| Title | SY STEMS OF NONLINEAR VARIATIONAL <br> INEQUALITIES ARISING FROM PHA SE TRA NSITION <br> PHENOMENA |
| :---: | :--- |
| Author（s） | KENMOCHI，N．；NIEZGODKA，M． |
| Citation | 数理解析研究所講究録（1992），812：18－28 |
| Issue Date | 1992－10 |
| URL | http：／hdl．handle．net／2433／83043 |
| Right | Departmental Bulletin Paper |
| Type | publisher |
| Textversion |  |

# SYSTEMS OF NONLINEAR VARIATIONAL INEQUALITIES ARISING FROM PHASE TRANSITION PHENOMENA 

N．KENMOCHI and M．NIEZGODKA

## 1．Introduction

We consider an evolution system，consisting of a nonlinear second－order parabolic PDE and a nonlinear fourth－order parabolic PDE with constraint，which is described as follows：

$$
\begin{gather*}
\rho(u)_{t}+\lambda(w)_{t}-\Delta u=h(t, x) \quad \text { in } Q:=(0, T) \times \Omega  \tag{1.1-1}\\
\frac{\partial u}{\partial n}+n_{o} u=h_{o}(t, x) \quad \text { on } \Sigma:=(0, T) \times \Gamma  \tag{1.1-2}\\
u(0, \cdot)=u_{o} \quad \text { in } \Omega  \tag{1.1-3}\\
w_{t}-\Delta\left(-\nu \Delta w+\xi+g(w)-\lambda_{o}(w) u\right)=0 \quad \text { in } Q  \tag{1.2-1}\\
\frac{\partial w}{\partial n}=0, \quad \frac{\partial}{\partial n}\left(-\nu \Delta w+\xi+g(w)-\lambda_{o}(w) u\right)=0 \quad \text { on } \Sigma  \tag{1.2-2}\\
\xi \in \beta(w) \quad \text { on } Q  \tag{1.2-3}\\
w(0, \cdot)=w_{o} \quad \text { in } \Omega \tag{1.2-4}
\end{gather*}
$$

Here $\Omega$ is a bounded domain in $\mathbf{R}^{N}(1 \leq N \leq 3)$ with smooth boundary $\Gamma=\partial \Omega ; \rho: \mathbf{R} \rightarrow$ $\mathbf{R}, g: \mathbf{R} \rightarrow \mathbf{R}$ and $\lambda: \mathbf{R} \rightarrow \mathbf{R}$ are given functions and $\lambda_{o}(r)=\lambda^{\prime}(r)$（＝the derivative of $\lambda$ ） for $r \in \mathbf{R} ; \nu>0$ and $n_{o} \geq 0$ are given constants，and $h$ and $h_{o}$ are given functions on $Q$ and $\Sigma$ ，respectively；$u_{\circ}$ and $w_{o}$ are initial data；$\beta$ is a given maximal monotone graph in $\mathbf{R} \times \mathbf{R}$ ．

The system（1．1）－（1．2）is interpreted as a simplified model for thermodynamical phase separation in which $w$ represents the order parameter，$\theta=-\frac{1}{u}$ the（Kelvin）temperature and the free energy functional $F(\theta, w)$ is supposed to be dependent upon the temperature $\theta$ and to be given by the formula

$$
\begin{gather*}
F(\theta, w):=\int_{\Omega} f(\theta, w, \nabla w) d x, \quad w \in H^{1}(\Omega),  \tag{1.3}\\
f(\theta, w, \nabla w)=\left\{\frac{1}{2}\left(\nu_{o}+\nu_{1} \theta\right)|\nabla w|^{2}+\tau(\theta)+\theta(\hat{\beta}(w)+\hat{g}(w))+\lambda(w)\right\},
\end{gather*}
$$

where $\hat{\beta}$ is a proper l．s．c．convex function such that $\partial \hat{\beta}=\beta$ in $\mathbf{R} \times \mathbf{R}, \hat{g}$ is a primitive of $g$ on $\mathbf{R}, \lambda$ is the same as above，$\nu_{o} \geq 0, \nu_{1}>0$ are constants and $\tau: \mathbf{R} \rightarrow \mathbf{R}$ is a smooth function．

In some general settings，various models for thermodynamical phase separation phenom－ ena have been proposed and studied for instance by Luckhaus－Visintin［11］and Alt－Pawlow $[1,2]$ ．However，in their models the constraint（1．2－3）is not taken account of．

To illustrate our system (1.1)-(1.2), for instance, consider a binary system of alloys with components $A$ and $B$ ocuppying $\Omega$; let $w:=w_{A}$ and $w_{B}$ be the local concentrations of $A$ and $B$, respectively, such that

$$
w_{A}+w_{B}=\text { const.; }
$$

suppose that the free energy functional $F(\theta, w)$ of the Ginzburg-Landau type is of the form (1.3). Then, according to the thermodynamics approach of DeGroot-Mazur [5] and AltPawlow [1,2], we can derive from (1.3) with transformation $u:=-1 / \theta$, the mass and energy balance equations:

$$
\begin{gather*}
\rho(u)_{t}+\lambda(w)_{t}+\left[\frac{1}{2} \nu_{o}|\nabla w|^{2}\right]_{t}+\nabla \cdot \mathbf{q}=h(t, x) \quad \text { in } Q  \tag{1.4}\\
w_{t}+\nabla \cdot \mathbf{j}=0 \quad \text { in } Q \tag{1.5}
\end{gather*}
$$

where $\rho(u)=\tau(\theta)-\theta \tau^{\prime}(\theta), \mathbf{q}$ is the energy flux due to heat and mass transfer, $\mathbf{j}$ is the mass flux of the component $A$ and $h$ is a given heat source. Now suppose further that the fluxes $q$ and $j$ are described by the following constitutive relations:

$$
\begin{gather*}
\mathbf{q}=\nabla\left(\frac{1}{\theta}\right)(=-\nabla u) \quad \text { in } Q  \tag{1.6}\\
\mathbf{j}=-\nabla\left(\frac{\mu}{\theta}\right)(=\nabla(u \mu)) \quad \text { in } Q \tag{1.7}
\end{gather*}
$$

where

$$
\begin{equation*}
\frac{\mu}{\theta}=\frac{\delta}{\delta w}\left[\int_{\Omega} \frac{f(\theta, w, \nabla w)}{\theta} d x\right] \tag{1.8}
\end{equation*}
$$

and $\frac{\delta}{\delta w}$ denotes the functional derivative with respect to $w$. Since $f(\theta, w, \nabla w)$ includes the non-smooth term $\hat{\beta}(w)$, the right hand side of (1.8) is here understood in the multivalued sense

$$
\begin{gather*}
\frac{\delta}{\delta w}\left[\int_{\Omega} \frac{f(\theta, w, \nabla w)}{\theta} d x\right] \\
=\left\{-\nabla \cdot\left(\frac{\nu_{o}}{\theta}+\nu_{1}\right) \nabla w+\xi+g(w)+\frac{\lambda^{\prime}(w)}{\theta} ; \xi \in L^{2}(\Omega), \xi \in \beta(w) \text { a.e. on } \Omega\right\}, \tag{1.9}
\end{gather*}
$$

Now, combine (1.4)-(1.5) with (1.6)-(1.9). Then we obtain

$$
\begin{equation*}
\rho(u)_{t}+\lambda(w)_{t}+\left[\frac{\nu_{o}}{2}|\nabla w|^{2}\right]_{t}-\Delta u=h \quad \text { in } Q \tag{1.10}
\end{equation*}
$$

and

$$
\begin{gather*}
w_{t}-\Delta\left(-\nabla \cdot\left(\nu_{1}-\nu_{o} u\right) \nabla w+\xi+g(w)-\lambda^{\prime}(w) u\right)=0 \quad \text { in } Q  \tag{1.11}\\
\xi \in \beta(w) \quad \text { in } Q . \tag{1.12}
\end{gather*}
$$

Therefore, if $\nu_{o}=0$ and $\nu_{1}=\nu$, or if in (1.10) the term $\left[\frac{\nu_{o}}{2}|\nabla w|^{2}\right]_{t}$ is experimentally allowed to be neglected and in (1.11) the coefficient $\left(\nu_{1}-\nu_{o} u\right)$ of $\nabla w$ replaced by a positive constant $\nu$, then system (1.1-1)-(1.2-i), $i=1,3$, is regarded as a simplified form of (1.10)-(1.12). System (1.1)-(1.2) consists of these equations and initial-boundary conditions (1.1-i), $i=2,3$, and (1.2-i), $i=2,4$.

The aim of this paper is to study a weak formulation for system (1.1)-(1.2) in the variational sense, taking advantage of subdifferential techniques in Hilbert spaces.

## 2. Main results

Throughout this note, for a general (real) Banach space $X$ we denote by $|\cdot|_{X}$ the norm in $X$ and by $X^{\star}$ the dual space of $X$.

For simplicity we use the notations:

$$
\begin{aligned}
(v, w) & :=\int_{\Omega} v w d x \\
(v, w)_{\Gamma} & :=\int_{\Gamma} v w d \Gamma(x) \\
a(v, w) & :=\int_{\Omega} \nabla v \cdot w, \in L^{2}(\Omega) \\
\nabla w d x & \text { for } v, w \in L^{2}(\Gamma)
\end{aligned}
$$

Moreover we put

$$
\begin{gathered}
H:=L^{2}(\Omega), \quad V:=H^{1}(\Omega), \\
H_{o}:=\left\{z \in H ; \int_{\Omega} z d x=0\right\}, \quad V_{o}:=V \cap H_{o},
\end{gathered}
$$

and denote by $\pi$ the projection from $H$ onto $H_{o}$, i.e.

$$
\pi(z)(x):=z(x)-\frac{1}{|\Omega|} \int_{\Omega} z(y) d y, \quad z \in H
$$

Also, $H_{o}$ is a Hilbert space with $|z|_{H_{o}}=|z|_{H}$ as well as $V_{o}$ with $|z|_{V_{o}}=|\nabla z|_{H}$; we use sometimes symbol $(\cdot, \cdot)_{o}$ for the inner product in $H_{o}$ and $\langle\cdot, \cdot\rangle_{o}$ for the duality pairing between $V_{o}^{\star}$ and $V_{o}$.

As usual, identifying $H$ with its dual, we have

$$
V \subset H \subset V^{\star}
$$

with dense and compact embeddings. Similarly, identifying $H_{o}$ with its dual, we have

$$
V_{o} \subset H_{o} \subset V_{o}^{\star}
$$

with dense and compact embeddings. Also, we denote by $J_{o}$ the duality mapping from $V_{o}$ onto $V_{o}^{\star}$ which is defined by the formula

$$
\left\langle J_{o} z, \eta\right\rangle_{o}=a(z, \eta) \quad \text { for all } z, \eta \in V_{o}
$$

Therefore, in particular, if $z^{\star}:=J_{o} z \in H_{o}$, then $z \in H^{2}(\Omega)$ and $z$ is the unique solution of the Neumann problem

$$
\begin{equation*}
-\Delta z=z^{\star} \quad \text { in } \Omega, \quad \frac{\partial z}{\partial n}=0 \quad \text { on } \Gamma, \quad \int_{\Omega} z d x=0 \tag{2.1}
\end{equation*}
$$

Accordingly, if $\eta \in H^{2}(\Omega)$ and $\frac{\partial \eta}{\partial n}=0$ a.e. on $\Gamma$, then $J_{o}[\pi(\eta)]=-\Delta \eta$.
Now, we denote by $(\mathrm{P})$ the system (1.1)-(1.2) mentioned in section 1 and discuss it under the following assumptions (A1)-(A6):
(A1) $\rho: \mathbf{R} \rightarrow \mathbf{R}$ is an increasing Lipschitz continuous function with Lipschitz continuous inverse $\rho^{-1}: \mathbf{R} \rightarrow \mathbf{R}$; we denote by $C_{\rho}$ a common Lipschitz constant of $\rho$ and $\rho^{-1}$.
(A2) $\lambda, \lambda_{o}: \mathbf{R} \rightarrow \mathbf{R}$ are Lipschitz continuous functions and $\lambda_{o}=\lambda^{\prime}$; we denote by $C_{\lambda}$ a common Lipschitz constant of $\lambda$ and $\lambda_{0}$.
(A3) $g: \mathbf{R} \rightarrow \mathbf{R}$ is a Lipschitz continuous function; we denote by $C_{g}$ the Lipschitz constant of $g$.
(A4) $\nu$ is a positive constant and $n_{o}$ is a non-negative constant.
(A5) $\beta$ is a maximal monotone graph in $\mathbf{R} \times \mathbf{R}$ with bounded and non-empty interior int. $D(\beta)$ of the domain $D(\beta)$ in $\mathbf{R}$; we put int. $D(\beta)=\left(\sigma_{\star}, \sigma^{\star}\right)$ for $-\infty<\sigma_{\star}<\sigma^{\star}<\infty$ and hence $\overline{D(\beta)}=\left[\sigma_{\star}, \sigma^{\star}\right]$, and we may assume that $\beta$ is the subdifferential of a nonnegative, proper, l.s.c. and convex function $\hat{\beta}$ on $\mathbf{R}$, since the range $R(\beta)$ of $\beta$ is the whole $\mathbf{R}$.
(A6) $0<T<\infty, h \in L^{2}(0, T ; H), h_{o} \in W^{1,2}\left(0, T ; L^{2}(\Gamma)\right)$ and $u_{o} \in H, w_{o} \in V$ with $\hat{\beta}\left(w_{o}\right) \in L^{1}(\Omega)$.

We introduce

$$
K(\hat{\beta}):=\left\{z \in H ; \hat{\beta}(z) \in L^{1}(\Omega)\right\}
$$

and

$$
K_{m}(\hat{\beta}):=\left\{z \in K(\hat{\beta}) ; \frac{1}{|\Omega|} \int_{\Omega} z d x=m\right\} \quad \text { for each } m \in \mathbf{R} .
$$

We next give the weak formulation for $(\mathrm{P})$.
Definition 2.1. A couple $\{u, w\}$ of functions $u:[0, T] \rightarrow V$ and $w:[0, T] \rightarrow H^{2}(\Omega)$ is called a (weak) solution of (P), if the following conditions (w1)-(w4) are satisfied:
(w1) $u \in L^{2}(0, T ; V) \cap L^{\infty}(0, T ; H), \rho(u) \in C_{w}([0, T] ; H), C_{w}([0, T] ; H)$ being the space of all weakly continuous functions from $[0, T]$ into $H, \rho(u)^{\prime}\left(=\frac{d}{d t} \rho(u)\right) \in L^{1}\left(0, T ; V^{\star}\right), w \in$ $L^{2}\left(0, T ; H^{2}(\Omega)\right) \cap L^{\infty}(0, T ; V), w^{\prime} \in L^{2}\left(0, T ; V^{\star}\right)$ and $\lambda(w)^{\prime} \in L^{1}\left(0, T ; V^{\star}\right) ;$
(w2) $\rho(u)(0)=\rho\left(u_{o}\right)$ and $w(0)=w_{o}$;
(w3) for a.e. $t \in[0, T]$ and all $z \in V$,

$$
\begin{equation*}
\frac{d}{d t}(\rho(u(t))+\lambda(w(t)), z)+a(u(t), z)+\left(n_{o} u(t)-h_{o}(t), z\right)_{\Gamma}=(h(t), z) \tag{2.2}
\end{equation*}
$$

(w4) for a.e. $t \in[0, T]$,

$$
\begin{equation*}
\frac{\partial}{\partial n} w(t)=0 \quad \text { a.e. on } \Gamma \tag{2.3}
\end{equation*}
$$

and there is a function $\xi \in L^{2}(0, T ; H)$ such that

$$
\begin{equation*}
\xi \in \beta(w) \quad \text { a.e. in } Q \tag{2.4}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{d}{d t}(w(t), \eta)+\nu(\Delta w(t), \Delta \eta)-\left(g(w(t))+\xi(t)-\lambda^{\prime}(w(t)) u(t), \Delta \eta\right)=0 \tag{2.5}
\end{equation*}
$$

for all $\eta \in H^{2}(\Omega)$ with $\frac{\partial \eta}{\partial n}$ a.e. on $\Gamma$, and a.e. $t \in[0, T]$.
When it is necessary to indicate the data $h, h_{o}, u_{o}, w_{o}$, we denote problem ( P ) by ( $\mathrm{P} ; h, h_{o}$, $u_{0}, w_{o}$ ).

Remark 2.1. Let $\{u, w\}$ be any solution of (P). Then it follows from (2.5) in (w4) that

$$
\frac{d}{d t}(w(t), 1)=0 \quad \text { for a.e. } t \in[0, T]
$$

whence

$$
\int_{\Omega} w(t, x) d x=\int_{\Omega} w_{o} d x \quad \text { for all } t \in[0, T] .
$$

Therefore, putting

$$
\begin{equation*}
m:=\frac{1}{|\Omega|} \int_{\Omega} w_{o} d x \tag{2.6}
\end{equation*}
$$

we observe that $w(t)-m \in V_{o}$ for all $t \in[0, T]$.
Our main results of this paper are stated as follows:
Theorem 2.1. Assume that $1 \leq N \leq 3$ and (A1)-(A6) hold, and assume with notation (2.6) that

$$
m \in \operatorname{int} . D(\beta), \text { i.e. } \sigma_{\star}<m<\sigma^{\star}
$$

Then $(P)$ has one and only one solution $\{u, w\}$. Moreover, the solution $\{u, w\}$ has the following bounds:

$$
\begin{align*}
& |u|_{L^{\infty}(0, T ; H)}+|u|_{L^{2}(0, T ; V)}+|w|_{L^{\infty}(0, T ; V)}+|\hat{\beta}(w)|_{L^{\infty}\left(0, T ; L^{1}(\Omega)\right)}+\left|w^{\prime}\right|_{L^{2}\left(0, T ; V^{*}\right)} \\
& \leq R_{o}\left(\left|u_{o}\right|_{H},\left|w_{o}\right|_{V},\left|\hat{\beta}\left(w_{o}\right)\right|_{L^{1}(\Omega)},|h|_{L^{2}(0, T ; H)},\left|h_{o}\right|_{L^{2}\left(0, T ; L^{2}(\Gamma)\right)}\right) \tag{2.7}
\end{align*}
$$

where $R_{o}: \mathbf{R}_{+}^{5} \rightarrow \mathbf{R}$ is a function which is bounded on each bounded subset of $\mathbf{R}_{+}^{5}$;

$$
\begin{gather*}
|w|_{L^{2}\left(0, T ; H^{2}(\Omega)\right)}+\left|\rho(u)^{\prime}\right|_{L^{1}\left(0, T ; V^{\star}\right)}+\left|\lambda(w)^{\prime}\right|_{L^{1}\left(0, T ; V^{\star}\right)} \\
\leq R_{1}\left(\frac{1}{\delta}, r(\delta),\left|u_{o}\right|_{H},\left|w_{o}\right|_{V},\left|\hat{\beta}\left(w_{o}\right)\right|_{L^{1}(\Omega)},|h|_{L^{2}(0, T ; H)},\left|h_{o}\right|_{L^{2}\left(0, T ; L^{2}(\Gamma)\right)}\right) \tag{2.8}
\end{gather*}
$$

where $R_{1}: \mathbf{R}_{+}^{7} \rightarrow \mathbf{R}_{+}$is a function which is bounded on each bounded subset of $\mathbf{R}_{+}^{7}, \delta$ is an arbitrary number satisfying

$$
\begin{equation*}
0<\delta<1, \quad \sigma_{\star}<m-\delta<m+\delta<\sigma^{\star} \tag{2.9}
\end{equation*}
$$

and

$$
\begin{equation*}
r(\delta)=\sup \left\{\left|r^{\prime}\right| ; r^{\prime} \in \beta(m-\delta) \cap \beta(m+\delta)\right\} \tag{2.10}
\end{equation*}
$$

Remark 2.2. In estimates (2.7) and (2.8), the dependence of the solution $\{u, w\}$ upon functions $\rho, \lambda, g$ and $\beta$ is not explicitly indicated. However, as will be able to be easily checked, the functions $R_{o}$ and $R_{1}$ are chosen so as to be independent of them, as long as the Lipschitz constants $C_{\rho}, C_{\lambda}, C_{g}$ and the length $\sigma^{\star}-\sigma_{\star}$ of $D(\beta)$ vary in a bounded subset of $\mathbf{R}_{+}$.

Theorem 2.2. Assume that $1 \leq N \leq 3$ and (A1)-(A5) hold. Let $\left\{h_{n}\right\},\left\{h_{o n}\right\},\left\{u_{o n}\right\}$ and $\left\{w_{o n}\right\}$ be bounded sequences in $L^{2}(0, T ; H), W^{1,2}\left(0, T ; L^{2}(\Gamma)\right), H$ and $V$, respectively, and assume that $\left\{\hat{\beta}\left(w_{o n}\right)\right\}$ is bounded in $L^{1}(\Omega)$. Further suppose that as $n \rightarrow \infty$

$$
h_{n} \rightarrow h \quad \text { in } L^{2}(0, T ; H), \quad h_{o n} \rightarrow h_{o} \quad \text { in } L^{2}\left(0, T ; L^{2}(\Gamma)\right)
$$

and

$$
u_{o n} \rightarrow u_{o} \quad \text { in } H, \quad w_{o n} \rightarrow w_{o} \quad \text { in } V .
$$

Then we have the following statements (i) and (ii):
(i) Suppose that

$$
\begin{equation*}
\sigma_{\star}<m_{n}:=\frac{1}{|\Omega|} \int_{\Omega} w_{o n} d x<\sigma^{\star} \quad \text { for all } n \tag{2.11}
\end{equation*}
$$

and

$$
\sigma_{\star}<m:=\frac{1}{|\Omega|} \int_{\Omega} w_{o} d x<\sigma^{\star} .
$$

Let $\left\{u_{n}, w_{n}\right\}$ be the solution of $\left(P_{n}\right):=\left(P ; h_{n}, h_{o n}, u_{o n}, w_{o n}\right)$ for each $n$ and $\{u, w\}$ be the solution of $(P):=\left(P ; h, h_{o}, u_{o}, w_{o}\right)$. Then, as $n \rightarrow \infty$,

$$
\begin{gathered}
u_{n} \rightarrow u \text { in } L^{2}(0, T ; H), \\
\rho\left(u_{n}\right) \rightarrow \rho(u) \quad \text { weakly in } H \text { and uniformly in } t \in[0, T], \\
w_{n} \rightarrow w \text { in } L^{2}(0, T ; V) \text { and weakly* in } L^{\infty}(0, T ; V)
\end{gathered}
$$

and

$$
w_{n}^{\prime} \rightarrow w^{\prime} \quad \text { weakly in } L^{2}\left(0, T ; V^{\star}\right)
$$

(ii) Suppose that (2.11) hlods and

$$
m=\sigma_{\star} \text { or } \sigma^{\star},
$$

Then, for the solution $\left\{u_{n}, w_{n}\right\}$ of $\left(P_{n}\right)$, we have as $n \rightarrow \infty$,

$$
\begin{array}{ll} 
& w_{n} \rightarrow m \quad \text { in } C([0, T] ; H) \\
u_{n} \rightarrow u \quad & \text { in } L^{2}(0, T ; H) \text { and weakly in } L^{2}(0, T ; V)
\end{array}
$$

and

$$
\rho\left(u_{n}\right) \rightarrow \rho(u) \quad \text { weakly in } H \text { and uniformly in } t \in[0, T],
$$

where $u \in C([0, T] ; H) \cap W_{\text {loc }}^{1,2}((0, T] ; H) \cap L_{\text {loc }}^{\infty}((0, T] ; V) \cap L^{2}(0, T ; V)$ is the unique solution of

$$
\frac{d}{d t}(\rho(u(t)), z)+a(u(t), z)+\left(n_{o} u(t)-h_{o}(t), z\right)_{\Gamma}=(h(t), z)
$$

$$
\begin{align*}
& \quad \text { for all } z \in V \text {, a.e. } t \in[0, T],  \tag{2.12}\\
& u(0)=u_{o} .
\end{align*}
$$

Remark 2.3. In (ii) of Theorem 2.2, moreover if $h \in L^{\infty}(Q), h_{o} \in L^{\infty}(\Sigma), u_{o} \in L^{\infty}(\Omega)$ and $m \in D(\beta)$, then the pair $\{u, w\}$, with the solution $u$ of $(2.12)$ and $w=m$, is the solution of $(\mathrm{P})$ in the sense of Definition 2.1. In fact, under such restrictions on the data we see that $u \in L^{\infty}(Q)$ and hence $\xi:=k-g(m)+\lambda^{\prime}(m) u \in \beta(m)$ on $Q$ for a certain constant $k$. Thus condition (w4) of Definition 2.1 is satisfied.

## 3. Sketch of proofs

(1) (Uniqueness) The uniqueness of the solution of ( P ) can be proved by using Gronwall's inequality with the help of the following embedding inequalities:

$$
|z|_{L^{q}(\Omega)} \leq C_{o}|\nabla z|_{H}, \quad|z|_{L^{q}(\Omega)} \leq \delta|\nabla z|_{H}+C_{\delta}|z|_{V_{o}^{*}}
$$

for all $z \in V_{o}$ and $1 \leq q<6$, where $C_{o}$ is a positive constant, and $\delta$ is an arbitrary positive constant with a constant $C_{\delta}$ dependent only on $\delta$.
(2) (Existence) For the construction of a solution of $(\mathrm{P})$ we consider the approximate problem $(\mathrm{P})_{\mu}\left(=\left(\mathrm{P}_{\mu} ; h, h_{o}, u_{o}, w_{o}\right)\right)$, with parameter $0<\mu \leq 1$, to find a pair of functions $u_{\mu}:[0, T] \rightarrow$ $V$ and $w_{\mu}:[0, T] \rightarrow H^{2}(\Omega)$ fulfilling the following conditions $(\mathrm{w} 1)_{\mu}-(\mathrm{w} 4)_{\mu}$ :
$(\mathrm{w} 1)_{\mu} \quad u_{\mu} \in W^{1,2}(0, T ; H) \cap L^{\infty}(0, T ; V), w_{\mu} \in W^{1,2}(0, T ; H) \cap L^{\infty}(0, T ; V) \cap L^{2}\left(0, T ; H^{2}(\Omega)\right) ;$
$(\mathbf{w} 2)_{\mu} u_{\mu}(0)=u_{o}$ and $w_{\mu}(0)=w_{o}$;
(w3) ${ }_{\mu}$ for a.e. $t \in[0, T]$ and all $z \in V$,

$$
\begin{equation*}
\left(\rho\left(u_{\mu}\right)^{\prime}(t)+\lambda\left(w_{\mu}\right)^{\prime}(t), z\right)+a\left(u_{\mu}(t), z\right)+\left(n_{o} u_{\mu}(t)-h_{o}(t), z\right)_{\Gamma}=(h(t), z) ; \tag{3.1}
\end{equation*}
$$

$(\mathrm{w} 4)_{\mu}$ for a.e. $t \in[0, T]$,

$$
\begin{equation*}
\frac{\partial w_{\mu}(t)}{\partial n}=0 \quad \text { a.e. on } \Gamma, \tag{3.2}
\end{equation*}
$$

and there is a function $\xi_{\mu} \in L^{2}(0, T ; H)$ such that

$$
\begin{equation*}
\xi_{\mu} \in \beta\left(w_{\mu}\right) \quad \text { a.e. on } Q \tag{3.3}
\end{equation*}
$$

and

$$
\begin{align*}
& \left(w_{\mu}^{\prime}(t), \eta\right)-\mu\left(w_{\mu}^{\prime}(t), \Delta \eta\right)+\nu\left(\Delta w_{\mu}(t), \Delta \eta\right) \\
& \quad-\left(g\left(w_{\mu}(t)\right)-\lambda^{\prime}\left(w_{\mu}(t)\right) u_{\mu}(t)+\xi_{\mu}(t), \Delta \eta\right)=0 \tag{3.4}
\end{align*}
$$

for all $\eta \in H^{2}(\Omega)$ with $\frac{\partial \eta}{\partial n}=0$ a.e. on $\Gamma$ and a.e. $t \in[0, T]$.

Besides we reformulate $(\mathrm{P})_{\mu}$ as a system of evolution equations including subdifferential operators. For this purpose, let us introduce convex functions $\varphi$ on $H_{o}$ and $\psi^{t}, t \leq t \leq T$, on $H$ as follows:

$$
\varphi(z):= \begin{cases}\frac{\nu}{2}|\nabla z|_{H}^{2}+\int_{\Omega} \hat{\beta}(z+m) d x & \text { if } z \in V_{o} \text { and } \hat{\beta}(z+m) \in L^{1}(\Omega)  \tag{3.5}\\ \infty & \text { otherwise }\end{cases}
$$

where

$$
m:=\frac{1}{|\Omega|} \int_{\Omega} w_{o} d x
$$

and

$$
\psi^{t}(z):= \begin{cases}\frac{1}{2}|\nabla z|_{H}^{2}+\frac{n_{o}}{2}|z|_{L^{2}(\Gamma)}^{2}-\left(h_{o}(t), z\right)_{\Gamma} & \text { if } z \in V  \tag{3.6}\\ \infty & \text { otherwise }\end{cases}
$$

We then consider the subdifferential $\partial \varphi$ of $\varphi$ in $H_{o}$ and the subdifferential $\partial \psi^{t}$ of $\psi$ in $H$. It is easy to see that
(i) $z^{\star} \in \partial \varphi(z)$ if and only if $z^{\star} \in H_{o}, z \in V_{o} \cap\left(K_{m}(\hat{\beta})-m\right)$ and

$$
\begin{gathered}
\left(z^{\star}, v-z\right)_{o} \leq \nu a(z, v-z)+\int_{\Omega} \hat{\beta}(v+m) d x-\int_{\Omega} \hat{\beta}(z+m) d x \\
\text { for all } v \in V_{o} \cap\left(K_{m}(\hat{\beta})-m\right)
\end{gathered}
$$

(ii) $\partial \psi^{t}$ is singlevalued, and $z^{\star}=\partial \psi^{t}(z)$ if and only if $z^{\star} \in H, z \in V$ and

$$
\left(z^{\star}, v\right)=a(z, v)+\left(n_{o} z-h_{o}(t), v\right)_{\Gamma} \quad \text { for all } v \in V
$$

For each $\mu \in(0,1]$, problem $(\mathrm{P})_{\mu}$ has at most one solution and we have:
Lemma 3.1. Let $\sigma_{\star}<m<\sigma^{\star}$, and $\lambda_{1}(r):=\lambda(r+m)$ and $g_{1}(r):=g(r+m)$ for $r \in \mathbf{R}$. Then a pair $\left\{u_{\nu}, w_{\mu}\right\}$ of functions is a solution of $(P)_{\mu}$ if and only if the pair $\left\{u_{\mu}, v_{\mu}\right\}$ with $v_{\mu}:=w_{\mu}-m$ is a solution of the problem $(P)_{\mu}^{\nu}$ defined below:
$(\mathrm{P})_{\mu}^{\prime}$ Find a pair $\left\{u_{\mu}, v_{\mu}\right\}$ of functions satisfying the following conditions $(w 1)_{\mu}^{\prime}-\left(w_{4}\right)_{\mu}^{\prime}$ :
$(\mathrm{w} 1)_{\mu}^{\prime} \quad u_{\mu} \in W^{1,2}(0, T ; H) \cap L^{\infty}(0, T ; V)$ and $v_{\mu} \in W^{1,2}\left(0, T ; H_{o}\right) \cap L^{\infty}\left(0, T ; V_{o}\right)$;
$(\mathbf{w} 2)_{\mu}^{\prime} u_{\mu}(0)=u_{o}$ and $v_{\mu}(0)=v_{o}:=w_{o}-m ;$
$(\mathrm{w} 3)_{\mu}^{\prime}$ for a.e. $t \in[0, T]$,

$$
\begin{equation*}
\rho\left(u_{\mu}\right)^{\prime}(t)+\lambda_{1}\left(v_{\mu}\right)^{\prime}(t)+\partial \psi^{t}\left(u_{\mu}(t)\right)=h(t) \tag{3.7}
\end{equation*}
$$

$(\mathrm{w} 4)_{\mu}^{\prime}$ for a.e. $t \in[0, T]$,

$$
\begin{equation*}
\left(J_{o}^{\star}+\mu I\right) v_{\mu}^{\prime}(t)+\partial \varphi\left(v_{\mu}(t)\right)+\pi\left[g_{1}\left(v_{\mu}(t)\right)-\lambda_{1}^{\prime}\left(v_{\mu}(t)\right) u_{\mu}(t)\right] \ni 0 \tag{3.8}
\end{equation*}
$$

We can prove Lemma 3.1 by using the following lemma which is concerned with the Lagrange multipliers of elliptic variational inequalities.

Lemma 3.2. Let $\sigma_{\star}<m<\sigma^{\star}$ and $\ell$ be any element of $H$. Consider the following two problems $\left(M_{m}\right)$ and $\left(M_{m}\right)^{\prime}$ :
$\left(\mathbf{M}_{m}\right)$ Find a function $z_{m} \in K_{m}(\hat{\beta}) \cap V$ such that

$$
\nu a\left(z_{m}, z_{m}-\eta\right)+\int_{\Omega} \hat{\beta}\left(z_{m}\right) d x \leq\left(\ell, z_{m}-\eta\right)+\int_{\Omega} \hat{\beta}(\eta) d x \quad \text { for all } \eta \in K_{m}(\hat{\beta}) \cap V
$$

$\left(\mathbf{M}_{m}\right)^{\prime}$ Find a function $z_{m} \in K_{m}(\hat{\beta}) \cap H^{2}(\Omega), \gamma_{m} \in \mathbf{R}$ and $\xi_{m} \in H$ such that

$$
-\nu \Delta z_{m}+\xi_{m}=\ell+\gamma_{m} \quad \text { in } \Omega
$$

and

$$
\xi_{m} \in \beta\left(z_{m}\right) \quad \text { a.e. on } \Omega, \quad \frac{\partial z_{m}}{\partial n}=0 \quad \text { a.e. on } \Gamma .
$$

Then $\left(M_{m}\right)^{\nu}$ has a solution $\left\{z_{m}, \xi_{m}, \gamma_{m}\right\}$ and the function $z_{m}$ is the unique solution of $\left(M_{m}\right)$. Moreover, $\gamma_{m}$ can be chosen so that

$$
\begin{equation*}
\left|\gamma_{m}\right| \leq 4 M^{5}\left(1+|\ell|_{H}\right), \tag{3.9}
\end{equation*}
$$

where $M=\max \left\{\frac{1}{\delta}, r(\delta), \sigma^{\star}-\sigma_{\star},|\Omega|, \frac{1}{|\Omega|}\right\}$ for $\delta$ and $r(\delta)$ satisfying (2.9) and (2.10); $z_{m}$ satisfies that

$$
\begin{equation*}
\left(-\Delta z_{m}, \xi_{m}\right) \geq 0 \tag{3.10}
\end{equation*}
$$

and

$$
\begin{equation*}
\nu\left|\Delta z_{m}\right|_{H} \leq|\ell|_{H}+\left|\gamma_{m}\right||\Omega|^{\frac{1}{2}} \tag{3.11}
\end{equation*}
$$

For the detail proof of Lemma 3.2 we refer to [9; Proposition 5.1]. Thanks to the additional term $\mu v_{\mu}^{\prime}$ problem $\left(P_{\mu}\right)^{\prime}$, hence $\left(P_{\mu}\right)$, is uniquely solved in the Hilbert spaces $H$ and $H_{o}$ by applying time-dependent subdifferential techniques evolved in [4, 10]. In fact, we have the following result.

Proposition 3.1. In addition to all the conditions of Theorem 2.1, assume that $u_{o} \in V$. Then, for each $\mu \in(0,1]$, problem $(P)_{\mu}$ has one and only one solution $\left\{u_{\mu}, w_{\mu}\right\}$. Moreover, the solution $\left\{u_{\mu}, w_{\mu}\right\}$ satisfies the bounds of the following type:

$$
\begin{gathered}
\left|u_{\mu}\right|_{C([0, T] ; H)}+\left|\nabla u_{\mu}\right|_{L^{2}(0, T ; H)}+\left|w_{\mu}^{\prime}\right|_{L^{2}\left(0, T ; V^{\star}\right)} \\
+\mu\left|w_{\mu}^{\prime}\right|_{L^{2}(0, T ; H)}^{2}+\left|w_{\mu}\right|_{L^{\infty}(0, T ; V)}+\left|\hat{\beta}\left(w_{\mu}\right)\right|_{L^{\infty}\left(0, T ; L^{1}(\Omega)\right)} \\
\leq \tilde{R}_{o}\left(\left|u_{o}\right|_{H},\left|w_{o}\right| V,\left|\hat{\beta}\left(w_{o}\right)\right|_{L^{1}(\Omega)},|h|_{L^{2}(0, T ; H)},\left|h_{o}\right|_{L^{2}\left(0, T ; L^{2}(\Gamma)\right)}\right),
\end{gathered}
$$

where $\tilde{R}_{o}: \mathbf{R}_{+}^{5} \rightarrow \mathbf{R}_{+}$is a function which is independent of $\mu$ and bounded on each bounded subset of $\mathbf{R}_{+}^{5}$;

$$
\left|w_{\mu}\right|_{L^{2}\left(0, T ; H^{2}(\Omega)\right)}+\left|\rho\left(u_{\mu}\right)^{\prime}\right|_{L^{1}\left(0, T ; V^{*}\right)}+\left|\lambda\left(w_{\mu}\right)^{\prime}\right|_{L^{1}\left(0, T ; V^{*}\right)}
$$

$$
\leq \tilde{R}_{1}\left(\frac{1}{\delta}, r(\delta),\left|u_{o}\right|_{H},\left|w_{o}\right|_{V},\left|\hat{\beta}\left(w_{o}\right)\right|_{L^{1}(\Omega)},|h|_{L^{2}(0, T ; H)},\left|h_{o}\right|_{L^{2}\left(0, T ; L^{2}(\Gamma)\right)}\right)
$$

where $\tilde{R}_{1}: \mathbf{R}_{+}^{7} \rightarrow \mathbf{R}_{+}$is a function which is independent of $\mu$ and bounded on each bounded subset of $\mathbf{R}_{+}^{7}, \delta$ is an arbitrary number satisfying (2.9) and $r(\delta)$ is a constant given by (2.10).

By the above proposition we obtain a solution $\{u, w\}$ of ( P ), passing to the limit in $\mu \rightarrow 0$, and see that the solution satisfies estimates (2.7) and (2.8).
(3) (Proof of Theorem 2.2) The assertions of Theorem 2.2 follow easily from estimates for the solution of $(\mathrm{P})$ in Theorem 2.1.

Remark 3.1. In this paper the domain $D(\beta)$ of $\beta$ is supposed to be bounded in R. However this is not essential for the assertions of Theorems 2.1 and 2.2. For instance, our results can be extended to the case when $\operatorname{int} . D(\beta) \neq 0$ and there are constants $k_{\beta}>0$ and $k_{\beta}^{\prime}>0$ such that

$$
|\beta(r)| \geq k_{\beta}|r|-k_{\beta}^{\prime} \quad \text { for all } r \in D(\beta) ;
$$

note that under this condition we may assume that

$$
\hat{\beta}(r) \geq \hat{k}_{\beta}|r|^{2} \quad \text { for all } r \in D(\hat{\beta})
$$

where $\hat{k}_{\beta}>0$ is a certain constant.
Application. As a typical example of maximal monotone graphs $\beta$ in $\mathbf{R} \times \mathbf{R}$ arising in the context of phase separation (cf. [3]), we consider an increasing smooth function $\beta^{c}:(0,1) \rightarrow$ $\mathbf{R}$ defined by

$$
\beta^{c}(w):=c \log \frac{w}{1-w}
$$

with positive real parameter $c$. Also, as an example of non-smooth $\beta$, we consider the subdifferential $\beta^{0}$ of the indicator function of the interval $[0,1]$ in $\mathbf{R}$, which is the limit of $\beta^{c}$ as $c \rightarrow 0$ in the sense of maximal monotone graphs in $\mathbf{R} \times \mathbf{R}$.

By virtue of Theorem 2.1, problem (P) with $\beta=\beta^{c}(c \geq 0)$ has one and only one solutioon $\left\{u^{c}, w^{c}\right\}$, provided that $u_{o} \in H, w_{o} \in V$ with $0<m<1$ and $\log \frac{w_{o}}{1-w_{o}} \in L^{1}(\Omega)$, $h \in L^{2}(0, T ; H)$ and $h_{o} \in W^{1,2}\left(0, T ; L^{2}(\Gamma)\right)$. Moreover, it easily follows from the estimates (2.7),(2.8) and the uniqueness of solutions to ( P ) that as $c \rightarrow 0$, the solution $\left\{u^{c}, w^{c}\right\}$ converges to the solution $\left\{u^{0}, w^{0}\right\}$ in the similar sense as in (i) of Theorem 2.2.

## References

[1] H. W. Alt and I. Pawlow, Dynamics of non-isothermal phase separation, in Free Boundary Value Problems, K.-H. Hoffmann and J. Sprekels ed., ISNM 95, Birkhäuser, Basel, 1990, pp. 1-26.
[2] H. W. Alt and I. Pawlow, Existence of solutions for non-isothermal phase separation, preprint.
[3] F. E. Blowey and C. M. Elliott, The Cahn-Hilliard gradient theory for phase separation with non-smooth free energy, Part I: Mathematical analysis, European J. Appl. Math. 2(1991), 233-280.
[4] P. Colli and A. Visintin, On a class of doubly nonlinear evolution equations, Comm. Partial Differential Equations 15(1990), 737-756.
[5] S. R. DeGroot and P. Mazur, Non-Equilibrium Thermodynamics, Dover Publ., New York, 1984.
[6] E. DiBenedetto and R. E. Showalter, A pseudo-parabolic variational inequality and Stefan problem, Nonlinear Anal.TMA 6(1982),279-291.
[7] C. E. Elliott, The Cahn-Hilliard model for the kinetics of phase separation, in Mathematical Models for Phase Change Problems, J. F. Rodrigues ed., ISNM 88, Birkhäuser, Basel, 1989, pp. 35-73.
[8] C. M. Elliott and S. Zheng, On the Cahn-Hilliard equation, Arch. Rat. Mech. Anal. 96(1986), 339-357.
[9] N. Kenmochi, M. Niezgódka and I. Pawlow, Subdifferential operator appraoch to the Cahn-Hilliard equation with constraint, preprint.
[10] N. Kenmochi and I. Pawlow, A class of nonlinear elliptic-parabolic equaions with timedependent constraints, Nonlinear Anal. TMA 10(1986), 1181-1202.
[11] S. Luckhaus and A. Visintin, Phase transition in multicomponent systems, Manuscripta Math. 43(1983), 261-288.
[12] Y. Oono and S. Puri, Study of the phase separation dynamics by use of cell dynamical systems, I. Modelling., Phys. Rev. A 38(1988), 434-453
[13] R. Temam, Infinite Dimensional Dynamical Systems in Mechanics and Physics,Springer Verlag, Berlin, 1988.
[14] W. von Wahl, On the Cahn-Hilliard equation $u^{\prime}+\Delta^{2} u-\Delta f(u)=0$, Delft Progress Report 10(1985), 291-310.
[15] S. Zheng, Asymptotic behaviour of the solution to the Cahn-Hilliard equation, Applicable Anal. 23(1986), 165-184.
N. Kenmochi: Department of Mathematics, Faculty of Education, Chiba University 1-33 Yayoi-cho, Chiba, 260 Japan
M. Niezgódka: Institute of Applied Mathematics and Mechanics, Warsaw University Banacha 2, 00-913 Warsaw, Poland

