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LETTER

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# Increased radon-222 in soil gas because of cumulative seismicity at active faults

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## Abstract

This study demonstrates how the radon-222 ( $^{222}\text{Rn}$ ) concentration of soil gas at an active fault is sensitive to cumulative recent seismicity by examining seven active faults in western Japan. The  $^{222}\text{Rn}$  concentration was found to correlate well with the total earthquake energy within a 100-km radius of each fault. This phenomenon can probably be ascribed to the increase of pore pressure around the source depth of  $^{222}\text{Rn}$  in shallow soil caused by frequently induced strain. This increase in pore pressure can enhance the ascent velocity of  $^{222}\text{Rn}$  carrier gas as governed by Darcy's law. Anomalous  $^{222}\text{Rn}$  concentrations are likely to originate from high gas velocities, rather than increased accumulations of parent nuclides. The high velocities also can yield unusual young gas under the radioactive nonequilibrium condition of short elapsed time since  $^{222}\text{Rn}$  generation. The results suggest that ongoing seismicity in the vicinity of an active fault can cause accumulation of strain in shallow fault soils. Therefore, the  $^{222}\text{Rn}$  concentration is a possible gauge for the degree of strain accumulation.

**Keywords:** Radon-222 concentration; Nonequilibrium condition; Gas velocity; Darcy's law; Total earthquake energy; Western Japan

## Findings

### Introduction

Radon-222 ( $^{222}\text{Rn}$ ) concentrations in soil gas and groundwater have been known to increase abruptly before and/or during large earthquakes in many parts of the world (Hauksson 1981; King 1984/85; Heinicke et al. 1995; Igarashi et al. 1995; Pérez et al. 2007; İnan et al. 2008; Kuo et al. 2010). The sensitivity of  $^{222}\text{Rn}$  concentration has been confirmed by its response to earthquakes a few hundred kilometers away (Fleischer 1981; Toutain and Baubron 1999; Cicerone et al. 2009; Ghosh et al. 2009).  $^{222}\text{Rn}$  is therefore thought to be a promising precursory gauge for large earthquakes. The observed spike-like increases in Rn concentration are similar to groundwater-level changes induced by earthquakes, and seismically enhanced crustal strain and pore pressure in saturated soils and rocks are the most plausible factors leading to this change (Manga et al. 2012).

Another widely recognized characteristic of  $^{222}\text{Rn}$  is that concentrations tend to be high at faults (King 1978;

King et al. 1996; Baubron et al. 2002; Walia et al. 2009; Neri et al. 2011). Therefore,  $^{222}\text{Rn}$  surveys have been used to map faults and, in particular, to detect active faults. However,  $^{222}\text{Rn}$  enhancement does not necessarily occur at all faults, and the mechanism for this enhancement is still disputed. Three possible causes are (i) the magnitude of fault displacement velocities during the late Quaternary, (ii) the accumulation of parent nuclides caused by historical large earthquakes, and (iii) recent seismicity at nearby faults. By undertaking a radioactivity survey at seven active fault areas in western Japan, this study aimed to identify the most rational cause of the  $^{222}\text{Rn}$  concentration enhancement. The results could enable discrimination between faults that are continually active and faults that are occasionally active, i.e., faults experiencing stationary fracturing at short time scales or faults experiencing intermittent large fracturing over long time intervals.

### Fault description

Seven active faults that have experienced historical and recent large earthquakes, the Atotsugawa (AT), Atera (AR), Nojima (NJ), Beppu-Haneyama (BH), Futagawa (FG), Hinagu (HN), and Izumi (IM), provide the data for this study. The hypocenters of earthquakes associated with

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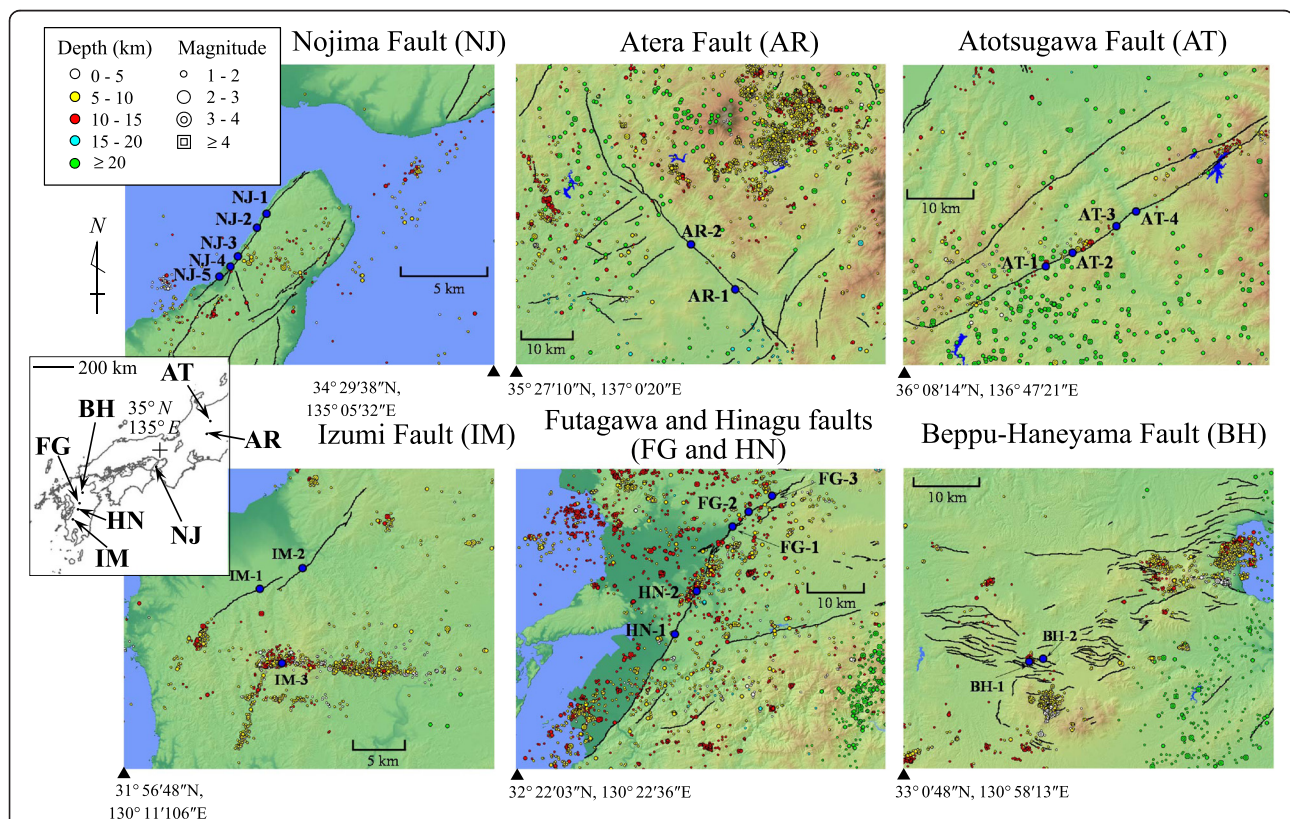
these faults from January 2001 to September 2010 are mostly located in the shallow crust at depths between 5 and 15 km at NJ, BH, FG, HN, and IM and at depths greater than 15 km at AT and AR (Figure 1). The density of earthquake epicenters is irregular and variable, even on the same fault, as can be seen clearly along HN. Earthquakes have been infrequent on AR, NJ, BH, and IM. Activity ranks are defined by the average slip per thousand years: 1 to 10 m for rank A and 0.1 to 1 m for rank B. Only AT and AR are evaluated as the most active rank A (The Research Group for Active Faults of Japan 1991). Table 1 provides a summary of the rank, length, movement pattern, radioactivity survey details, and surface geology for each fault. The lines shown in Figure 1 are perpendicular to fault strike (Figure 2). Surveys were undertaken on days without rain.

The latest large earthquakes and fault movements for the faults studied are well constrained by the evaluation reports of active faults (The Headquarters for Earthquake Research Promotion 2013). An 1858 earthquake (magnitude  $M$  7.0 to 7.1) has been ascribed to movement of a

portion of AT. The latest large earthquake on AR, which occurred in 1586, is thought to have been about  $M$  7.8. A destructive earthquake of  $M$  7.3 occurred in January 1995 on the northern part of NJ at a depth of 16 km. The most recent movements on BH, which is composed of a set of normal sub-faults forming grabens in an active volcanic area, range from 3.9 ka to the sixth century AD in the north and from after the thirteenth century AD in the south. The latest remarkable movements of FG, HN, and IM are estimated from trenching investigations to be in the ranges of 6.9 to 2.4 ka, 8.4 to 2.2 ka, and 7.3 to 2.4 ka, respectively. A large earthquake of  $M$  6.4 in 1997 was located on the south end of IM, where many small earthquakes had occurred along two orthogonal segments (Figure 1). The IM-3 survey line was positioned in this zone.

#### Methods of radioactivity survey

Measurement lines were established near trench investigation sites (FG-1, HN-1, and IM-1 and IM-2), in concentrated epicenter zones (all AT lines, all BH lines, FG-2 and



**Figure 1** Location maps of seven active faults in western Japan and the associated distribution of earthquake epicenters. Fault traces are after a digital fault map (Nakata and Imaizumi 2002). Locations of measurement lines and the epicenters of earthquakes from January 2001 to September 2010 in the vicinity of the faults as recorded in the Seismological and Volcanological Bulletin of Japan (CD-ROM) and based on the observations of the Japan Meteorological Agency are depicted. The beginning of the earthquake data observation period was selected to comport with the commencement of Hi-net (the High Sensitivity Seismograph Network Japan, National Research Institute for Earth Science and Disaster Prevention), October 2000, for hypocenter determinations in the study areas. The color and shape of the epicenter indicators correspond to hypocenter depth and earthquake magnitude, respectively.

**Table 1 Details of the active faults studied**

Fault name	Atotsugaw (AT)	Atera (AR)	Nojima (NJ)	Beppu-Haneyama (BH)	Futagawa (FG)	Hinagu (HN)	Izumi (IM)
Rank of activity	A	A	B	B	B	B	B
Length (km)	70	66	15	30 to 40	29	70	20
Movement pattern	Right-lateral slip	Left-lateral slip	Right-lateral slip	A set of E-W normal sub-faults	Right-lateral slip	Right-lateral slip	Normal fault with right-lateral slip
Number of lines	4	2	5	2	3	2	3
Number of survey points at each line	AT-1 (9), AT-2 (5), AT-3 (5), AT-4 (4)	AR-1 (15), AR-2 (13)	NJ-1 (26), NJ-2 (12), NJ-3 (7), NJ-4 (4), NJ-5 (4)	BH-1 (9), BH-2(8)	FG-1 (10), FG-2 (4), FG-3 (7)	HN-1 (25), HN-2 (33)	IM-1 (9), IM-2 (10), IM-3 (11)
Total survey points	23	28	53	17	21	58	30
Number of soil samples at each line	AT-2 (2)	AR-1 (10), AR-2 (8)	NJ-1 (4), NJ-2 (3), NJ-3 (2), NJ-4 (1), NJ-5 (2)	BH-1 (9)	FG-1 (10), FG-2 (4), FG-3 (7)	HN-1 (11), HN-2 (33)	IM-1 (8), IM-2 (3), IM-3 (5)
Total soil samples	2	18	12	9	21	44	16
Survey period	Nov. to Dec. 2008	Dec. 2006	Dec. 2007	Dec. 2009	Oct. to Dec. 2004	Oct. 2005, Jan. and Oct. 2006, Oct. 2009	Dec. 2010
Surface geology	M	S	S	V	S (FG-2), V	M (HN-1), V (HN-2)	G (IM-3), S
Number of nonequilibrium points (%)	6 (26)	5 (18)	0 (0)	6 (35)	6 (29)	17 (29)	8 (27)

The following details are included: the rank of fault activity (The Research Group for Active Faults of Japan 1991); fault length; fault movement pattern (The Headquarters for Earthquake Research Promotion 2013); radioactivity survey specifications (number of measurement lines, line code, number of survey points, number of soil samples on each line, and survey period); surface geology classified into granitic (G), volcanic (V), sedimentary (S), and metamorphic (M) rock types based on a digital geological map (Geological Survey of Japan 2013); and the number of nonequilibrium points with the percentage of total survey points.

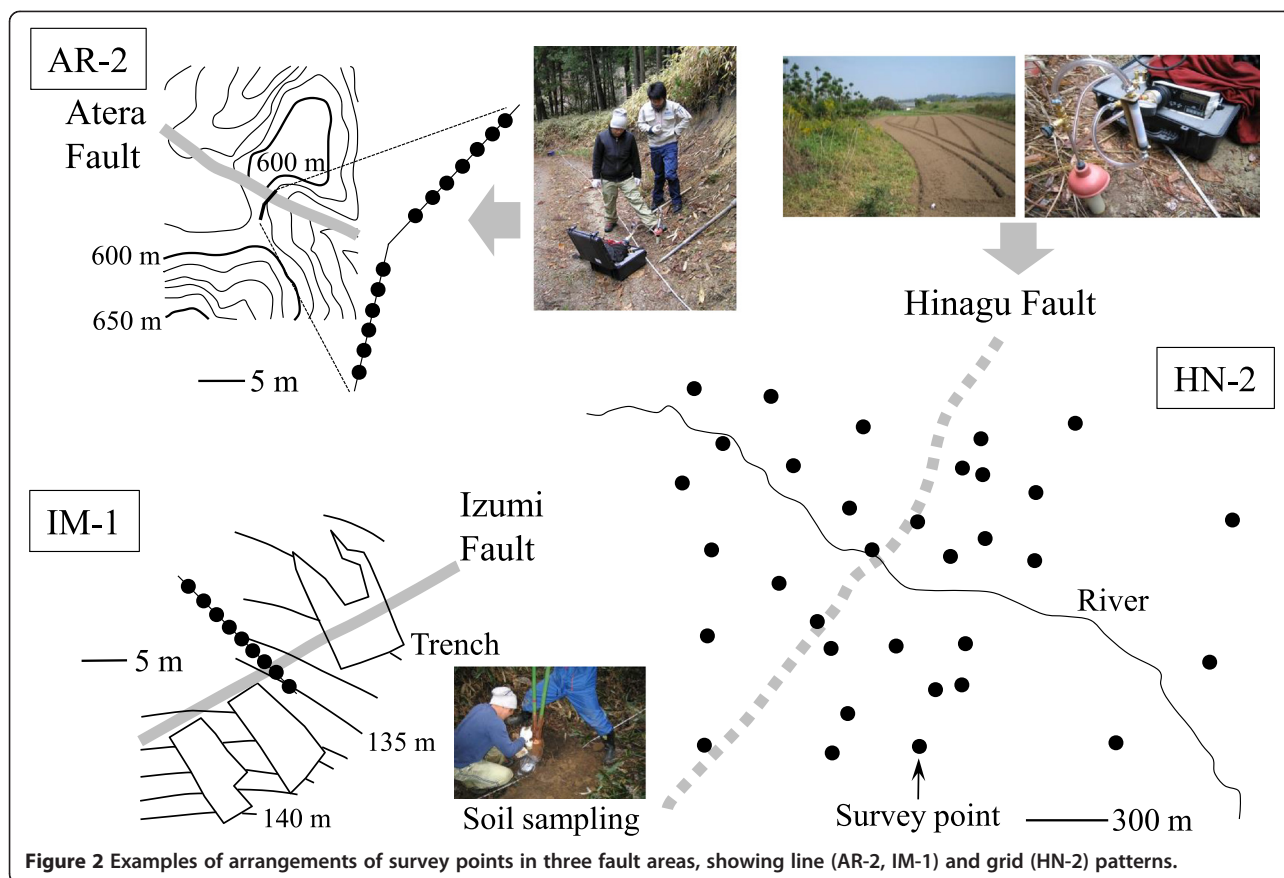
FG-3, HN-2, and IM-3), near fault outcrops (all AR lines), and near the surface displaced by the latest large earthquake (all NJ lines). Survey points were placed at 1- to 10-m intervals along a line or at 100- or 500-m intervals on a grid pattern applied to HN-2 (Figure 2). Because it was difficult to identify the position of the Hinagu Fault in the flat terraced landscape of the HN-2 area, the epicenters of  $M \geq 2$  earthquakes were used to establish the survey points there. At each survey point, soil gas was sampled from a hole with a depth of 60 cm and a diameter of 3 cm. The total number of  $\alpha$  particles per minute (cpm) resulting from the decay of  $^{222}\text{Rn}$  and  $^{220}\text{Rn}$  and their daughter nuclides,  $^{218}\text{Po}$  and  $^{216}\text{Po}$ , were counted successively for 10 to 20 min using a portable  $\alpha$ -scintillation detector, RDA-200 (Scintrex, Vaughan, Canada) or AB-5 (Pylon Electrics, Ottawa, Canada). The long count time was used to ascertain whether the gas was at nonequilibrium or equilibrium when judgment was difficult.

To estimate the true  $^{222}\text{Rn}$  concentration from the total cpm, the CRAS method (calculation method of numbers of Rn isotope atoms and their daughters at the start of measurement: Koike et al. 2009), which considers the radioactive equilibrium relationship between  $^{222}\text{Rn}$  and  $^{218}\text{Po}$  in soil gas, was used. Numbers of  $^{222}\text{Rn}$  ( $N_{222}$ ),

$^{220}\text{Rn}$ ,  $^{218}\text{Po}$ , and  $^{216}\text{Po}$  were determined. Only  $N_{222}$  was used to characterize the faults studied because the half-life of  $^{222}\text{Rn}$  (3.823 days) is much longer than that of the other isotopes. One noteworthy feature is that the soil gases at 21% of the points exhibited nonequilibrium conditions, i.e., the cpm increased with time due to an increase in  $^{218}\text{Po}$  decays (Figure 3). This type of gas is young (i.e., not much time has elapsed after the generation of  $^{222}\text{Rn}$  atoms) and observed typically in fumarole gases with high ascent velocities (Koike et al. 2009, 2014). Another feature is that the percentages of nonequilibrium points are similar for each fault except NJ (Table 1). The other gases exhibited equilibrium conditions in which the cpm mostly decreased in the first few minutes due to a decrease in decays of  $^{220}\text{Rn}$  (half-life 53 s; Figure 3), while the rates of decay of  $^{222}\text{Rn}$  and  $^{218}\text{Po}$  remained constant. Usual soil gases show this cpm pattern. The predicted total decays agree well with the measurement data for both conditions (Figure 3), demonstrating the validity of the CRAS approach.

Soil samples were taken from the bottom of the borehole after the soil gas survey, dried at 80°C for 24 h, and arranged for samples with particle sizes smaller than 0.25 mm and mass of 150 g. A  $\gamma$ -ray spectrometer using





a Ge semiconductor detector (GMX-25190-P: EG&G ORTEC, Oak Ridge, USA) was used to measure the  $\gamma$ -ray intensity (counts) of the 12  $\gamma$ -decay nuclides in the uranium and thorium series. The concentration of nuclides was assumed to be proportional to the  $\gamma$ -ray intensity. Only the  $\gamma$ -ray intensity of  $^{226}\text{Ra}$  ( $I_{226}$  122 data in total) was used in conjunction with  $N_{222}$  (230 data in total) because of the isotope's long half-life (1,601 years) and because it is a direct parent nuclide of  $^{222}\text{Rn}$ . We used the conventional definition of  $I_{226}$  in becquerels per gram from the original  $\gamma$ -ray channel data in counts by the following transformation. First, the net spectral peak count of  $^{226}\text{Ra}$  per second was measured using Covell's method (Covell 1959), and then it was divided by three factors: the detector efficiency of  $^{226}\text{Ra}$  (0.065%), sample weight (g), and  $\gamma$ -ray yield from a standard source.

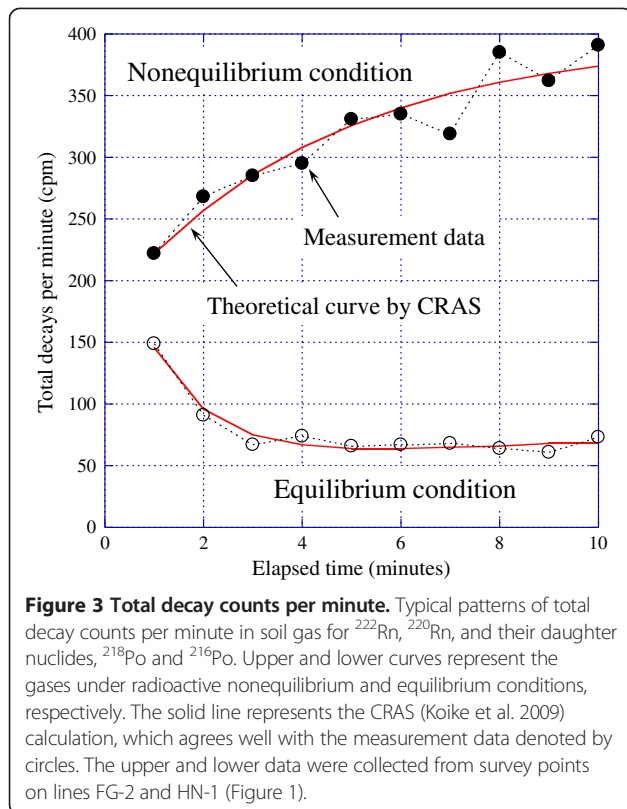
## Results and discussion

### Identifying the controlling factor for $^{222}\text{Rn}$ concentration

Because gas ascent velocity generally is slow in the study area, the  $^{222}\text{Rn}$  source must be located in shallow soils. The average length by molecular diffusion of  $^{222}\text{Rn}$  in soils,  $L$  (m), is given by a function of the diffusion coefficient,  $D$  ( $\text{m}^2/\text{s}$ ), and decay constant,  $\lambda_{222}$  ( $\text{s}^{-1}$ ), as

$\sqrt{D/\lambda_{222}}$  (Lehmann et al. 2000). The assignment of a general value for dry soil,  $D = 5.0 \times 10^{-6} \text{ m}^2/\text{s}$  (Megumi and Mamuro 1973; Holford et al. 1993), derives  $L = 1.6$  m. First, the values of  $N_{222}$  and  $I_{226}$  were analyzed from their relationship with the surface geology, which was classified roughly into granitic (G), volcanic (V), sedimentary (S), and metamorphic (M) rock types using the seamless digital geological map of Japan (Geological Survey of Japan 2013). The soils originated from weathering of these rocks. Summary statistics of the  $N_{222}$  and  $I_{226}$  data shown with box plots (Figure 4a,b) reveal that the magnitude orders by rock type of the upper quartile and median are the same ( $G > V > S > M$ ) in both cases. This implies that surface geology has some effect on the radioactive nuclide concentrations. However, focusing on the median values, the differences in the values are slight for both  $N_{222}$  and  $I_{226}$ . Therefore, some other stronger factors must cause the large concentration variability.

A scattergram of the  $N_{222}$  and  $I_{226}$  data at the same points highlights the absence of a clear correlation (Figure 5: correlation coefficient  $R = 0.17$ ). Such an absence also has been observed in other case studies (Ball et al. 1991; Cecil et al. 1991; Ioannides et al. 2003; Appleton et al. 2011). This feature is probably the result



of the difference in state, i.e., gas ( $^{222}\text{Rn}$ ) and solid ( $^{226}\text{Ra}$ ). Accordingly, it is most pertinent to interpret that high  $^{222}\text{Rn}$  concentrations do not originate from accumulated parent nuclides in soils but rather from the high ascent velocity of the carrier gas ( $\nu$ ). This interpretation is evidenced by the simple kinetic transport of  $^{222}\text{Rn}$  in the vertical direction in shallow soils. The transport can be expressed by the following mass balance equation,

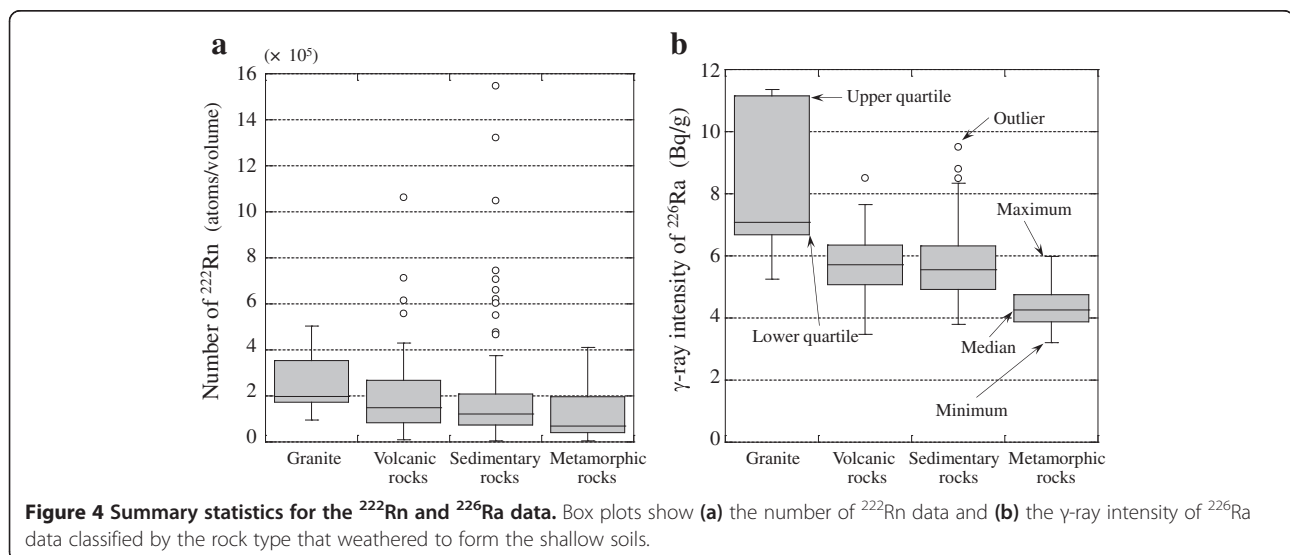
which combines molecular diffusion and advection of the carrier gas flow (Kristiansson and Malmqvist 1982; Ioannides et al. 2003):

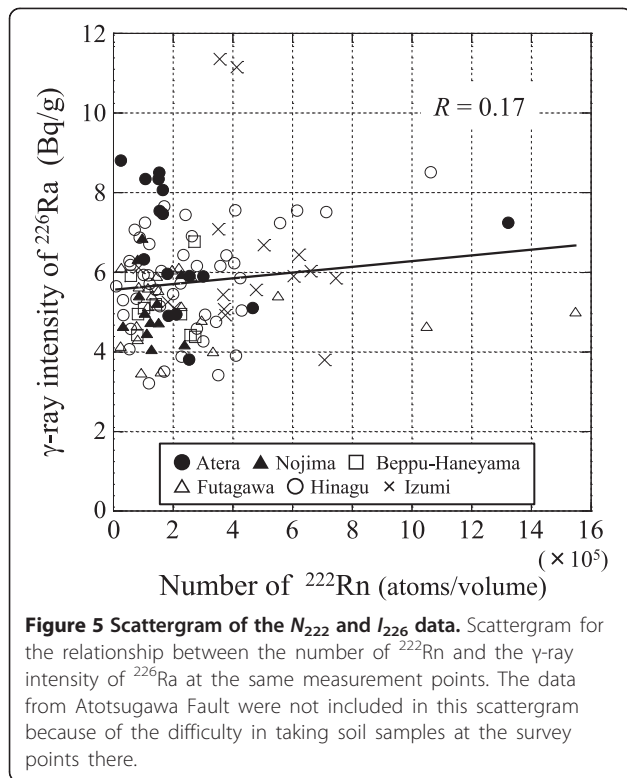
$$\frac{\partial C_{222}(z)}{\partial t} = \phi - \lambda_{222}C(z) + \frac{D}{\varepsilon} \frac{\partial^2 C_{222}(z)}{\partial z^2} - \nu \frac{\partial C_{222}(z)}{\partial z} \quad (1)$$

where  $C_{222}(z)$  is the  $^{222}\text{Rn}$  concentration at depth  $z$  (positive downward from the surface),  $\phi$  is the  $^{222}\text{Rn}$  production rate per unit time, and  $\varepsilon$  is the soil porosity. Under the steady state condition  $\partial C_{222}(z)/\partial t = 0$  and the  $^{222}\text{Rn}$  concentration at the ground surface  $C_{222}(z=0) = 0$ , a solution for the equation is the following (Kristiansson and Malmqvist 1982; Ioannides et al. 2003).

$$C_{222}(z) = \frac{\phi}{\lambda_{222}} \left\{ 1 - \exp \left[ - \left( \sqrt{\left( \frac{\nu \varepsilon}{2D} \right)^2 + \frac{\lambda_{222} \varepsilon}{D}} + \frac{\nu \varepsilon}{2D} \right) z \right] \right\} \quad (2)$$

If we assume the equilibrium condition between  $^{222}\text{Rn}$  and  $^{226}\text{Ra}$ , then  $\lambda_{222}C_{222}(z)$  should be equal to  $\lambda_{226}C_{226}(z)$ . It follows that  $\phi = \lambda_{226}C_{226}(z)$ , where  $C_{226}(z)$  and  $\lambda_{226}$  are the  $^{226}\text{Ra}$  concentration at  $z$  and the decay constant of  $^{226}\text{Ra}$ , respectively.  $C_{222}(z)$  at  $z = 60$  cm (the bottom depth of the borehole) is calculated by assigning the above value  $D = 5 \times 10^{-6}$  m<sup>2</sup>/s and a fixed  $\varepsilon$  of 0.3. In normal soils,  $\nu$  is smaller than  $1 \times 10^{-5}$  m/s (Kristiansson and Malmqvist 1982; Schery and Siegel 1986). The value  $1 \times 10^{-3}$  m/s is its upper limit according to the study of Brown (2000). Using two values for high and low velocity, the ratio of  $C_{222}(z)$  at  $\nu = 5 \times 10^{-5}$  and  $5 \times 10^{-6}$  (m/s) is 3.1 for the same  $C_{226}(z)$ . This means that the  $^{222}\text{Rn}$  concentration is highly variable by a factor of





**Figure 5 Scattergram of the  $N_{222}$  and  $I_{226}$  data.** Scattergram for the relationship between the number of  $^{222}\text{Rn}$  and the  $\gamma$ -ray intensity of  $^{226}\text{Ra}$  at the same measurement points. The data from Atotsugawa Fault were not included in this scattergram because of the difficulty in taking soil samples at the survey points there.

3 in this case even though the  $^{226}\text{Ra}$  concentration is constant. For a pair of velocities with less contrast,  $v = 5 \times 10^{-5}$  and  $1 \times 10^{-5}$  (m/s), the ratio is 2.3. In loose deposits,  $D$  ranges from  $1.5 \times 10^{-6}$  to  $6.8 \times 10^{-6}$   $\text{m}^2/\text{s}$  (Gasparini and Mantovani 1978). When  $D$  is halved ( $2.5 \times 10^{-6}$   $\text{m}^2/\text{s}$ ), the ratios for the above two  $v$  pairs become 2.4 and 1.7. Consequently, the absence of a correlation between  $^{222}\text{Rn}$  and  $^{226}\text{Ra}$  can be attributed to the changeability in gas ascent velocity in the soils, not to the heterogeneous distribution of  $^{226}\text{Ra}$  concentration.

#### Effect of earthquake energy and mechanism of $^{222}\text{Rn}$ enhancement

Next, we evaluate the effect of earthquake energy,  $E(J)$ , on  $N_{222}$ .  $E(J)$  is related to  $M$  by the Gutenberg-Richter law,  $\log E = 1.5M + 4.8$  or  $E = 10^{1.5M + 4.8}$ . Considering the attenuation of  $E$  with distance from the hypocenter,  $d$ , the  $E$  at  $d$  per unit area was set to be proportional to the inverse of  $d^2$ , based on the spherical divergence of a seismic wave propagating in a homogeneous medium. In each active fault area, the distance from a measurement point location to a hypocenter was defined as  $d$ , and all  $E/d^2$  values were summed. This summation is termed the total earthquake energy (TEE).

Among the earthquakes shown in Figure 1, those occurring within a circle of a given radius were selected to

clarify the relationship between TEE and radius. Therefore, TEE is expressed as

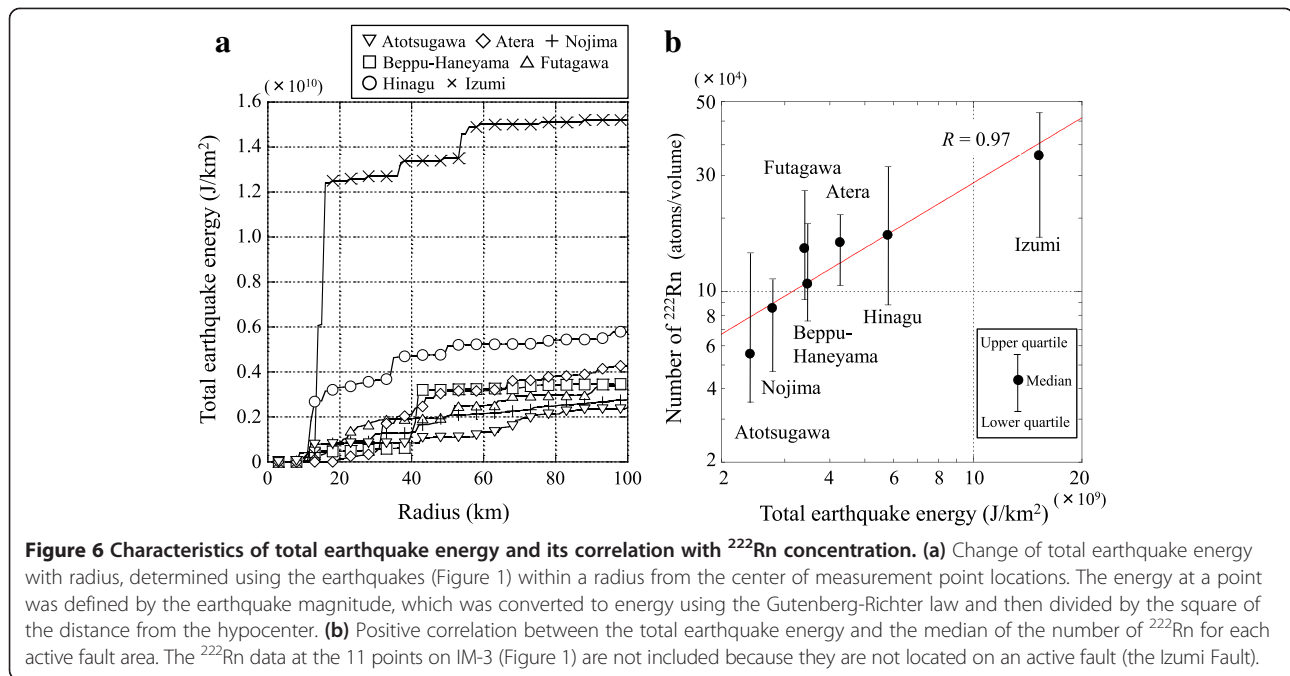
$$\text{TEE} = \sum_{i=1}^n \frac{E_i}{d_i^2} \quad (3)$$

where  $n$  is the total number of earthquakes from January 2001 to September 2010 with  $d \leq$  the radius. TEE essentially increases as the radius increases, but the TEE values of all faults converge at a radius of approximately 100 km (Figure 6a). The  $^{222}\text{Rn}$  concentration anomaly can occur at an epicentral distance of 100 km for an approximately  $M$  4.7 earthquake and a crustal strain of  $10^{-8}$  from an empirical relationship derived by using worldwide earthquake data (Fleischer 1981; Toutain and Baubron 1999). Thus, TEE for each fault area was defined by this radius. Because the  $N_{222}$  values are dispersed widely, which is common to all faults, their median value was used for characterizing TEE. The median values have an obviously strong positive correlation with TEE ( $R = 0.97$ ), as shown by the regression line in Figure 6b. It is proved, therefore, that the  $^{222}\text{Rn}$  concentration in an active fault area is affected by the summation of recent earthquakes within a 100-km radius. Our findings show that for an active fault with historically large earthquakes but no recent seismicity, the  $^{222}\text{Rn}$  concentration will tend to be low, as evidenced by the case of AT.

The effect of the length of time on TEE was examined by fixing the radius at 100 km and using 10 periods from 2 to 20 years with 2-year intervals. The  $R$  values between TEE and the median of the  $N_{222}$  values for the seven faults were calculated for each time period. As shown in Figure 7, the values approximately show a convex curve:  $R$  increases as the period becomes longer, reaches a peak at the 10-year period, and then decreases toward the 20-year period. Thus, the TEE analysis indicates that the 10-year period from January 2001 to September 2010 is the most suitable.

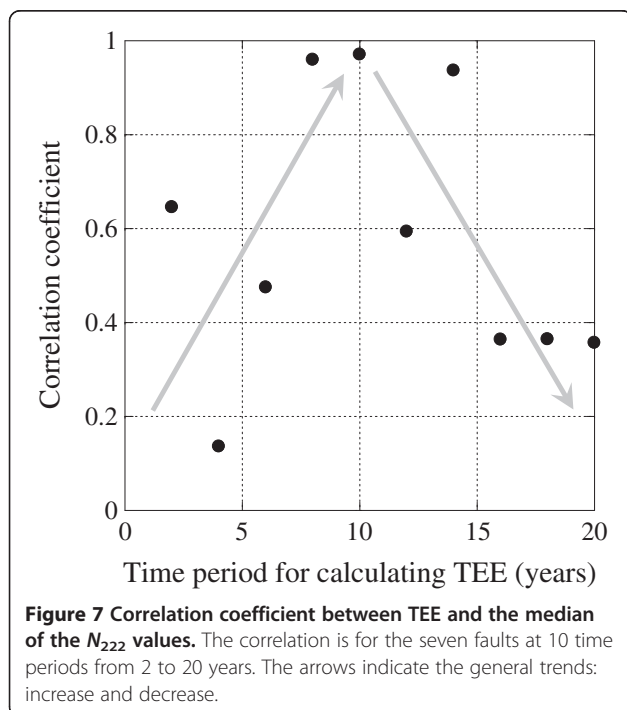
The mechanism causing the strong correlation described above can be interpreted by a simplified dislocation model (Fleischer 1981). This model assumes the Earth to be an infinite isotropic elastic solid containing a circular dislocation loop of radius  $r$  and slip vector  $b$ . As long as the ratio  $r/b$  is constant, the shear strain change  $\Delta\gamma$  at  $d$  can be related to  $M$  as  $\Delta\gamma \propto 10^{1.44M}/d^3$ . For the same  $d$ , this equation derives  $\Delta\gamma \propto 10^M$  and, further,  $\Delta\gamma \propto E$  from the Gutenberg-Richter law. Therefore, the shear strain change becomes proportional to the earthquake energy, and the increase in strain causes a corresponding increase in stress,  $\Delta\sigma$ .

In shallow soils, the carrier gas velocity of  $^{222}\text{Rn}$  (the above  $v$ ) is thought to be governed by Darcy's law (Kristiansson and Malmqvist 1982; Schery and Siegel 1986; Ioannides et al. 2003). Assuming that the intrinsic



permeability of soil and gas viscosity are constant with time and the gas flow is laminar with Reynolds number  $<1$ , then  $v$  becomes proportional to the pressure gradient,  $v \propto \partial p / \partial z$ . If the gradient is constant over the depth range, then  $v$  can be further simplified as  $v \propto p(s) - p(g)$ , where  $p(s)$  and  $p(g)$  denote the pore pressures at the  $^{222}\text{Rn}$  source depth (the maximum depth from which  $^{222}\text{Rn}$  can reach the detector)

and the ground surface, respectively;  $p(g)$  is almost constant ( $=1$  atm) despite small fluctuations with weather conditions, and  $p(s)$  is variable. It is reasonable to consider that the  $\Delta\sigma$  can enhance  $p(s)$ . This enhancement induces an increase in  $v$  and consequently an increase in  $^{222}\text{Rn}$  concentration because of the sensitivity of the concentration to  $v$ . At active fault areas with many recent earthquakes (e.g., IM, HN, and AR),  $p(s)$  is considered to be at a higher level than in inactive seismic areas because of frequent ground motions that enhance strain and stress. Additionally, the motions may have generated preferred pathways in shallow soils for the carrier gas, such as cracks and connections between pores. This pathway generation can also enhance  $v$  and cause a nonequilibrium condition of soil gas. These states essentially enhance  $^{222}\text{Rn}$  concentration without the accumulation of  $^{226}\text{Ra}$ .



### Conclusion

Through the radioactivity survey at seven active fault areas in western Japan and the data analyses, three points were clarified by this study: the radioactive equilibrium condition for soil gas, the reason for enhanced  $^{222}\text{Rn}$  concentrations in soil gas, and the differences in  $^{222}\text{Rn}$  concentrations in soil gas associated with active faults. The cause of the nonequilibrium gas patterns at 21% of the points, indicating young gas of short elapsed time since  $^{222}\text{Rn}$  generation, was interpreted as high ascent velocity of  $^{222}\text{Rn}$  carrier gas. Therefore, the radioactive equilibrium condition of the soil gas can help to specify the locations at which the formation of passages along which the gas can ascend rapidly has occurred.



The high velocity can generate anomalous  $^{222}\text{Rn}$  concentrations without accumulations of its parent nuclides.

A new finding of this study is that the observed concentrations are correlated with the summation of recent earthquake energy within a 100-km radius of the fault. From Darcy's law and the dislocation model, this phenomenon was interpreted by the increase of pore pressure around the source depth of  $^{222}\text{Rn}$  in shallow soils caused by frequently induced strain. The increase in pore pressure can enhance the gas ascent velocity and, consequently, the  $^{222}\text{Rn}$  concentration. This correlation suggests that terrestrial gas is sensitive to small changes in strain. Despite being associated with typical active faults with historically large earthquakes, the  $^{222}\text{Rn}$  concentrations measured at the study sites in Japan are not necessary high if they are not accompanied by recent seismicity. Carrier gas velocity is the dominant factor affecting concentration.

In conclusion,  $^{222}\text{Rn}$  concentration in an active fault area, in addition to being an earthquake precursor, acts as a capable gauge for evaluating the conditions of steady seismicity and strain accumulation around the fault. The  $^{222}\text{Rn}$  concentration may contribute to the identification of regions within which the crust is repeatedly ruptured.

#### Competing interests

The authors declare that they have no competing interests.

#### Authors' contributions

KK designed the research plan, led the development of this study, and wrote the manuscript with contributions from TY, TU, and HA. TY was responsible for all radioactive surveys and sample analyses. TU contributed to the data analyses, mechanism clarification, and some of the surveys. HA assisted with the surveys, data analyses, and interpretations. All authors read and approved the final manuscript.

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