

1 Combined Effects of Nanoparticles and Surfactants upon Controlling Foam

2 Stability Synergy Between Nanoparticles and Surfactants in Controlling Foam Stability

3 Mohammad Javad Shojaei¹, Yves Méheust², Abdulkadir Osman³, Paul Grassia⁴, and Nima

4 Shokri*⁵

5¹Department of Earth Science and Engineering, Imperial College London, London, United

6 Kingdom

7²Université de Rennes, CNRS, Géosciences Rennes, UMR 6118, 35000 Rennes, France

8³Department of Chemical Engineering and Analytical Science, The University of Manchester,

9 Manchester, UK

10⁴Department of Chemical and Process Engineering, University of Strathclyde, Glasgow G1

11 1XJ, United Kingdom

12⁵Hamburg University of Technology, Institute of Geo-Hydroinformatics, Am

13 Schwarzenberg-Campus 3 (E), 21073 Hamburg, Germany

14 *Corresponding Author

15 Prof. Nima Shokri

16 Email: nima.shokri@tuhh.de

17 Website: www.tuhh.de/ghi

18Abstract

19We investigate the effect of surfactants with different charges (anionic, cationic, and non-
20ionic) on foam stability in the presence of charge-stabilized silica (SiO₂) nanoparticles.
21Toward this aim, a comprehensive series of experiments on a Hele-Shaw cell and a foam
22column is conducted at bubble and bulk-scale respectively, that is, investigating
23phenomenologies of foam coarsening separately by gas diffusion and bubble coalescence,
24and by gravitational drainage. Our results show nanoparticles, despite their ability to position
25themselves at liquid-gas interfaces and thus limit the resulting surface tension coefficient, do
26not necessarily have a positive effect on foam stability; the nature and magnitude of this
27effect depends strongly on the nature of the surfactant, its concentration and the concentration
28of nanoparticles. In less stable systems, significant coarsening ~~by bubble coalescence~~ occurs.
29Both results from bubble-scale and the bulk-scale experiments suggest that compatibility
30experiments are pre-requisite to foam stability analysis to test the compatibility between
31surfactants and nanoparticles.

32**Keywords:** Foam stability, Diffusional coarsening, Bubble coalescence, Gravitational
33drainage, Nanoparticle, Surfactant, ~~Bubble and Bulk scale~~.

34 Introduction

35 Gas injection into subsurface reservoirs is a common practice in many industrial and
36 engineering processes such as enhanced oil recovery (EOR), CO₂ sequestration and soil
37 remediation (Benson and Cole, 2008; Blunt et al., 1993; Feng et al., 2012; Kantzas et al.,
38 1988). In most cases, viscous fingering and gravity override due to unfavourable viscosity
39 and density ratios between the gas and the resident liquid(s), and preferential flow of gas due
40 to reservoirs heterogeneity, are responsible for low sweep efficiency (Chang et al., 1994;
41 Garcia and Pruess, 2003). Foams, which are dispersions of a large volume of gas in a liquid
42 such that the gas phase is made discontinuous by films of the liquid phase denoted lamellae
43 (Hirasaki et al., 1997a; Kam and Rossen, 2003; Shojaei et al., 2018a), are a promising
44 potential remedy to these complications (Hirasaki and Lawson, 1985; Shojaei et al., 2019).
45 The apparent viscosity of foam can be up to 1000 times higher than that of its constituents,
46 which makes foams ideal for fluid displacement (Hirasaki et al., 1997b; Shojaei et al.,
47 2018b).

48 In general, foams are classified into two categories, which are typically known as *bulk foam*
49 and *confined foam* (Rossen, 1996), based on the size of bubbles relative to the typical length
50 scale of the confined media (e. g., the average pore size or channel width). foam can be
51 considered a *bulk foam* when the dimension of the confining space is significantly larger than
52 the typical bubble size. On the other hand, the foam is *confined foam* when the bubbles have
53 the same size or are larger than the characteristic length scale of the confining space. Foams
54 exhibit two different geometries depending on their quality, i. e., their gas content (Ma et al.,
55 2012). In wet systems (i.e., at low foam qualities), the lamellae are thick, the foam bubbles
56 have a quasi-spherical shape, and the foams are fine-textured, whereas, at higher foam
57 quality, the lamellae are thinner and foam bubbles tend to have a more polyhedral shape.

58The stability of a foam refers to its capacity to retain its geometry/topology over a significant
59amount of time despite not being stable thermodynamically. In porous media applications
60involving non-aqueous phase liquids (NAPLs), such as foam EOR or foam-based remediation
61of NAPL-contaminated aquifers and soils, the foam's texture evolves irreversibly in time as a
62consequence of four different processes: (1) gas diffusion (2), liquid drainage (3), interaction
63with oil/NAPL and (4) capillary suction (Ma et al., 2012; Osei-Bonsu et al., 2015; Rossen,
641996){Rossen, 1996 #2512}. In the capillary suction mechanism, when the capillary pressure
65(the pressure across the interface between the gas and the surfactant solution) increases, the
66lamellae thickness decreases, eventually causing it to break if a threshold in capillary pressure
67is exceeded. That threshold is called the maximum capillary pressure beyond which
68coalescence (i.e., appearance of a larger bubble as a result of the breakage of film between
69two smaller bubbles) occurs.

70Adjacent foam bubbles do not have the same size, and hence the gas is at different pressures
71inside the bubbles. The gas in smaller foam bubbles is at a higher pressure than the gas in
72larger ones. Indeed, the bubble radius controls the pressure inside the bubbles as a
73consequence of the Young–Laplace equation, which relates the pressure difference ΔP across
74a fluid interface to the surface tension coefficient σ and the principal radii of curvature r_1 and

75 r_2 according to $\Delta P = \sigma \left(\frac{1}{r_1} + \frac{1}{r_2} \right)$ (Lemlich, 1978). Gas thus diffuses from the small bubbles

76with higher pressure to larger bubbles with lower pressure, which eventually causes the
77disappearance of neighbouring small bubbles (Blijdenstein et al., 2010; Maestro et al., 2014;
78Saint-Jalmes, 2006). This phenomenon is called gas diffusion ~~on~~ coarsening.

79Liquid drainage is a multistage process consisting of (a) liquid flow from the lamellae to the
80Plateau border (which are the lamellae's intersections) due to capillary suction, (b) liquid

81 release from the coalescence of foam bubbles, and (c) downward liquid drainage along
82 Plateau borders under the effect of gravity, resulting in accumulation of liquid in the lower
83 layer of the foam (Exerowa and Kruglyakov, 1997). The entire process is mainly controlled
84 by gravity and capillary suction and eventually leads to film breakage ruptures (and, hence,
85 bubble coalescence) as the thickness of lamellae falls below a certain value (Bhakta and
86 Ruckenstein, 1997). Drainage, therefore, presents a challenge for foam-based displacement
87 processes.

88 Another major challenge to the effective utilization of foam application in oil displacement is
89 the adverse effect of oil on foam stability as a result of direct surface interactions between oil
90 and foam, which leads to aqueous film thinning and breakage rupture (Koczo et al., 1992;
91 Nikolov et al., 1986; Osei-Bonsu et al., 2018). The negative effect of oil on foam stability
92 depends on the properties of the surfactant and oil. Light oil (small hydrocarbon chains) has
93 been found to be more detrimental to foam stability than heavier oil (long hydrocarbon
94 chains) (Lobo et al., 1989; Talebian et al., 2013).

95 In view of the above-mentioned challenges to foam stability, in recent years, there has been a
96 growing interest in the joint utilization of nanoparticles and surfactant to stabilize foams
97 (Karakashev et al., 2011; Kumar and Mandal, 2017; Maestro et al., 2014; Yekeen et al., 2018;
98 Yu et al., 2012b). The effective contribution of nanoparticles to foam stability is attributed to
99 the adsorption and accumulation of nanoparticles at the gas-liquid interfaces of foam bubbles
100 and Plateau borders (Yekeen et al., 2018). Nanoparticles reduce the direct contact between
101 the fluids, which decelerates the gas diffusion rate and bubble bursting (Karakashev et al.,
102 2011; Maestro et al., 2014; Yu et al., 2012b), and film drainage is slowed as well due to the
103 presence of the nanoparticles. The lower tendency of nanoparticles (compared to the
104 surfactant) to adsorb on reservoir rocks is another reason that makes them a desirable foam
105 stabilizer (Yekeen et al., 2018). Nanoparticles are well suited to subsurface applications.

106 Their small size limits the possibility of pore plugging as they pass through the pore throats in
107 porous media (Yu et al., 2012b). Their solid nature also makes them highly resistant to the
108 harsh condition of reservoirs such as high pressure and temperature, high salinity and the
109 presence of oil (Yusuf et al., 2013; Zhang et al., 2009). Also, nanoparticles can be
110 functionalized with different chemical groups to improve their aqueous stability and tune the
111 wettability of the solutions, or coated for different purposes such as increasing their CO₂
112 solvation capability and capability to adhere to the fluid-gas interface, which contributes to
113 improving the foam's stability (Panthi et al., 2017; Singh and Mohanty, 2017).

114 In addition to the decrease in gas diffusion and liquid drainage, the main proposed causes for
115 the increase in foam stability when using nanoparticles are an increase in particle detachment
116 energy and in the maximum capillary pressure for [bubblefoam](#) coalescence (Yekeen et al.,
117 2018). The particle detachment energy is the energy required for the removal of individual
118 nanoparticles from lamellae (Singh and Mohanty, 2015). The adsorption of nanoparticles at
119 the interface is thus considered irreversible due to their large detachment energy, while other
120 conventional foaming agents can easily adsorb and desorb from the gas-liquid interface of
121 foam bubbles. Therefore, the presence of the adsorbed SiO₂ nanoparticles, by increasing the
122 lamellar stability, increases the maximum capillary pressure beyond which coalescence
123 occurs (Yekeen et al., 2018). This increase in maximum capillary pressure depends on
124 nanoparticle concentration and on how they agglomerate at the gas-liquid interface.

125 The presence of nanoparticles at the gas-liquid interface decreases the [surface](#)
126 ~~tension~~[interfacial tension](#) of foam bubbles with respect to conventional foams (Kantzas et al.,
127 1988), and hence decreases the capillary pressure. Consequently, the pressure differences
128 between adjacent bubbles decrease in the presence of nanoparticle and gas diffusion
129 decreases accordingly, while the permeability of the film to gas decreases also. [Also, the fact](#)
130 [that the surface tension of the surfactant solution decreases](#)~~Also, the surface tension of the~~

131 ~~surfactant solution decreases~~ in the presence of silica nanoparticle (Jia et al., 2020;
132 Vatanparast et al., 2018) that could potentially improve foam stability and foam generation.

133 Nanoparticles can be arranged at the gas-liquid interface as a monolayer, bilayer, or a
134 network of particles based on their surface wettability (Bi et al., 2004; Horozov, 2008). The
135 resistance of nanoparticles to exit the interface controls the stability of a monolayer
136 nanoparticle arrangement (Kantzas et al., 1988), while the stability of a bilayer and network
137 of nanoparticles arrangement is influenced by interfacial rheological properties and by the
138 capillary pressure (Kantzas et al., 1988). Generally, a network of nanoparticles provides
139 higher stability by forming thick solid lamellae that prevent film thinning and gas diffusion
140 more effectively by increasing the surfactant solution's viscosity and decreasing gas
141 diffusivity. In addition, liquid drainage and gravitational drainage could be decelerated in the
142 presence of nanoparticles. Hence the arrangement of nanoparticles at the interface during
143 liquid drainage is a key control parameter in foam stability enhancement by nanoparticles.

144 It has been claimed based on experimental data that in any given system, there is an optimal
145 concentration of nanoparticles that improves foam stability to the largest extent (Espinoza et
146 al., 2010). At low concentration, the presence of nanoparticles at the gas-liquid interface is
147 not sufficient to achieve high stability. As the nanoparticle concentration is increased, more
148 nanoparticles find themselves at the gas-liquid interfaces, which enhances foam stability by
149 reducing foam drainage and liquid film thinning. However, foam stability either remains
150 constant or decreases when the concentration passes a critical value (AttarHamed et al., 2014;
151 Chen et al., 2014). It has been established that nanoparticles, irrespective of the type, have a
152 significant influence on static and dynamic stability of foam (Yekeen et al., 2018). What is
153 not yet clearly understood is how the nature of the surfactant affects foam stability in the
154 presence of nanoparticles. [To improve our physical understanding of the interaction between](#)
155 [nanoparticles and surfactants in determining foam stability.](#) ~~To improve our physical~~

156 ~~understanding of the synergy between nanoparticles and surfactants on foam stability~~, in this
 157 study we investigate the impact of nanoparticles in the presence of surfactants with varying
 158 charges (anionic, cationic, and non-ionic) on foam stability, using column experiments and
 159 Hele-Shaw cell experiments. The experiments performed in a horizontal Hele-Shaw cell
 160 provide information about foam coarsening in the absence of gravity drainage, while the
 161 column experiments allow quantifying the magnitude of gravitational drainage. In particular
 162 we present the first investigation of the synergy between non-ionic surfactant and
 163 nanoparticles.

164 **Materials and Methods**

165 **Foaming suspensions:**

166 All of the foam experiments were prepared using deionized water in ambient conditions
 167 ($T \sim 23^\circ\text{C}$, $\text{RH} \sim 36\%$). We used deionized water to keep the chemistry as simple as possible
 168 although oil reservoir conditions could be saline. Three surfactants of different natures
 169 (respectively anionic, cationic, and non-ionic) were used in this study; sodium dodecyl sulfate
 170 (SDS) (Sigma, UK), dodecyl trimethyl ammonium bromide (DTAB) (Sigma, UK) and Triton
 171 X100 (Sigma, UK), respectively. The properties of these surfactants used in this work are
 172 summarised in Table 1. The surfactants were used at their CMC, (unless otherwise specified).

173 ~~Charge-stabilized dispersions of spherical colloidal silica particles~~ ~~Spherical charge-stabilized~~
 174 ~~dispersions of colloidal silica particles~~ (Ludox HS, Grace) with a diameter of 16 nm were
 175 added to the surfactant solutions.

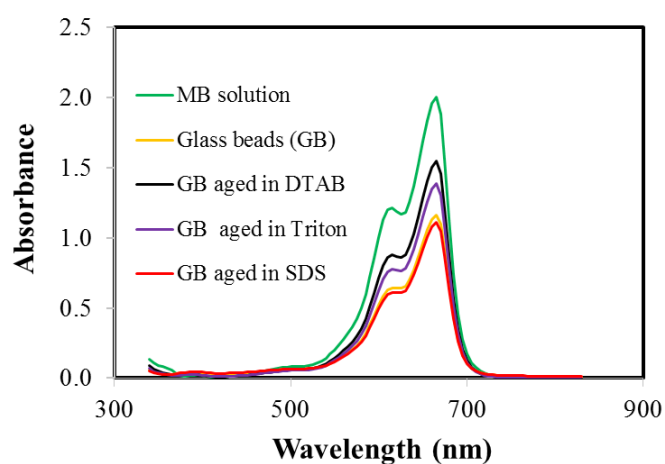
176 Table 1 Properties of the surfactants (Lin et al., 1999; Yu et al., 2012a), including the critical
 177 micellar concentration (CMC).

Surfactant	Charge	CMC (mM)	CMC (%w/ wv)
Sodium dodecyl sulphate (SDS)	Anionic	8	0.23

Dodecyl trimethyl ammonium bromide (DTAB)	Cationic	11	0.46
Triton X100	Non-ionic	0.24	0.02

178

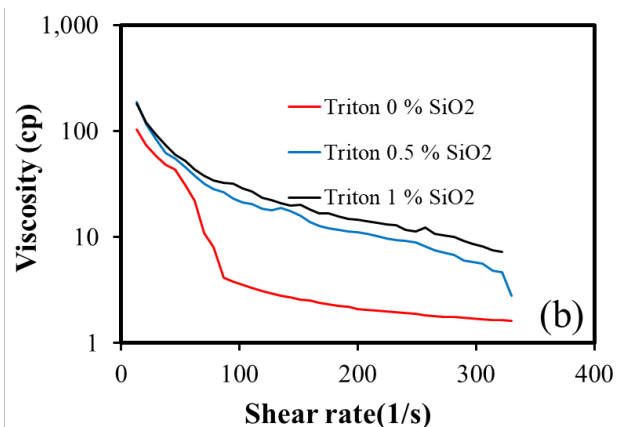
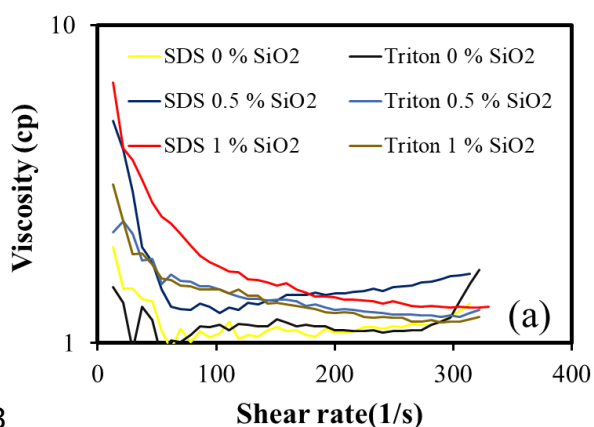
179The interaction of the surfactants with the silica particles was characterized qualitatively
180using ultraviolet-visible (UV-vis) spectroscopy, with an analysis based on the method
181described by Desarnaud et al. (Desarnaud et al., 2016). It is based on measuring the
182~~decolourization~~ ~~decolorization~~ of a dye solution (here a cationic dye: methylene blue (MB))
183due to the adsorption of the dye on the oppositely charged surface (i.e., the silica particles).
184Here, one would expect that the formation of a silica-surfactant complex, due to charged
185interactions, would reduce the ~~decolourization~~ ~~decolorisation~~ of the MB solution due to the
186surface of the silica particles being essentially covered by adsorbed surfactant molecules,
187which limits absorption of the dye onto the particles' surface. To confirm this, silica particles
188were mixed in each of the prepared surfactant solutions. The solutions were then filtered and
189left to dry. The obtained dried particles were then placed in an MB solution, and using
190UV/Vis-spectrometer, the ~~decolourization~~ ~~decolorisation~~ of each dye solution was measured.



191

192Figure 1: UV-vis absorption spectra of methylene blue solution (MB) after the addition of
193silica beads and silica beads treated in SDS, DTAB, and Triton X100 solution.

194 In the case of the negatively charged SDS surfactant, the measurements show that the anionic
 195 surfactant hardly adsorbs onto the negatively charged surface of the silica particles, presented
 196 in Figure 1, since the reduction of absorbance (indicative of the [decolourization](#)
 197 ~~decolorization~~ of the solution) is nearly identical when silica particles and silica particles
 198 treated in SDS solution are added to the methylene blue solution. ~~This observation is~~
 199 ~~consistent with the recent work of Osman et al. (Osman et al., 2018).~~ Conversely, in the case
 200 of silica particles treated in DTAB solutions, the decolourization is measured to be
 201 significantly less intense due to adsorption of the cationic surfactants onto the oppositely
 202 charged silica surfaces, which minimize the interactions between the MB dye and silica.
 203 Similarly, but to a lesser extent, adsorption of Triton X100 onto the silica particles also
 204 occurs, as seen in Figure 1. Figure 13 in Appendix A shows an image of the solutions
 205 containing silica nanoparticles. In the case of DTAB (Figure 13b), flocculation occurs due to
 206 the strong interaction of the cationic surfactants with the anionic silica particles. The
 207 adsorption of DTAB onto the surface of the particles tunes the DLVO barrier, which
 208 describes the balance between charge-induced repulsive forces together with the attraction
 209 induced by van der Waals forces at a short-range (Derjaguin and Landau, 1993; Verwey and
 210 Overbeek, 1955). In this case, adsorption leads to a decrease in the electrostatic repulsion
 211 between the nanoparticles, and consequently, the van der Waals attraction become dominant,
 212 thus contributing to flocculation of the suspension.



213

214 Figure 2: Dependence of the apparent viscosity on the shear rate of the different foaming
215 suspensions at different silica nanoparticles concentrations. Both surfactants are at their
216 CMC in (a). The concentration of Triton X100 in (b) is 1 % (is over 50 times CMC~~6.7 times~~
217 ~~more than its CMC~~). The vertical axis is in log scale. As the DTAB precipitated and we had
218 two phases, we did not present its viscosity data in (a).

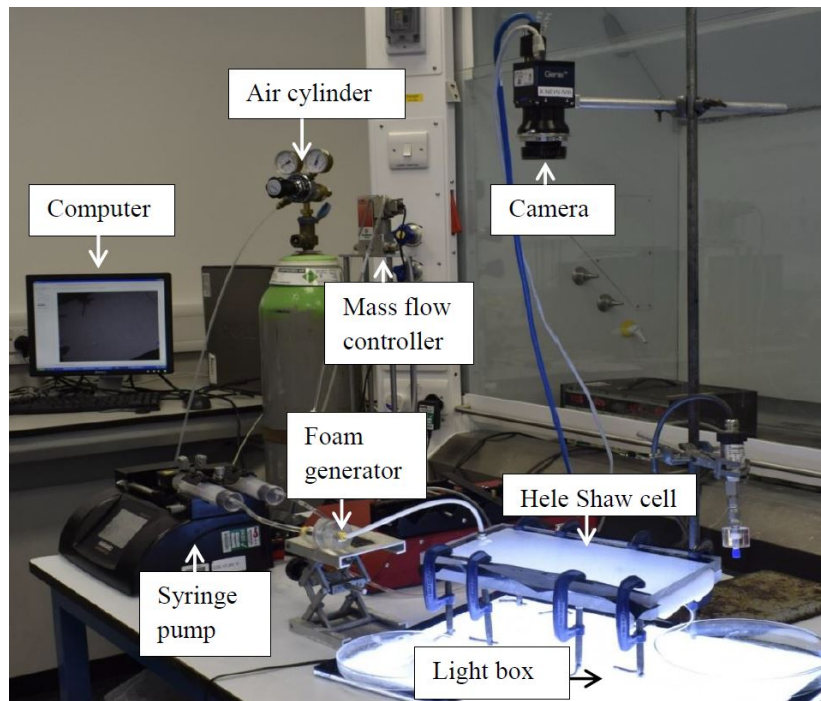
219

220 The rheology of the foaming suspensions was measured using a rheometer (Rotational DV3T
221 Rheometer, Brookfield) in the plate-plate configuration. Figure 2 presents the viscosity's
222 dependence on the shear rate for surfactant solutions at different nanoparticle concentrations
223 for surfactants at the CMC (Figure 2a) and at 1% (Figure 2b). The plots do not exhibit
224 monotonicity in shear rate as ~~expected~~, most probably due to heterogeneity in nanoparticle
225 density within the samples during measurements. However, they show that the addition of
226 silica nanoparticles tends to increase the viscosity of the solution: the larger the concentration
227 in particles, the larger the viscosity. This effect is especially significant for Triton X100, as
228 viscosity increases by up to one order of magnitude when particles are added for Triton at
229 concentration of was 1% rather than as well as at ~~This effect is especially significant for Triton~~
230 ~~X100, as viscosity increases by up to one order of magnitude when particles are added and~~
231 ~~surfactant concentration is increased from the CMC to 1%.~~ This ~~can~~ could be due to the
232 interaction between nanoparticles and surfactants at high concentrations of Triton X100- and
233 also probably partly due to the fact that Triton X100 is a viscous liquid, which would
234 explains the higher measured viscosities at low shear rates and 0wt% NP nanoparticles. Note
235 also that rheometry measurements could not be performed with DTAB together with
236 nanoparticles, due to the flocculation of the nanoparticles in DTAB-based suspensions.

237 **Experimental Set-up and Procedure**

238A series of foam stability experiments were conducted using a Hele-Shaw cell (Figure 3) and
239a column cell (Figure 3) to investigate the synergy between nanoparticles and surfactants in
240impacting foam stability at bubble and bulk-scale, respectively.

241**Bubble-scale experiments:** The Hele-Shaw cell consisted of two plexiglass plates of
242dimensions $30 \times 17 \times 0.5 \text{ cm}^3$. The plexiglass plates were tightened using medium-duty
243clamps in all experiments. A gasket of thickness 1 mm was clamped between the two plates
244to impose a constant distance between them and prevent leakage. ~~Two~~One holes (1 mm
245diameter) ~~were~~was drilled on ~~opposite~~the sides on the top of the plexiglass plate to act as
246inlet ~~and outlet channels~~ for the flow of foam through the Hele-Shaw cell. Foam was
247generated by injecting both compressed air and the surfactant solution simultaneously into a
248foam generator fitted with a sintered glass disc (Scientific Glass, UK) with a pore size
249distribution between 40 and 60 μm . The flows of gas and surfactant were set to 10 ml/min
250and 1.11 ml/min respectively to achieve a 90 % foam quality for all the foam displacement
251experiments. The pressure was measured at the inlet of the Hele-Shaw cell via a pressure
252transducer, while the outlet was connected to the atmosphere. The Hele-Shaw cell was
253initially fully saturated by air.



254

255

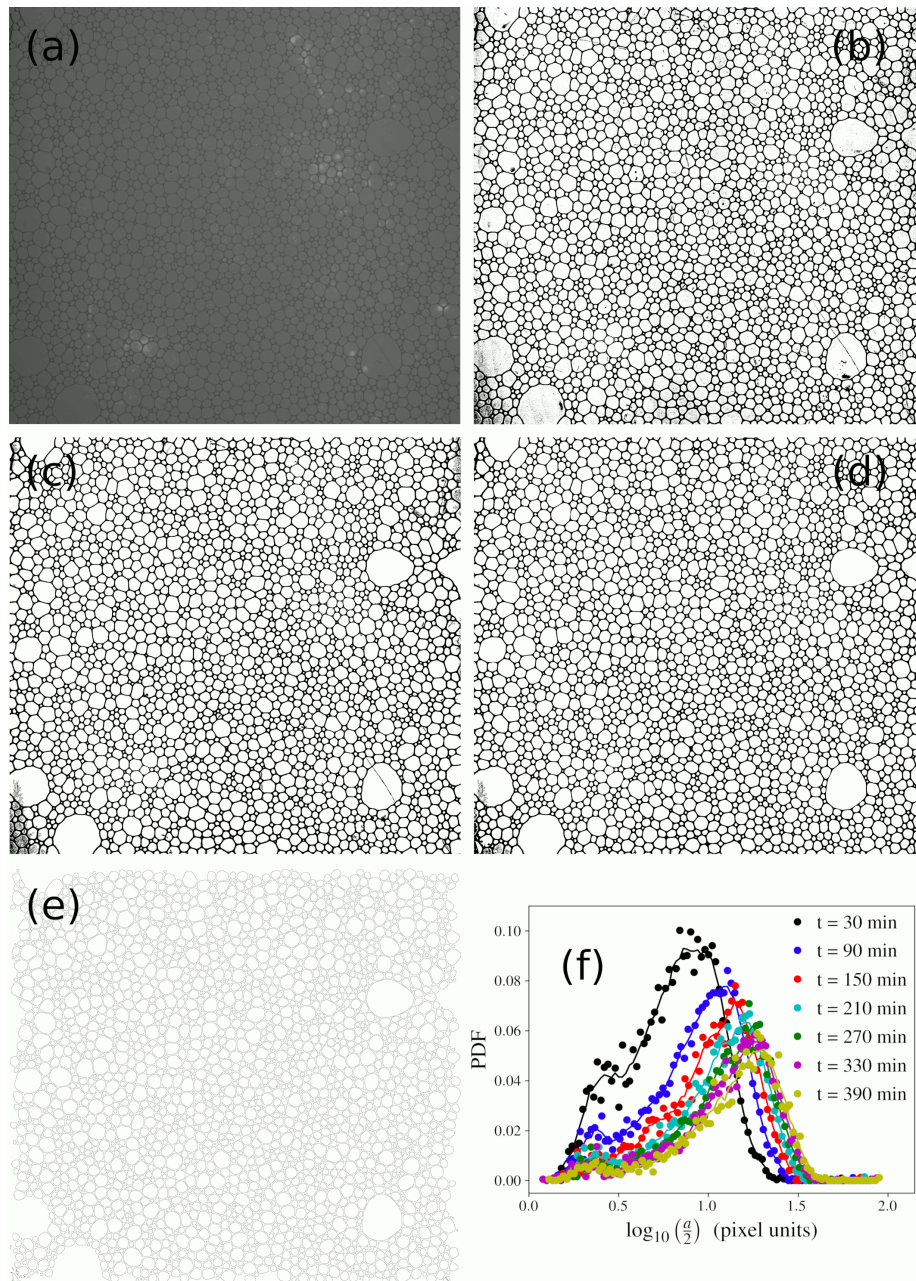
Figure 3: Hele-Shaw cell experimental set-up.

256 A high-resolution camera (Teledyne DALSA genie) was placed above the micro model and
 257 captured a snapshot of the ageing process every 30 minutes for a total of at 6 hrs or more. The
 258 images produced were 8-bit grey levels with a resolution of 2560 x 2048 pixels. The contrast
 259 of the images was improved by the use of a lightbox placed underneath the model.

260 The images were treated using ImageJ and Matlab in order to identify individual bubbles and
 261 measure their apparent area. The procedure is illustrated in Figure 4 using the image recorded
 262 6.5 h after the start of the experiment performed with a suspension containing SDS at its
 263 CMC and 1% of SiO₂ nanoparticles. The raw image (Figure 4a) is first segmented using
 264 ImageJ's "local thresholding" procedure, with a local threshold value obtained from running
 265 a moving average filter with a window of linear size 500 pixels; the resulting image is shown
 266 in Figure 4b. From this image, a better image (Figure 4c) is obtained by removing all
 267 connected black regions except the largest one, which runs between the bubbles; this removes
 268 black spots which are seen inside bubbles in Figure 4b. Other black spots are removed from

269bubbles in the image of Figure 4d, which has been obtained from that of Figure 4c by
270replacing each white connected region by its filled convex hull (that is, the filled convex
271polygon that covers the region the most closely). These two steps are done with custom-made
272MATLAB scripts. For some of the data sets this last step of the treatment is not necessary.
273Individual bubble areas are then measured from analysing the connected white regions in the
274image of Figure 4d, disregarding those of these regions which are in contact with an image
275boundary (and therefore, which correspond to bubbles that are not entirely captured inside the
276image). An outline of the corresponding bubbles is shown in Figure 4e.

277



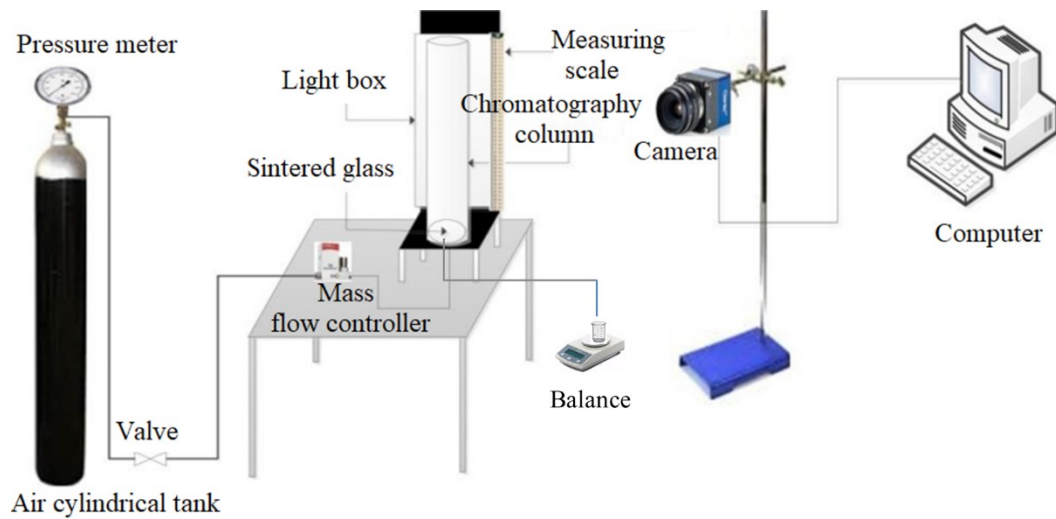
278

279 Figure 4: Various steps of the image treatment procedure: (a) Raw image; (b) Image
 280 segmented using the ImageJ “local thresholding” function based on local average; (c) image
 281 obtained from (b) by only keeping in the image the largest connected black region; (d) image
 282 obtained from (c) by replacing each connected white region by its filled complex hull; (e)
 283 Outlines of the resulted connected white regions which are identified as whole bubbles; (f)
 284 Corresponding temporal evolution of the PDF for the decimal logarithm of the equivalent
 285 radius, $a/2$ (equivalent particle diameter a being the equivalent particle size).

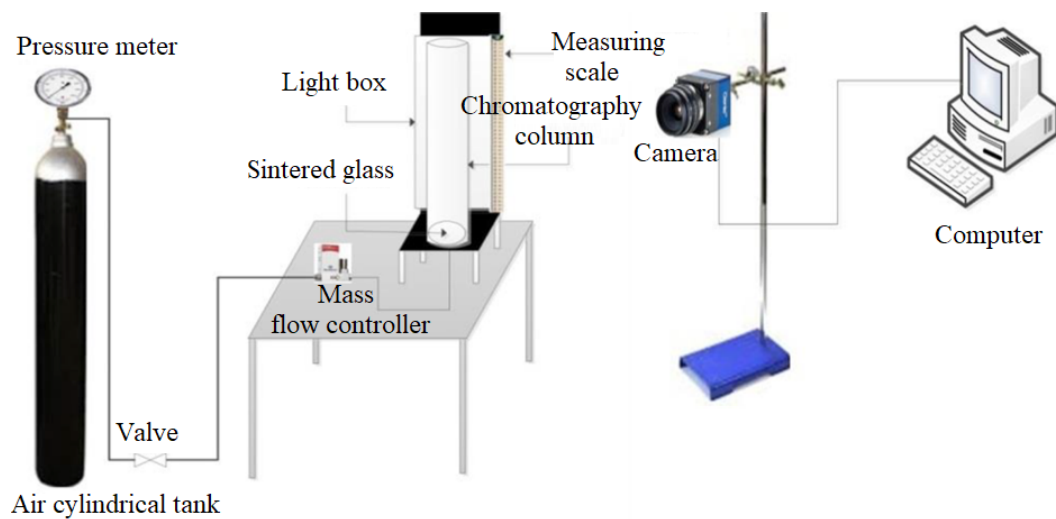
286 From the list of bubble areas, a list of equivalent radii is then obtained as the radii of the disks
287 that have the same area as the bubbles. Statistical measures such as the mean and median
288 values and the standard deviations are computed from these statistics, as well as the
289 probability density functions of the equivalent radii, which is obtained from a histogram.
290 Figure 4f shows the time evolution of the probability distribution function (PDF) of
291 equivalent radii for the experimental run corresponding to Figure 4a-e. Note that since we
292 consider the logarithm of the equivalent radius here, the distribution is becoming wider with
293 time, and this to a considerable extent. The normalization of the PDF takes this into account,
294 which is why its peak decreases with time. The visual impression that the area below the
295 curves is not conserved with time is due to this log-binning of the equivalent radius. In fact, it
296 is conserved, equal to unity at all times.

297 **Bulk-scale experiments:** The column experiments were conducted in a chromatography
298 column (Scientific glass, UK) with an inner diameter of 4 cm and a height of 80 cm,
299 respectively. Figure 5 shows a schematic diagram of the column used in this study. A sintered
300 glass disc with a pore size distribution between 40 and 60 μm was placed at the bottom of the
301 column as a foam generator. The liquid phase for each surfactant was prepared by adding the
302 surfactants to deionized water at their CMC (Table 1) and then mixing using a stirrer (Fisher
303 Scientific, UK) for 2 hours. Silica nanoparticles were added to the solution at various
304 concentrations (0-1%) and mixed for an additional 30 minutes. The experiments were
305 conducted immediately after the solution was prepared to prevent hydrolysis of the
306 surfactants. Air was injected through a tube with an inner diameter of 0.5 cm into the column
307 through the sintered glass using a mass flow controller at 100 ml/min flow rate. The gas flow
308 rate was adjusted using the Flow View and Flow DDE (Bronkhorst, UK) software. Flow
309 [DDE View](#) provides an interface between the computer and the mass flow controller while
310 Flow [DDE View](#) provides the user with manual control of the desired flow rate. The injection

311 was stopped when the column had wholly filled with foam, that is when the foam inside the
312 column had reached a height 80 cm. The liquid then drained from the column by gravity. The
313 drained ~~mass of~~ liquid flowed to a reservoir placed on a balance and its mass was recorded
314 ~~every minute for each experiment with a balance placed close to the column~~ placed closely to
315 ~~the column~~ at different times during each experiment. The drained liquid transferred over the
316 balance for the measurement.



317



318

319

Figure 5: A schematic diagram of the bulk-scale experiment set-up.

320 Each experiment was repeated three times to check the reproducibility. The results presented

321 in the next section will thus be an average of all three tests unless specified otherwise. [All the](#)

322 concentrations of surfactants and nanoparticles are percentage by weight percent (wt%), if
323 not otherwise mentioned.

324 **Results and discussion**

325 Interaction between nanoparticles and surfactants affecting foam stability at bubble

326 scale Synergistic effect of nanoparticles and surfactants on foam stability at bubble scale

327 A series of Hele-Shaw cell experiments were conducted to investigate the synergy between
328 nanoparticles and surfactants in impacting foam properties at bubble scale. In these
329 experiments, gravity-driven foam drainage was negligible since we were working with a
330 quasi-two-dimensional model positioned horizontally.

331 Figure 6 shows the pressure drop measured during foam flooding of the Hele-Shaw cell at
332 different experimental conditions. The pressure drop enables determining the apparent

333 viscosity (Pa.s) of the foam based on the Darcy law: $(\mu_{app} = \frac{KA}{q} \frac{\Delta P}{L})$, where K (m^2) is the

334 permeability of the Hele-Shaw cell, which is close to the theoretical value for infinite planes

335 separated by a distance e , $e^2/12$, e (m)

336 being here the smallest dimension of the cell, q (m^3/s) is the flow rate, A (m^2) is the cross-

337 sectional area, and L (m) is the length of the system). As the imposed flow rate (10 ml/min)

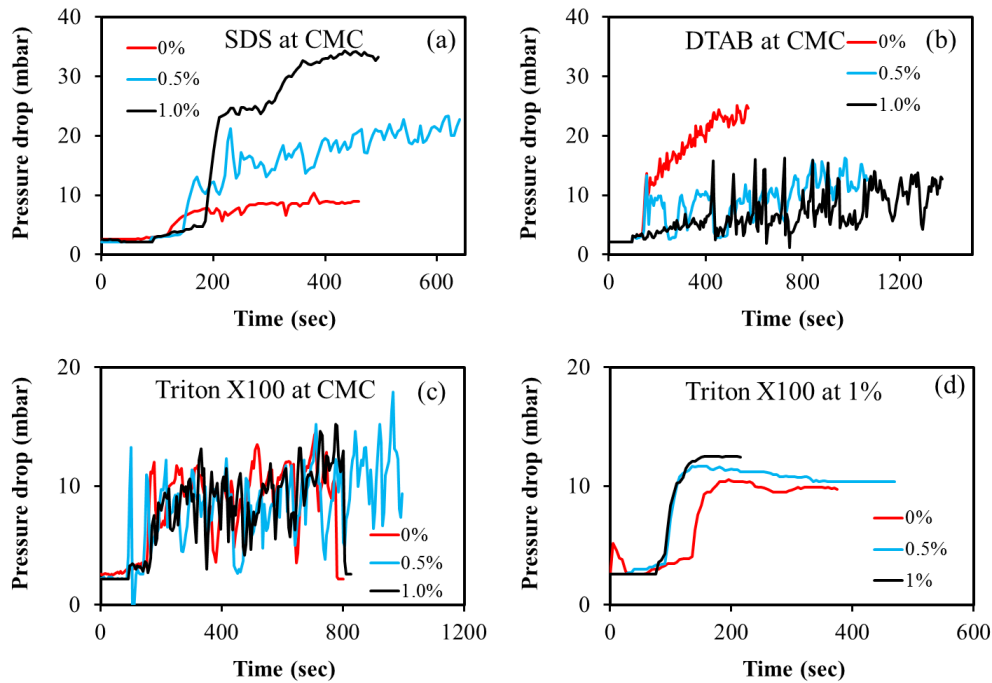
338 and permeability are the same for all measurements here the imposed flow rate is the same

339 for all measurements, we use the pressure drop as a proxy for the apparent/effective

340 viscosity of the foam. Hence the effective viscosity of foam is since the two quantities only

341 differ by a factor which is identical for all measurements proportional to the pressure drop

342 (all other things i.e. permeability, flow rate) being equal.



343

344 Figure 6: Measured pressure drop versus time during foam injection during foam injecting at
 345 different experimental conditions in the Hele-Shaw cell. The legends indicate nanoparticle
 346 concentrations.

347 Figure 6 is already a surrogate for coalescence data. Systems that exhibit lots of coalescences
 348 tend to have unstable pressure drop signals. On the other hand, systems with limited
 349 coalescence have less noisy pressure signals that oscillate less. Figure 6a shows the pressure
 350 drop for the SDS surfactant in the presence of silica nanoparticles. It can be seen that an
 351 increase in silica concentration leads to a larger pressure drop, which can be interpreted as the
 352 generation of a stronger foam. The presence of nanoparticles decreases the surface tension
 353 (Jia et al., 2020; Vatanparast et al., 2018) and increases the strength of the generated foam. In
 354 addition, Figure 6 also gives indication about the level of bubble bursting occurring in the
 355 foam. Given the simple geometry of the Hele-Shaw cell, pressure fluctuations are unlikely to
 356 be related to geometry-related capillary fluctuations, such as could occur inside complex
 357 porous media (Cox et al., 2004; Rossen, 1990). Unstable pressure drop signals are then likely

358related to bursting events. Conversely, systems with limited bursting are expected to exhibit
359less “noisy” pressure signals, with less oscillations.

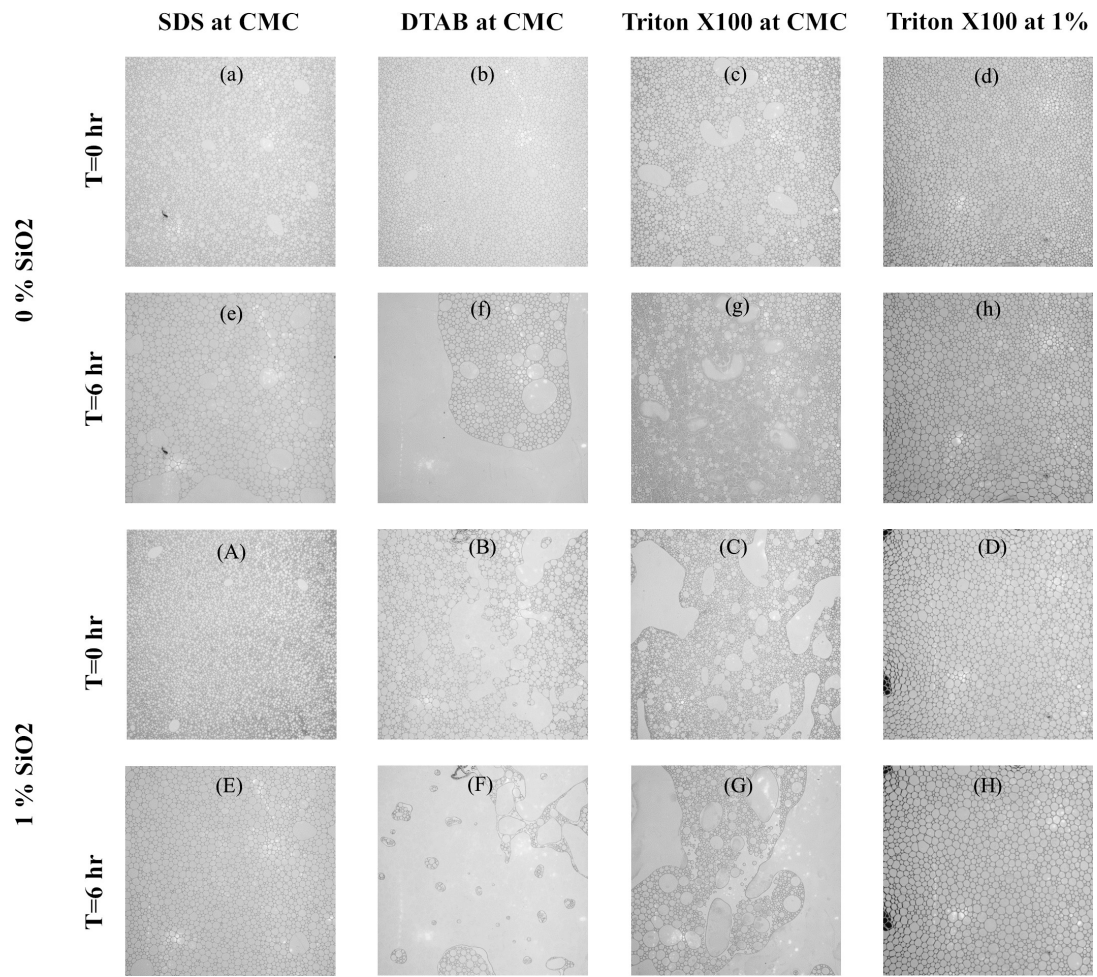
360We stopped the injection of gas and liquid solution at the end of the experiments, and the
361structural evolution of foam bubbles was then monitored in time over the Hele-Shaw cell.
362Although these results were thereby obtained, ~~Although, these results were obtained under~~at
363static conditions, ~~but they could~~an be used as screening experiments prior to flow
364testsscreening experiments for flow tests (Jones et al., 2016; Nasr et al., 2019). Figure 7
365qualitatively shows foam coarsening in the Hele-Shaw cell for eight different foaming
366suspensions, prepared with the three different types surfactants at a concentration equal to
367their CMC (and one surfactant well above CMC)~~for six different foaming suspensions,~~
368prepared with the three different types surfactants at a concentration equal to their CMC and
369with two different concentrations of nanoparticles (0 and 1%). The pictures taken at time t=0
370after the end of the injection show the foam structure at the end of foam generation while
371comparing the picture at t=0, and t=6 h provides information about foam coarsening.
372Comparing Figure 7a with Figure 7Ag, show that the SDS foam has a finer texture in the
373presence of silica nanoparticles. ~~Although, t~~ The foam generator was fixed, i.e. had fixed pore
374sizesthe same for all generated foams (hence, with the same pore size distribution), but the
375bubble size/~~bubble~~ texture produced by a given generator (for specified gas and liquid flow
376rates) could vary from foaming suspension liquid to foaming suspensionliquid. Indeed thea
377finer texture foam could relate toresult from an increase in the maximum capillary pressure of
378coalescence due to the presence of silica nanoparticles, thereby leading to less film breakage
379during foam generationOne reason for smaller bubbles mightcan be attributed to less film
380breakage that leads to smaller bubbles. Another possible reason for finer texture foam could
381be related to an increase in the maximum capillary pressure of coalescence due to the
382presence of silica nanoparticles leading to less film breakage. A finer textured foam when

383 ~~flowing provides a higher pressure drop when flowing through a permeable medium. A finer~~
384 ~~textured foam provides a higher pressure drop~~ and hence results in a higher apparent viscosity
385 of the foam. ~~The other possible reason for finer texture foam could be related to an increase~~
386 ~~in the maximum capillary pressure of coalescence due to the presence of silica nanoparticles.~~

387 In the case of the DTAB surfactant, on the contrary, the presence of nanoparticles results in a
388 decrease of the foam's apparent viscosity, as shown in Figure 6b. ~~This is due, as we have~~
389 ~~said discussed above, to~~ ~~This is due to~~ the adsorption of the cationic surfactants onto the silica
390 nanoparticles, which promotes flocculation of the suspension, that is, phase separation of the
391 solution between the flocculated/sedimented phase and the liquid phase, as seen in Figure 13b
392 in the appendix. Consequently, less surfactant will be available in the solution for strong
393 foam generation. ~~Comparing~~ Figure 7b with Figure 7hB shows coarser foam bubbles were
394 generated in the presence of SiO₂. Higher fluctuation in pressure drop curves for a larger
395 concentration of nanoparticles is also an indication of a more marked instability of the foam.
396 ~~This is due to It could be due to the existence of a flocculated phase of colloids which does~~
397 ~~not easily contribute to foam formation when in contact with the surfactant solution and~~
398 ~~gas, the existence of a flocculated phase of colloids that sedimented out of the surfactant~~
399 ~~solution and did not enter the foam generator. This leads to that does not easily contact foam~~
400 ~~generator and leads to low foam generation does not occupy the volume of liquid phase~~
401 ~~homogeneously.~~

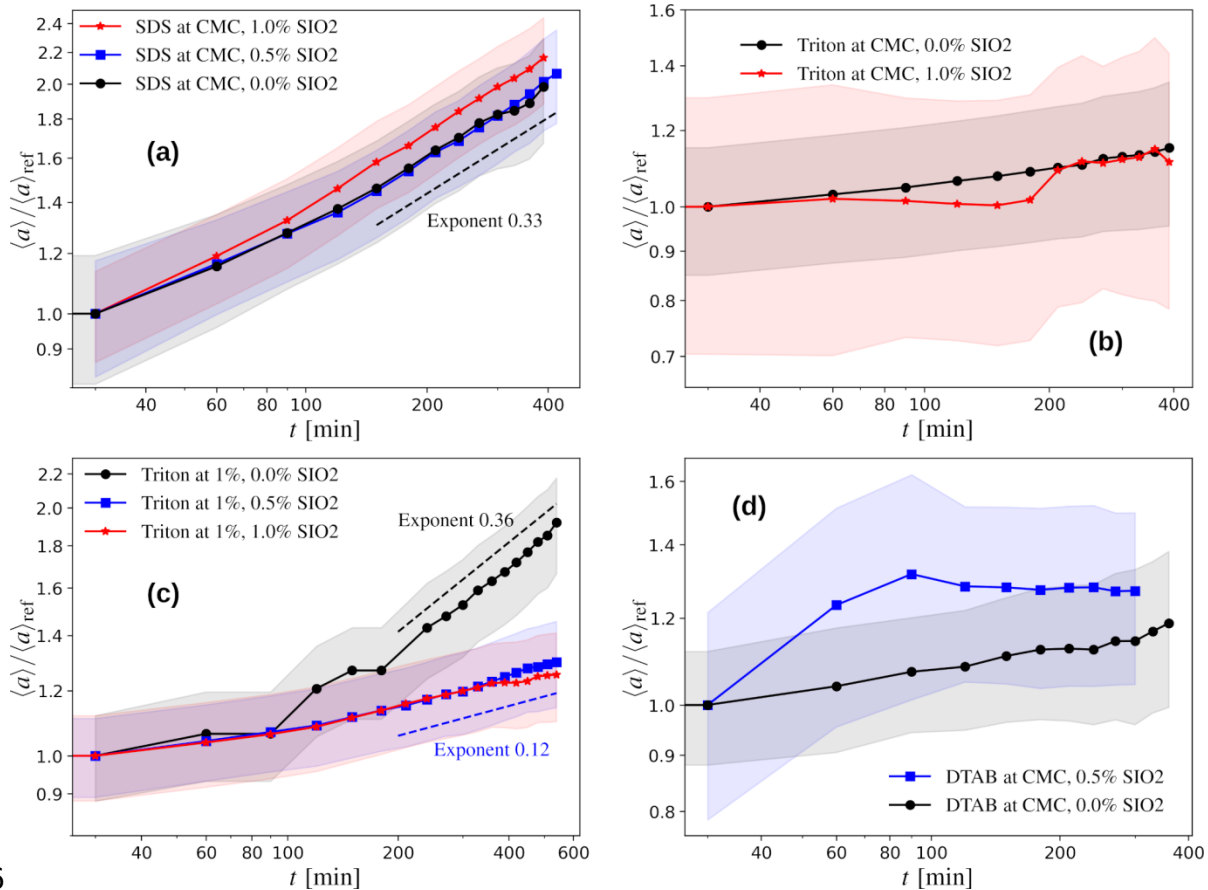
402 ~~Figure 6c indicates that for Triton X100 at CMC~~ ~~Figure 6e~~ ~~the foam is generally unstable~~
403 ~~both in the absence and presence of nanoparticles~~ ~~indicates that the foam is generally unstable~~
404 ~~in the absence and presence of nanoparticles for Triton X100~~, as indicated by the strong
405 fluctuations, whose amplitude is not impacted by the concentration in SiO₂ nanoparticles.
406 This might be due to the low molar concentration of Triton X100 at its CMC (Table 1).
407 ~~Generation of even a limited amount of foam can then reduce the concentration in the bulk~~

408 solution below the CMC, which can affect ongoing foam generation (Boos et al., 2012). This
409 effect can be intensified in the presence of nanoparticles given the low CMC value molar
410 concentration of Triton X100 at the CMC (see Table 1). ~~The issue is that the CMC~~
411 ~~corresponds to a low molar concentration and mass concentration (compared to the other~~
412 ~~surfactant), meaning that "losing" a given mass of Triton to adsorption on the particles, can~~
413 significantly impact a great deal on ~~significantly~~ the concentration remaining in the bulk
414 solution. ~~To validate this point, we performed experiments at a higher concentration (1.0 %)~~
415 ~~of Triton X100 with varying concentrations of silica nanoparticle~~ Consequently, sw We also
416 performed experiments at a higher concentration (1.0 %) of Triton X100 with varying
417 concentrations of silica nanoparticles to investigate the impact of the concentration of Triton
418 on the foam's stability ~~validate this point~~. The corresponding temporal evolution of the
419 pressure drop across the flow cell is presented in Figure 6d. Comparisons between Figure 6c
420 and Figure 6d, with the plots in Figure 6d appearing much smoother than those in Figure 6c,
421 shows that an increase in the surfactant concentration improved foam stability ~~shows that an~~
422 ~~increase in the surfactant~~ concentration improved foam stability tremendously. These
423 findings suggest that surface tension (which tends to remain fixed above the CMC) is not the
424 only physical quantity- controlling foam stability and foam generation and that CMC may not
425 be the optimal concentration to generate the most stable foams in the case of surfactants with
426 an extremely low CMC, in particular in the case of attractive interaction between the
427 nanoparticles and the surfactant. Note also that once the quantity of surfactant available for
428 the fluid-gas interfaces is sufficiently large (Figure 6d), the addition of SiO₂ nanoparticles
429 slightly enhances the foam's apparent viscosity, but to a significantly lesser extent than what
430 is observed in Figure 6a for the SDS-based foaming solutions.



431

432 *Figure 7: Bubbles ~~observed for different foaming suspensions~~ observed for six different*
 433 *foaming suspensions, corresponding to the three types of surfactants and two different SiO₂*
 434 *concentrations of 0 and 1%, at two successive times 0 hr and 6 hr after the flow has been*
 435 *stopped in the Hele-Shaw cell. ~~All surfactants are at a concentration equal to their CMC.~~*



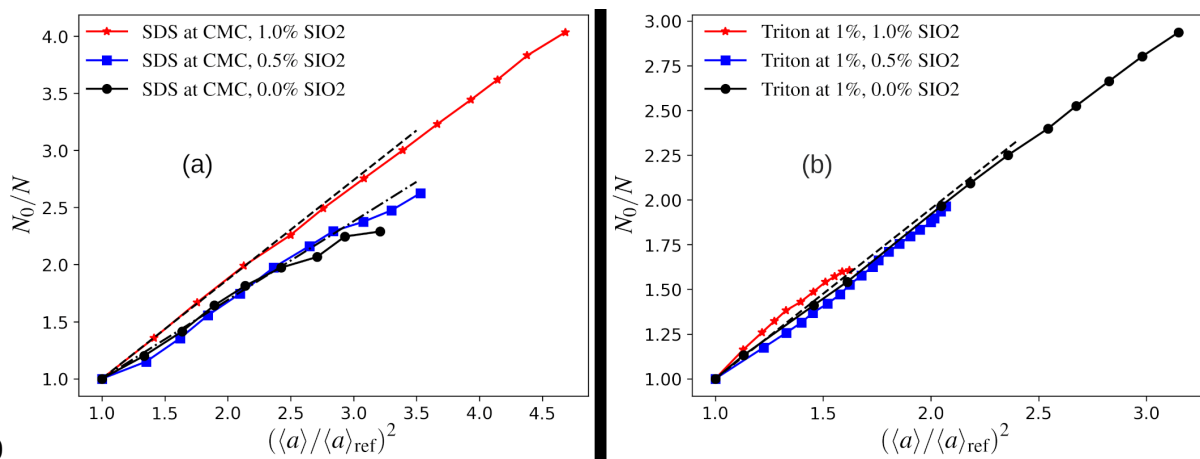
436

437 Figure 8: Ratio of the mean *initial*-equivalent bubble radius, $\langle a \rangle / 2$, to its value at time $t = 30$
 438 min, plotted as a function of time for different foaming suspensions prepared with the three
 439 types of surfactant and different concentrations of SiO_2 nanoparticles, as presented in the
 440 legends. Each coloured area corresponds to the tolerance interval of the curve of identical
 441 colour, defined as having a vertical extent equal to half the standard deviation of $\langle a \rangle / \langle a \rangle_{\text{ref}}$.

442 Image analysis performed from images such as presented in Figure 7 allowed us to extract the
 443 probability density functions (PDFs) of bubble sizes, as explained in the “Methods” section
 444 above. Figure 8 summarizes the behaviour observed for the mean *and standard deviations of*
 445 these PDFs, for four types of foaming solutions and, depending on the case
 446 considered depending on the foaming solution, for two or three concentrations of the
 447 nanoparticles. definition, for two or three concentrations of the nanoparticles. This quantity
 448 is mainly a measure of (likely to be associated with bubble coarsening by both gas diffusion

449 and film breakage insofar as the latter results in a coalescence of two bubbles into a larger
 450 one. Bubble bursting would be expected to lead to size distributions where a few larger
 451 bubbles exist in the midst of smaller bubbles (see for example *Figure 7c* and *7g*), but
 452 distinguishing between coarsening by gas diffusion and by film breakage was made difficult
 453 by the initial polydispersity of the foams. Note also that in our image treatment (see *Figure*
 454 44), the bubbles that touch the boundaries of the field of view are not taken into account in the
 455 statistics, because their real size may not be (and in most cases, is not) captured entirely in
 456 our image. For the DTAB-based foams and those based on Triton at CMC, the foam is
 457 observed to burst over large areas starting from the boundaries of the domain until a large
 458 part of the domain corresponds to the result of that “catastrophic” bursting (see *Figure 6f*, *6F*,
 459 *6g* and *6G*). The measure of the mean bubble size is insensitive to large scale bubble bursting
 460 is blind to they may include some bubble coalescence. However, large-scale this large scale
 461 bubble bursting since is not included in this large boundary-touching voids are excluded from
 462 the calculation calculation. The total number of bubbles N The total number of bubbles
 463 measured by the image treatment, on the contrary, decreases strongly due to that large scale
 464 bursting. For foams in which no such catastrophic bursting is visible, the squared mean
 465 bubble size $\langle \sigma^2 \rangle$ (which does not differ much from $\langle \sigma^2 \rangle$) and the inverse number of bubbles,
 466 when normalized by their initial value, are supposed to be more or less equal to each other
 467 since the sum of all bubble areas is not very different from the total domain area. Indeed,
 468 *Figure 9* confirms a linear relationship between the two quantities for the SDS-based foams
 469 and the foams based on Triton at 1%. The slope is not exactly 1, probably because the
 470 apparent area of the lamellae has been neglected in the above argument, but one can safely
 471 conclude that, for these foams, the two quantities (mean bubble size and number of bubbles)
 472 contain the same information. For foams based on Triton at CMC with nanoparticles and
 473 those based on DTAB, on the contrary, the information on the mean bubble size and standard

474 deviation of the bubble size PDFs is not conclusive without additional information on the
 475 number of bubbles.
 476 coalescence may be going on also, but ~~note~~ also that to a certain extent, the pressure
 477 fluctuations in Figure 6 are already indicative of large scale bursting-proxy measurement for
 478 it, which complement the present ~~we look instead at diffusive coarsening in what follows~~
 479 measurement of bubble coarsening.



480

481 Figure 9: Relationship between N_0/N (N_0 being the number of bubbles at $t=30$ min) and $\langle \rangle$
 482 normalized by its value at $t=30$ min, for (a) SDS-based foaming suspensions, and (b) foaming
 483 suspensions based on Triton at 1%.

484

485 Returning to consider Figure 8a, this presents the results ~~Figure 8a presents the results~~ for a
 486 foaming solution consisting of SDS at the critical micelle concentration (CMC). For the three
 487 nanoparticle concentrations (0.0, 0.5 and 1.0%), the evolution in time of the bubbles' mean
 488 radius $\langle a \rangle / 2$, recorded over a duration of about 400 min and normalized by its value at $t=30$
 489 min, shows that all three curves tend to follow power laws of exponent 0.33 after $t=150$ min,
 490 but with a prefactor which is about 10% larger for the largest concentration in SiO₂. The
 491 dispersion of the PDF around the mean follows a similar behaviour, proportional to the mean.

492 This seems to indicate that adding nanoparticles at these concentrations does not provide any
493 limiting effect on foam coarsening for SDS-based foams.

494

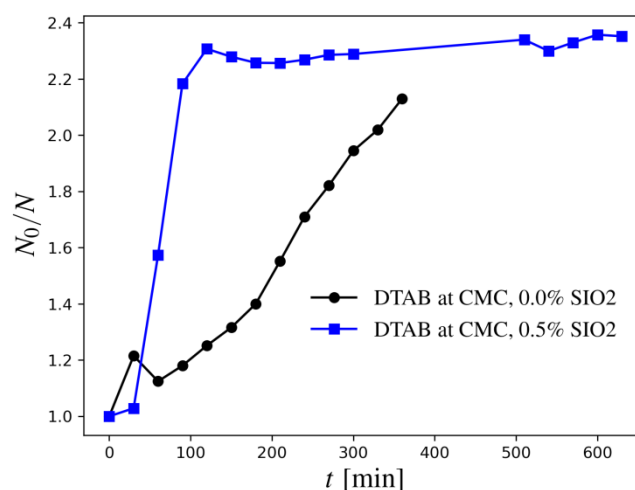
495 In the case of Triton X100, Figure 8b and Figure 8c reveal that a change in the concentration
496 of the surfactants can radically modify their impact on foam coarsening and coalescence. ~~their~~
497 ~~impact on foam coarsening~~. This can clearly be associated with the combined effects of the
498 extremely low CMC value of Triton X100 and of its interaction with the silica nanoparticles
499 in line with the result of another study (Martinez et al., 2008). Indeed, at low concentration of
500 Triton X100, ~~most-much of the surfactant adsorb~~~~of the surfactants adsorb~~ onto the surface of
501 the silica particles, leaving little surfactant deposition at the gas-liquid interface. This is
502 detrimental to foam stability and foam generation, as discussed above in relation to Figure 6
503 and Figure 7. ~~Note however~~ the mean bubble radius evolves in a similar manner in the
504 absence of SiO₂ nanoparticles as in their presence at 1.0% for Triton X100 at CMC, though
505 the visual observation of the two bubble populations shows two very different behaviours: in
506 the presence of the nanoparticles, the aforementioned many bubbles burst, leading to the
507 appearance of liquid films on the horizontal glass plates, which hide part of the bubbles
508 population on the image large scale bursting from the domain boundaries occurs; on the
509 contrary, in the absence of nanoparticles, the bubbles evolve through diffusion-controlled
510 coarsening, with bubble sizes evolving in time but few of them bursting within the
511 experimental time duration. Figure 9(b) simply means confirms, as discussed above, that the
512 mean bubble size of surviving bubbles is not, when comparing these cases, a good measure of
513 the foam stability when large scale bursting occurs. Accordingly, the N_0/N plots (not shown
514 here), show a much steeper increase in the presence of SiO₂ nanoparticles, as a consequence
515 of the large scale bursting, than in the absence of SiO₂. We can conclude from this data that
516 the affinity of Triton with the nanoparticles renders its use at CMC ineffective to study the

517 impact of added nanoparticles on the foam's stability. ~~It worth mentioning, Figure 8~~
518 ~~focuses presents diffusive coarsening behaviour (not coalescence) on the diffusive coarsening~~
519 ~~behaviour, because of the particular way we defined the bubble radius $\langle a \rangle / 2$ (i.e. deliberately~~
520 ~~omitting large voids), what we are presenting with Figure 8 is the diffusive coarsening (not~~
521 ~~coalescence) behaviour. Regarding presenting foam coalescence, Figure 6 already gives a~~
522 ~~"surrogate" for coalescence data. Systems that exhibit lots of coalescences tend to have~~
523 ~~unstable pressure drop signals. On the other hand, systems with limited coalescence have less~~
524 ~~noisy pressure signals that oscillate less, and we highlighted this point in the manuscript.~~

525 For Triton X100 at 1%, on the contrary, there is enough surfactant for it to be present at
526 liquid-gas interfaces while also adsorbing onto the nanoparticles, as discussed above in
527 relation to Figure 6c and Figure 6d. Foam coarsening is then observed to be strongly
528 impacted by the addition of SiO₂ nanoparticles. Coarsening of the foam prepared with the
529 suspension devoid of nanoparticles exhibits a power-law growth of the mean bubble size, of
530 exponent 0.36. If nanoparticles are added to the foaming suspensions, this power-law
531 behaviour has an exponent 0.12, which is identical for concentrations of 0.5% and 1.0% in
532 SiO₂. This is consistent with the observations of Figure 6d, showing that the foam's effective
533 viscosity is larger as the concentration in SiO₂ nanoparticles is larger. ~~However, the impact~~
534 ~~on foam stability against of diffusive coarsening and bubble coalescence over time~~ However,
535 ~~the impact on foam stability over time~~, demonstrated here, is more spectacular than the
536 impact on its effective viscosity.

537 In contrast, according to Figure 8d, foam coarsening becomes faster in DTAB-based
538 suspensions as the concentration in silica nanoparticles is larger. This is likely due to the
539 interaction between DTAB and SiO₂, which results in flocculation of the suspensions, and
540 therefore in a decrease of the number of nanoparticles available for the liquid-gas interfaces,
541 as discussed earlier. The curve for 0.51% SiO₂ shows a fast initial increase of the mean

542 bubble size, followed by a plateau. This plateau is somewhat misleading as it results from the
 543 disappearance of larger bubbles by bursting, which leaves only smaller bubbles whose size
 544 does not evolve much to contribute to the mean bubble size. This is confirmed by the time
 545 evolution of the normalized inverse number of bubbles, N_0/N (Figure 10), in which the initial
 546 rise is much steeper than for the foam prepared without SiO_2 nanoparticles, but soon reached
 547 a plateau. This corresponds to extremely fast catastrophic bursting from the boundaries,
 548 leading to a configuration where the bubbles left are essentially round and isolated (which
 549 removes the possibility of coarsening by gas diffusion. This plateau value in Figure 8d); the
 550 bubble number then slowly evolves under additional slow bursting of these isolated bubbles,
 551 which explains the plateau in Figure 10, but bubble sizes hardly change any more, which
 552 explains the plateau value in Figure 8d. This plateau is, however, corresponds to a bubble
 553 size larger than the mean bubble size measured during the evolution of the foam which is
 554 devoid of SiO_2 particles. In a 2-d Hele-Shaw cell geometry, as gas escapes from (but liquid is
 555 retained by) a foam that is no longer connected to the cell walls, thus the effective liquid
 556 fraction of the foam rises over time, and diffusive coarsening is expected to slow as a result
 557 (Furuta et al., 2016).

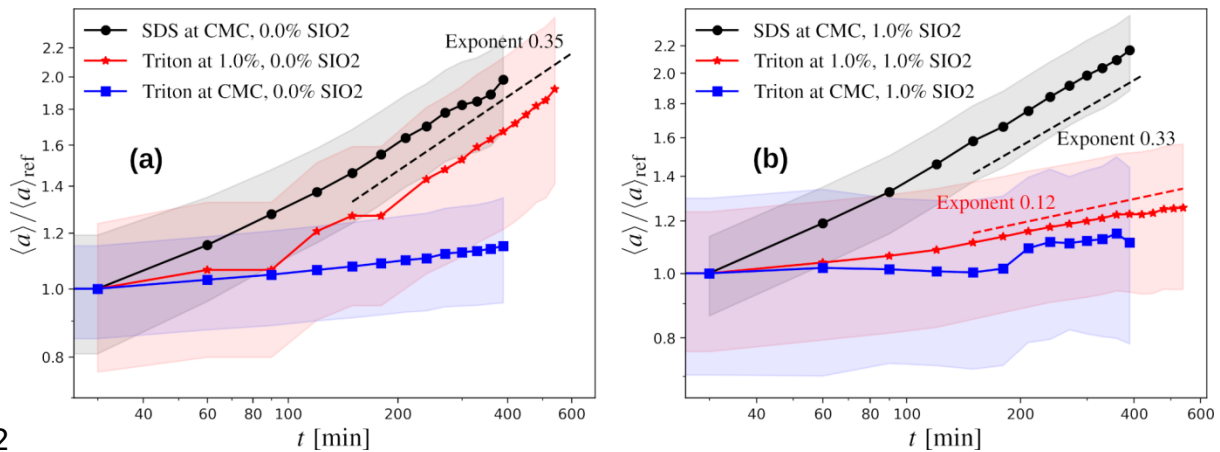


558

559 Figure 10: Time evolution of the ratio of the initial number of bubbles to the current one, for the foams prepared with DTAB
 560 at CMC.

561

562 In *Figure 11* we have plotted the time evolution of the mean bubble radius for SDS and
563 Triton-based foams, grouping in *Figure 11a* all the data obtained with foams devoid of SiO₂
564 nanoparticles, and in *Figure 11b* all the data obtained with foaming solutions containing the
565 nanoparticles at a 1.0% concentration in weight. *Figure 11a* shows that, in the absence of
566 SiO₂ particles, the foams prepared with SDS at CMC and Triton at 1.0% behave in the same
567 manner, while that prepared with Triton at CMC ages more slowly. When SiO₂ nanoparticles
568 are present at a concentration of 1% in weight, the aging of the foams based on SDS ~~and~~
569 ~~Triton at CMC are not~~ is not much impacted, while that of the foam prepared with Triton at
570 1% is slowed down considerably. For Triton at CMC, the mean bubble size does not vary
571 much, but -the discussion above has shown that that it is actually an artefact of the measure
572 used to characterize bubble size during ageing, since bubble bursting is not accounted for this
573 quantity is simply not a relevant measure of foam aging in this configuration, since the
574 number of bubbles decreases dramatically due to large scale bubble bursting. ~~However,~~ for
575 a sufficiently large concentration of Triton (such as 1%wt), however, the addition of
576 nanoparticles improves the foam stability, while it has little impact on an SDS-based foam.
577 Recall that adding nanoparticles to a 1% Triton solution led to a large increase in viscosity
578 (see Figure 2). This is expected to reduce gas diffusivity through films (hence reducing
579 diffusive coarsening) and also to reduce the rate at which films break (hence reducing
580 coalescence). This may explain the slower coarsening seen when nanoparticles are added to
581 1% Triton solution.



582

583 Figure 11: Ratio of the mean initial equivalent bubble radius, $\langle a \rangle / 2$, to its value at time $t = 30$
 584 min, plotted as a function of time for different foaming suspensions prepared with the three
 585 types of surfactant and either 0% (a), or 1% (b) of SiO_2 nanoparticles, as presented in the
 586 legends. Each coloured area corresponds to the tolerance interval of the curve of identical
 587 colour, defined as having a vertical extent equal to half the standard deviation of $\langle a \rangle / \langle a \rangle_{ref}$.

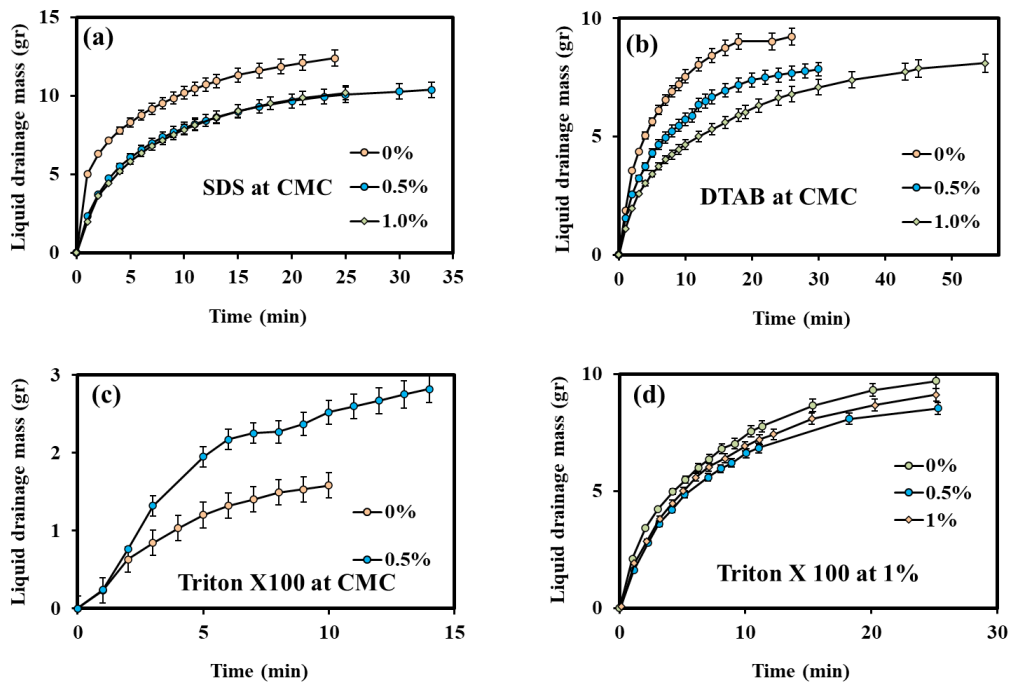
588

589 Interaction between nanoparticles and surfactants affecting liquid drainage Synergistic 590 impact of Nanoparticles and Surfactants on Liquid Drainage

591 The duration of the column experiments is between 15 min and 1hr (see [Figure 12](#) Figure 12).
 592 This duration is nearly one order of magnitude smaller than the time scales which are
 593 characteristic of foam coarsening, as probed by the Hele-Shaw experiments. [Hence the](#)
 594 [column experiments investigate mostly the effect of gravitational drainage on foam stability,](#)
 595 [rather than that of and not diffusive coarsening and bubble coalescence](#) Hence the column
 596 [experiments investigate mostly the effect of gravitational drainage on foam stability.](#)

597 [Figure 12](#) Figure 12-presents the mass of drained liquid measured at the bottom of the column
 598 for 11 different foaming suspensions. [The end of the experiments was](#) ~~ere defined~~ [terminated at](#)
 599 [the state when either there was no further change in the liquid drainage was visible, or all the](#)
 600 [foam bubbles inside the column had collapsed. Typically the former situation was observed](#)

601 [with experiments corresponding to Figure 12s \(a\), \(1012b\), 1012and \(d, whereas\) and the](#)
 602 [latter situation was observed with experiments corresponding to Figure 12c.](#)



603

604 Figure 12: Liquid drainage over time for different foaming solutions based on SDS, DTAB,
 605 and Triton X100 surfactants (the latter at two different concentrations). The legends indicate
 606 nanoparticles concentration. Foam generation did not occur for Triton X100 at its CMC and 1
 607 % concentration of SiO₂.

608 Figure 12 [Figure 12 a](#) for SDS shows that the drained liquid mass at any given time decreases
 609 with the addition of the SiO₂ nanoparticles at 0.5 wt% compared to the same surfactant
 610 solution devoid of nanoparticles. This is probably due in part to the increase in the solution's
 611 viscosity resulting from the presence of the nanoparticles. However, a further increase in
 612 nanoparticle concentration from 0.5 to 1% results in no significant changes in the liquid
 613 drainage rate: the effect saturates. [This indicates that the slower drainage is also related to](#)
 614 [the occupation of the gas-liquid interfaces by the nanoparticles, an effect that is likely to](#)
 615 [saturate at sufficiently high enough concentration of NPs. Covering interfaces with NPs](#)

616 ~~decreases their surface tension~~ This indicates that the slower drainage is also related to the
617 ~~occupation of the gas-liquid interfaces by the nanoparticles, which decreases the surface~~
618 ~~tension~~ (Jia et al., 2020; Vatanparast et al., 2018) thus rendering them more stable and
619 delaying their bursting due to lamella-thinning. Another possible reason for stabilising
620 behaviour is that the nanoparticles can also slow the drainage down by blocking rendering the
621 Plateau borders less permeable to liquid rendering ~~les permeable~~ (Carn et al., 2009).

622 In the case of DTAB, addition of the nanoparticles lowers the rate of liquid drainage
623 significantly, and liquid drainage becomes slower as the concentration in silica nanoparticles
624 is larger (see *Figure 12* ~~Figure 12~~ 12b). This is believed to be due ultimately to the adsorption
625 of DTAB ~~This is due to the adsorption of DTAB~~ surfactants on the silica particles, as shown
626 in *Figure 1*, which leads to a decrease in the electrostatic repulsion between the nanoparticles,
627 and thus to flocculation, as discussed earlier. Hence, after the solution was poured into the
628 column, the particle-surfactant complex precipitated at the bottom of the column on the foam
629 generator. This interaction between DTAB and silica nanoparticle increases the viscosity of
630 the complex fluid considerably. This flocculated part of the mixture contributes to the largest
631 part of the foam generation since it is where air first contacts the solution. This highly viscous
632 solution present in the lamellae and Plateau borders decelerates liquid drainage.

633 In the case of Triton X100, *Figure 12* Figure 12c and *Figure 12* Figure 12 d suggest that the
634 effectiveness of silica nanoparticles to generate foams which are less prone to collapsing
635 under gravitational drainage depends on the concentration of the surfactant, as was the case
636 for the Hele-Shaw cell experiments. *Figure 12* Figure 12c for Triton at CMC shows that the
637 drained liquid mass measured at any given time increases with the concentration of
638 nanoparticles. For the 1% SiO₂ concentration, foam generation hardly occurred due to
639 adsorption of ~~most~~ much of the surfactant of the nanoparticles, as discussed at length above
640 in relation to Figures 6, 7, and 8 ~~and 9~~; hence we have not included the corresponding data in

641 *Figure 12* [Figure12c](#). At a concentration of Triton X100 of 1%, *Figure 12* [Figure12d](#)
642 suggests that a 0.5% concentration in nanoparticles provide higher stability against
643 gravitational drainage than 0 and 1% concentrations. This means that a further increase in
644 silica nanoparticle concentration from 0.5% led to faster liquid drainage, possibly due to a
645 saturation of the effect related to occupation of liquid-gas interfaces by nanoparticles. [Note](#)
646 [that in the case of Interestingly foam is found to have a higher drainage rate with a higher](#)
647 [concentration of Triton X 100s at CMC \(eComparing Figure12-\(c\), comparatively little drains](#)
648 [out -with-\(d\)- because the initial volume of generated foam is far from reaching the entire](#)
649 [volume of the cylindrical cell unlike the other cases.at the higher concentration was much](#)
650 [larger than the lower concentration. In other words, despite its high viscosity, Triton drains](#)
651 [faster at a higher concentration as we already explained compared to Triton at CMC and](#)
652 [more liquid enters to the foam network in the first place. Hence, as Triton at CMC forms a](#)
653 [relatively small amount of little foam, there is relatively little liquid to drain out of the foam.](#)

654 **Summary and conclusion**

655 We have presented an investigation of foam stability using surfactants with different charges
656 (anionic, cationic and non-ionic) in the presence of charge-stabilized silica (SiO_2)
657 nanoparticles. A comprehensive series of experiments were conducted using a horizontal
658 Hele-Shaw cell and columnar flow cells. Hele-Shaw cell experiments are typically termed
659 'bubble scale' experiments in the literature (Osei-Bonsu et al., 2016); in our study, they
660 mostly probed the foam's [instability by coarsening through gas diffusion or else by bubble](#)
661 [bursting instability by coarsening through gas diffusion and bubble bursting](#). Columns
662 experiments are typically termed 'bulk scale' experiments; more importantly, they mostly
663 probe the foam instability by gravitational drainage.

664 For foams prepared with the anionic surfactants SDS (which do not adsorb on SiO_2
665 nanoparticles), the presence of the nanoparticles increased foam stability with respect to

666foam apparent viscosity (i.e., measured pressure drop during flow) and with respect to
667gravitational drainage but had little impact on foam coarsening by diffusion. In the case of a
668foaming suspension prepared with the cationic surfactant DTAB, the presence of oppositely
669charged nanoparticles leads to flocculation and sedimentation of the nanoparticles, which
670removes surfactant adsorbed on the particles from the solution. Consequently, the foam is
671less stable, ~~both in terms of gravitational drainage in the columns and at least in terms of~~
672~~coarsening in the Hele-Shaw cells. Apparent viscosity, i.e. which is inferred from the~~
673~~measured pressure drop signal during flow through the Hele-Shaw cells, also fluctuates a~~
674~~great deal for DTAB with nanoparticles, suggesting poor foam stability in terms of~~
675~~coarsening in the Hele-Shaw cells.~~

676For foaming suspensions prepared with the surfactant Triton X100, which adsorbs on the
677SiO₂ nanoparticles but to a lesser extent than DTAB, the concentration in surfactant should
678significantly exceed the CMC so that enough surfactant is present at the liquid-gas interfaces
679to generate stable foam. Of course, the amount of surfactant needed depends on the
680concentration of nanoparticles. Once this requirement is met, our findings suggest that there
681exists a concentration of nanoparticles that allows slowing down gravitational drainage
682optimally, whereas the addition of even more nanoparticles is all the more beneficial in terms
683of limiting foam coarsening by diffusion when the concentration in nanoparticles is larger.
684Therefore, finding the formulation of the foaming suspension, which is optimal in terms of
685global stability of the foam is not straightforward. Compatibility experiments between
686surfactant and nanoparticles are pre-requisite to optimizing foam stability.

687The prospects of this study include similar experiments performed within porous media. In
688addition, since in deep geological formation solutions are often strongly [alkalinesaline](#), one
689can wonder how these results would be impacted when considering foaming suspensions in
690saline solutions. An increase in salt concentration will shrink the electrical double layer's

691thickness and thus favour nanoparticle attractive interactions and flocculation. Hence, the
692balance between the various forces at play will be displaced when increasing the salt
693concentration, but we expect most of the phenomenology to be similar. We shall test these
694hypotheses on the effect of salinity in future studies, as well as investigate the impact of high
695temperatures.

696**Acknowledgements**

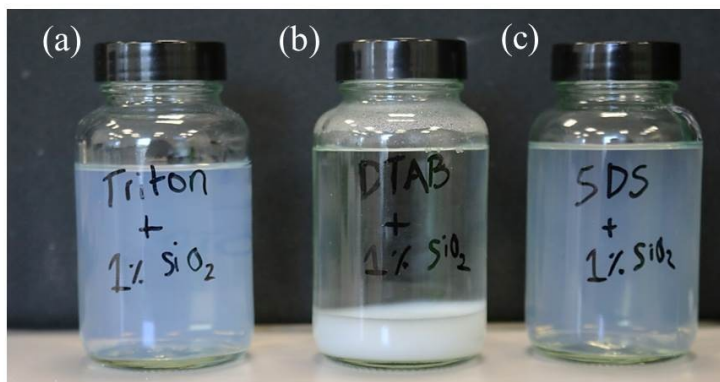
697We would like to acknowledge the UK Engineering and Physical Sciences Research Council
698(EPSC) for providing the Ph.D. studentship for Mohammad Javad Shojaei. All the
699experiments related to the analysis of foam stability at the bubble- and bulk-scale were
700conducted at The University of Manchester. The authors are grateful to Dr. Sajad Kiani for
701conducting rheology measurements.

702

703-

704 Appendix A: Visual aspect of the various foaming solutions with nanoparticles

705 Three bottles containing foaming suspensions consisting of SiO_2 nanoparticles at 1% by
706 weight suspended in solutions of the surfactants Triton, DTAB and SDS, respectively, are
707 shown in *Figure 13*.



708

709 *Figure 13: Visual aspect of surfactant solutions (a) Triton X100 (b) DTAB and (c) SDS*
710 *containing silica particles (1%).*

711References

712Uncategorized References

- 713AttarHamed, F., Zoveidavianpoor, M., Jalilavi, M., 2014. The incorporation of silica
714nanoparticle and alpha olefin sulphonate in aqueous CO₂ foam: Investigation of
715foaming behavior and synergistic effect. *Petroleum science and technology* 32,
7162549-2558.
- 717Benson, S.M., Cole, D.R., 2008. CO₂ sequestration in deep sedimentary
718formations. *Elements* 4, 325-331.
- 719Bhakta, A., Ruckenstein, E., 1997. Decay of standing foams: drainage,
720coalescence and collapse. *Advances in Colloid and Interface Science* 70, 1-124.
- 721Bi, Z., Liao, W., Qi, L., 2004. Wettability alteration by CTAB adsorption at surfaces
722of SiO₂ film or silica gel powder and mimic oil recovery. *Applied Surface Science*
723221, 25-31.
- 724Blijdenstein, T., De Groot, P., Stoyanov, S., 2010. On the link between foam
725coarsening and surface rheology: why hydrophobins are so different. *Soft Matter*
7266, 1799-1808.
- 727Blunt, M., Fayers, F.J., Orr Jr, F.M., 1993. Carbon dioxide in enhanced oil recovery.
728*Energy Conversion and Management* 34, 1197-1204.
- 729Boos, J., Drenckhan, W., Stubenrauch, C., 2012. On how surfactant depletion
730during foam generation influences foam properties. *Langmuir* 28, 9303-9310.
- 731Carn, F., Colin, A., Pitois, O., Vignes-Adler, M., Backov, R., 2009. Foam drainage in
732the presence of nanoparticle– surfactant mixtures. *Langmuir* 25, 7847-7856.
- 733Chang, Y.-B., Lim, M., Pope, G., Sepehrnoori, K., 1994. CO₂ flow patterns under
734multiphase flow: heterogeneous field-scale conditions. *SPE Reservoir Engineering*
7359, 208-216.
- 736Chen, S., Hou, Q., Zhu, Y., Wang, D., Li, W., 2014. On the origin of foam stability:
737understanding from viscoelasticity of foaming solutions and liquid films. *Journal*
738*of dispersion science and technology* 35, 1214-1221.
- 739Cox, S., Neethling, S., Rossen, W., Schleifenbaum, W., Schmidt-Wellenburg, P.,
740Cilliers, J., 2004. A theory of the effective yield stress of foam in porous media:
741the motion of a soap film traversing a three-dimensional pore. *Colloids and*
742*Surfaces A: Physicochemical and Engineering Aspects* 245, 143-151.
- 743Derjaguin, B., Landau, L., 1993. Theory of the stability of strongly charged
744lyophobic sols and of the adhesion of strongly charged particles in solutions of
745electrolytes. *Progress in Surface Science* 43, 30-59.
- 746Desarnaud, J., Bonn, D., Shahidzadeh, N., 2016. The pressure induced by salt
747crystallization in confinement. *Scientific reports* 6, 1-8.
- 748Espinoza, D.A., Caldelas, F.M., Johnston, K.P., Bryant, S.L., Huh, C., 2010.
749Nanoparticle-stabilized supercritical CO₂ foams for potential mobility control
750applications, *SPE Improved Oil Recovery Symposium*. Society of Petroleum
751Engineers.
- 752Exerowa, D., Kruglyakov, P.M., 1997. *Foam and foam films: theory, experiment,*
753*application*. Elsevier.
- 754Feng, C., Kong, Y., Jiang, G., Yang, J., Pu, C., Zhang, Y., 2012. Wettability
755modification of rock cores by fluorinated copolymer emulsion for the
756enhancement of gas and oil recovery. *Applied Surface Science* 258, 7075-7081.
- 757Furuta, Y., Oikawa, N., Kurita, R., 2016. Close relationship between a dry-wet
758transition and a bubble rearrangement in two-dimensional foam. *Scientific*
759*reports* 6, 1-8.
- 760Garcia, J.E., Pruess, K., 2003. *Flow instabilities during injection of CO₂ into*
761*salineaquifers*. Ernest Orlando Lawrence Berkeley NationalLaboratory, Berkeley,
762CA (US).

763Hirasaki, G., Miller, C., Szafranski, R., Lawson, J., Akiya, N., 1997a.
764Surfactant/foam process for aquifer remediation, International symposium on
765oilfield chemistry. Society of Petroleum Engineers.
766Hirasaki, G., Miller, C., Szafranski, R., Tanzil, D., Lawson, J., Meinardus, H., Jin, M.,
767Londergan, J., Jackson, R., Pope, G., 1997b. Field demonstration of the surfactant/
768foam process for aquifer remediation, SPE Annual Technical Conference and
769Exhibition. Society of Petroleum Engineers.
770Hirasaki, G.J., Lawson, J., 1985. Mechanisms of foam flow in porous media:
771apparent viscosity in smooth capillaries. Society of Petroleum Engineers Journal
77225, 176-190.
773Horozov, T.S., 2008. Foams and foam films stabilised by solid particles. Current
774Opinion in Colloid & Interface Science 13, 134-140.
775Jia, H., Huang, W., Han, Y., Wang, Q., Wang, S., Dai, J., Tian, Z., Wang, D., Yan, H.,
776Lv, K., 2020. Systematic investigation on the interaction between SiO₂
777nanoparticles with different surface affinity and various surfactants. Journal of
778Molecular Liquids, 112777.
779Jones, S., Van Der Bent, V., Farajzadeh, R., Rossen, W., Vincent-Bonnieu, S.,
7802016. Surfactant screening for foam EOR: Correlation between bulk and core-
781flood experiments. Colloids and Surfaces A: Physicochemical and Engineering
782Aspects 500, 166-176.
783Kam, S., Rossen, W., 2003. A model for foam generation in homogeneous media.
784Spe Journal 8, 417-425.
785Kantzas, A., Chatzis, I., Dullien, F., 1988. Enhanced oil recovery by inert gas
786injection, SPE enhanced oil recovery symposium. Society of Petroleum Engineers.
787Karakashev, S.I., Ozdemir, O., Hampton, M.A., Nguyen, A.V., 2011. Formation and
788stability of foams stabilized by fine particles with similar size, contact angle and
789different shapes. Colloids and Surfaces A: Physicochemical and Engineering
790Aspects 382, 132-138.
791Koczko, K., Lobo, L., Wasan, D., 1992. Effect of oil on foam stability: aqueous
792foams stabilized by emulsions. Journal of colloid and interface science 150, 492-
793506.
794Kumar, S., Mandal, A., 2017. Investigation on stabilization of CO₂ foam by ionic
795and nonionic surfactants in presence of different additives for application in
796enhanced oil recovery. Applied Surface Science 420, 9-20.
797Lemlich, R., 1978. Prediction of changes in bubble size distribution due to
798interbubble gas diffusion in foam. Industrial & Engineering Chemistry
799Fundamentals 17, 89-93.
800Lin, C.E., Wang, T.Z., Chiu, T.C., Hsueh, C.C., 1999. Determination of the critical
801micelle concentration of cationic surfactants by capillary electrophoresis. Journal
802of High Resolution Chromatography 22, 265-270.
803Lobo, L., Nikolov, A., Wasan, D., 1989. Foam stability in the presence of oil: on
804the importance of the second virial coefficient. JOURNAL OF DISPERSION SCIENCE
805AND TECHNOLOGY 10, 143-161.
806Ma, K., Lontas, R., Conn, C.A., Hirasaki, G.J., Biswal, S.L., 2012. Visualization of
807improved sweep with foam in heterogeneous porous media using microfluidics.
808Soft Matter 8, 10669-10675.
809Maestro, A., Rio, E., Drenckhan, W., Langevin, D., Salonen, A., 2014. Foams
810stabilised by mixtures of nanoparticles and oppositely charged surfactants:
811relationship between bubble shrinkage and foam coarsening. Soft Matter 10,
8126975-6983.
813Martinez, A.C., Rio, E., Delon, G., Saint-Jalmes, A., Langevin, D., Binks, B.P., 2008.
814On the origin of the remarkable stability of aqueous foams stabilised by
815nanoparticles: link with microscopic surface properties. Soft Matter 4, 1531-1535.

816Nasr, N.H., Mahmood, S.M., Hematpur, H., 2019. A rigorous approach to analyze
817bulk and coreflood foam screening tests. *Journal of Petroleum Exploration and*
818*Production Technology* 9, 809-822.

819Nikolov, A., Wasan, D., Huang, D., Edwards, D., 1986. The effect of oil on foam
820stability: mechanisms and implications for oil displacement by foam in porous
821media, SPE annual technical conference and exhibition. Society of Petroleum
822Engineers.

823Osei-Bonsu, K., Grassia, P., Shokri, N., 2018. Effects of pore geometry on flowing
824foam dynamics in 3D-printed porous media. *Transport in Porous Media* 124, 903-
825917.

826Osei-Bonsu, K., Shokri, N., Grassia, P., 2015. Foam stability in the presence and
827absence of hydrocarbons: From bubble-to bulk-scale. *Colloids and Surfaces A:*
828*Physicochemical and Engineering Aspects* 481, 514-526.

829Osei-Bonsu, K., Shokri, N., Grassia, P., 2016. Fundamental investigation of foam
830flow in a liquid-filled Hele-Shaw cell. *Journal of colloid and interface science* 462,
831288-296.

832Panthi, K., Singh, R., Mohanty, K.K., 2017. Microencapsulation and stimuli-
833responsive controlled release of particles using water-in-air powders. *Langmuir*
83433, 3998-4010.

835Rossen, W.R., 1990. Theory of mobilization pressure gradient of flowing foams in
836porous media: I. Incompressible foam. *Journal of colloid and interface science*
837136, 1-16.

838Rossen, W.R., 1996. Foams in enhanced oil recovery. *Foams: theory,*
839*measurements and applications* 57, 413-464.

840Saint-Jalmes, A., 2006. Physical chemistry in foam drainage and coarsening. *Soft*
841*Matter* 2, 836-849.

842Shojaei, M.J., de Castro, A.R., Méheust, Y., Shokri, N., 2019. Dynamics of foam
843flow in a rock fracture: Effects of aperture variation on apparent shear viscosity
844and bubble morphology. *Journal of colloid and interface science* 552, 464-475.

845Shojaei, M.J., Osei-Bonsu, K., Grassia, P., Shokri, N., 2018a. Foam flow
846investigation in 3D-printed porous media: fingering and gravitational effects.
847*Industrial & Engineering Chemistry Research* 57, 7275-7281.

848Shojaei, M.J., Osei-Bonsu, K., Richman, S., Grassia, P., Shokri, N., 2018b. Foam
849stability influenced by displaced fluids and by pore size of porous media.
850*Industrial & Engineering Chemistry Research* 58, 1068-1074.

851Singh, R., Mohanty, K.K., 2015. Synergy between nanoparticles and surfactants
852in stabilizing foams for oil recovery. *Energy & Fuels* 29, 467-479.

853Singh, R., Mohanty, K.K., 2017. Nanoparticle-stabilized foams for high-
854temperature, high-salinity oil reservoirs, SPE Annual Technical Conference and
855Exhibition. Society of Petroleum Engineers.

856Talebian, S.H., Masoudi, R., Tan, I.M., Zitha, P.L., 2013. Foam assisted CO₂-EOR;
857concepts, challenges and applications, SPE Enhanced Oil Recovery Conference.
858Society of Petroleum Engineers.

859Vatanparast, H., Shahabi, F., Bahramian, A., Javadi, A., Miller, R., 2018. The role
860of electrostatic repulsion on increasing surface activity of anionic surfactants in
861the presence of hydrophilic silica nanoparticles. *Scientific reports* 8, 1-11.

862Verwey, E.J.W., Overbeek, J.T.G., 1955. Theory of the stability of lyophobic
863colloids. *Journal of Colloid Science* 10, 224-225.

864Yekeen, N., Manan, M.A., Idris, A.K., Padmanabhan, E., Junin, R., Samin, A.M.,
865Gbadamosi, A.O., Oguamah, I., 2018. A comprehensive review of experimental
866studies of nanoparticles-stabilized foam for enhanced oil recovery. *Journal of*
867*Petroleum Science and Engineering* 164, 43-74.

868Yu, D., Huang, F., Xu, H., 2012a. Determination of critical concentrations by
869synchronous fluorescence spectrometry. *Analytical Methods* 4, 47-49.

870Yu, J., An, C., Mo, D., Liu, N., Lee, R.L., 2012b. Foam mobility control for
871nanoparticle-stabilized supercritical CO₂ foam, SPE improved oil recovery
872symposium. Society of Petroleum Engineers.
873Yusuf, S., Manan, M., Zaidi Jaafar, M., 2013. Aqueous foams stabilized by
874hydrophilic silica nanoparticles via in-situ physisorption of nonionic TX100
875Surfactant. Iranian Journal of Energy and Environment 4, 0-0.
876Zhang, T., Roberts, M., Bryant, S.L., Huh, C., 2009. Foams and emulsions
877stabilized with nanoparticles for potential conformance control applications, SPE
878international symposium on oilfield chemistry. Society of Petroleum Engineers.

879