve V1 (see Fig 6.6) was opened, and the desired transpiration flow set.

After opening the by-pass, the fan was started. Once running, the trap door was closed, and the baffle at the fan inlet set for the required air velocity in the working

section. This was checked with a pitot tube in the entry length. The free stream temperature was monitored, and once steady it was assumed that the wind tunnel had reached operating conditions. Ten to fifteen minutes was usually required for steady state to be attained.

The hot wire transducer was then positioned at X1,  $y=1$ , and the anemometer switched on. The anemometer output required about two minutes to settle, during which time the parameters required for measuring the secondary flow rate were noted. These were : pressure drop across the orifice plate, and the upstream static pressure and temperature,  $\Delta h$ ,  $p_1$  and  $t_1$  respectively.

$V_{LIN}$  was read on the DVM, and  $y$  on the dial gauge, and then the hot wire was moved to the next  $y$  point. This was continued to the end of the velocity traverse. Finally the centre line velocity was measured, which was assumed to be the local  $U_\infty$ , and the free stream temperature,  $t_2$ .

The traversing rig was shifted to X2, and the transducer was traversed to  $y=1$ . The entire operation was repeated. Velocity profiles were measured at X1, X2, X3 and X4.

Special care was taken to seal the plenum chamber. Prestik was used at the porous plate edges, to ensure that the secondary flow did pass through the porous plate, and not form jets at the edges. This was more important when testing the A grade plate, which had very low porosity.

TABLE 6.1 :  $x$ -co-ordinate positions

Position denoted by :-	Distance from leading edge (mm)
X1	780
X2	856
X3	945
X4	1045

## CHAPTER 7

### PRESENTATION AND DISCUSSION OF RESULTS

#### 7.1 Introduction

This chapter is divided into three main sections : the proving of the tunnel, the experimental results and the computer runs.

A spanwise velocity traverse is presented, and some difficulties with regard to the tunnel are discussed. In the first part, turbulence intensity and the spectrum analysis are given.

Secondly, the velocity traverses are shown in the popular non-dimensional form,  $u/u_0$  vs  $y/\theta$ . Graphs containing information more compactly are given to illustrate trends when the lower porous surface was partially blocked. The bilog plotting technique of Black and Sarnecki (1959) is used to show the development of the profile with injection ratio, and also to compare the present results with those of Mickley and Davis (1956). The possibility of a new two-parameter family of curves is discussed, which appear to be independent of  $u_w^*$  and  $F$ .

Plots obtained from the computer runs are included in the third section. Comparing these results with experimental work proved more problematic than was expected. For low injection ratio, it was possible to compare one of the profiles with that of Mickley and Davis (1959).

## 7.2 Initial Tests

### 7.2.1 Transverse velocity profile

The transverse velocity profile was measured to ensure that two dimensional flow was approximated at the centre line. The transducer was fixed at a particular value of  $y$  and  $x$ . The traversing rig was moved across the tunnel, taking readings at various intervals. One of these profiles is reproduced in Fig 7.2.

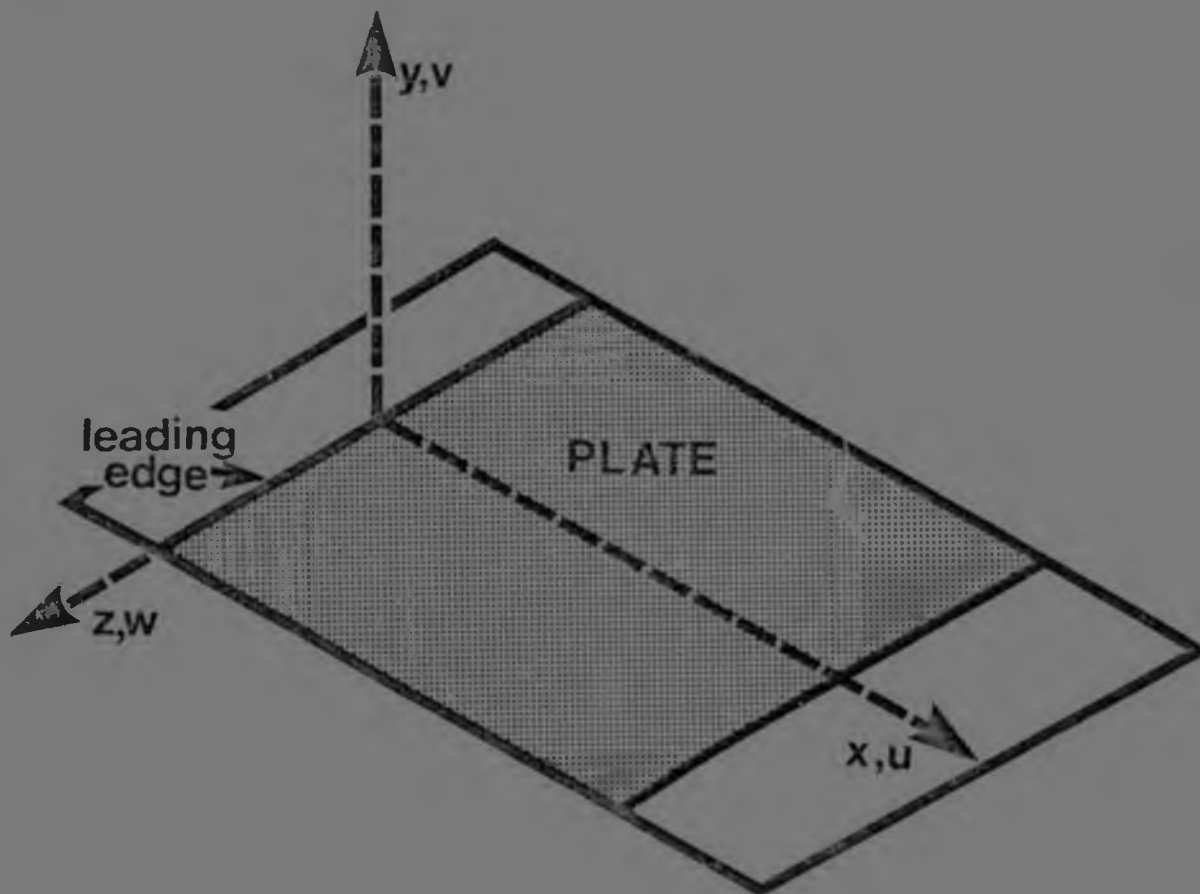


FIG 7.1 : DEFINITION OF AXES : THE X-AXIS LIES ON THE CENTRE-LINE OF THE WORKING SECTION, COINCIDING WITH THE FLOOR, Y WAS NORMAL TO THE FLOOR, AND Z TRANSVERSELY COINCIDENT WITH THE FLOOR. THE VELOCITY COMPONENTS ARE INDICATED BY  $u$ ,  $v$  AND  $w$  RESPECTIVELY



FIG 7.2 : TRANSVERSE VELOCITY PROFILE

The standard deviation of the transverse profile is  $s = 0,048$  m/s, which is within the expected accuracy of the anemometer. The wind tunnel thus produced essentially two dimensional flow in the Working Section. All profiles were measured on the centre line, ie. at  $z = 0$ .

### 7.2.2 Turbulence and turbulence intensity

The turbulent fluctuations were observed by connecting the output of the lineariser to a Tektronix oscilloscope. This output was recorded and is presented in Figs 7.3 to 7.5. The first two traces show that the fluctuations were low in the main-stream. The trace depicted in Fig 7.5 was obtained by placing the probe 1 mm from the working section floor. It should be noted that the single wire sensor necessarily registered  $u'$ ,  $v'$  and a small percentage of  $w'$ . The probe was placed normal to the flow for all tests. Consideration of the results of Klebanoff, reproduced in Schlichting (1968), showed  $u'$ ,  $v'$  and  $w'$  all increased near the wall.

Figs 7.3 and 7.4 gave an indication of the autocorrelation, because the two superimposed traces were taken at time  $t$  and  $t + \tau$ . This was not quantified.

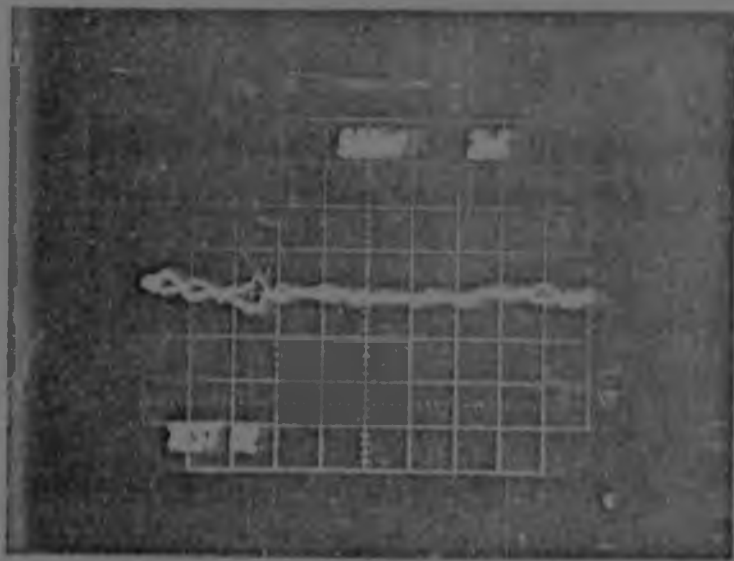


FIG 7.3 :  
TWO TRACES FOR  
 $U_{\infty} = 16$  m/s, PROBE  
AT CENTRE LINE

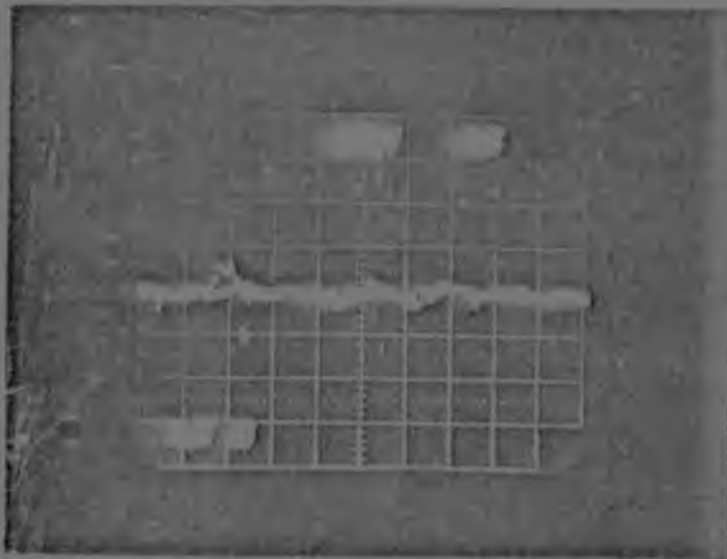


FIG 7.4 :  
TWO TRACES FOR  
 $U_{\infty}=12\text{m/s}$ , PROBE  
AT CENTRE LINE



FIG 7.5 :  
SINGLE TRACE FOR  
 $U_{\infty}=16\text{ m/s}$ , PROBE  
1 mm ABOVE SOLID  
WALL

A series of tests were performed to ascertain the turbulence intensity of the wind tunnel. By definition,

$$T = \sqrt{1/3(\overline{u'^2} + \overline{v'^2} + \overline{w'^2})}/U_{\infty} \quad \dots\dots\dots [7.2.1]$$

For isotropic flow (produced in wind tunnels behind a grid or screen)  $\overline{u'^2} = \overline{v'^2} = \overline{w'^2}$ , so for wind tunnels this equation becomes (Schlichting 1968)

$$T = \sqrt{\overline{u'^2}}/U_{\infty} \quad \dots\dots\dots [7.2.2]$$

The transducer was subject to  $u'$  and  $v'$ . Because  $w'$  was parallel to the sensor wire, this contribution should have been small. If  $\overline{r'^2} = \overline{u'^2} + \overline{v'^2}$ , then define a wind tunnel turbulence intensity

$$T = \sqrt{\overline{r'^2}}/U_{\infty} \dots\dots [7.2.3]$$

The results of a number of runs at different velocities are given in Table 7.1. The integrator time constant was 10 sec. The probe was positioned at  $x = 800$  mm,  $y = 112,5$  mm,  $z = 0$ .

TABLE 7.1

Test	$V_{\infty}$ (volts)	$U_{\infty}$ (m/s)	$\sqrt{\overline{r'^2}}$ (volts)	$T\%$
A	3,190	15,359	0,891	0,59
B	2,558	12,335	0,640	0,53
C	2,191	10,579	0,504	0,49
D	1,835	8,876	0,457	0,53
E	1,392	6,756	0,368	0,56

$U_{\infty}$  was varied by opening the by-pass. Considering Fig 7.6, it would be expected that  $T$  would have been lower had a more sophisticated method for decreasing  $U_{\infty}$  been employed. In the main test, this was done by closing the fan intake with a baffle plate. The increased angle of the trap door may explain why  $T$  dropped initially, but then increased in the last two tests.



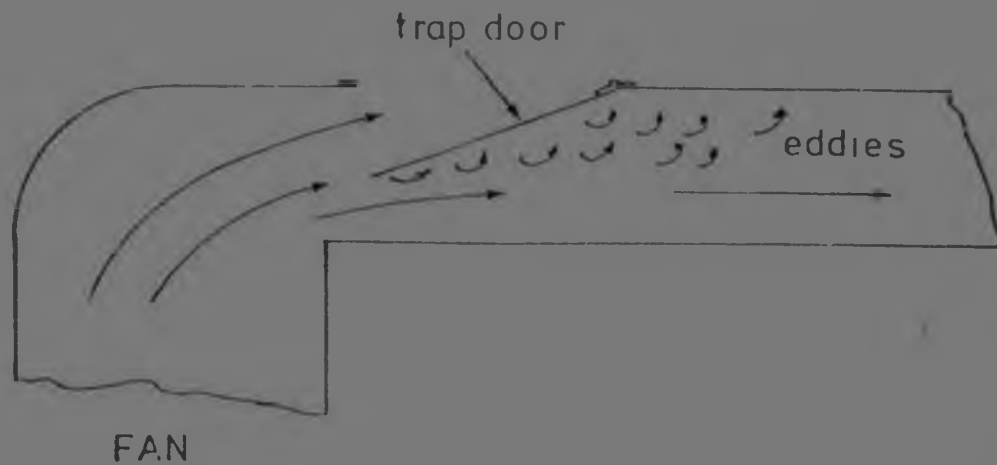


FIG 7.6 : DETAIL OF BY-PASS, SHOWING POSSIBLE EFFECTS OF OPENING THE TRAP DOOR

### 7.2.3 The energy spectrum

The energy spectrum was measured in tests A through E. The results are reproduced in Figs 7.7a to 7.7e. These are plots of  $\log E_1(n)$  vs  $\log n$ . The inertial subrange was established i.e. the portion of the curve which is fitted by the  $-5/3$  law. This is a range in which vortices were neither generated or dissipated. It has been proposed that the large, low frequency eddies are formed by drawing energy from the main stream, creating the turbulent stresses. Because of the instability of the flow, smaller and smaller eddies appear until they are dissipated as heat (see Schlichting, 1968).

These plots show that a full spectral distribution of the energy of fluctuation was present in the tunnel, throughout the range of free stream velocities of interest.

In each spectrum the inertial, eddy generating region can be seen. Furthermore, the function  $E_1(n) \propto n^{-5/3}$  fitted the data in the inertial subrange. Noise from the anemometer and analyser disguised the high frequency range.

Fig 7.7b to Fig 7.7e show an energy bulge at high frequency. This occurred at approximately 6 and 10 Hz and is thought to have been caused by vortex shedding behind the trap door. This was discussed in section 7.2.2 in relation to Fig 7.6, and resulted in the discontinued use of the by-pass as a method of flow control.

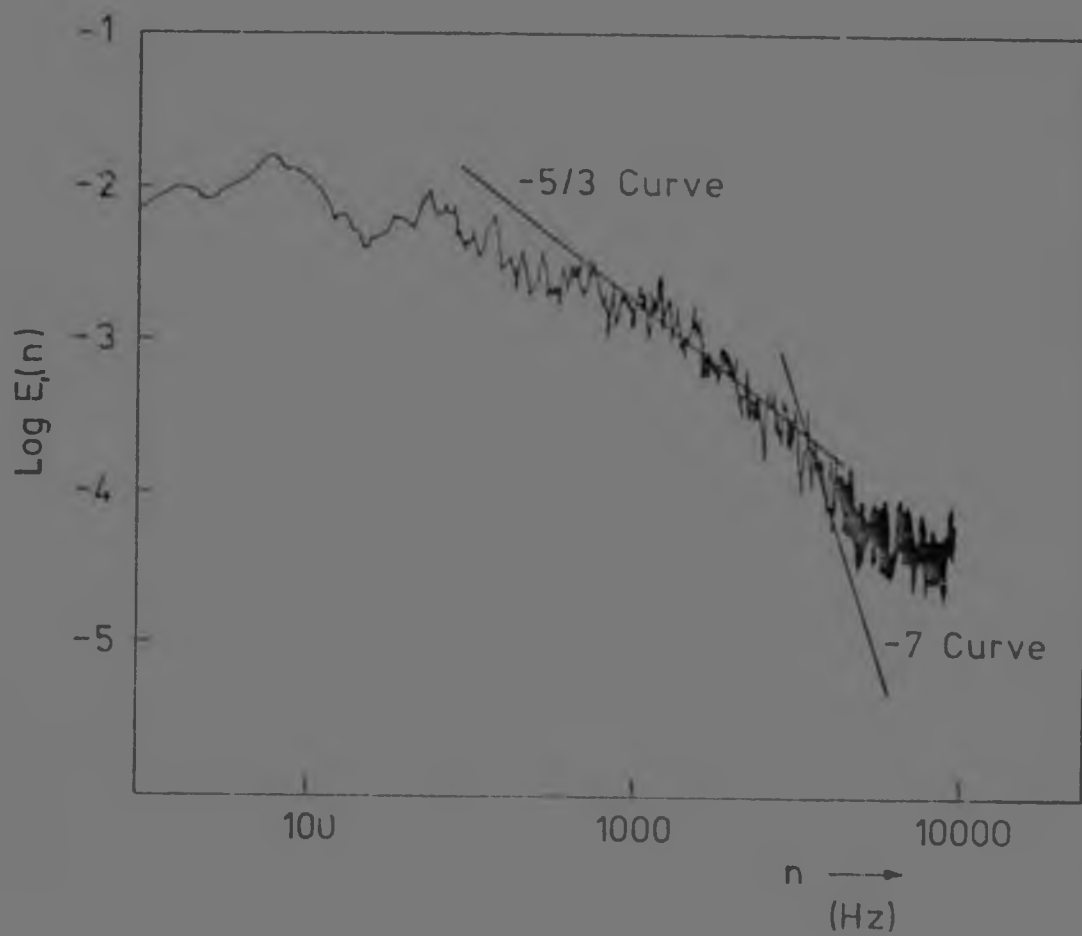


FIG 7.7a : TEST A

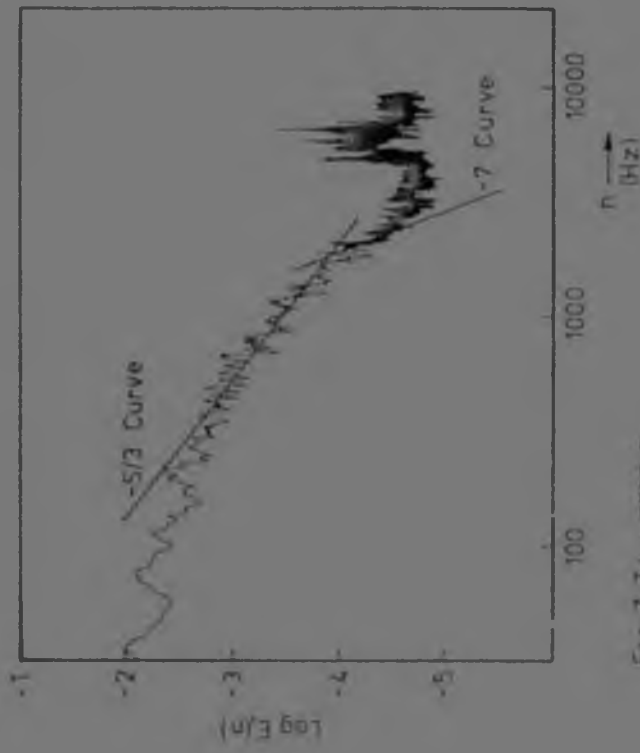


Fig 7.7d : TEST D

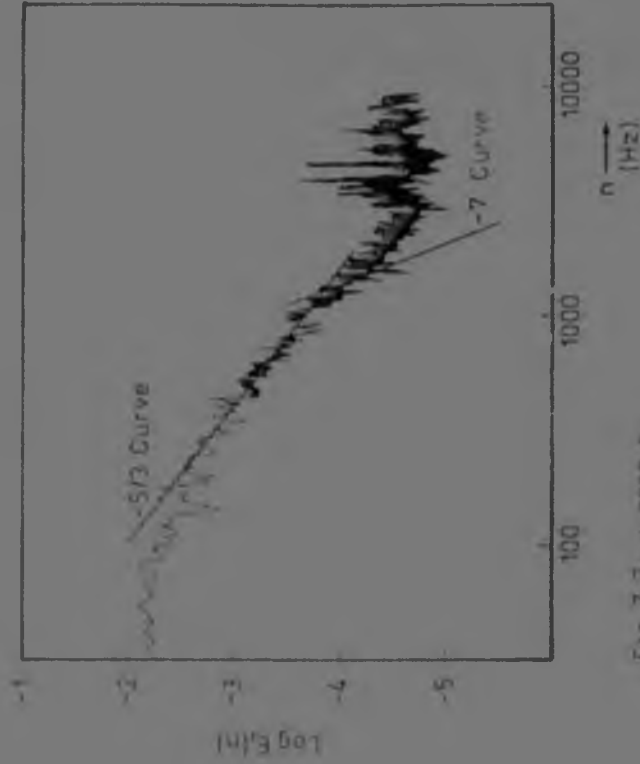


Fig 7.7e : TEST E

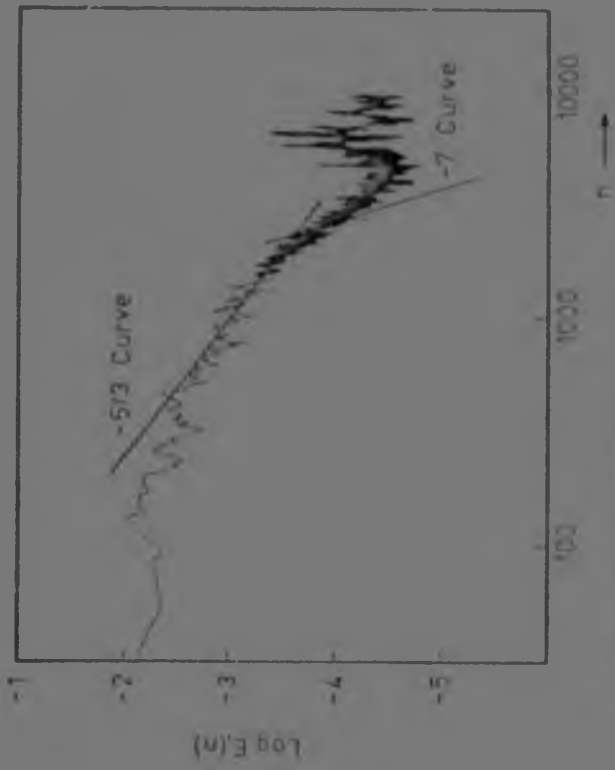


Fig 7.7a : TEST A

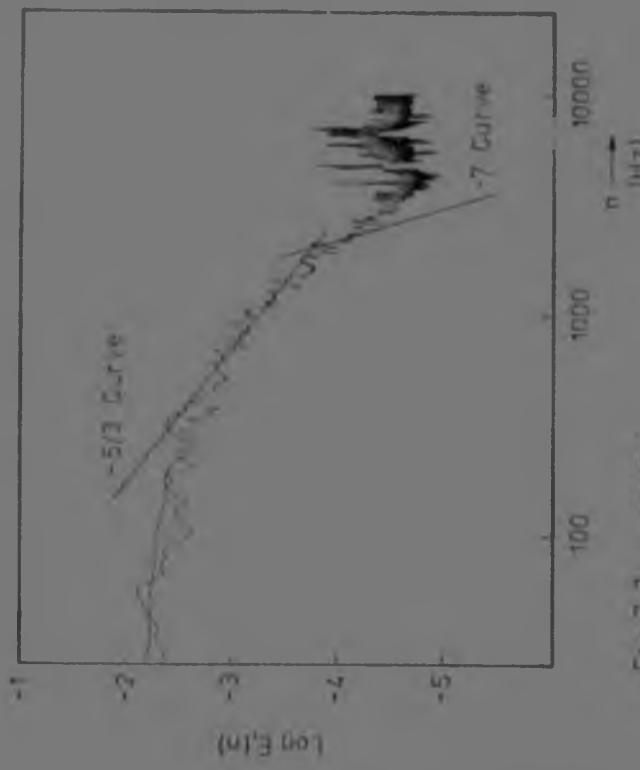


Fig 7.7c : TEST C

### 7.3 Experimental Work

In each test, the profile at a station just prior to the start of the porous section was measured. This station was denoted X1. The shape factor for these profiles ranged between 1,27 and 1,30, which is the correct value of  $H$  for the unblown, zero pressure gradient turbulent boundary layer. The simple Blasius profile was found not to fit the experimental data. Theoretical curves were constructed from Thompson's two parameter family (Thompson, 1965). Near the wall agreement was not as good, but it should be borne in mind that the given families had to be interpolated to obtain the approximate value of  $u/\theta$ . Two of these profiles are given in Figs 7.8 and 7.9.

The reason for this result is thought to be as follows. The contraction was in the XZ-plane, and thus the transverse profile was exceptionally flat. In the YZ-plane, the tunnel height remained constant from the start of the diffuser. This resulted in a very thick boundary layer arriving in the entry section. It was therefore not possible to draw the entire boundary layer off through the boundary layer bleed. Almost duct flow resulted, rather than the ideal "flat plate in a free stream". Large  $\delta$  did have the advantage that the  $y$  measurement was less critical, and the rate of change of  $u$  with  $y$  was much smaller. The disturbance of a transducer in the shear layer was thus decreased.

It must be pointed out that numerous investigators used a pitot tube for measuring the velocity. Because this device is larger than the miniature hot wire transducer used in the present work, the disturbance to the flow in the region of measurement was necessarily greater. In addition, it was found to be essential that the pitot tube was exposed to the correct static pressure. When a pitot tube was used in the tunnel running at a pressure slightly higher than atmospheric, incorrect readings were obtained. The same comments would apply to a pitot rake.

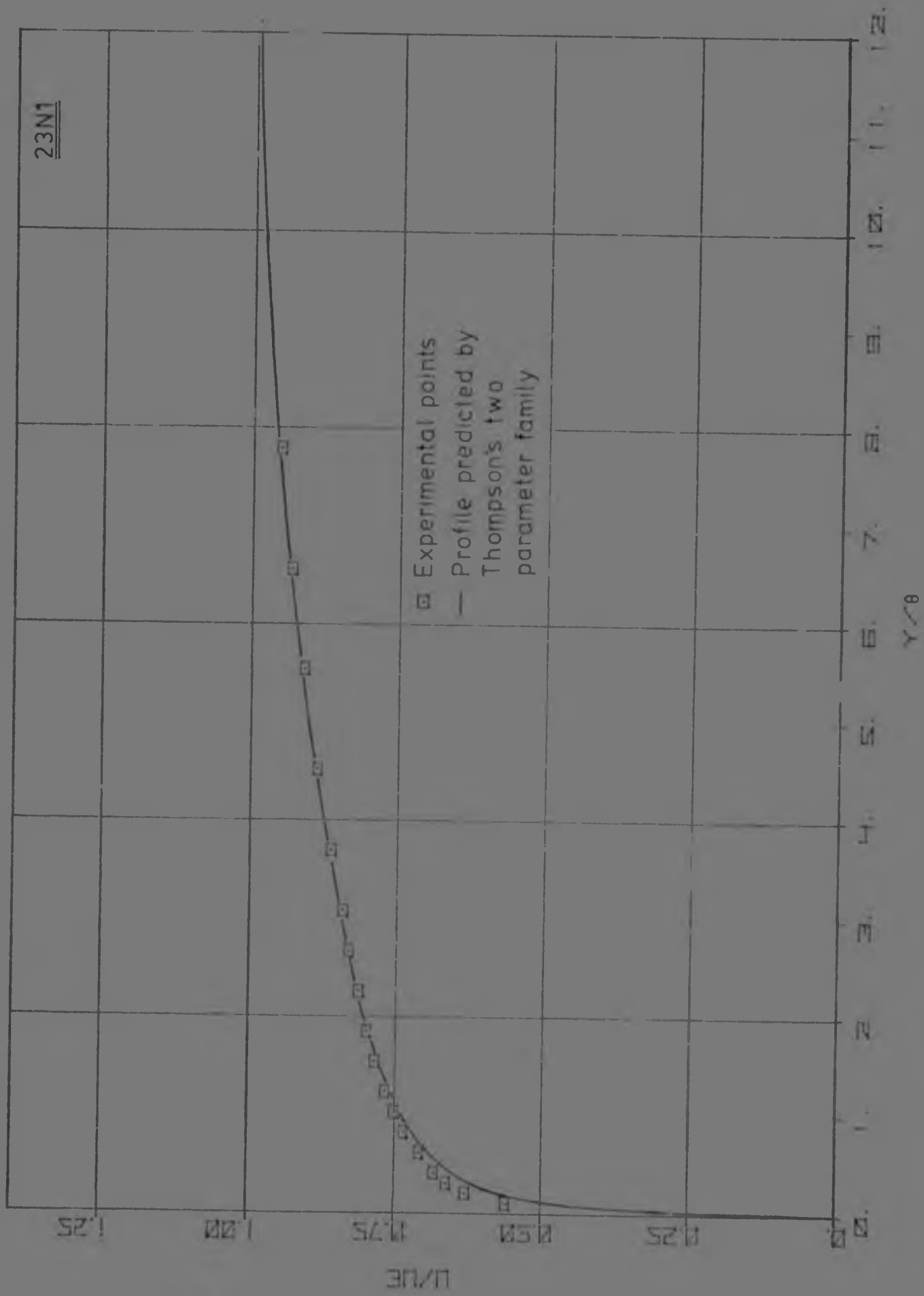


FIG 7.8 : EXPERIMENTAL PROFILE MEASURED ON THE SOLID PLATE FOR TEST 23N1

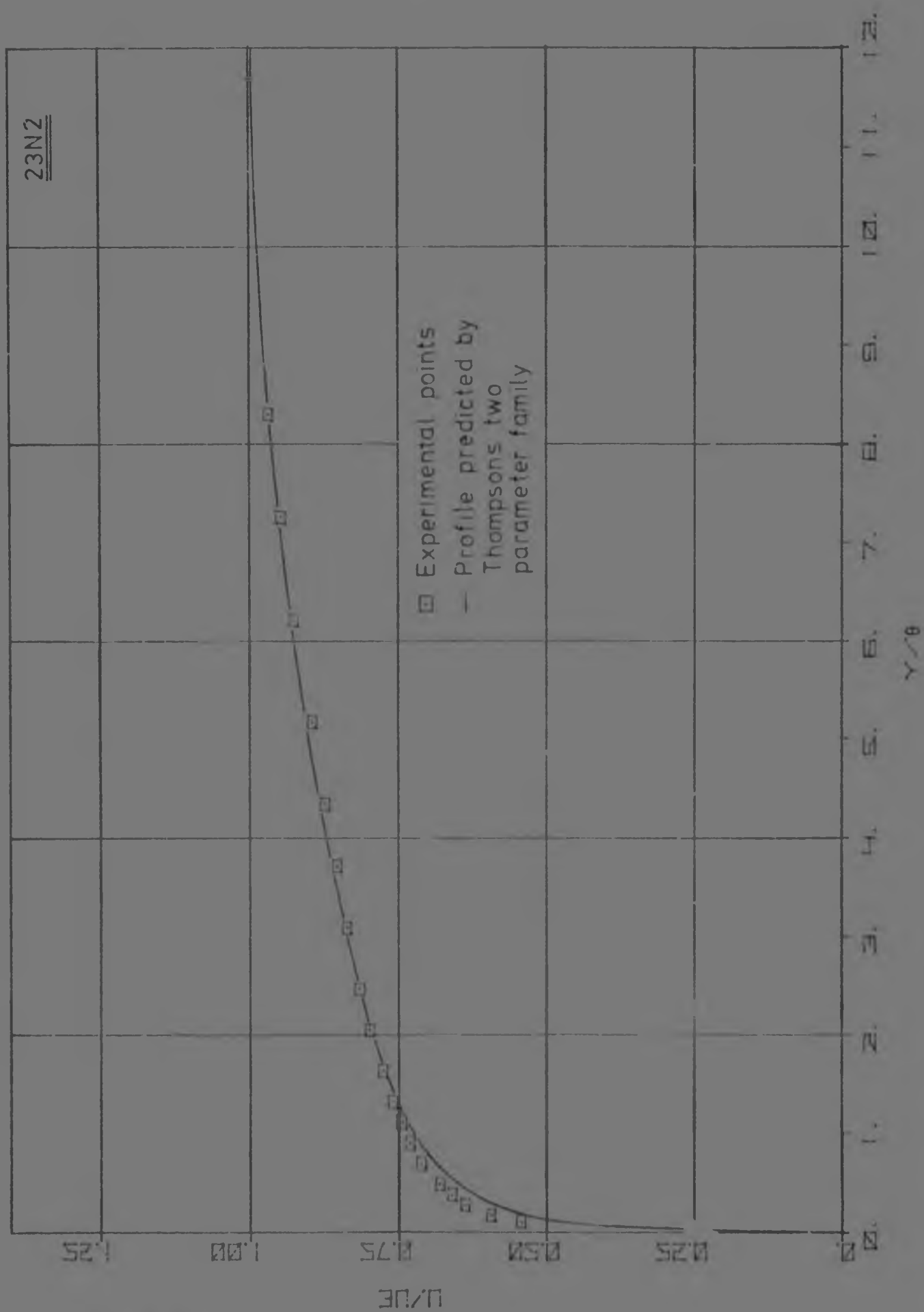


Fig 7.9 : EXPERIMENTAL PROFILE MEASURED ON THE SOLID PLATE FOR TEST 23N2

Great care was taken in measuring the profiles for the present work. The hot wire was re-calibrated for each run. A difficulty always present when using hot wires is that the transducer is not to be traversed to  $y = 0$  to determine that point. A shield was attached to the probe holder with a broken probe. This probe was traversed down until it just touched the floor, and the dial gauge reading noted. The calibrated transducer was then placed in the holder. Since the two transducers were of equal length,  $y$  was calculated.

The results were processed on an IBM 370 computer, which printed and plotted the data.

#### 7.3.1 The experimental results

The basic data is given in Appendix E1. The non-dimensionalized plots for all the tests, and baseline data for each test are given in Fig 7.10. The velocity profiles are followed by graphs of shape factor,  $H$  vs physical  $x$  for the maximum main stream velocity tests, and the tests with blockage. In these graphs the change in profile caused by blockage is indicated.

In the tables of Fig 7.10 the heading includes the date, e.g. test 28S1 was run on the 28th September, the 1 stands for the number of the test. The second part indicates the grade of porous plate, then  $F_{av} \times 10^{-3}$  and finally the main stream velocity,  $V_{CL}$  in m/s.  $Re_m$  was based on the distance measured from the leading edge. These were X1 = 0,7805 m, X2 = 0,8565, X3 = 0,9455, X4 = 1,0455. X1 was on the solid plate, X2 = 18,66%, X3 = 49,13% and X4 = 83,36% along the porous plate.  $F_{av}$  was based on the inlet velocity.

In Fig 7.10, the symbols used for plotting are :- A for X1, + for X2, X for X3 and  $\odot$  for X4. The co-ordinates used are  $u/u_e$  vs  $y/\theta$ . This is preferred to  $y/\delta$ , because  $\delta$  is less easy to determine.

The first five tests were run at maximum free stream velocity, with an E grade plate installed. Results are shown in Fig 7.10a to e. In all tests, the increasing 'laminarisation' with  $x$  is evident, the curves coinciding at approximately  $y/\theta = 3,5 \rightarrow 5$ , this value increasing with  $F$ . For high  $F$ , the profiles at X3 and X4 tended to be very similar (see Fig 7.10e).

In test 28S3, the profile at X3 showed a strange step. This profile was obtained with probe No.4, which appeared to be faulty, and was not used again. Although not broken, the calibration curve for this probe changed during a test. Possibly a particle in the stream collided with the transducer.

Five tests, very similar to the above, were performed on the E grade plate, but at the reduced free stream velocity;  $U_{\infty} = 13$  m/s at the Working Section inlet. Very similar results to the first set were obtained, but results were not identical. This indicates that the profile was not only dependent on  $F$ . Clearly  $Re_{\infty}$  was different, but probably also the effective roughness,  $k^+$ . These tests are shown in Fig 7.10f to j.

Tests 0701 and 0601 were also run over the unblocked E grade plate, but at  $U_{CL} = 9$  m/s, and therefore at much lower  $Re_x$ .

The next set of tests, 09N1 to 11N2, shown in Fig 7.10m to q, were a repeat of the first 12 runs, but on the C grade, unblocked plate. The results were very similar, but showed a shift which was thought to have resulted from the change in roughness,  $k^+$ .

Finally, a sample test was performed, using the A grade plate. Because of the low porosity, and the resultant high pressure drop, great care was taken when sealing the plenum chamber. Powder metallurgy was used to make this plate, whereas the



PLOT OF U/UE VS Y/TTA

RN 28-9 1,4-7,3-17

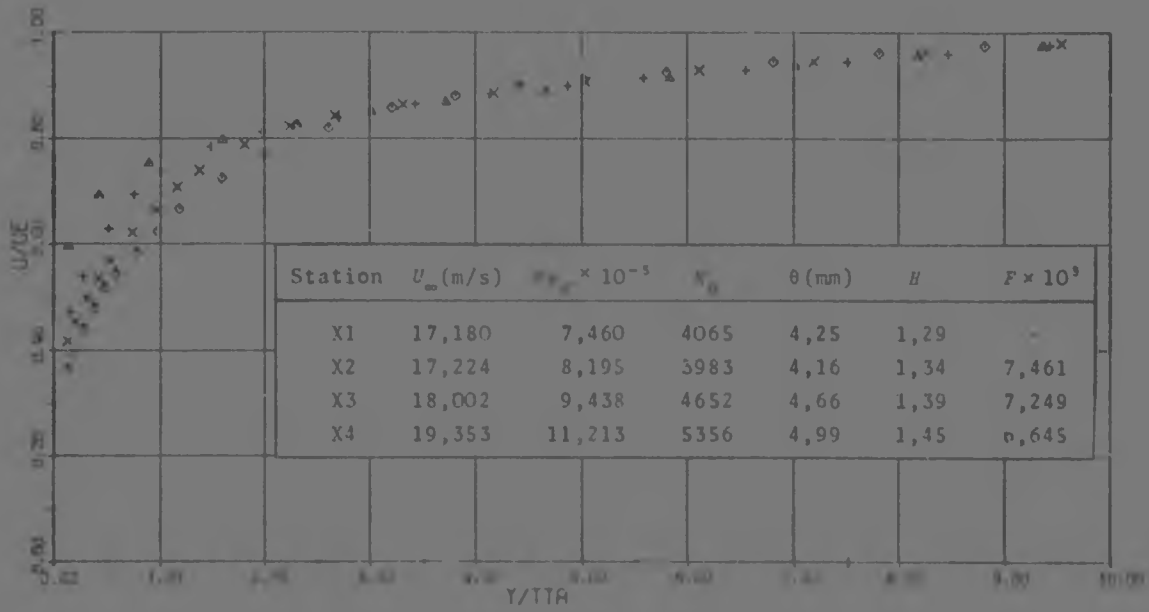


FIG 7.10a : TEST 28S1 E-7,58-17

PLOT OF U/UE VS Y/TTA

RN 28-8/1,4-10,1-18

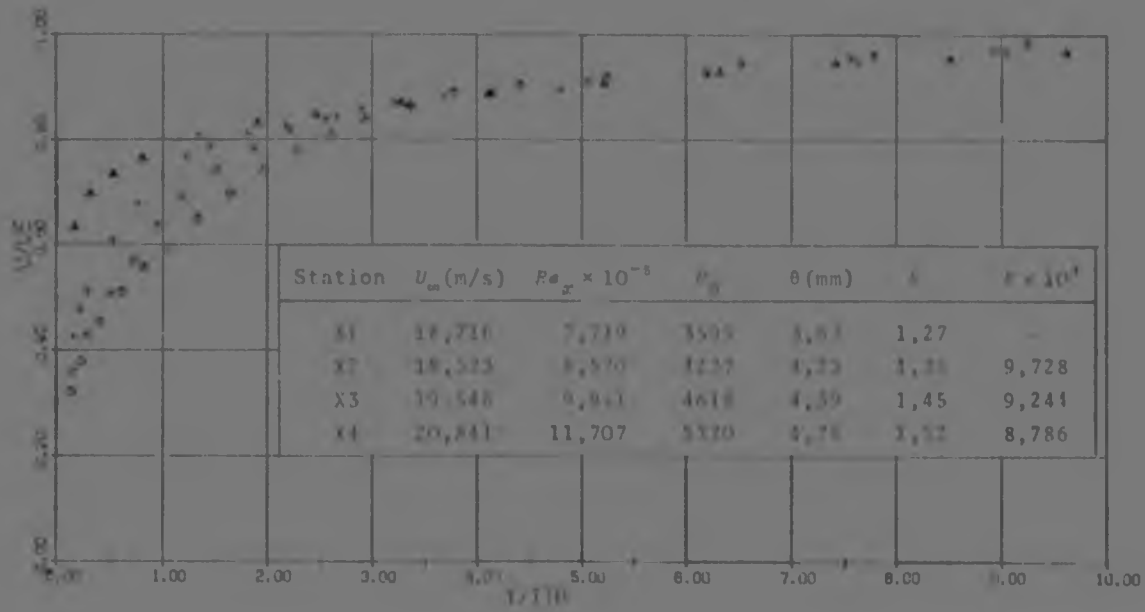


FIG 7.10b : TEST 28S1 E-7,58-17

PLOT OF U/UE VS Y/TTA

PLN 28-9/5,8-12,0-17

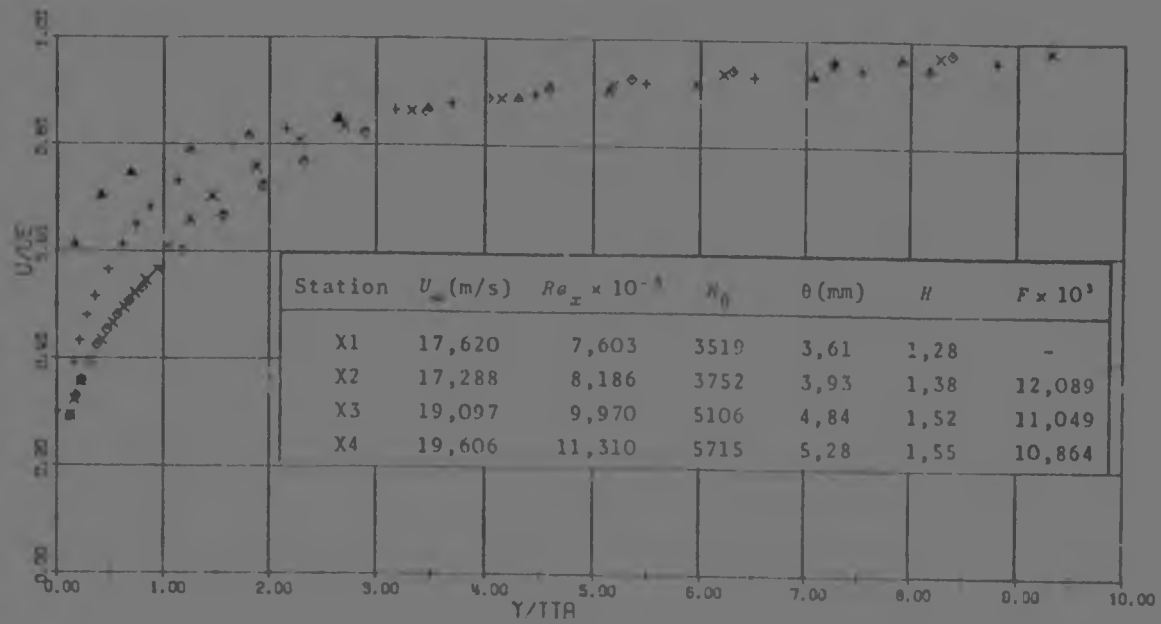


FIG 7.10c : TEST 2852 E-12,41-17

PLOT OF U/UE VS Y/TTA

PLN 30-9/5,8-12,1-17

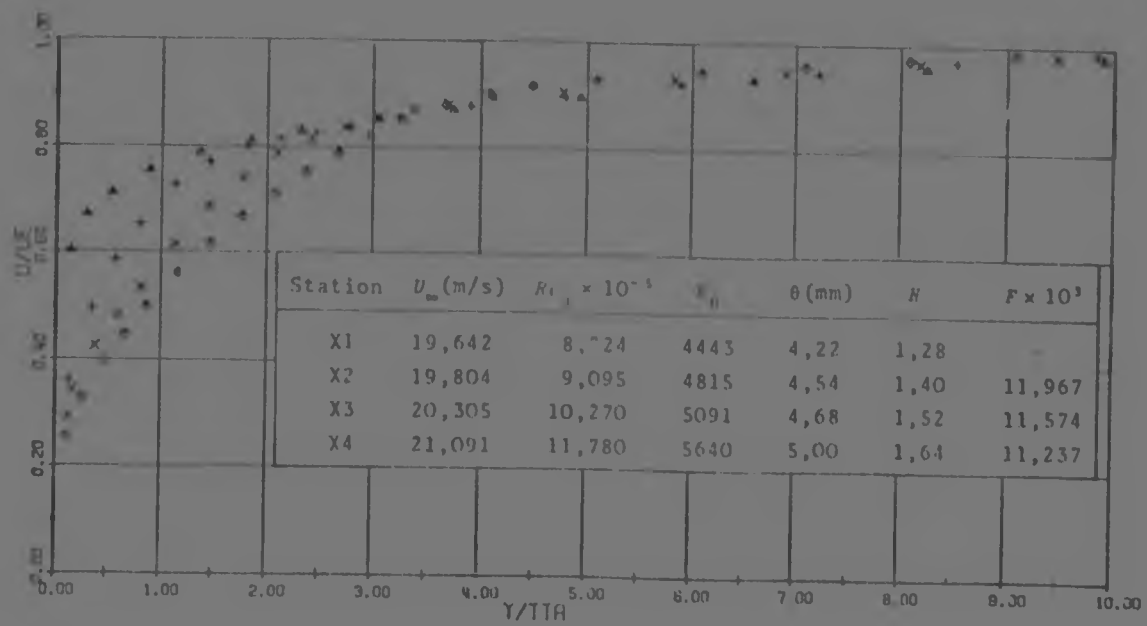


FIG 7.10d : TEST 3052 E-13,79-17

PLOT OF U/UE VS Y/TA

RN 28-9/9,12-15.5-17

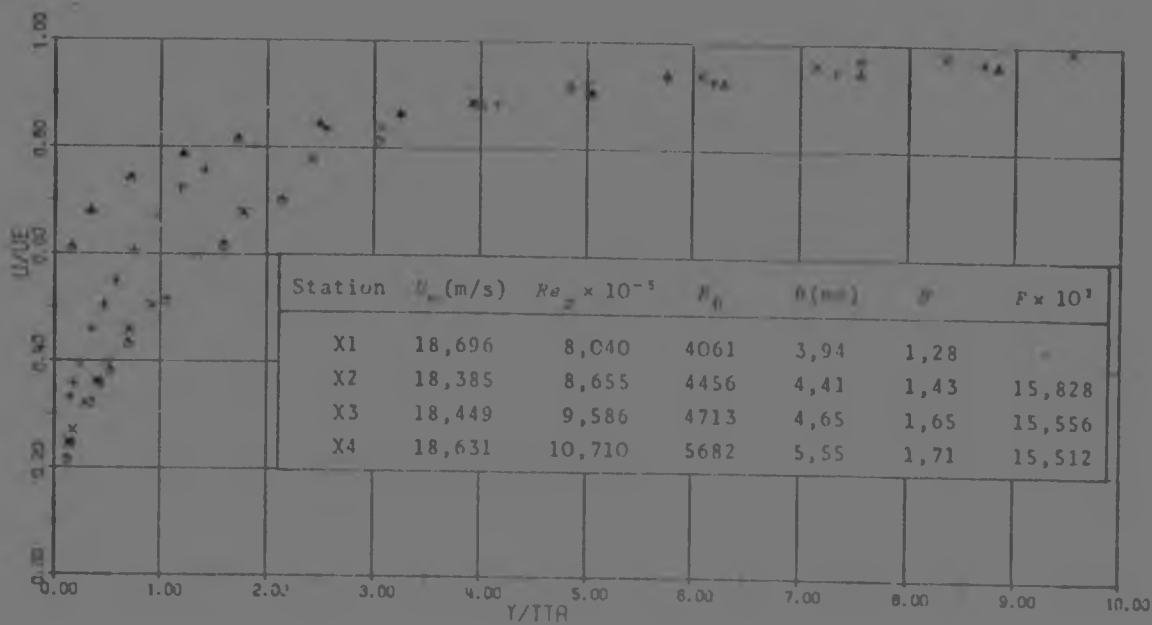


FIG 7.10e : TEST 2853 E-17,0-17

PLOT OF U/UE VS Y/TA

RN 03-10/5,8-5.3-13

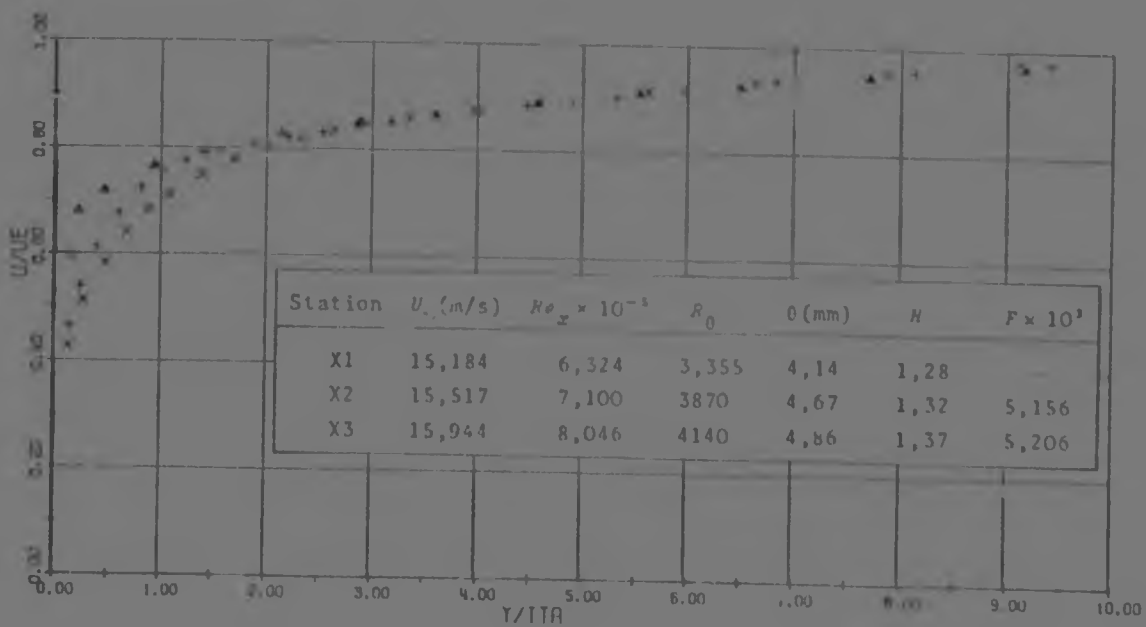


FIG 7.10f : TEST 0302 E-6,43-13

PLOT OF U/UE VS Y/TTA

MM 04-10/5, D-9, 65-13

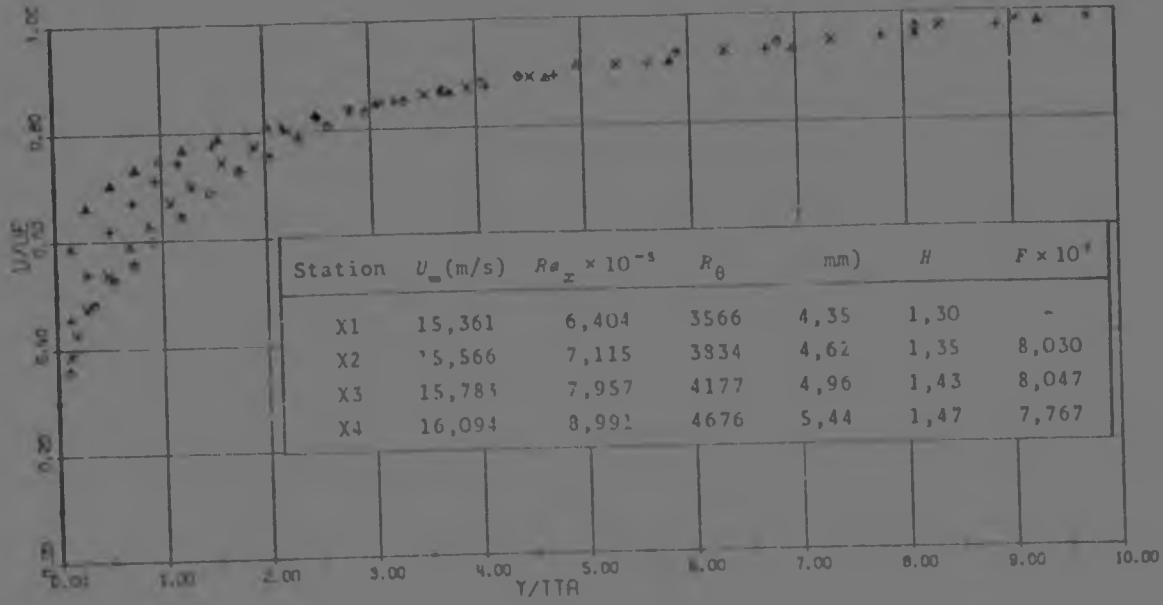


FIG 7.10g : TEST 0402 E-9,69-13

PLOT OF U/UE VS Y/TTA

MM 05-10/1, N-10, 13

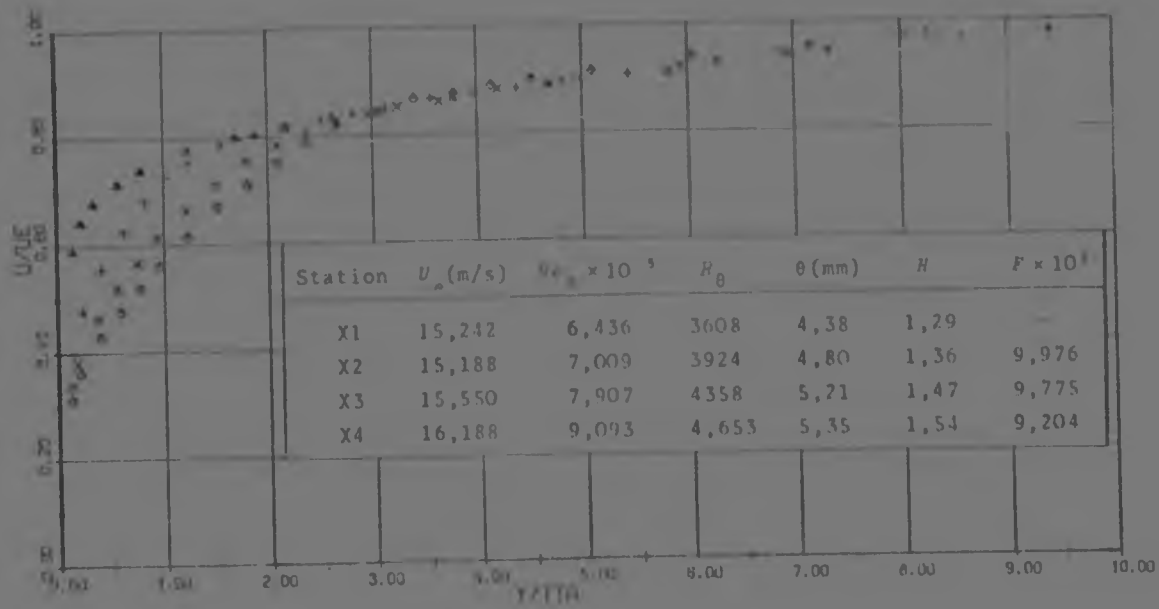


FIG 7.10h : TEST 0301 E-11,59-13

PLOT OF U/UE VS Y/TTA

MM 4-10/1, 4-13, 0-13

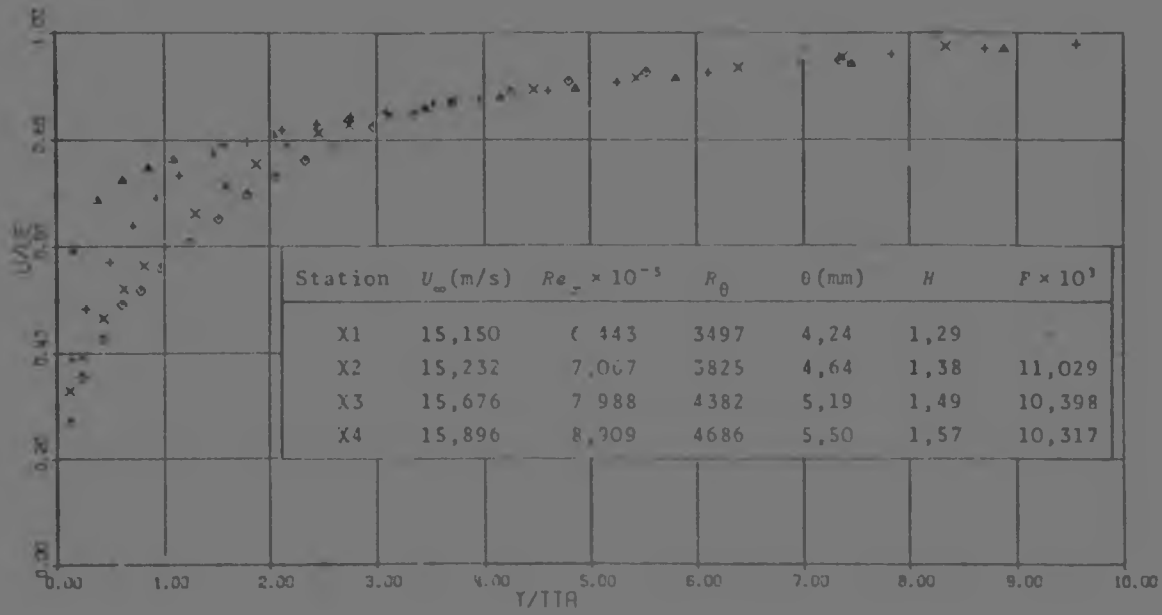


FIG 7.10i : 0401 E-12,79-13

PLOT OF U/UE VS Y/TTA

MM 5-05/1, 5-14, 3-13

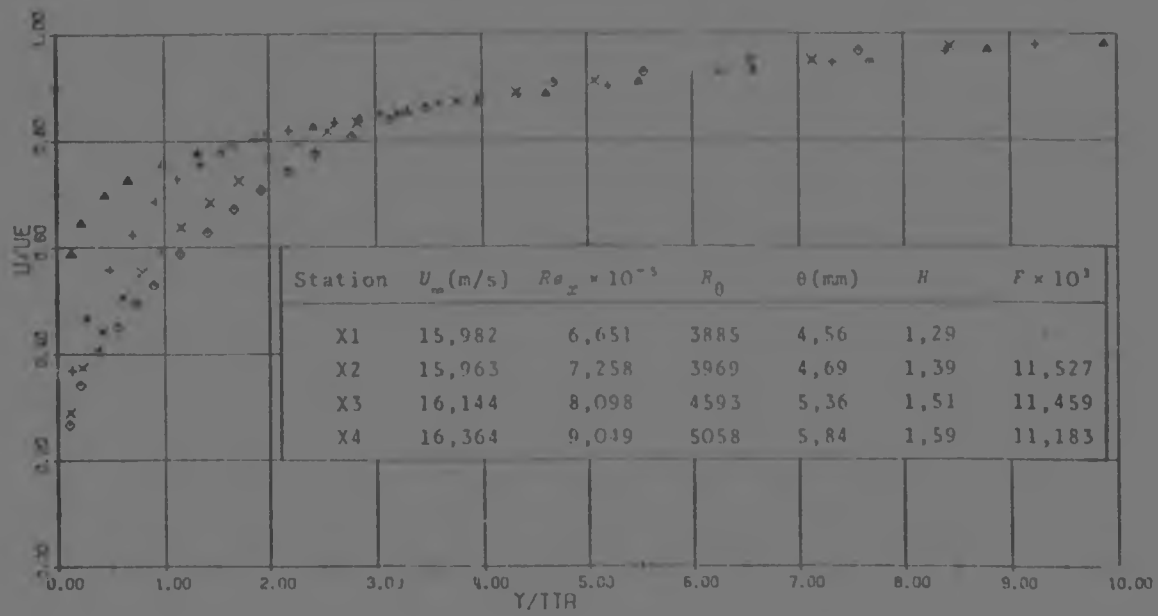


FIG 7.10j : TEST 0501 E-14,13-13

PLOT OF U/UE VS Y/TTA

REN 7-10/1.4-7.1-9.8

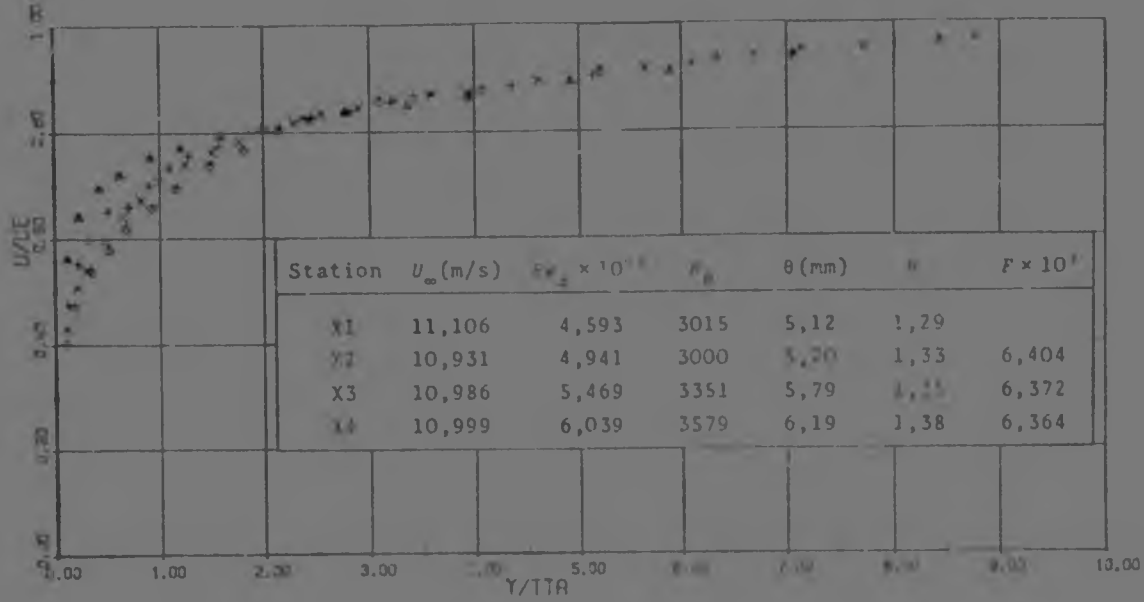


FIG 7.10k : TEST 0701 E-7,74-9

PLOT OF U/UE VS Y/TTA

REN 8-10/1.4-10.5-9.8

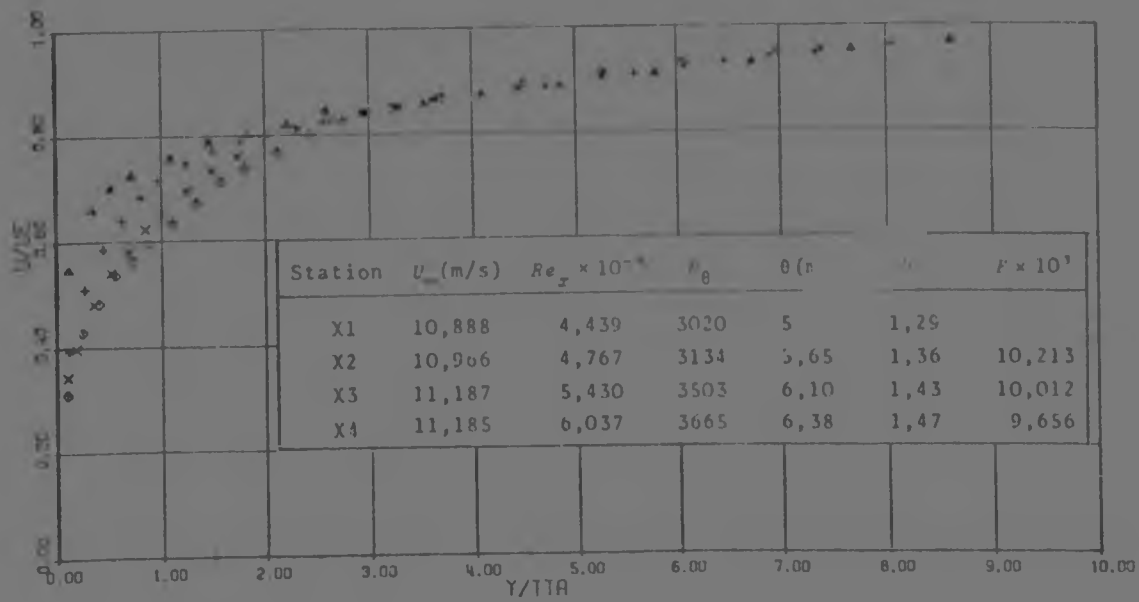


FIG 7.10l : TEST 0601 E-12,03-9

PLOT OF U/UE VS Y/TTA

REF: 11-11/1.4-17

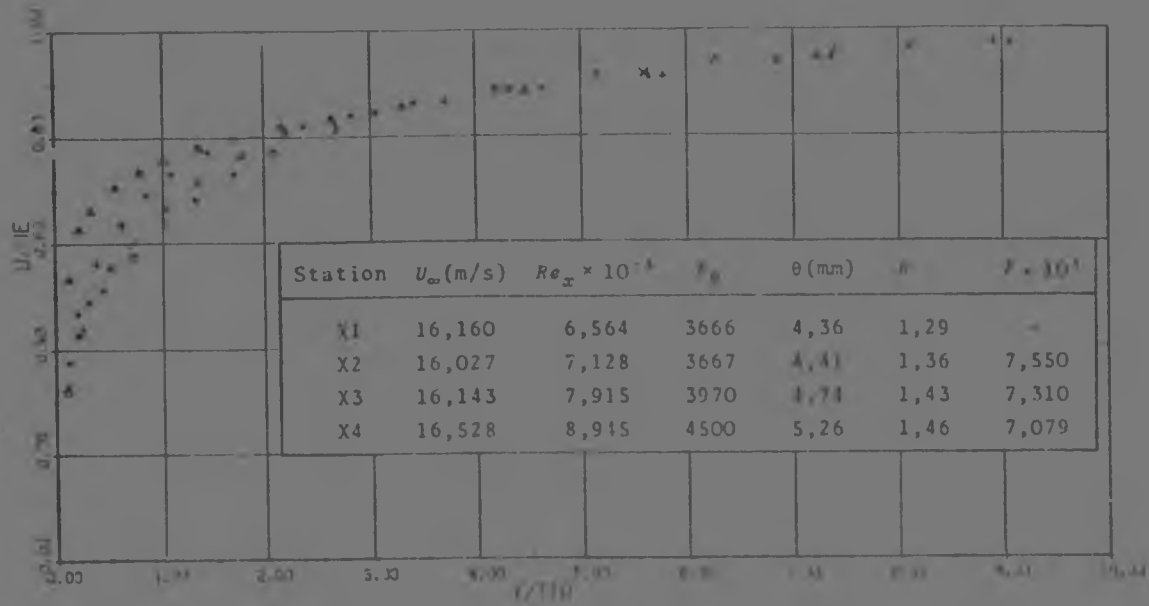


FIG 7.10m : TEST 09N1 C-7,07-17

PLOT OF U/UE VS Y/TTA

REF: 11-11/1.4-17

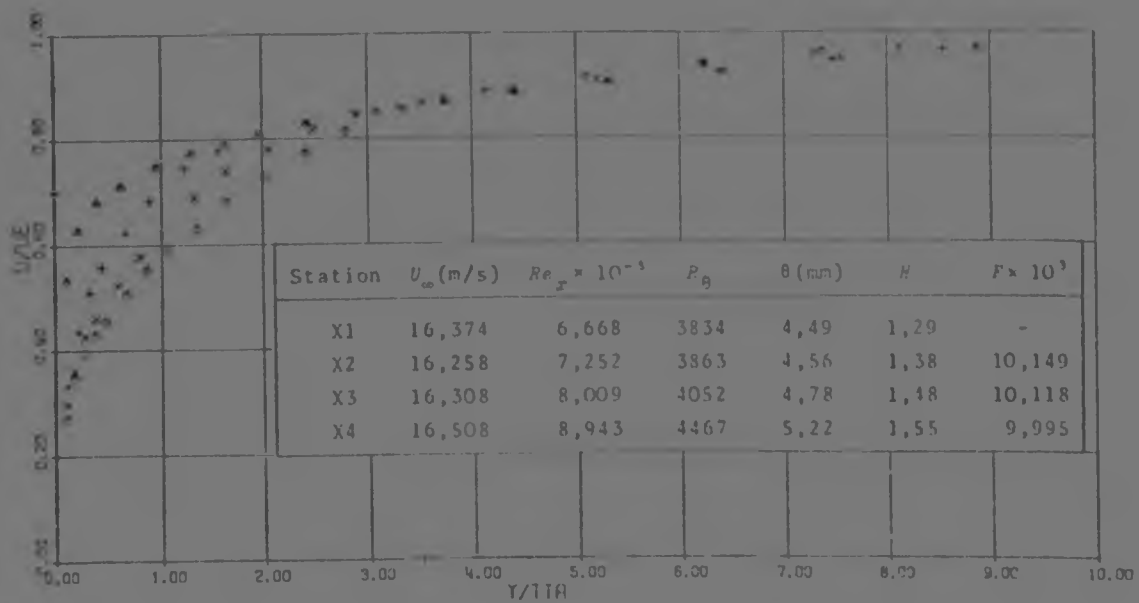


FIG 7.10n : TEST 10N1 C-9,82-17

PLOT OF  $U/U_e$  VS  $Y/\delta^*$

RN 13-11/5, 8-0, 02-17

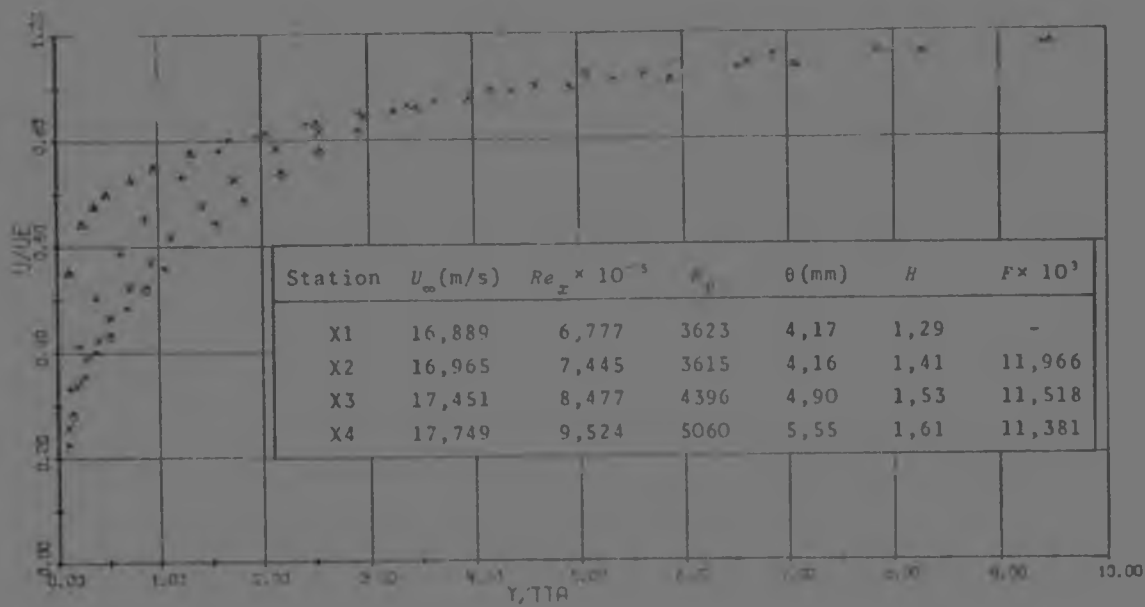


FIG 7.10o : TEST 10N2 C-12,02-17

PLOT OF  $U/U_e$  VS  $Y/\delta^*$

RN 11-11/5, 1-0, 15-13

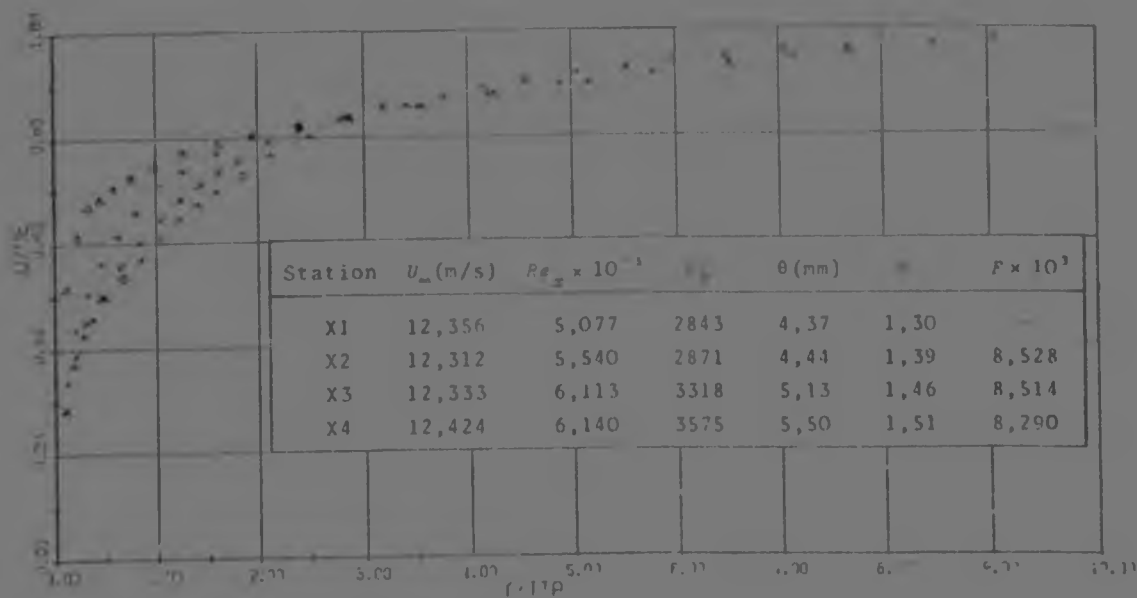


FIG 7.10p : TEST 11N1 C-8,39-13



PLOT OF U/UE VS Y/TTA

MM 11N2 C-12,22-13

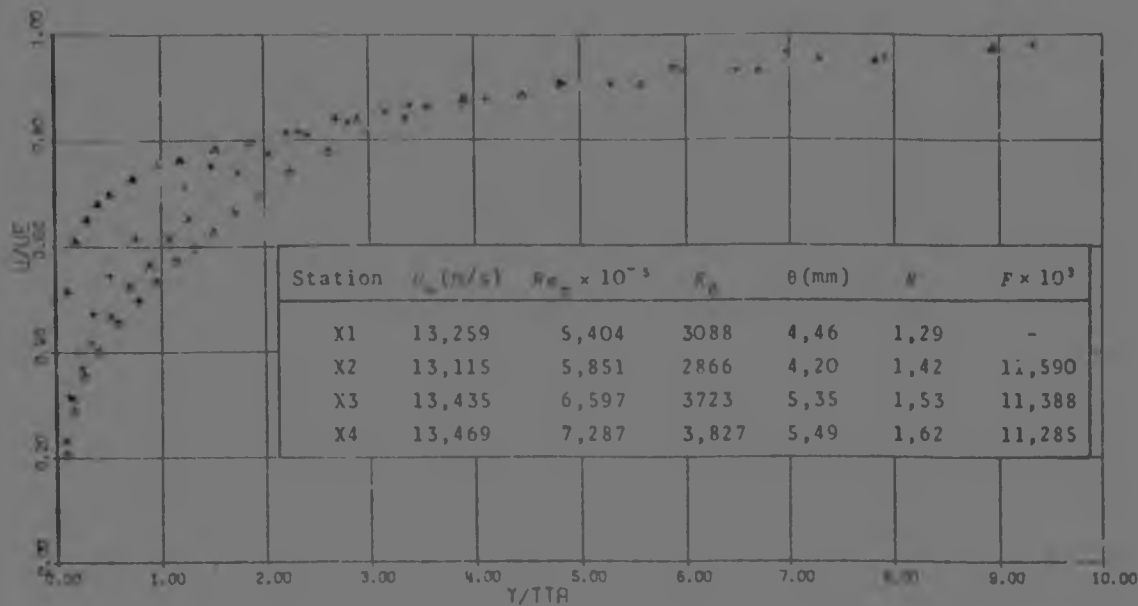


FIG 7.10q : TEST 11N2 C-12,22-13

PLOT OF U/UE VS Y/TTA

MA 16N1 A-7,55-17

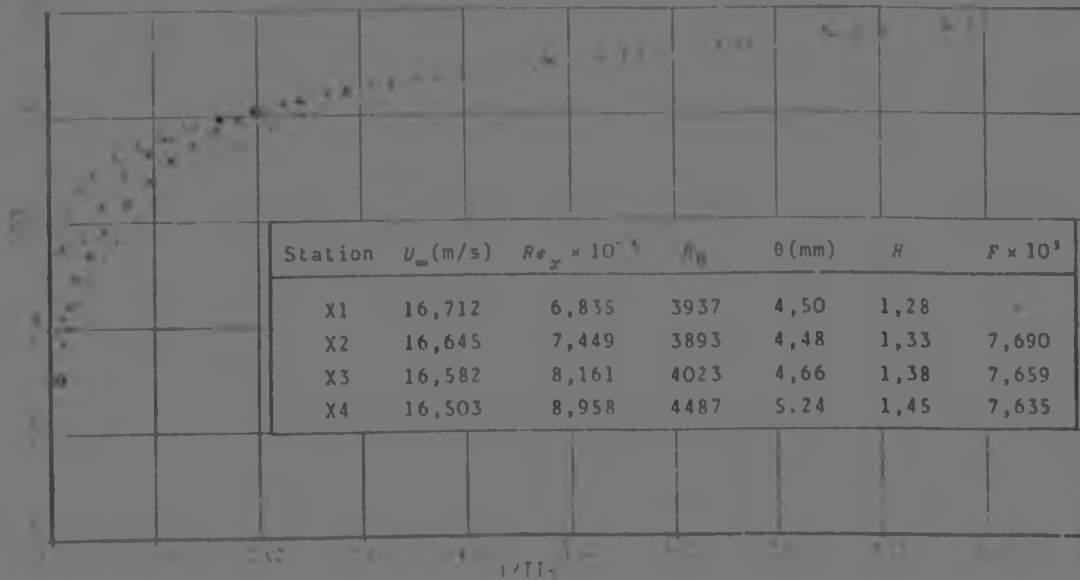


FIG 7.10r : TEST 16N1 A-7,55-17

PLOT OF U/UE VS Y/TTA

RUN 23 11/1,4-8-17E

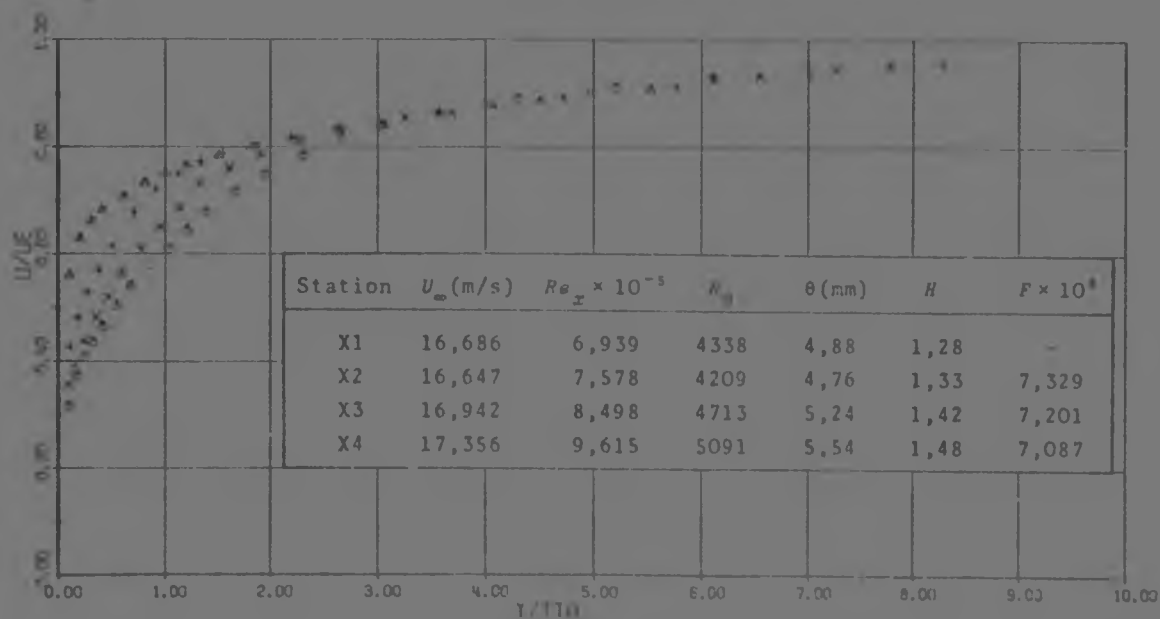


FIG 7.10s : TEST 23N1 E-7,27-17 50% BLOCKED

PLOT OF U/UE VS Y/TTA

RUN 23 11/5,8-8-17E

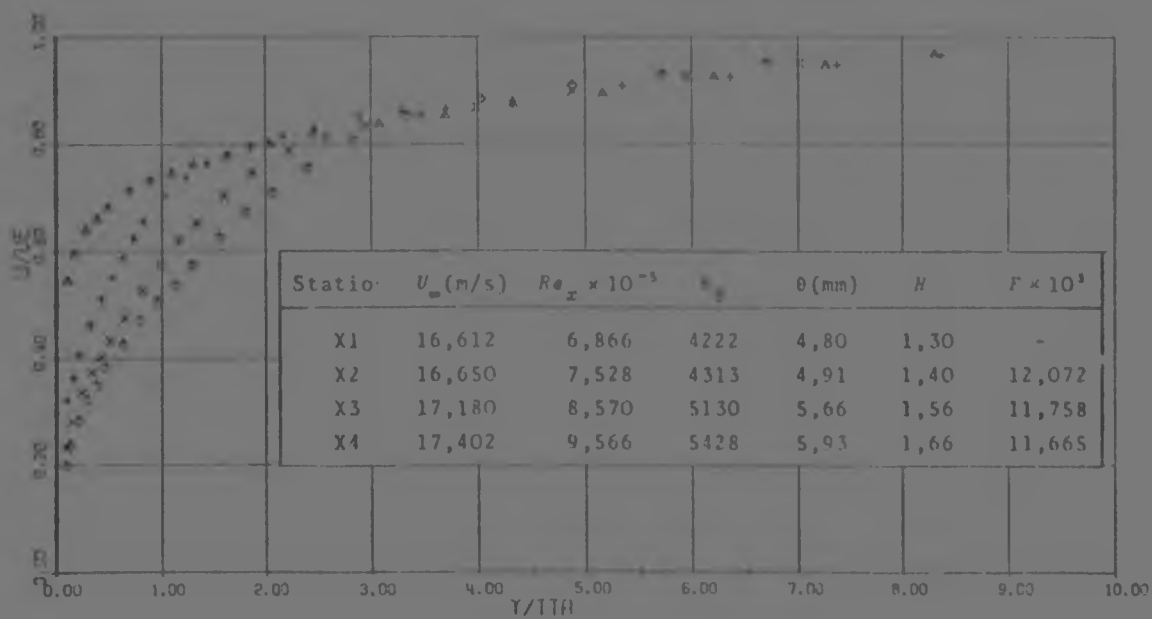


FIG 7.10t : TEST 23N2 E-12,0-17 50% BLOCKED

PLOT OF U/UE VS Y/TTA

RUN 24-11/1,4-8-17E

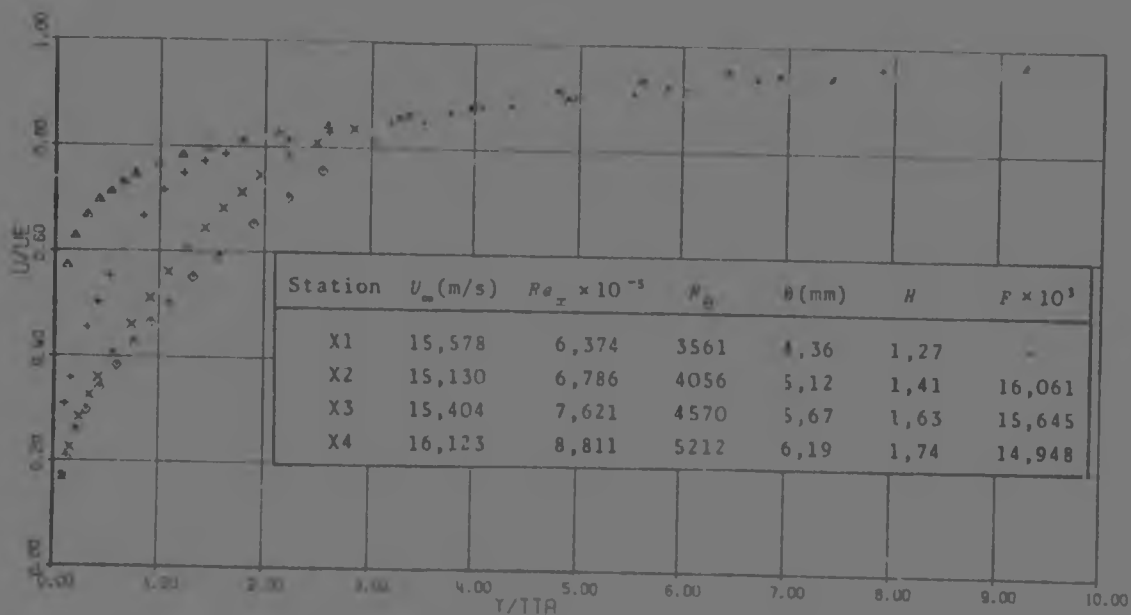


FIG 7.10u : TEST 24N1 E-14,51-17 50% BLOCKED

PLOT OF U/UE VS Y/TTA

RUN 28-11/1,4-8-17E

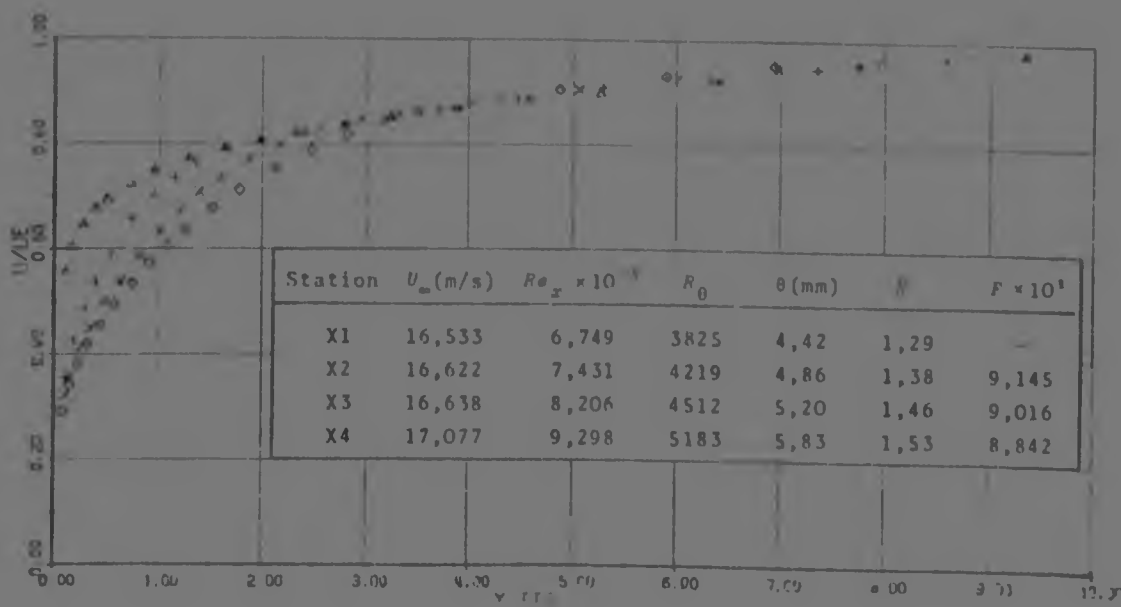


FIG 7.10v : TEST 28N1 E-9,04-17 67% BLOCKED

PLOT OF U/UE VS Y/TTA

RM 23 11/1, 4 B 17E

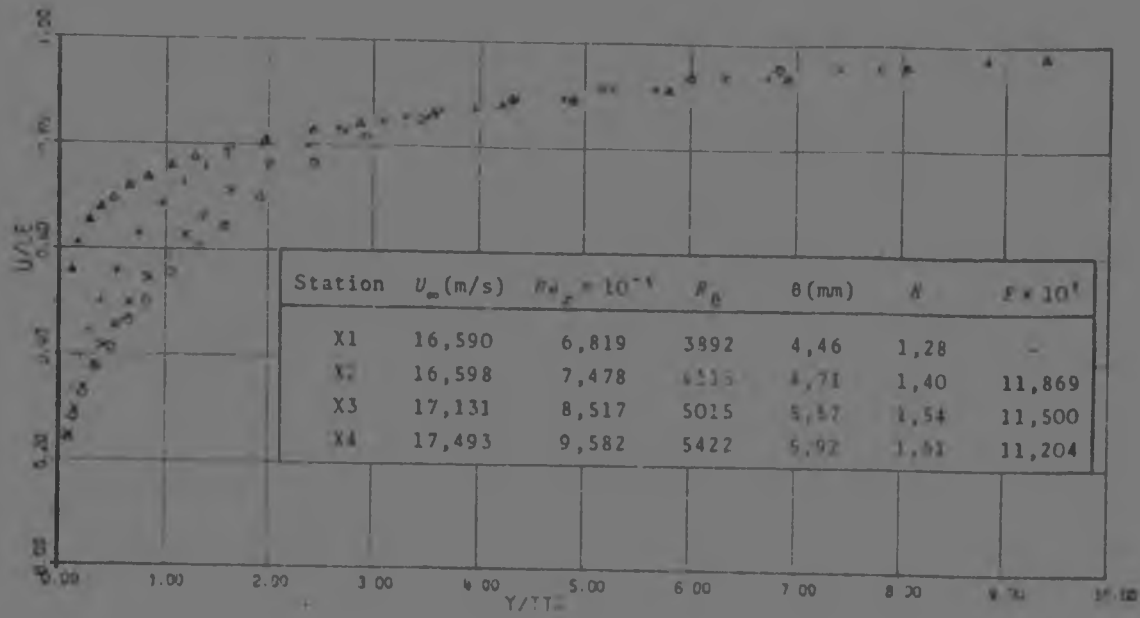


FIG 7.10w : TEST 29N1 E-11,77-17 67% BLOCKED

PLOT OF U/UE VS Y/TTA

RM 30 11/1, 4 B 17E

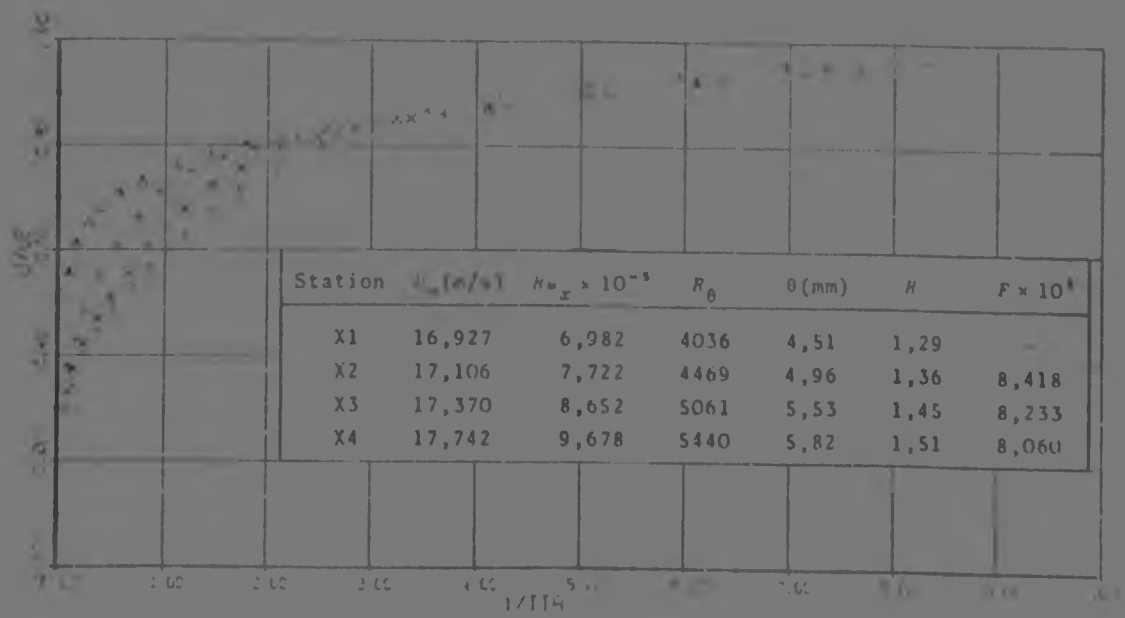


FIG 7.10x : TEST 30N1 E-8,56-17 80% BLOCKED

PLOT OF U/V<sub>w</sub> VS Y/T<sub>TA</sub>

Run 01 12/1, 4, 5, 17

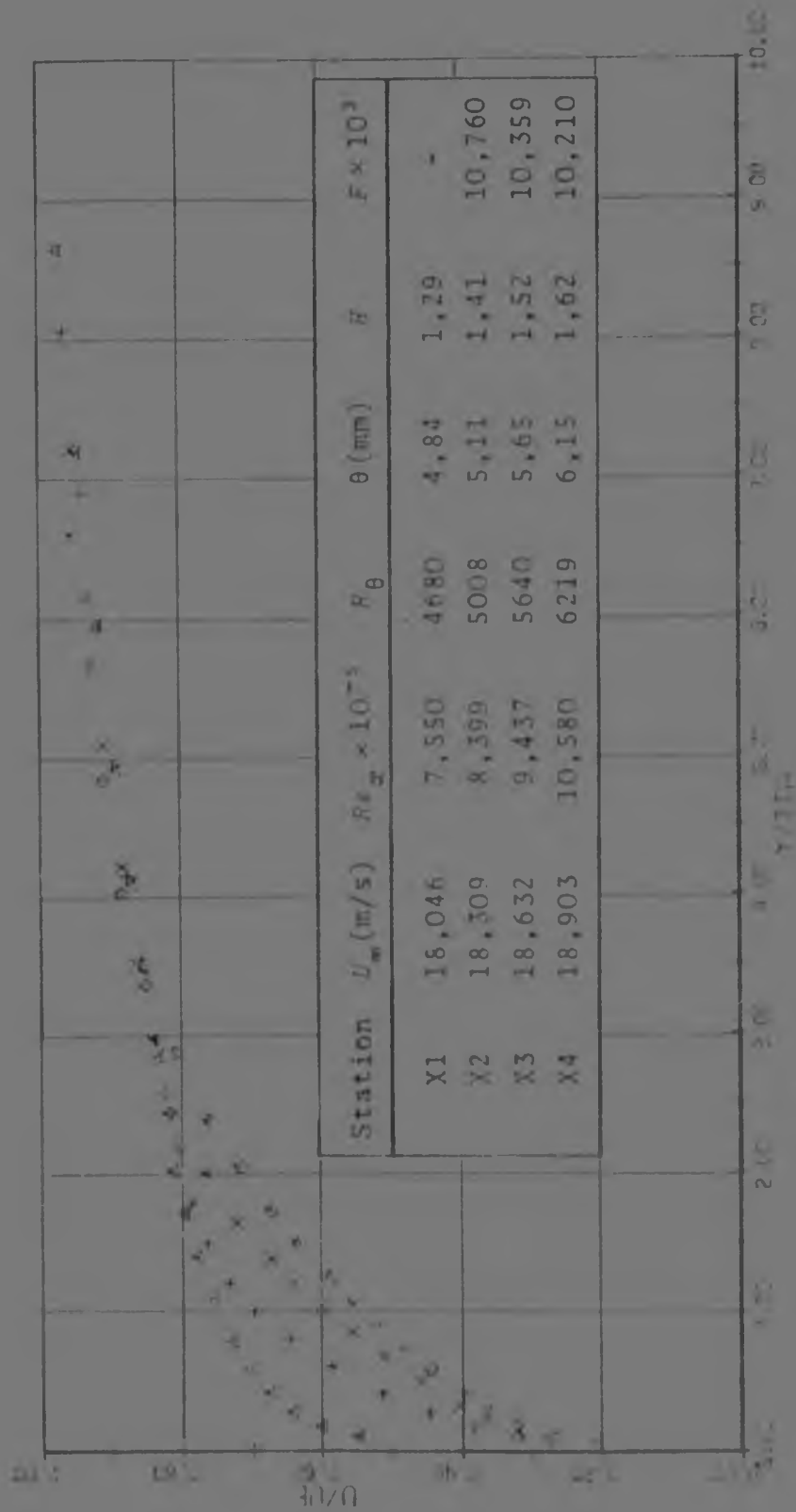


FIG 7.10y : TEST 01D1 E-11,66-17 80% BLOCKED

E and C grade plates were made by sintering spherical particles. The surface of the A grade plate was therefore much smoother, and the injection pattern was expected to be different from that experienced using grades E and C. (See Fig 7.11 and Appendix B).

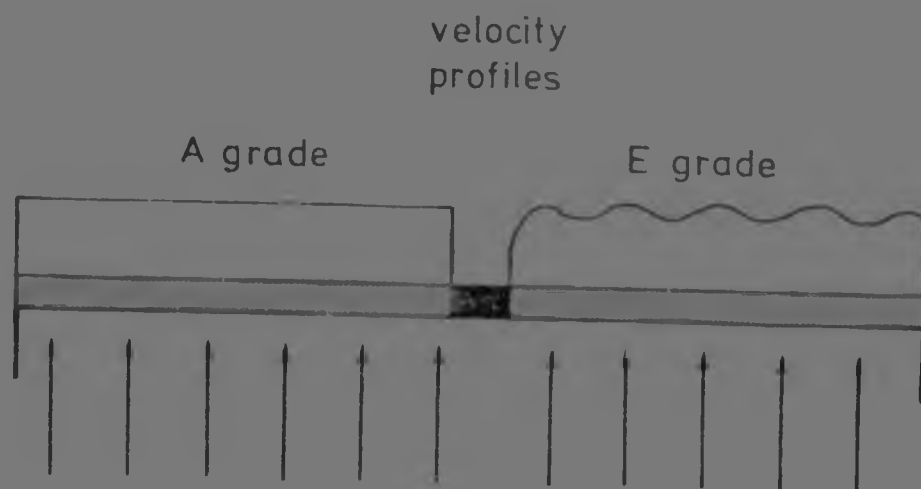


FIG 7.11 :  $v_w$  PROFILES EXPECTED FOR A AND E GRADE PLATES

Profiles obtained for the A grade plate and for test 28S1 were virtually identical. Note that for the former,  $F_{av} = 0,00755$ , for the latter  $F_{av} = 0,00758$ . The input profiles were the same as well. It was concluded that the velocity profile was mainly a function of injection ratio.

Having compiled a set of data for the unblocked porous plate, a few tests were done in which the lower porous surface was partially sealed with masking tape. The first tests were run with 50% blockage, as is shown in Fig 7.12.

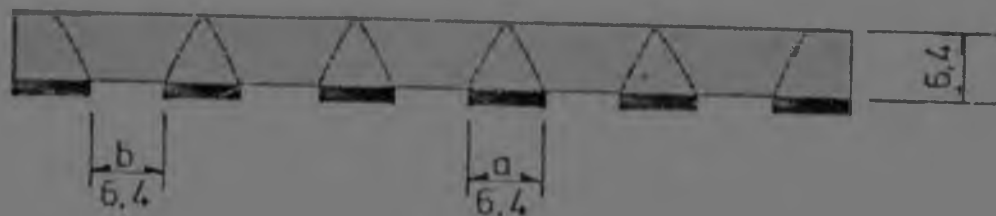


FIG 7.12 : SKETCH SHOWING THE PARTIALLY BLOCKED POROUS PLATE, 50% BLOCKED

$u_e = 17 \text{ M/S}$

E GRADE ( 80 PERC. BLOCKED )

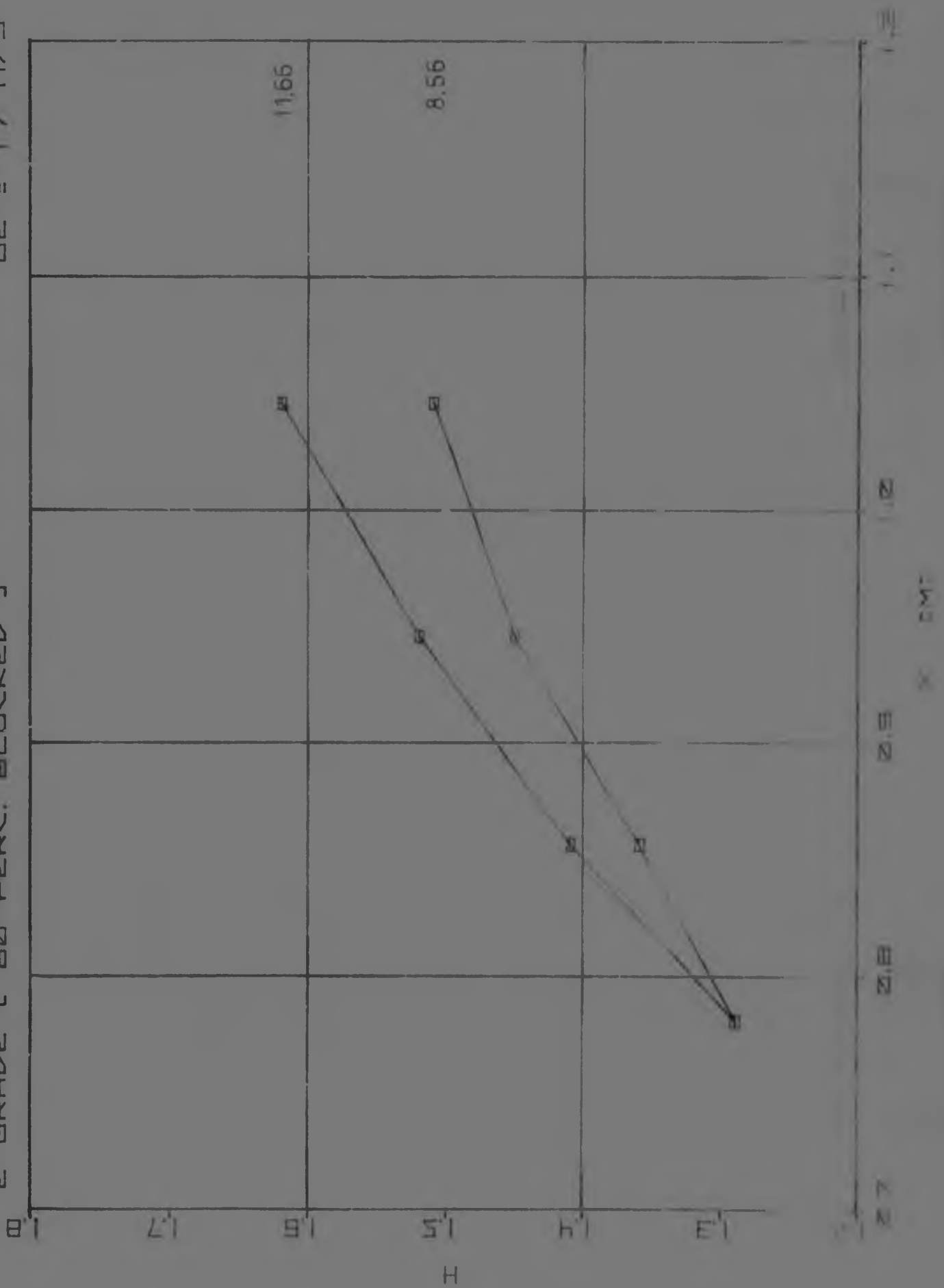


FIG 7.13d : PLOT OF  $H$  VS  $x$  FOR  $u_e = 17 \text{ m/s}$ , 80% BLOCKED

U<sub>∞</sub> = 17 M/S

U<sub>∞</sub> = 17 M/S

3 GRADE ( 80 PERC. BLOCKED )

3 GRADE ( 67 PERC. BLOCKED )

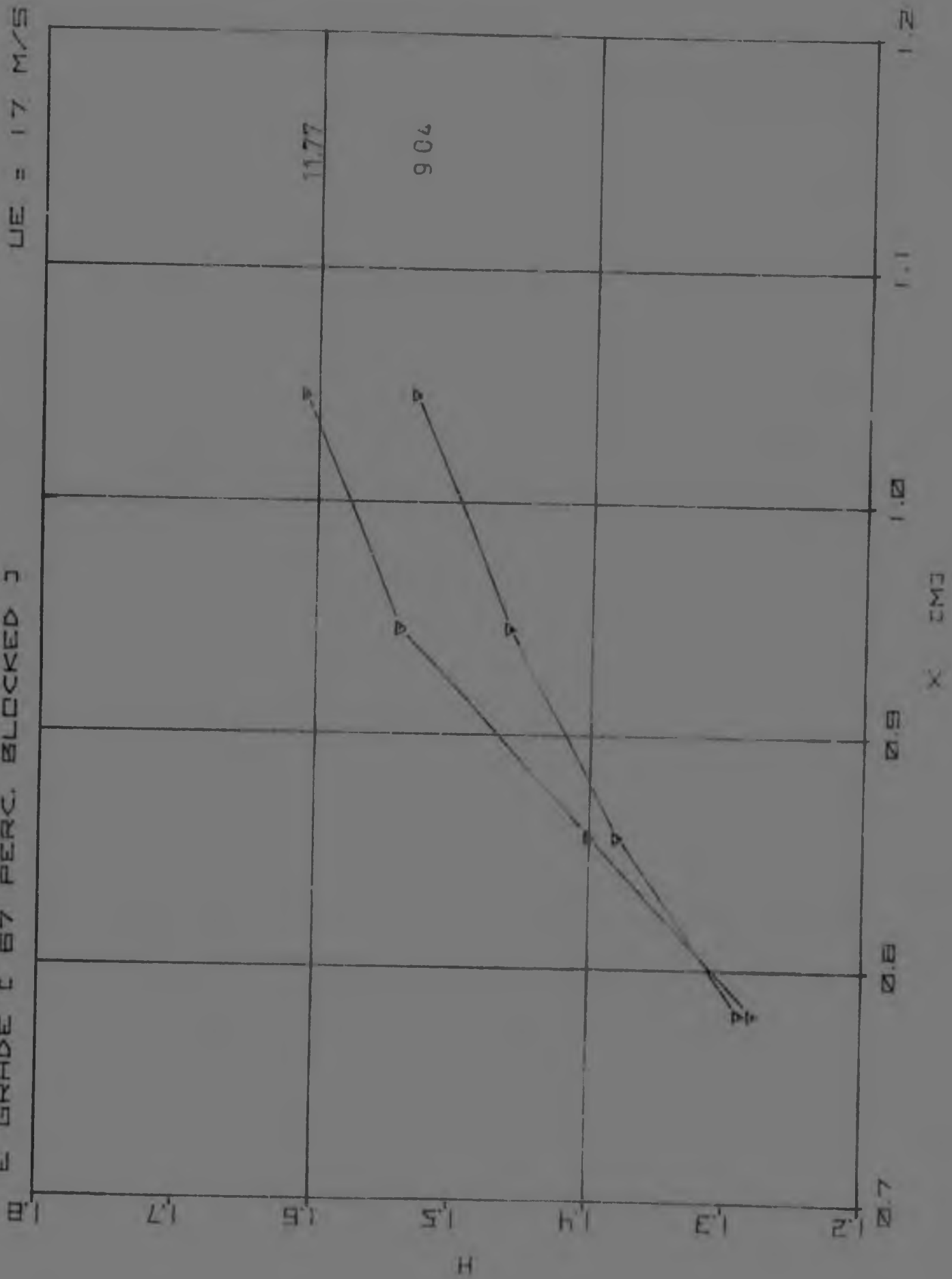


FIG 7.13c : PLOT OF H vs x FOR u<sub>∞</sub> = 17 m/s, 67% BLOCKED

Fig 7



U<sub>10</sub> = 17 m/s

U<sub>10</sub> = 17 m/s

50% GRADE (50 PERC. BLOCKED)

H VS X

50% GRADE (50 PERC. BLOCKED)

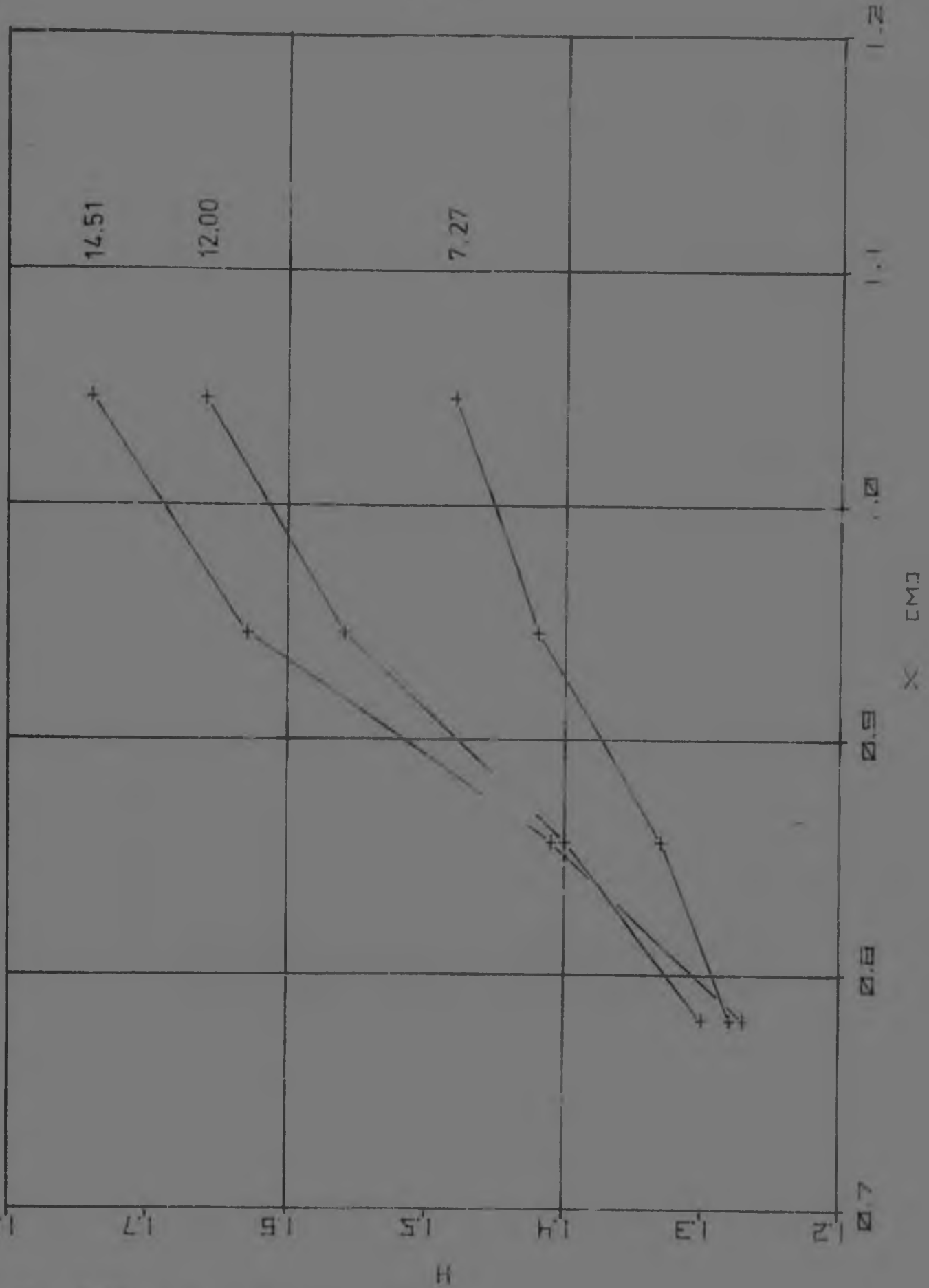


FIG 7.13b : PLOT OF H VS X FOR U<sub>10</sub> = 17 m/s, 50% BLOCKED

$u_e = 17 \text{ m/s}$

$H$  vs  $x$

UNBLOCKED

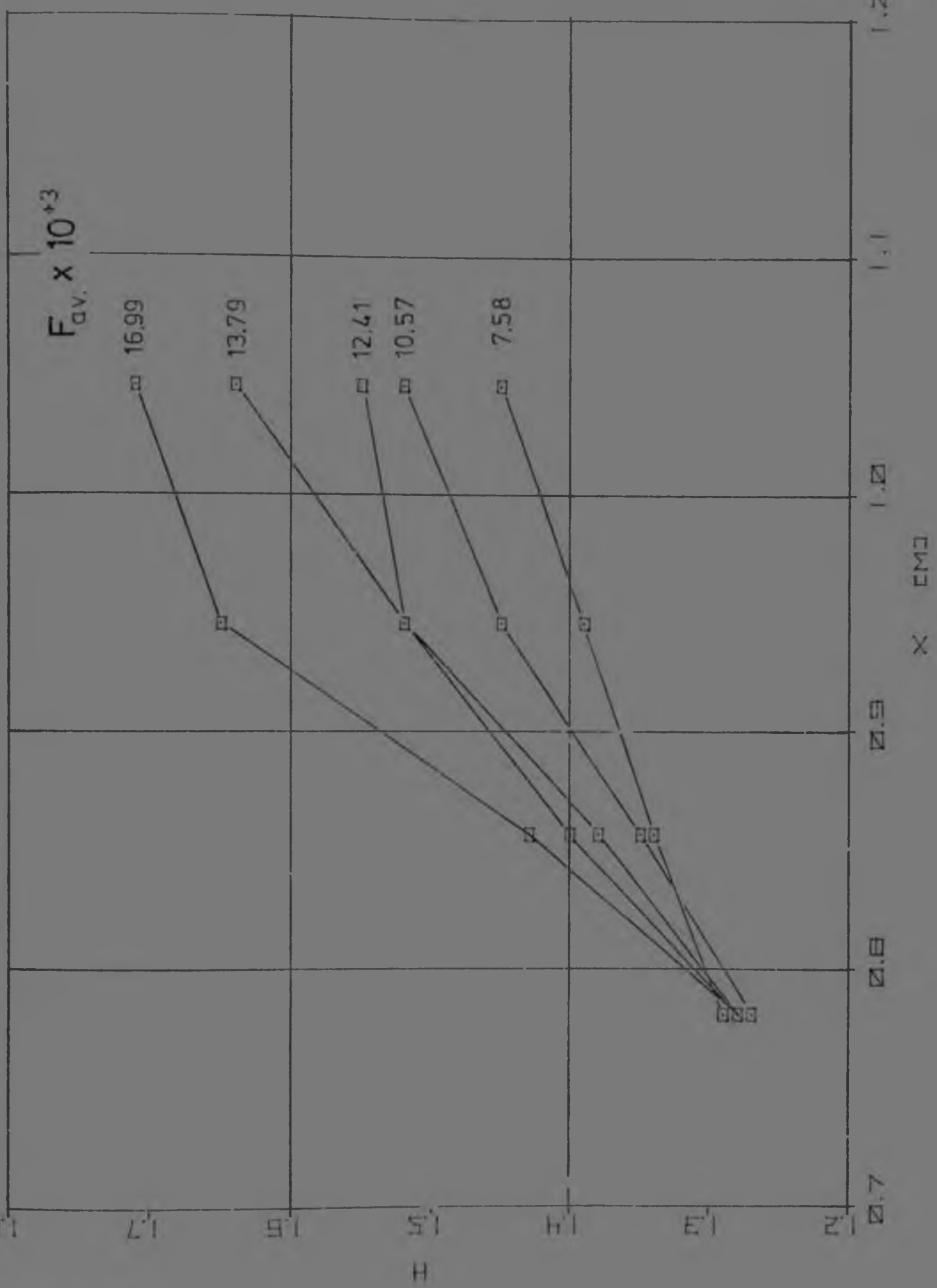


FIG 7.13a : PLOT OF  $H$  vs  $x$  FOR  $u_e = 17 \text{ m/s}$ , UNBLOCKED

Tests 23N1 to 24N1 were run at  $U_m = 17$  m/s, 50% blockage of the E grade plate. See Fig 7.10s to u.

With the tape covering 67% of the plate,  $a = 12,7$  mm and  $b = 6,35$  mm, tests 28N1 and 29N1 were performed. These profiles are presented in Fig 7.10v and w.

Two tests were run with  $a = 12,7$  mm,  $b = 3,18$  mm; this resulted in 80% blockage. Fig 7.10x and y are the profiles for these tests.

The effect of blockage is shown in Fig 7.13. This plot of shape factor,  $H$  vs  $x$ -station, was thought to present the change in a more succinct way than the usual velocity profile of Fig 7.10. The injection ratio,  $F$ , became a very average parameter. The definition of  $F$  now becomes :

$$F = \frac{Q_{\text{trans}}}{A_{\text{upper}} U_{\infty}} \quad \text{for } \rho \text{ constant} \quad \dots [7.3.1]$$

Where  $Q_{\text{trans}}$  was the transpiration flow rate and  $A_{\text{upper}}$  the area of porous plate exposed to the main stream.

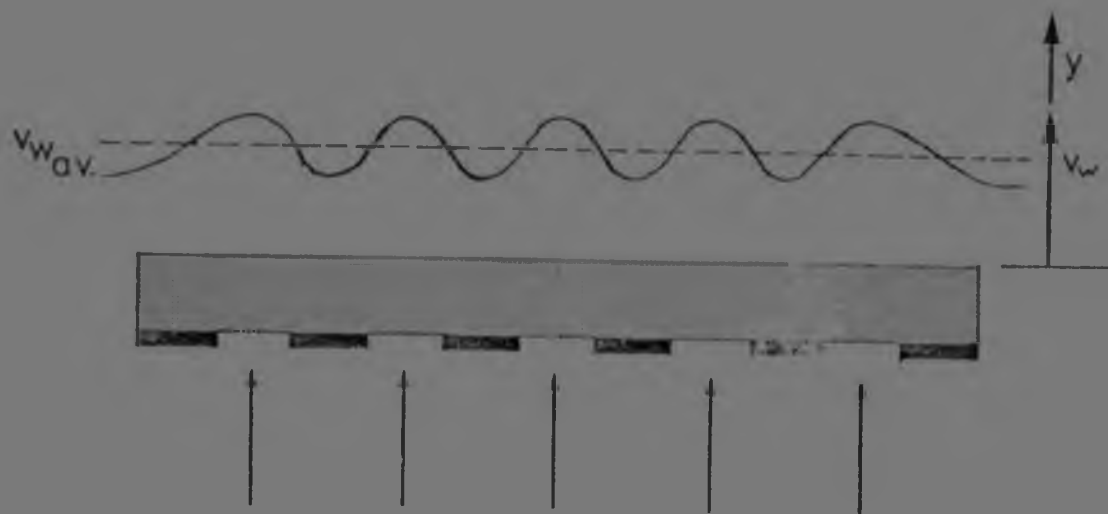


FIG 7.14 : DISTRIBUTION OF  $v_w$  WITH  $x$  FOR THE PARTIALLY BLOCKED POROUS PLATE, AND THE SPREAD OF THE FLOW IN THE POROUS MEDIUM WITH  $y$

By considering Fig 7.13, it is seen that the curves do not increase with  $F$  as they do for the unblocked plate. The profile is developing for an apparently higher  $F$ . Clearly  $v_w$  was alternately above and below the average transpiration velocity. The average momentum in the  $y$ -direction must remain the same as for the unblocked porous plate, however, some of the injected flow penetrates deeper into the shear layer. Thus it is concluded that with the same transpiration flow rate, blockage results in an apparent increase in  $F$  and  $v_w$ . If the Stanton number is expected to be a function of  $v_w$  (see Moffat and Kays, 1968), a decreased heat transfer coeff. would become possible with blockage.\* Intuitively, the opposite may be expected, however, only experiment is expected to resolve the matter.

It would therefore appear that the blockage caused by strengthening ribs below the porous surface, could result in improved performance.

7.3.2 The bilog law plotting technique

This method for determining  $v_w$  from the velocity profile was discussed in Chapter 2, section 3, at some length. It was developed by Black and Sarnecki (1958), and yielded an effective imaginary shear velocity at the wall. This was merely a technique, and did not imply reverse flow, as may be encountered for flows in an adverse pressure gradient. The results are given in Fig 7.15a to c, Fig 7.16 and Fig 7.17. These graphs are presented because (i) Fig 7.15 showed the trend in the profiles for increasing  $F$ , and (ii) the graphs allowed an interesting comparison of the present results with those of Mickley and Davis (1956).

Fig 7.15a to c are graphs of  $u/u_e - y^2$  vs  $y$ , for the basic data of Mickley and Davis (1956), and the present work. Fig 7.15a is for  $U_e = 17$  m/s, E grade. This graph also shows the limiting curves, ie. for  $u/u_e = 0$  and  $u/u_e = 1$ , ie. the edge of the boundary layer. Plots shown in

\* see Addendum for further explanation.

$U_i = 17 \text{ M/S}$      $C_x \ U_i = 13 \text{ M/S}$      $U/U_i - Y_i/Y_i \text{ VS } Y_i$      $C \text{ AND } A \text{ GRADE UNBLOCKED}$

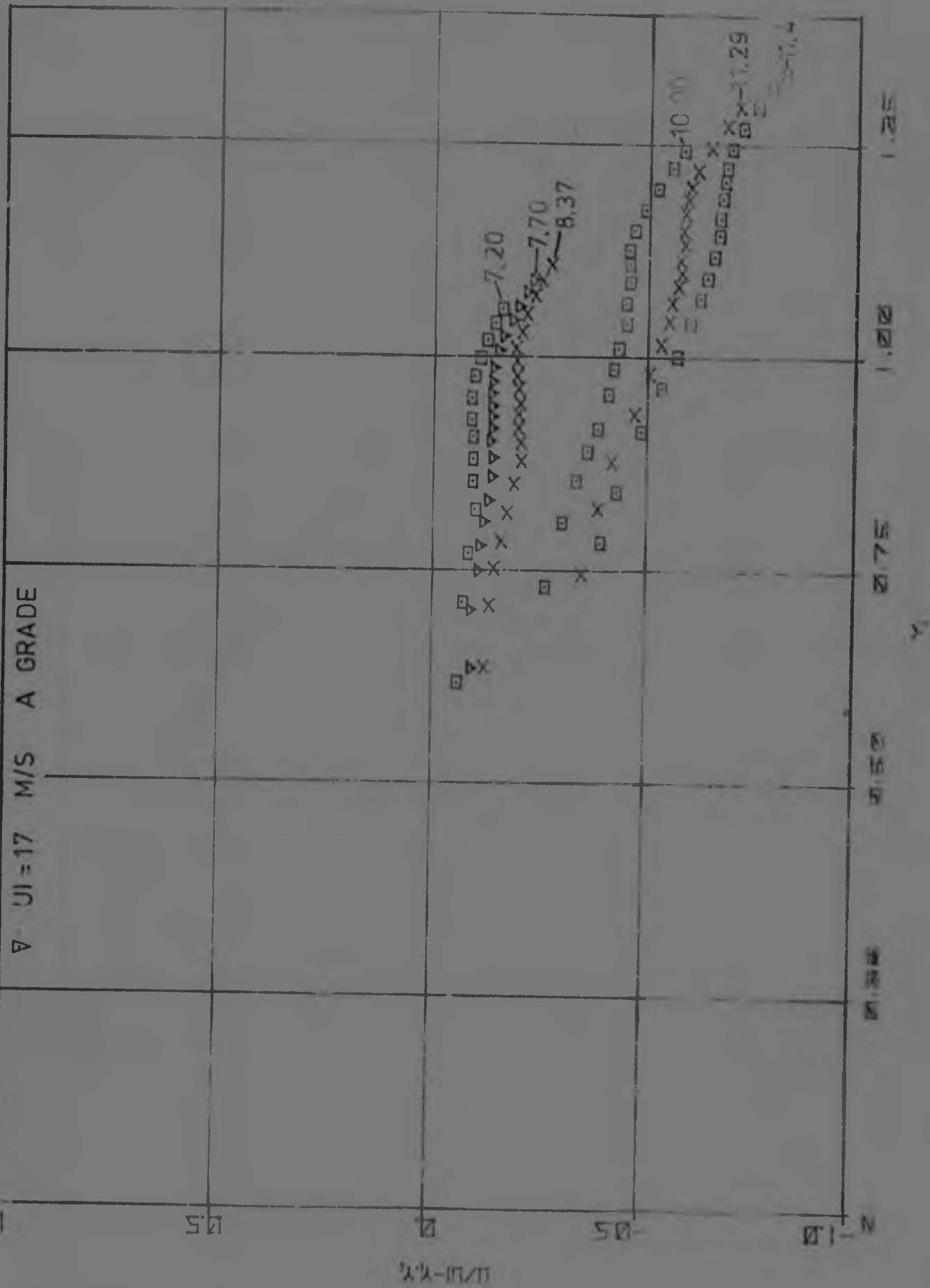


Fig 7.15c : SOME EXPERIMENTAL RESULTS PLOTTED ON THE BILOG CO-ORDINATES

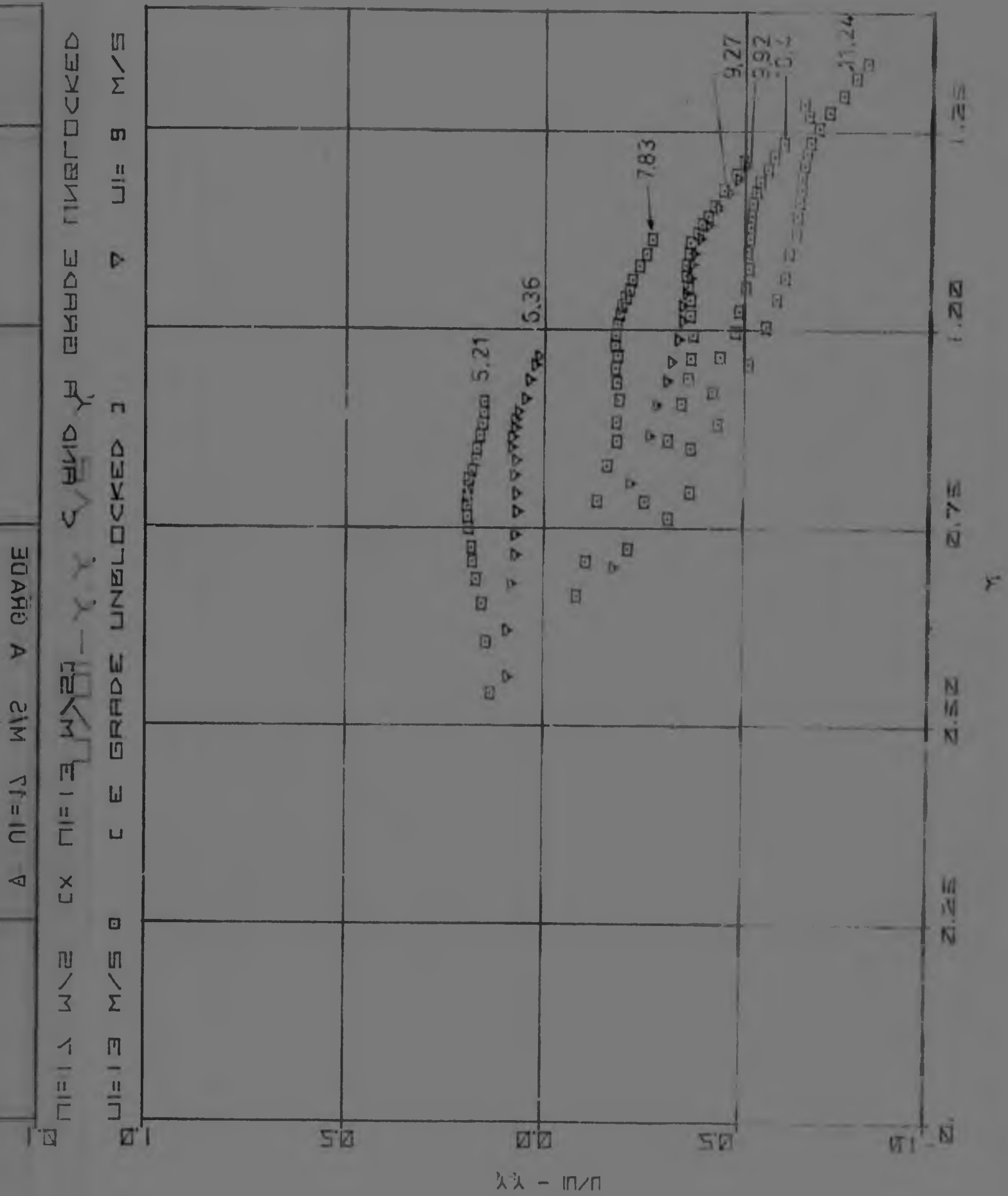


FIG 7.15B : SOME EXPERIMENTAL RESULTS PLOTTED ON THE BILOG CO-ORDINATES

017

A. NI = 1.2 M/S V GRADE

NI = 1.3 W/2 0.2 M/S GRADE UNBLOCKED 3 + MICKLEY

NI = 1.7 M/S

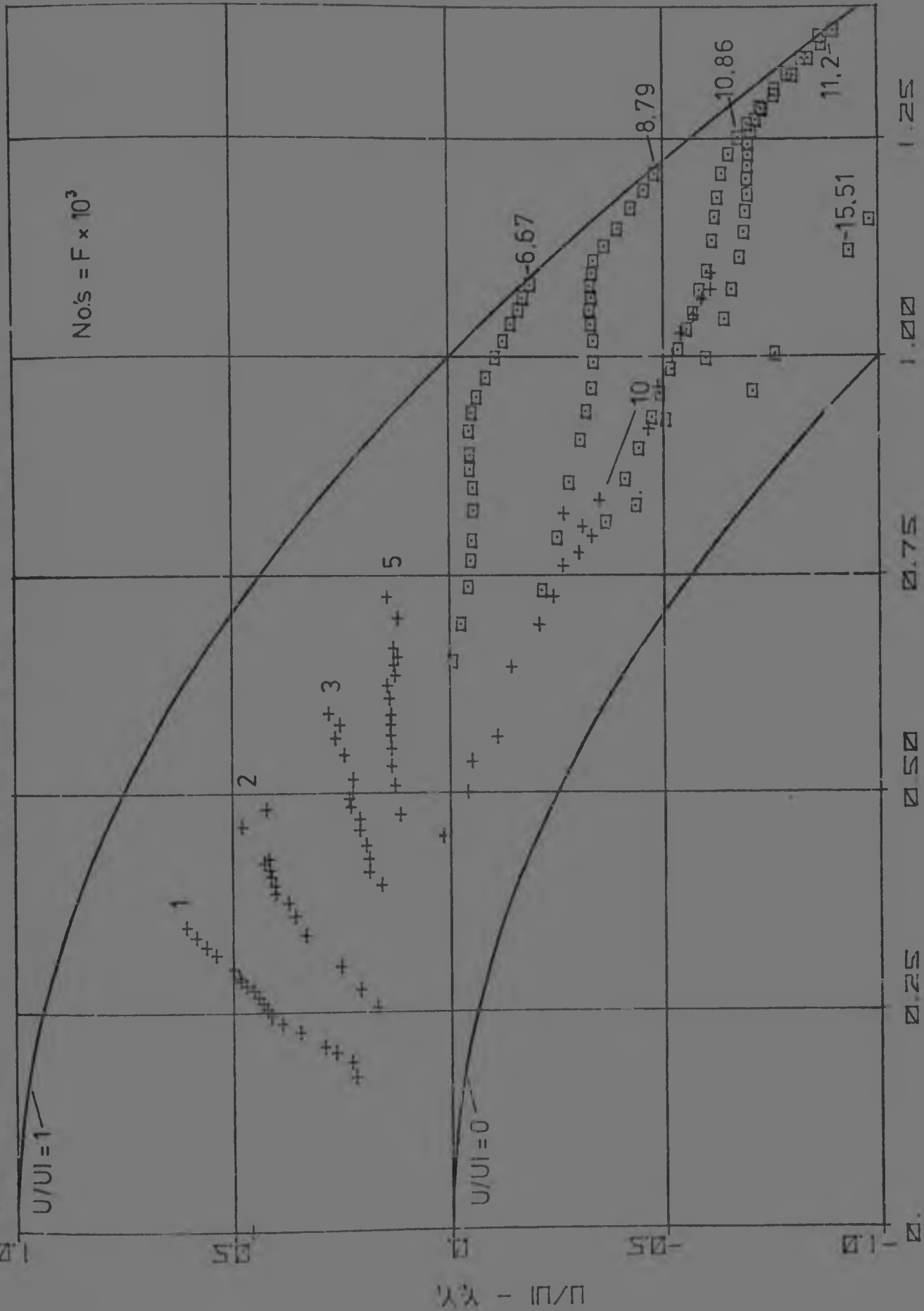


Fig 7.15A : PRESENT DATA AND THAT OF MICKLEY AND DAVIS (1956) PLOTTED ACCORDING TO THE BILOG LAW TECHNIQUE

Fig 7.15A

PLOT OF  $U^+$  VS  $\log Y^+$

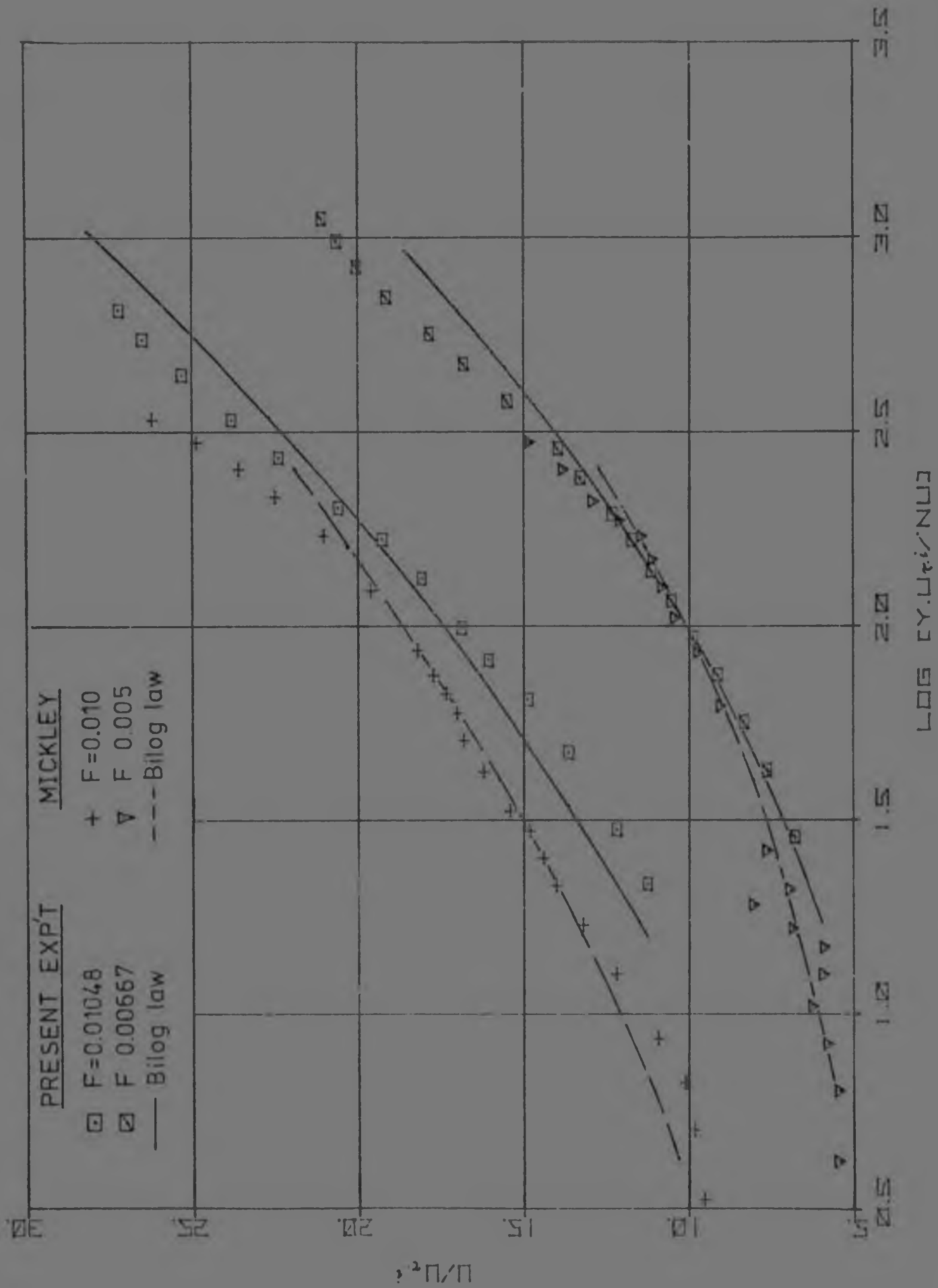


FIG 7.16 : EXPERIMENTAL DATA AND BILOG PREDICTED PROFILE



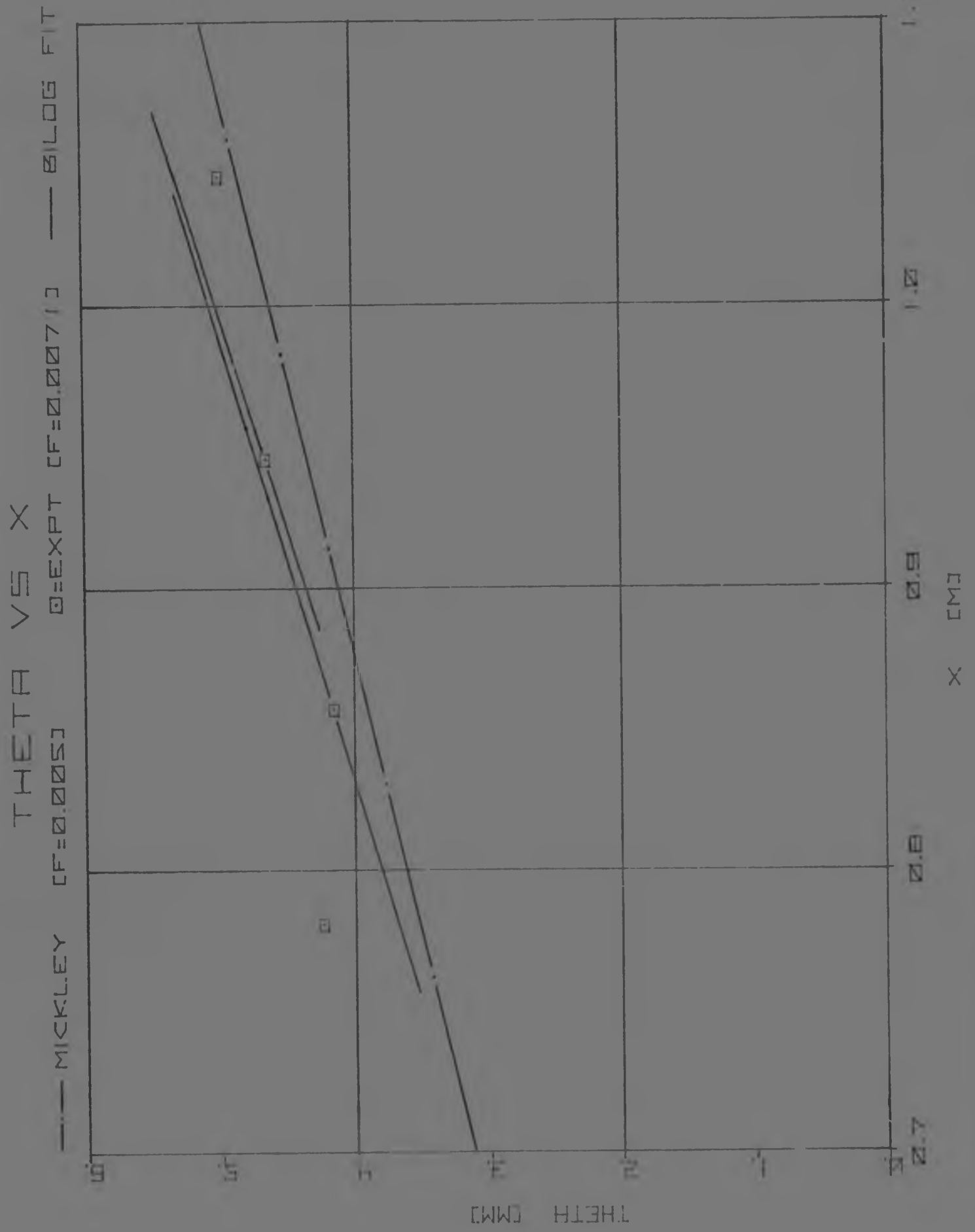


FIG 7.17 : GRAPH OF  $\theta$  vs  $x$  - EXPERIMENTAL AND BILOG PREDICTED SLOPE

Fig 7.15b are for  $U_\infty = 13$  m/s ( $\ominus$ ) and  $U_\infty = 9$  m/s ( $\nabla$ ), both sets on the E grade plate. Results obtained for  $U_\infty = 17$  m/s ( $\omin�$ ),  $U_\infty = 13$  m/s ( $\times$ ) on the C grade plate, and for the single run using the A grade plate ( $U_\infty = 17$  m/s ( $\nabla$ )) are plotted on Fig 7.15c.

By fitting a straight line to the experimental points of a run (in the bilog region),  $u_\tau$  could be calculated from the slope and ordinate intercept of that line (see equations 3.3.21 to 3.3.23). If this line did not pass through the limiting line of  $u/u_\tau = 0$ ,  $u_\tau$  would be imaginary. This was the case for all values of  $F \geq 5$ . The uncertainty in fitting the straight line can be seen from the scatter of the points. Two examples are shown in the non-dimensional plot of  $u/u_\tau$  vs  $\log y \cdot u_\tau / \nu$  in Fig 7.16. Again comparable results of Mickley and Davis (1956) are included. The lines represent the profiles predicted by the bilog law.

Because it was of importance to know  $d\theta/dx$ , it was thought useful to show that the bilog technique could predict this derivative, but gave no indication of the origin of the boundary layer. Fig 7.17 shows a typical example for  $F = 0,0071$ . The slope was calculated at a point, and a straight line with that gradient drawn through the point. Good agreement was again obtained with the result of Mickley and Davis (1956).

It is felt that as this method yielded imaginary shear, it is only useful as a plotting technique, and is inadequate for estimating  $c_f$ . The graphs did allow comparison of the present work with that of other researchers.

### 7.3.3 A two-parameter family of curves

The two-parameter family of curves of Thompson (1965) was used to construct the mean velocity profile ahead of the porous wall. McQuaid (1967) suggested a family of curves for the transpired boundary layer. This set of curves

predicted the velocity profile and the skin friction. Intermittency was included in the formulation (by the term  $\delta_s$ , see eqn.[3.3.30]). This author showed that  $u/U_\infty$  was a function of  $y/\delta_s$ ,  $\alpha_f$ ,  $R\delta_s$ ,  $v_w/U_\infty$ , which was transformed to eqn.[3.3.29]. This was a function of  $y/\theta$ ,  $H$ ,  $R_\theta$ ,  $v_w/U_\infty$ .

A cross-plot of the present results was attempted, and is shown in Fig 7.18. These values were obtained from the velocity profiles (Fig 7.10) by keeping  $u/U_\infty$  constant, and reading  $y/\theta$  for the various  $x$  stations (effectively  $Re_x$ , or  $Re_\theta$ ). It was thought that for constructing mean velocity profiles, it would be more useful to have  $u/u_0$  as one of the axes. Thus  $y/\theta$  was kept constant, and the value of  $u/u_0$  read from Fig 7.10. These values were plotted, and a number of points from the experimental work of Mickley and Davis (1956) and McQuaid (1967) added. A series of curves were obtained, apparently independent of  $v_w/U_\infty$ , and are presented as Fig 7.19 a to d.

This result may be stated :

$$H = H(R_\theta, v_w/U_\infty) \dots\dots [7.3.2]$$

$$u/U_\infty = f(y/\theta, H) \dots\dots [7.3.3]$$

A limitation of this family is that  $e_f$  is not accounted for. McQuaid (1967) presented a skin friction law, for  $v_w/U_\infty$  up to 0,010 :

$$e_f = e_f(H, R_\theta, v_w/U_\infty) \dots\dots [7.3.4]$$

These curves were assumed to hold for a porous wall, continuous from  $x = 0$ . It was thought that the values of  $e_f$  would be very different for discontinuous injection. However, it must again be emphasised that only direct measurement of  $v_w$  would give reliable values of  $e_f$ . This would be necessary in each individual test, for its particular input profile, porous plate configuration, surface roughness, pressure gradient, etc.

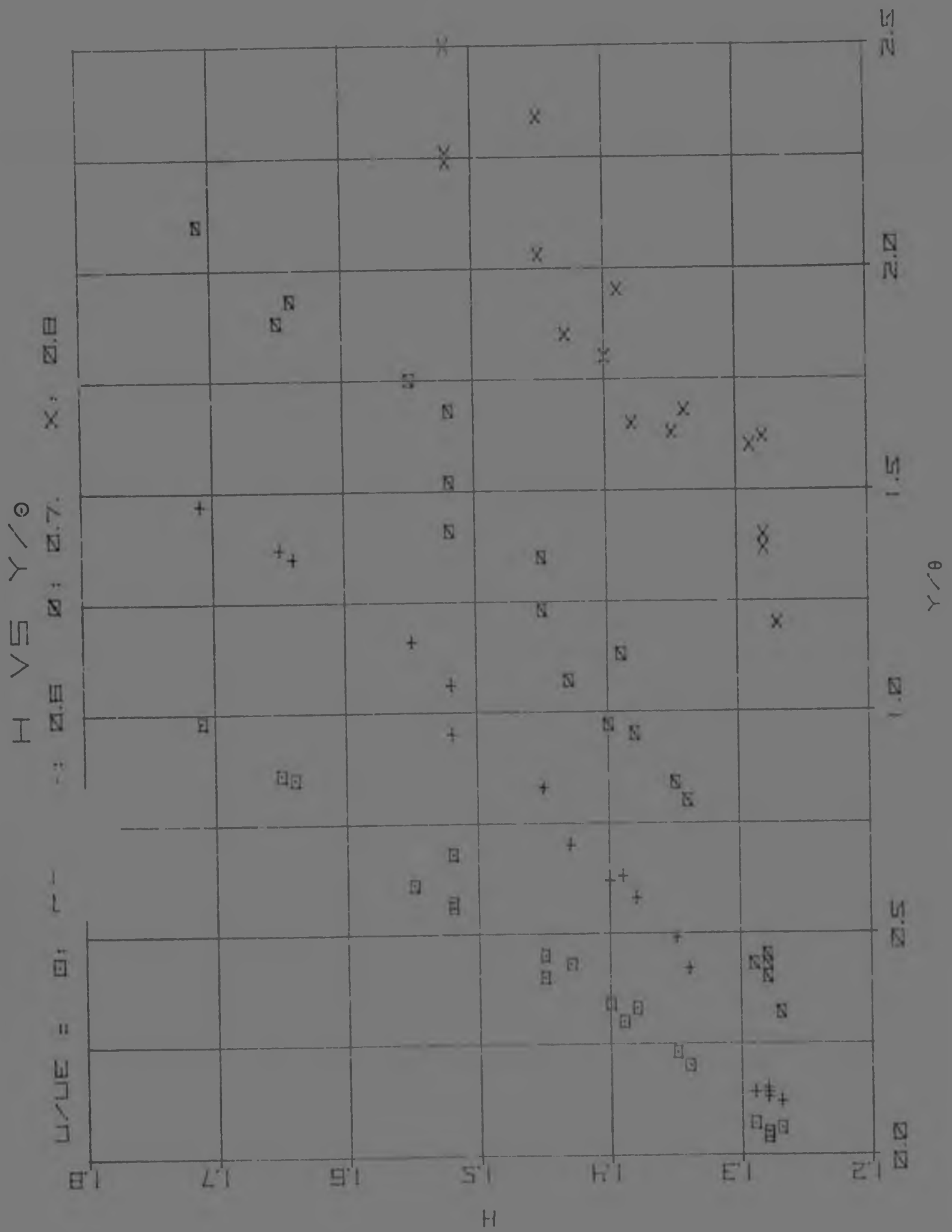


FIG 7.18 : CROSS PLOT OF PRESENT RESULTS,  $u/u_e$  CONSTANT

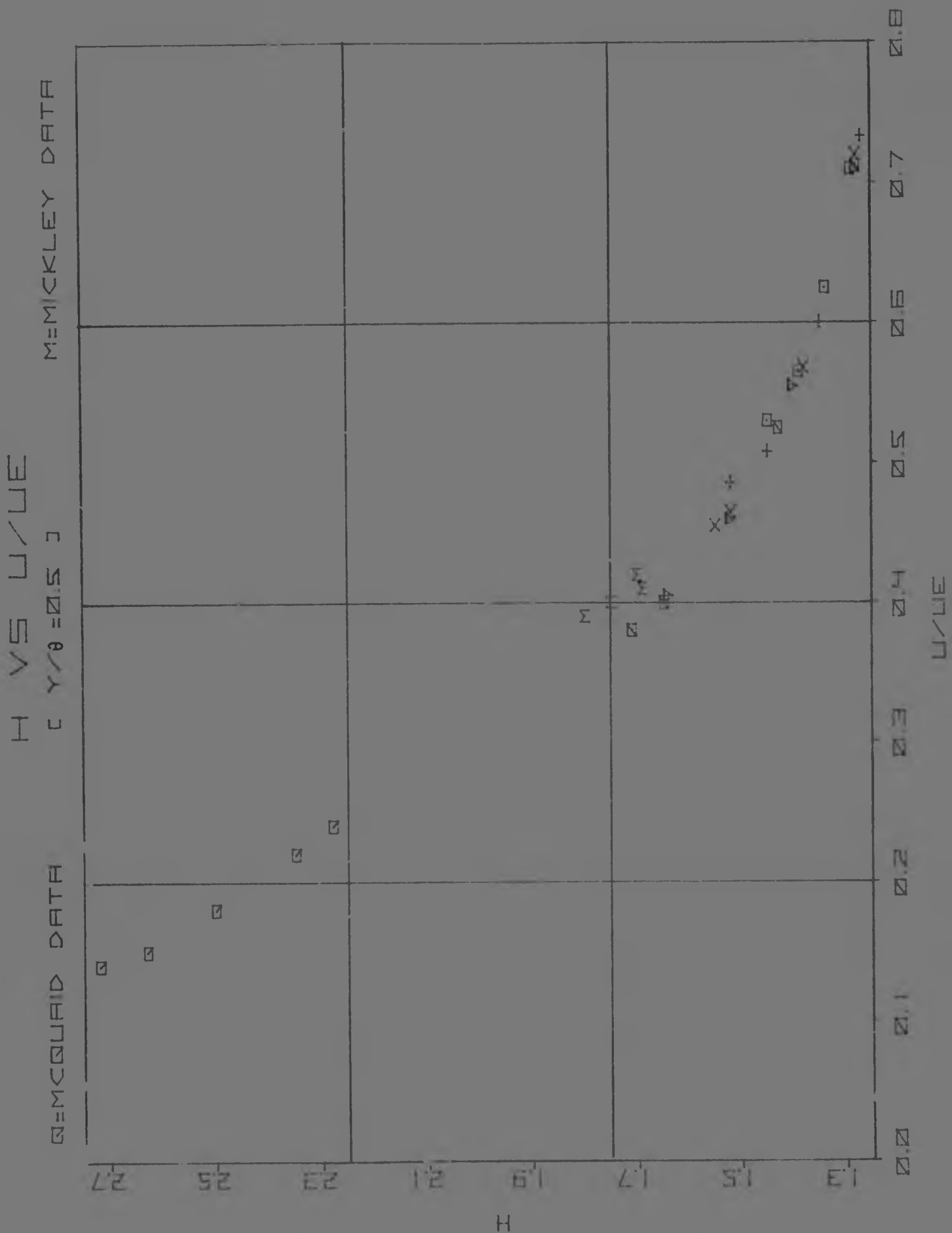


FIG 7.19A : CROSS PLOT OF EXPERIMENTAL DATA,  
FOR  $y/\theta$  CONSTANT

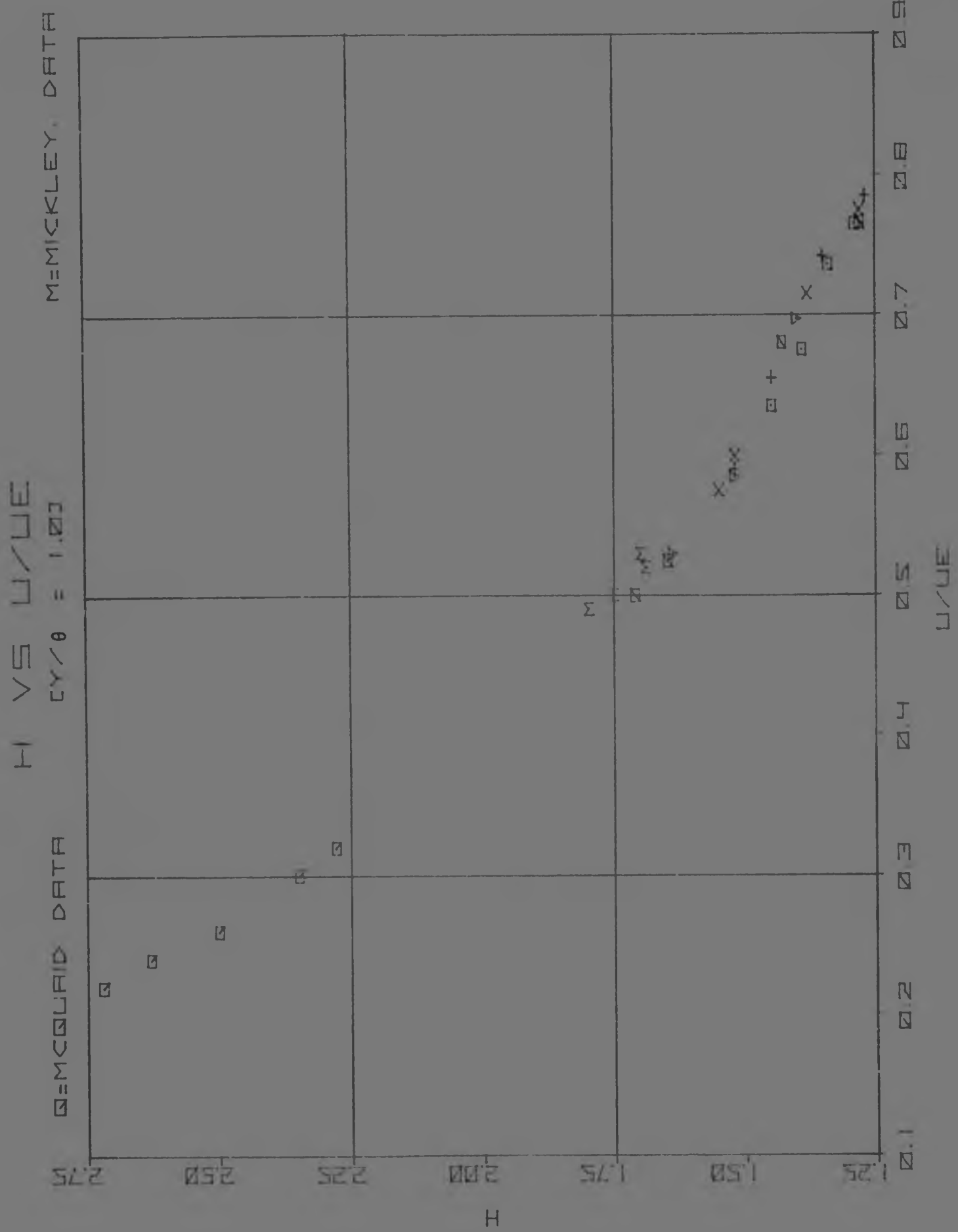


FIG 7.10B : CROSS PLOT OF EXPERIMENTAL DATA,  
FOR  $\gamma/\theta$  CONSTANT



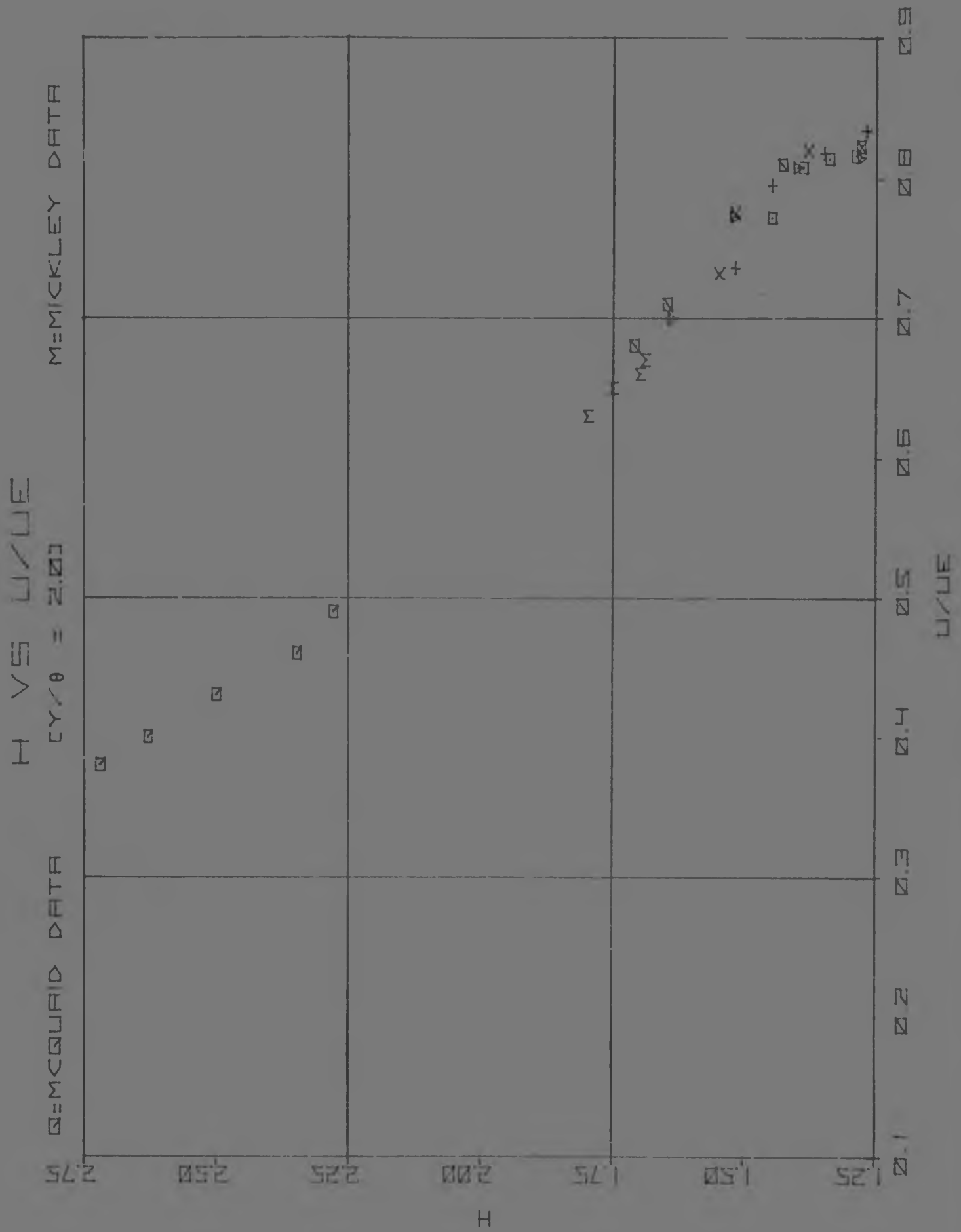


FIG 7.19D : CROSS PLOT OF EXPERIMENTAL DATA,  
FOR  $\gamma/\theta$  CONSTANT



#### 7.4 Computer Programme Results

As discussed previously, the programme required two initial profiles to start the solution. These could be either laminar or turbulent, but with  $v_w = 0$ . Starting the solution with transpiration from the boundary layer origin was therefore not possible. By installing a trip wire, it is possible to obtain a turbulent profile for very low values of  $Re_x$  experimentally. However, the equation used to generate the turbulent profile was valid for  $Re_x > 5 \times 10^4$ . Thus it was not possible to simulate experimental conditions very effectively.

The boundary layer thickness was over-estimated because the equation used was valid only to  $y^+ = 10^4$  (depending on Reynolds number). This is illustrated below in Fig 7.20.

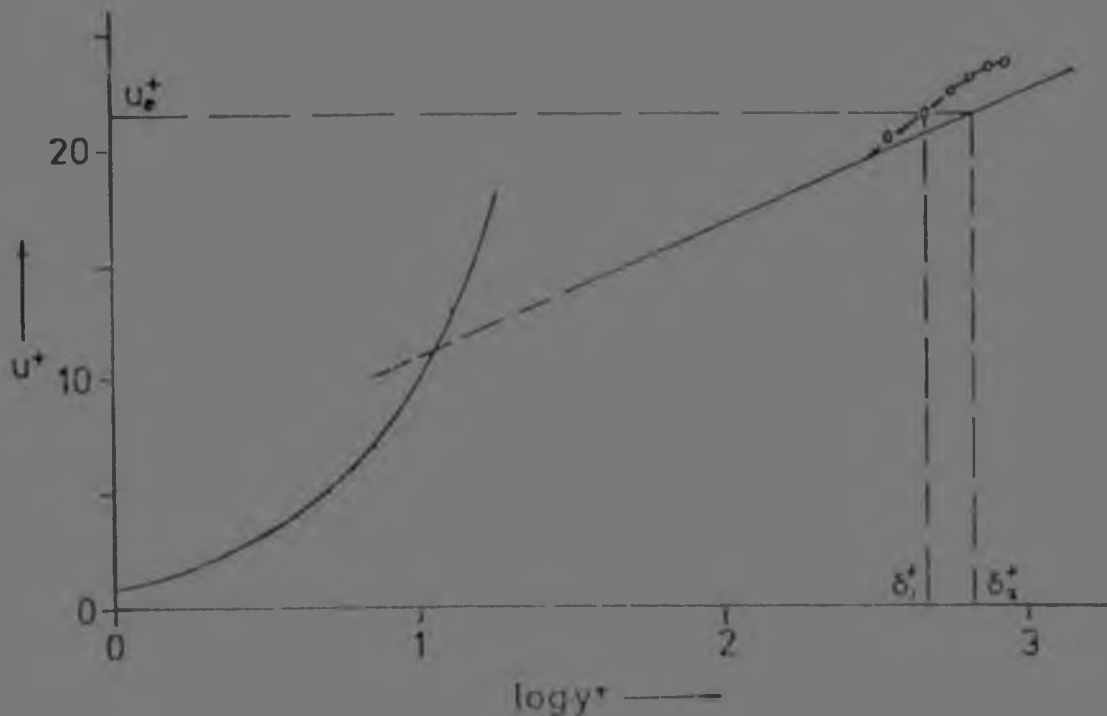


FIG 7.20 : DETAIL OF SEMI-LOG PLOT OF TURBULENT PROFILE, SHOWING ERRONEOUS VALUE OF  $\delta$

Clearly, if the straight line was used,  $\delta^+ = \delta_2^+$ , instead of  $\delta_1^+$ , obtained from experimental points. If a wake law was not used in this region, a power law would probably be the easiest method of generating  $u^+$  near the edge of the layer.

#### 7.4.1 Main computer runs

The results of the main computer runs are shown graphically in Fig 7.21 a to j. See also Appendix F3 for tabulated values of  $R_{\theta}$ ,  $H$ , etc.

By generating the input profiles on a flat plate, the programme was a close simulation of the experimental work. It was, however, not possible to obtain the identical input profile that was measured experimentally without extensive re-programming - the Van Driest profile was used instead. Furthermore, the programme as written is for the 'flat plate in an infinite stream', whereas almost channel flow was obtained in the tunnel, ie. the boundary layer growth in the tunnel was limited to the centre line of the tunnel.

Examination of Fig 7.21 shows that the profiles generated are very similar to those depicted in Fig 7.10. The trends with  $x$  and  $F$  were identical, ie. increased 'laminarisation' with  $x$  and  $F$ .

The values of  $\sigma_f$  appear to be highly erroneous. On the other hand, there is no experimental data for the step change in transpiration and surface roughness with which these values could be compared. If  $\phi''$  is obtained from  $\phi$ , negative  $\sigma_f$  results, and hence imaginary  $u_0^+$ ,  $p^+$  and  $v_0^+$ . Although this is predicted by the bilog law, this is not of much use, as the programme cannot solve the equations for imaginary shear velocity.  $\phi''$  was therefore obtained from  $\phi'$ , and it is expected that the error in  $\phi''$  does not affect the solution significantly (Note:-  $\phi''$  affects the inner eddy-viscosity equation). Clearly the finite difference equation which yields satisfactory values of  $\sigma_e$  for the unblown layer was inadequate for the blown boundary layer. By decreasing the step size more accurate results would be obtained, but this would increase the cost of running the programme significantly. In this work,  $H=0,065$  was common, whereas other investigators have reported  $H=0,001$ , with 300 steps.

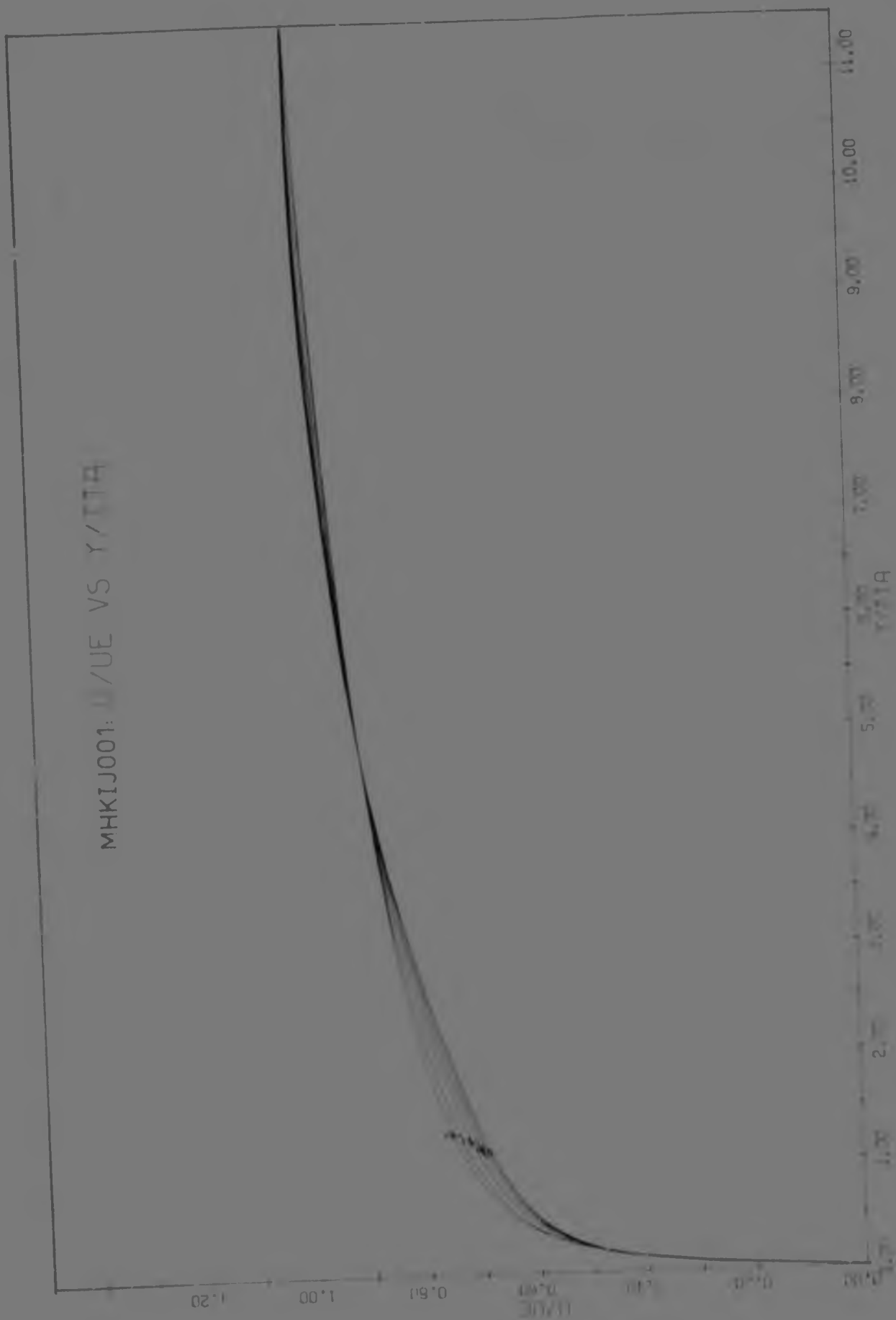


FIG 7.21A : TURBULENT PROFILES WITH INCREASING INJECTION,  $U_w = 17$  m/s,  $F = 0,001$

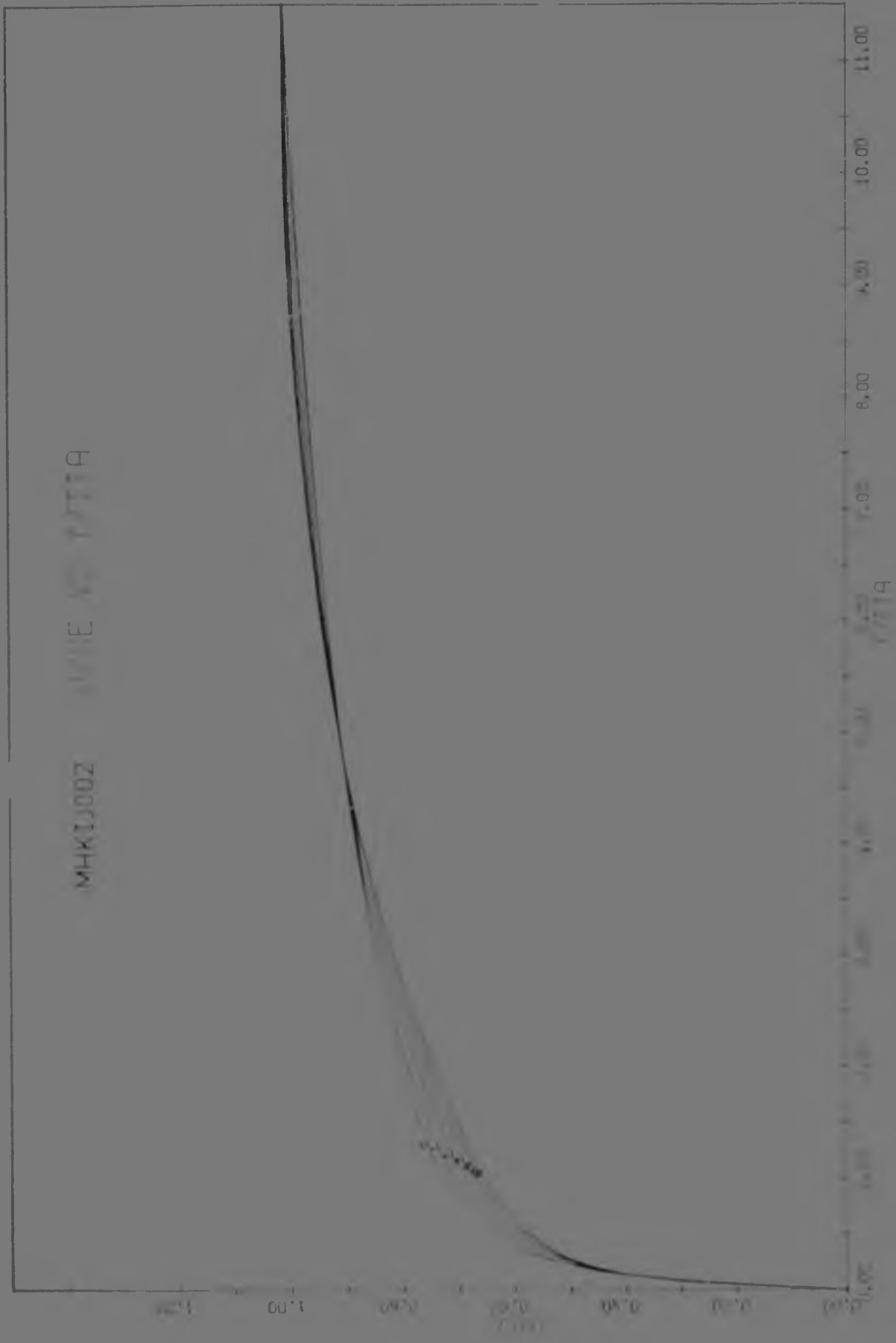


FIG 7.21B : INJECTED TURBULENT PROFILES FOR  
 $U_w = 17 \text{ m/s}$ ,  $F = 0,002$

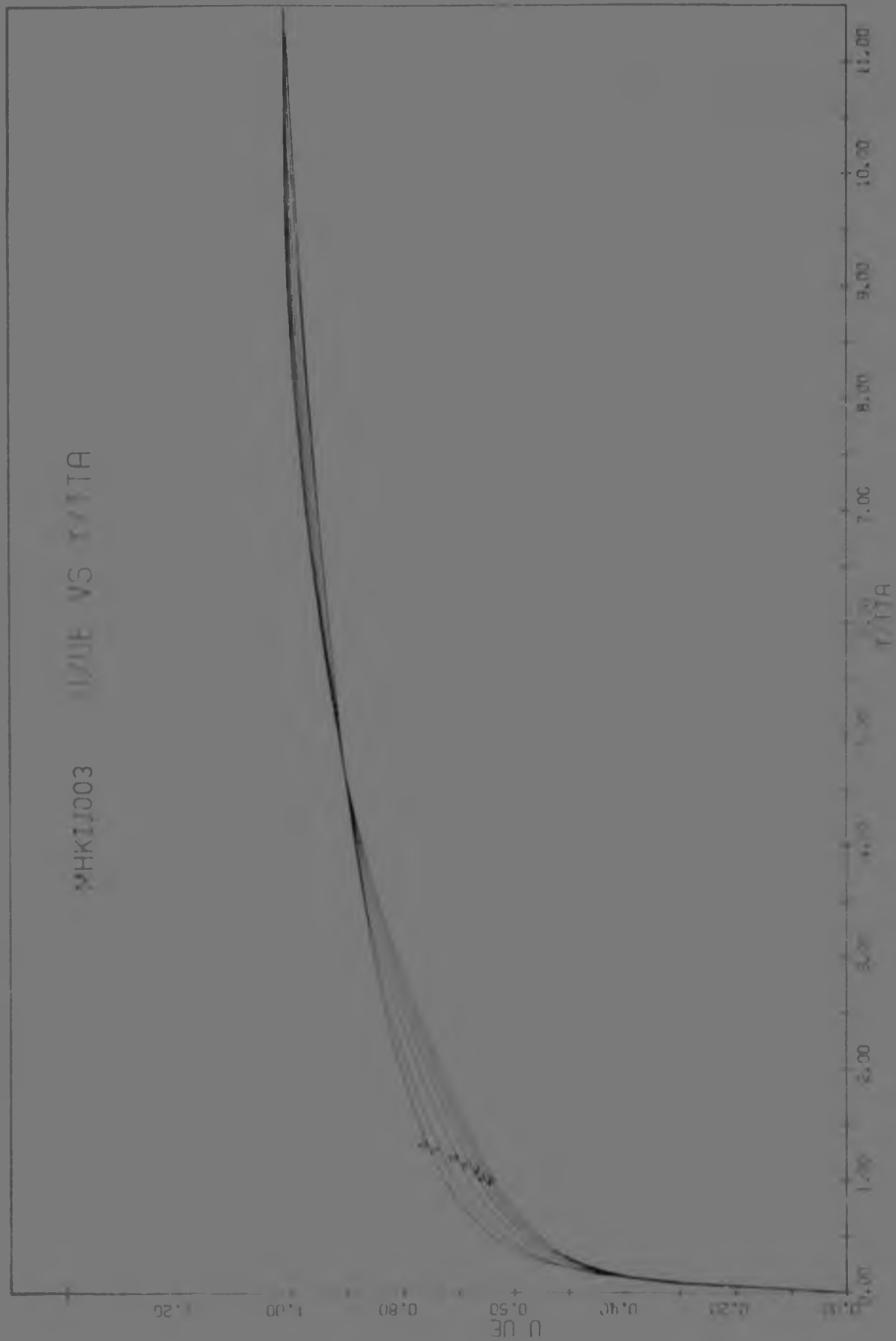


FIG 7.2]c : INJECTED TURBULENT PROFILES FOR  
 $U_e = 17 \text{ m/s}$ ,  $F = 0,003$

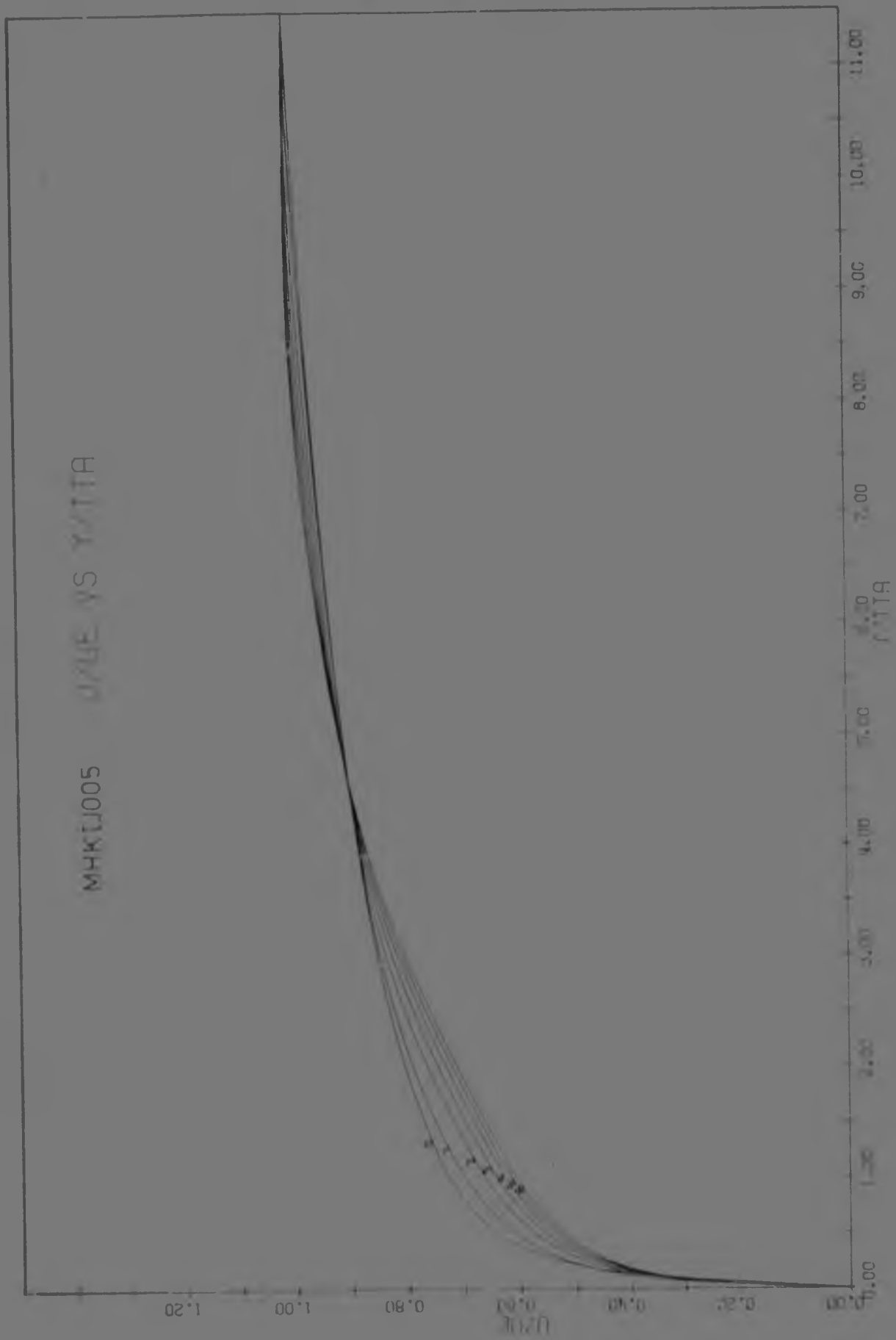


FIG 7.2]D : INJECTED TURBULENT PROFILES FOR  
 $U_m = 17 \text{ m/s}$ ,  $F = 0,005$

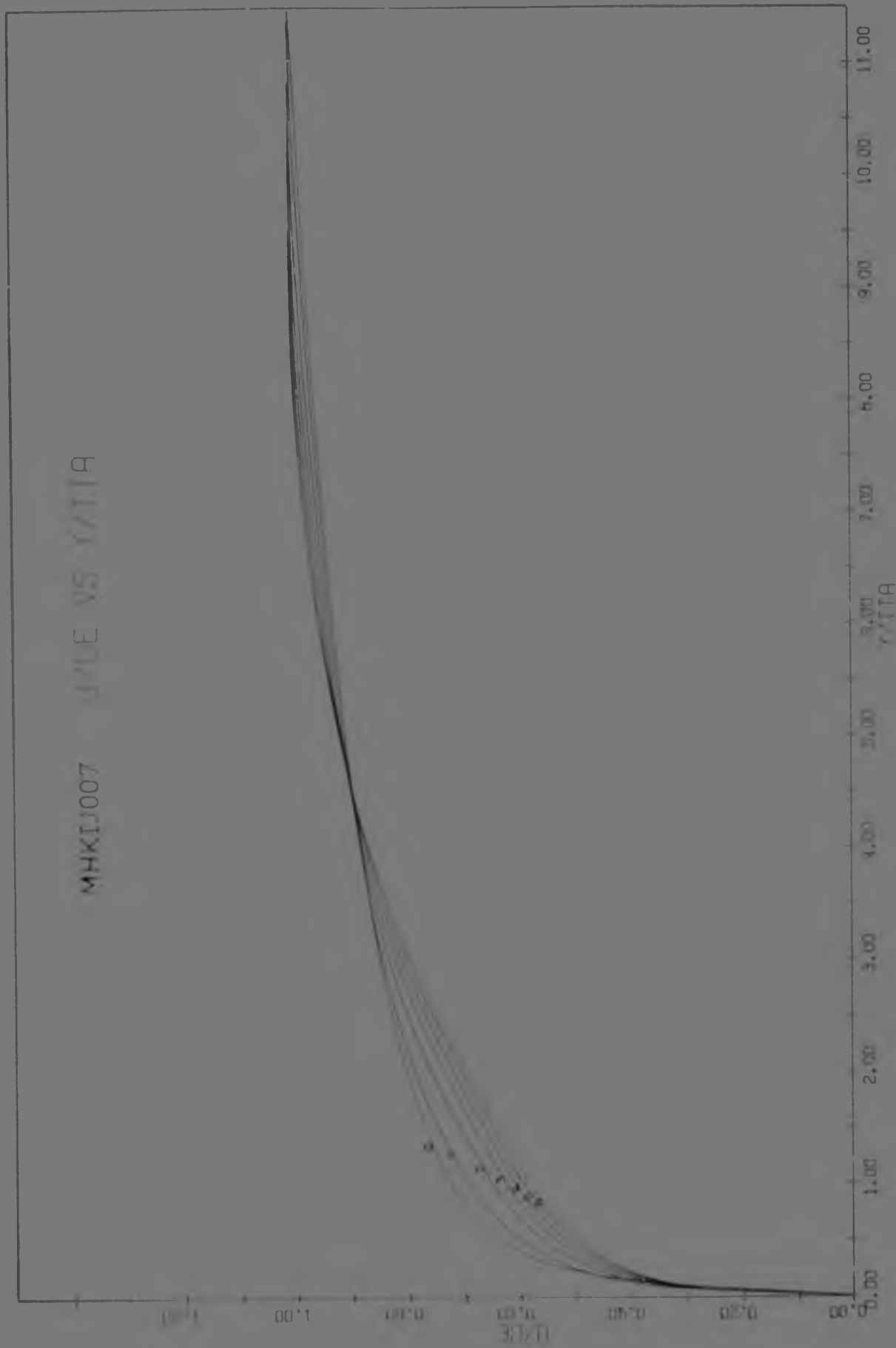


FIG 7.21E : INJECTED TURBULENT PROFILES FOR  
 $U_{in} = 17 \text{ m/s}$ ,  $F = 0,007$

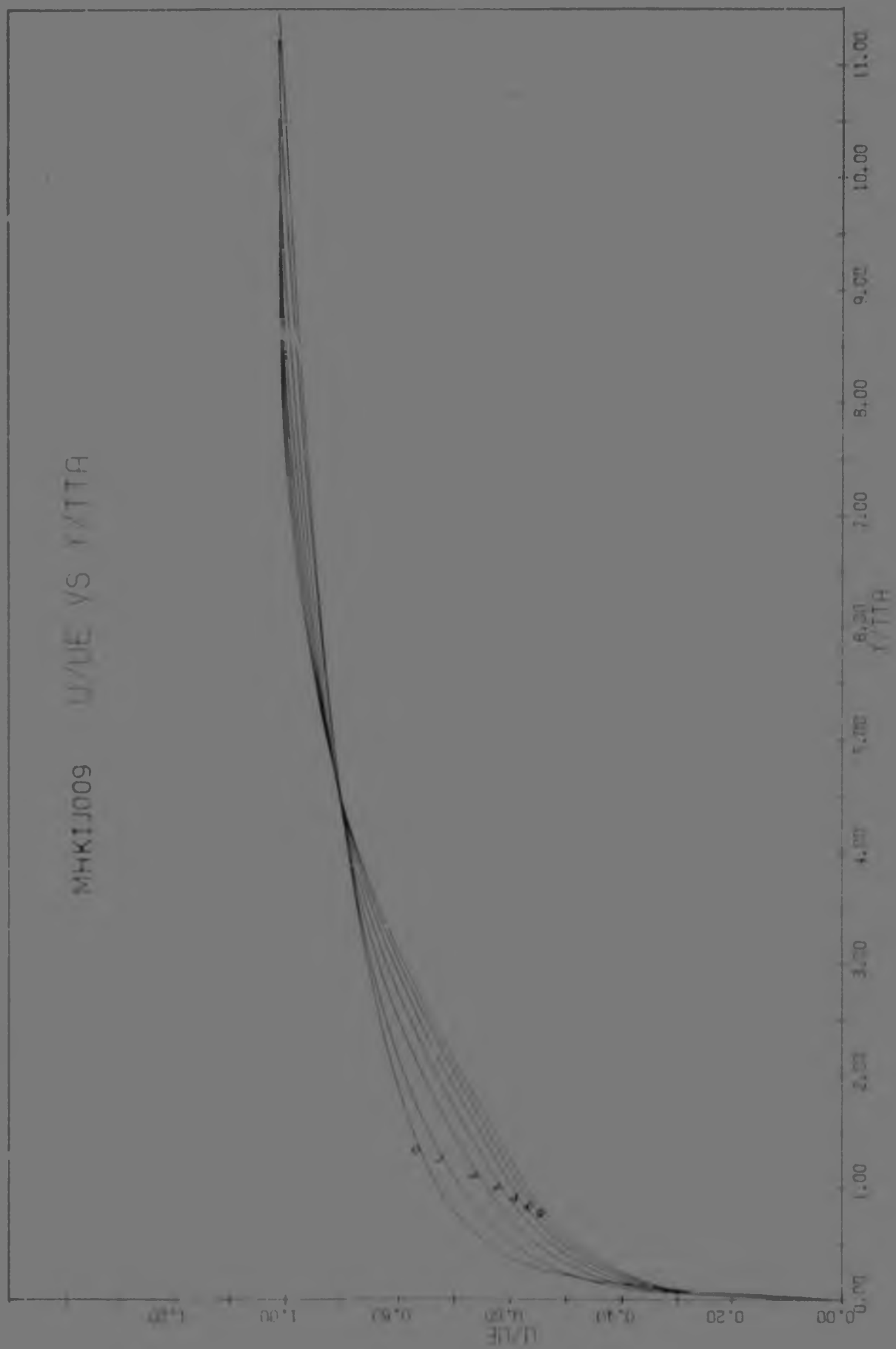


FIG 7.21F : INJECTED TURBULENT PROFILES FOR  
 $V_{in} = 17 \text{ m/s}$ ,  $F = 0.005$



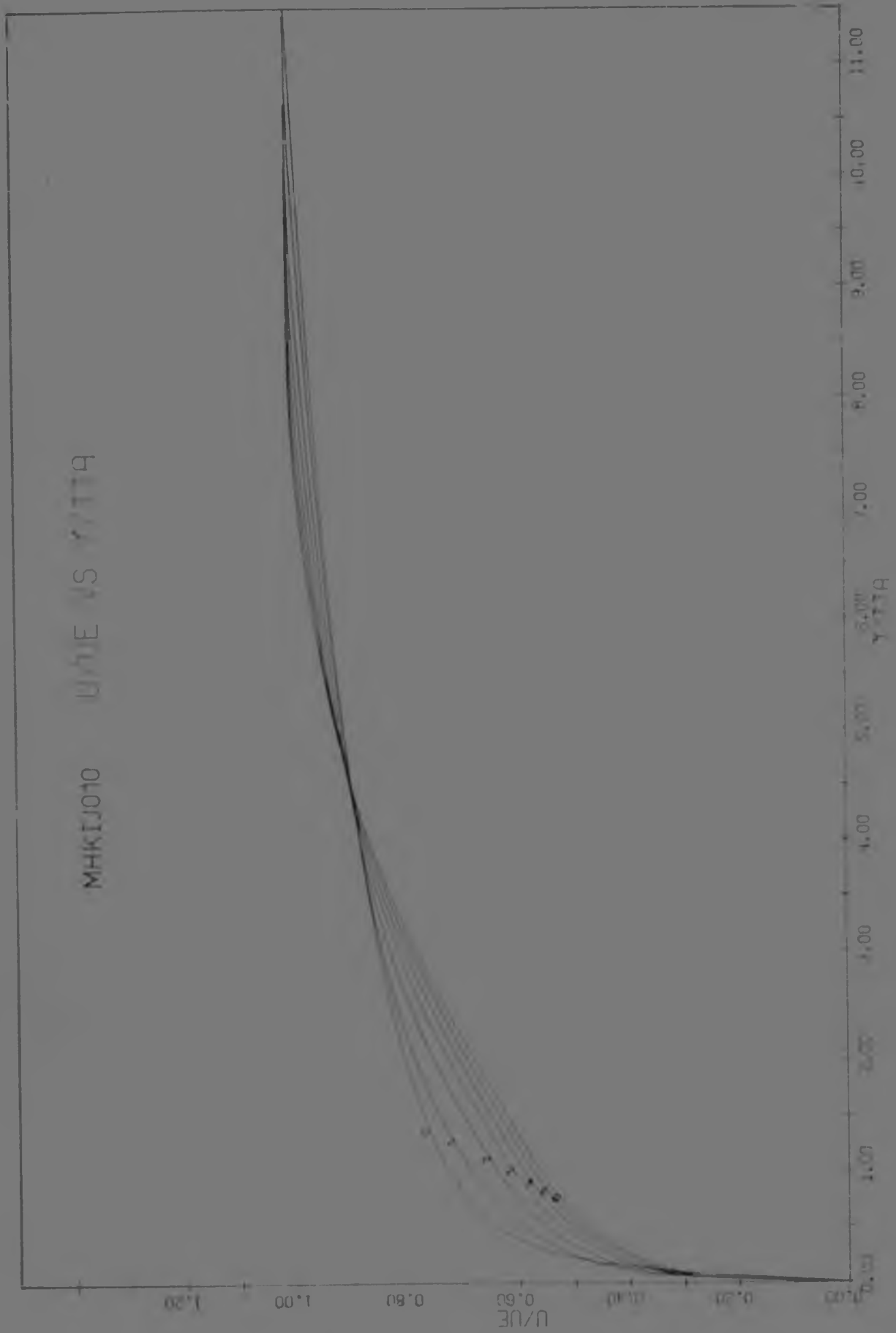


FIG 7.21G : INJECTED TURBULENT PROFILES FOR  
 $U_{in} = 17 \text{ m/s}$ ,  $F = 0,010$

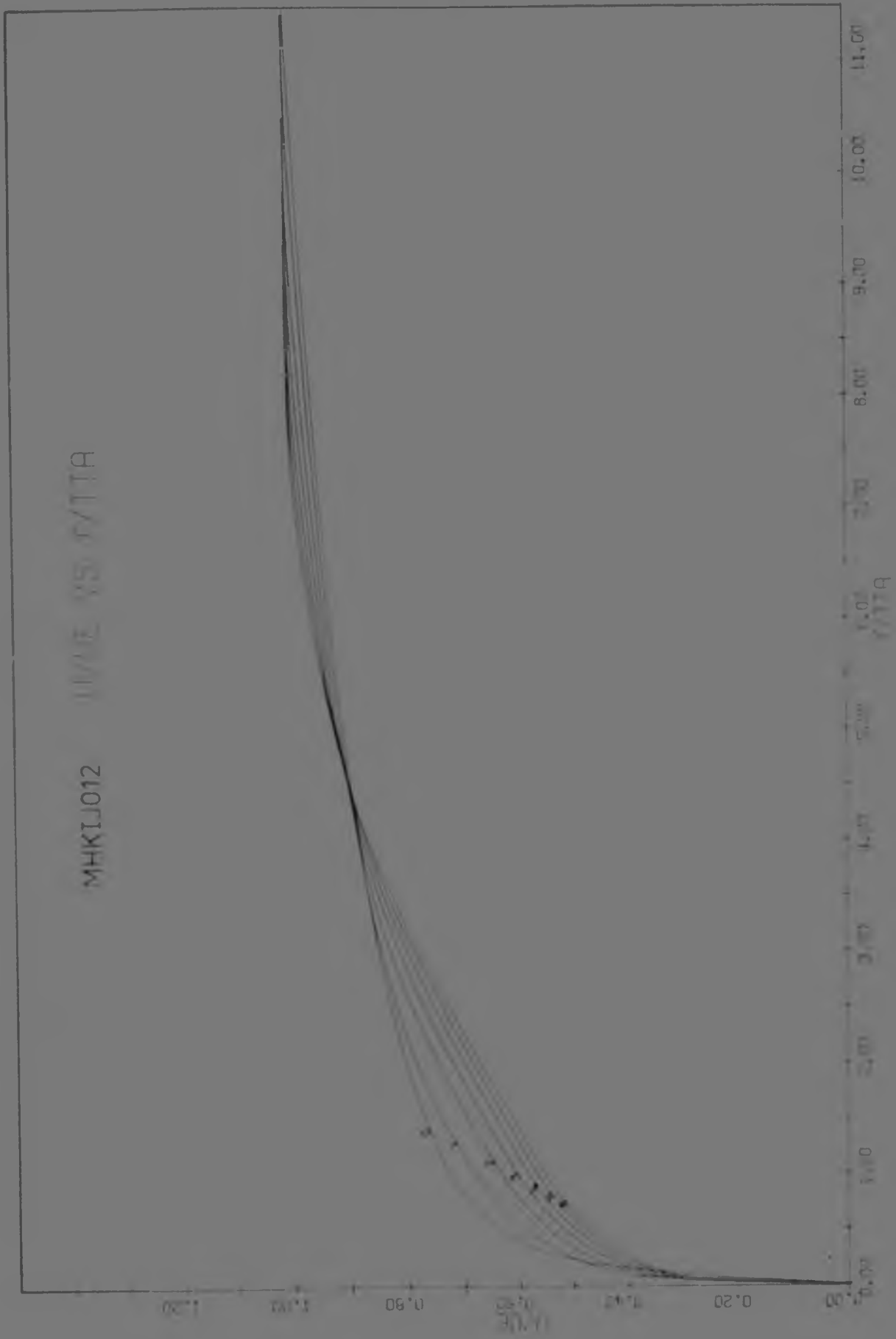


FIG 7.21H : INJECTED TURBULENT PROFILES FOR  
 $U_{in} = 17 \text{ m/s}$ ,  $t = 0,012$

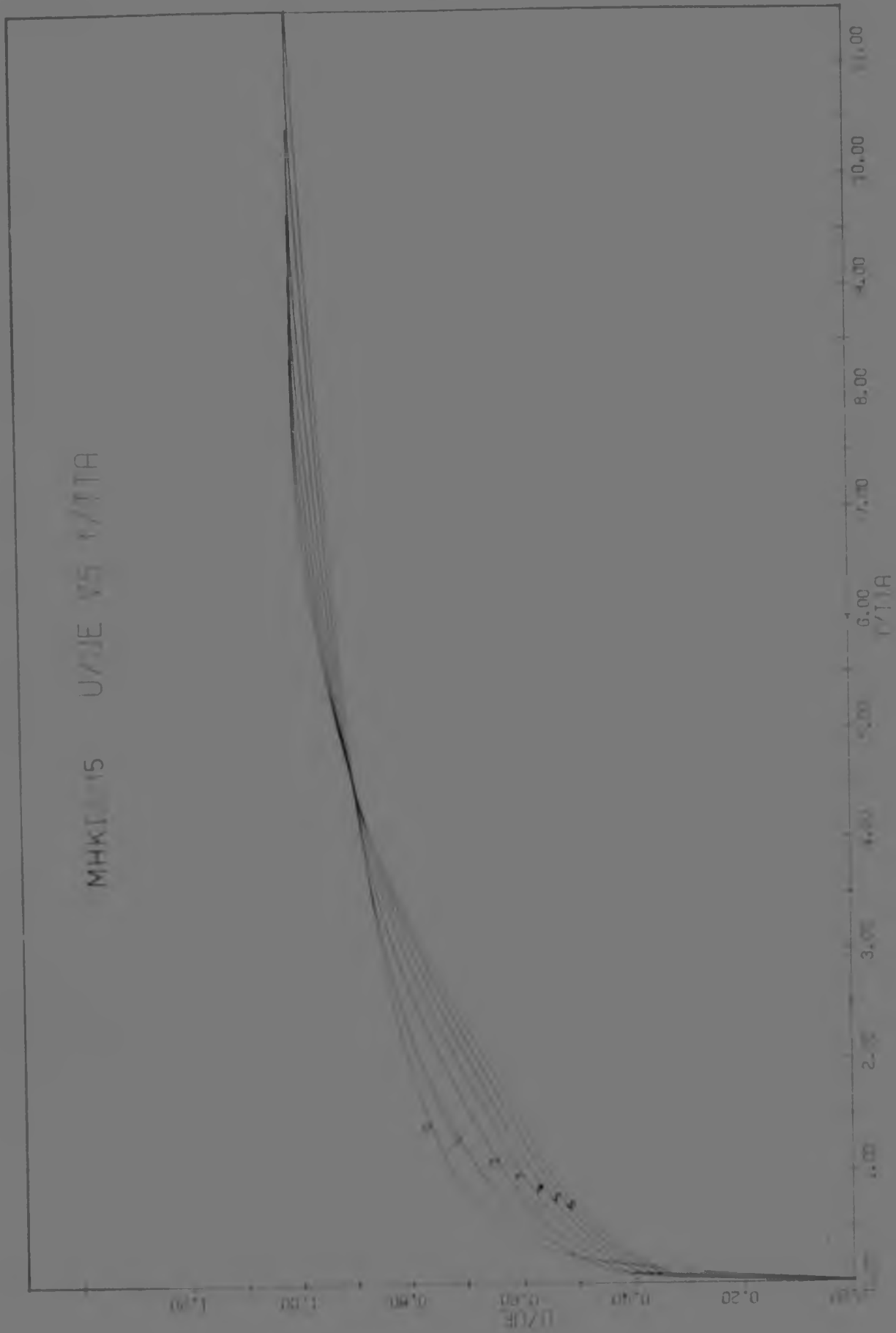


FIG 7.2]1 : INJECTED TURBULENT PROFILES FOR  
 $U_{in} = 17 \text{ m/s}$ ,  $F = 0,015$

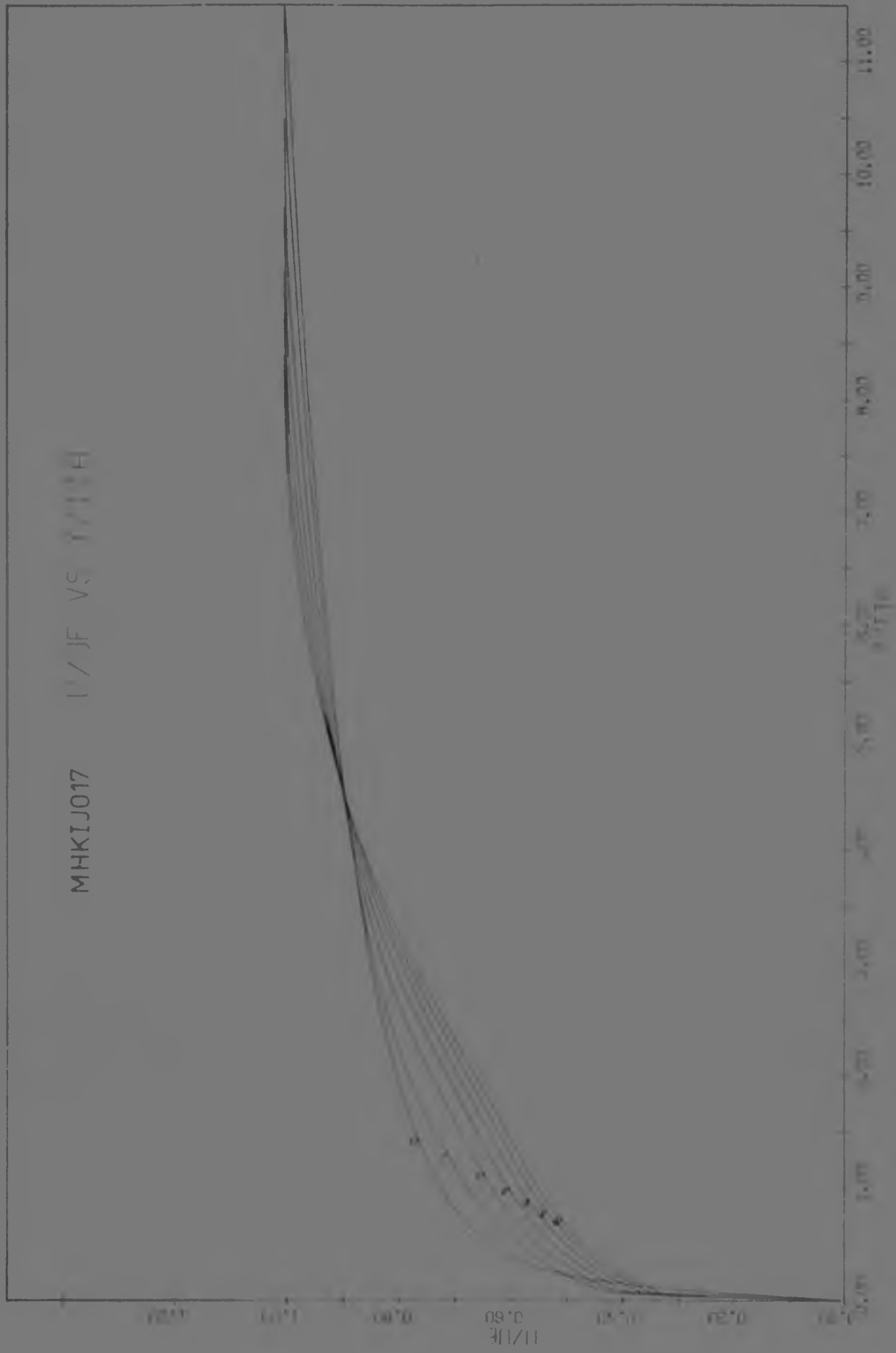


FIG 7.21J : INJECTED TURBULENT PROFILES FOR  
 $U_{\infty} = 17 \text{ m/s}$ ,  $F = 0,017$

By using the approximate "  $\phi$  " as described, the programme was remarkably stable, solving for injection ratios up to  $F = 0,017$ . These high injection ratios have not been reported in the literature, and are approaching, if not exceeding, the limit of the boundary layer analysis.

The configuration and condition of the fluid for the main computer runs are reported in Table 7.2. Fig 7.22 shows the eddy viscosity distribution across the boundary layer.

TABLE 7.2 : Key to  $x$ -Stations used in Main Runs

Station	$x$ (m)	$\xi \times 10^4$	$Re_x \times 10^5$	Plot symbol	
3	0,790	2,416	7,274	0	$U = 17 \text{ m/s}$
10	0,825	2,523	7,596	1	$du/dx = 0 \text{ 1/s}$
20	0,875	2,676	8,056	2	$\rho = 0,9971 \text{ kg/m}^3$
30	0,925	2,829	8,516	3	$\mu = 1,825 \text{ kg/ms}$
40	0,975	2,982	8,977	4	XTR = 0,8 m
50	1,025	3,135	9,437	5	IY = 100
60	1,075	3,288	9,898	6	K = 1,025
					H = 0,065
					DX = 0,005

#### 7.4.2 Some additional runs

A run with injection and pressure gradient was performed - the results presented in Fig 7.23. MIKIJ003 was the comparable run with  $du_e/dx = 0 \text{ 1/s}$ , and considering the shape factors at the same  $x$  stations show the difference very clearly, e.g. Station 60 :  $H_o = 1,45$ ,  $H_{du/dx=10} = 1,36$ .  $\delta$  grew at a much slower rate (32,4 mm and 25,2 mm respectively), which is as expected.

MIKLONG3 and MIKLONG5 were runs in which an uninterrupted porous plate was simulated. The porous plate started at  $x = 0,312 \text{ m}$ . This was improved in MIKTEST3, which is compared with the Mickley and Davis results for their run C-3-50. See Fig 7.26.

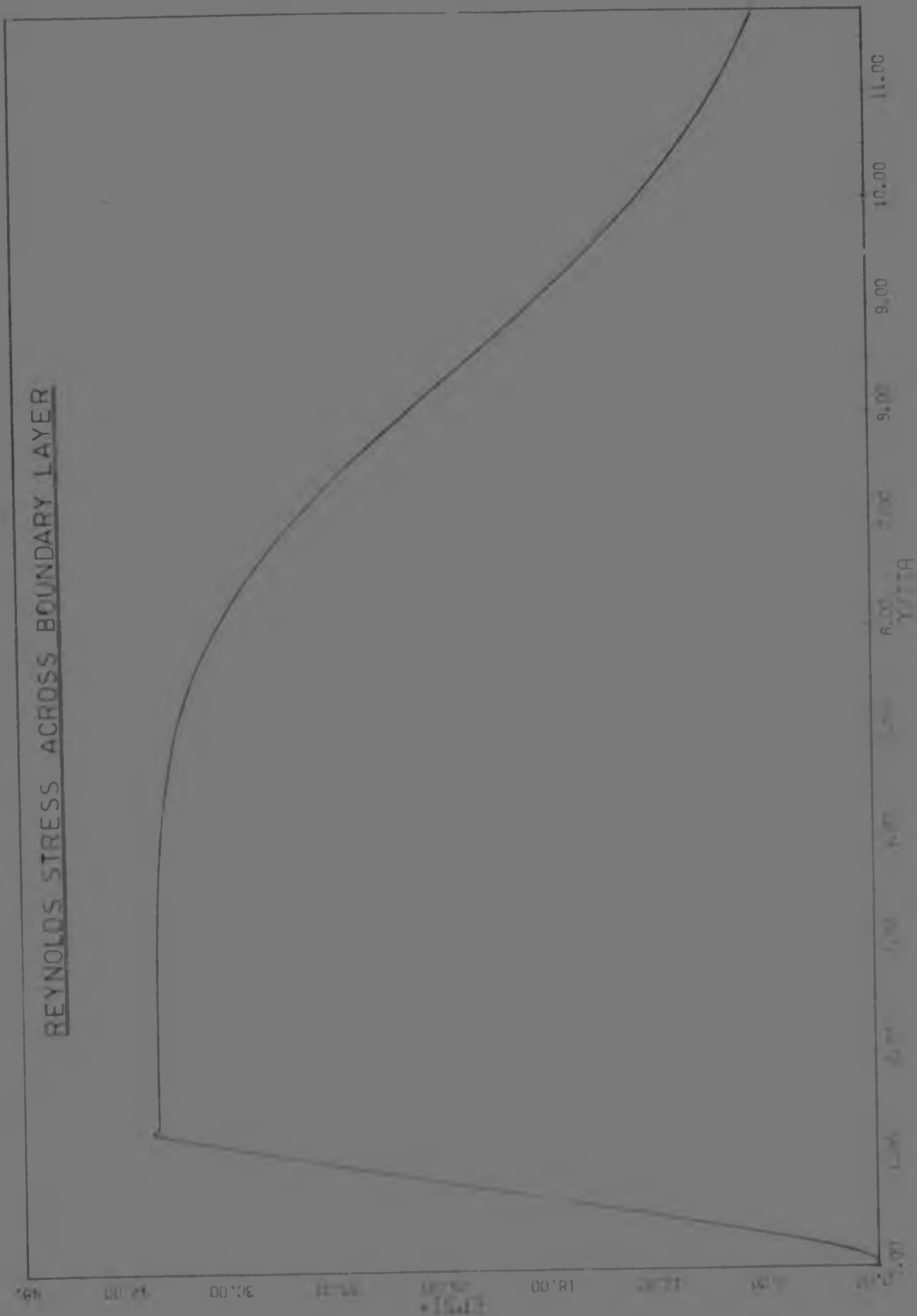


FIG 7.22 : EDDY-VISCOSITY ACROSS THE BOUNDARY LAYER  
AT  $x = 0,790$  m,  $U_m = 17$  m/s

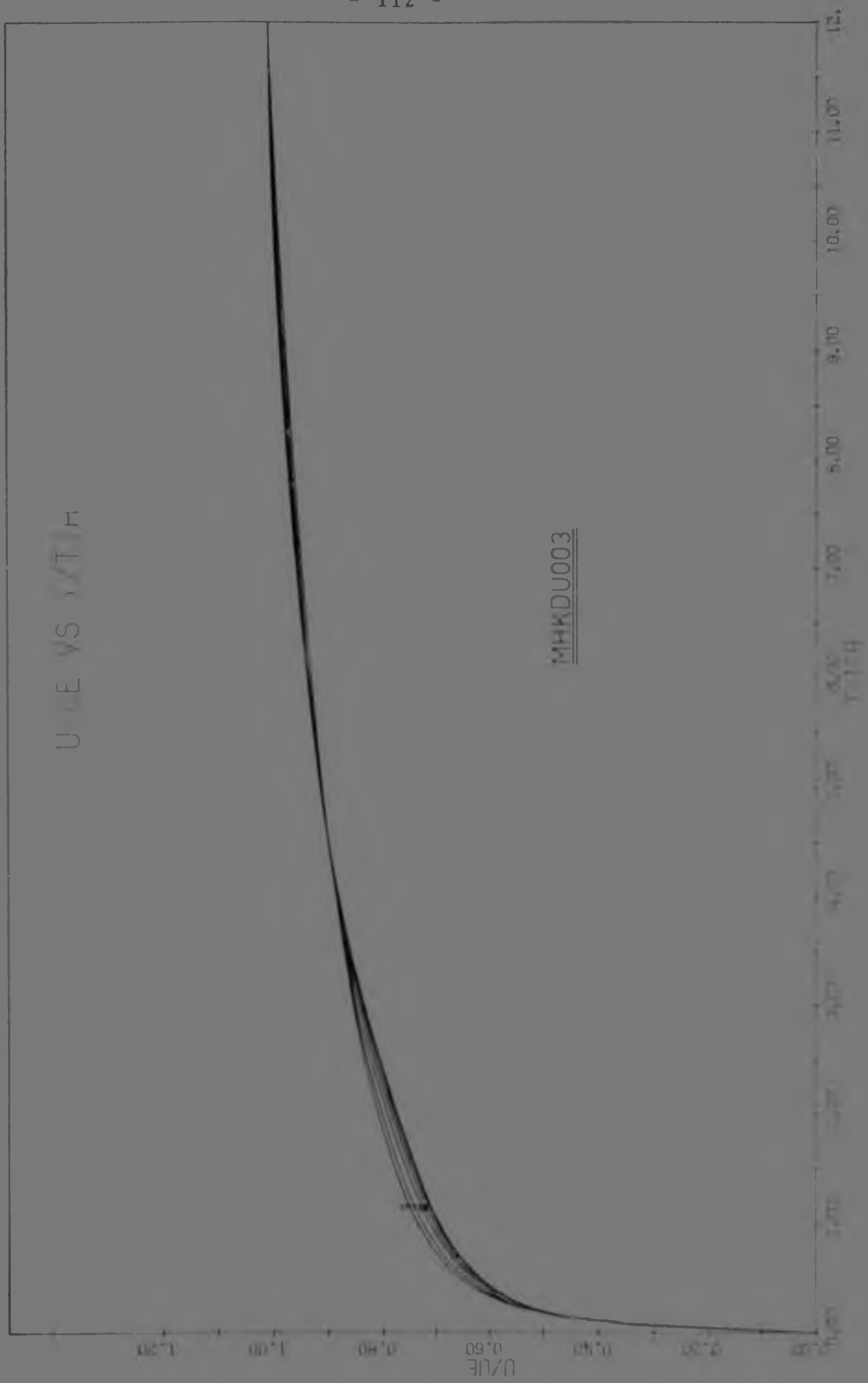


FIG 7.23 : VELOCITY PROFILE WITH INJECTION AND PRESSURE GRADIENT

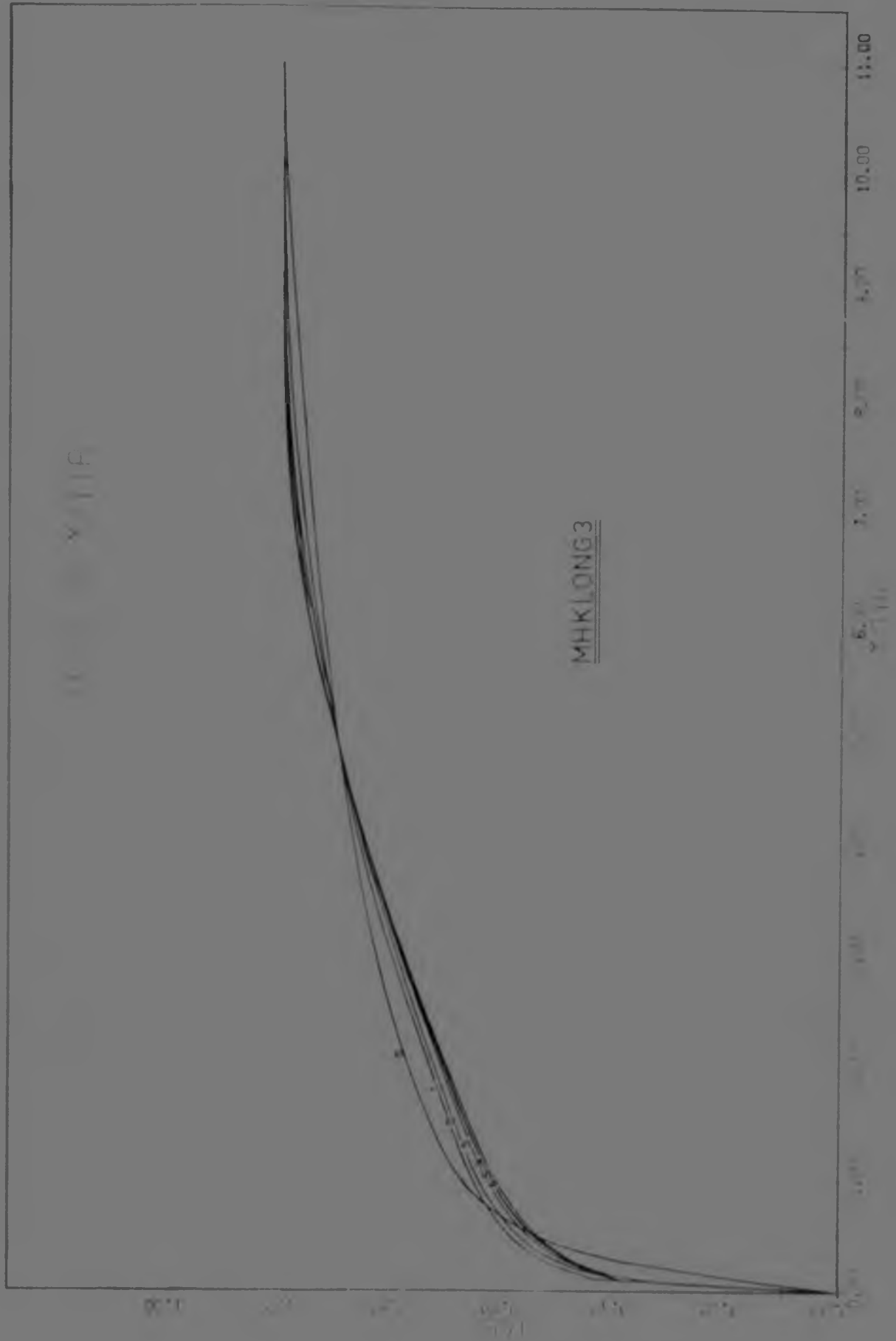


FIG 7.24 : VELOCITY PROFILE WITH INJECTION  
STARTING AT  $x = 0,312$  m



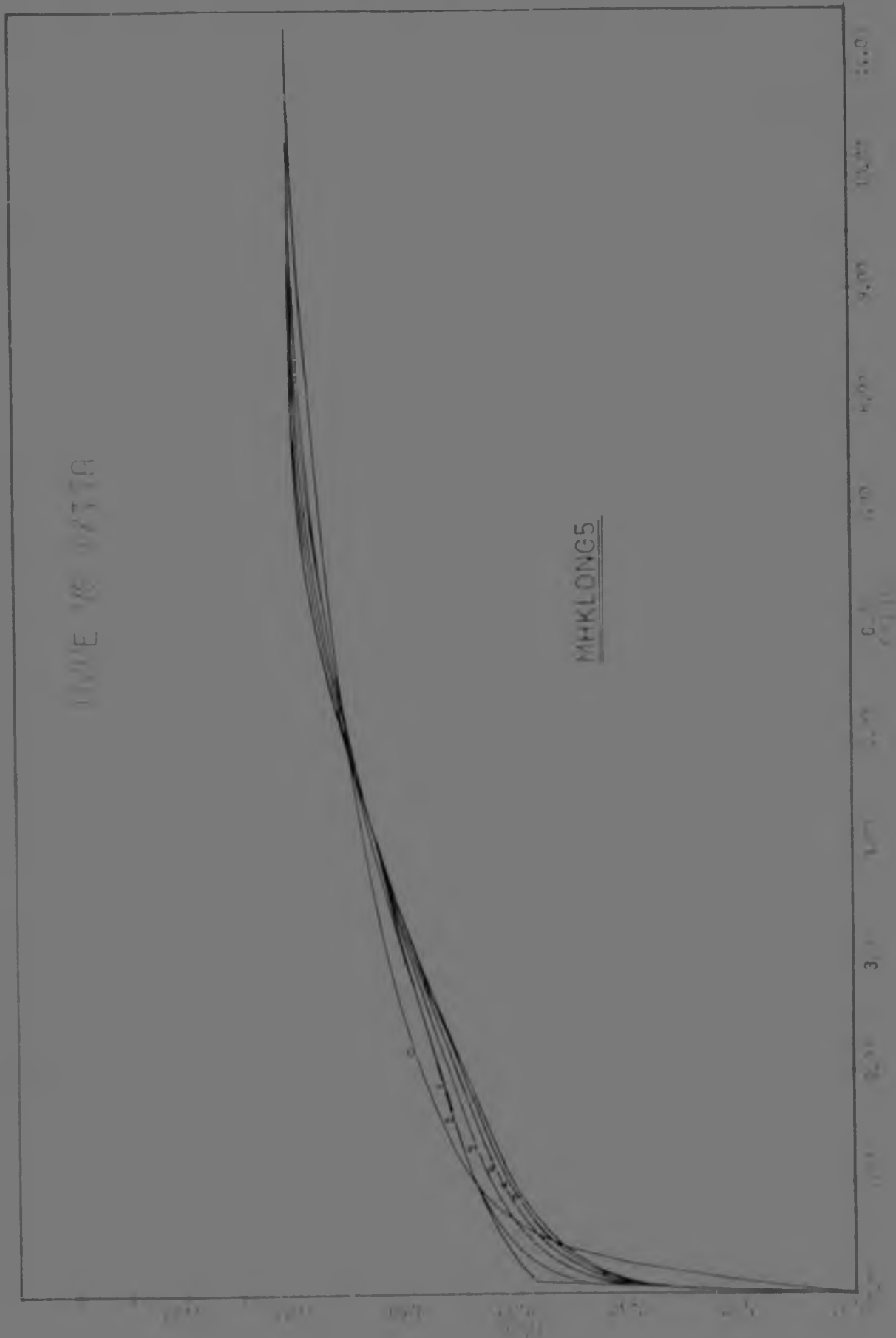


FIG 7.25 : VELOCITY PROFILE WITH INJECTION  
STARTING AT  $x = 0,312$  m

TABLE 7.3 : Comparison Table

MHKTEST 3 STATION 60				MICKLEY AND DAVIS C-3-50 STATION H			
$Re_{\theta}$	$\delta$ (mm)	$\theta$ (mm)	$H$	$Re_{\theta}$	$\delta$	$\theta$ (mm)	$H$
2687	25,87	2,94	1,50	2830	22,86	2,85	1,46

This is for the final station of run MHKTEST3, and  $Re_{\theta}$  was still somewhat below that of the Mickley and Davis experiment.  $\delta$  was again high, and hence also  $\theta$ . The profiles compare well, and it is expected that if  $Re_{\theta}$  were more accurately matched, the correspondence would be even better.

The programme was thus shown to produce results which compared well with experiment. Qualifying the programme more extensively would require many more runs, changing the free stream conditions and injection ratio, and finding similar experimental profiles. Note that the computer was using local (Johannesburg) density for air, which was why  $Re_{\theta}$  was lower than that of Mickley and Davis given in Table 7.3, for similar values of  $\theta$  and  $u_{\infty}$ .

To allow smaller step lengths to be used, a thinner boundary layer was necessary. For this reason the main runs were performed from  $x=0,790$  m, i.e. fully turbulent flow, but  $\delta=23$  mm, whereas the experiment produced  $\delta=45$ mm. The latter would require a much coarser grid, and hence loss of accuracy. It was therefore not possible to compare the main runs with the experimental results. However, the trend of the 'transition' region from impermeable wall to injected flow obtained experimentally was clearly reflected by the computer results.

Skin friction results were thought to be in error by orders of magnitude. As presented, the programme calculated  $\phi''$  from a second order fit, which was very satisfactory for the unblown profile. If Fig 5.9 is a correct representation of  $u$ , then a linear fit with very small step length at the wall would produce better results.

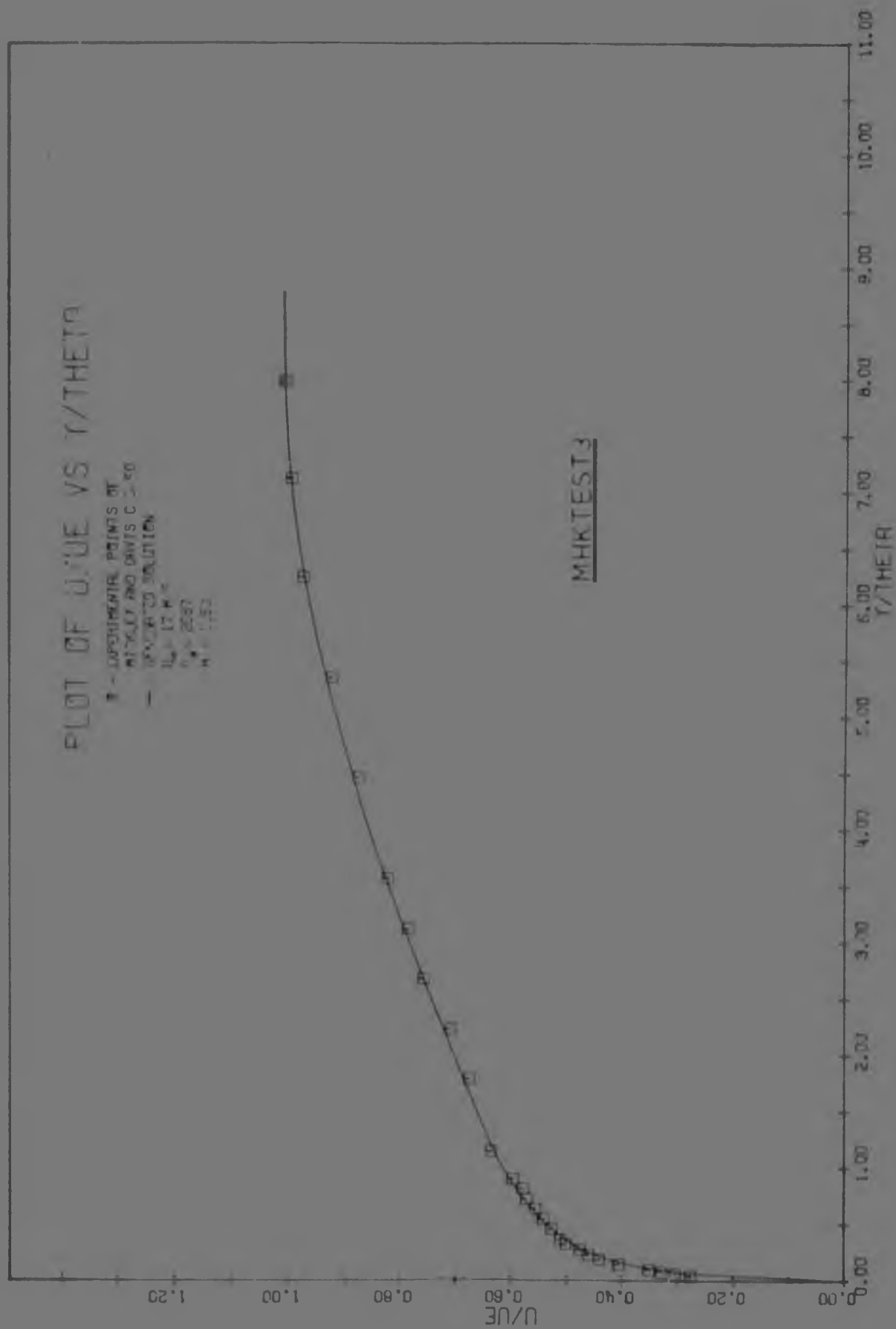


FIG 7.26 : COMPUTER GENERATED PROFILE COMPARED TO EXPERIMENTAL RESULTS OF MICKLEY AND DAVIS (1956)

## CHAPTER 8

### CONCLUSIONS

#### 8.1 Experimental Features

A two dimensional, open circuit wind tunnel was designed, built and commissioned. A free stream velocity of 21 m/s could be obtained in the working section. To simulate the conditions described on page 7, a turbulent boundary layer was produced on a flat plate, which was interspersed with a porous plate. The boundary layer thus encountered injection, which was introduced through the porous plate.

Using the hot wire anemometer, a transverse velocity profile was obtained, and showed that the flow was two dimensional. The turbulence intensity was better than 0.5%, which is less than the expected level for this type of tunnel (ref. Schlichting, 1968). Measured energy spectra at a number of free stream velocities compared favourably with accepted theoretical predictions and experimental data.

Velocity profiles on the smooth flat plate were ascertained and shown to follow the profiles predicted by the Thompson two parameter family of curves.

To further qualify the performance of the tunnel, profiles with injection were plotted on the Bilog axes suggested by Black and Sarnecki, together with experimental data of Mickley and Davis (1956). This allowed direct comparison with other authors work, and also showed the development of the profile with increasing injection. Furthermore, the need for accurate skin friction information was highlighted.

## 8.2 Velocity Profiles and General Experimental Results

Velocity profiles were measured for a range of free stream velocities and injection ratio. By using different grades of porous plate, the roughness of the wall was also varied. A set of data for the turbulent boundary layer with uniform injection was established. By partially blocking the lower surface of the porous plate, velocity profiles with non-uniform injection were ascertained, and compared with the initial set of results.

The experimental results gathered in the region in which the flow encountered the abrupt change in injection and roughness showed that the average velocity in the shear layer decreased with increasing  $x$ . The shape parameter  $H$  can be used to quantify this change in momentum. Plotting  $H$  vs  $x$  showed that  $H$  increased rapidly with  $x$ , with a decrease in slope for larger  $x$ . It also showed that  $H$  increased with injection ratio. Consideration of the results of Mickley and Davis confirms this result, and indicates that  $H$  may asymptote to a constant value for a fixed injection ratio.

When the results for the partially blocked porous plate were compared with the initial set of data, it was found that  $H$  was higher for a particular  $x$  station at the same average injection ratio. As is explained in the Addendum this result could indicate a lower wall temperature in the case of a heated stream and cold air injection (assuming constant heat transfer). This would be a very desirable result for the applications discussed on page 1, however, this can only be confirmed with extensive additional experimentation.

To the author's knowledge, these tests constituted original work for the following reasons :-

- i) The turbulent boundary layer encountered a step change in surface roughness and injection ratio, an intended feature of the test rig. Previous investigators had considered the fully developed turbulent boundary layer with uniform injection.

- ii) Partial blockage of the porous plate, thus altering the injection profile, constituted a further controlled parameter. This feature simulated the attachment of the porous matrix to a rigid base, which would thus enable these materials to be used for effusion cooling in highly stressed materials.

A technique for predicting the turbulent boundary layer velocity profile with injection was developed. A two parameter family of curves was presented, which was independent of injection ratio. The results of Mickley and Davis as well as McQuaid were added to these plots, and indicate the existence of a correlation equation over a significant range of injection ratios. It was thought that more experimental data was required before a correlation equation could confidently be deduced.

The lowest point is  $H = 1.3$ , which is the value for the un-injected turbulent boundary layer. As the layer encountered injection, the value of the shape parameter increased, while  $u/u_0$  decreased for a fixed value of  $y/\delta$ . This trend continued until the value at  $X1$  was reached which was very close to the results of Mickley and Davis, whose tests were done in the 'fully developed turbulent boundary layer'.

### 8.3 Features Pertinent to the Numerically Obtained Results

The momentum equations were solved numerically, and the velocity profiles thus generated were plotted, reflecting the trends determined experimentally, i.e. the decrease of velocity in the shear layer with increased injection ratio and distance along the plate. To qualify the programme, various test runs were performed. Firstly, the laminar profile generated was compared with analytical results given in Schlichting. Then a turbulent profile was plotted on semi log axes and compared with established experimental and empirical data. Various test cases were investigated, in which  $Re_x$  was varied with and without pressure gradient, and also one in which the computer solution was allowed to march into the transition region.

Finally, runs with injection were performed. One of these was compared with the experimental results of Mickley and Davis. In all these comparisons, correlation was excellent, and the programme was considered to yield accurate velocity profiles.

The programme was run for a range of injection ratios of  $F = 0,017$ . This was the highest ratio investigated experimentally, and was thought to be at the limit of the assumption of the analysis. It should be noted that  $F = 0,017$  is significantly higher than any injection ratio encountered in the literature survey, both in experimental or numerical work. The programme remained exceptionally stable, even at high values of  $F$ .

## CHAPTER 9

### SUGGESTIONS FOR FURTHER WORK

1. Modifications to the wind tunnel, which would result in a thin boundary layer arriving in the entry section, allowing greater control.
2. Measurement of the local skin friction by direct means, so that this matter may be cleared up satisfactorily. Initially, the effect of surface roughness with transpiration would have to be determined. If a floating element balance was used, the element must be made of porous material with the same roughness as the plate; and with injection at its surface.
3. To establish the heat transfer characteristics of step change, distributed injection, with and without blockage, tests must be done with a heated mainstream. This latter provision is intended to obviate the inverted temperature profile which results from heating the transpiration flow. The inverted temperature profile acts as a turbulence inhibitor, thus affecting the heat transfer mechanism.
4. The computer programme requires an improved finite difference technique for obtaining the second derivative, especially at the wall. From this, more realistic values of  $c_f$  may be expected. A more sophisticated turbulence transport equation will be necessary to account for surface roughness.
5. Alterations to the programme could be designed to allow for enclosed flow, and for more complicated pressure gradient situations.



6. By including the Energy Equation, the heat transfer to the plate could be calculated.
7. A deeper understanding of the flow in the porous matrix should be obtained, mainly from experiment. The spread of the jet, and the exit velocity of the air at the exposed surface should be studied.
8. Turbine blades in cascade with porous inserts should be manufactured and tests on heat transfer characteristics done. These could be compared with tests for blades with discrete hole injection.

APPENDIX A

A1 Transformation of the Momentum Equation

For convenience, the continuity and momentum equations of section 4.2 are restated here :-

Continuity :  $\frac{\delta u}{\delta x} + \frac{\delta v}{\delta y} = 0$  ..... [A1.1]

Momentum :  $u \frac{\delta u}{\delta x} + v \frac{\delta u}{\delta y} = \mu \frac{\delta^2 u}{\delta x^2} + \frac{\delta}{\delta y} (v \frac{\delta u}{\delta y} - u'v')$  [A1.2]

Numerous transformations exist for equation A1.2, a popular one is the transformation to the Von Mises plane. The equation takes on the form of an unsteady heat-conduction equation. The disadvantage in solving an equation which is pseudo-elliptic (ie. the entire flow-field has to be solved simultaneously), is the excessive computer time used. A transformation similar to the Falkner-Skan solution could be used in a finite difference scheme that marches in the direction of predominant flow. In this way, conditions at each successive step could be set as required, e.g. surface roughness, injection (steady or varying with  $x$ ), and also step length.

Following the treatment of Cebeci and Smith (1970a), the Levy-Lees transformation was used. This resulted in combining two independent variables,  $x$  and  $y$ , shortened the abscissa and lengthened the ordinate.

The Levy-Lees transformation for flat-plate, incompressible flow is :-

$d\eta = \rho \mu u_e dx; d\eta = \left(\frac{\rho u_e}{2\xi}\right)^{1/2} d\eta$  ..... [A1.3]

Define a dimensionless stream function,  $f$

$\psi(x,y) = (2\xi)^{1/2} f(\xi,\eta)$  ..... [A1.4]

$\psi$  is the stream function which satisfies the continuity equation :-

$$\frac{\delta\psi}{\delta y} = \rho u ; \quad \frac{\delta\psi}{\delta x} = -\rho v \quad \dots\dots\dots [A1.5]$$

Substituting A1.5 into A1.2, for  $\mu$  and  $\rho$  constant :

$$\frac{1}{\rho^2} \frac{\delta\psi}{\delta y} \cdot \frac{\delta^2\psi}{\delta x \delta y} - \frac{1}{\rho^2} \frac{\delta\psi}{\delta x} \cdot \frac{\delta^2\psi}{\delta y^2} = u_e \frac{du_e}{dx} + \frac{\nu}{\rho} \frac{\delta}{\delta y} [(1 + \epsilon^+) \frac{\delta^2\psi}{\delta y^2}] \quad [A1.6]$$

With  $\epsilon^+ = \epsilon/\nu$

Hence :-

$$\frac{\delta\psi}{\delta y} \cdot \frac{\delta^2\psi}{\delta x \delta y} - \frac{\delta\psi}{\delta x} \cdot \frac{\delta^2\psi}{\delta y^2} = \rho^2 u_e \frac{du_e}{dx} + \nu \frac{\delta}{\delta y} [(1 + \epsilon^+) \frac{\delta^2\psi}{\delta y^2}] \quad \dots [A1.7]$$

From A1.5 (see Cebeci and Smith, 1970a) :-

$$\left. \begin{aligned} \frac{\delta}{\delta x} &= \rho \mu u_e \left( \frac{\delta}{\delta \xi} + \frac{\delta \eta}{\delta \xi} \cdot \frac{\delta}{\delta \eta} \right) \\ \frac{\delta}{\delta y} &= \frac{\rho \mu u_e}{(2\xi)^{1/2}} \frac{\delta}{\delta \eta} \end{aligned} \right\} \dots\dots\dots [A1.8]$$

From A1.4 :-

$$\left. \begin{aligned} \frac{\delta\psi}{\delta \xi} &= (2\xi)^{1/2} \left( \frac{f}{2\xi} + \frac{\delta f}{\delta \xi} \right) \\ \frac{\delta\psi}{\delta \eta} &= (2\xi)^{1/2} f' \end{aligned} \right\} \dots\dots\dots [A1.9]$$

Using [A1.9] and  $\beta = (2\xi) \cdot (du_e/d\xi)/u_e$ , eqn.[A1.7] becomes :

$$[(1 + \epsilon^+) f'']' + f f'' + \beta(1 - f'^2) = (2\xi) \left[ f' \frac{\delta f'}{\delta \xi} - f'' \frac{\delta f}{\delta \xi} \right] \quad \dots\dots\dots [A1.10]$$

where the prime denotes  $\partial/\partial\eta$ , and  $f' = \delta f/\delta\eta = u/u_e$ .

The left hand side of [A1.10] will be recognised as the solution obtained by Falkner and Skan (1931) for the laminar boundary layer with  $u_e(x) \propto x^{1/2}$  (ie. the potential flow was

proportional to a power of the length measured from the stagnation point).

Now introduce a translated stream function,  $\phi$ . (This was done to improve the numerical computation.)

$$\left. \begin{aligned} f &= \phi + \eta \\ f' &= \phi' + 1 \\ f'' &= \phi'' \end{aligned} \right\} \dots\dots\dots [A1.11]$$

Substitute [A1.11] into [A1.10]:-

$$\begin{aligned} [(1 + \epsilon^+) \phi'']' + (\phi + \eta) \phi' - \beta \phi' (\phi' + 2) &= \\ 2\xi [(\phi' + 1) \delta \phi' / \delta \xi - \phi'' \delta \phi / \delta \xi] &\dots\dots\dots [A1.12] \end{aligned}$$

This equation was also given as [4.6.1], and is an o.d.e. in  $\eta$  at  $\xi = \xi_n$ .

A2 The Finite Difference Approximations

Consider a function  $f$  at grid points  $i-2, i-1, i, \dots$  etc., and let the distance between these points be  $h_1, h_2, h_3$  and  $h_4$  :-

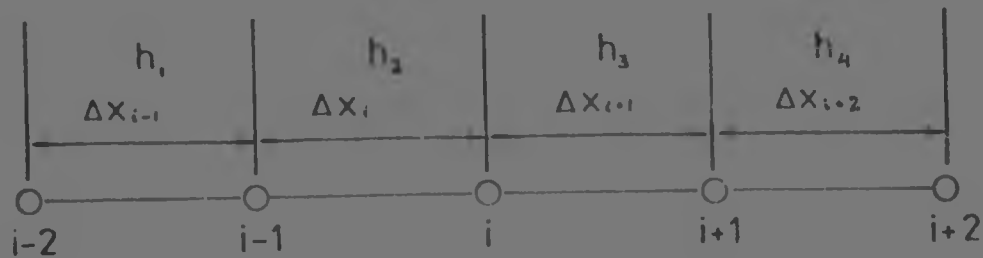


FIG A2.1 : SKETCH OF GRID POINTS, SHOWING NOMENCLATURE

Using Taylor's series to express :-

$$f(x_{i+1}) = f(x_i) + h_3 f'(x_i) + \frac{h_3^2 f''(x_i)}{2!} + O(h_3^3) \quad [A2.1]$$

$$f(x_{i-1}) = f(x_i) - h_2 f'(x_i) + \frac{h_2^2}{2!} f''(x_i) + O(h_2^3) \dots [A2.2]$$

By doing the following manipulation : [A2.1] x  $h_2^2$  -  
[A2.2] x  $h_3^2$  obtain :

$$f'(x_i) = \frac{-h_2 f(x_{i-1})}{h_2(h_2 + h_3)} - \frac{(h_2 - h_3) f(x_i)}{h_2 h_3} + \frac{h_3 f(x_{i+1})}{h_3(h_2 + h_3)} + O(h^2) \dots [A.2.3]$$

This equation was written for  $h_2 \neq h_3$ , which creates some difficulty in expressing the order of error. In [A2.3] it is approximately  $O(h_2 \cdot h_3)$ , and if  $h_3 + h_2 = h$ ,  $O(h^2)$ .

Eliminating  $f'$  from [A2.1] and [A2.2] :

$$f''(x_i) = \frac{2f(x_{i-1})}{h_2(h_2 + h_3)} - \frac{2f(x_i)}{h_2 h_3} + \frac{2f(x_{i+1})}{h_3(h_2 + h_3)} \dots [A2.4]$$

This equation has an error  $O(h_2)$  or  $O(h_3)$ , ie. much less accurate than [A2.3].

Equations derived in a similar manner are :-

$$f''(x_{i+1}) = \frac{2f(x_i)}{h_3(h_3 + h_4)} - \frac{2f(x_{i+1})}{h_3 h_4} + \frac{2f(x_{i+2})}{h_4(h_3 + h_4)} \dots [A2.5]$$

$$f''(x_{i-1}) = \frac{2f(x_{i-2})}{h_1(h_1 + h_2)} - \frac{2f(x_{i-1})}{h_1 h_2} + \frac{2f(x_i)}{h_2(h_1 + h_2)} \dots [A2.6]$$

Considering Fig A2.1 again, and using eqn.[5.3.3]

$$\left. \begin{aligned} h_1 &= \Delta x_{i-1} = h \cdot K^{i-3} \\ h_2 &= \Delta x_i = h \cdot K^{i-2} \\ h_3 &= \Delta x_{i+1} = h \cdot K^{i-1} \\ h_4 &= \Delta x_{i+2} = h \cdot K^i \end{aligned} \right\} \dots [A2.7]$$

Where  $h$  is the first step of the  $n$  grid.

Substituting [A2.7] into [A2.3] and [A2.4] :-

$$f'(x_i) = \frac{-f(x_{i-1})}{hK^{i-3}(1+K)} - \frac{(1-K)f(x_i)}{hK^{i-1}} + \frac{f(x_{i+1})}{hK^{i-1}(1+K)} \dots [A2.8]$$

$$f''(x_i) = \frac{2f(x_{i-1})}{h^2K^{2i-4}(1+K)} - \frac{2f(x_i)}{h^2K^{2i-3}} + \frac{2f(x_{i+1})}{h^2K^{2i-3}(1+K)} [A2.9]$$

Each term in the last two equations could be divided into three parts, namely  $f$ , a constant in  $K$  and  $h$ , and a coefficient in  $i$ . This fact was used to simplify the equations. The first term of equation [A2.8] thus becomes :

$$\frac{-f(x_{i-1})}{hK^{i-3}(1+K)} = -f(x_{i-1}) \cdot A = -f(x_{i-1}) \cdot \frac{A1}{K^i} \dots [A2.10]$$

So  $A1 = -K^3/h(1+K)$ , and all other coefficients dependent on  $h$  and  $K$  were calculated only once, ie. in subroutine MK002.

Because the  $x$ -stations were not restrained as the  $\eta$  points were, coefficients required for derivatives in the  $\xi$ -direction were calculated at each new station.

The coefficients that appeared in equation [5.4.2], are listed below :-

$$\left. \begin{aligned} A &= \frac{-K^3}{h(1+K)} \cdot \frac{1}{K^i} = \frac{A1}{K^i} \\ B &= \frac{-(1-K)K}{h} \cdot \frac{1}{K^i} = \frac{B1}{K^i} \\ C &= \frac{K}{h(1+K)} \cdot \frac{1}{K^i} = \frac{C1}{K^i} \\ D &= \frac{2K^4}{h^2(1+K)} \cdot \frac{1}{K^{2i}} = \frac{D1}{K^{2i}} \\ E &= \frac{-2K^3}{h^2} \cdot \frac{1}{K^{2i}} = \frac{E1}{K^{2i}} \\ F &= \frac{2K^3}{h^2(1+K)} \cdot \frac{1}{K^{2i}} = \frac{F1}{K^{2i}} \end{aligned} \right\} \dots [A2.11]$$

$$\begin{aligned}
 X1 &= \frac{2K^6}{h^2(1+K)} \cdot \frac{1}{K^2} i \\
 X2 &= \frac{-2K^5}{h^2} \cdot \frac{1}{K^2} i \\
 X3 &= \frac{2K^5}{h^2(1+K)} \cdot \frac{1}{K^2} i \\
 X4 &= \frac{2K^2}{h^2(1+K)} \cdot \frac{1}{K^2} i \\
 X5 &= \frac{-2K}{h^2} \cdot \frac{1}{K^2} i \\
 X6 &= \frac{2K}{h^2(1+K)} \cdot \frac{1}{K^2} i
 \end{aligned}
 \left. \begin{array}{l} \\ \\ \\ \\ \\ \\ \end{array} \right\} \dots\dots\dots [A2.11]$$

Coefficients starting with X were not split up into two parts. This had been done initially, in a different form, and is the reason for real G1, H1, I1, J1, K1 being defined in subroutine MK002.

The streamwise derivatives were calculated as follows :-

$$\frac{\delta \phi_i}{\delta \xi} \Big|_n = L \phi_{n-2, i} + M \phi_{n-1, i} + N \phi_{n, i} \quad \dots\dots [A2.12]$$

$$\frac{\delta \phi'_i}{\delta \xi} \Big|_n = L \phi'_{n-2, i} + M \phi'_{n-1, i} + N \phi'_{n, i} \quad \dots\dots [A2.13]$$

These equations were obtained in a manner very similar to the derivation of [A2.3]. They were backward difference equations, and the step length remained variable.

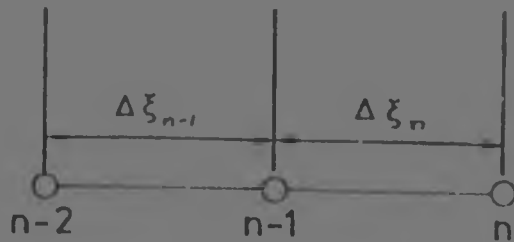


FIG A2.2 : SKETCH OF ξ GRID POINTS

$$\frac{\delta f(\xi_n)}{\delta \xi} = \frac{\Delta \xi_n f(\xi_{n-1})}{\Delta \xi_{n-1} (\Delta \xi_n + \Delta \xi_{n-1})} - \frac{(\Delta \xi_{n-1} + \Delta \xi_n) f(\xi_{n-1})}{\Delta \xi_{n-1} \cdot \Delta \xi_n} + \frac{(\Delta \xi_{n-1} + 2\Delta \xi_n) f(\xi_n)}{\Delta \xi_n (\Delta \xi_{n-1} + \Delta \xi_n)} \dots\dots\dots [A2.14]$$

Thus :-

$$L = \frac{\Delta \xi_n}{\Delta \xi_{n-1} (\Delta \xi_n + \Delta \xi_{n-1})}$$

$$M = \frac{-(\Delta \xi_{n-1} + \Delta \xi_n)}{\Delta \xi_{n-1} \cdot \Delta \xi_n} \dots\dots\dots [A2.15]$$

$$N = \frac{(\Delta \xi_{n-1} + 2\Delta \xi_n)}{\Delta \xi_n (\Delta \xi_{n-1} + \Delta \xi_n)}$$

To calculate  $\phi''$ , Newton's Divided Difference method was used. The following equation was derived (see Milne-Thomson, 1933), from a second order fit :-

$$\phi_1'' = \frac{\phi_1' - \phi_2'}{\eta_1 - \eta_2} + (\eta_1 - \eta_2) \frac{\frac{\phi_2' - \phi_3'}{\eta_2 - \eta_3}}{\eta_1 - \eta_3} \dots\dots [A2.16]$$

From Chapter 5,  $\phi_1$  and  $\phi_2$  were to be determined from boundary conditions. An equation was given for  $\phi_1$ , but  $\phi_2$  still remained to be calculated. From eqn. [4.6.3], one of the boundary conditions is -

$$\phi'(\xi, 0) = \phi_1' = -1 \dots\dots\dots [A2.17]$$

Using a very simple finite difference equation -

$$\phi_1' = \frac{\phi_2 - \phi_1}{\eta_2 - \eta_1} = -1 \dots\dots\dots [A2.18]$$

$$\therefore \phi_2 = \phi_1 - 1(\eta_2 - \eta_1)$$

[A2.18] did not give very satisfactory results, but by using the divided difference technique, the following relationship was derived :-

$$\phi_2 = \frac{\phi_1 \times K \times (K+2) + \phi_1 - K \times \eta_1 \times (K+1)}{(K+1)^2} \dots\dots [A2.19]$$



The unblown profile was checked, and [A2.19] was found to agree to 6 figures. After the experience gained with calculating  $\phi_1$  from a second order fit, it was felt that the plot of  $\phi$  vs  $\eta$  should be reconsidered when injected, and a better fit developed for this curve. Although not tested, it was thought that a 4th order fit would be an improvement.

Except at the boundaries, central difference equations were used throughout. Note that these equations were derived to allow calculation of a function at  $i$ , and not mid-way between  $i-1$  and  $i+1$ , which would have occurred had the function been calculated from a standard central difference equation for equal step length.

The equations derived were relatively simple, but produced a banded matrix with 5 columns. More complex equations, which become possible for equal step length, would have resulted in more columns in the banded matrix. This would require more core storage and more calculations, with related round-off errors, when solving  $[A]\{\phi\} = \{B\}$ . For generating the initial turbulent profile, it was found that the unequal step length improved the accuracy significantly. Because  $\delta u/\delta y$  was much lower for laminar profiles, equal steps were more suitable for these calculations. This could be obtained by setting  $K = 1$ .

### A3 Flow in Porous Media

One of the important aspects of using porous materials in transpiration cooling, is the need for strengthening the structure. As already suggested, this could be done by attaching the porous matrix to a rigid base with metering holes or slots.

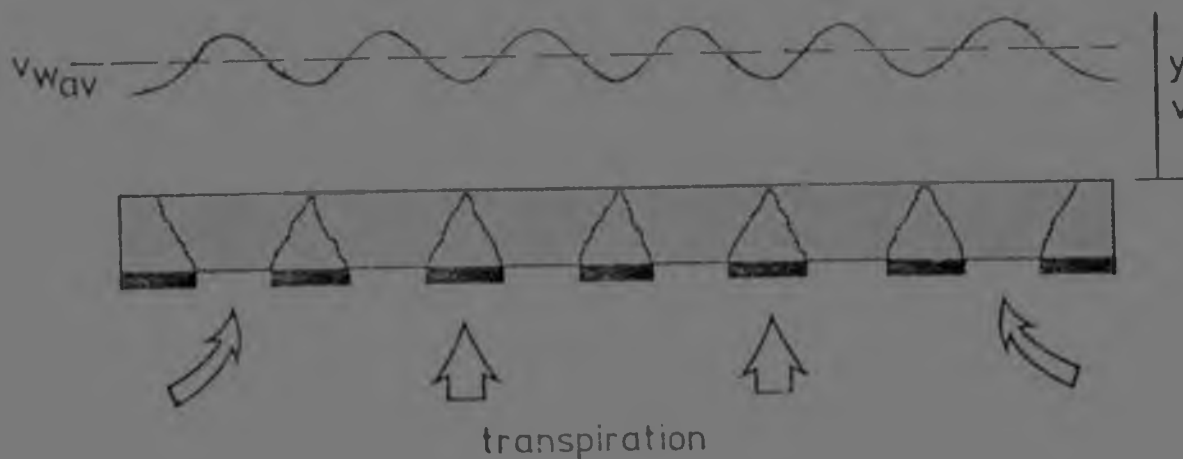


FIG A3.1 : THE POROUS MATRIX (  $\square$  ) ATTACHED TO A RIGID BASE (  $\blacksquare$  ) WITH METERING SLOTS

In the past, all results for pressure drop through the matrix have been obtained for the unrestricted case, ie, a one-dimensional flow situation. The average velocity at any  $y$  value in the porous material was constant. The pressure drop was found to be proportional to the superficial velocity,  $v_{av}$ , for the so-called laminar region, and to  $v_{av}^2$  for the 'turbulent' region. This was found experimentally, and is useful if  $v_{av}$  is constant.

The configuration shown in Fig A3.1 was expected to behave very differently. Because the jet is spreading in the matrix, the velocity is changing, and hence also  $\Delta p$ . To design a porous cover for transpiration cooling, it is necessary to know : (i) the spread of the jet with  $y$ , and from that, the distance between metering slots or holes to ensure a suitable insulating layer of cool air on the upper surface of the porous material, and (ii) the pressure drop for the two-dimensional flow situation

Two approaches to the problem are possible. The microscopic, analytic approach, or the macroscopic approach, discussed in Appendix A4.

Because a porous matrix is made up of randomly sintered particles with sizes spanning a given tolerance, it is not meaningful to consider the flow at a point. Following

Slattery (1972) the momentum equation was volume averaged, in much the same way that turbulent flows were averaged on a time base.

Two theorems, developed by Slattery (1972), were used, as were the Reynolds rules of averages. The theorems were for the volume average of gradient, and of the divergence, of a second order tensor or a spatial vector field, called  $\underline{B}$ .

$$\overline{\nabla \underline{B}} = \frac{1}{V} \int_{V_f} \nabla \underline{B} dV = \nabla \underline{\overline{B}} + \frac{1}{V} \int_{S_w} \underline{B} \cdot \underline{n} dS \quad \dots\dots [A3.1]$$

$$\overline{\text{div} \underline{B}} = \frac{1}{V} \int_{V_f} \text{div} \underline{B} dV = \text{div} \underline{\overline{B}} + \frac{1}{V} \int_{S_w} \underline{B} \cdot \underline{n} dS \quad \dots\dots [A3.2]$$

where  $V_f$  is the volume,  $S_w$  the surface area.

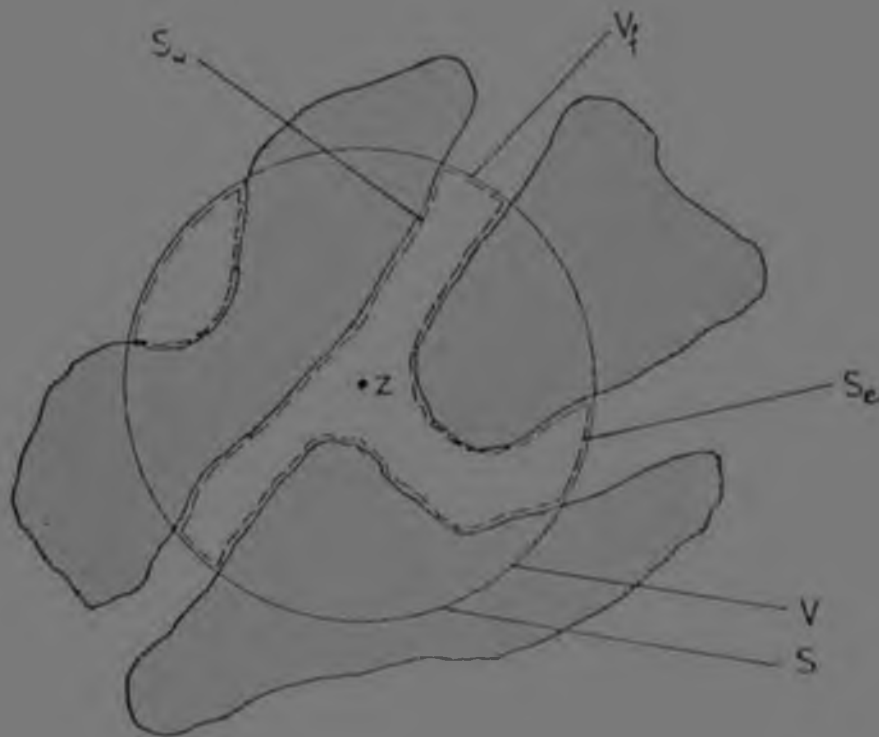


FIG A3.2 : THE POROUS MEDIUM, SHOWING VOLUME  $V_f$ , AND THE ENCLOSING SURFACE  $S$ . NOTE THAT  $S$  MUST ENCLOSE SOLID AND FLUID.  $S_w$  COINCIDES WITH THE PORE WALLS.  $\underline{n}$  IS THE NORMAL VECTOR.

The Navier-Stokes equation for the steady flow of an incompressible newtonian fluid may be derived from Cauchy's first law of motion (assume no external forces were present) :

$$\text{div}(\underline{v} \underline{v}^1) = -\nabla p + \mu \text{div}(\nabla \underline{v}) \quad \dots\dots\dots [A3.3]$$

Taking the local volume average of eqn.[A3.3] :-

$$\frac{1}{V} \int_{V_f} \rho \text{div}(\underline{v} \underline{v}) dV = -\frac{1}{V} \int_{V_f} \nabla p dV + \frac{1}{V} \int_{V_f} \mu \text{div}(\nabla \underline{v}) dV \quad \dots [A3.4]$$

Noting that  $(\underline{v} \underline{v} \cdot \underline{n}) = 0$  on the wall,  $S_w$ , and introducing a quantity  $v'$ , representing the deviation of the local point velocity from the local volume average velocity :-

$$\underline{v} = \underline{\bar{v}} + \underline{v}' \quad \dots\dots\dots [A3.5]$$

Note  $\underline{\bar{v}} = \underline{\bar{v}} \quad \therefore \underline{\bar{v}}' = 0$

Equation [A3.3] become :-

$$\rho \text{div}(\underline{v} \underline{v}) = -\nabla \bar{p} + \mu (\nabla^2 \underline{v}) - \rho \text{div}(\underline{\bar{v}}' \underline{v}') + \frac{1}{V} \int_{S_w} p n dS \quad [A3.6]$$

The interesting feature of [A3.6] was the additional term,  $-\rho \text{div}(\underline{\bar{v}}' \underline{v}')$  which is very similar to the Reynolds stress term produced in the time based averaging process used for turbulent flows.

Various techniques have been suggested to simplify [A3.6], and for one dimensional creep flow (ie. delete inertial terms) the Darcy law can be derived. Another simplification for two dimensional, inertial flow, suggested by Whitaker (1969), introduced a total resistance tensor  $R$ . The nine terms in this tensor would have to be evaluated experimentally, and were expected to be proportional to  $\mu$ , some length scale and the permeability of the porous medium.

---

1  $\underline{v} \underline{v}$  is known as the dyadic product, and was treated as a second order tensor.

At the present stage it was felt that the above line of thought was not worth pursuing further : techniques for determining  $R$  were not known to the author.

#### A4 Tests on the Porous Matrix

At the present state of the art, it was felt that a macroscopic approach could be expected to yield results more quickly. A preliminary investigation was done, starting with a dimensional analysis.

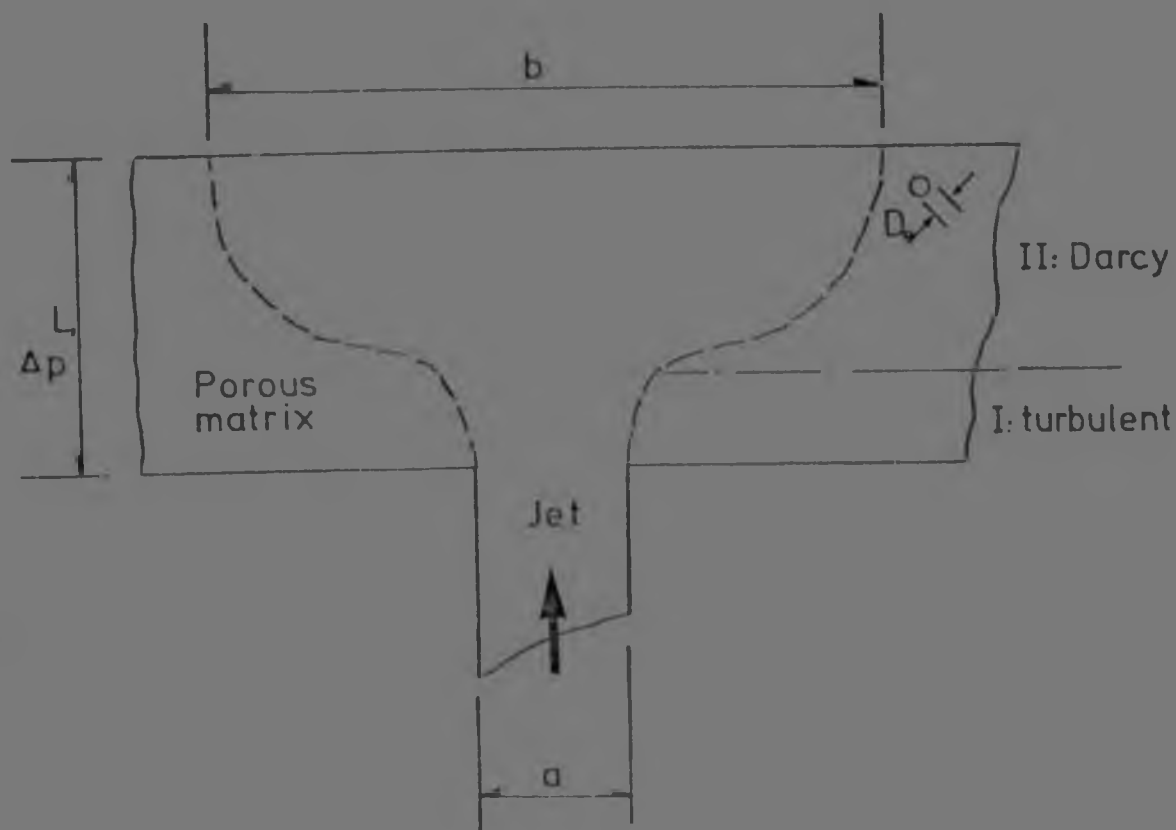


FIG A4.1 : SKETCH OF POROUS MATRIX WITH JET AT LOWER SURFACE

Because designers of transpiration cooled turbine nozzle guide vanes have a very limited pressure differential to work with,  $\Delta p$  was considered the most important quantity.

The variation of  $b$  with  $a$ ,  $L$  and  $V$  is required if a uniform cold layer is to be present on the outer surface.

$$\Delta p = f(L, a, D, b, V, \mu, \rho) \quad \dots\dots\dots [A4.1]$$

Typically, the following relationship was produced :-

$$\frac{\Delta p}{\mu^2/\rho D^2} = (Re_D)^u \left(\frac{a}{D}\right)^v \left(\frac{b}{D}\right)^w \left(\frac{L}{D}\right)^x$$

Because 3 length scales were present, these could be interchanged at will. The first dimensionless group was particularly interesting :-

$$\frac{\Delta p}{\mu^2/\rho D^2} = \frac{\Delta p}{\rho V^2} \cdot Re_D^2 \quad \dots\dots\dots [A4.2]$$

$\Delta p/\rho V^2$  is the well known pressure drop coefficient  $C_p$ .

Define

$$C_{pD}' = C_p \cdot Re_D^2 \quad \dots\dots\dots [A4.3]$$

Some tests were performed on a matrix made of plastic spheres,  $D_o = 37,2$  mm above an axisymmetric jet from a 80 mm pipe. The flow rate was measured with an orifice plate, and  $V$  was the average velocity in the pipe. A static pressure tapping was placed at the jet outlet.

A steam generator was placed upstream in the pipe, and the air passed over dry ice, thus producing a visible flow at the porous matrix surface. (See Fig A4.2).  
It was not possible to measure  $b$  with any reliability.



FIG A4.2 : LOW VELOCITY JET,  $v = 2,72$  m/s,  
EMERGING FROM MATRIX

It was observed that at high velocity, the smoke formed a jet above the matrix, whereas at low velocity, the smoke filtered slowly through the matrix, and had spread much more. The conclusion reached was that for large  $L$ , an initial 'turbulent' region existed, causing a high pressure drop. This caused the jet to spread, hence slowing down to the laminar Reynolds number range. Creep flow resulted in this region - the loss of inertia being clearly visible when the smoke merely filtered through the last row of spheres. The 3 short lines on Fig A4.5 diverged from the solid line for increasing  $Re_L$ . This was thought to have been caused by the small additional pressure drop through the 'Darcy' region. Fig A4.3 shows turbulent flow emerging from the porous matrix.



FIG A4.3 : TURBULENT JET :  $v = 6,5$  m/s,  
EMERGING FROM MATRIX

The full results were not reported, but two of the graphs plotted are presented in Figs A4.4 and A4.5.



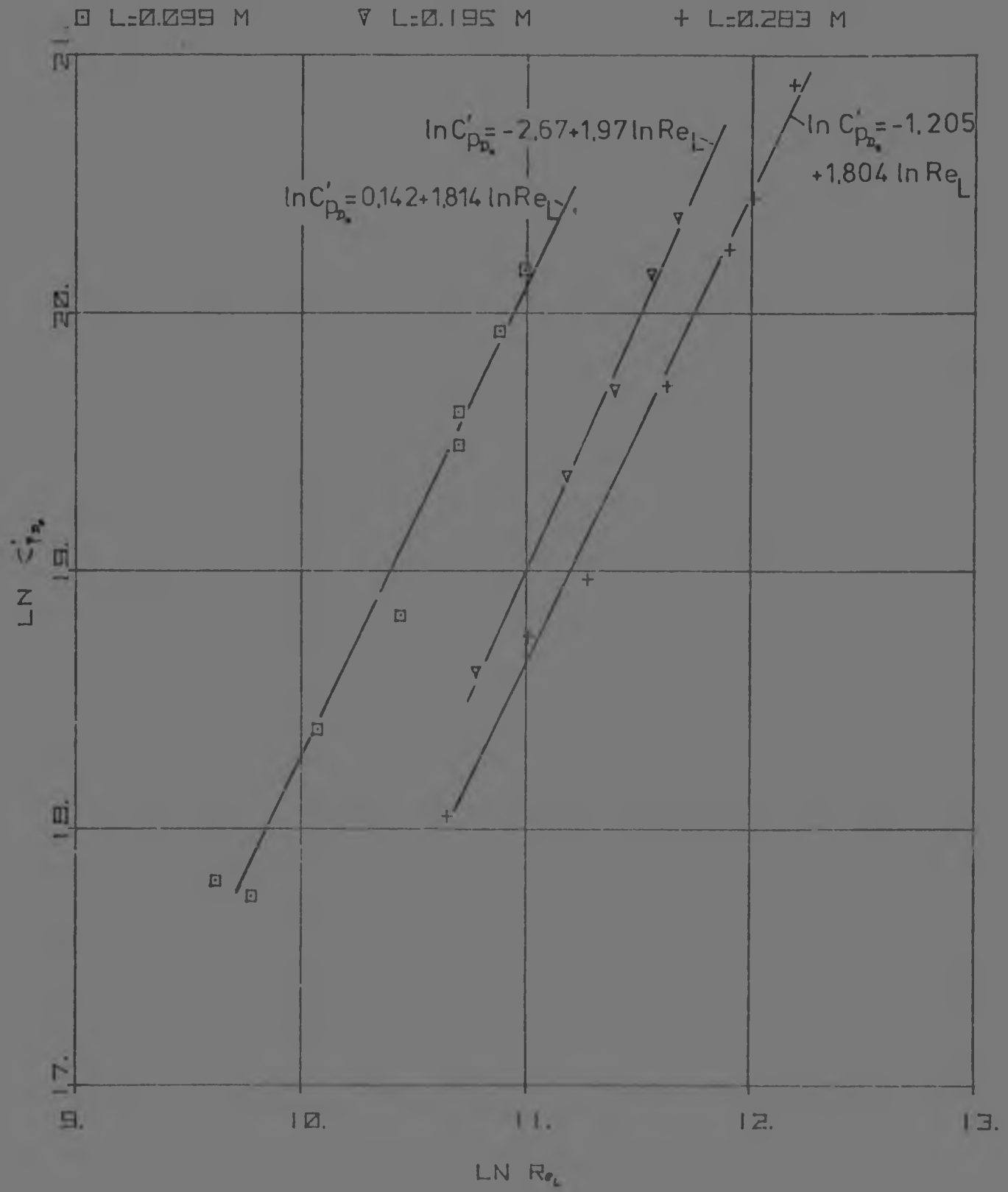


FIG A4.4 : PLOT OF  $C'_{p_{D_0}}$  vs  $Re_L$

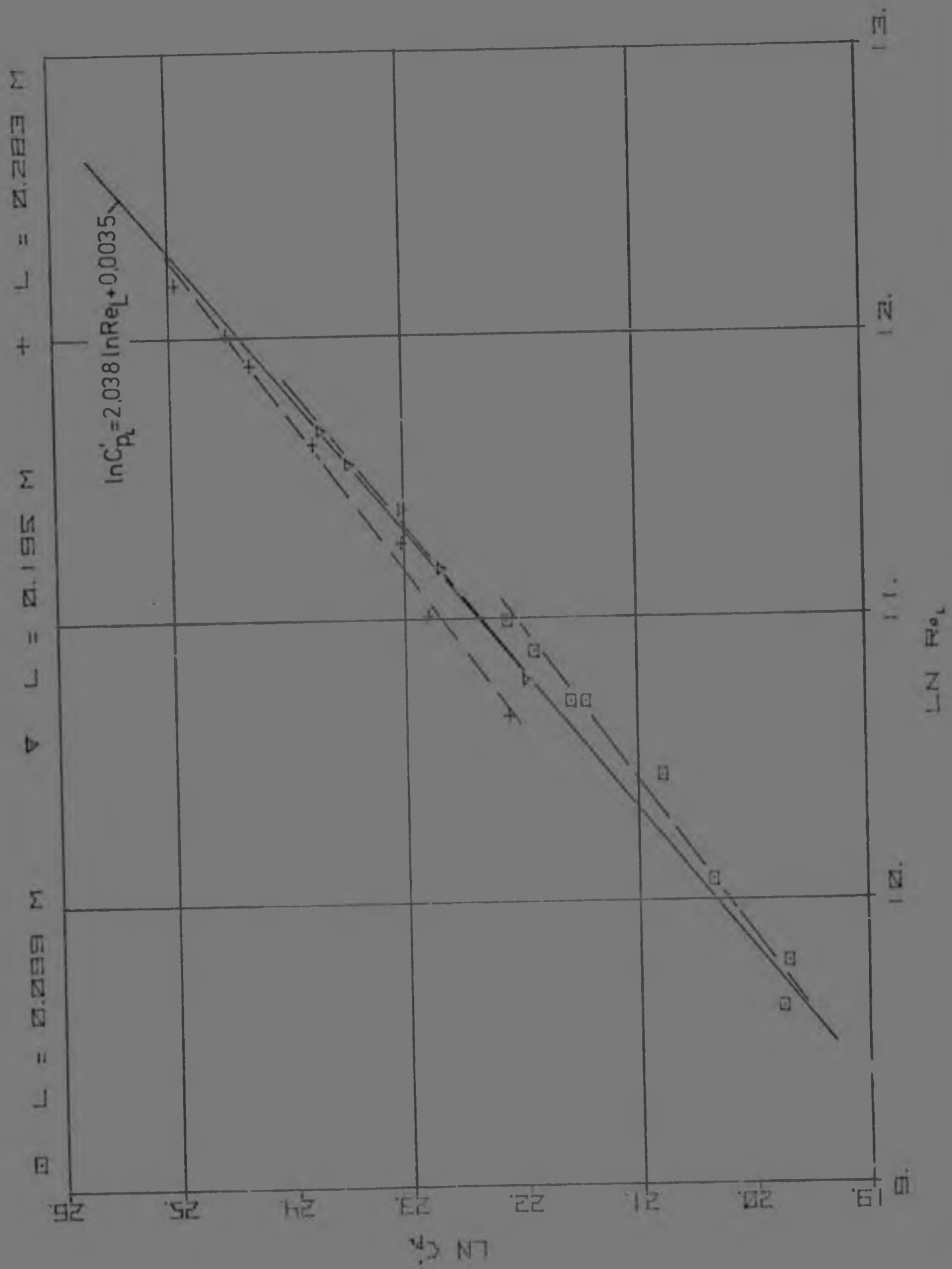


FIG A4.5 : PLOT OF  $C'_{PL}$  vs  $Re_L$

## APPENDIX B

### B1 Design of the Wind Tunnel, Fan and Contraction

Subsonic wind tunnel design has become a well documented exercise. Basic decisions such as whether the tunnel should be open or closed circuit, and type of fan are very much cost dependent.

The tunnel built for this project was to form part of an existing unit. After certain sections had been constructed, it was found to be preferable to be independent of other research programmes, and a tunnel was designed around the parts already built. This was unfortunate, because the tunnel could have been better designed.

The open circuit tunnel consisted of a centrifugal fan, diffuser, settling chamber, contraction, entry section, working section and an exit diffuser.

Materials used were wood strips (100 mm wide), banded to 3 mm thick masonite, the smooth side of which formed the inside walls of the tunnel. The edges were supported with slotted angle iron. An outside company made the settling chamber and contraction of mild steel sheet. Details of the working section appear in Appendix B2. The entire tunnel is supported at shoulder height on wooden legs.

A centrifugal fan, somewhat overpowered by a 22 kW squirrel cage induction motor, was made available for the project. This unit was used to blow air down the tunnel, rather than sucking the air through the tunnel. If the ducting was attached to the fan inlet, it was expected that swirl due to the rotation of the fan would be transmitted back to the Working Section. The removal of such a vortex may have required straightener vanes, a complication which was to be avoided. The profile at the fan outlet was highly non-

uniform, but this was to be controlled with screens. In the event, this proved more difficult than was expected. The profile in the entry section was much improved by the addition of honeycomb in the settling chamber. Aluminium honeycomb, 50 mm thick with a cell size of 6 mm, was used.

The contraction had an area ratio of 3:1. Theoretically this should have reduced the velocity fluctuations by a ninth. A potential flow analysis (see Cheers, 1945) was used to design the contraction. The equations were handled on a Hewlett-Packard 9810 calculator. The streamline denoted as contraction wall was plotted by the calculator (Fig B1.2).



Fig B1.1A : PLAN VIEW OF THE WIND TUNNEL, SHOWING  
(A) THE SETTLING CHAMBER, (B) THE CONTRAC-  
TION, (C) THE ENTRY SECTION, (D) THE  
WORKING SECTION

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uniform, but this was to be controlled with screens. In the event, this proved more difficult than was expected. The profile in the entry section was much improved by the addition of honeycomb in the settling chamber. Aluminium honeycomb, 50 mm thick with a cell size of 6 mm, was used.

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FIG B1.1A : PLAN VIEW OF THE WIND TUNNEL, SHOWING  
(A) THE SETTLING CHAMBER, (B) THE CONTRAC-  
TION, (C) THE ENTRY SECTION, (D) THE  
WORKING SECTION

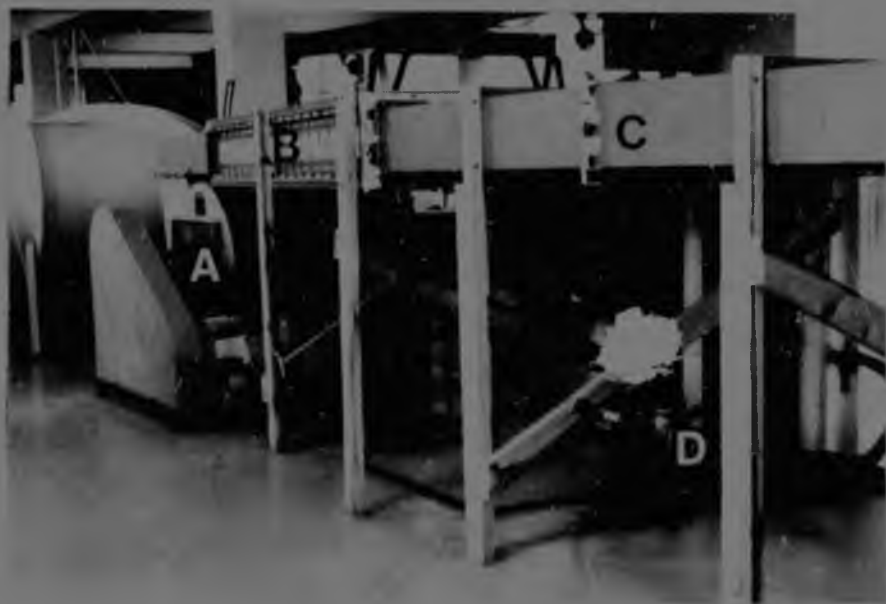


FIG B1.1B : SIDE VIEW OF THE TUNNEL WITH (A) THE FAN UNIT, (B) THE DIFFUSER, (C) SETTLING CHAMBER AND CONTRACTION, AND (D) THE COMPRESSED AIR SUPPLY

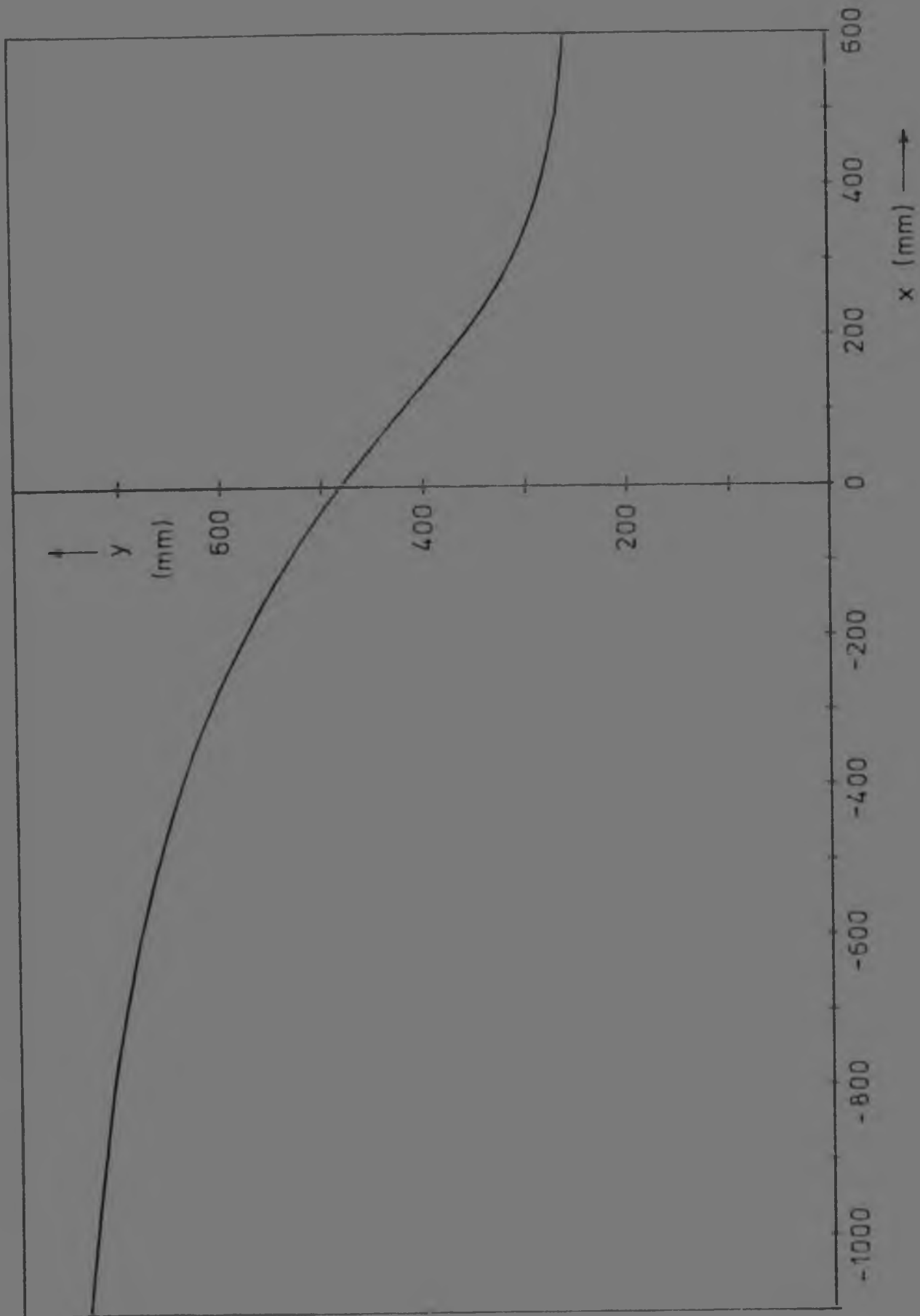


FIG B1.2 : CONTOUR OF WIND TUNNEL CONTRACTION

### 625mm DIA SI CD50 DUST FAN NO 32

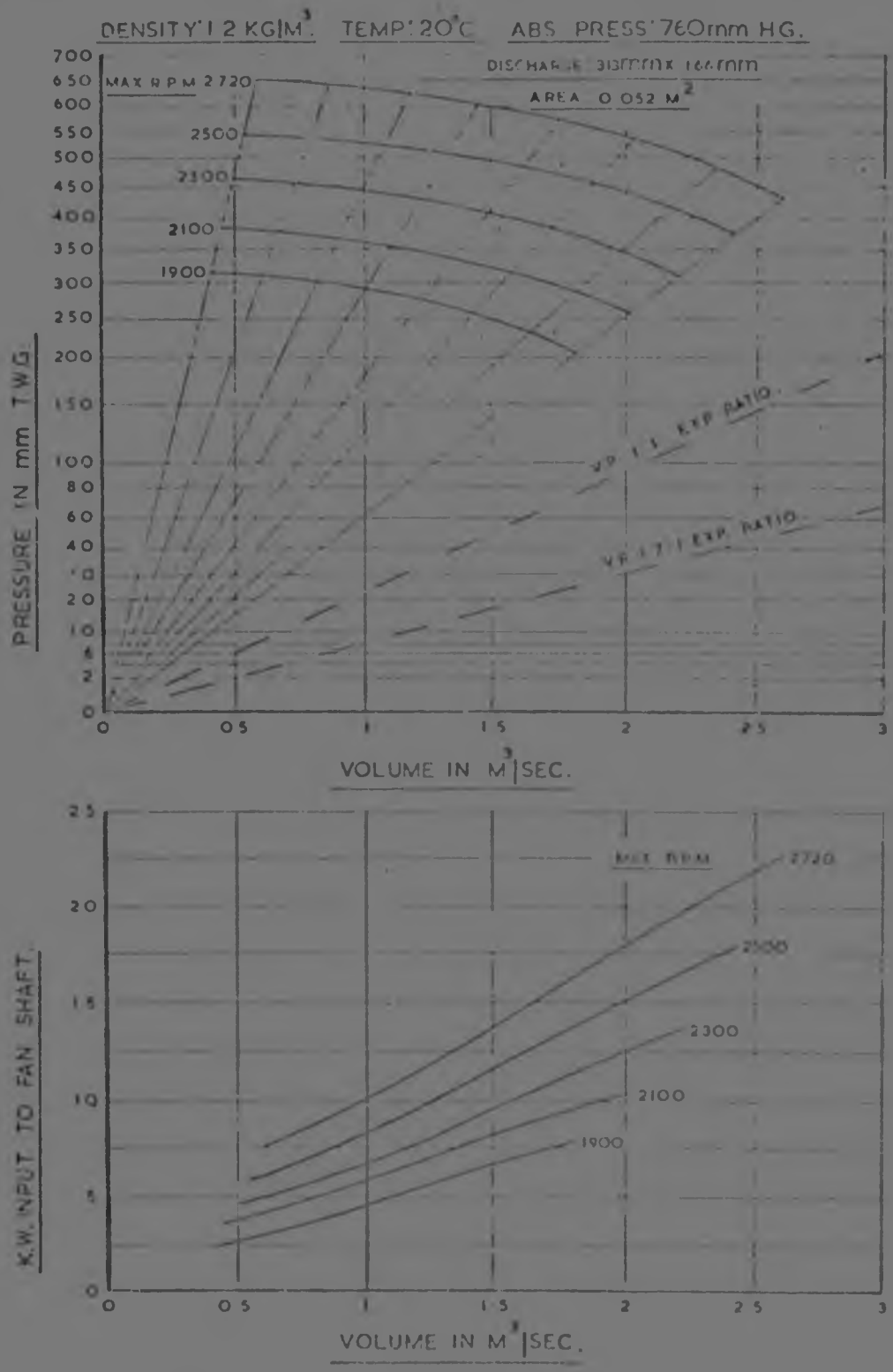


FIG B1.3 : FAN PERFORMANCE CURVE



## B2 The Working Section

A working section producing two dimensional flow with  $Re_L = 1,2 \times 10^6$  was required. The plate was to be solid, interspersed with a porous section through which a secondary flow could be introduced.

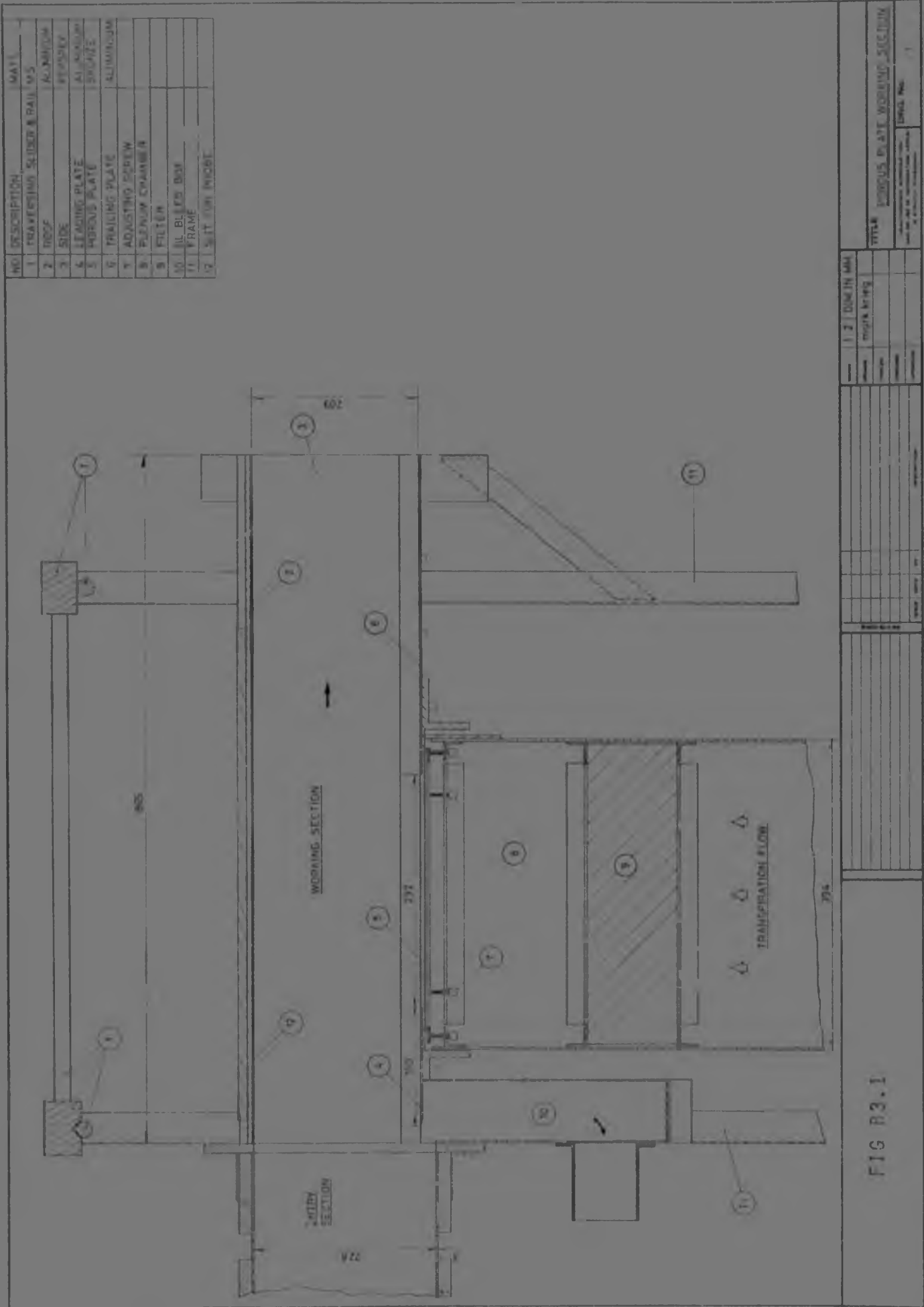
The working section was made rectangular, with dimensions 0,496 x 0,209 m. At a flowrate of 1,5 m<sup>3</sup>/s, a velocity of 17 m/s was expected in the jet. With the 1 m long, parallel duct ahead of the working section, various plates could be attached to the leading plate, allowing variation of  $Re_L$ . With a leading plate 0,9 m long, as used in most of the tests,  $Re_L = 1,2 \times 10^6$  was achieved at the end of the working section, which was 0,865 m long. A boundary layer bleed drew air off below the leading plate.

The plates and the roof of the working section were made of 3 mm aluminium sheet. Perspex (6,35 mm thick) was used for side walls. The porous plates had dimensions 500 x 244 x 6,35 mm. The porous plate leading and trailing edges were milled at 45°, and was pulled up to the aluminium plates with adjusting screws inside the plenum chamber. Filters were incorporated in the plenum chambers, to avoid the lower porous surface from becoming clogged with dirt or grease.

An exit diffuser was attached to the end of the working section. It was 1 m long, and prevented downstream conditions from affecting the flow in the jet. This section was on wheels, and could be removed to allow work inside the working section when setting up the hot wire transducer. The perspex sides were also removable.

## B3 Engineering Drawing

Fig B3.1 is a reduction of the assembly drawing of the Working Section.



## APPENDIX C

### C1 Calibration of the Orifice Plate

The orifice plate was calibrated according to BS 1042.

### C2 Calibration of Thermocouple

Copper-constantan thermocouples were used. The calibration curve is shown in Fig C2.1 below.

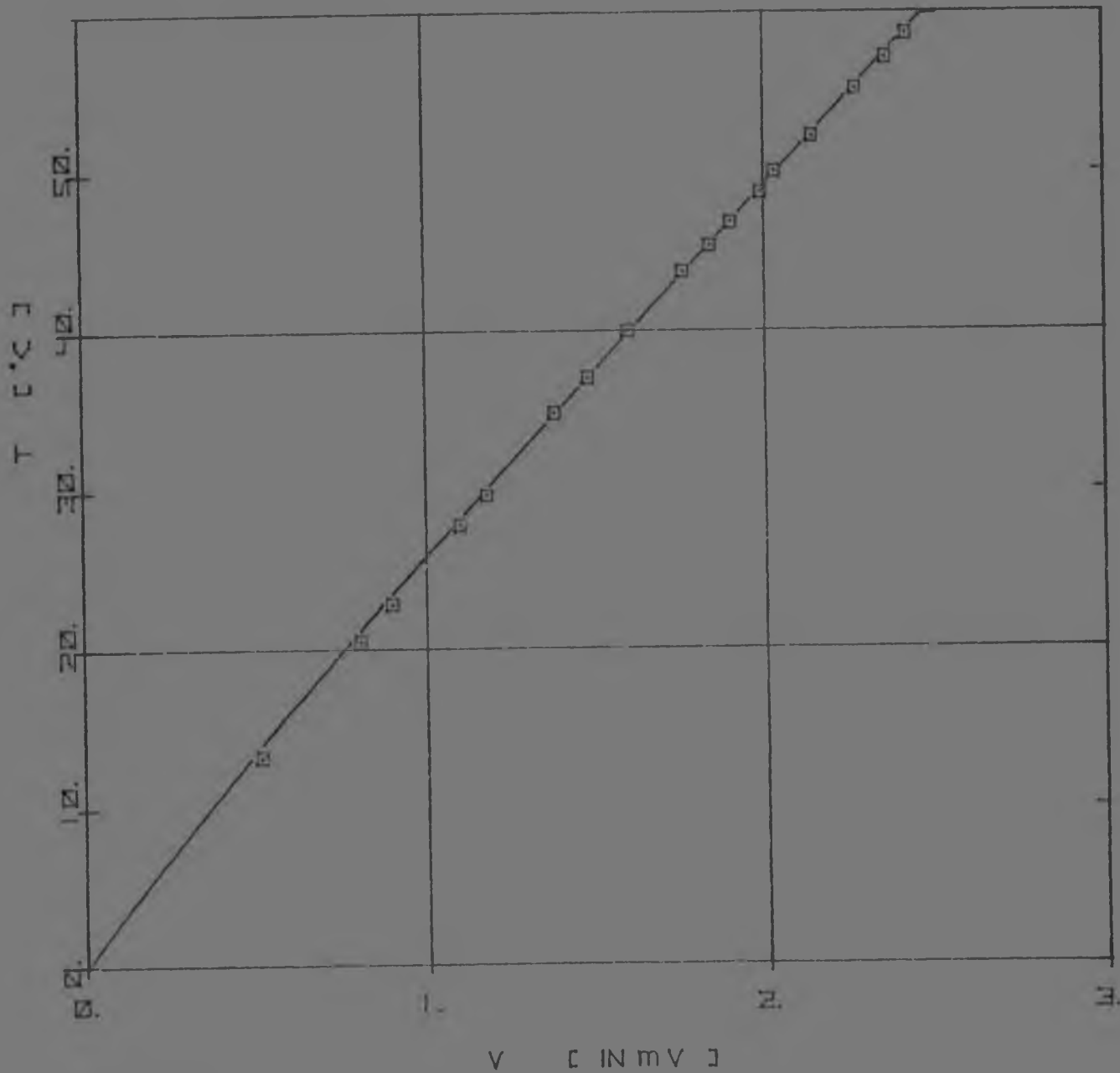


FIG C2.1

C3 The Calibration of the Hot Wires

The single sensor hot wires were calibrated using the Disa calibrating rig. This unit consisted of the Disa Pressure Control Unit, type 55 D 44, the Pressure Converter, type 55 D 46, and the nozzle unit, type 55 D 45. A Van Essen manometer model 6750 was used to measure the pressure drop across the nozzle.

The anemometer gain was adjusted using a square wave generator, and an oscilloscope to see when the amplifier became unstable. The zero flow voltage was determined, after amplification by the Auxilliary Unit, as was the hot wire resistance at the calibrating temperature.

Each hot wire (there were 4) was calibrated using the anemometer and the Auxilliary Unit.

King's Law was used to describe the calibration curves :-

$$V^2 = V_o'^2 + BU^n \quad \dots\dots [C3.1]$$

To obtain B and n, the data was plotted on log-log axes in the following form :-

$$\log\left\{\left(\frac{V}{V_o'}\right)^2 - 1\right\} = n \log U$$

Clearly, if the lhs was plotted against log U, the slope was n. Note that  $V_o' = \sigma V_o$ , where  $\sigma < 1$ . The value of  $\sigma$  was varied to find the value which gave the best straight line through the points.

The lineariser used the function

$$V_{LIN} = K(V^2 - V_o'^2)^{1/n} \quad \dots\dots [C3.2]$$

$$\therefore V_{LIN} = K(V_o'^2 - BU^n - V_o'^2)^{1/n} \quad \dots\dots [C3.3]$$

$$V_{LIN} = KBU = \sigma U \quad \dots\dots [C3.4]$$

to convert the non-linear anemometer output to a linear voltage. The exponent was set on the unit, and the hot wire was calibrated with the lineariser.

#### C4 Temperature Correction for the Anemometer Output

The heat loss from a hot-wire is influenced by the air velocity and by the air temperature. The anemometer output for a specific velocity must be corrected if the air temperature is different to the calibration temperature. This correction was performed from the following considerations :

The heat balance equation for a hot wire is :-

$$\frac{E^2}{R_w} = (T_w - T)h \quad \dots\dots [C4.1]$$

- $E$  = voltage drop across the hot wire
- $R_w$  = working resistance of the hot wire
- $T_w$  = working temperature of the hot wire
- $T$  = fluid temperature
- $h$  = heat transfer coefficient.

For constant  $h$ , the expression for the correction of the anemometer output for a change in  $T$  :-

$$E_n^2 = E^2 \frac{(T_w - T_n)}{(T_w - T)} \quad \dots\dots [C4.2]$$

- $E_n$  = corrected voltage
- $T_n$  = calibration temperature

The resistance of the hot wire changes with temperature according to :

$$R = R_n \{1 + \alpha(T - T_n)\} \quad \dots\dots [C4.3]$$

$R$  is the resistance at temperature  $T$ ,  $\alpha$  is the thermal resistance coefficient :-

$$\alpha = 0,0036 \text{ 1/}^\circ\text{C} \quad \dots\dots\dots [\text{C4.4}]$$

Equation C4.2 may be rewritten :-

$$E_n^2 = E^2 \frac{(R_{11} - R_{12})}{(R_{12} - R)} \quad \dots\dots\dots [\text{C4.5}]$$

R was calculated from [C4.3], and the corrected anemometer reading from [C4.5].

C5 Correction for Proximity of the Wall

It was expected that the heat loss from the hot wire would increase very near a wall ( $y < 2\text{mm}$ ). Various attempts have been made to correct for this apparent increase in velocity, see Dryden (1936).

The change in voltage  $\Delta V$ , can be calculated from

$$\Delta V = \sigma n \frac{(T_{11} - T)^2}{y^2}$$

which was given by Horsley (1975). The value of  $\sigma$  was estimated to be  $2,53 \times 10^{-6}$ , for  $y$  in mm.  $n$  was the exponent for the hot wire, discussed in C4.

## APPENDIX D

### D1 The Experimental Results

Presented here are the results of the velocity profiles measured in the boundary layer. The velocity measurements were corrected for temperature and the proximity of the wall.

#### Notes on experimental results

A power law was fitted to the input profile. It was given in the first line of printout. FD was  $u/u_0$ . The coefficient of the first right hand term was the inverse of the power calculated.

A virtual origin for the first profile was calculated from

$$\delta = 0,383x/Re_x^{1/5} \quad \dots\dots [D1.1]$$

$Re_x$  was based on this virtual origin. The difference in  $x$  length between subsequent stations was added to the first value.

The programme, DP2, performed the integrations for  $\delta^*$  and  $\theta$ , and plotted the results. The graphs are given in Fig 7.10.

\*\*\*\*\*  
 STATION 3  
 \*\*\*\*\*  
 UINF= 19.002 A/S X= 2.245  
 REX= 0.223E+07 RMF= 0.9985 KG/CM  
 TH= 17.40 DEG. NU= 0.1803E-04 MM/S  
 MU= 0.1803E-04 KG/MS

Y(MM)	U (A/S)	V(TTA)	U/UE	Y/DCL
0.0124	7.996	0.173	0.4178	0.0124
0.0124	8.125	0.174	0.4178	0.0124
0.0124	8.254	0.175	0.4178	0.0124
0.0124	8.383	0.176	0.4178	0.0124
0.0124	8.512	0.177	0.4178	0.0124
0.0124	8.641	0.178	0.4178	0.0124
0.0124	8.770	0.179	0.4178	0.0124
0.0124	8.899	0.180	0.4178	0.0124
0.0124	9.028	0.181	0.4178	0.0124
0.0124	9.157	0.182	0.4178	0.0124
0.0124	9.286	0.183	0.4178	0.0124
0.0124	9.415	0.184	0.4178	0.0124
0.0124	9.544	0.185	0.4178	0.0124
0.0124	9.673	0.186	0.4178	0.0124
0.0124	9.802	0.187	0.4178	0.0124
0.0124	9.931	0.188	0.4178	0.0124
0.0124	10.060	0.189	0.4178	0.0124
0.0124	10.189	0.190	0.4178	0.0124
0.0124	10.318	0.191	0.4178	0.0124
0.0124	10.447	0.192	0.4178	0.0124
0.0124	10.576	0.193	0.4178	0.0124
0.0124	10.705	0.194	0.4178	0.0124
0.0124	10.834	0.195	0.4178	0.0124
0.0124	10.963	0.196	0.4178	0.0124
0.0124	11.092	0.197	0.4178	0.0124
0.0124	11.221	0.198	0.4178	0.0124
0.0124	11.350	0.199	0.4178	0.0124
0.0124	11.479	0.200	0.4178	0.0124
0.0124	11.608	0.201	0.4178	0.0124
0.0124	11.737	0.202	0.4178	0.0124
0.0124	11.866	0.203	0.4178	0.0124
0.0124	11.995	0.204	0.4178	0.0124
0.0124	12.124	0.205	0.4178	0.0124
0.0124	12.253	0.206	0.4178	0.0124
0.0124	12.382	0.207	0.4178	0.0124
0.0124	12.511	0.208	0.4178	0.0124
0.0124	12.640	0.209	0.4178	0.0124
0.0124	12.769	0.210	0.4178	0.0124
0.0124	12.898	0.211	0.4178	0.0124
0.0124	13.027	0.212	0.4178	0.0124
0.0124	13.156	0.213	0.4178	0.0124
0.0124	13.285	0.214	0.4178	0.0124
0.0124	13.414	0.215	0.4178	0.0124
0.0124	13.543	0.216	0.4178	0.0124
0.0124	13.672	0.217	0.4178	0.0124
0.0124	13.801	0.218	0.4178	0.0124
0.0124	13.930	0.219	0.4178	0.0124
0.0124	14.059	0.220	0.4178	0.0124
0.0124	14.188	0.221	0.4178	0.0124
0.0124	14.317	0.222	0.4178	0.0124
0.0124	14.446	0.223	0.4178	0.0124
0.0124	14.575	0.224	0.4178	0.0124
0.0124	14.704	0.225	0.4178	0.0124
0.0124	14.833	0.226	0.4178	0.0124
0.0124	14.962	0.227	0.4178	0.0124
0.0124	15.091	0.228	0.4178	0.0124
0.0124	15.220	0.229	0.4178	0.0124
0.0124	15.349	0.230	0.4178	0.0124
0.0124	15.478	0.231	0.4178	0.0124
0.0124	15.607	0.232	0.4178	0.0124
0.0124	15.736	0.233	0.4178	0.0124
0.0124	15.865	0.234	0.4178	0.0124
0.0124	15.994	0.235	0.4178	0.0124
0.0124	16.123	0.236	0.4178	0.0124
0.0124	16.252	0.237	0.4178	0.0124
0.0124	16.381	0.238	0.4178	0.0124
0.0124	16.510	0.239	0.4178	0.0124
0.0124	16.639	0.240	0.4178	0.0124
0.0124	16.768	0.241	0.4178	0.0124
0.0124	16.897	0.242	0.4178	0.0124
0.0124	17.026	0.243	0.4178	0.0124
0.0124	17.155	0.244	0.4178	0.0124
0.0124	17.284	0.245	0.4178	0.0124
0.0124	17.413	0.246	0.4178	0.0124
0.0124	17.542	0.247	0.4178	0.0124
0.0124	17.671	0.248	0.4178	0.0124
0.0124	17.800	0.249	0.4178	0.0124
0.0124	17.929	0.250	0.4178	0.0124
0.0124	18.058	0.251	0.4178	0.0124
0.0124	18.187	0.252	0.4178	0.0124
0.0124	18.316	0.253	0.4178	0.0124
0.0124	18.445	0.254	0.4178	0.0124
0.0124	18.574	0.255	0.4178	0.0124
0.0124	18.703	0.256	0.4178	0.0124
0.0124	18.832	0.257	0.4178	0.0124
0.0124	18.961	0.258	0.4178	0.0124
0.0124	19.090	0.259	0.4178	0.0124
0.0124	19.219	0.260	0.4178	0.0124
0.0124	19.348	0.261	0.4178	0.0124
0.0124	19.477	0.262	0.4178	0.0124
0.0124	19.606	0.263	0.4178	0.0124
0.0124	19.735	0.264	0.4178	0.0124
0.0124	19.864	0.265	0.4178	0.0124
0.0124	19.993	0.266	0.4178	0.0124
0.0124	20.122	0.267	0.4178	0.0124
0.0124	20.251	0.268	0.4178	0.0124
0.0124	20.380	0.269	0.4178	0.0124
0.0124	20.509	0.270	0.4178	0.0124
0.0124	20.638	0.271	0.4178	0.0124
0.0124	20.767	0.272	0.4178	0.0124
0.0124	20.896	0.273	0.4178	0.0124
0.0124	21.025	0.274	0.4178	0.0124
0.0124	21.154	0.275	0.4178	0.0124
0.0124	21.283	0.276	0.4178	0.0124
0.0124	21.412	0.277	0.4178	0.0124
0.0124	21.541	0.278	0.4178	0.0124
0.0124	21.670	0.279	0.4178	0.0124
0.0124	21.799	0.280	0.4178	0.0124
0.0124	21.928	0.281	0.4178	0.0124
0.0124	22.057	0.282	0.4178	0.0124
0.0124	22.186	0.283	0.4178	0.0124
0.0124	22.315	0.284	0.4178	0.0124
0.0124	22.444	0.285	0.4178	0.0124
0.0124	22.573	0.286	0.4178	0.0124
0.0124	22.702	0.287	0.4178	0.0124
0.0124	22.831	0.288	0.4178	0.0124
0.0124	22.960	0.289	0.4178	0.0124
0.0124	23.089	0.290	0.4178	0.0124
0.0124	23.218	0.291	0.4178	0.0124
0.0124	23.347	0.292	0.4178	0.0124
0.0124	23.476	0.293	0.4178	0.0124
0.0124	23.605	0.294	0.4178	0.0124
0.0124	23.734	0.295	0.4178	0.0124
0.0124	23.863	0.296	0.4178	0.0124
0.0124	23.992	0.297	0.4178	0.0124
0.0124	24.121	0.298	0.4178	0.0124
0.0124	24.250	0.299	0.4178	0.0124
0.0124	24.379	0.300	0.4178	0.0124
0.0124	24.508	0.301	0.4178	0.0124
0.0124	24.637	0.302	0.4178	0.0124
0.0124	24.766	0.303	0.4178	0.0124
0.0124	24.895	0.304	0.4178	0.0124
0.0124	25.024	0.305	0.4178	0.0124
0.0124	25.153	0.306	0.4178	0.0124
0.0124	25.282	0.307	0.4178	0.0124
0.0124	25.411	0.308	0.4178	0.0124
0.0124	25.540	0.309	0.4178	0.0124
0.0124	25.669	0.310	0.4178	0.0124
0.0124	25.798	0.311	0.4178	0.0124
0.0124	25.927	0.312	0.4178	0.0124
0.0124	26.056	0.313	0.4178	0.0124
0.0124	26.185	0.314	0.4178	0.0124
0.0124	26.314	0.315	0.4178	0.0124
0.0124	26.443	0.316	0.4178	0.0124
0.0124	26.572	0.317	0.4178	0.0124
0.0124	26.701	0.318	0.4178	0.0124
0.0124	26.830	0.319	0.4178	0.0124
0.0124	26.959	0.320	0.4178	0.0124
0.0124	27.088	0.321	0.4178	0.0124
0.0124	27.217	0.322	0.4178	0.0124
0.0124	27.346	0.323	0.4178	0.0124
0.0124	27.475	0.324	0.4178	0.0124
0.0124	27.604	0.325	0.4178	0.0124
0.0124	27.733	0.326	0.4178	0.0124
0.0124	27.862	0.327	0.4178	0.0124
0.0124	27.991	0.328	0.4178	0.0124
0.0124	28.120	0.329	0.4178	0.0124
0.0124	28.249	0.330	0.4178	0.0124
0.0124	28.378	0.331	0.4178	0.0124
0.0124	28.507	0.332	0.4178	0.0124
0.0124	28.636	0.333	0.4178	0.0124
0.0124	28.765	0.334	0.4178	0.0124
0.0124	28.894	0.335	0.4178	0.0124
0.0124	29.023	0.336	0.4178	0.0124
0.0124	29.152	0.337	0.4178	0.0124
0.0124	29.281	0.338	0.4178	0.0124
0.0124	29.410	0.339	0.4178	0.0124
0.0124	29.539	0.340	0.4178	0.0124
0.0124	29.668	0.341	0.4178	0.0124
0.0124	29.797	0.342	0.4178	0.0124
0.0124	29.926	0.343	0.4178	0.0124
0.0124	30.055	0.344	0.4178	0.0124
0.0124	30.184	0.345	0.4178	0.0124
0.0124	30.313	0.346	0.4178	0.0124
0.0124	30.442	0.347	0.4178	0.0124
0.0124	30.571	0.348	0.4178	0.0124
0.0124	30.700	0.349	0.4178	0.0124
0.0124	30.829	0.350	0.4178	0.0124
0.0124	30.958	0.351	0.4178	0.0124
0.0124	31.087	0.352	0.4178	0.0124
0.0124	31.216	0.353	0.4178	0.0124
0.0124	31.345	0.354	0.4178	0.0124
0.0124	31.474	0.355	0.4178	0.0124
0.0124	31.603	0.356	0.4178	0.0124
0.0124	31.732	0.357	0.4178	0.0124
0.0124	31.861	0.358	0.4178	0.0124
0.0124	31.990	0.359	0.4178	0.0124
0.0124	32.119	0.360	0.4178	0.0124
0.0124	32.248	0.361	0.4178	0.0124
0.0124	32.377	0.362	0.4178	0.0124
0.0124	32.506	0.363	0.4178	0.0124
0.0124	32.635	0.364	0.4178	0.0124
0.0124	32.764	0.365	0.4178	0.0124
0.0124	32.893	0.366	0.4178	0.0124
0.0124	33.022	0.367	0.4178	0.0124
0.0124	33.151	0.368	0.4178	0.0124
0.0124	33.280	0.369	0.4178	0.0124
0.0124	33.409	0.370	0.4178	0.0124
0.0124	33.538	0.371	0.4178	0.0124
0.0124	33.667	0.372	0.4178	0.0124
0.0124	33.796	0.373	0.4178	0.0124
0.0124	33.925	0.374	0.4178	0.0124
0.0124	34.054	0.375	0.4178	0.0124
0.0124	34.183	0.376	0.4178	0.0124
0.0124	34.312	0.377	0.4178	0.0124
0.0124				



CURVE 3 SITE LOGICR3\* 3.689E-01\*LOG(Y/3)\* 6.105E-02

UINF= 18.216 M/S  
RHO= 0.1758E+01  
T= 2.91 DEG C  
MU= 0.1817E-04 KG/MS

E-10.57-17  
3051

DEL= 48.84 MM  
RDEL= 0.5409E+05

Y(MM)	U (M/S)	Y/TTA	L/UE	Y/DEL
0.00	11.674	0.171	0.6135	0.0125
0.05	12.574	0.222	0.6072	0.0236
0.10	13.474	0.273	0.6009	0.0347
0.15	14.374	0.324	0.5946	0.0458
0.20	15.274	0.375	0.5883	0.0569
0.25	16.174	0.426	0.5820	0.0680
0.30	17.074	0.477	0.5757	0.0791
0.35	17.974	0.528	0.5694	0.0902
0.40	18.874	0.579	0.5631	0.1013
0.45	19.774	0.630	0.5568	0.1124
0.50	20.674	0.681	0.5505	0.1235
0.55	21.574	0.732	0.5442	0.1346
0.60	22.474	0.783	0.5379	0.1457
0.65	23.374	0.834	0.5316	0.1568
0.70	24.274	0.885	0.5253	0.1679
0.75	25.174	0.936	0.5190	0.1790
0.80	26.074	0.987	0.5127	0.1901
0.85	26.974	1.038	0.5064	0.2012
0.90	27.874	1.089	0.5001	0.2123
0.95	28.774	1.140	0.4938	0.2234
1.00	29.674	1.191	0.4875	0.2345
1.05	30.574	1.242	0.4812	0.2456
1.10	31.474	1.293	0.4749	0.2567
1.15	32.374	1.344	0.4686	0.2678
1.20	33.274	1.395	0.4623	0.2789
1.25	34.174	1.446	0.4560	0.2900
1.30	35.074	1.497	0.4497	0.3011
1.35	35.974	1.548	0.4434	0.3122
1.40	36.874	1.599	0.4371	0.3233
1.45	37.774	1.650	0.4308	0.3344
1.50	38.674	1.701	0.4245	0.3455
1.55	39.574	1.752	0.4182	0.3566
1.60	40.474	1.803	0.4119	0.3677
1.65	41.374	1.854	0.4056	0.3788
1.70	42.274	1.905	0.3993	0.3899
1.75	43.174	1.956	0.3930	0.4010
1.80	44.074	2.007	0.3867	0.4121
1.85	44.974	2.058	0.3804	0.4232
1.90	45.874	2.109	0.3741	0.4343
1.95	46.774	2.160	0.3678	0.4454
2.00	47.674	2.211	0.3615	0.4565
2.05	48.574	2.262	0.3552	0.4676
2.10	49.474	2.313	0.3489	0.4787
2.15	50.374	2.364	0.3426	0.4898
2.20	51.274	2.415	0.3363	0.5009
2.25	52.174	2.466	0.3300	0.5120
2.30	53.074	2.517	0.3237	0.5231
2.35	53.974	2.568	0.3174	0.5342
2.40	54.874	2.619	0.3111	0.5453
2.45	55.774	2.670	0.3048	0.5564
2.50	56.674	2.721	0.2985	0.5675
2.55	57.574	2.772	0.2922	0.5786
2.60	58.474	2.823	0.2859	0.5897
2.65	59.374	2.874	0.2796	0.6008
2.70	60.274	2.925	0.2733	0.6119
2.75	61.174	2.976	0.2670	0.6230
2.80	62.074	3.027	0.2607	0.6341
2.85	62.974	3.078	0.2544	0.6452
2.90	63.874	3.129	0.2481	0.6563
2.95	64.774	3.180	0.2418	0.6674
3.00	65.674	3.231	0.2355	0.6785
3.05	66.574	3.282	0.2292	0.6896
3.10	67.474	3.333	0.2229	0.7007
3.15	68.374	3.384	0.2166	0.7118
3.20	69.274	3.435	0.2103	0.7229
3.25	70.174	3.486	0.2040	0.7340
3.30	71.074	3.537	0.1977	0.7451
3.35	71.974	3.588	0.1914	0.7562
3.40	72.874	3.639	0.1851	0.7673
3.45	73.774	3.690	0.1788	0.7784
3.50	74.674	3.741	0.1725	0.7895
3.55	75.574	3.792	0.1662	0.8006
3.60	76.474	3.843	0.1599	0.8117
3.65	77.374	3.894	0.1536	0.8228
3.70	78.274	3.945	0.1473	0.8339
3.75	79.174	3.996	0.1410	0.8450
3.80	80.074	4.047	0.1347	0.8561
3.85	80.974	4.098	0.1284	0.8672
3.90	81.874	4.149	0.1221	0.8783
3.95	82.774	4.200	0.1158	0.8894
4.00	83.674	4.251	0.1095	0.9005
4.05	84.574	4.302	0.1032	0.9116
4.10	85.474	4.353	0.0969	0.9227
4.15	86.374	4.404	0.0906	0.9338
4.20	87.274	4.455	0.0843	0.9449
4.25	88.174	4.506	0.0780	0.9560
4.30	89.074	4.557	0.0717	0.9671
4.35	89.974	4.608	0.0654	0.9782
4.40	90.874	4.659	0.0591	0.9893
4.45	91.774	4.710	0.0528	1.0004
4.50	92.674	4.761	0.0465	1.0115
4.55	93.574	4.812	0.0402	1.0226
4.60	94.474	4.863	0.0339	1.0337
4.65	95.374	4.914	0.0276	1.0448
4.70	96.274	4.965	0.0213	1.0559
4.75	97.174	5.016	0.0150	1.0670
4.80	98.074	5.067	0.0087	1.0781
4.85	98.974	5.118	0.0024	1.0892
4.90	99.874	5.169	0.0000	1.1003
5.00	100.000	5.200	0.0000	1.1114

DEL= 48.84 MM  
RDEL= 0.5409E+05  
SHAPE FACTOR= 1.27

UINF= 18.216 M/S  
RHO= 0.1758E+01  
T= 2.91 DEG C  
MU= 0.1817E-04 KG/MS

DEL= 48.84 MM  
RDEL= 0.5409E+05  
SHAPE FACTOR= 1.27

Y(MM)	U (M/S)	Y/TTA	L/UE	Y/DEL
0.00	7.755	0.146	0.6278	0.0125
0.05	8.255	0.197	0.6215	0.0236
0.10	8.755	0.248	0.6152	0.0347
0.15	9.255	0.299	0.6089	0.0458
0.20	9.755	0.350	0.6026	0.0569
0.25	10.255	0.401	0.5963	0.0680
0.30	10.755	0.452	0.5900	0.0791
0.35	11.255	0.503	0.5837	0.0902
0.40	11.755	0.554	0.5774	0.1013
0.45	12.255	0.605	0.5711	0.1124
0.50	12.755	0.656	0.5648	0.1235
0.55	13.255	0.707	0.5585	0.1346
0.60	13.755	0.758	0.5522	0.1457
0.65	14.255	0.809	0.5459	0.1568
0.70	14.755	0.860	0.5396	0.1679
0.75	15.255	0.911	0.5333	0.1790
0.80	15.755	0.962	0.5270	0.1901
0.85	16.255	1.013	0.5207	0.2012
0.90	16.755	1.064	0.5144	0.2123
0.95	17.255	1.115	0.5081	0.2234
1.00	17.755	1.166	0.5018	0.2345
1.05	18.255	1.217	0.4955	0.2456
1.10	18.755	1.268	0.4892	0.2567
1.15	19.255	1.319	0.4829	0.2678
1.20	19.755	1.370	0.4766	0.2789
1.25	20.255	1.421	0.4703	0.2900
1.30	20.755	1.472	0.4640	0.3011
1.35	21.255	1.523	0.4577	0.3122
1.40	21.755	1.574	0.4514	0.3233
1.45	22.255	1.625	0.4451	0.3344
1.50	22.755	1.676	0.4388	0.3455
1.55	23.255	1.727	0.4325	0.3566
1.60	23.755	1.778	0.4262	0.3677
1.65	24.255	1.829	0.4199	0.3788
1.70	24.755	1.880	0.4136	0.3899
1.75	25.255	1.931	0.4073	0.4010
1.80	25.755	1.982	0.4010	0.4121
1.85	26.255	2.033	0.3947	0.4232
1.90	26.755	2.084	0.3884	0.4343
1.95	27.255	2.135	0.3821	0.4454
2.00	27.755	2.186	0.3758	0.4565
2.05	28.255	2.237	0.3695	0.4676
2.10	28.755	2.288	0.3632	0.4787
2.15	29.255	2.339	0.3569	0.4898
2.20	29.755	2.390	0.3506	0.5009
2.25	30.255	2.441	0.3443	0.5120
2.30	30.755	2.492	0.3380	0.5231
2.35	31.255	2.543	0.3317	0.5342
2.40	31.755	2.594	0.3254	0.5453
2.45	32.255	2.645	0.3191	0.5564
2.50	32.755	2.696	0.3128	0.5675
2.55	33.255	2.747	0.3065	0.5786
2.60	33.755	2.798	0.3002	0.5897
2.65	34.255	2.849	0.2939	0.6008
2.70	34.755	2.900	0.2876	0.6119
2.75	35.255	2.951	0.2813	0.6230
2.80	35.755	3.002	0.2750	0.6341
2.85	36.255	3.053	0.2687	0.6452
2.90	36.755	3.104	0.2624	0.6563
2.95	37.255	3.155	0.2561	0.6674
3.00	37.755	3.206	0.2498	0.6785
3.05	38.255	3.257	0.2435	0.6896
3.10	38.755	3.308	0.2372	0.7007
3.15	39.255	3.359	0.2309	0.7118
3.20	39.755	3.410	0.2246	0.7229
3.25	40.255	3.461	0.2183	0.7340
3.30	40.755	3.512	0.2120	0.7451
3.35	41.255	3.563	0.2057	0.7562
3.40	41.755	3.614	0.1994	0.7673
3.45	42.255	3.665	0.1931	0.7784
3.50	42.755	3.716	0.1868	0.7895
3.55	43.255	3.767	0.1805	0.8006
3.60	43.755	3.818	0.1742	0.8117
3.65	44.255	3.869	0.1679	0.8228
3.70	44.755	3.920	0.1616	0.8339
3.75	45.255	3.971	0.1553	0.8450
3.80	45.755	4.022	0.1490	0.8561
3.85	46.255	4.073	0.1427	0.8672
3.90	46.755	4.124	0.1364	0.8783
3.95	47.255	4.175	0.1301	0.8894
4.00	47.755	4.226	0.1238	0.9005
4.05	48.255	4.277	0.1175	0.9116
4.10	48.755	4.328	0.1112	0.9227
4.15	49.255	4.379	0.1049	0.9338
4.20	49.755	4.430	0.0986	0.9449
4.25	50.255	4.481	0.0923	0.9560
4.30	50.755	4.532	0.0860	0.9671
4.35	51.255	4.583	0.0797	0.9782
4.40	51.755	4.634	0.0734	0.9893
4.45	52.255	4.685	0.0671	1.0004
4.50	52.755	4.736	0.0608	1.0115
4.55	53.255	4.787	0.0545	1.0226
4.60	53.755	4.838	0.0482	1.0337
4.65	54.255	4.889	0.0419	1.0448
4.70	54.755	4.940	0.0356	1.0559
4.75	55.255	4.991	0.0293	1.0670
4.80	55.755	5.042	0.0230	1.0781
4.85	56.255	5.093	0.0167	1.0892
4.90	56.755	5.144	0.0104	1.1003
4.95	57.255	5.195	0.0041	1.1114
5.00	57.500	5.200	0.0000	1.1114

DEL= 48.84 MM  
RDEL= 0.5409E+05  
SHAPE FACTOR= 1.27

UINF= 19.746 M/S  
RHO= 0.199E+07  
T= 12.85 DEG C  
MU= 0.1825E-04 KG/MS







STATION 2  
 X= 2.005  
 RHO= 0.9779 KG/CM<sup>3</sup>  
 NU= 0.1871E-04 M<sup>2</sup>/S

Y (MM)	U (M/S)	V (M/S)	U/V	V/U	Y/DEL
0.62	7.150	0.133	0.0185	0.0540	0.0117
1.25	6.402	0.225	0.0351	0.0934	0.0229
1.75	5.710	0.375	0.0657	0.1536	0.0341
2.25	5.075	0.509	0.1004	0.2140	0.0453
2.75	4.490	0.673	0.1499	0.2844	0.0565
3.25	3.950	0.875	0.2191	0.3648	0.0677
3.75	3.450	1.100	0.3143	0.4552	0.0789
4.25	2.990	1.350	0.4478	0.5556	0.0901
4.75	2.570	1.620	0.6226	0.6750	0.1013
5.25	2.190	1.910	0.8611	0.8143	0.1125
5.75	1.850	2.220	1.2054	0.9736	0.1237
6.25	1.540	2.550	1.6519	1.1529	0.1349
6.75	1.260	2.900	2.2937	1.3522	0.1461
7.25	1.010	3.270	3.2471	1.5715	0.1573
7.75	0.790	3.660	4.6341	1.8108	0.1685
8.25	0.600	4.070	6.7000	2.0701	0.1797
8.75	0.440	4.500	10.2271	2.3494	0.1909
9.25	0.310	4.950	15.8054	2.6487	0.2021
9.75	0.210	5.420	25.7341	2.9680	0.2133
10.25	0.140	5.910	42.1529	3.3073	0.2245
10.75	0.090	6.420	71.0715	3.6666	0.2357
11.25	0.060	6.950	117.4898	4.0459	0.2469
11.75	0.040	7.500	197.4079	4.4452	0.2581
12.25	0.030	8.070	337.8259	4.8645	0.2693
12.75	0.020	8.660	587.8438	5.3038	0.2805
13.25	0.015	9.270	1017.3617	5.7631	0.2917
13.75	0.010	9.900	1807.3796	6.2424	0.3029
14.25	0.007	10.550	3157.3975	6.7417	0.3141
14.75	0.005	11.220	5467.4154	7.2610	0.3253
15.25	0.003	11.910	9617.4333	7.8003	0.3365
15.75	0.002	12.620	16767.4512	8.3596	0.3477
16.25	0.001	13.350	29917.4691	8.9389	0.3589
16.75	0.000	14.100	52067.4870	9.5382	0.3701
17.25	0.000	14.870	91217.5049	10.1575	0.3813
17.75	0.000	15.660	159767.5228	10.7968	0.3925
18.25	0.000	16.470	280317.5407	11.4561	0.4037
18.75	0.000	17.300	475867.5586	12.1354	0.4149
19.25	0.000	18.150	840417.5765	12.8347	0.4261
19.75	0.000	19.020	1474967.5944	13.5540	0.4373
20.25	0.000	19.910	2599517.6123	14.2933	0.4485
20.75	0.000	20.820	4524067.6302	15.0526	0.4597
21.25	0.000	21.750	7848617.6481	15.8319	0.4709
21.75	0.000	22.700	13773167.6660	16.6312	0.4821
22.25	0.000	23.670	24097717.6839	17.4505	0.4933
22.75	0.000	24.660	41842267.7018	18.2898	0.5045
23.25	0.000	25.670	73086817.7197	19.1491	0.5157
23.75	0.000	26.700	12653237.7376	19.9284	0.5269
24.25	0.000	27.750	21997787.7555	20.7277	0.5381
24.75	0.000	28.820	39242337.7734	21.5470	0.5493
25.25	0.000	29.910	67386887.7913	22.3863	0.5605
25.75	0.000	31.020	11943237.8092	23.2456	0.5717
26.25	0.000	32.150	21047787.8271	24.1249	0.5829
26.75	0.000	33.300	37292337.8450	25.0242	0.5941
27.25	0.000	34.470	64636887.8629	25.9435	0.6053
27.75	0.000	35.660	11408237.8808	26.8828	0.6165
28.25	0.000	36.870	20412787.8987	27.8421	0.6277
28.75	0.000	38.100	36557337.9166	28.8214	0.6389
29.25	0.000	39.350	64801887.9345	29.8207	0.6501
29.75	0.000	40.620	11524737.9524	30.8399	0.6613
30.25	0.000	41.910	20769287.9703	31.8692	0.6725
30.75	0.000	43.220	37213837.9882	32.9185	0.6837
31.25	0.000	44.550	64858387.0061	33.9878	0.6949
31.75	0.000	45.900	11570287.0240	35.0771	0.7061
32.25	0.000	47.270	20914837.0419	36.1864	0.7173
32.75	0.000	48.660	37559387.0598	37.3157	0.7285
33.25	0.000	50.070	66503937.0777	38.4650	0.7397
33.75	0.000	51.500	11884843.0956	39.6343	0.7509
34.25	0.000	52.950	21429293.1135	40.8236	0.7621
34.75	0.000	54.420	38373843.1314	42.0329	0.7733
35.25	0.000	55.910	66718393.1493	43.2622	0.7845
35.75	0.000	57.420	11806289.1672	44.5115	0.7957
36.25	0.000	58.950	21350739.1851	45.7808	0.8069
36.75	0.000	60.500	38295289.2030	47.0701	0.8181
37.25	0.000	62.070	66739839.2209	48.3894	0.8293
37.75	0.000	63.660	11818439.2388	49.7387	0.8405
38.25	0.000	65.270	21362889.2567	51.1180	0.8517
38.75	0.000	66.900	38307439.2746	52.5273	0.8629
39.25	0.000	68.550	66751989.2925	53.9666	0.8741
39.75	0.000	70.220	11819639.3104	55.4359	0.8853
40.25	0.000	71.910	21364089.3283	56.9352	0.8965
40.75	0.000	73.620	38308639.3462	58.4645	0.9077
41.25	0.000	75.350	66753189.3641	59.9238	0.9189
41.75	0.000	77.100	11819769.3820	61.4131	0.9301
42.25	0.000	78.870	21364219.4000	62.9324	0.9413
42.75	0.000	80.660	38308769.4179	64.4817	0.9525
43.25	0.000	82.470	66753319.4358	66.0610	0.9637
43.75	0.000	84.300	11819819.4537	67.6703	0.9749
44.25	0.000	86.150	21364369.4716	69.3096	0.9861
44.75	0.000	88.020	38308919.4895	70.9789	0.9973
45.25	0.000	89.910	66754069.5074	72.6782	1.0085
45.75	0.000	91.820	11819969.5253	74.4075	1.0197
46.25	0.000	93.750	21364519.5432	76.1668	1.0309
46.75	0.000	95.700	38309069.5611	77.9561	1.0421
47.25	0.000	97.670	66754819.5790	79.7754	1.0533
47.75	0.000	99.660	11820119.5969	81.6247	1.0645
48.25	0.000	101.670	21364669.6148	83.5040	1.0757
48.75	0.000	103.700	38309219.6327	85.4133	1.0869
49.25	0.000	105.750	66755069.6506	87.3526	1.0981
49.75	0.000	107.820	11820269.6685	89.3219	1.1093
50.25	0.000	109.910	21364819.6864	91.3212	1.1205
50.75	0.000	112.020	38309369.7043	93.3505	1.1317
51.25	0.000	114.150	66755819.7222	95.4098	1.1429
51.75	0.000	116.300	11820419.7401	97.4991	1.1541
52.25	0.000	118.470	21364969.7580	99.6184	1.1653
52.75	0.000	120.660	38309519.7759	101.7677	1.1765
53.25	0.000	122.870	66756069.7938	103.9470	1.1877
53.75	0.000	125.100	11820569.8117	106.1563	1.1989
54.25	0.000	127.350	21365119.8296	108.3956	1.2101
54.75	0.000	129.620	38309669.8475	110.6649	1.2213
55.25	0.000	131.910	66756319.8654	112.9642	1.2325
55.75	0.000	134.220	11820719.8833	115.2935	1.2437
56.25	0.000	136.550	21365269.9012	117.6528	1.2549
56.75	0.000	138.900	38309819.9191	120.0421	1.2661
57.25	0.000	141.270	66756569.9370	122.4614	1.2773
57.75	0.000	143.660	11820869.9549	124.9107	1.2885
58.25	0.000	146.070	21365419.9728	127.3899	1.2997
58.75	0.000	148.500	38309969.9907	129.9092	1.3109
59.25	0.000	150.950	66756819.0086	132.4685	1.3221
59.75	0.000	153.420	11821019.0265	135.0678	1.3333
60.25	0.000	155.910	21365569.0444	137.7071	1.3445
60.75	0.000	158.420	38310119.0623	140.3864	1.3557
61.25	0.000	160.950	66757069.0802	143.1057	1.3669
61.75	0.000	163.500	11821169.0981	145.8650	1.3781
62.25	0.000	166.070	21365719.1160	148.6643	1.3893
62.75	0.000	168.660	38310269.1339	151.5036	1.4005
63.25	0.000	171.270	66757319.1518	154.3829	1.4117
63.75	0.000	173.900	11821269.1697	157.3022	1.4229
64.25	0.000	176.550	21365869.1876	160.2615	1.4341
64.75	0.000	179.220	38310419.2055	163.2608	1.4453
65.25	0.000	181.910	66757569.2234	166.2999	1.4565
65.75	0.000	184.620	11821369.2413	169.3792	1.4677
66.25	0.000	187.350	21366019.2592	172.4985	1.4789
66.75	0.000	190.100	38310569.2771	175.6578	1.4901
67.25	0.000	192.870	66757819.2950	178.8571	1.5013
67.75	0.000	195.660	11821469.3129	182.0964	1.5125
68.25	0.000	198.470	21366169.3308	185.3757	1.5237
68.75	0.000	201.300	38310719.3487	188.6950	1.5349
69.25	0.000	204.150	66758069.3666	192.0543	1.5461
69.75	0.000	207.020	11821569.3845	195.4536	1.5573
70.25	0.000	209.910	21366319.4024	198.8929	1.5685
70.75	0.000	212.820	38310869.4203	202.3722	1.5797
71.25	0.000	215.750	66758319.4382	205.8915	1.5909
71.75	0.000	218.700	11821669.4561	209.4508	1.6021
72.25	0.000	221.670	21366469.4740	213.0501	1.6133
72.75	0.000	224.660	38311019.4919	216.6894	1.6245
73.25	0.000	227.670	66758569.5098	220.3687	1.6357
73.75	0.000	230.700	11821769.5277	224.0880	1.6469
74.25	0.000	233.750	21366619.5456	227.8473	1.6581
74.75	0.000	236.820	38311169.5635	231.6466	1.6693
75.25	0.000	239.910	66758819.5814	235.4859	1.6805
75.75	0.000	243.020	11821869.6013	239.3652	1.6917
76.25	0.000	246.150	21366769.6192	243.2845	1.7029
76.75	0.000	249.300	38311319.6371	247.2438	1.7141
77.25	0.000	252.470	66759069.6550	251.2431	1.7253
77.75	0.000	255.660	11821969.6729	255.2824	1.7365
78.25	0.000	258.870	21366919.6908	259.3617	1.7477
78.75	0.000	262.100	38311469.7087	263.4810	1.7589
79.25	0.000	265.350	66759319.7266	267.6403	1.7701
79.75	0.000	268.620	11822069.7445	271.8396	1.7813
80.25	0.000	271.910	21367069.7624	276.07	

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0.002

LINE# 15.761 4/5

REX# 0.1800000000

T# 23.09 DEG.S

MJM# 0.1832E+04 KG/M5

X# 2.221

RHO# 0.9757 KG/CM.M

NU# 0.1812E+04 M.M/5

DEL# 0.002 MM

POST# 0.0000000000

THETA# 0.00 MM

RDT# 0.0000000000

RTTA# 0.0000000000

START# 0.0000000000

STOP# 0.0000000000

LINE# 15.761 4/5

REX# 0.1800000000

T# 23.09 DEG.S

MJM# 0.1832E+04 KG/M5

X# 2.221

RHO# 0.9757 KG/CM.M

NU# 0.1812E+04 M.M/5

DEL# 0.002 MM

POST# 0.0000000000

THETA# 0.00 MM

RDT# 0.0000000000

RTTA# 0.0000000000

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STOP# 0.0000000000

LINE# 15.761 4/5

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MJM# 0.1832E+04 KG/M5

X# 2.221

RHO# 0.9757 KG/CM.M

NU# 0.1812E+04 M.M/5

DEL# 0.002 MM

POST# 0.0000000000

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RHO# 0.9757 KG/CM.M

NU# 0.1812E+04 M.M/5

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NU# 0.1812E+04 M.M/5

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RHO# 0.9757 KG/CM.M

NU# 0.1812E+04 M.M/5

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THETA# 0.00 MM

RDT# 0.0000000000

RTTA# 0.0000000000

START# 0.0000000000

STOP# 0.0000000000

VIEW# 0.002

U (M/S) 6.34

Y/TTA 0.175

U/UE 0.3121

Y/DFL 0.3134

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VIEW# 0.002

U (M/S) 6.34

Y/TTA 0.175

U/UE 0.3121

Y/DFL 0.3134

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VIEW# 0.002

U (M/S) 6.34

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U/UE 0.3121

Y/DFL 0.3134

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VIEW# 0.002

U (M/S) 6.34

Y/TTA 0.175

U/UE 0.3121

Y/DFL 0.3134

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STATION 3  
 UINF= 10.144 M/S  
 R-X= 1.2000E+07  
 T= 20.80 DEG.C  
 MJ= 0.1836E-04 KG/MS  
 X= 2.369  
 RHO= 0.0737 KG/CM<sup>3</sup>  
 NU= 0.1888E-04 MM<sup>2</sup>/S  
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C-7.07-17

09M

OTW 10.143 WFF  
EPR 5.107647  
MIP 8.10247-54 WGMF

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OTW 10.143 WFF  
EPR 5.107647  
MIP 8.10247-54 WGMF

STATION 1  
Y = 2.327  
RHT = 0.0033 KG/CU.M  
N = 0.1025E+04 MM/S

UINF = 16.328 M/S  
RFX = 0.1071E+07  
TX = 24.31 DEG  
MUF = 0.1851E-14 KG/MS

Y(M)	U (M/S)	Y(TTA)	U(UF)	Y(DEL)
0.55	0.739	0.100	0.2966	0.3136
1.10	0.731	0.180	0.3511	0.3181
1.65	0.726	0.201	0.4251	0.3208
2.20	0.721	0.205	0.4613	0.3254
2.75	0.716	0.204	0.5220	0.3358
3.30	0.711	0.201	0.5734	0.3481
3.85	0.706	0.196	0.6194	0.3620
4.40	0.701	0.190	0.6600	0.3773
4.95	0.696	0.183	0.6953	0.3939
5.50	0.691	0.176	0.7255	0.4116
6.05	0.686	0.167	0.7507	0.4303
6.60	0.681	0.157	0.7710	0.4499
7.15	0.676	0.146	0.7864	0.4704
7.70	0.671	0.134	0.7969	0.4918
8.25	0.666	0.121	0.8025	0.5141
8.80	0.661	0.107	0.8033	0.5373
9.35	0.656	0.093	0.8003	0.5614
9.90	0.651	0.078	0.7937	0.5864
10.45	0.646	0.062	0.7835	0.6123
11.00	0.641	0.045	0.7698	0.6391
11.55	0.636	0.027	0.7526	0.6668
12.10	0.631	0.009	0.7320	0.6954
12.65	0.626	0.000	0.7081	0.7249
13.20	0.621	0.000	0.6808	0.7553
13.75	0.616	0.000	0.6501	0.7866
14.30	0.611	0.000	0.6161	0.8188
14.85	0.606	0.000	0.5788	0.8520
15.40	0.601	0.000	0.5381	0.8862
15.95	0.596	0.000	0.4939	0.9214
16.50	0.591	0.000	0.4472	0.9576
17.05	0.586	0.000	0.3981	0.9948
17.60	0.581	0.000	0.3466	1.0330
18.15	0.576	0.000	0.2927	1.0722
18.70	0.571	0.000	0.2374	1.1124
19.25	0.566	0.000	0.1808	1.1536
19.80	0.561	0.000	0.1229	1.1958
20.35	0.556	0.000	0.0637	1.2390
20.90	0.551	0.000	0.0033	1.2832
21.45	0.546	0.000	0.0000	1.3284
22.00	0.541	0.000	0.0000	1.3746
22.55	0.536	0.000	0.0000	1.4218
23.10	0.531	0.000	0.0000	1.4700
23.65	0.526	0.000	0.0000	1.5192
24.20	0.521	0.000	0.0000	1.5694
24.75	0.516	0.000	0.0000	1.6206
25.30	0.511	0.000	0.0000	1.6728
25.85	0.506	0.000	0.0000	1.7260
26.40	0.501	0.000	0.0000	1.7802
26.95	0.496	0.000	0.0000	1.8354
27.50	0.491	0.000	0.0000	1.8916
28.05	0.486	0.000	0.0000	1.9488
28.60	0.481	0.000	0.0000	2.0070
29.15	0.476	0.000	0.0000	2.0662
29.70	0.471	0.000	0.0000	2.1264
30.25	0.466	0.000	0.0000	2.1876
30.80	0.461	0.000	0.0000	2.2508
31.35	0.456	0.000	0.0000	2.3150
31.90	0.451	0.000	0.0000	2.3802
32.45	0.446	0.000	0.0000	2.4464
33.00	0.441	0.000	0.0000	2.5136
33.55	0.436	0.000	0.0000	2.5818
34.10	0.431	0.000	0.0000	2.6510
34.65	0.426	0.000	0.0000	2.7212
35.20	0.421	0.000	0.0000	2.7924
35.75	0.416	0.000	0.0000	2.8646
36.30	0.411	0.000	0.0000	2.9378
36.85	0.406	0.000	0.0000	3.0120
37.40	0.401	0.000	0.0000	3.0872
37.95	0.396	0.000	0.0000	3.1634
38.50	0.391	0.000	0.0000	3.2406
39.05	0.386	0.000	0.0000	3.3188
39.60	0.381	0.000	0.0000	3.3980
40.15	0.376	0.000	0.0000	3.4782
40.70	0.371	0.000	0.0000	3.5594
41.25	0.366	0.000	0.0000	3.6416
41.80	0.361	0.000	0.0000	3.7248
42.35	0.356	0.000	0.0000	3.8090
42.90	0.351	0.000	0.0000	3.8942
43.45	0.346	0.000	0.0000	3.9804
44.00	0.341	0.000	0.0000	4.0676
44.55	0.336	0.000	0.0000	4.1558
45.10	0.331	0.000	0.0000	4.2450
45.65	0.326	0.000	0.0000	4.3352
46.20	0.321	0.000	0.0000	4.4264
46.75	0.316	0.000	0.0000	4.5186
47.30	0.311	0.000	0.0000	4.6118
47.85	0.306	0.000	0.0000	4.7060
48.40	0.301	0.000	0.0000	4.8012
48.95	0.296	0.000	0.0000	4.8974
49.50	0.291	0.000	0.0000	4.9946
50.05	0.286	0.000	0.0000	5.0928
50.60	0.281	0.000	0.0000	5.1920
51.15	0.276	0.000	0.0000	5.2922
51.70	0.271	0.000	0.0000	5.3934
52.25	0.266	0.000	0.0000	5.4956
52.80	0.261	0.000	0.0000	5.5988
53.35	0.256	0.000	0.0000	5.7030
53.90	0.251	0.000	0.0000	5.8082
54.45	0.246	0.000	0.0000	5.9144
55.00	0.241	0.000	0.0000	6.0216
55.55	0.236	0.000	0.0000	6.1308
56.10	0.231	0.000	0.0000	6.2410
56.65	0.226	0.000	0.0000	6.3522
57.20	0.221	0.000	0.0000	6.4644
57.75	0.216	0.000	0.0000	6.5776
58.30	0.211	0.000	0.0000	6.6918
58.85	0.206	0.000	0.0000	6.8070
59.40	0.201	0.000	0.0000	6.9232
59.95	0.196	0.000	0.0000	7.0404
60.50	0.191	0.000	0.0000	7.1586
61.05	0.186	0.000	0.0000	7.2778
61.60	0.181	0.000	0.0000	7.3980
62.15	0.176	0.000	0.0000	7.5192
62.70	0.171	0.000	0.0000	7.6414
63.25	0.166	0.000	0.0000	7.7646
63.80	0.161	0.000	0.0000	7.8888
64.35	0.156	0.000	0.0000	8.0140
64.90	0.151	0.000	0.0000	8.1402
65.45	0.146	0.000	0.0000	8.2674
66.00	0.141	0.000	0.0000	8.3956
66.55	0.136	0.000	0.0000	8.5248
67.10	0.131	0.000	0.0000	8.6550
67.65	0.126	0.000	0.0000	8.7862
68.20	0.121	0.000	0.0000	8.9184
68.75	0.116	0.000	0.0000	9.0516
69.30	0.111	0.000	0.0000	9.1858
69.85	0.106	0.000	0.0000	9.3210
70.40	0.101	0.000	0.0000	9.4572
70.95	0.096	0.000	0.0000	9.5944
71.50	0.091	0.000	0.0000	9.7326
72.05	0.086	0.000	0.0000	9.8718
72.60	0.081	0.000	0.0000	10.0120
73.15	0.076	0.000	0.0000	10.1532
73.70	0.071	0.000	0.0000	10.2954
74.25	0.066	0.000	0.0000	10.4386
74.80	0.061	0.000	0.0000	10.5828
75.35	0.056	0.000	0.0000	10.7280
75.90	0.051	0.000	0.0000	10.8742
76.45	0.046	0.000	0.0000	11.0214
77.00	0.041	0.000	0.0000	11.1696
77.55	0.036	0.000	0.0000	11.3188
78.10	0.031	0.000	0.0000	11.4690
78.65	0.026	0.000	0.0000	11.6202
79.20	0.021	0.000	0.0000	11.7724
79.75	0.016	0.000	0.0000	11.9256
80.30	0.011	0.000	0.0000	12.0808
80.85	0.006	0.000	0.0000	12.2370
81.40	0.001	0.000	0.0000	12.3942
81.95	0.000	0.000	0.0000	12.5524
82.50	0.000	0.000	0.0000	12.7116
83.05	0.000	0.000	0.0000	12.8718
83.60	0.000	0.000	0.0000	13.0330
84.15	0.000	0.000	0.0000	13.1952
84.70	0.000	0.000	0.0000	13.3584
85.25	0.000	0.000	0.0000	13.5226
85.80	0.000	0.000	0.0000	13.6878
86.35	0.000	0.000	0.0000	13.8540
86.90	0.000	0.000	0.0000	14.0212
87.45	0.000	0.000	0.0000	14.1894
88.00	0.000	0.000	0.0000	14.3586
88.55	0.000	0.000	0.0000	14.5288
89.10	0.000	0.000	0.0000	14.6990
89.65	0.000	0.000	0.0000	14.8702
90.20	0.000	0.000	0.0000	15.0424
90.75	0.000	0.000	0.0000	15.2156
91.30	0.000	0.000	0.0000	15.3898
91.85	0.000	0.000	0.0000	15.5650
92.40	0.000	0.000	0.0000	15.7412
92.95	0.000	0.000	0.0000	15.9184
93.50	0.000	0.000	0.0000	16.0966
94.05	0.000	0.000	0.0000	16.2758
94.60	0.000	0.000	0.0000	16.4560
95.15	0.000	0.000	0.0000	16.6372
95.70	0.000	0.000	0.0000	16.8194
96.25	0.000	0.000	0.0000	17.0026
96.80	0.000	0.000	0.0000	17.1868
97.35	0.000	0.000	0.0000	17.3720
97.90	0.000	0.000	0.0000	17.5582
98.45	0.000	0.000	0.0000	17.7454
99.00	0.000	0.000	0.0000	17.9336
99.55	0.000	0.000	0.0000	18.1228
100.10	0.000	0.000	0.0000	18.3130
100.65	0.000	0.000	0.0000	18.5042
101.20	0.000	0.000	0.0000	18.6964
101.75	0.000	0.000	0.0000	18.8896
102.30	0.000	0.000	0.0000	19.0838
102.85	0.000	0.000	0.0000	19.2790
103.40	0.000	0.000	0.0000	19.4752
103.95	0.000	0.000	0.0000	19.6724
104.50	0.000	0.000	0.0000	19.8706
105.05	0.000	0.000	0.0000	20.0698
105.60	0.000	0.000	0.0000	20.2700
106.15	0.000	0.000	0.0000	20.4712
106.70	0.000	0.000	0.0000	20.6734
107.25	0.000	0.000	0.0000	20.8766
107.80	0.000	0.000	0.0000	21.0808
108.35	0.000	0.000	0.0000	21.2860
108.90	0.000	0.000	0.0000	21.4922
109.45	0.000	0.000	0.0000	21.6994
110.00	0.000	0.000	0.0000	21.9076
110.55	0.000	0.000	0.0000	22.1168
111.10	0.000	0.000	0.0000	22.3270
111.65	0.000	0.000	0.0000	22.5382
112.20	0.000	0.000	0.0000	22.7504
112.75	0.000	0.000	0.0000	22.9636
113.30	0.000	0.000	0.0000	23.1778
113.85	0.000	0.000	0.0000	23.3930
114.40	0.000	0.000	0.0000	23.6092
114.95	0.000	0.000	0.0000	23.8264
115.50	0.000	0.000	0.0000	24.0446
116.05				

CUMULATIVE LOG(FDI) = 0.0205E+00 LOG(V/D) = 7.3278E-02

C-12.02-17  
ION2

UNIT = 17.431 M/S  
RMS = 7.1025E-07  
T = 30.22 D 3.000  
M = 0.1057E+04 KG/MS  
X = 2.117  
MAY 04 09 53 KG/CM  
NU = 2.1187E+04 M/M/S

TIME	U (M/S)	V/T/A	U/V/E	Y/D/F/L
1.00	1.00	0.111	0.051	0.0114
2.00	2.00	0.0224	0.0254	0.0102
3.00	3.00	0.0355	0.0474	0.0303
4.00	4.00	0.0475	0.0715	0.0428
5.00	5.00	0.0591	0.0951	0.0400
6.00	6.00	0.0704	0.1177	0.0324
7.00	7.00	0.0814	0.1396	0.0200
8.00	8.00	0.0921	0.1609	0.0133
9.00	9.00	0.1025	0.1816	0.0077
10.00	10.00	0.1126	0.2017	0.0114
11.00	11.00	0.1224	0.2212	0.0170
12.00	12.00	0.1319	0.2402	0.0215
13.00	13.00	0.1411	0.2587	0.0275
14.00	14.00	0.1500	0.2767	0.0340
15.00	15.00	0.1586	0.2942	0.0411
16.00	16.00	0.1669	0.3112	0.0487
17.00	17.00	0.1749	0.3277	0.0568
18.00	18.00	0.1826	0.3437	0.0654
19.00	19.00	0.1900	0.3592	0.0745
20.00	20.00	0.1971	0.3743	0.0841
21.00	21.00	0.2039	0.3889	0.0942
22.00	22.00	0.2104	0.4031	0.1048
23.00	23.00	0.2166	0.4169	0.1159
24.00	24.00	0.2225	0.4303	0.1275
25.00	25.00	0.2281	0.4433	0.1396
26.00	26.00	0.2334	0.4559	0.1521
27.00	27.00	0.2384	0.4681	0.1651
28.00	28.00	0.2431	0.4799	0.1786
29.00	29.00	0.2475	0.4913	0.1925
30.00	30.00	0.2516	0.5023	0.2068
31.00	31.00	0.2554	0.5129	0.2215
32.00	32.00	0.2589	0.5231	0.2366
33.00	33.00	0.2621	0.5329	0.2521
34.00	34.00	0.2650	0.5423	0.2680
35.00	35.00	0.2676	0.5513	0.2843
36.00	36.00	0.2700	0.5600	0.3010
37.00	37.00	0.2721	0.5683	0.3181
38.00	38.00	0.2739	0.5762	0.3356
39.00	39.00	0.2754	0.5837	0.3534
40.00	40.00	0.2766	0.5908	0.3715
41.00	41.00	0.2775	0.5975	0.3900
42.00	42.00	0.2781	0.6038	0.4088
43.00	43.00	0.2784	0.6097	0.4280
44.00	44.00	0.2784	0.6152	0.4475
45.00	45.00	0.2781	0.6203	0.4673
46.00	46.00	0.2775	0.6250	0.4874
47.00	47.00	0.2766	0.6293	0.5078
48.00	48.00	0.2754	0.6332	0.5284
49.00	49.00	0.2739	0.6367	0.5492
50.00	50.00	0.2721	0.6398	0.5702
51.00	51.00	0.2700	0.6425	0.5914
52.00	52.00	0.2676	0.6448	0.6128
53.00	53.00	0.2650	0.6467	0.6344
54.00	54.00	0.2621	0.6482	0.6561
55.00	55.00	0.2589	0.6493	0.6780
56.00	56.00	0.2554	0.6499	0.7001
57.00	57.00	0.2516	0.6499	0.7224
58.00	58.00	0.2475	0.6493	0.7449
59.00	59.00	0.2431	0.6481	0.7675
60.00	60.00	0.2384	0.6463	0.7902
61.00	61.00	0.2334	0.6440	0.8130
62.00	62.00	0.2281	0.6413	0.8359
63.00	63.00	0.2225	0.6381	0.8589
64.00	64.00	0.2166	0.6344	0.8820
65.00	65.00	0.2104	0.6302	0.9052
66.00	66.00	0.2039	0.6255	0.9285
67.00	67.00	0.1971	0.6203	0.9519
68.00	68.00	0.1900	0.6146	0.9754
69.00	69.00	0.1826	0.6083	0.9990
70.00	70.00	0.1749	0.6015	1.0227
71.00	71.00	0.1669	0.5942	1.0465
72.00	72.00	0.1586	0.5864	1.0704
73.00	73.00	0.1500	0.5781	1.0944
74.00	74.00	0.1411	0.5693	1.1185
75.00	75.00	0.1319	0.5600	1.1427
76.00	76.00	0.1224	0.5503	1.1670
77.00	77.00	0.1126	0.5402	1.1914
78.00	78.00	0.1025	0.5297	1.2159
79.00	79.00	0.0921	0.5188	1.2405
80.00	80.00	0.0814	0.5075	1.2652
81.00	81.00	0.0704	0.4958	1.2900
82.00	82.00	0.0591	0.4837	1.3149
83.00	83.00	0.0475	0.4712	1.3400
84.00	84.00	0.0355	0.4583	1.3652
85.00	85.00	0.0224	0.4450	1.3905
86.00	86.00	0.0091	0.4313	1.4159
87.00	87.00	0.0000	0.4172	1.4414
88.00	88.00	0.0000	0.4027	1.4670
89.00	89.00	0.0000	0.3879	1.4927
90.00	90.00	0.0000	0.3728	1.5185
91.00	91.00	0.0000	0.3574	1.5444
92.00	92.00	0.0000	0.3417	1.5704
93.00	93.00	0.0000	0.3257	1.5965
94.00	94.00	0.0000	0.3094	1.6227
95.00	95.00	0.0000	0.2928	1.6490
96.00	96.00	0.0000	0.2759	1.6754
97.00	97.00	0.0000	0.2587	1.7019
98.00	98.00	0.0000	0.2413	1.7285
99.00	99.00	0.0000	0.2236	1.7552
100.00	100.00	0.0000	0.2056	1.7820

DEL = 42.68 MM DSTP = 5.37 MM THETA = 4.12 MM  
RDEL = 0.479E+05 RSTP = 0.1623E+04 RTTA = 0.3623E+04  
SHAPE FACTOR = 1.20

UNIT = 17.431 M/S  
RMS = 7.1025E-07  
T = 30.22 D 3.000  
M = 0.1057E+04 KG/MS  
X = 2.117  
MAY 04 09 53 KG/CM  
NU = 2.1187E+04 M/M/S

TIME	U (M/S)	V/T/A	U/V/E	Y/D/F/L
1.00	1.00	0.111	0.051	0.0114
2.00	2.00	0.0224	0.0254	0.0102
3.00	3.00	0.0355	0.0474	0.0303
4.00	4.00	0.0475	0.0715	0.0428
5.00	5.00	0.0591	0.0951	0.0400
6.00	6.00	0.0704	0.1177	0.0324
7.00	7.00	0.0814	0.1396	0.0200
8.00	8.00	0.0921	0.1609	0.0133
9.00	9.00	0.1025	0.1816	0.0077
10.00	10.00	0.1126	0.2017	0.0114
11.00	11.00	0.1224	0.2212	0.0170
12.00	12.00	0.1319	0.2402	0.0215
13.00	13.00	0.1411	0.2587	0.0275
14.00	14.00	0.1500	0.2767	0.0340
15.00	15.00	0.1586	0.2942	0.0411
16.00	16.00	0.1669	0.3112	0.0487
17.00	17.00	0.1749	0.3277	0.0568
18.00	18.00	0.1826	0.3437	0.0654
19.00	19.00	0.1900	0.3592	0.0745
20.00	20.00	0.1971	0.3743	0.0841
21.00	21.00	0.2039	0.3889	0.0942
22.00	22.00	0.2104	0.4031	0.1048
23.00	23.00	0.2166	0.4169	0.1159
24.00	24.00	0.2225	0.4303	0.1275
25.00	25.00	0.2281	0.4433	0.1396
26.00	26.00	0.2334	0.4559	0.1521
27.00	27.00	0.2384	0.4681	0.1651
28.00	28.00	0.2431	0.4799	0.1786
29.00	29.00	0.2475	0.4913	0.1925
30.00	30.00	0.2516	0.5023	0.2068
31.00	31.00	0.2554	0.5129	0.2215
32.00	32.00	0.2589	0.5231	0.2366
33.00	33.00	0.2621	0.5329	0.2521
34.00	34.00	0.2650	0.5423	0.2680
35.00	35.00	0.2676	0.5513	0.2843
36.00	36.00	0.2700	0.5600	0.3010
37.00	37.00	0.2721	0.5683	0.3181
38.00	38.00	0.2739	0.5762	0.3356
39.00	39.00	0.2754	0.5837	0.3534
40.00	40.00	0.2766	0.5908	0.3715
41.00	41.00	0.2775	0.5975	0.3900
42.00	42.00	0.2781	0.6038	0.4088
43.00	43.00	0.2784	0.6097	0.4280
44.00	44.00	0.2784	0.6152	0.4475
45.00	45.00	0.2781	0.6203	0.4673
46.00	46.00	0.2775	0.6250	0.4874
47.00	47.00	0.2766	0.6293	0.5078
48.00	48.00	0.2754	0.6332	0.5284
49.00	49.00	0.2739	0.6367	0.5492
50.00	50.00	0.2721	0.6398	0.5702
51.00	51.00	0.2700	0.6425	0.5914
52.00	52.00	0.2676	0.6448	0.6128
53.00	53.00	0.2650	0.6467	0.6344
54.00	54.00	0.2621	0.6482	0.6561
55.00	55.00	0.2589	0.6493	0.6780
56.00	56.00	0.2554	0.6499	0.7001
57.00	57.00	0.2516	0.6499	0.7224
58.00	58.00	0.2475	0.6493	0.7449
59.00	59.00	0.2431	0.6481	0.7675
60.00	60.00	0.2384	0.6463	0.7902
61.00	61.00	0.2334	0.6440	0.8130
62.00	62.00	0.2281	0.6413	0.8359
63.00	63.00	0.2225	0.6381	0.8589
64.00	64.00	0.2166	0.6344	0.8820
65.00	65.00	0.2104	0.6302	0.9052
66.00	66.00	0.2039	0.6255	0.9285
67.00	67.00	0.1971	0.6203	0.9519
68.00	68.00	0.1900	0.6146	0.9754
69.00	69.00	0.1826	0.6083	0.9990
70.00	70.00	0.1749	0.6015	1.0227
71.00	71.00	0.1669	0.5942	1.0465
72.00	72.00	0.1586	0.5864	1.0704
73.00	73.00	0.1500	0.5781	1.0944
74.00	74.00	0.1411	0.5693	1.1185
75.00	75.00	0.1319	0.5600	1.1427
76.00	76.00	0.1224	0.5503	1.1670
77.00	77.00	0.1126	0.5402	1.1914
78.00	78.00	0.1025	0.5297	1.2159
79.00	79.00	0.0921	0.5188	1.2405
80.00	80.00	0.0814	0.5075	1.2652
81.00	81.00	0.0704	0.4958	1.2900
82.00	82.00	0.0591	0.4837	1.3149
83.00	83.00	0.0475	0.4712	1.3400
84.00	84.00	0.0355	0.4583	1.3652
85.00	85.00	0.0224	0.4450	1.3905
86.00	86.00	0.0091	0.4313	1.4159
87.00	87.00	0.0000	0.4172	1.4414
88.00	88.00	0.0000	0.4027	1.4670
89.00	89.00	0.0000	0.3879	1.4927
90.00	90.00	0.0000	0.3728	1.5185
91.00	91.00	0.0000	0.3574	1.5444
92.00	92.00	0.0000	0.3417	1.5704
93.00	93.00	0.0000	0.3257	1.5965
94.00	94.00	0.0000	0.3094	1.6227
95.00	95.00	0.0000	0.2928	1.6490
96.00	96.00	0.0000	0.2759	1.6754
97.00	97.00	0.0000	0.2587	1.7019
98.00	98.00	0.0000	0.2413	1.7285
99.00	99.00	0.0000	0.2236	1.7552
100.00	100.00	0.0000	0.2056	1.7820

DEL = 51.32 MM DSTP = 5.93 MM THETA = 3.5



CURVE 1 FIT: LOG(FCT) = 5.1283E+05 \* LOG(V/D) + 3.1137E-02

UINF = 12.35J 4/S  
REX = 0.127E+07  
T = 26.0J DEGR  
NU = 0.1842E-04 KG/MS

X = 1.554  
RHO = 0.9688 KG/CM<sup>3</sup>  
NU = 0.1842E-04 M/M/S

C-539-13  
11M1

Y (MM)	U (M/S)	V/TA	U/VE	V/DEL
0.51	6.478	0.117	0.5142	0.0000
1.02	7.076	0.234	0.0123	0.0000
1.53	8.000	0.351	0.0718	0.0000
2.04	9.000	0.468	0.0755	0.0000
2.55	10.000	0.585	0.0792	0.0000
3.06	11.000	0.702	0.0829	0.0000
3.57	12.000	0.819	0.0866	0.0000
4.08	13.000	0.936	0.0903	0.0000
4.59	14.000	1.053	0.0940	0.0000
5.10	15.000	1.170	0.0977	0.0000
5.61	16.000	1.287	0.1014	0.0000
6.12	17.000	1.404	0.1051	0.0000
6.63	18.000	1.521	0.1088	0.0000
7.14	19.000	1.638	0.1125	0.0000
7.65	20.000	1.755	0.1162	0.0000
8.16	21.000	1.872	0.1199	0.0000
8.67	22.000	1.989	0.1236	0.0000
9.18	23.000	2.106	0.1273	0.0000
9.69	24.000	2.223	0.1310	0.0000
10.20	25.000	2.340	0.1347	0.0000
10.71	26.000	2.457	0.1384	0.0000
11.22	27.000	2.574	0.1421	0.0000
11.73	28.000	2.691	0.1458	0.0000
12.24	29.000	2.808	0.1495	0.0000
12.75	30.000	2.925	0.1532	0.0000
13.26	31.000	3.042	0.1569	0.0000
13.77	32.000	3.159	0.1606	0.0000
14.28	33.000	3.276	0.1643	0.0000
14.79	34.000	3.393	0.1680	0.0000
15.30	35.000	3.510	0.1717	0.0000
15.81	36.000	3.627	0.1754	0.0000
16.32	37.000	3.744	0.1791	0.0000
16.83	38.000	3.861	0.1828	0.0000
17.34	39.000	3.978	0.1865	0.0000
17.85	40.000	4.095	0.1902	0.0000
18.36	41.000	4.212	0.1939	0.0000
18.87	42.000	4.329	0.1976	0.0000
19.38	43.000	4.446	0.2013	0.0000
19.89	44.000	4.563	0.2050	0.0000
20.40	45.000	4.680	0.2087	0.0000
20.91	46.000	4.797	0.2124	0.0000
21.42	47.000	4.914	0.2161	0.0000
21.93	48.000	5.031	0.2198	0.0000
22.44	49.000	5.148	0.2235	0.0000
22.95	50.000	5.265	0.2272	0.0000
23.46	51.000	5.382	0.2309	0.0000
23.97	52.000	5.499	0.2346	0.0000
24.48	53.000	5.616	0.2383	0.0000
24.99	54.000	5.733	0.2420	0.0000
25.50	55.000	5.850	0.2457	0.0000
26.01	56.000	5.967	0.2494	0.0000
26.52	57.000	6.084	0.2531	0.0000
27.03	58.000	6.201	0.2568	0.0000
27.54	59.000	6.318	0.2605	0.0000
28.05	60.000	6.435	0.2642	0.0000
28.56	61.000	6.552	0.2679	0.0000
29.07	62.000	6.669	0.2716	0.0000
29.58	63.000	6.786	0.2753	0.0000
30.09	64.000	6.903	0.2790	0.0000
30.60	65.000	7.020	0.2827	0.0000
31.11	66.000	7.137	0.2864	0.0000
31.62	67.000	7.254	0.2901	0.0000
32.13	68.000	7.371	0.2938	0.0000
32.64	69.000	7.488	0.2975	0.0000
33.15	70.000	7.605	0.3012	0.0000
33.66	71.000	7.722	0.3049	0.0000
34.17	72.000	7.839	0.3086	0.0000
34.68	73.000	7.956	0.3123	0.0000
35.19	74.000	8.073	0.3160	0.0000
35.70	75.000	8.190	0.3197	0.0000
36.21	76.000	8.307	0.3234	0.0000
36.72	77.000	8.424	0.3271	0.0000
37.23	78.000	8.541	0.3308	0.0000
37.74	79.000	8.658	0.3345	0.0000
38.25	80.000	8.775	0.3382	0.0000
38.76	81.000	8.892	0.3419	0.0000
39.27	82.000	9.009	0.3456	0.0000
39.78	83.000	9.126	0.3493	0.0000
40.29	84.000	9.243	0.3530	0.0000
40.80	85.000	9.360	0.3567	0.0000
41.31	86.000	9.477	0.3604	0.0000
41.82	87.000	9.594	0.3641	0.0000
42.33	88.000	9.711	0.3678	0.0000
42.84	89.000	9.828	0.3715	0.0000
43.35	90.000	9.945	0.3752	0.0000
43.86	91.000	10.062	0.3789	0.0000
44.37	92.000	10.179	0.3826	0.0000
44.88	93.000	10.296	0.3863	0.0000
45.39	94.000	10.413	0.3900	0.0000
45.90	95.000	10.530	0.3937	0.0000
46.41	96.000	10.647	0.3974	0.0000
46.92	97.000	10.764	0.4011	0.0000
47.43	98.000	10.881	0.4048	0.0000
47.94	99.000	11.000	0.4085	0.0000
48.45	100.000	11.117	0.4122	0.0000

DEL = 45.91 MM DSTP = 8.17 MM THETA = 4.48 MM  
RDEL = 0.2870E+05 ROSTR = 0.2801E+04 RTA = 0.2971E+04  
SHAPE FACTOR = 1.29

UINF = 12.35J 4/S  
REX = 0.1313E+07  
T = 26.33 DEGR  
NU = 0.1842E-04 KG/MS

X = 2.230  
RHO = 0.9687 KG/CM<sup>3</sup>  
NU = 0.1842E-04 M/M/S

Y (MM)	U (M/S)	V/TA	U/VE	V/DEL
0.51	6.478	0.117	0.3370	0.0117
1.02	7.076	0.234	0.0123	0.0117
1.53	8.000	0.351	0.0718	0.0117
2.04	9.000	0.468	0.0755	0.0117
2.55	10.000	0.585	0.0792	0.0117
3.06	11.000	0.702	0.0829	0.0117
3.57	12.000	0.819	0.0866	0.0117
4.08	13.000	0.936	0.0903	0.0117
4.59	14.000	1.053	0.0940	0.0117
5.10	15.000	1.170	0.0977	0.0117
5.61	16.000	1.287	0.1014	0.0117
6.12	17.000	1.404	0.1051	0.0117
6.63	18.000	1.521	0.1088	0.0117
7.14	19.000	1.638	0.1125	0.0117
7.65	20.000	1.755	0.1162	0.0117
8.16	21.000	1.872	0.1199	0.0117
8.67	22.000	1.989	0.1236	0.0117
9.18	23.000	2.106	0.1273	0.0117
9.69	24.000	2.223	0.1310	0.0117
10.20	25.000	2.340	0.1347	0.0117
10.71	26.000	2.457	0.1384	0.0117
11.22	27.000	2.574	0.1421	0.0117
11.73	28.000	2.691	0.1458	0.0117
12.24	29.000	2.808	0.1495	0.0117
12.75	30.000	2.925	0.1532	0.0117
13.26	31.000	3.042	0.1569	0.0117
13.77	32.000	3.159	0.1606	0.0117
14.28	33.000	3.276	0.1643	0.0117
14.79	34.000	3.393	0.1680	0.0117
15.30	35.000	3.510	0.1717	0.0117
15.81	36.000	3.627	0.1754	0.0117
16.32	37.000	3.744	0.1791	0.0117
16.83	38.000	3.861	0.1828	0.0117
17.34	39.000	3.978	0.1865	0.0117
17.85	40.000	4.095	0.1902	0.0117
18.36	41.000	4.212	0.1939	0.0117
18.87	42.000	4.329	0.1976	0.0117
19.38	43.000	4.446	0.2013	0.0117
19.89	44.000	4.563	0.2050	0.0117
20.40	45.000	4.680	0.2087	0.0117
20.91	46.000	4.797	0.2124	0.0117
21.42	47.000	4.914	0.2161	0.0117
21.93	48.000	5.031	0.2198	0.0117
22.44	49.000	5.148	0.2235	0.0117
22.95	50.000	5.265	0.2272	0.0117
23.46	51.000	5.382	0.2309	0.0117
23.97	52.000	5.499	0.2346	0.0117
24.48	53.000	5.616	0.2383	0.0117
24.99	54.000	5.733	0.2420	0.0117
25.50	55.000	5.850	0.2457	0.0117
26.01	56.000	5.967	0.2494	0.0117
26.52	57.000	6.084	0.2531	0.0117
27.03	58.000	6.201	0.2568	0.0117
27.54	59.000	6.318	0.2605	0.0117
28.05	60.000	6.435	0.2642	0.0117
28.56	61.000	6.552	0.2679	0.0117
29.07	62.000	6.669	0.2716	0.0117
29.58	63.000	6.786	0.2753	0.0117
30.09	64.000	6.903	0.2790	0.0117
30.60	65.000	7.020	0.2827	0.0117
31.11	66.000	7.137	0.2864	0.0117
31.62	67.000	7.254	0.2901	0.0117
32.13	68.000	7.371	0.2938	0.0117
32.64	69.000	7.488	0.2975	0.0117
33.15	70.000	7.605	0.3012	0.0117
33.66	71.000	7.722	0.3049	0.0117
34.17	72.000	7.839	0.3086	0.0117
34.68	73.000	7.956	0.3123	0.0117
35.19	74.000	8.073	0.3160	0.0117
35.70	75.000	8.190	0.3197	0.0117
36.21	76.000	8.307	0.3234	0.0117
36.72	77.000	8.424	0.3271	0.0117
37.23	78.000	8.541	0.3308	0.0117
37.74	79.000	8.658	0.3345	0.0117
38.25	80.000	8.775	0.3382	0.0117
38.76	81.000	8.892	0.3419	0.0117
39.27	82.000	9.009	0.3456	0.0117
39.78	83.000	9.126	0.3493	0.0117
40.29	84.000	9.243	0.3530	0.0117
40.80	85.000	9.360	0.3567	0.0117
41.31	86.000	9.477	0.3604	0.0117
41.82	87.000	9.594	0.3641	0.0117
42.33	88.000	9.711	0.3678	0.0117
42.84	89.000	9.828	0.3715	0.0117
43.35	90.000	9.945	0.3752	0.0117
43.86	91.000	10.062	0.3789	0.0117
44.37	92.000	10.179	0.3826	0.0117
44.88	93.000	10.296	0.3863	0.0117
45.39	94.000	10.413	0.3900	0.0117
45.90	95.000	10.530	0.3937	0.0117
46.41	96.000	10.647	0.3974	0.0117
46.92	97.000	10.764	0.4011	0.0117
47.43	98.000	10.881	0.4048	0.0117
47.94	99.000	11.000	0.4085	0.0117
48.45	100.000	11.117	0.4122	0.0117

DEL = 45.91 MM DSTP = 8.17 MM THETA = 4.48 MM  
RDEL = 0.2870E+05 ROSTR = 0.2801E+04 RTA = 0.2971E+04  
SHAPE FACTOR = 1.29

UINF = 12.35J 4/S  
REX = 0.1373E+07  
T = 26.70 DEGR  
NU = 0.1842E-04 KG/MS

X = 2.119  
RHO = 0.9675 KG/CM<sup>3</sup>  
NU = 0.1907E-04 M/M/S

Y (MM)	U (M/S)	V/TA	U/VE	V/DEL
0.51	6.478	0.117	0.2910	0.0000
1.02	7.076	0.234	0.0123	0.0000
1.53	8.000	0.351	0.0718	0.000

C-12.22-13

11N2

STATION 7

DATE 12-22-13 TIME 14:00

WIND 10-15 KTS DIRECTION 090

SEA 1-2 FT WAVE PERIOD 8-10 SEC

TEMP AIR 50 F SURFACE 48 F

WIND 10-15 KTS DIRECTION 090

WIND 10-15 KTS DIRECTION 090

SEA 1-2 FT WAVE PERIOD 8-10 SEC

DATE 12-22-13 TIME 14:00

WIND 10-15 KTS DIRECTION 090

SEA 1-2 FT WAVE PERIOD 8-10 SEC

TEMP AIR 50 F SURFACE 48 F

WIND 10-15 KTS DIRECTION 090

WIND 10-15 KTS DIRECTION 090

SEA 1-2 FT WAVE PERIOD 8-10 SEC







E-120-17

LOG LOG(=) 2.1264E+00 LOG(Y/D) = 0.135NE-01

UINF= 15.410 V/S  
RHO= 0.9733 KG/CM  
NU= 0.1889E-04 MM/S

X= 2.339  
RHO= 0.9733 KG/CM  
NU= 0.1889E-04 MM/S

Y/DEL  
U (M/S)  
Y/DEL  
U (M/S)

22.62	19.181	3.098	0.8091	0.4442
27.62	15.160	4.802	0.8949	0.5424
33.62	11.160	6.041	0.9272	0.6403
39.62	7.198	7.031	0.9524	0.7781
45.62	3.192	8.002	0.9714	0.9566
51.62	1.192	9.002	0.9854	1.1751
57.62	0.192	10.002	0.9944	1.4336
63.62	0.092	11.002	0.9984	1.7321
69.62	0.042	12.002	0.9994	2.0706
75.62	0.022	13.002	0.9998	2.4491
81.62	0.012	14.002	0.9999	2.8676
87.62	0.007	15.002	0.9999	3.3261
93.62	0.004	16.002	0.9999	3.8246
99.62	0.002	17.002	0.9999	4.3631
105.62	0.001	18.002	0.9999	4.9416
111.62	0.000	19.002	0.9999	5.5601
117.62	0.000	20.002	0.9999	6.2186
123.62	0.000	21.002	0.9999	6.9171
129.62	0.000	22.002	0.9999	7.6456
135.62	0.000	23.002	0.9999	8.4041
141.62	0.000	24.002	0.9999	9.1926
147.62	0.000	25.002	0.9999	10.0111
153.62	0.000	26.002	0.9999	10.8596
159.62	0.000	27.002	0.9999	11.7381
165.62	0.000	28.002	0.9999	12.6466
171.62	0.000	29.002	0.9999	13.5851
177.62	0.000	30.002	0.9999	14.5536
183.62	0.000	31.002	0.9999	15.5521
189.62	0.000	32.002	0.9999	16.5806
195.62	0.000	33.002	0.9999	17.6391
201.62	0.000	34.002	0.9999	18.7276
207.62	0.000	35.002	0.9999	19.8461
213.62	0.000	36.002	0.9999	20.9946
219.62	0.000	37.002	0.9999	22.1731
225.62	0.000	38.002	0.9999	23.3816
231.62	0.000	39.002	0.9999	24.6201
237.62	0.000	40.002	0.9999	25.8886
243.62	0.000	41.002	0.9999	27.1871
249.62	0.000	42.002	0.9999	28.5156
255.62	0.000	43.002	0.9999	29.8741
261.62	0.000	44.002	0.9999	31.2626
267.62	0.000	45.002	0.9999	32.6811
273.62	0.000	46.002	0.9999	34.1296
279.62	0.000	47.002	0.9999	35.6081
285.62	0.000	48.002	0.9999	37.1166
291.62	0.000	49.002	0.9999	38.6551
297.62	0.000	50.002	0.9999	40.2236
303.62	0.000	51.002	0.9999	41.8221
309.62	0.000	52.002	0.9999	43.4506
315.62	0.000	53.002	0.9999	45.1091
321.62	0.000	54.002	0.9999	46.7976
327.62	0.000	55.002	0.9999	48.5161
333.62	0.000	56.002	0.9999	50.2646
339.62	0.000	57.002	0.9999	52.0431
345.62	0.000	58.002	0.9999	53.8516
351.62	0.000	59.002	0.9999	55.6901
357.62	0.000	60.002	0.9999	57.5586
363.62	0.000	61.002	0.9999	59.4571
369.62	0.000	62.002	0.9999	61.3856
375.62	0.000	63.002	0.9999	63.3441
381.62	0.000	64.002	0.9999	65.3326
387.62	0.000	65.002	0.9999	67.3511
393.62	0.000	66.002	0.9999	69.3996
399.62	0.000	67.002	0.9999	71.4781
405.62	0.000	68.002	0.9999	73.5866
411.62	0.000	69.002	0.9999	75.7251
417.62	0.000	70.002	0.9999	77.8936
423.62	0.000	71.002	0.9999	80.0921
429.62	0.000	72.002	0.9999	82.3206
435.62	0.000	73.002	0.9999	84.5791
441.62	0.000	74.002	0.9999	86.8676
447.62	0.000	75.002	0.9999	89.1861
453.62	0.000	76.002	0.9999	91.5346
459.62	0.000	77.002	0.9999	93.9131
465.62	0.000	78.002	0.9999	96.3216
471.62	0.000	79.002	0.9999	98.7601
477.62	0.000	80.002	0.9999	101.2286
483.62	0.000	81.002	0.9999	103.7271
489.62	0.000	82.002	0.9999	106.2556
495.62	0.000	83.002	0.9999	108.8141
501.62	0.000	84.002	0.9999	111.4026
507.62	0.000	85.002	0.9999	114.0211
513.62	0.000	86.002	0.9999	116.6696
519.62	0.000	87.002	0.9999	119.3481
525.62	0.000	88.002	0.9999	122.0566
531.62	0.000	89.002	0.9999	124.7951
537.62	0.000	90.002	0.9999	127.5636
543.62	0.000	91.002	0.9999	130.3621
549.62	0.000	92.002	0.9999	133.1906
555.62	0.000	93.002	0.9999	136.0491
561.62	0.000	94.002	0.9999	138.9376
567.62	0.000	95.002	0.9999	141.8561
573.62	0.000	96.002	0.9999	144.8046
579.62	0.000	97.002	0.9999	147.7831
585.62	0.000	98.002	0.9999	150.7916
591.62	0.000	99.002	0.9999	153.8291
597.62	0.000	100.002	0.9999	156.8956
603.62	0.000	101.002	0.9999	159.9911
609.62	0.000	102.002	0.9999	163.1156
615.62	0.000	103.002	0.9999	166.2691
621.62	0.000	104.002	0.9999	169.4516
627.62	0.000	105.002	0.9999	172.6631
633.62	0.000	106.002	0.9999	175.9046
639.62	0.000	107.002	0.9999	179.1761
645.62	0.000	108.002	0.9999	182.4776
651.62	0.000	109.002	0.9999	185.8091
657.62	0.000	110.002	0.9999	189.1706
663.62	0.000	111.002	0.9999	192.5621
669.62	0.000	112.002	0.9999	195.9836
675.62	0.000	113.002	0.9999	199.4351
681.62	0.000	114.002	0.9999	202.9166
687.62	0.000	115.002	0.9999	206.4281
693.62	0.000	116.002	0.9999	210.9696
699.62	0.000	117.002	0.9999	215.5411
705.62	0.000	118.002	0.9999	220.1426
711.62	0.000	119.002	0.9999	224.7741
717.62	0.000	120.002	0.9999	229.4356
723.62	0.000	121.002	0.9999	234.1271
729.62	0.000	122.002	0.9999	238.8486
735.62	0.000	123.002	0.9999	243.5901
741.62	0.000	124.002	0.9999	248.3516
747.62	0.000	125.002	0.9999	253.1331
753.62	0.000	126.002	0.9999	257.9346
759.62	0.000	127.002	0.9999	262.7561
765.62	0.000	128.002	0.9999	267.5976
771.62	0.000	129.002	0.9999	272.4591
777.62	0.000	130.002	0.9999	277.3406
783.62	0.000	131.002	0.9999	282.2421
789.62	0.000	132.002	0.9999	287.1636
795.62	0.000	133.002	0.9999	292.1051
801.62	0.000	134.002	0.9999	297.0666
807.62	0.000	135.002	0.9999	302.0481
813.62	0.000	136.002	0.9999	307.0496
819.62	0.000	137.002	0.9999	312.0711
825.62	0.000	138.002	0.9999	317.1126
831.62	0.000	139.002	0.9999	322.1741
837.62	0.000	140.002	0.9999	327.2556
843.62	0.000	141.002	0.9999	332.3571
849.62	0.000	142.002	0.9999	337.4786
855.62	0.000	143.002	0.9999	342.6201
861.62	0.000	144.002	0.9999	347.7816
867.62	0.000	145.002	0.9999	352.9631
873.62	0.000	146.002	0.9999	358.1646
879.62	0.000	147.002	0.9999	363.3861
885.62	0.000	148.002	0.9999	368.6276
891.62	0.000	149.002	0.9999	373.8891
897.62	0.000	150.002	0.9999	379.1706
903.62	0.000	151.002	0.9999	384.4721
909.62	0.000	152.002	0.9999	389.7936
915.62	0.000	153.002	0.9999	395.1351
921.62	0.000	154.002	0.9999	400.4966
927.62	0.000	155.002	0.9999	405.8781
933.62	0.000	156.002	0.9999	411.2796
939.62	0.000	157.002	0.9999	416.7011
945.62	0.000	158.002	0.9999	422.1426
951.62	0.000	159.002	0.9999	427.6041
957.62	0.000	160.002	0.9999	433.0856
963.62	0.000	161.002	0.9999	438.5871
969.62	0.000	162.002	0.9999	444.1086
975.62	0.000	163.002	0.9999	449.6501
981.62	0.000	164.002	0.9999	455.2116
987.62	0.000	165.002	0.9999	460.7931
993.62	0.000	166.002	0.9999	466.3946
999.62	0.000	167.002	0.9999	472.0161
1005.62	0.000	168.002	0.9999	477.6576
1011.62	0.000	169.002	0.9999	483.3191
1017.62	0.000	170.002	0.9999	488.9906
1023.62	0.000	171.002	0.9999	494.6721
1029.62	0.000	172.002	0.9999	500.3736
1035.62	0.000	173.002	0.9999	506.0951
1041.62	0.000	174.002	0.9999	511.8366
1047.62	0.000	175.002	0.9999	517.5981
1053.62	0.000	176.002	0.9999	523.3796
1059.62	0.000	177.002	0.9999	529.1811
1065.62	0.000	178.002	0.9999	535.0026
1071.62	0.000	179.002	0.9999	540.8441
1077.62	0.000	180.002	0.9999	546.7056
1083.62	0.000	181.002	0.9999	552.5871
1089.62	0.000	182.002	0.9999	558.4886
1095.62	0.000	183.002	0.9999	564.4101
1101.62	0.000	184.002	0.9999	570.3516
1107.62	0.000	185.002	0.9999	576.3131
1113.62	0.000	186.002	0.9999	582.2946
1119.62	0.000	187.002	0.9999	588.2961
1125.62	0.000	188.002	0.9999	594.3176
1131.62	0.000	189.002	0.9999	600.3591
1137.62	0.000	190.002	0.9999	606.4206
1143.62	0.000	191.002	0.9999	612.4921
1149.62	0.000	192.002	0.9999	618.5836
1155.62	0.000	193.002	0.9999	624.6951
1161.62	0.000	194.002	0.9999	630.8266
1167.62	0.000	195.002	0.9999	636.9781
1173.62	0.000	196.002	0.9999	643.1496
1179.62	0.000	197.002	0.9999	649.3411
1185.62	0.000	198.002	0.9999	655.5526
1191.62	0.000	199.002	0.9999	661.7841
1197.62	0.000	200.002	0.9999	668.0356

E-14.51-17

24NT

50% Blocked

CUMV I FIT: L00(FD) = 0.1135E+00 \* L05(Y/D) + -0.3739E-02

UINF = 15.578 M/S  
REX = 0.1865E+07  
T = 26.78 DEG.C  
MU = 0.1865E-04 KG/MS

X = 2.062  
RHO = 0.9578 KG/CU.M  
NU = 0.1909E-04 MM/S

Y(MM)	U (M/S)	V(TTA)	U(U/E)	Y(DEL)
0.52	0.006	0.110	0.5712	0.01306
1.04	0.012	0.197	0.6247	0.02612
1.56	0.018	0.312	0.6693	0.03918
2.08	0.024	0.427	0.6972	0.05224
2.60	0.030	0.542	0.7125	0.06530
3.12	0.036	0.657	0.7115	0.07836
3.64	0.042	0.772	0.7462	0.09142
4.16	0.048	0.887	0.7674	0.10448
4.68	0.054	1.002	0.7823	0.11754
5.20	0.060	1.117	0.7924	0.13060
5.72	0.066	1.232	0.8025	0.14366
6.24	0.072	1.347	0.8126	0.15672
6.76	0.078	1.462	0.8227	0.16978
7.28	0.084	1.577	0.8328	0.18284
7.80	0.090	1.692	0.8429	0.19590
8.32	0.096	1.807	0.8530	0.20896
8.84	0.102	1.922	0.8631	0.22202
9.36	0.108	2.037	0.8732	0.23508
9.88	0.114	2.152	0.8833	0.24814
10.40	0.120	2.267	0.8934	0.26120
10.92	0.126	2.382	0.9035	0.27426
11.44	0.132	2.497	0.9136	0.28732
11.96	0.138	2.612	0.9237	0.30038
12.48	0.144	2.727	0.9338	0.31344
13.00	0.150	2.842	0.9439	0.32650
13.52	0.156	2.957	0.9540	0.33956
14.04	0.162	3.072	0.9641	0.35262
14.56	0.168	3.187	0.9742	0.36568
15.08	0.174	3.302	0.9843	0.37874
15.60	0.180	3.417	0.9944	0.39180
16.12	0.186	3.532	0.9995	0.40486
16.64	0.192	3.647	0.9996	0.41792
17.16	0.198	3.762	0.9997	0.43098
17.68	0.204	3.877	0.9998	0.44404
18.20	0.210	3.992	0.9999	0.45710
18.72	0.216	4.107	0.9999	0.47016
19.24	0.222	4.222	0.9999	0.48322
19.76	0.228	4.337	0.9999	0.49628
20.28	0.234	4.452	0.9999	0.50934
20.80	0.240	4.567	0.9999	0.52240
21.32	0.246	4.682	0.9999	0.53546
21.84	0.252	4.797	0.9999	0.54852
22.36	0.258	4.912	0.9999	0.56158
22.88	0.264	5.027	0.9999	0.57464
23.40	0.270	5.142	0.9999	0.58770
23.92	0.276	5.257	0.9999	0.60076
24.44	0.282	5.372	0.9999	0.61382
24.96	0.288	5.487	0.9999	0.62688
25.48	0.294	5.602	0.9999	0.63994
26.00	0.300	5.717	0.9999	0.65300
26.52	0.306	5.832	0.9999	0.66606
27.04	0.312	5.947	0.9999	0.67912
27.56	0.318	6.062	0.9999	0.69218
28.08	0.324	6.177	0.9999	0.70524
28.60	0.330	6.292	0.9999	0.71830
29.12	0.336	6.407	0.9999	0.73136
29.64	0.342	6.522	0.9999	0.74442
30.16	0.348	6.637	0.9999	0.75748
30.68	0.354	6.752	0.9999	0.77054
31.20	0.360	6.867	0.9999	0.78360
31.72	0.366	6.982	0.9999	0.79666
32.24	0.372	7.097	0.9999	0.80972
32.76	0.378	7.212	0.9999	0.82278
33.28	0.384	7.327	0.9999	0.83584
33.80	0.390	7.442	0.9999	0.84890
34.32	0.396	7.557	0.9999	0.86196
34.84	0.402	7.672	0.9999	0.87502
35.36	0.408	7.787	0.9999	0.88808
35.88	0.414	7.902	0.9999	0.90114
36.40	0.420	8.017	0.9999	0.91420
36.92	0.426	8.132	0.9999	0.92726
37.44	0.432	8.247	0.9999	0.94032
37.96	0.438	8.362	0.9999	0.95338
38.48	0.444	8.477	0.9999	0.96644
39.00	0.450	8.592	0.9999	0.97950
39.52	0.456	8.707	0.9999	0.99256
40.04	0.462	8.822	0.9999	1.00562
40.56	0.468	8.937	0.9999	1.01868
41.08	0.474	9.052	0.9999	1.03174
41.60	0.480	9.167	0.9999	1.04480
42.12	0.486	9.282	0.9999	1.05786
42.64	0.492	9.397	0.9999	1.07092
43.16	0.498	9.512	0.9999	1.08398
43.68	0.504	9.627	0.9999	1.09704
44.20	0.510	9.742	0.9999	1.11010
44.72	0.516	9.857	0.9999	1.12316
45.24	0.522	9.972	0.9999	1.13622
45.76	0.528	10.087	0.9999	1.14928
46.28	0.534	10.202	0.9999	1.16234
46.80	0.540	10.317	0.9999	1.17540
47.32	0.546	10.432	0.9999	1.18846
47.84	0.552	10.547	0.9999	1.20152
48.36	0.558	10.662	0.9999	1.21458
48.88	0.564	10.777	0.9999	1.22764
49.40	0.570	10.892	0.9999	1.24070
49.92	0.576	11.007	0.9999	1.25376
50.44	0.582	11.122	0.9999	1.26682
50.96	0.588	11.237	0.9999	1.27988
51.48	0.594	11.352	0.9999	1.29294
52.00	0.600	11.467	0.9999	1.30600
52.52	0.606	11.582	0.9999	1.31906
53.04	0.612	11.697	0.9999	1.33212
53.56	0.618	11.812	0.9999	1.34518
54.08	0.624	11.927	0.9999	1.35824
54.60	0.630	12.042	0.9999	1.37130
55.12	0.636	12.157	0.9999	1.38436
55.64	0.642	12.272	0.9999	1.39742
56.16	0.648	12.387	0.9999	1.41048
56.68	0.654	12.502	0.9999	1.42354
57.20	0.660	12.617	0.9999	1.43660
57.72	0.666	12.732	0.9999	1.44966
58.24	0.672	12.847	0.9999	1.46272
58.76	0.678	12.962	0.9999	1.47578
59.28	0.684	13.077	0.9999	1.48884
59.80	0.690	13.192	0.9999	1.50190
60.32	0.696	13.307	0.9999	1.51496
60.84	0.702	13.422	0.9999	1.52802
61.36	0.708	13.537	0.9999	1.54108
61.88	0.714	13.652	0.9999	1.55414
62.40	0.720	13.767	0.9999	1.56720
62.92	0.726	13.882	0.9999	1.58026
63.44	0.732	13.997	0.9999	1.59332
63.96	0.738	14.112	0.9999	1.60638
64.48	0.744	14.227	0.9999	1.61944
65.00	0.750	14.342	0.9999	1.63250
65.52	0.756	14.457	0.9999	1.64556
66.04	0.762	14.572	0.9999	1.65862
66.56	0.768	14.687	0.9999	1.67168
67.08	0.774	14.802	0.9999	1.68474
67.60	0.780	14.917	0.9999	1.69780
68.12	0.786	15.032	0.9999	1.71086
68.64	0.792	15.147	0.9999	1.72392
69.16	0.798	15.262	0.9999	1.73698
69.68	0.804	15.377	0.9999	1.75004
70.20	0.810	15.492	0.9999	1.76310
70.72	0.816	15.607	0.9999	1.77616
71.24	0.822	15.722	0.9999	1.78922
71.76	0.828	15.837	0.9999	1.80228
72.28	0.834	15.952	0.9999	1.81534
72.80	0.840	16.067	0.9999	1.82840
73.32	0.846	16.182	0.9999	1.84146
73.84	0.852	16.297	0.9999	1.85452
74.36	0.858	16.412	0.9999	1.86758
74.88	0.864	16.527	0.9999	1.88064
75.40	0.870	16.642	0.9999	1.89370
75.92	0.876	16.757	0.9999	1.90676
76.44	0.882	16.872	0.9999	1.91982
76.96	0.888	16.987	0.9999	1.93288
77.48	0.894	17.102	0.9999	1.94594
78.00	0.900	17.217	0.9999	1.95900
78.52	0.906	17.332	0.9999	1.97206
79.04	0.912	17.447	0.9999	1.98512
79.56	0.918	17.562	0.9999	1.99818
80.08	0.924	17.677	0.9999	2.01124
80.60	0.930	17.792	0.9999	2.02430
81.12	0.936	17.907	0.9999	2.03736
81.64	0.942	18.022	0.9999	2.05042
82.16	0.948	18.137	0.9999	2.06348
82.68	0.954	18.252	0.9999	2.07654
83.20	0.960	18.367	0.9999	2.08960
83.72	0.966	18.482	0.9999	2.10266
84.24	0.972	18.597	0.9999	2.11572
84.76	0.978	18.712	0.9999	2.12878
85.28	0.984	18.827	0.9999	2.14184
85.80	0.990	18.942	0.9999	2.15490
86.32	0.996	19.057	0.9999	2.16796
86.84	1.002	19.172	0.9999	2.18102
87.36	1.008	19.287	0.9999	2.19408
87.88	1.014	19.402	0.9999	2.20714
88.40	1.020	19.517	0.9999	2.22020
88.92	1.026	19.632	0.9999	2.23326
89.44	1.032	19.747	0.9999	2.24632
89.96	1.038	19.862	0.9999	2.25938
90.48	1.044	19.977	0.9999	2.27244
91.00	1.050	20.092	0.9999	2.28550
91.52	1.056	20.207	0.9999	2.29856
92.04	1.062	20.322	0.9999	2.31162
92.56	1.068	20.437	0.9999	2.32468
93.08	1.074	20.552	0.9999	2.33774
93.60	1.080	20.667	0.9999	2.35080
94.12	1.086	20.782	0.9999	2.36386
94.64	1.092	20.897	0.9999	2.37692
95.16	1.098	21.012	0.9999	2.39000
95.68	1.104	21.127	0.9999	2.40306
96.20	1.110	21.242	0.9999	2.41612
96.72	1.116	21.357	0.9999	2.42918
97.24	1.122	21.472	0.9999	2.44224
97.76	1.128	21.587	0.9999	2.45530
98.28	1.134	21.702	0.9999	2.46836
98.80	1.140	21.817	0.9999	2.48142
99.32	1.146	21.932	0.9999	2.49448
99.84	1.152	22.047	0.9999	2.50754
100.36	1.158	22.162	0.9999	2.52060
100.88	1.164	22.277	0.9999	2.53366
101.40	1.170	22.392	0.9999	2.54672
101.92	1.176	22.507	0.9999	2.55978
102.44	1.182	22.622	0.9999	2.57284
102.96	1.188	22.737	0.9999	2.58590
103.48	1.194	22.852	0.9999	2.59896
104.00	1.200	22.967	0.9999	2.61202
104.52	1.206	23.082	0.9999	2.62508
105.04	1.212	23.197	0.9999	2.63814
105.56	1.218	23.312	0.9999	2.65120
106.08	1.224	23.427	0.9999</	



STATION 3  
 UINF= 17.131 M/S  
 REF= 2089E+07  
 T= 20.18 DEG.C  
 MU= 0.1843E-04 KG/MS

X= 2.319  
 RHC= 0.9572 KG/CU.M  
 NU= 0.1901E-04 M=MS

Y (MM) U (M/S) Y/TTA U/UE  
 0.000 4.1350 0.000 0.2470  
 0.000 5.0177 0.197 0.2490  
 0.000 5.7447 0.310 0.2481  
 0.000 6.3237 0.399 0.2471  
 0.000 6.7714 0.457 0.2452  
 0.000 7.1180 0.504 0.2437  
 0.000 7.3722 0.548 0.2427  
 0.000 7.5411 0.587 0.2420  
 0.000 7.6287 0.620 0.2415  
 0.000 7.6341 0.647 0.2411  
 0.000 7.5566 0.667 0.2408  
 0.000 7.3973 0.681 0.2405  
 0.000 7.1681 0.689 0.2402  
 0.000 6.8733 0.691 0.2399  
 0.000 6.5277 0.687 0.2396  
 0.000 6.1371 0.677 0.2393  
 0.000 5.7062 0.661 0.2390  
 0.000 5.2391 0.639 0.2387  
 0.000 4.7307 0.611 0.2384  
 0.000 4.1850 0.577 0.2381  
 0.000 3.6062 0.538 0.2378  
 0.000 2.9981 0.494 0.2375  
 0.000 2.3657 0.446 0.2372  
 0.000 1.7141 0.393 0.2369  
 0.000 1.0483 0.336 0.2366  
 0.000 0.3724 0.275 0.2363  
 0.000 0.0000 0.210 0.2360

DELTA= 0.57 MM  
 HDSTR= 0.772E+04  
 RHYE= 0.61E+04

STATION A

UINF= 17.455 M/S  
 REF= 2217E+07  
 T= 20.32 DEG.C  
 MU= 0.1846E-04 KG/MS

X= 2.415  
 RHC= 0.9871 KG/CU.M  
 NU= 0.1939E-04 M=MS

Y (MM) U (M/S) Y/TTA U/UE  
 0.000 4.1666 0.000 0.2418  
 0.000 5.0500 0.211 0.2410  
 0.000 5.7800 0.331 0.2409  
 0.000 6.3633 0.407 0.2408  
 0.000 6.8055 0.465 0.2407  
 0.000 7.1174 0.505 0.2406  
 0.000 7.3000 0.532 0.2405  
 0.000 7.3722 0.550 0.2404  
 0.000 7.3448 0.558 0.2403  
 0.000 7.2174 0.555 0.2402  
 0.000 6.9900 0.541 0.2401  
 0.000 6.6633 0.519 0.2400  
 0.000 6.2467 0.493 0.2399  
 0.000 5.7450 0.463 0.2398  
 0.000 5.2533 0.429 0.2397  
 0.000 4.7767 0.391 0.2396  
 0.000 4.3100 0.350 0.2395  
 0.000 3.8583 0.307 0.2394  
 0.000 3.4267 0.261 0.2393  
 0.000 2.9981 0.212 0.2392  
 0.000 2.5767 0.161 0.2391  
 0.000 2.1667 0.108 0.2390  
 0.000 1.7733 0.053 0.2389  
 0.000 1.3983 0.000 0.2388  
 0.000 1.0442 0.000 0.2387  
 0.000 0.7142 0.000 0.2386  
 0.000 0.4100 0.000 0.2385  
 0.000 0.1350 0.000 0.2384

DELTA= 0.51 MM  
 HDSTR= 0.8713E+04  
 RHYE= 0.5422E+04

STATION A

E-1177-17

29N1

67% Blocked

CURVE 1 FIT LCG(FDI)= 0.1209E+00 LOG(Y/DI)= 0.5621E-02

UINF= 16.500 M/S  
 REF= 1189E+07  
 T= 20.40 DEG.C  
 MU= 0.1842E-04 KG/MS

X= 2.154  
 RHC= 0.9659 KG/CU.M  
 NU= 0.1899E-04 M=MS

Y (MM) U (M/S) Y/TTA U/UE  
 0.000 4.1112 0.000 0.2387  
 0.000 5.0000 0.195 0.2387  
 0.000 5.7447 0.314 0.2387  
 0.000 6.3433 0.394 0.2387  
 0.000 6.8055 0.452 0.2387  
 0.000 7.1300 0.497 0.2387  
 0.000 7.3233 0.530 0.2387  
 0.000 7.3850 0.552 0.2387  
 0.000 7.3122 0.562 0.2387  
 0.000 7.0989 0.560 0.2387  
 0.000 6.8300 0.547 0.2387  
 0.000 6.5141 0.524 0.2387  
 0.000 6.1541 0.492 0.2387  
 0.000 5.7541 0.451 0.2387  
 0.000 5.3141 0.400 0.2387  
 0.000 4.8841 0.349 0.2387  
 0.000 4.4541 0.298 0.2387  
 0.000 4.0241 0.247 0.2387  
 0.000 3.5941 0.196 0.2387  
 0.000 3.1641 0.145 0.2387  
 0.000 2.7341 0.094 0.2387  
 0.000 2.3041 0.043 0.2387  
 0.000 1.8741 0.000 0.2387  
 0.000 1.4441 0.000 0.2387  
 0.000 1.0141 0.000 0.2387  
 0.000 0.5841 0.000 0.2387  
 0.000 0.1541 0.000 0.2387

DELTA= 0.46 MM  
 HDSTR= 0.3892E+04  
 RHYE= 0.61E+04

STATION 2

UINF= 16.700 M/S  
 REF= 1047E+07  
 T= 20.13 DEG.C  
 MU= 0.1847E-04 KG/MS

X= 2.220  
 RHC= 0.9611 KG/CU.M  
 NU= 0.1901E-04 M=MS

Y (MM) U (M/S) Y/TTA U/UE  
 0.000 4.1000 0.000 0.2375  
 0.000 5.0000 0.187 0.2375  
 0.000 5.7447 0.307 0.2375  
 0.000 6.3433 0.387 0.2375  
 0.000 6.8055 0.445 0.2375  
 0.000 7.1300 0.490 0.2375  
 0.000 7.3233 0.523 0.2375  
 0.000 7.3850 0.545 0.2375  
 0.000 7.3122 0.555 0.2375  
 0.000 7.0989 0.543 0.2375  
 0.000 6.8300 0.530 0.2375  
 0.000 6.5141 0.507 0.2375  
 0.000 6.1541 0.475 0.2375  
 0.000 5.7541 0.443 0.2375  
 0.000 5.3141 0.411 0.2375  
 0.000 4.8841 0.379 0.2375  
 0.000 4.4541 0.347 0.2375  
 0.000 4.0241 0.315 0.2375  
 0.000 3.5941 0.283 0.2375  
 0.000 3.1641 0.251 0.2375  
 0.000 2.7341 0.219 0.2375  
 0.000 2.3041 0.187 0.2375  
 0.000 1.8741 0.155 0.2375  
 0.000 1.4441 0.123 0.2375  
 0.000 1.0141 0.091 0.2375  
 0.000 0.5841 0.059 0.2375  
 0.000 0.1541 0.027 0.2375

DELTA= 0.57 MM  
 HDSTR= 0.5749E+04  
 RHYE= 0.4112E+04

STATION 2









APPENDIX E

E1 Accuracy of Experimental Equipment

Table E1.1 lists the equipment used, and the errors involved.

TABLE E1.1

INSTRUMENT, MAKE & MODEL	ACCURACY OR ERROR
Cambridge potentiometer, type 44228	$\pm 0,01$ millivolt
Doric DVM, type DS-100	$\pm 0,001$ volt
Riken Denshi X-Y Recorder, type F-3D	Within $\pm 0,3\%$
Positioning of hot-wire	$\pm 0,5$ mm
Water U-tubes	+ 0,4 mm H <sub>2</sub> O
Van Essen manometer, model 6750	$\pm 0,002''$ H <sub>2</sub> O
Fortin barometer	$\pm 0,001$ mm Hg
Disa anemometer, type 55D01	Output read on DVM
Disa lineariser, type 55D10	Output read on DVM
Disa RMS voltmeter type 55D35	Output read on DVM
Disa Auxilliary Unit, type 55D25	Output read on DVM
Hewlett-Packard Spectrum analyser, type 3580 A	Output plotted by the X-Y recorder

TABLE E1.1 (cont.)

INSTRUMENT, MAKE & MODEL	ACCURACY OR ERROR
Hewlett-Packard Log Voltmeter/ Converter, type 7562	Output plotted by the X-Y recorder
Tektronix oscilloscope model with Polaroid camera	-

The values presented in Table E1.1 are the maximum errors expected.

When using the Van Essen manometer with the nozzle unit of the hot-wire calibration rig, the error in velocity may be estimated from :-

$$\frac{dV}{V} = \frac{1}{2} \left( \frac{dh}{h} - \frac{\delta\rho}{\rho} \right) \dots\dots\dots [E1.1]$$

Assuming no error in  $\rho$ ,

V(m/s)	h("H <sub>2</sub> O)	% error
2,2	0,01	10
5,0	0,05	2
16,0	0,50	0,2

Clearly, the expected error at low velocities is very large. A linear regression analysis was used to fit a curve to the calibration values, and an error of 0,1 m/s was typical.

The error in measuring the transpiration flow rate was estimated as follows (using an ordinary water manometer) :-

Basic error for orifice plate : 2,3% (BS 1042)

$$\frac{dQ}{Q} = \frac{1}{2} \left[ \frac{dh}{h} + \frac{dP_1}{P_1} + \frac{dT}{T} \right] + 2 \frac{dL}{L} \dots\dots\dots [E1.2]$$

$$\frac{dQ}{Q} = \frac{1}{2} \left[ \frac{0,5}{20} + \frac{0,5}{8534} + (0,01 + 0,004) \right] + 2 \frac{0,01}{38}$$

$$\frac{dQ}{Q} = 2\% + \text{basic error}$$

$$\frac{dQ}{Q} = 4,3\%$$

As  $v_w$  was directly proportional to  $Q$ ,

$$\frac{dv_w}{v_w} = 4,3\%$$

This could be even greater if an error had been made in calculating the surface area of the porous plate. However, as  $h$  increased, the large term,  $dh/h$ , in E1.2 decreased.

In measuring the power spectrum, the emphasis was on the shape of the curve, and the occurrence of any peaks. The exact frequency and power of these peaks was not required to be known, and therefore an error analysis was not attempted for these measurements.

The main stream velocities measured with the anemometer did not drop below 5 m/s, which was within 2% and considered acceptable.

## APPENDIX F

### F1 Some Details Pertaining to the Programme

#### F1.1 The pressure gradient term

A pressure gradient would clearly alter the main stream velocity, and therefore  $u_e$ ,  $\beta$  and  $p^+$ . It would, however, also affect the  $\xi$  grid spacing, even if constant  $x$  spacing was used in the physical plane. As a preliminary test, a constant pressure gradient was assumed. It was thought that the small increase in  $u_m$  caused by the growth of the boundary layer in an enclosed flow situation could be fitted with a straight line.

$$\text{Then } \quad du_e/dx = c \quad \dots\dots\dots \text{ [F1.1]}$$

This changes the  $x$ -axis transformation :

$$\int_0^\xi d\xi = \nu \rho \int_0^x (x_0 + u_{e0}) dx \quad \dots\dots\dots \text{ [F1.2]}$$

For condition [F1.1],

$$\beta = 1 - \left( \frac{u_{e0}}{u_e} \right), \quad \dots\dots\dots \text{ [F1.3]}$$

where  $u_{e0} = 0,99 \times U_0$ .

With injection, assumption [F1.1] proved inadequate.

#### F1.2 Numerical Integration

Certain quantities were obtained by integrating.  $\delta^*$  was obtained as follows.

$$\delta^* = - \left( \frac{2\xi}{\rho u_e} \right)^{1/2} \int_0^{\eta_\infty} \phi'(\xi, \eta) d\eta$$

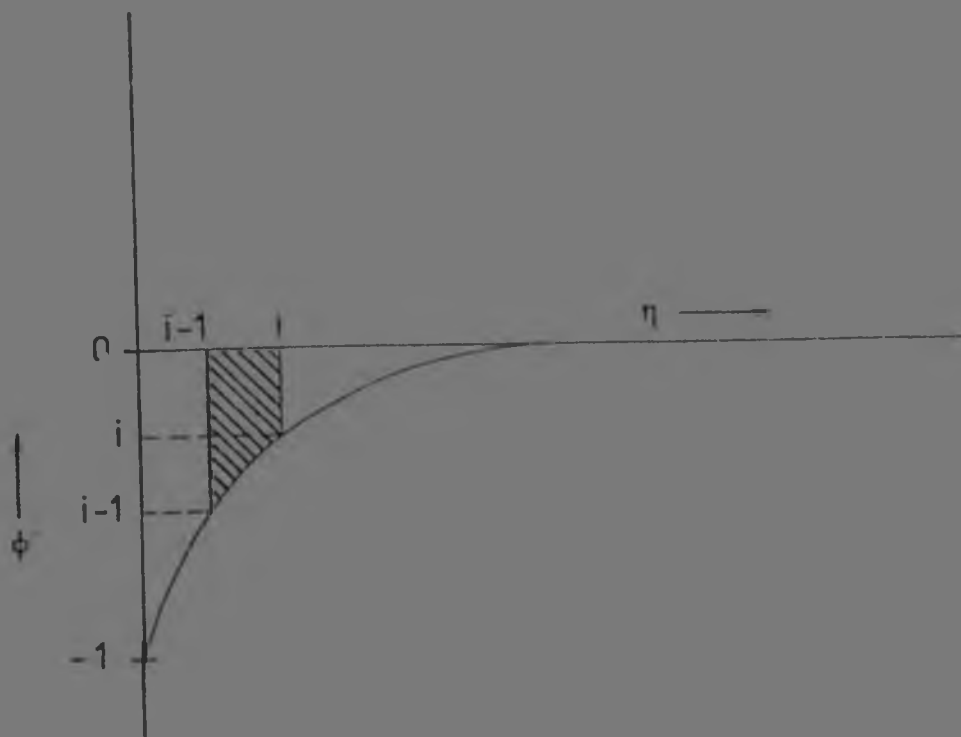


FIG F1.1 : SKETCH SHOWING TRAPEZOIDAL RULE

$$\text{Integral} = \sum_{i=2}^N (\phi_i + \phi_{i-1}) \times (\eta_i - \eta_{i-1}) \times \frac{1}{2} \dots \text{ [F1.4]}$$

F2 Initial Computer Runs

A few trial runs were attempted without transpiration. A laminar profile was generated near the point of transition, assumed to occur at  $Re_x = 3,2 \times 10^5$ . From this point  $\epsilon^+ \neq 0$ , and the programme marched into the transition region. The result obtained is shown in Fig F2.1. The uninjected laminar boundary layer with a small, linear pressure gradient ( $du/dx = 10^{1/2}$ ) was attempted next. The profiles, shown in Fig F2.2, varied slightly from the first profile in Fig F2.1 (laminar, zero pressure gradient). Finally, a series of turbulent profiles were generated by the programme, starting at  $Re_x = 3,24 \times 10^5$  with fully turbulent profiles. Values of  $c_f$  were not in satisfactory agreement with theory, thus the forward difference equation for  $\phi_w''$  was altered.

TABLE F2.1 Results of Initial Computer Runs

$x$ (m)	$\xi \times 10^4$	$A^+$	$Re_x \times 10^{-5}$	$Re_\theta$	$\theta$ (mm)	$H$	$\delta$ (mm)	$c_f \times 10^3$
MHKTRANS L-0-10 ( $x_{trip} = 0,591$ m)								
0,588	1,0578	26,0	3,185	330	0,609	2,56	5,04	1,41
0,597	1,0740	26,0	3,233	328	0,606	2,15	5,08	2,35
0,606	1,0902	26,0	3,232	333	0,616	1,90	5,23	4,18
0,615	1,1064	26,0	3,331	343	0,633	1,81	5,63	4,46
MHKDUX10 L-0-10								
0,076	0,1420	28,99	0,443	125	0,214	2,48	1,85	4,23
0,085	0,1595	29,27	0,500	132	0,224	2,47	1,96	4,01
0,094	0,1771	29,43	0,557	138	0,233	2,47	2,05	3,87
0,103	0,1949	29,55	0,616	144	0,242	2,45	2,13	3,75
MHKFLOW1 T-0-10								
0,598	1,0758	26,0	3,239	881	1,627	1,41	18,20	5,06
0,607	1,0920	26,0	3,287	893	1,649	1,41	18,70	5,16
0,616	1,1082	26,0	3,336	905	1,670	1,41	19,21	5,15
0,625	1,1244	26,0	3,385	916	1,691	1,41	19,35	5,11

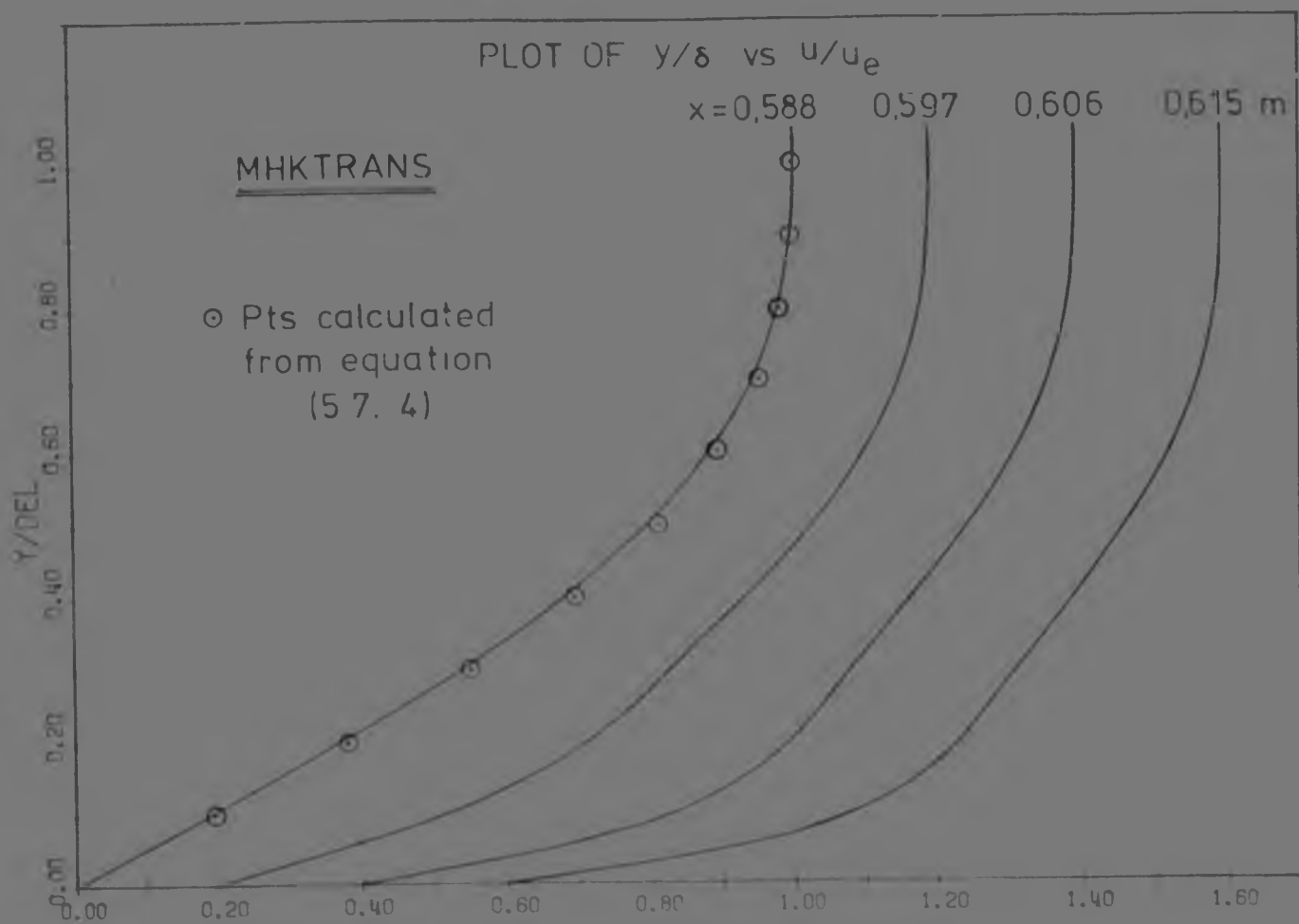


FIG F2.1 : LAMINAR PROFILES BEFORE AND DURING TRANSITION

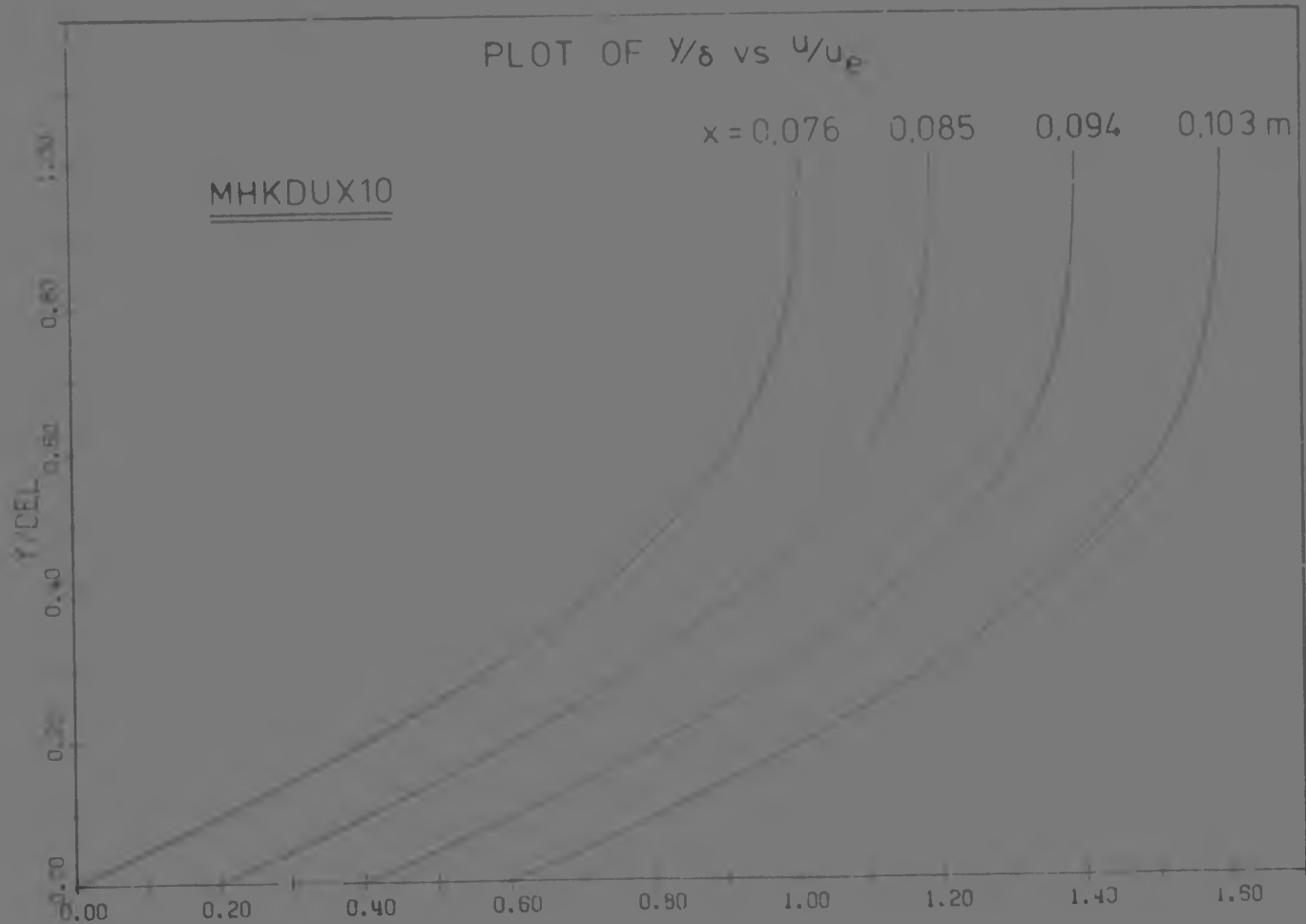


FIG F2.2 : LAMINAR PROFILES WITH PRESSURE GRADIENT

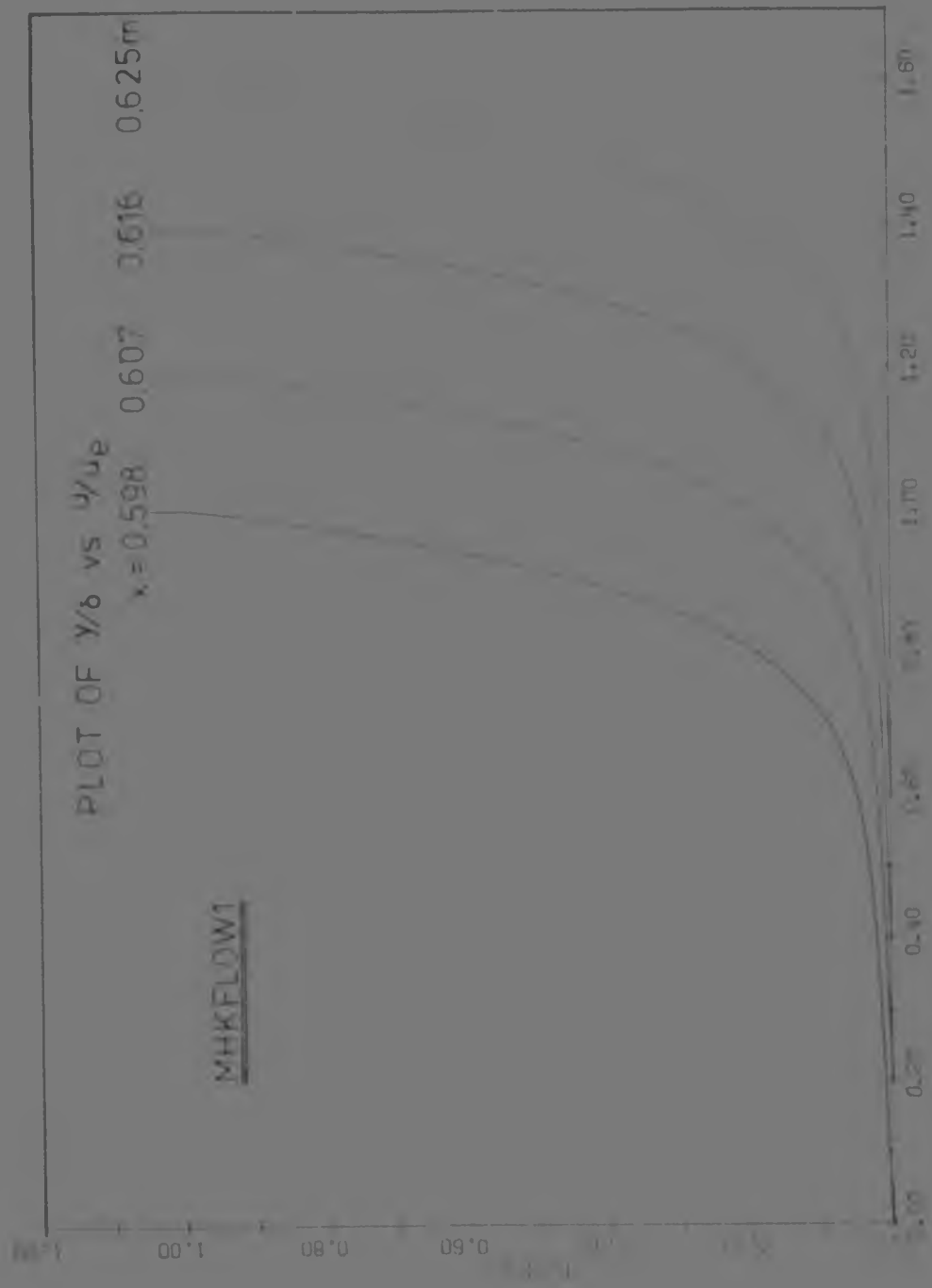


FIG F2.3 : TURBULENT INPUT PROFILES



F3 Results of Main Computer Runs

TABLE F3.1

Station	$A^+$	$R_{w\theta}$	$t$ (mm)	$H$	$\delta$ (mm)	$c_f \times 10^3$
MHKIJO01 T-1-17						
3	26,00	1812	1,97	1,33	23,14	4,085
10	23,11	1894	2,06	1,34	25,01	5,120
20	23,03	2019	2,19	1,36	26,48	4,818
30	22,95	2139	2,32	1,38	27,23	4,560
40	22,89	2257	2,45	1,39	28,74	4,357
50	22,83	2372	2,58	1,39	29,47	4,190
60	22,78	2486	2,70	1,40	31,03	4,049
MHKIJO02 T-2-17						
3	26,00	1812	1,97	1,33	23,14	4,085
10	20,93	1910	2,07	1,34	25,01	6,054
20	20,76	2065	2,24	1,37	26,48	5,580
30	20,59	2216	2,41	1,39	27,99	5,197
40	20,44	2364	2,57	1,41	28,74	4,896
50	20,31	2509	2,72	1,42	30,30	4,648
60	20,19	2652	2,88	1,43	31,89	4,438
MHKIJO03 T-3-17						
3	26,00	1811	1,97	1,33	23,53	4,088
10	19,15	1925	2,09	1,35	24,73	6,845
20	18,89	2109	2,29	1,38	26,92	6,229
30	18,64	2288	2,49	1,41	28,46	5,743
40	18,42	2465	2,68	1,43	29,22	5,361
50	18,22	2639	2,87	1,44	30,80	5,046
60	18,05	2812	3,05	1,45	32,42	4,779

Station	$A^+$	$Re_\theta$	$\theta$ (mm)	$H$	$\delta$ (mm)	$\sigma_f \times 10^3$
MHKIJ005 T-5-17						
3	26,00	1811	1,97	1,33	23,53	4,088
10	16,47	1953	2,12	1,35	24,73	8,507
20	16,09	2189	2,38	1,40	26,92	7,652
30	15,72	2421	2,63	1,43	28,46	6,980
40	15,40	2653	2,88	1,46	30,04	6,446
50	15,11	2883	3,13	1,48	31,66	6,003
60	14,84	3114	3,38	1,50	33,33	5,626
MHKIJ007 T-7-17						
3	26,00	1811	1,97	1,33	23,25	4,091
10	14,38	1978	2,15	1,36	25,13	9,909
20	13,94	2260	2,45	1,41	27,37	8,897
30	13,52	2542	2,76	1,46	28,93	8,085
40	13,14	2825	3,07	1,49	31,39	7,432
50	12,79	3109	3,38	1,52	33,08	6,886
60	12,47	3394	3,69	1,54	34,84	6,420
MHKIJ009 T-9-17						
3	26,00	1811	1,97	1,33	23,25	4,091
10	12,81	2000	2,17	1,36	25,13	11,501
20	12,39	2322	2,52	1,42	27,37	10,390
30	11,94	2646	2,87	1,47	29,74	9,450
40	11,55	2975	3,23	1,51	32,27	8,683
50	11,18	3307	3,59	1,54	34,00	8,034
60	10,84	3643	3,96	1,57	36,77	7,476
MHKIJ010 T-10-17						
3	26,00	1811	1,97	1,33	23,15	4,095
10	11,97	2010	2,18	1,36	25,74	11,853
20	11,56	2353	2,56	1,43	28,03	10,727
30	11,11	2700	2,93	1,48	29,63	9,761
40	10,70	3052	3,31	1,52	32,15	8,969
50	10,33	3409	3,70	1,56	34,82	8,298
60	9,98	3771	4,10	1,58	36,65	7,719

Station	A <sup>+</sup>	$Re_{\theta}$	$\theta$ (mm)	H	$\delta$ (mm)	$c_f \times 10^3$
MHKIJ012 T-12-17						
3	26,00	1811	1,97	1,33	23,33	4,096
10	10,78	2029	2,20	1,36	25,23	13,279
20	10,42	2403	2,61	1,43	28,25	12,142
30	9,98	2786	3,03	1,48	30,70	11,092
40	9,58	3177	3,45	1,53	33,30	10,216
50	9,22	3576	3,88	1,56	35,09	9,467
60	8,88	3982	4,32	1,60	37,96	8,817
MHKIJ015 T-15-17						
3	26,00	1810	1,97	1,33	23,70	4,098
10	9,34	2051	2,23	1,35	25,62	15,419
20	9,07	2461	2,67	1,42	27,90	14,318
30	8,66	2886	3,13	1,47	31,18	13,155
40	8,29	3325	3,61	1,52	33,82	12,162
50	7,94	3776	4,10	1,56	36,63	11,301
60	7,61	4239	4,60	1,59	39,61	10,547
MHKIJ017 T-17-17						
3	26,00	1810	1,97	1,33	23,22	4,100
10	9,27	2052	2,23	1,35	25,82	15,201
20	8,99	2464	2,68	1,42	28,12	14,105
30	8,59	2891	3,14	1,47	30,56	12,957
40	8,21	3332	3,62	1,52	34,08	11,977
50	7,87	3785	4,11	1,56	36,91	11,129
60	7,54	4249	4,61	1,59	39,92	10,386

TABLE F3.2 Turbulent Shear Layer with Injection and Pressure Gradient

Test run MHKDUX003,  $du/dx = 10^{1/3}$ ,  $F = 0,003$ ,  $XTR = 0,8$  m

$x$ (m)	$Re_x \times 10^6$	$\delta$ (mm)	$\theta$ (mm)	$Re_\theta$	$H$	$p^+ \times 10^3$	$A^+$	$c_f \times 10^3$
0,790	1,069	22,05	1,78	2406	1,31	3,49	26,55	3,842
0,825	1,132	22,32	1,80	2464	1,31	1,81	21,03	6,003
0,875	1,224	23,30	1,84	2578	,33	1,89	21,00	5,584
0,925	1,320	23,63	1,88	2685	1,34	2,00	20,95	5,250
0,975	1,418	24,61	1,92	2788	1,35	2,09	20,93	4,979
1,025	1,518	24,91	1,95	2888	1,36	2,16	20,91	4,759
1,075	1,622	25,19	1,98	2986	1,36	2,21	20,91	4,573

TABLE F3.3 Computer runs with XTR small,  $U_\infty = 17$  m/s,  $du/dx = 0^{1/3}$

$x$ (m)	$Re_x \times 10^5$	$\delta$ (mm)	$\theta$	$Re_\theta$	$H$	$A^+$	$c_f \times 10^3$
MHKLONG3 T-3 XTR = 0,312 m							
0,254	2,325	8,27	0,73	665	1,46	26,00	5,001
0,359	3,286	10,79	1,02	934	1,40	21,97	22,759
0,509	4,659	14,52	1,45	1328	1,43	21,23	15,222
0,659	6,032	18,08	1,92	1755	1,46	20,52	11,231
0,809	7,405	21,89	2,40	2200	1,47	19,94	8,978
0,959	8,778	26,02	2,89	2644	1,48	19,45	7,542
1,109	10,151	30,52	3,38	3095	1,49	18,91	6,550
MHKLONG5 T-4,765-17 XTR = 0,312 m							
0,254	3,325	8,27	0,73	665	1,46	26,00	5,001
0,359	3,286	10,80	1,01	925	1,35	20,72	31,735
0,509	4,659	14,53	1,43	1313	1,37	19,75	20,820
0,659	6,032	18,65	1,95	1786	1,42	18,78	14,980
0,809	7,405	23,26	2,51	2300	1,45	17,97	11,685
0,959	8,778	27,64	3,10	2841	1,48	17,29	9,594
1,109	10,151	32,41	3,71	3398	1,49	16,70	8,156

F4 Input Data for the Programme

NX, the number of  $\xi(x)$  stations, including the first two  
DX, the  $\xi$  step length  $\Delta\xi$ , for equal steps (as noted before,  
these steps need not be equal)  
X(1), the value of  $x$  at the first  $\xi$  station, in m  
XTR, the value of  $x$  at the start of the porous section for  
injected flow  
IY, the number of  $\eta$  grid points  
K, the ratio of adjacent  $\eta$  step lengths. If  $K = 1$ , equal  
step length results  
H, the first  $\eta$  step length  
UINF, mainstream velocity, in m/s  
DUDX, a pressure gradient term, in  $s^{-1}$   
PSTAT, static pressure, in Pa  
T, temperature of mainstream, in  $^{\circ}C$   
VW, injection velocity, in m/s  
IJOB, 0 - no plot output  
1 - plot velocity profile } For IJOB = 1, the JCL had  
to set up another job

EQUIVALENT NOTATION USED IN PROGRAMME

ROMAN

$A^+$	AP	$U_{CL}$	UCL
$c_f$	CF	$u_w$	JE
$du_e/dx$	DUDX	$u_w^*, u_\tau$	USTR
$f''$	FDD	$u^+$	UP
$h$	H	$U_w$	UINF
$H$	H12	$v_w$	VW, VW1
$K$	K	$v_w^+$	VWP
$p^+$	PP	$x$	X
$Re_x$	REX	$y$	Y
$Re_{\delta^+}$	RDSTR	$y^+$	YP
$Re_\theta$	RTTA	$y/\theta$	YOTTA
$u$	U	$y/\delta$	YODEL

GREEK

$\beta$	BETA	$\mu$	MU
$\delta$	DEL	$\nu$	NU
$\delta^+$	DSTR	$\xi$	XI
$\epsilon^+$	EP	$\rho$	RHO
$\epsilon_i^+$	EPI	$\tau$	
$\eta$	ETA	$\phi$	P
$\theta$	TTA	$\phi'$	PD
$\kappa$	KL	$\phi''$	PDD













IONS IN EFFECT: NAME (PAINT OPTIMIZE) L1Z (COUNT) L2Z (MAX AUTODIFF) SOURCE F (SIC) NDLIST (M) (C) (D) (E) (F) (G) (H) (I) (J) (K) (L) (M) (N) (O) (P) (Q) (R) (S) (T) (U) (V) (W) (X) (Y) (Z) (AA) (AB) (AC) (AD) (AE) (AF) (AG) (AH) (AI) (AJ) (AK) (AL) (AM) (AN) (AO) (AP) (AQ) (AR) (AS) (AT) (AU) (AV) (AW) (AX) (AY) (AZ) (BA) (BB) (BC) (BD) (BE) (BF) (BG) (BH) (BI) (BJ) (BK) (BL) (BM) (BN) (BO) (BP) (BQ) (BR) (BS) (BT) (BU) (BV) (BW) (BX) (BY) (BZ) (CA) (CB) (CC) (CD) (CE) (CF) (CG) (CH) (CI) (CJ) (CK) (CL) (CM) (CN) (CO) (CP) (CQ) (CR) (CS) (CT) (CU) (CV) (CW) (CX) (CY) (CZ) (DA) (DB) (DC) (DD) (DE) (DF) (DG) (DH) (DI) (DJ) (DK) (DL) (DM) (DN) (DO) (DP) (DQ) (DR) (DS) (DT) (DU) (DV) (DW) (DX) (DY) (DZ) (EA) (EB) (EC) (ED) (EE) (EF) (EG) (EH) (EI) (EJ) (EK) (EL) (EM) (EN) (EO) (EP) (EQ) (ER) (ES) (ET) (EU) (EV) (EW) (EX) (EY) (EZ) (FA) (FB) (FC) (FD) (FE) (FF) (FG) (FH) (FI) (FJ) (FK) (FL) (FM) (FN) (FO) (FP) (FQ) (FR) (FS) (FT) (FU) (FV) (FW) (FX) (FY) (FZ) (GA) (GB) (GC) (GD) (GE) (GF) 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15N 0002 SUBROUTINE M0004 (A, TAIL, AL, H, H0, P, PD, Y, CUM, V, INC, FDD, NDEL,
      PUL, PP, AP, XI, U, TR, DSTR, DEL, GLE)
      PROGRAMME TO GENERATE INITIAL VELOCITY PROFILES.
      REAL*4 SR4
      REAL*4 TAIL, AL, H0, HND, DSTR, V, USTR, UC, PP, AP, KU, FA1, SFA1, SR2, SR1,
      SR3, URC, J, DNTI (4), DIF (1), DIF (1), DIF (1), DIF (1)
      REAL*4 G (4, 100), YP (100), UP (100), F (100), P (4, 100), PD (4, 100),
      G (100), X (100)
      REAL*4 X (100)
      REAL*4 P (100)
      REAL*4 DIL
      USTR = MU * FDD * U * 0.500 * (2.000 / XI (INC)) * 0.500
      UST = USTR * 0.400
      NU = MU / RHO
      FA1 = (2.000 * XI (INC)) * 0.500
      AF X1 = FA1 * USTR / (MU * U)
      DJ 1 = 1.1Y
      YP (1) = AF X1 * ETAL (1)
      CONTINUE
      G (1) = 1.000
      DO 2 1=2, 1Y
      SR2 = YP (1) / ZAP
      IF (SR2 .LT. 0.0100) GO TO 300
      P = 0.000
      IF (SR2 .LT. 100) P = 1.000 / DEXP (SR2)
      SR4 = 1.000 - P * W
      GO TO 301
      SR4 = 0.000
      SR1 = 7.000 * XI * YP (1) * SR4
      SR3 = SR1 * 0.000
      G (1) = 2.000 / (1.000 + DEXP (1.000 + SR3))
      CONTINUE
      INTEGRATION
      UP (1) = 0.000
      DO 3 1=2, 1Y
      J = 1 - 1
      UP (1) = (G (1) * G (J)) * (YP (1) - YP (J)) / 2.000 * UP (J)
      CONTINUE
      USE = USTR / UP
      DO 4 1=1, 1Y
      PD (1) = UP (1) / USE
      CONTINUE
      DETERMINE EDGE OF BOUNDARY LAYER
      NDEL = 0
      NDEL = NDEL + 1
      DIF = 1.1, 0.000 - PD (NDEL)
      IF (DIF .LT. 0.000100) GO TO 9
      IF (NDEL .EQ. 1Y) GO TO 9

```

EVEL 2.2 (SEPT 76) HWTB 05/360 FORTMAN H EXTENDED DATE 78.027/16.01.49 PAGE 2

```

15N 0048 GO TO 8
15N 0049 NDEL = NDEL
15N 0050 DO 41 1=NDEL, 1Y
15N 0051 PD (1) = 1.000
15N 0052 CONTINUE
      INTEGRATE TO OBTAIN F
      F (1) = 0.000
      DO 5 1=2, 1Y
      J = 1 - 1
      F (1) = (PD (1) * PD (J)) * (ETAL (1) - ETAL (J)) / 2.000 * F (J)
      CONTINUE
      GENERATE P AND PD
      DO 6 1=1, 1Y
      P (INC, 1) = F (1) - ETAL (1)
      CONTINUE
      DO 7 1=1, 1Y
      PD (INC, 1) = F (1) - 1.000
      CONTINUE
      CALCULATE DSTR
      TEMP = 0.000
      DO 10 1=2, NDEL
      J = 1 - 1
      TEMP = TEMP + (PD (INC, 1) * PD (INC, J)) * (ETAL (1) - ETAL (J)) / 2.000
      CONTINUE
      DSTR (INC) = F (1) * TEMP / (H * H0)
      FUEL = ETAL (NDEL)
      DEL = FUEL * DEL / (H * H0 * UP)
      RETURN
      END

```

IONS IN EFFECT: NAME (PAINT OPTIMIZE) L1Z (COUNT) L2Z (MAX AUTODIFF) SOURCE F (SIC) NDLIST (M) (C) (D) (E) (F) (G) (H) (I) (J) (K) (L) (M) (N) (O) (P) (Q) (R) (S) (T) (U) (V) (W) (X) (Y) (Z) (AA) (AB) (AC) (AD) (AE) (AF) (AG) (AH) (AI) (AJ) (AK) (AL) (AM) (AN) (AO) (AP) (AQ) (AR) (AS) (AT) (AU) (AV) (AW) (AX) (AY) (AZ) (BA) (BB) (BC) (BD) (BE) (BF) (BG) (BH) (BI) (BJ) (BK) (BL) (BM) (BN) (BO) (BP) (BQ) (BR) (BS) (BT) (BU) (BV) (BW) (BX) (BY) (BZ) (CA) (CB) (CC) (CD) (CE) (CF) (CG) (CH) (CI) (CJ) (CK) (CL) (CM) (CN) (CO) (CP) (CQ) (CR) (CS) (CT) (CU) (CV) (CW) (CX) (CY) (CZ) (DA) (DB) (DC) (DD) (DE) (DF) (DG) (DH) (DI) (DJ) (DK) (DL) (DM) (DN) (DO) (DP) (DQ) (DR) (DS) (DT) (DU) (DV) (DW) (DX) (DY) (DZ) (EA) (EB) (EC) (ED) (EE) (EF) (EG) (EH) (EI) (EJ) (EK) (EL) (EM) (EN) (EO) (EP) (EQ) (ER) (ES) (ET) (EU) (EV) (EW) (EX) (EY) (EZ) (FA) (FB) (FC) (FD) (FE) (FF) (FG) (FH) (FI) (FJ) (FK) (FL) (FM) (FN) (FO) (FP) (FQ) (FR) (FS) (FT) (FU) (FV) (FW) (FX) (FY) (FZ) (GA) (GB) (GC) (GD) (GE) (GF) (GG) (GH) (GI) (GJ) (GK) (GL) (GM) (GN) (GO) (GP) (GQ) (GR) (GS) (GT) (GU) (GV) (GW) (GX) (GY) (GZ) (HA) (HB) (HC) (HD) (HE) (HF) (HG) (HH) (HI) (HJ) (HK) (HL) (HM) (HN) (HO) (HP) (HQ) (HR) (HS) (HT) (HU) (HV) (HW) (HX) (HY) (HZ) (IA) (IB) (IC) (ID) (IE) (IF) (IG) (IH) (II) (IJ) (IK) (IL) (IM) (IN) (IO) (IP) (IQ) (IR) (IS) (IT) (IU) (IV) (IW) (IX) (IY) (IZ) (JA) (JB) (JC) (JD) (JE) (JF) (JG) (JH) (JI) (JJ) (JK) (JL) (JM) (JN) (JO) (JP) (JQ) (JR) (JS) (JT) (JU) (JV) (JW) (JX) (JY) (JZ) (KA) (KB) (KC) (KD) (KE) (KF) (KG) (KH) (KI) (KJ) (KL) (KM) (KN) (KO) (KP) (KQ) (KR) (KS) (KT) (KU) (KV) (KW) (KX) (KY) (KZ) (LA) (LB) (LC) (LD) (LE) (LF) (LG) (LH) (LI) (LJ) (LK) (LM) (LN) (LO) (LP) (LQ) (LR) (LS) (LT) (LU) (LV) (LW) (LX) (LY) (LZ) (MA) (MB) (MC) (MD) (ME) (MF) (MG) (MH) (MI) (MJ) (MK) (ML) (MN) (MO) (MP) (MQ) (MR) (MS) (MT) (MU) (MV) (MW) (MX) (MY) (MZ) (NA) (NB) (NC) (ND) (NE) (NF) (NG) (NH) (NI) (NJ) (NK) (NL) (NM) (NO) (NP) (NQ) (NR) (NS) (NT) (NU) (NV) (NW) (NX) (NY) (NZ) (OA) (OB) (OC) (OD) (OE) (OF) (OG) (OH) (OI) (OJ) (OK) (OL) (OM) (ON) (OO) (OP) (OQ) (OR) (OS) (OT) (OU) (OV) (OW) (OX) (OY) (OZ) (PA) (PB) (PC) (PD) (PE) (PF) (PG) (PH) (PI) (PJ) (PK) (PL) (PM) (PN) (PO) (PP) (PQ) (PR) (PS) (PT) (PU) (PV) (PW) (PX) (PY) (PZ) (QA) (QB) (QC) (QD) (QE) (QF) (QG) (QH) (QI) (QJ) (QK) (QL) (QM) (QN) (QO) (QP) (QQ) (QR) (QS) (QT) (QU) (QV) (QW) (QX) (QY) (QZ) (RA) (RB) (RC) (RD) (RE) (RF) (RG) (RH) (RI) (RJ) (RK) (RL) (RM) (RN) (RO) (RP) (RQ) (RS) (RT) (RU) (RV) (RW) (RX) (RY) (RZ) (SA) (SB) (SC) (SD) (SE) (SF) (SG) (SH) (SI) (SJ) (SK) (SL) (SM) (SN) (SO) (SP) (SQ) (SR) (SS) (ST) (SU) (SV) (SW) (SX) (SY) (SZ) (TA) (TB) (TC) (TD) (TE) (TF) (TG) (TH) (TI) (TJ) (TK) (TL) (TM) (TN) (TO) (TP) (TQ) (TR) (TS) (TT) (TU) (TV) (TW) (TX) (TY) (TZ) (UA) (UB) (UC) (UD) (UE) (UF) (UG) (UH) (UI) (UJ) (UK) (UL) (UM) (UN) (UO) (UP) (UQ) (UR) (US) (UT) (UU) (UV) (UW) (UX) (UY) (UZ) (VA) (VB) (VC) (VD) (VE) (VF) (VG) (VH) (VI) (VJ) (VK) (VL) (VM) (VN) (VO) (VP) (VQ) (VR) (VS) (VT) (VU) (VV) (VW) (VX) (VY) (VZ) (WA) (WB) (WC) (WD) (WE) (WF) (WG) (WH) (WI) (WJ) (WK) (WL) (WM) (WN) (WO) (WP) (WQ) (WR) (WS) (WT) (WU) (WV) (WW) (WX) (WY) (WZ) (XA) (XB) (XC) (XD) (XE) (XF) (XG) (XH) (XI) (XJ) (XK) (XL) (XM) (XN) (XO) (XP) (XQ) (XR) (XS) (XT) (XU) (XV) (XW) (XX) (XY) (XZ) (YA) (YB) (YC) (YD) (YE) (YF) (YG) (YH) (YI) (YJ) (YK) (YL) (YM) (YN) (YO) (YP) (YQ) (YR) (YS) (YT) (YU) (YV) (YW) (YX) (YY) (YZ) (ZA) (ZB) (ZC) (ZD) (ZE) (ZF) (ZG) (ZH) (ZI) (ZJ) (ZK) (ZL) (ZM) (ZN) (ZO) (ZP) (ZQ) (ZR) (ZS) (ZT) (ZU) (ZV) (ZW) (ZX) (ZY) (ZZ)

```

15N 0002 SUBROUTINE M0005 (G, G0, K, H, Y, INC, ETAL)
      GENERAL DIFFERENTIAL EQUATION
      REAL*8 G (4, 100), G0 (4, 100), K,
      ETAL (100)
      REAL*8 H, X (100), YP (100), AS, BS, CS
      H = INC
      H1 = 1.0 * H
      KPI = K / H
      YP (1) = 0
      AS = Y * AS / (H * YP (1))
      HS = Y * BS / H
      CS = K / (H * KPI)
      DO 1 1=2, 1Y
      G (1, 1) = (AS * G (H, 1 - 1) + BS * G (H, 1) + CS * G (H, 1 + 1)) / 4.0
      CONTINUE
      TO OBTAIN G (H, 1) AND G (H, Y), A FIRST ORDER POLYNOMIAL HAS
      TO BE USED
      G (H, 1) = G (H, 1) * Y + (Y + 2.0) * G (H, 2) * YP (1) * YP (1) - G (H, 2)
      G (H, 1) = G (H, 1) / (H * YP (1))
      YH2 = Y - 2
      G (H, 1Y) = YP (1) * G (H, 1YH2) - Y * G (H, 1YH2)
      RETURN
      END

```



LEVEL 2.2 (SEPT 76)

05/360 FORTRAN M EXTENDED

DATE 78.027/16.02.10

PAGE 1

REQUESTED OPTIONS:

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(1) LINECOUNT(0) SIZE(MAX) AUTOBOL(NONE) SOURCE(FORTRAN) NOLIST NOCHECK OBJECT NOMAP NOFORMAT NOGOINT NOXREF NOALC NOANSF TERM FLAG(1)

```

15N 0002      SUBROUTINE MK009(X1,P,PD,PPD,NCA,NCA1,NCA2,NC,IY,NITEM,U)
PROGRAMME TO CALCULATE VECTOR R
15N 0003      REAL*8 LL,MH,DH1,DH2,T1,T2,T3,FR,T4
15N 0004      REAL*8 X1(10)
15N 0005      REAL*8 P(4,100),PD(4,100),PPD(4,100),D(96,1)
15N 0006      IY=1,DD=1(1)
15N 0007      IY=1
15N 0008      IF(NITEM.GT.0) GO TO 10
15N 0009      NT=NCA1
15N 0010      DO 1 J=1,NT
15N 0011      DH1=X1(NC-1)-X1(NC-2)
15N 0012      DH2=X1(NC)-X1(NC-1)
15N 0013      LL=DD*(DH1+DH2)/(DH1+DH2)
15N 0014      MH=- (DH1+DD*(DH1+DH2))/(DH1+DH2)
15N 0015      DO 1 J=1,NT
15N 0016      FR=LL*PD(NCA2,1)+MH*PPD(NCA1,1)
15N 0017      T1=(PD(NT,1)+1.C)/T1
15N 0018      T2=(PD(NT,1)+1.C)/T1
15N 0019      T3=LL*PD(NCA2,1)+MH*PPD(NCA1,1)
15N 0020      T4=PPD(1,1)+T3
15N 0021      D(J,1)=FR*(T2-T4)
15N 0022      CONTINUE
15N 0023      RETURN
15N 0024      END
15N 0025

```

LEVEL 2.2 (SEPT 76)

05/360 FORTRAN M EXTENDED

DATE 78.027/16.02.14

PAGE 1

REQUESTED OPTIONS:

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(1) LINECOUNT(0) SIZE(MAX) AUTOBOL(NONE) SOURCE(FORTRAN) NOLIST NOCHECK OBJECT NOMAP NOFORMAT NOGOINT NOXREF NOALC NOANSF TERM FLAG(1)

```

15N 0002      SUBROUTINE MK010(H,P,IY)
PROGRAMME TO WRITE B AS W
15N 0003      REAL*8 B(96,1),P(4,100)
15N 0004      IY=1,DD=1
15N 0005      DO 1 J=1,IF
15N 0006      I=J*2
15N 0007      P(A,1)=B(I,1)
15N 0008      CONTINUE
15N 0009      FIND P(4,IY-1) AND P(4,IY)
15N 0010      P(A,IY-1)=P(A,IY-2)
15N 0011      P(A,IY)=P(A,IY-2)
15N 0012      RETURN
15N 0013      END

```

LEVEL 2.2 (SEPT 75)

05/360 FORTRAN M EXTENDED

DATE 78.027/16.02.19

PAGE 1

REQUESTED OPTIONS:

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(1) LINECOUNT(0) SIZE(MAX) AUTOBOL(NONE) SOURCE(FORTRAN) NOLIST NOCHECK OBJECT NOMAP NOFORMAT NOGOINT NOXREF NOALC NOANSF TERM FLAG(1)

```

15N 0002      SUBROUTINE MK011(P,M,IY)
PROGRAMME TO CALCULATE PD WITH BOUNDARY CONDITIONS
15N 0003      REAL*8 P(4,100),M(4,100)
15N 0004      REAL*8 H,M,TB,TC,K1,KP1,KIM
15N 0005      PD(A,1)=1.000
15N 0006      KP1=M*1.000
15N 0007      KIM=1.000-K
15N 0008      IY=1
15N 0009      TA=-K*0.3.C/D/(H*KP1)
15N 0010      TB=-K*M*K/H
15N 0011      TC=K/(H*KP1)
15N 0012      DO 1 I=2,IY
15N 0013      P(I,1)=P(A,I-1)+TA/K+P(A,I-1)*TB/K+P(A,I-1)*TC/K
15N 0014      CONTINUE
15N 0015      PD(A,IY)=0.000
15N 0016      RETURN
15N 0017      END

```

LEVEL 2.2 (SEPT 76)

05/360 FORTRAN M EXTENDED

DATE 78.027/16.02.24

PAGE 1

REQUESTED OPTIONS:

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(1) LINECOUNT(0) SIZE(MAX) AUTOBOL(NONE) SOURCE(FORTRAN) NOLIST NOCHECK OBJECT NOMAP NOFORMAT NOGOINT NOXREF NOALC NOANSF TERM FLAG(1)

```

15N 0002      SUBROUTINE MK012(P,T1A,PD,P,TA,K,TA2)
PROGRAMME TO CALCULATE PDD FOR 14 W
15N 0003      REAL*8 P(4,100),T1A,PD,P,TA,K,TA2
15N 0004      REAL*8 P(4,100),T1A,TA2,T3A,K1,KF
15N 0005      REAL*8 M,T1A,T2A,T3A,K1,KF
15N 0006      IY=1
15N 0007      M=1.000
15N 0008      K1=M*1.000
15N 0009      T1A=-M*0.3.C/D/(H*KP1)
15N 0010      T2A=-M*M/K/H
15N 0011      T3A=K/(H*KP1)
15N 0012      DO 1 I=2,IY
15N 0013      K1=M*1
15N 0014      PDD(A,1)=(P(A,I-1)+T1A+P(A,I)+T2A+P(A,I+1)+T3A)/IY
15N 0015      CONTINUE
15N 0016      DE=1.000
15N 0017      T1A=1.000
15N 0018      T2A=1.000
15N 0019      T3A=1.000
15N 0020      PDD(A,1)=P(A,1)+T1A*(1-T1A)
15N 0021      IY=1
15N 0022      PDD(A,IY)=(PDD(A,IY)-PDD(A,IY+1))/(PDD(A,IY))
15N 0023      RETURN
15N 0024      END

```

CTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(1) LIN(COUNT(10) SIZE(MAX) AUTODEL(INCN)) SOURCE(FORCIC NOLIST NODEC) OBJECT(NOMAP NOFORAT NOGOSTMT NOXREF NOALC NOANSF TERM FLAG(1))

```

15N 0002 C SUBROUTINE MKD1(PD,NDEL,ETA,NCA,LY)
PROGRAMME TO CALCULATE THE VALUE OF THE BOUNDARY LAYER
REAL*8 PD(4,100),ETA(100)
REAL*8 P(4,1),D(4,1)
NDEL=0
DO 1 I=1,NDEL
  D(I,1)=PD(4,I)
  IF (D(I,1).GT.0.000100) GO TO 2
  IF (D(I,1).GT.1) GO TO 1
  GO TO 1
2 D(I,1)=PD(4,I)
WRITE(6,100) D(I,1),NDEL
FORMAT(7F10.4,1P10.4) AND NDEL=,14.7)
RETURN
END

```

CTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(1) LIN(COUNT(10) SIZE(MAX) AUTODEL(INCN)) SOURCE(FORCIC NOLIST NODEC) OBJECT(NOMAP NOFORAT NOGOSTMT NOXREF NOALC NOANSF TERM FLAG(1))

```

15N 0002 C SUBROUTINE MKD2(PD,NDEL,K,H,DSTR,NC,ETA)
PROGRAMME TO CALCULATE THE VALUE OF THE BOUNDARY LAYER
REAL*8 PD(4,100),ETA(100)
REAL*8 P(4,1),D(4,1),ATP,PX
REAL*8 X(100)
15N 0003 C ATP=0.000
15N 0004 DO 1 I=1,NDEL
15N 0005 J=1
15N 0006 ATP=ATP+(PD(4,I)+PD(4,J))*((TA(1)-ETA(J))/2.000)
15N 0007 CONTINUE
15N 0008 PX=ATP*(NC)*2.000+0.500
15N 0009 DSTR=PX/ATP
15N 0010 RETURN
15N 0011 END

```

EVEL 2.2 (SEPT 76) OS/J60 FORTRAN M EXTENDED DATE 78.027/16.02.39 PAGE 1

WJESTED OPTIONS:

CTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(1) LIN(COUNT(10) SIZE(MAX) AUTODEL(INCN)) SOURCE(FORCIC NOLIST NODEC) OBJECT(NOMAP NOFORAT NOGOSTMT NOXREF NOALC NOANSF TERM FLAG(1))

```

15N 0002 C SUBROUTINE MKD3(PD,PDD,OS,NCA,LY)
PROGRAMME TO WRITE THE PRESENT ITERATION INTO THE PRESENT STATION
15N 0003 C REAL*8 P(4,100),PD(4,100),PDD(4,100),DSTR(4)
15N 0004 DO 1 I=1,LY
15N 0005 P(INCA,1)=PD(4,1)
15N 0006 CONTINUE
15N 0007 DO 2 I=1,LY
15N 0008 PD(4,CA,1)=PD(4,1)
15N 0009 CONTINUE
15N 0010 DO 3 I=1,LY
15N 0011 PDD(INCA,1)=PDD(4,1)
15N 0012 CONTINUE
15N 0013 DSTR(INCA) DSTR(4)
15N 0014 RETURN
15N 0015 END

```

CTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(1) LIN(COUNT(10) SIZE(MAX) AUTODEL(INCN)) SOURCE(FORCIC NOLIST NODEC) OBJECT(NOMAP NOFORAT NOGOSTMT NOXREF NOALC NOANSF TERM FLAG(1))

```

15N 0002 C SUBROUTINE MKD4(PD,PDD,OS,NCA,LY,ATP,PX,PX1)
PROGRAMME TO CALCULATE THE TA
15N 0003 C REAL*8 PD(4,100),PDD(4,100),ETA,ATP,PX,PX1
15N 0004 REAL*8 P(4,1),D(4,1),ATP1,PX1
15N 0005 ATP=0.000
15N 0006 DO 1 I=1,NDEL
15N 0007 J=1
15N 0008 T1=PD(INCA,1)*(1.000+D(INCA,1))
15N 0009 T2=PD(INCA,1)*(1.000+D(INCA,1))
15N 0010 ATP=ATP+(T1+T2)*((TA(1)-ETA(J))/2.000)
15N 0011 CONTINUE
15N 0012 PX1=ATP*(NC)*2.000+0.500
15N 0013 TA=PX1/ATP
15N 0014 RETURN
15N 0015 END

```

CTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(1) LIN(COUNT(10) SIZE(MAX) AUTODEL(INCN)) SOURCE(FORCIC NOLIST NODEC) OBJECT(NOMAP NOFORAT NOGOSTMT NOXREF NOALC NOANSF TERM FLAG(1))

```

15N 0002 C SUBROUTINE MKD5(PD,PDD,OS,NCA,LY,ATP,PX,PX1,TAU,PK)
PROGRAMME TO CALCULATE VARIABLES NEEDED TO PRINTING
15N 0003 C REAL*8 P(4,100),PDD(4,100),PDD1(4,100),UV(100),V(100)
15N 0004 PDDTA(100),X(100),X1(100)
15N 0005 REAL*8 CF,UTM,DSTR(4),UE,TA,PK,ETA,PD(4,100),MU,DEL,D,L,TA
15N 0006 P(4,1)
15N 0007 P(4,1)=PD(4,1)
15N 0008 PX1=(X(100)+2.000)*0.500
15N 0009 CF=MUP(1.000/1.000)*0.500+PD(INCA,1)
15N 0010 TAU=(PX1/PX1)*0.500
15N 0011 U(1)=TAU/PHU
15N 0012 U(2)=(TAU/PHU)*0.500
15N 0013 DO 1 I=1,LY
15N 0014 V(I)=PDDTA(I)/(PHU*CF)
15N 0015 CONTINUE
15N 0016 DEL=V(1)
15N 0017 PK=PHU*U(1)
15N 0018 D(4,1)=PDD(4,1)
15N 0019 D(4,1)=PDD(4,1)
15N 0020 D(4,1)=PDD(4,1)
15N 0021 D(4,1)=PDD(4,1)
15N 0022 D(4,1)=PDD(4,1)
15N 0023 D(4,1)=PDD(4,1)
15N 0024 D(4,1)=PDD(4,1)
15N 0025 D(4,1)=PDD(4,1)
15N 0026 D(4,1)=PDD(4,1)
15N 0027 D(4,1)=PDD(4,1)
15N 0028 D(4,1)=PDD(4,1)
15N 0029 D(4,1)=PDD(4,1)
15N 0030 D(4,1)=PDD(4,1)
15N 0031 D(4,1)=PDD(4,1)
15N 0032 D(4,1)=PDD(4,1)
15N 0033 D(4,1)=PDD(4,1)
15N 0034 D(4,1)=PDD(4,1)
15N 0035 D(4,1)=PDD(4,1)
15N 0036 D(4,1)=PDD(4,1)
15N 0037 D(4,1)=PDD(4,1)
15N 0038 D(4,1)=PDD(4,1)
15N 0039 D(4,1)=PDD(4,1)
15N 0040 D(4,1)=PDD(4,1)
15N 0041 D(4,1)=PDD(4,1)
15N 0042 D(4,1)=PDD(4,1)
15N 0043 D(4,1)=PDD(4,1)
15N 0044 D(4,1)=PDD(4,1)
15N 0045 D(4,1)=PDD(4,1)
15N 0046 D(4,1)=PDD(4,1)
15N 0047 D(4,1)=PDD(4,1)
15N 0048 D(4,1)=PDD(4,1)
15N 0049 D(4,1)=PDD(4,1)
15N 0050 D(4,1)=PDD(4,1)
15N 0051 D(4,1)=PDD(4,1)
15N 0052 D(4,1)=PDD(4,1)
15N 0053 D(4,1)=PDD(4,1)
15N 0054 D(4,1)=PDD(4,1)
15N 0055 D(4,1)=PDD(4,1)
15N 0056 D(4,1)=PDD(4,1)
15N 0057 D(4,1)=PDD(4,1)
15N 0058 D(4,1)=PDD(4,1)
15N 0059 D(4,1)=PDD(4,1)
15N 0060 D(4,1)=PDD(4,1)
15N 0061 D(4,1)=PDD(4,1)
15N 0062 D(4,1)=PDD(4,1)
15N 0063 D(4,1)=PDD(4,1)
15N 0064 D(4,1)=PDD(4,1)
15N 0065 D(4,1)=PDD(4,1)
15N 0066 D(4,1)=PDD(4,1)
15N 0067 D(4,1)=PDD(4,1)
15N 0068 D(4,1)=PDD(4,1)
15N 0069 D(4,1)=PDD(4,1)
15N 0070 D(4,1)=PDD(4,1)
15N 0071 D(4,1)=PDD(4,1)
15N 0072 D(4,1)=PDD(4,1)
15N 0073 D(4,1)=PDD(4,1)
15N 0074 D(4,1)=PDD(4,1)
15N 0075 D(4,1)=PDD(4,1)
15N 0076 D(4,1)=PDD(4,1)
15N 0077 D(4,1)=PDD(4,1)
15N 0078 D(4,1)=PDD(4,1)
15N 0079 D(4,1)=PDD(4,1)
15N 0080 D(4,1)=PDD(4,1)
15N 0081 D(4,1)=PDD(4,1)
15N 0082 D(4,1)=PDD(4,1)
15N 0083 D(4,1)=PDD(4,1)
15N 0084 D(4,1)=PDD(4,1)
15N 0085 D(4,1)=PDD(4,1)
15N 0086 D(4,1)=PDD(4,1)
15N 0087 D(4,1)=PDD(4,1)
15N 0088 D(4,1)=PDD(4,1)
15N 0089 D(4,1)=PDD(4,1)
15N 0090 D(4,1)=PDD(4,1)
15N 0091 D(4,1)=PDD(4,1)
15N 0092 D(4,1)=PDD(4,1)
15N 0093 D(4,1)=PDD(4,1)
15N 0094 D(4,1)=PDD(4,1)
15N 0095 D(4,1)=PDD(4,1)
15N 0096 D(4,1)=PDD(4,1)
15N 0097 D(4,1)=PDD(4,1)
15N 0098 D(4,1)=PDD(4,1)
15N 0099 D(4,1)=PDD(4,1)
15N 0100 D(4,1)=PDD(4,1)

```

EVEL 2.2 (SEPT 76)

05/360 FORTRAN H EXTENDED

DATE 76.027/16.02.54

PAGE 1

LISTED OPTIONS:

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(1) LINE COUNT(60) SIZE(MAX) AUTODIAG(MORE)  
SOURCE(FRCDIC) NODLIST NODCHK OBJECT NODMAP NODFORMAT NODOSTMT NODREF NODJALC NODANSF TERM FLAG(1)

```

15N 0002      SUBROUTINE AAZM(AA,0,N)
15N 0003      1=1/CIT PLAL03 (A-H,0-7)
15N 0004      DIM N (ON A(06,07),X(96),AA(06,07),J(96,11)
15N 0005      J1=1
15N 0006      M1=N*1
15N 0007      DO 1 J=1,N
15N 0008      DO 2 I=1,N
15N 0009      2 CONTINUE
15N 0010      1 CONTINUE
15N 0011      J=3
15N 0012      DO 1000 I=1,J
15N 0013      A(I,1)=AA(I,J)
15N 0014      J=J-1
15N 0015      1000 CONTINUE
15N 0016      J=4
15N 0017      DO 1001 I=1,4
15N 0018      A(I,2)=AA(I,J)
15N 0019      J=J-1
15N 0020      1001 CONTINUE
15N 0021      C
15N 0022      M4=N-4
15N 0023      DO 1002 IS=1,M4
15N 0024      J=5
15N 0025      JJ=IS+2
15N 0026      IT=IS+4
15N 0027      DO 1003 I=1,IT
15N 0028      A(I,JJ)=AA(I,J)
15N 0029      J=J-1
15N 0030      1003 CONTINUE
15N 0031      1002 CONTINUE
15N 0032      C
15N 0033      J=5
15N 0034      M3=N-3
15N 0035      M1=N-1
15N 0036      DO 1004 I=N-3,N
15N 0037      A(I,M1)=AA(I,J)
15N 0038      J=J-1
15N 0039      1004 CONTINUE
15N 0040      J=5
15N 0041      M2=N-2
15N 0042      DO 1005 I=N-2,N
15N 0043      A(I,1)=AA(I,J)
15N 0044      J=J-1
15N 0045      1005 CONTINUE
15N 0046      DO 1006 I=1,N
15N 0047      A(I,NP1)=0(I,1)
15N 0048      1006 CONTINUE
15N 0049      M7=1
15N 0050      N=JJ-1
15N 0051      IF (N.E.0) STOP
15N 0052      DEL=1.

```

EVEL 2.2 (SEPT 76)

AAZM

05/360 FORTRAN H EXTENDED

DATE 76.027/16.02.54

PAGE 2

```

15N 0053      L=N-1
15N 0054      NP=N*1
15N 0055      DO 36 K=1,L
15N 0056      M=K*1
15N 0057      ADD=DARS(A(K,K))
15N 0058      IR=K
15N 0059      DO 38 J=K,N
15N 0060      ATT=DARS(A(I,K))
15N 0061      IF (ADD-ATT)37,30,38
15N 0062      37 ADD=ATT
15N 0063      IR=J
15N 0064      38 CONTINUE
15N 0065      IF (IR=K)39,41,JJ
15N 0066      39 DET=-DET
15N 0067      DO 40 J=K,ND
15N 0068      S=A(IR,J)
15N 0069      A(IR,J)=A(K,J)
15N 0070      A(K,J)=S
15N 0071      DO 34 I=K,N
15N 0072      IF (A(I,K).E.0.0160)
15N 0073      S=A(I,K)/A(K,K)
15N 0074      DO 36 J=K,ND
15N 0075      A(I,J)=A(I,J)-S*A(K,J)
15N 0076      36 CONTINUE
15N 0077      X(N)=1/(N,ND)*A(N,N)
15N 0078      DO 43 I=1,I
15N 0079      K=N-1
15N 0080      M=K*1
15N 0081      S=0.
15N 0082      DO 42 J=K,N
15N 0083      S=S+A(K,J)*X(J)
15N 0084      X(K)=(A(K,ND)-S)/A(K,K)
15N 0085      42 CONTINUE
15N 0086      IF (M7.E.0)GOTO 30
15N 0087      33 CONTINUE
15N 0088      DO 1007 I=1,N
15N 0089      P(I,1)=X(I)
15N 0090      1007 CONTINUE
15N 0091      M1=N
15N 0092      END

```

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(1) LINE COUNT(60) SIZE(MAX) AUTODIAG(MORE)

OPTIONS IN EFFECT: SOURCE(FRCDIC) NODLIST NODCHK OBJECT NODMAP NODFORMAT NODOSTMT NODREF NODJALC NODANSF TERM FLAG(1)

STATISTICS\* SOURCE STATEMENTS = 93, PROGRAM SIZE = 76, K, SUBPROGRAM NAME = AAZM

STATISTICS\* NO DIAGNOSTICS GENERATED

\*\*\* END OF COMPILE \*\*\*

375K BYTES OF CODE NOT USED

STATISTICS\* NO DIAGNOSTICS THIS STEP



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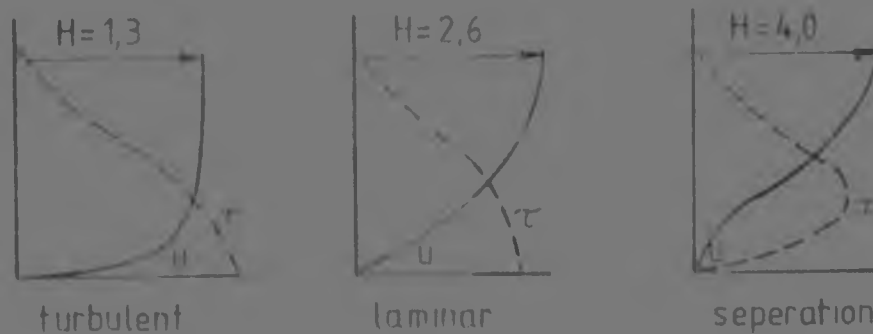
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## ADDENDUM

The ratio of displacement thickness to momentum thickness,  $H$ , numerically quantifies the shape of the boundary layer. The "fuller" the profile, the smaller  $H$ .

The following admittedly over-simplified analysis suggests the trend discussed. A consideration of the following sketches indicates the typical values of  $H$ , and its relationship with other properties of the boundary layer.



ie. as  $H$  increases,  $\tau$  decreases.

Thus  $H = f(1/\tau)$

From Reynolds Analogy, and experimental results for heat transfer,

$$St = \tau / 2 \quad (\text{Bird, Stewart and Lightfoot, 1960})$$

And with injection :-

$$St = \tau / 2 \times 1,16 \quad \mu = 0,004 \quad (\text{Moffat and Kays, 1968})$$

Thus  $St = f(1/H)$

However,  $\sigma_p$  decreases with injection ratio, as does  $St$  (see Moffat and Kays, 1968).

Thus, as  $H$  increases,  $St$  decreases. From Fig 7.13, it is clear that  $H$  increased with  $F$ , and for the restricted flow case,  $H$  was higher for a given  $x$  station and injection ratio, for fixed primary and secondary flow rates.

By definition,  $St = \frac{h}{\rho c_p V_c}$        $h =$  heat transfer coefficient

Heat transfer,  $Q = h(t_e - t_w)$

If  $St$  decreases, so does  $h$ , and for hypothetically constant  $Q$  and  $t_e$ ,  $t_w$  must decrease. Hence a lower wall temperature results, which is the desired effect.

Considering the injection velocity profile with blockage shown in Fig 7.14, leads to comparing it with the turbulent velocity, varying in space rather than with time. Although the average injected momentum is unaltered for a fixed secondary flowrate, it may be more meaningful to use some form of root mean square value for  $v$ , with blockage. The greater penetration of  $y$  momentum clearly has a greater decelerating effect on the boundary layer - it may also decrease the heat transfer coefficient. However, only extensive experimentation can conclusively verify this qualitative speculation.

**Author** Krieg M

**Name of thesis** An investigation into practical aspects of transplantation flows 1978

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