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CONTACT PROBLEMS FOR INTERFACE CRACKS UNDER HARMONIC SHEAR LOADING

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Key words: Interface Crack Closure, Boundary Integrals, Friction.

Summary. The linear crack between two dissimilar elastic isotropic half-spaces under normal harmonic shear loading is considered. To take the crack faces interaction into account we assumed that the contact satisfies the Signorini constraints and the Coulomb friction law. The problem is solved numerically using the iterative process – the solution changes until the distribution of physical values satisfying the contact constraints is found. The numerical convergence of the method with respect to the number of the Fourier coefficients and mesh size is analysed. The effects of material properties and values of the friction coefficient on the distribution of displacements and contact forces are presented and analysed. Special attention is paid to the size of the contact zone and the results are compared with the classical model solutions obtained for the static problems with and without friction.

1 INTRODUCTION

Since all existing engineering materials contain cracks, notches, delaminations and other crack-like defects, a considerable body of work is devoted to the solution of two- and three-dimensional fracture mechanics problems for cracked materials under static and dynamic loading [1–12].

It shall be specifically noted that the opposite faces of the existing cracks almost always interact with each other under deformation, significantly changing the solution near the crack [13–23]. The nature of the dynamic contact interaction between opposite crack faces is very complex. Under deformation of the material the initial contact region changes in time; its shape is unknown beforehand and must be determined as a part of the solution. The complexity of the problem is further compounded by the fact that the contact behaviour is very sensitive to the material properties of the two contacting surfaces and parameters of the external loading. It is

only possible to solve these problems using advanced numerical methods, since the analytical solutions are limited to a relatively small number of idealized model problems corresponding to very special geometrical configurations and loading conditions. Consequently, due to the nonlinearity of the problem and substantial computational difficulties, researchers often neglect effects of crack faces contact interaction, for comprehensive reviews please see [16, 17].

The current study is devoted to the analysis of the size of the contact zone for the interface crack under shear loading. The results are obtained taking friction into account and completely disregarding it using the iterative algorithm [24], and the solution is compared with the classical model solutions for the static problems with and without friction [25].

2 STATEMENT OF THE PROBLEM

Let us consider an elastic bimaterial consisting of two homogeneous isotropic half-spaces, $\Omega^{(1)}$ and $\Omega^{(2)}$, with the crack of finite length located at the interface between half-spaces. The crack has no initial opening, and the normal harmonic shear loading with the frequency of $\omega = 2\pi/T$ is applied.

In both half-spaces the equation of motion and the generalized Hooke's law lead to the linear Lamé equations of elastodynamics for the displacement field with the standard initial and boundary conditions for displacements and stresses (namely, no initial deformations; given initial load at the crack faces, $\Gamma^{(1)cr}$ and $\Gamma^{(2)cr}$; continuity conditions at the bonding interface, $\Gamma^* = \Gamma^{(1)} \cap \Gamma^{(2)}$; and the Sommerfeld radiation-type condition at the infinity). Furthermore, the components of the displacement could be represented in terms of the boundary displacements and tractions using the Somigliana dynamic identity with the appropriate fundamental solutions $U_{ij}^{(m)}(\mathbf{x}, \mathbf{y}, t - \tau)$ and $W_{ij}^{(m)}(\mathbf{x}, \mathbf{y}, t - \tau)$ [1, 7, 12, 16, 17, 20]:

$$u_j^{(m)}(\mathbf{x}, t) = \int_{\Gamma^{(m)}} \int_{\Gamma^{(m)}} (p_i^{(m)}(\mathbf{y}, \tau) U_{ij}^{(m)}(\mathbf{x}, \mathbf{y}, t - \tau) - u_i^{(m)}(\mathbf{y}, \tau) W_{ij}^{(m)}(\mathbf{x}, \mathbf{y}, t - \tau)) dy d\tau, \quad (1)$$

$$\mathbf{x} \in \Omega^{(m)}, \quad t \in T, \quad j=1,2.$$

As it was mentioned before, due to the crack's closure the traction vector at the crack surface is the superposition of the initial traction caused by the incident load, $\mathbf{g}(\mathbf{x}, t)$, and the contact force, $\mathbf{q}(\mathbf{x}, t)$, that arises in the contact region, which is unknown beforehand, depends on the properties of the bimaterial and the friction coefficient and must be determined as a part of the solution.

To include the contact interaction into account, the Signorini unilateral constraints and the Coulomb friction law with the friction coefficient k_τ must be imposed for the normal and tangential components of the displacement discontinuity, $[\mathbf{u}(\mathbf{x}, t)] = \mathbf{u}^{(1)}(\mathbf{x}, t) - \mathbf{u}^{(2)}(\mathbf{x}, t)$, and contact forces:

$$[u_n(\mathbf{x}, t)] \geq 0, \quad q_n(\mathbf{x}, t) \geq 0, \quad [u_n(\mathbf{x}, t)] q_n(\mathbf{x}, t) = 0, \quad (2)$$

$$|q_\tau(\mathbf{x}, t)| < k_\tau q_n(\mathbf{x}, t) \Rightarrow \frac{\partial [u_\tau(\mathbf{x}, t)]}{\partial t} = 0, \quad (3)$$

$$|q_\tau(\mathbf{x}, t)| = k_\tau q_n(\mathbf{x}, t) \Rightarrow \frac{\partial [u_\tau(\mathbf{x}, t)]}{\partial t} = - \frac{q_\tau(\mathbf{x}, t)}{|q_\tau(\mathbf{x}, t)|} \left| \frac{\partial [u_\tau(\mathbf{x}, t)]}{\partial t} \right|. \quad (4)$$

The contact constraints (2)–(4) ensure that there is no interpenetration of the opposite crack faces, the normal component of the contact force is unilateral; and the opposite crack faces remain immovable with respect to each other in tangential direction while they are held by the friction force before the slipping occurs [16, 17, 20, 24].

Contact interaction makes the problem non-linear and the solution even for the case of harmonic loading becomes cyclic, but non-monoharmonic. That is why the normal and tangential components of the displacement discontinuity and the traction at the crack surface can be approximated by the following exponential Fourier series with respect to the time [16]:

$$f(\bullet, t) = \operatorname{Re} \left\{ \sum_{k=-\infty}^{+\infty} f^k(\bullet) e^{i\omega_k t} \right\}, \quad f^k(\bullet) = \frac{\omega}{2\pi} \int_0^T f(\bullet, t) e^{-i\omega_k t} dt, \quad (5)$$

where $\omega_k = 2\pi k/T$ and i is the imaginary unit.

Thus, the system of boundary integral equations in the frequency domain can be represented as follows [20, 24]:

$$-\int_{\Gamma^{(1)\text{cr}}} \tilde{p}_i^{k,(1)}(\mathbf{y}) U_{ij}^{(1)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y} = -\frac{1}{2} u_i^{k,(1)}(\mathbf{x}) - \int_{\Gamma^{(1)\text{cr}}} u_i^{k,(1)}(\mathbf{y}) W_{ij}^{(1)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y} + \int_{\Gamma^*} u_i^{k,*}(\mathbf{y}) W_{ij}^{(1)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y} - \int_{\Gamma^*} p_i^{k,*}(\mathbf{y}) U_{ij}^{(1)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y}, \quad \mathbf{x} \in \Gamma^{(1)\text{cr}}, \quad (6)$$

$$-\int_{\Gamma^{(2)\text{cr}}} \tilde{p}_i^{k,(2)}(\mathbf{y}) U_{ij}^{(2)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y} = -\frac{1}{2} u_i^{k,(2)}(\mathbf{x}) - \int_{\Gamma^{(2)\text{cr}}} u_i^{k,(2)}(\mathbf{y}) W_{ij}^{(2)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y} - \int_{\Gamma^*} u_i^{k,*}(\mathbf{y}) W_{ij}^{(2)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y} + \int_{\Gamma^*} p_i^{k,*}(\mathbf{y}) U_{ij}^{(2)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y}, \quad \mathbf{x} \in \Gamma^{(2)\text{cr}}, \quad (7)$$

$$-\int_{\Gamma^{(1)\text{cr}}} \tilde{p}_i^{k,(1)}(\mathbf{y}) U_{ij}^{(1)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y} = -\frac{1}{2} u_i^{k,*}(\mathbf{x}) - \int_{\Gamma^{(1)\text{cr}}} u_i^{k,(1)}(\mathbf{y}) W_{ij}^{(1)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y} + \int_{\Gamma^*} u_i^{k,*}(\mathbf{y}) W_{ij}^{(1)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y} - \int_{\Gamma^*} p_i^{k,*}(\mathbf{y}) U_{ij}^{(1)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y}, \quad \mathbf{x} \in \Gamma^*, \quad (8)$$

$$-\int_{\Gamma^{(2)\text{cr}}} \tilde{p}_i^{k,(2)}(\mathbf{y}) U_{ij}^{(2)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y} = -\frac{1}{2} u_i^{k,*}(\mathbf{x}) - \int_{\Gamma^{(2)\text{cr}}} u_i^{k,(2)}(\mathbf{y}) W_{ij}^{(2)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y} - \int_{\Gamma^*} u_i^{k,*}(\mathbf{y}) W_{ij}^{(2)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y} + \int_{\Gamma^*} p_i^{k,*}(\mathbf{y}) U_{ij}^{(2)}(\mathbf{x}, \mathbf{y}, \omega_k) d\mathbf{y}, \quad \mathbf{x} \in \Gamma^*, \quad (9)$$

where $\tilde{p}_i^{k,(m)}(\mathbf{x})$, $p_i^{k,*}(\mathbf{x})$, $u_i^{k,(m)}(\mathbf{x})$ and $u_i^{k,*}(\mathbf{x})$ are the Fourier coefficients of the tractions and displacements at the crack's faces and the bonding interface.

Fundamental solutions in the frequency domain $U_{ij}^{(m)}(\mathbf{x}, \mathbf{y}, \omega_k)$ and $W_{ij}^{(m)}(\mathbf{x}, \mathbf{y}, \omega_k)$ have the following form [1, 9, 12, 16]:

$$U_{ij}^{(m)}(\mathbf{x}, \mathbf{y}, \omega_k) = \frac{1}{2\pi\mu^{(m)}} \left(\psi^{(m)} \delta_{ij} - \chi^{(m)} \frac{(y_i - x_i)(y_j - x_j)}{r} \right), \quad (10)$$

$$W_{ij}^{(m)}(\mathbf{x}, \mathbf{y}, \omega_k) = \lambda^{(m)} n_i^{(m)}(\mathbf{y}) \frac{\partial}{\partial y_k} U_{kj}^{(m)}(\mathbf{x}, \mathbf{y}, \omega_k) + \mu^{(m)} n_k^{(m)}(\mathbf{y}) \left[\frac{\partial}{\partial y_k} U_{ij}^{(m)}(\mathbf{x}, \mathbf{y}, \omega_k) + \frac{\partial}{\partial y_i} U_{kj}^{(m)}(\mathbf{x}, \mathbf{y}, \omega_k) \right]. \quad (11)$$

Here δ_{ij} is the Kronecker delta, $\lambda^{(m)}$ and $\mu^{(m)}$ are the Lamé coefficients, $r = |x_1 - y_1|$ is the distance between the observation and load points. Functions $\psi^{(m)}$ and $\chi^{(m)}$ for the harmonic loading in 2D case are given as:

$$\psi^{(m)} = K_0(l_2^{(m)}) + \frac{1}{l_2^{(m)}} \left[K_1(l_2^{(m)}) - \frac{c_2^{(m)}}{c_1^{(m)}} K_1(l_1^{(m)}) \right], \quad \chi^{(m)} = K_2(l_2^{(m)}) - \left(\frac{c_2^{(m)}}{c_1^{(m)}} \right)^2 K_2(l_1^{(m)}), \quad (12)$$

where $l_1^{(m)} = i\omega_k r/c_1^{(m)}$, $l_2^{(m)} = i\omega_k r/c_2^{(m)}$; $K_n(\bullet)$ is the modified Bessel function of the second kind and order n ; and $c_1^{(m)} = \sqrt{(\lambda^{(m)} + 2\mu^{(m)})/\rho^{(m)}}$ and $c_2^{(m)} = \sqrt{\mu^{(m)}/\rho^{(m)}}$ are the velocities of the longitudinal and transversal waves in the upper and lower half-spaces.

The appropriate system of linear algebraic equations (similar to the one presented in [7]) can be obtained from the system of boundary integral equations (6)–(9) and solved numerically using the iterative algorithm. At the first step the monoharmonic solution of the problem neglecting the effects of the crack closure and friction is obtained. Then the correction of the solution is performed applying the contact constraints (2)–(4) and the Fourier coefficients are gradually changed until the cyclic solution satisfying the constraints is found. The detailed analysis of the algorithm convergence for homogeneous and layered materials can be found in [18, 24]. Note also that the appearing divergent integrals that depend on the type and order of the used space approximation shall be regularised and calculated. In the current study the simplest piecewise-constant approximation was used (as it successfully proved its efficiency for 2D problems comparing, for example, with the Galerkin method [19]).

The detailed investigation of the solution's convergence with respect to the number of Fourier coefficients used in series (5), the number of the boundary elements (and size of the elements in the vicinity of the crack's tips) and the number of time intervals has been presented in [18] and [24] for linear cracks in homogeneous and layered materials under tension-compression loading, thus in this study we followed the recommendations given in [24].

3 NUMERICAL RESULTS

For the validation of the numerical model the crack of the length $2L$ under the normal shear loading of amplitude σ_0 and the frequency closed to zero (with the normalised wave number $k_2^{(2)}L = \omega L/c_2^{(2)} = 0.01$) was considered. The mechanical properties of the material ($\nu^{(1)} = 0.1$, $E^{(1)} = 29\text{GPa}$, and $\nu^{(2)} = 0.49$, $E^{(2)} = 400\text{GPa}$) were chosen to satisfy the following ratio [25]:

$$\beta = \frac{\mu^{(2)}(\kappa^{(1)} - 1) - \mu^{(1)}(\kappa^{(2)} - 1)}{\mu^{(2)}(\kappa^{(1)} + 1) + \mu^{(1)}(\kappa^{(2)} + 1)} = 0.5, \quad (13)$$

where $E^{(m)}$ is the Young's modulus, $\kappa^{(m)} = 3 - 4\nu^{(m)}$, and $\nu^{(m)}$ is the Poisson's coefficient.

The distributions of the normal components of the displacements, $2\mu_0 u_n/\sigma_0 L$, at the crack faces and the bonding interface at a half-period of oscillation are given in Figures 1–4 for different numbers of the correction steps in the absence of the friction ($k_\tau = 0.0$), note that [9]:

$$\mu_0 = \mu^{(1)} \frac{1 - \gamma_2}{1 + \kappa^{(1)}}, \quad \gamma_2 = \frac{a_1}{2} - a_2, \quad a_1 = \frac{\mu^{(1)} - \mu^{(2)}}{\mu^{(1)} + \kappa^{(1)}\mu^{(2)}}, \quad a_2 = \frac{\kappa^{(1)}\mu^{(2)} - \kappa^{(2)}\mu^{(1)}}{2(\mu^{(2)} + \kappa^{(2)}\mu^{(1)})}. \quad (14)$$

Please note that after the correction (see Figure 4) the Signorini constraints (2) are satisfied on the entire surface of the crack, and the Sommerfeld conditions are satisfied at the infinity, so the iterative process effectively corrected the solution.

The normal components of the displacement jump and contact forces at the crack surface are presented in Figures 5 and 6 for different values of the friction coefficient. The contact forces

and the size of the contact zone are compared with the model static solution by Comninou and Dundurs [25]. As one can see the results are in a very good agreement. It should be noted that the friction significantly affected the solution, changing the size of the contact zone and the distribution of the displacements and forces (even for the “quasi-static” case considered for the validation purposes).

Furthermore, increase in the frequency of the loading changes the solution even more (for comparison, please see [26]). As an example, the time distributions of the normal and tangential components of crack faces displacements are given in Figures 7–10 for $k_2^{(2)}L = 0.1$, $k_\tau = 1.0$. Note the significant difference in the displacement amplitudes caused by the shear loading.

4 CONCLUSIONS

Computing and full parametric analysis of the normal opening and shear modes of the stress intensity factors for different friction coefficients, will be the next stage of this study. Please also note that the suggested boundary integral approach has been successfully applied for two-dimensional impact contact problems for homogeneous and layered materials [27–29]. In the future it can be extended to three-dimensional fracture mechanics problems for cracked materials under arbitrary dynamic loading, and the special attention shall be paid to the coupling oscillation singularities in the vicinity of the interface crack front [1–3, 13, 23].

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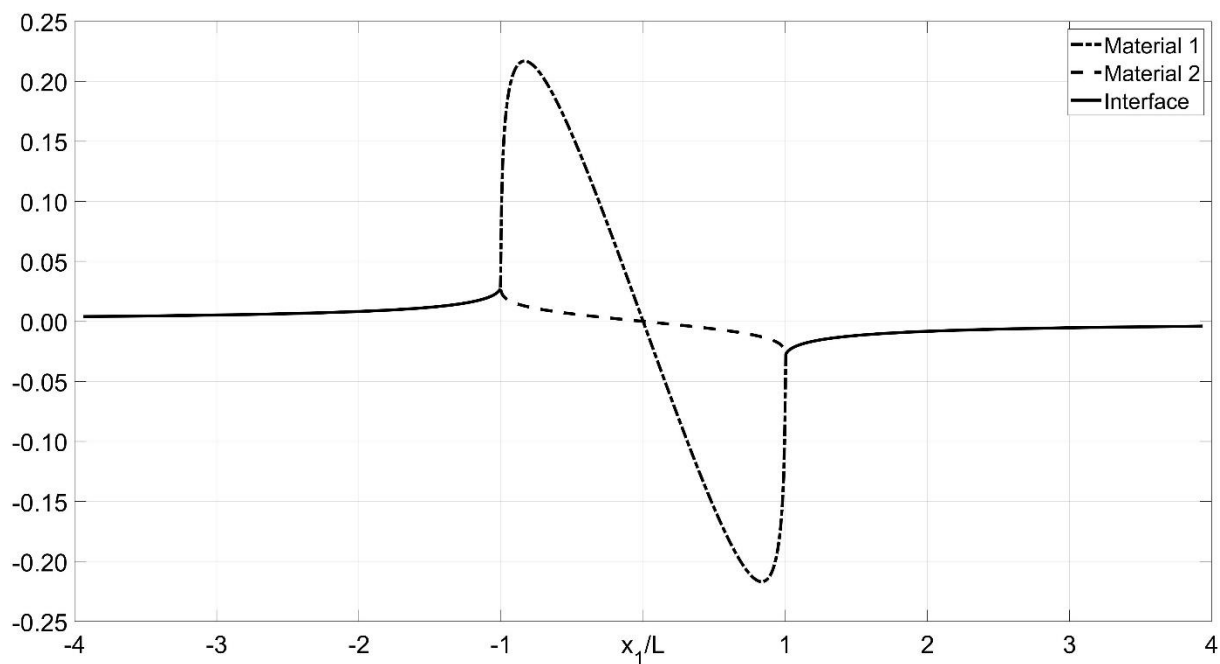


Figure 1: Normal displacements, $2\mu_0 u_n / \sigma_0 L$, before iterative corrections

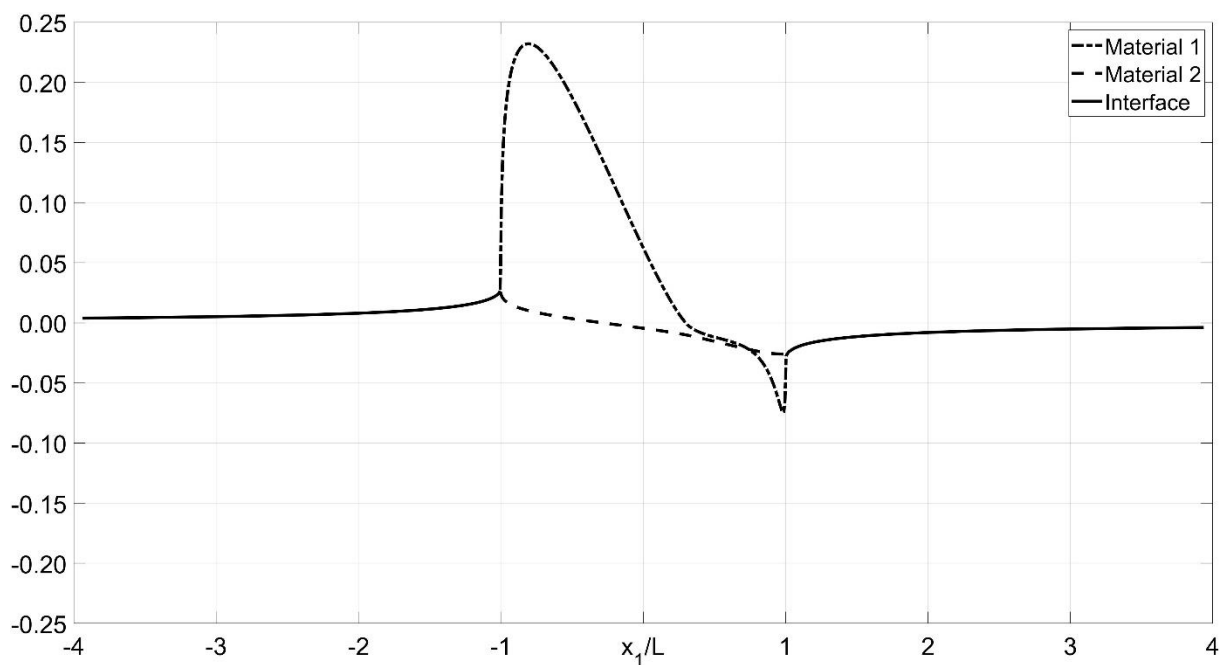


Figure 2: Normal displacements, $2\mu_0 u_n / \sigma_0 L$, after 25 iterations

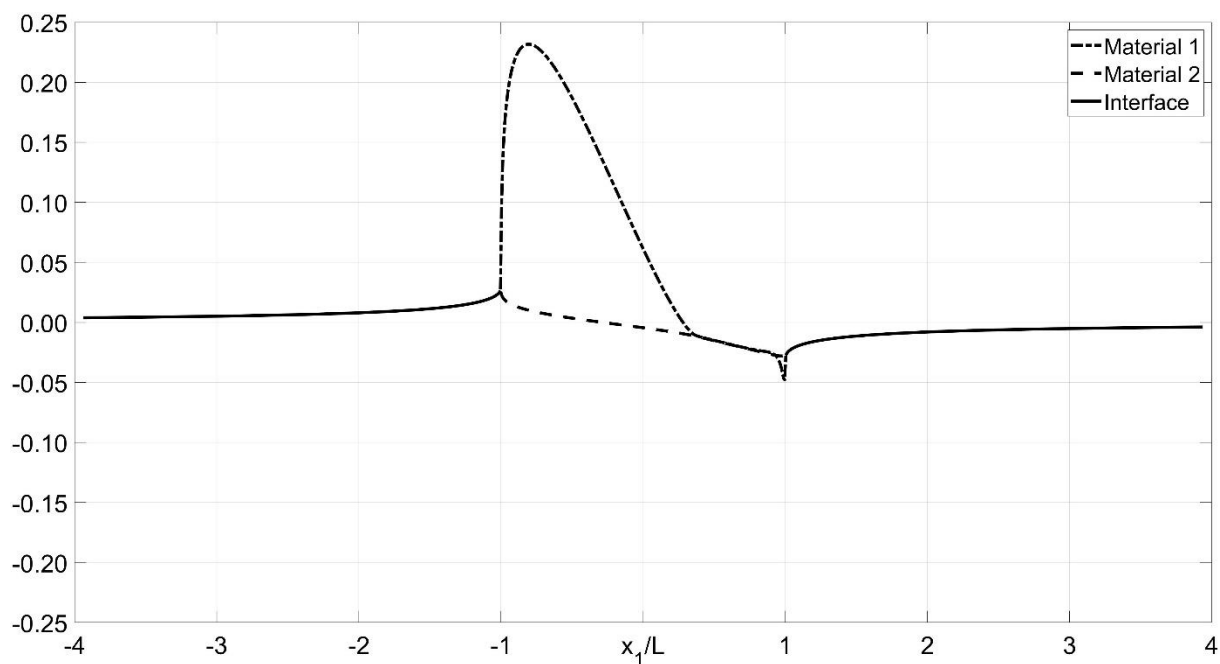


Figure 3: Normal displacements, $2\mu_0 u_n / \sigma_0 L$, after 100 iterations

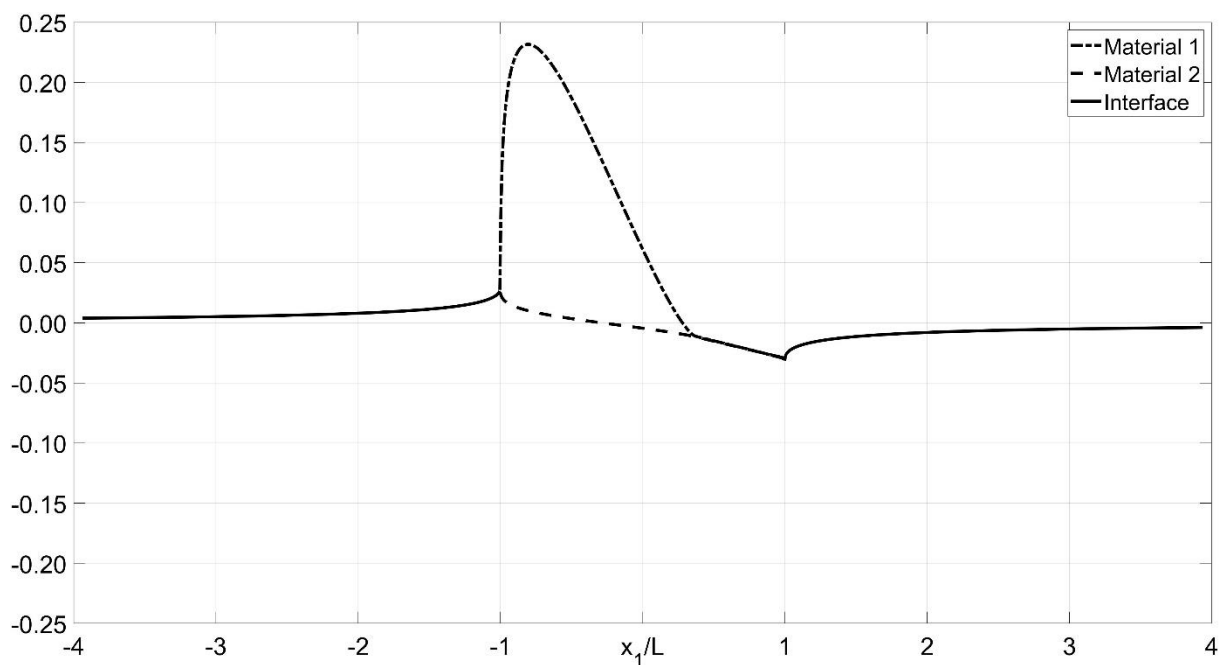


Figure 4: Normal displacements, $2\mu_0 u_n / \sigma_0 L$, after 1000 iterations (correction is completed)

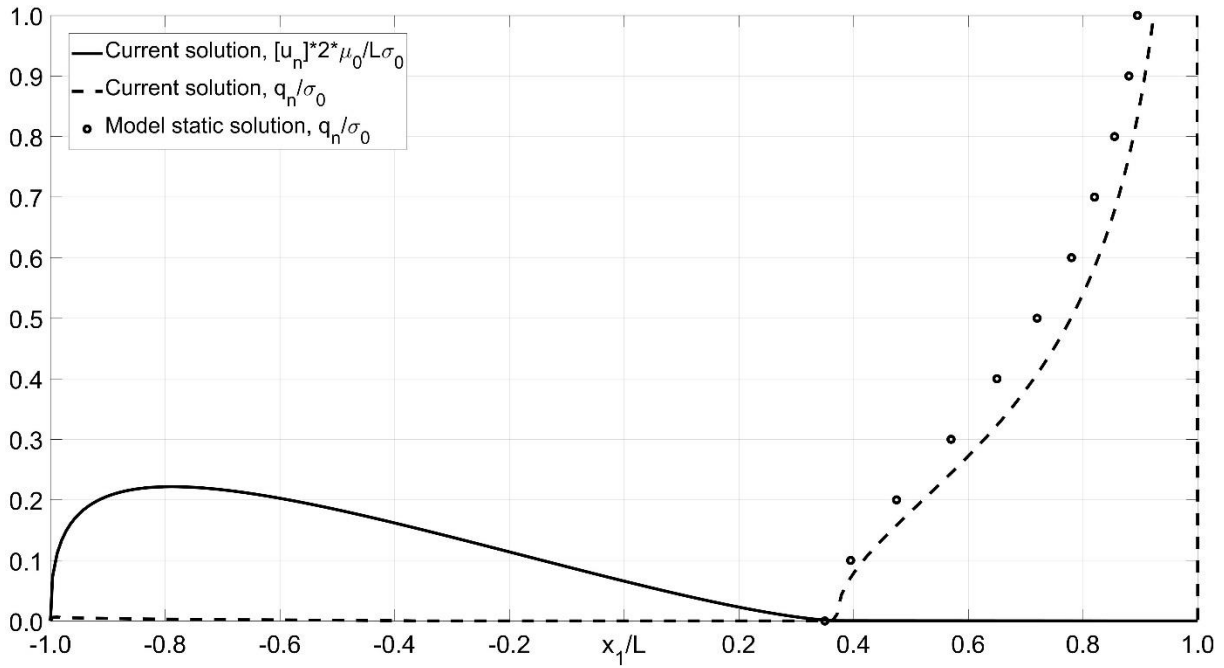


Figure 5: Normal components of the displacement jump and contact forces without friction ($k_\tau = 0.0$)

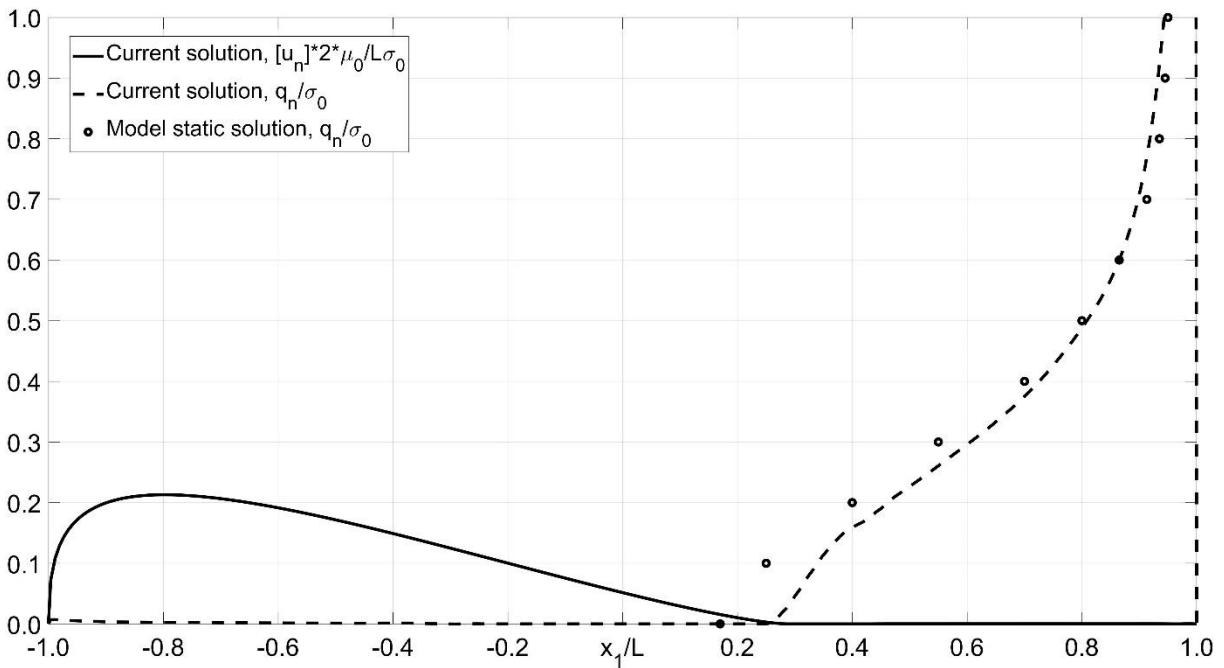


Figure 6: Normal components of the displacement jump and contact forces with friction ($k_\tau = 1.0$)

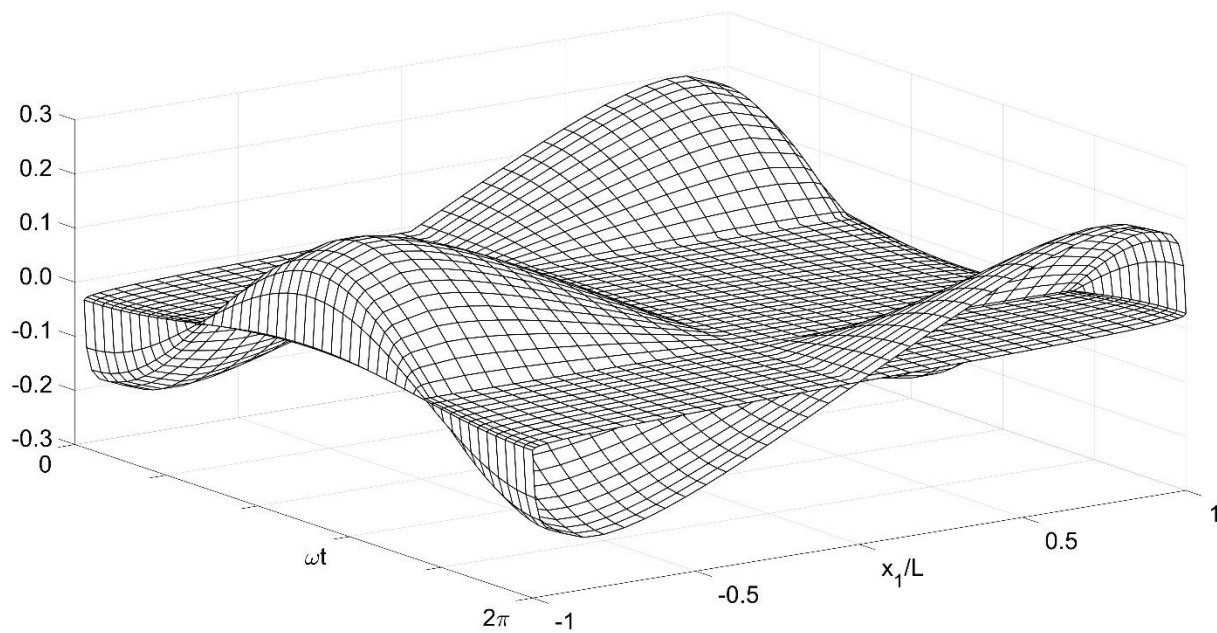


Figure 7: Normal components of the crack faces displacements before correction, $2\mu_0 u_n / \sigma_0 L$

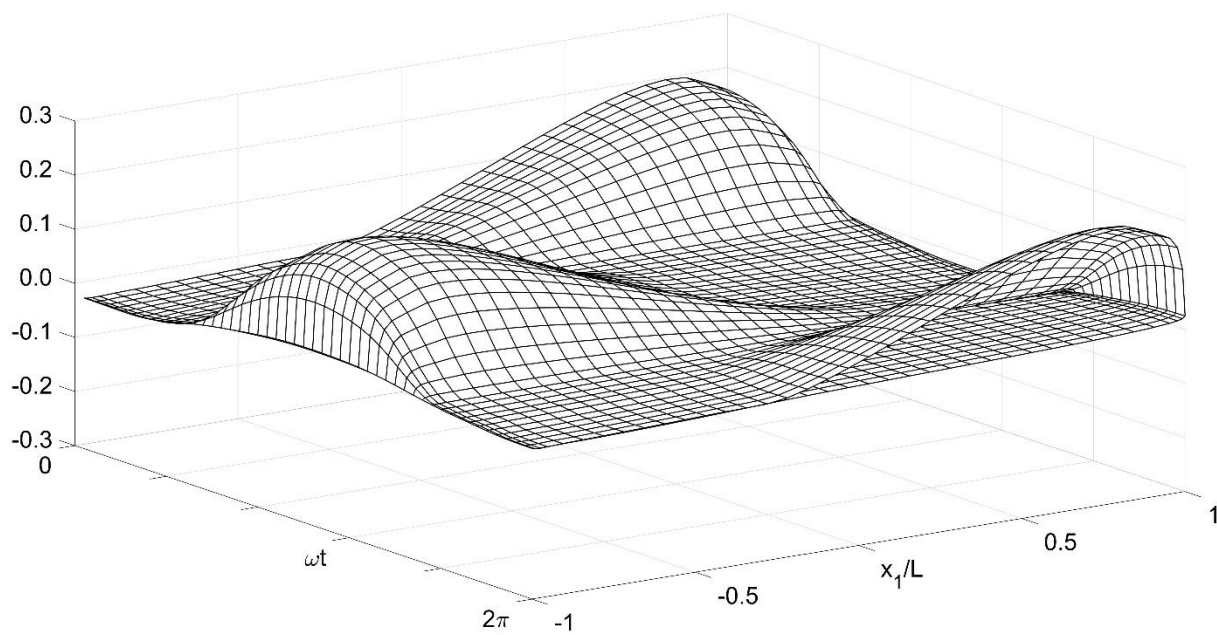


Figure 8: Normal components of the crack faces displacements after correction, $2\mu_0 u_n / \sigma_0 L$

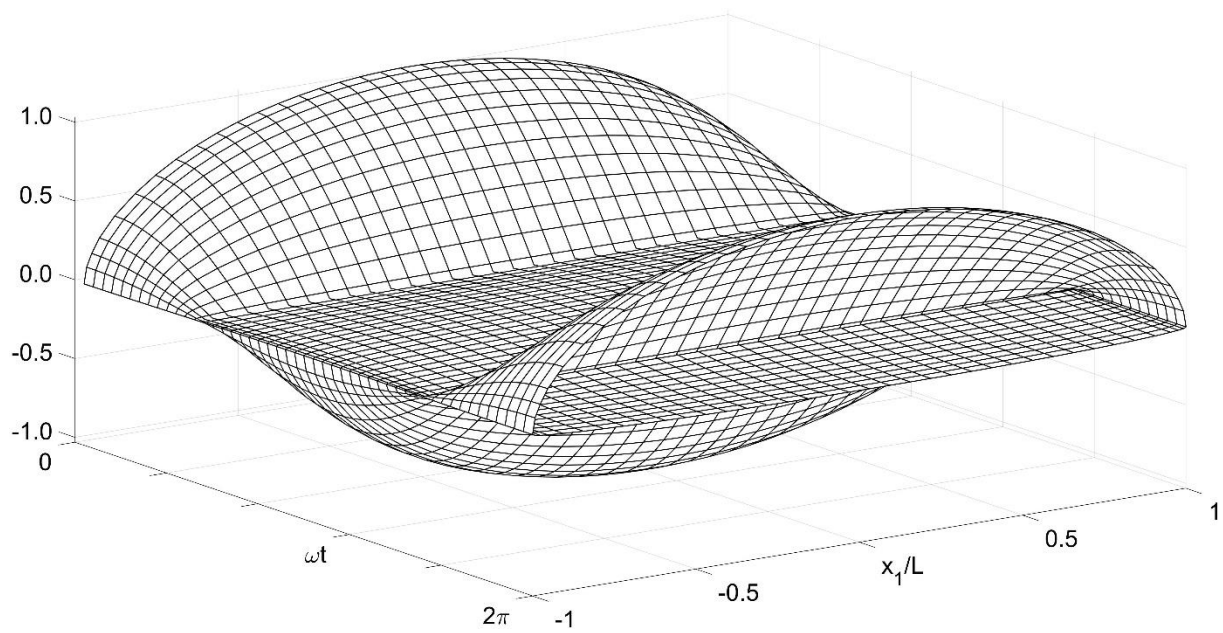


Figure 9: Tangential components of the crack faces displacements before correction, $2\mu_0 u_\tau / \sigma_0 L$

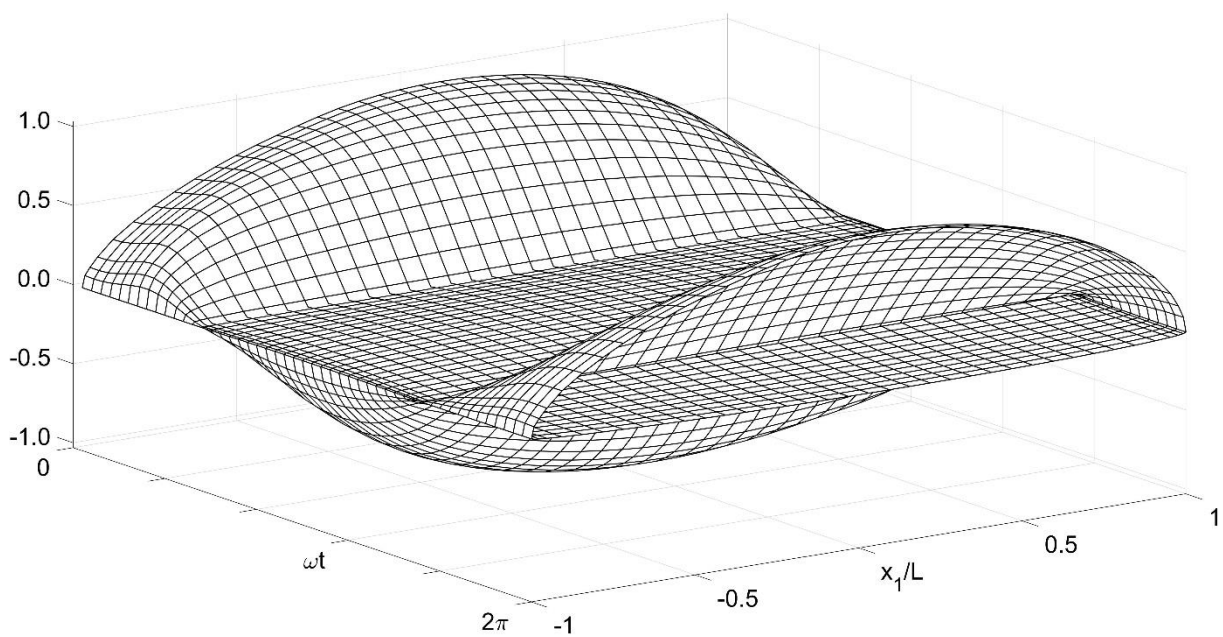


Figure 10: Tangential components of the crack faces displacements after correction, $2\mu_0 u_\tau / \sigma_0 L$