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1	Improved models of the piezomagnetic field for the 2011 Mw 9.0
2	Tohoku-oki earthquake
3	
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10	

13	ABSTRACT

To assess the feasibility of observing changes in the magnetic field produced by the 14piezomagnetic effect, an improved model of the piezomagnetic field corresponding to 15the Mw 9.0 Tohoku-oki earthquake is presented. In contrast to an earlier study, the 16 proposed model explicitly considers the spatial distribution of slip on the seismic fault, 17and the results from this new model differ significantly from those of the previous 18 model where slip distributions were ignored. Quantitative aspects of the piezomagnetic 19 effect are discussed through comparisons of data and models. One feature clarified is 20that, because the fault rupture is so far offshore, the expected amplitudes are quite small 2122at onshore existing observation sites; consequently, there would have been little chance 23of observing sizable piezomagnetic signals at inland sites during the Tohoku-oki earthquake. Nevertheless, piezomagnetic signals were reportedly detected at a few sites, 24possibly indicating that the stress sensitivity or the initial magnetization was larger (by 25several factors) than assumed. On the other hand, relatively large variations in the 26magnetic field of up to 10 nT may have occurred offshore. This means that if 2728ocean-bottom sensors had been installed, larger piezomagnetic signals would have been detected. Moreover, the piezomagnetic field in offshore areas is sensitive to the detailed 29

- 30 slip distribution, suggesting that observations of the magnetic field at ocean-bottom
- 31 sites might provide important constraints on determination of slip models.
- 32
- 33 Keywords: piezomagnetic effect; stress sensitivity; 2011 Tohoku-oki earthquake; slip
- 34 distribution
- 35

36

The piezomagnetic effect, which describes changes in magnetization caused by mechanical stress, predicts changes in the Earth's magnetic field following a major earthquake. In earlier studies (e.g. Sasai, 1991, 1994, and references therein), a constitutive law of the relation between stress changes and magnetization changes has been proposed, as follows:

44

45
$$\Delta M_i = \frac{3}{2} \beta \Delta T_{ij} M_j, (i, j = x, y, z)$$
(1)

46

where ΔM_i is the change in remanent and induced magnetization, ΔT_{ij} is the deviatoric stress tensor, M_j is the initial total magnetization (i.e. the sum of induced and remanent magnetization without changes in stress), and β is a proportional coefficient that is usually referred to as the (piezomagnetic) stress sensitivity. Because of the piezomagnetic effect, it should be possible to monitor changes in stress in the Earth's crust with geomagnetic observations. Changes in the magnetic field that arise from the piezomagnetic effect are referred to as piezomagnetic fields. These fields are inverted to

54	changes in magnetization in terms of the magnetic Coulomb's law, and further inverted
55	to changes in stress in terms of eq. (1).
56	However, the usefulness of the piezomagnetic effect as a tool for monitoring
57	changes in stress is still not clear. Whether or not the piezomagnetic field can be
58	observed depends on the spatial distribution of the piezomagnetic field which, in turn,
59	depends on source type, depth and distance. If the piezomagnetic field has a detectable
60	magnitude that is restricted to just a narrow area, then its detection will be difficult with
61	poorly or sparsely spaced arrays of instruments. Numerical examinations of realistic
62	source models need to be performed in order to assess the detectability of the
63	piezomagnetic field.
64	The piezomagnetic stress sensitivity is another uncertain factor that determines
65	whether or not the piezomagnetic field is detectable. While the proportional relation (i.e.
66	eq. 1) is partially inferred from considerations based on thermodynamics (e.g.
67	Nakamura and Nagahama, 1997), and the values of stress sensitivity can be determined
68	by theoretical considerations (Stacey and Johnston, 1972), the actual magnetization
69	fraction and type varies from rock to rock and representative values for a particular
70	region must be determined from magnetic anomaly maps, geology and laboratory

72	stress sensitivities are on the order of 10^{-9} Pa^{-1} . A stress sensitivity of this order is
73	usually assumed when the piezomagnetic effect is considered in studies of volcanoes
74	(e.g. Currenti et al., 2005) and earthquakes (e.g. Okubo et al., 2011). However, these
75	values are sometimes too small to explain the observed offsets in the magnetic field
76	associated with stress changes (e.g. Nishida et al., 2004; Oshiman et al., 1990; Zhan,
77	1989). The effective values of the stress sensitivity on the geophysical scale (i.e. larger
78	than the laboratory scale) should be evaluated by comparing observational and
79	theoretical models.
80	The 2011 Mw 9.0 Tohoku-chihou Taiheiyou-oki earthquake (herein referred to as
81	the Tohoku-oki earthquake), which occurred on the boundary between the Pacific and
82	Eurasian plates, is one event for which the magnitudes of the piezomagnetic field can be
83	examined. The Tohoku-oki earthquake is the largest seismic event to have been
84	observed with a dense network of modern geophysical instruments. Along with
85	seismological and geodetic data, geomagnetic data were obtained for this extreme event.
86	Utada et al. (2011) presented a prompt and comprehensive report on observed variations
87	in the geomagnetic field associated with the Tohoku-oki earthquake. Together with
88	several types of geomagnetic variations that followed the earthquake, they also reported
89	that magnetic field offsets, which probably arose from the piezomagnetic effect, are

90	actually observed, but they are only up to 1.0 nT at the observation sites. In their
91	conclusions, Utada et al. (2011) presented a negative view on the detectability of the
92	piezomagnetic field.
93	Although the observations reported in Utada et al. (2011) provide constraints on the
94	phenomena that actually occurred at the time of the earthquake, their conclusions about
95	the piezomagnetic effect need to be reconsidered because they are based on
96	oversimplified source models that ignore the spatial distribution of slip on the fault. Any
97	reconsideration should incorporate improved piezomagnetic field models in the hope of
98	clarifying the quantitative nature of the piezomagnetic field and evaluating the
99	usefulness of observing it.
100	The aims of this study are to: (1) provide constraints on the piezomagnetic stress
101	sensitivity around the Tohoku region, near the seismic fault of the Mw 9.0 Tohoku-oki
102	earthquake; and (2) assess the usefulness of the magnetic observations as a tool for
103	detecting stress changes. To these ends, improved models of the piezomagnetic field are
104	presented, and the various models are compared and assessed using the data presented
105	by Utada et al. (2011).
106	

2. Procedures for modeling the piezomagnetic field

109	This study presents a new model of the piezomagnetic field in which the spatial
110	distribution of slip on the plate-boundary fault is explicitly considered. In many studies
111	of the piezomagnetic field in relation to earthquakes (e.g. Nishida et al., 2007), uniform
112	slip models are employed. However, the slip on a fault is generally heterogeneous, and
113	it is therefore preferable to explicitly consider the spatial distribution of slip. In a
114	uniform slip model, the fault parameters are averaged to obtain a single slip parameter.
115	Such a simplification is valid if the spatial scale of the slip on the fault is smaller than
116	the distance between the observation site and the fault. This criterion is not satisfied in
117	the case of the Tohoku-oki earthquake, for which the spatial scale of the slip and the
118	distance from the surface expression of fault rupture to observational sites are of the
119	order of 100 km. Therefore, the assumption of uniform fault slip used previously by
120	Utada et al. (2011) is likely inadequate in the calculations of the piezomagnetic fields
121	produced by the earthquake.
122	This study considers the heterogeneous fault slip models of Hayes (2011) and Shac
123	et al. (2011) (herein, referred to as the USGS and UCSB models, respectively), which
124	are derived from seismic wave inversion. Averaged slips from these models were used

125	in the calculations of Utada et al. (2011). These models were published immediately
126	following the Tohoku-oki earthquake and we can expect further refinements as more
127	complete geophysical data sets are taken in account (e.g. Koketsu et al., 2011).
128	Nevertheless, important characteristics of the Tohoku-oki fault slip are apparent in both
129	the USGS and UCSB models. For example, both models indicate large slip near the
130	trench axis. With these slip distributions, it is possible to evaluate the importance of
131	heterogeneous slip distributions and to compare our results with those of Utada et al.
132	(2011).
133	Piezomagnetic fields that correspond to a heterogeneous slip model are calculated
134	according to the following procedure. The fault plane (i.e. plate boundary) is divided
135	into sub-faults, and on each sub-fault the slip is assumed to be uniform. The total
136	piezomagnetic signal at any point is the sum of the contributions from all sub-faults, and
137	each contribution can be calculated by analytical formulae (Utsugi et al., 2000). Using
138	this procedure, we can calculate the piezomagnetic signals that correspond to
139	heterogeneous slip models for the Tohoku-oki earthquake.
140	Because the above procedure involves formulae derived by Utsugi et al. (2000), all
141	the assumptions involved in the formulations of Utsugi et al. (2000) are also used in the
142	present models. The initial total magnetization of the crust is assumed to be uniform

143	between the ground surface and a constant Curie point depth. The Earth's crust is
144	approximated by elastic half-space, surface of which locates at sea level. These
145	assumptions are not satisfied in reality, thus producing some uncertainty in the models
146	obtained for the piezomagnetic field. Errors should also be included in the slip models,
147	as mentioned above. To estimate the importance of uncertainly in the slip models and
148	Curie point depths, we calculate the piezomagnetic fields that correspond to the two slip
149	models (USGS and UCSB) with two values of H (15 and 30 km). The values for H used
150	here are the same as those used by Utada et al. (2011), and they provide reasonable
151	estimates for the island arc of the Tohoku district and for the subduction zone east of the
152	Tohoku district of Japan (Tanaka et al., 1999). The effects of heterogeneities in the
153	initial magnetization will be discussed separately, later.
154	
155	3. Features of the new piezomagnetic models
156	
157	Using the above procedures, and the parameters listed in Table 1, models of the
158	piezomagnetic field have been constructed. The spatial distribution of the expected
159	amplitude of the piezomagnetic effect is shown in Fig. 1. Observations are assumed to
160	have been made at sea level, i.e. the surface of a uniform elastic half-space. Below, I

161 enumerate the features that are commonly observed in the results and which correspond 162to all sets of parameters. It should be noted that the absolute values given in the results 163are strongly dependent on the assumed sets of parameters. For this reason, this analysis 164 focuses on relative rather than absolute values. 165Relatively large signals of the piezomagnetic field are expected to occur in offshore areas in all cases. For the UCSB slip model, piezomagnetic fields larger than 3 nT are 166167predicted in offshore areas. For the USGS slip models, the predicted piezomagnetic 168 fields are smaller than those for the UCSB model, yet changes larger than 1 nT are 169predicted. In contrast, the expected amplitudes of piezomagnetic signals over the more distant 170171onshore, including the sites of observation, are rather small. Over a vast part of the land

area, the predicted amplitudes of the piezomagnetic field are up to 0.4 nT. Precise

173 values of the expected changes at the observation sites are listed in Table 2, together

174 with the observed changes reported by Utada et al. (2011). In some models, the changes

175 predicted at some locations are as large as 0.6 nT. For example, the predicted change at

176 the ESA site is 0.6 nT for model b (i.e. UCSB slip model with H = 30 km) whereas the

177 predicted change at the same site is zero for other models. There is no location where all

the models predict changes greater than 0.4 nT.

179	It should be pointed out that these models are quite different from those that use
180	uniform slip distribution. Figure 2 shows the calculated piezomagnetic field intensities
181	that correspond to uniform slip models where the slip parameters are averaged over the
182	fault plane. Numerous differences can be observed between Figs. 1 and 2. For example,
183	the amplitudes of the signals predicted with the uniform slip model do not exceed 1.0 nT,
184	except for some localized areas. If we focus on this result, the impression is that the
185	detection of coseismic piezomagnetic signals is a hopeless task, even if the
186	observational area is extended to the seafloor. However, the amplitudes of signals
187	predicted by the heterogeneous slip model are larger than 1–2 nT across a wide area of
188	ocean. In this case, the amplitudes of the coseismic piezomagnetic signals would have
189	been detected, if suitable magnetometers had been installed in the region.
190	The large differences that exist between the uniform and heterogeneous slip models
191	highlight the importance of considering the heterogeneous model for the Tohoku-oki
192	earthquake. The large differences also indicate that many of the conclusions about
193	piezomagnetic signals by Utada et al. (2011) need to be reconsidered and probably
194	changed.

1 96 4	. Discu	ission
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198	The goals of this study were to provide constraints on stress sensitivity, and to
199	assess the usefulness of magnetic observations as tools for monitoring stress. The
200	former can be accomplished by comparing the data with the models. The latter can be
201	achieved by analyzing the constructed model. These matters are further discussed
202	below.
203	
204	4.1. Possible values of the piezomagnetic stress sensitivity
205	
206	To provide constraints on the stress sensitivity using the results of piezomagnetic
207	modeling, I make reference to the data presented by Utada et al. (2011). The idea is as
208	follows. In the proposed model, the stress sensitivity (β) and the intensity of the initial
209	total magnetization ($M = (M_x^2 + M_y^2 + M_z^2)^{1/2}$) are assumed to be given as in Table 1.
210	The assumed values are possibly different from the actual values. As the calculated
211	value of F_p (denoted by $F_p^{\text{calculated}}$) is proportional to the assumed value of βM ,
212	$[(\beta M)^{\text{assumed}}]$, the difference between $(\beta M)^{\text{assumed}}$ and the actual value of $\beta M [(\beta M)^{\text{actual}}]$
213	yields the disparity between the observed value of F_p ($F_p^{observed}$) and the calculated F_p

214 $(F_p^{\text{calculated}})$. The value of $(\beta M)^{\text{actual}}$ is given by

215

216
$$(\beta M)^{\text{actual}} = \frac{F_{\text{p}}^{\text{observed}}}{F_{\text{p}}^{\text{calculated}}} \times (\beta M)^{\text{assumed}}.$$
 (2)

217

218 If F_p^{observed} and $F_p^{\text{calculated}}$ correlate well, it is possible to determine a plausible value of 219 $(\beta M)^{\text{actual}}$.

220	Regrettably, the correlation between observed and calculated signals of the
221	piezomagnetic field is not good (Fig. 3). This means that assumptions of a uniform
222	Curie point depth, a uniform initial total magnetization, and/or an assumed slip model,
223	are inadequate. In particular, ignoring the heterogeneity of the initial magnetization is
224	possibly problematic because it is known to enhance the piezomagnetic field (e.g.
225	Oshiman, 1990). Aeromagnetic surveys over the Tohoku region have shown that
226	magnetic anomalies in this region are rather strong (i.e. 10-100 nT) (Fig. 4), raising the
227	possibility of a strong heterogeneity in the initial total magnetization. Consideration of
228	the heterogeneity of the initial magnetization is clearly important if we are to calculate
229	the piezomagnetic field accurately. However, an accurate determination of the structure
230	of the initial magnetization is generally laborious and full of possible errors; hence, an
231	accurate determination of a generated piezomagnetic field is difficult in the presence of

a strong heterogeneity in the initial magnetization (e.g. Yamazaki, 2011).

233	Nevertheless, we can attempt to provide constraints on the possible values of the
234	stress sensitivity using data just from sites KAK and KTR. Around these sites, the
235	gradient of the magnetic anomaly is relatively small (Fig. 4), and we can therefore
236	anticipate that the model with uniform initial magnetization will provide reasonable
237	calculated results. The amplitude of the piezomagnetic signal observed at KTR was -0.8
238	\pm 0.2 nT, whereas those predicted in the theoretical models (Fig. 1a–d) are between –0.2
239	and -0.3 nT. The piezomagnetic signal observed at KAK was -0.22, whereas those
240	predicted by theoretical models are between -0.07 and -0.22 . To explain the
241	observations at KAK and KTR, the actual value of βM needs to be larger than the value
242	assumed in the present calculation (i.e. $1.0 \text{ Pa}^{-1}\text{Am}^{-1}$) by factors of 2–3. Provided that
243	the assumption of $M = 1.0$ A/m (Table 1) is correct, the above result means that the
244	stress sensitivity is about $2.0-3.0 \times 10^{-9} \text{ Pa}^{-1}$. This value is on the same order as that
245	assumed in many piezomagnetic models (e.g. Johnston et al., 1989).
246	
247	4.2 Potential usefulness of seafloor magnetic observations
248	

In all the models of the piezomagnetic field shown in Fig. 2, the amplitudes of the

250	piezomagnetic signals are expected to be small on land and relatively large offshore.
251	Because the actual value of βM is possibly larger than assumed, as discussed in the
252	previous subsection, the actual changes in the magnetic field are possibly larger than
253	those shown in Fig. 1. In particular, large offsets in the magnetic field are expected near
254	the trench axis. Although this result is obtained for an assumption that observations are
255	made at sea level, piezomagnetic signals are also expected to be large at seafloor
256	because the seafloor is rather closer to the rupture. If offshore ocean-bottom
257	magnetometers had been installed, they would have detected significant amplitudes of
258	piezomagnetic signals corresponding to the Tohoku-oki earthquake.
259	Observing piezomagnetic signals would not be very useful if the piezomagnetic
260	signals were insensitive to details of the fault parameters, but the results of the
261	piezomagnetic models demonstrate that this is not the case. Indeed, the spatial
262	distributions of the piezomagnetic field are strongly dependent on the slip model that is
263	adopted. It is possible that we could have improved the determination of earthquake
264	source parameters if data from ocean-bottom magnetometers had been available, instead
265	of relying solely on the results of inversions of seismic and geodetic data.
266	In general, it is difficult to measure accurately the distributions of slip for
267	earthquakes that occur on an offshore plate boundary, because geodetic measurements

268	are made mainly on land. In the case of the Tohoku-oki earthquake, extremely large
269	slips near the trench axis have been suggested by inversions of the seismic data (e.g.
270	Hayes, 2011; Shao et al., 2011), but better constraints on slip distributions could have
271	been obtained from seafloor geodetic measurements (Sato et al., 2011). Given that
272	seafloor geodetic equipment is costly and difficult to manage, geomagnetic observations
273	might provide useful additional data for monitoring interplate earthquakes along
274	subduction zones. This solution may still apply, even when we consider magnetic
275	anomalies on the seafloor, because heterogeneities in the magnetization of the crust may
276	possibly enhance the generated piezomagnetic field (e.g. Oshiman, 1990).
277	Regrettably, there are also drawbacks to making seafloor magnetic observations.
278	First, seafloor observations are quite costly. Second, it may be difficult to keep the
279	sensors stably located during quakes, and if a sensor is displaced during a quake, an
280	apparent change in the magnetic field will be recorded. Even if it were possible to
281	obtain accurate data of the geomagnetic field at a certain point, it would be necessary to
282	consider heterogeneities of the initial magnetization and ocean-bottom topography when
283	converting the observed changes in the magnetic field to fault source parameters. For
284	these reasons, the usefulness of observing the piezomagnetic field remains uncertain.
285	However, similar difficulties also exist with respect to making ocean-bottom geodetic

286	observations. Not only are they are extremely costly, but monument stability during
287	earthquakes is also a problem. It is also difficult to process the observations correctly
288	and obtain precise geodetic information. A decision on prioritizing geodetic and
289	geomagnetic techniques should be based on which drawbacks are most easily overcome.
290	If costs allow, an integrated use of both techniques is most desirable because they
291	independently bring useful information to bear on these tectonic phenomena.
292	
293	5. Conclusions
294	
295	To calculate the piezomagnetic field that corresponds to the 2011 Mw 9.0
296	Tohoku-oki earthquake, it is necessary to consider the best representation of the spatial
297	distribution of slip along the fault, and in this paper, I demonstrate the importance of
298	such a consideration, and construct an appropriate slip model. Although this model still
299	cannot entirely explain the observed distribution of piezomagnetic signals, constraints
300	
	are obtained from data at two onshore sites, where the model seems to provide adequate
301	are obtained from data at two onshore sites, where the model seems to provide adequate results. Comparisons between the data and the model show the stress sensitivity to be
301 302	are obtained from data at two onshore sites, where the model seems to provide adequate results. Comparisons between the data and the model show the stress sensitivity to be about $2-3 \times 10^{-9}$ Pa ⁻¹ , which is on the same order as that assumed in many

304	geomagnetic total forces, due to the piezomagnetic effect, will be relatively large in
305	offshore areas closer to the rupture, and relatively small onshore, far from the rupture.
306	Because the expected magnitudes of the piezomagnetic signals are small at existing sites,
307	stress sensitivity of the piezomagnetic effect is likely to be on the order of 10^{-9} , though
308	this is not tightly constrained. Nevertheless, the possibility of a large piezomagnetic
309	field occurring at ocean-bottom stations is not excluded. Details of the spatial
310	distribution of the piezomagnetic field in oceanic areas are highly dependent on the slip
311	model used. The implication is that detection of the piezomagnetic field with
312	ocean-bottom magnetometers might have provided constraints on the slip models of the
313	Tohoku-oki earthquake, if such observations had been available.
314	
315	Acknowledgments
316	
317	Slip models proposed by Shao et al. (2011) and Hayes (2011), and a program coded
318	by M. Utsugi, were used to calculate the piezomagnetic fields. Numerical data on the
319	geomagnetic anomalies around Japan were provided by the Geospatial Information
320	Authority of Japan through their Web site. Comments and advice from anonymous
321	reviewers and the journal editor (P. Shearer) helped to improve an earlier version of the

322 manuscript. The software Generic Mapping Tools (Wessel and Smith, 1998) was used to

323 prepare some of the figures, including maps.

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326

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388 Figure captions

390	Fig. 1. Models of the piezomagnetic field corresponding to the UCSB and USGS slip
391	models. (a) UCSB slip model with $H = 15$ km. (b) UCSB slip model with $H = 30$
392	km. (c) USGS slip model with $H = 15$ km. (d) USGS slip model with $H = 30$ km.
393	The rectangle represents the fault plane on which heterogeneous slip is considered.
394	
395	Fig. 2. Models of the piezomagnetic field corresponding to averaged versions of the
396	UCSB and USGS slip models. (a) UCSB slip model with $H = 15$ km. (b) UCSB
397	slip model with $H = 30$ km. (c) USGS slip model with $H = 15$ km. (d) USGS slip
398	model with $H = 30$ km. The rectangle represents the fault plane on which uniform
399	slip is considered.
400	
401	Fig. 3. Comparisons between observed and calculated piezomagnetic signals.
402	Calculated values in this figure are the averages of four piezomagnetic models.
403	Error bars in calculated values represent maximum and minimum values for the
404	four models. Error bars in observed values are from Utada et al. (2011). Open
405	circles indicate the results at sites KAK and KTR, where the magnetic anomalies

406	are rather small,	while solid	circles in	ndicate the	results at	other sites.

408	Fig. 4. The magnetic anomaly over the Tohoku region as observed by an aeromagnetic
409	survey at a height of 5000 m. Contour intervals are 10 nT. Observations were
410	conducted by the Geographical Survey Institute (predecessor of the Geospatial
411	Information Authority) of Japan in 1990, and the data are available on their Web
412	site, in Japanese
413	(http://vldb.gsi.go.jp/sokuchi/geomag/menu_03/aeromag_data.html; last access: 26
414	October 2012).
415	
416	
417	



Figure 1 Models of the piezomagnetic field corresponding to the UCSB and USGS slip models. (a) UCSB slip model with H = 15 km. (b) UCSB slip model with H = 30 km. (c) USGS slip model with H = 15 km. (d) USGS slip model with H = 30 km. The rectangle represents the fault plane on which heterogeneous slip is considered.



Figure 2 Models of the piezomagnetic field corresponding to averaged versions of the UCSB and USGS slip models. (a) UCSB slip model with H = 15 km. (b) UCSB slip model with H = 30 km. (c) USGS slip model with H = 15 km. (d) USGS slip model with H = 30 km. The rectangle represents the fault plane on which uniform slip is considered.



Figure 3 Comparisons between observed and calculated piezomagnetic signals. Calculated values in this figure are the averages of four piezomagnetic models. Error bars in calculated values represent maximum and minimum values for the four models. Error bars in observed values are from Utada et al. (2011). Solid circles indicate the results at sites KAK and KTR, where the magnetic anomalies are rather small, while open circles indicate the results at other sites.



Figure 4 The magnetic anomaly over the Tohoku region as observed by an aeromagnetic survey at a height of 5000 m. Contour intervals are 10 nT. Observations were conducted by the Geographical Survey Institute (predecessor of the Geospatial Information Authority) of Japan in 1990, and the data are available on their Web site, in Japanese (http://vldb.gsi.go.jp/sokuchi/geomag/menu_03/aeromag_data.html; last access: 26 October 2012).

Table 1

Parameters assumed in the modeling of the piezomagnetic field.

Parameter	Value
Rigidity	$57(53) \times 10^9$ Pa
Poisson's ratio	0.25
Magnetization	1.0 A m^{-1}
Piezomagnetic	$1 \times 10^{-9} \mathrm{Pa}^{-1}$
stress sensitivity	
Curie point depth	15 and 30 km
Inclination of the ambient	51.0 degree
geomagnetic field	
Declination of the ambient	-7.5 degree
geomagnetic field	
Observation altitude	0 m

Table 2

Comparisons of coseismic changes in the geomagnetic total intensity, as predicted by the piezomagnetic models versus those calculated from data reported in Utada et al. (2011). Piezomagnetic field models are determined for Curie point depths of 30 and 15 km together with two slip models (USGS and UCSB).

Station code	USGS slip model		UCSB slip model		Observed (error)	
	30 km	15 km	30 km	15 km		
ESA	+0.05	+0.05	+0.55	+0.49	-0.09	(0.46)
HAR	-0.56	-0.37	-0.40	-0.06	-0.43	(0.36)
IWK	-0.48	-0.28	-0.44	-0.14	+0.21	(0.42)
KTR	-0.18	-0.10	-0.29	-0.18	-0.83	(0.20)
OTA	+0.14	+0.16	-0.08	-0.03	-0.01	(0.28)
KAK	-0.22	-0.14	-0.19	-0.07	-0.22	(0.10)