A piece-wise affine contracting map with positive entropy

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Abstract

We construct the simplest chaotic system with a two-point attractor.

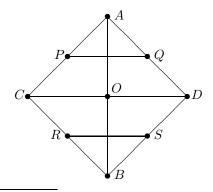
If $f: X \to X$ is an isometry of the metric space (X, d), then the topological entropy vanishes: $h_{top}(f) = 0$ (for definitions and notations consult e.g. [4]).

This follows from the fact, that the iterated distance $d_n^f = \max_{0 \le i < n} (f^i)^*(d)$ equals d. If f is distance non-increasing, the same equality holds and again $h_{\text{top}}(f) = 0$.

Whenever f can have discontinuities of some tame nature, so that f is piece-wise continuous, even the isometry result becomes difficult. In dimension 2 for invertible maps it was proven by Gutkin and Haydn [3]. In arbitrary dimension Buzzi proved that piece-wise affine isometries have zero topological entropy [2].

In the same paper after the theorem (remark 4) it is claimed that the result holds for arbitrary piece-wise (non-strictly) contracting maps. This latter is however wrong and the goal of this note is to present a counter-example.

Example: Let X be a rhombus ADBC with vertices $(\pm 1, 0), (0, \pm 1)$, see the figure below. Let O be its center and P, Q, R, S be on the sides as is shown.



¹Keywords: Piecewise affine maps, topological entropy.

Let f be partially defined on the interior of four big triangles forming the rhombus. These triangles are bijectively mapped by f as follows:

 $ACO \longrightarrow APQ, ADO \longrightarrow BRS, BCO \longrightarrow AQP, BDO \longrightarrow BSR.$

Thus the piece-wise affine map is defined.

If P, Q, R, S are middle-points of the intervals AC, AD and BC, BD, then the map is not strictly contracting. But if they are closer to the vertices A and B respectively than to C, D, then f is strictly contracting. In any case, the attractor of the system is the 2-point set $\{A, B\}$. Notice that the points belong to the singularity set, where the map f is not (uniquely) defined.

Taking $\varepsilon = \frac{1}{2}$ we observe that the cardinality of minimal (n, ε) -spanning set satisfies: $2^{n+2} \leq N(f, n, \varepsilon) \leq 2^{n+3}$. In fact, if we partition CD into 2^n equal intervals $Z_i Z_{i+1}$, then every $d_n^f \varepsilon$ -ball is contained in some triangle $AZ_i Z_{i+1}$ or $BZ_i Z_{i+1}$ and every such a triangle is covered by two $d_n^f \varepsilon$ -balls.

Therefore the topological entropy $h_{top}(f) = \log 2$ is positive. In addition, the Lyapunov spectrum is strictly negative at each point (for strict contractions), no invariant measure exists and so the variational principle breaks.

The result of Buzzi [2] generalizes however in the following fashion:

Theorem. Let f be a piece-wise affine map with restriction to each continuity component being conformal (non-strict) contraction. Then $h_{top}(f) = 0$.

Now we can repeat Buzzi's remark 4 [2]: The proof of his theorem 3 applies almost literally to the above case of piece-wise affine conformal contracting maps. Therefore we omit the proof.

Remark. It is obvious that if the attractor consists of one point only, then $h_{top}(f) = 0$. If the phase space $X \subset \mathbb{R}^1$ is one-dimensional and the map is (non-strictly) contracting, then again $h_{top}(f) = 0$. We don't even need to require piece-wise affine property. This follows from the Buzzi proposition 4 [1], yielding $h_{top}(f) \leq h_{mult}(f)$, where $h_{mult}(f)$ is the multiplicity entropy, because the latter always vanishes in dimension one.

Thus our example with 2 points attractor and 2-dimensional phase-space X is the simplest possible example with positive topological entropy.

References

- [1] J. Buzzi, Intrinsic ergodicity of affine maps in $[0,1]^d$, Mh. Math. **124** (1997), 97–118.
- J. Buzzi, Piecewise isometries have zero topological entropy, Ergod. Th. & Dynam. Sys. 21 (2001), 1371–1377.
- [3] E. Gutkin, N. Haydn, Topological entropy of polygon excange transformations and polygonal billiards, Ergod. Th. & Dynam. Sys. **17** (1997), 849–867.
- [4] A. Katok, B. Hasselblatt, Introduction to the Modern Theory of Dynamical Systems, Cambridge University Press (1995).