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First measurement of electron neutrino scattering cross section on argon

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We report the first electron neutrino cross section measurements on argon, based on data collected by the ArgoNeuT experiment running in the GeV-scale NuMI beamline at Fermilab. A flux-averaged $\nu_e + \bar{\nu}_e$ total and a lepton angle differential cross section are extracted using 13 ν_e and $\bar{\nu}_e$ events identified with fully automated selection and reconstruction. We employ electromagnetic-induced shower characterization and analysis tools developed to identify $\nu_e/\bar{\nu}_e$ -like events among complex interaction topologies present in ArgoNeuT data ($\langle E_{\bar{\nu}_e} \rangle = 4.3$ GeV and $\langle E_{\nu_e} \rangle = 10.5$ GeV). The techniques are widely applicable to searches for electron-flavor appearance at short and long baseline using liquid argon time projection chamber technology. Notably, the data-driven studies of GeV-scale $\nu_e/\bar{\nu}_e$ interactions presented here probe an energy regime relevant for future DUNE oscillation physics.

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While neutrino mass and mixing has enjoyed a bounty of rich discoveries over the past few decades, a number of questions remain. Most notably, the ordering of the neutrino mass states, the value of the CP -violating phase (δ_{CP}), and the possibility of new degrees of freedom driving oscillations (e.g., $\nu_{e,\mu,\tau} \rightarrow \nu_s$) remain open questions. Electron neutrino identification and characterization is essential to the $\nu_\mu \rightarrow \nu_e$ and $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ appearance-based short- and long-baseline experiments seeking answers to these questions [1–5]. Many current and upcoming neutrino oscillation experiments will benefit from the

high-resolution detection capabilities of liquid argon time projection chamber (LArTPC) technology. In particular, the SBN Program at Fermilab [1] studies the possibility of a sterile flavor participating in oscillations, and DUNE [2] seeks to determine the neutrino mass hierarchy and extract δ_{CP} , both using LArTPCs.

Exploring these physics topics with LArTPC experiments requires careful reconstruction of ν_e and $\bar{\nu}_e$ interactions, often difficult to identify with automated methods. Even with LArTPC technology and its ability to provide mm-scale-resolution pictures of the events in question, hit and cluster finding, shower formation, and finally, neutrino energy reconstruction and flavor identification, remain challenging. Algorithms for effectively interpreting the abundance of information provided in LArTPC data are critical for extracting physics results. This is particularly true for DUNE, which will rely on the inclusive set of all $\nu_e/\bar{\nu}_e$ charged current (CC) interactions in the few-GeV energy range [2] rather than an exclusive

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CC quasielastic-like signal channel. The selection must accommodate substantial contributions from the varying event topologies associated with quasielastic, resonant, and deep inelastic scattering, and significant effects from nuclear physics, including multinucleon correlations and final state interactions [6]. Background events in a $\nu_e/\bar{\nu}_e$ search in DUNE also present challenges; even for underground detectors with low or negligible cosmic contamination, the electromagnetic (EM) showers characteristic of $\nu_e/\bar{\nu}_e$ events are readily mimicked by numerous neutrino-induced background processes, especially $\nu_\mu/\bar{\nu}_\mu$ CC and neutral current (NC) interactions featuring $\pi^0 \rightarrow \gamma\gamma$ and non-negligible $\Delta \rightarrow N\gamma$ contributions. These energy reconstruction and background issues directly affect oscillation measurements. For example, while DUNE is expected to be statistics limited early on with exposures less than 100 kt · MW · year, energy-scale, flux, and interaction model systematic uncertainties will quickly take the lead in the δ_{CP} measurement uncertainty budget [2].

On the way to efficient electron-flavor reconstruction with minimal background in LArTPC neutrino experiments, data-based studies at the GeV scale are largely absent. These are the first $\nu_e/\bar{\nu}_e$ measurements extracted from GeV-scale neutrino beam data using automated methods.

Previously, ArgoNeuT demonstrated that topological information alone could be used to identify electron neutrino candidates by rejecting gamma backgrounds based on the characteristic gap expected between the neutrino interaction vertex and the beginning of a gamma-induced shower due to the large (relative to LArTPC spatial resolution) conversion length of 18 cm in liquid argon [7]. It was further shown, using samples of events selected by visual scanning methods containing either an electron or gamma candidate, that vertex dE/dx could be used to separate electrons from gammas, a notable milestone in LArTPC reconstruction for exploiting the wealth of charge and spatial detail provided by the technology. However, these strategies are quickly complicated by interactions with high multiplicity where hadronic overlap with EM showers can obscure the essential gap and dE/dx information close to the vertex.

Toward the total $\nu_e + \bar{\nu}_e$ CC cross section reported in the later part of this report, we first provide a short description of the ArgoNeuT detector and a detailed explanation of the EM shower reconstruction, background and systematics estimation, and signal extraction procedures employed, providing an analysis framework for future LArTPC-based $\nu_e/\bar{\nu}_e$ appearance searches. While previous $\nu_e/\bar{\nu}_e$ CC studies in LArTPCs have relied only on topological and calorimetric information specific to the neutrino interaction vertex and hand scans [7,8], this study broadens the scope of classification tools to take advantage of the entire EM shower topology, a necessary step toward developing inclusive $\nu_e/\bar{\nu}_e$ CC selection strategies for GeV-scale

neutrino interactions in the presence of significant background.

The ArgoNeuT LArTPC experiment at Fermilab collected data in the NuMI beamline just upstream of the MINOS near detector [9] in 2009–2010, with the vast majority (1.25×10^{20} protons on target) taken in “low-energy antineutrino mode” ($\langle E_{\bar{\nu}_e} \rangle = 4.3$ GeV with 68% falling between 1.0 and 6.5 GeV and $\langle E_{\nu_e} \rangle = 10.5$ GeV with 68% falling between 2.5 and 21.5 GeV) [10]. ArgoNeuT featured a $40 \times 47 \times 90$ cm³ [vertical, drift, horizontal (beam)] TPC at 481 V/cm with 240 induction and 240 collection wires separated by 4 mm and sampled at a rate of 5 MHz by the readout electronics. The detector is described in detail in Ref. [11].

Neutrino interactions in ArgoNeuT are simulated using the GENIE [12] neutrino event generator in combination with GEANT4-based [13] detector and particle propagation models. Neutrino and antineutrino fluxes from the NuMI beam are provided by the MINERνA Collaboration [10]. After event simulation, interactions in the ArgoNeuT detector are first reconstructed using the LArSoft software package [14]. The algorithms, described in detail in Ref. [15], proceed with the following steps: (i) noise removal and deconvolution of raw wire signals to correct for electronics and field response, (ii) hit finding, (iii) clustering of hits on each plane based on proximity to one another, (iv) reconstruction of three-dimensional (3D) tracks by matching clusters across wire planes with temporal consistency, and (v) calorimetric reconstruction.

Custom reconstruction tools then use the output of the standard software package to build candidate EM showers. The shower reconstruction algorithm relies on two objects produced by the standard tools: (i) 3D tracks with associated vertex and direction information and (ii) clusters of hits on each plane tagged as “showerlike” or “tracklike” based on the measure of multiple coulomb scattering along the clustered hits, the size of the cluster, and its proximity to other clusters.

The shower reconstruction algorithm used in this analysis is designed to reconstruct electrons in ArgoNeuT. While the subsequently described selection procedure could be applied to any population of electron neutrino candidate events, we focus on reconstructing only the leading shower in each neutrino interaction. Particularly, we make use of the reliable 3D track reconstruction for defining the vertex and direction of a shower. The shower reconstruction builds candidate electron showers around reconstructed 3D tracks by looking for showerlike clusters of hits in close proximity to the track axis; the hit and proximity thresholds for finding candidate showers were optimized to maximize reconstruction completeness, purity, and efficiency for electrons specifically.

After reconstruction, a set of quality cuts is applied to the data. First, a filter is applied to reject events with a muon reconstructed in the downstream MINOS near detector that

when projected backward crosses the ArgoNeuT detector active volume, indicative of a background $\nu_\mu/\bar{\nu}_\mu$ CC interaction. The vertex of each candidate signal electron must lie within the fiducial volume defined to be 3 cm inside the anode and cathode planes, 4 cm from the top and bottom boundaries of the TPC, 6 cm from the upstream face of the detector, and 20 cm from the downstream face of the detector. The vertex must be at least 20 cm from the back of the detector to give candidate electrons enough space to begin exhibiting showerlike qualities, motivated by the radiation length ($X_0 = 14$ cm) in argon. Additionally, reconstructed showers must have $\cos(\theta_z) > 0.05$ (where θ_i are the opening angles with respect to the detector axes), resulting in a negligible loss of phase space in favor of removing backward-going failures of reconstruction more prominent in data due to electronics noise. To remove events with through-going muons from neutrino interactions upstream of the TPC, we reject any event with a reconstructed 3D cluster that falls within 2.5 cm of the upstream face. Finally, we require that the closest hit in a candidate shower on each plane be within 2 cm of the reconstructed vertex to remove track reconstruction failures.

The calorimetric discrimination techniques described here could be applied to any reconstructed shower object, independent of the reconstruction algorithm. We simply assume a reconstructed shower possesses (i) a collection of hits on at least one plane, (ii) a vertex, and (iii) a direction. In the small ArgoNeuT detector, the rarity of complete shower containment prohibits the use of total charge (e.g., for complete shower characterization and energy reconstruction). We instead use charge ratios constructed from topological regions of charge, shown in Fig. 1, to characterize the shape and evolution of each candidate electron shower [6]. For example, longitudinal development

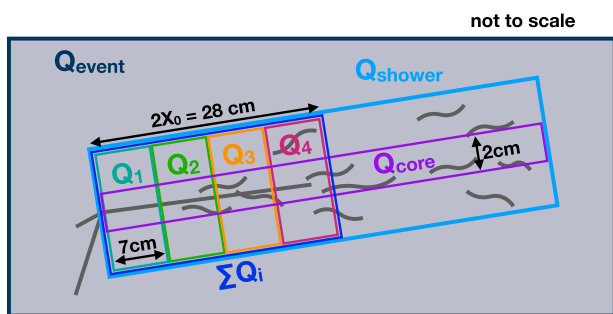


FIG. 1. The topological charge regions defined for each reconstructed shower, intended to characterize the transverse and longitudinal development of the candidate shower. The cells defined Q_1 through Q_4 extend in the transverse direction to include all hits identified as part of the candidate shower. Q_{core} extends to include all shower hits in the longitudinal direction. Q_{shower} extends in both directions to include all hits identified as part of the candidate shower. The distances defined are in two dimensions.

of the shower is modeled by defining the ratios $Q_n/\sum_i Q_i$ where $n = 1, 2, 3, 4$, and transverse shower development is characterized using $Q_{\text{core}}/Q_{\text{shower}}$. These topologically motivated charge ratios are powerful discriminators, along with vertex dE/dx , for selecting signal $\nu_e/\bar{\nu}_e$ CC events among backgrounds involving a variety of event classes. These include difficult-to-reconstruct background deep-inelastic events often characterized by multiple overlapping tracks and EM activity. The following variables are defined for $\nu_e/\bar{\nu}_e$ CC classification using a boosted decision tree (BDT): three angles [$\cos(\theta_x)$, $\cos(\theta_y)$, and $\cos(\theta_z)$], $Q_{\text{shower}}/Q_{\text{event}}$, $\sum_i Q_i/Q_{\text{shower}}$, $Q_{\text{core}}/Q_{\text{shower}}$, $Q_n/\sum_i Q_i$, and vertex dE/dx , calculated by taking the median charge in the first 4 cm of the track [7]. All charge variables are defined using the collection plane only. The output of the BDT trained using these quantities is shown in Fig. 2. The three most important inputs for separating signal and background, all with approximately equal impact, are $Q_{\text{shower}}/Q_{\text{event}}$, $Q_{\text{core}}/Q_{\text{shower}}$, and vertex dE/dx . The distance between the neutrino vertex and EM shower start is not used in this analysis for signal identification; the high neutrino energies and resulting large track multiplicities complicate automated gap reconstruction, yielding weak separation power between electrons and gammas.

The purpose of the BDT is twofold: it (i) distinguishes well-reconstructed signal interactions from well-reconstructed backgrounds and (ii) rejects reconstruction failures. Given the difficulty of reconstructing EM activity and deep inelastic scattering interactions in particular, the shower reconstruction was designed to maximize the efficiency of reconstructing the signal channel, with little concern for the number of reconstruction failures generated. This strategy is predicated on the fact that these failures can be subsequently rejected using the BDT, which

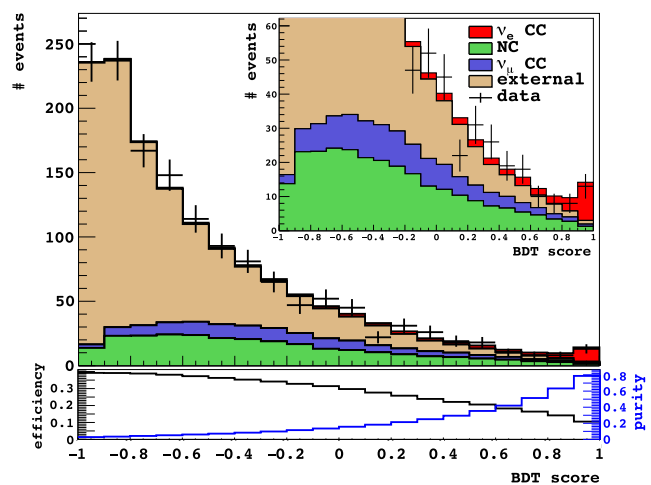


FIG. 2. The distribution of BDT scores for data and simulation. The signal selection in this analysis uses events with BDT score > 0.9 . The inset shows the same information zoomed in to better show this signal region.

considers the complete topology of the reconstructed object, rather than relying solely on vertex dE/dx .

One difficulty in this analysis is the small size of the ArgoNeuT detector. We find that a significant background comes from EM-like activity in the detector produced by interactions originating outside of the detector active volume. This is a complication unique to ArgoNeuT, where it is impossible to move sufficiently far from the edge of the detector to reject a significant fraction of these outside backgrounds while simultaneously maintaining satisfactory signal statistics. Additionally, we find that the external background is underestimated in the ArgoNeuT simulation, which only generates neutrino interactions that occur with and inside the cryostat. While the simulation reproduces the energy and topological characteristics of external EM-like backgrounds in the detector, it misrepresents the total quantity of these backgrounds. To correct for this deficit and constrain the external background contribution in the strict $\nu_e/\bar{\nu}_e$ selection region, the external background is scaled as a function of BDT score based on a data-simulation comparison sideband with score <0 . To determine the scale factor, a line was fit to the ratio of data to simulation after subtracting out internal background contributions (ν_μ CC and ν NC) in the background-only sideband. The scale factor is ~ 3.5 at BDT score = -1 and ~ 1.0 at BDT score = 1 . The data-driven function is motivated by the fact that external backgrounds tend to look topologically distinct from signal, a characteristic which is quantitatively described by decreasing BDT score, a proxy for event topology. The scaling was validated with a hand scan of events in data at low- and mid-range BDT scores in the background-only sideband.

To reduce the impact of the uncertainty associated with the background scaling on the final selection, we have limited our signal definition to events with BDT score >0.9 , a bin with low external and total background that yields the most significant signal selection. A conservative 100% uncertainty on the quantity of external background is included in the systematic error, which encompasses the range of multiple scaling methods that were considered. Other systematic uncertainties considered include those associated with the neutrino interaction model, found by varying a set of relevant parameters in GENIE independently according to Ref. [12], in addition to uncertainties in the integrated flux, collected protons on target, and number of target argon nuclei. Given the low statistics of our measurement, statistical uncertainties dominate the results reported here.

The reconstructed vertex dE/dx distribution for data after the final selection is shown in Fig. 3. One of the candidate electron neutrino interactions, among the 13 selected, is shown in Fig. 4. All 13 selected interactions are presented in the Supplemental Material [16]. The topology of selected events is consistent with expectation and the strict cut on BDT score means the selection is more

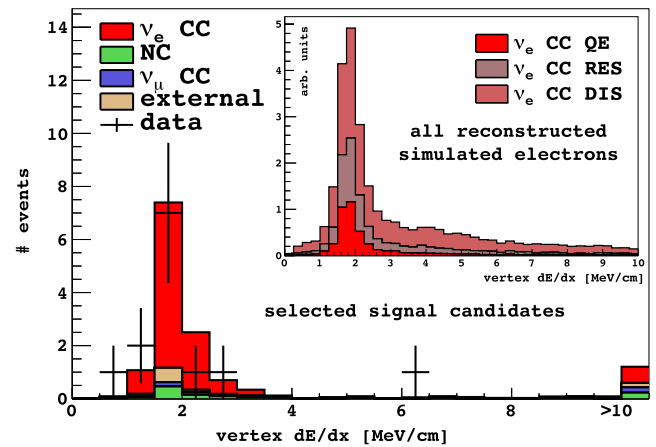


FIG. 3. Reconstructed vertex dE/dx for data and simulation after selection. The inset shows the vertex dE/dx distribution for electrons reconstructed from a sample of simulated $\nu_e/\bar{\nu}_e$ events broken up by interaction mechanism, demonstrating that the vertex dE/dx tail is mainly from deep inelastic scattering.

efficient for quasielastic events than deep inelastic scattering (DIS) events. We expect 3.6 quasielastic events, 3.4 resonant events, and 4.1 DIS events in the final sample, and note that many of the interactions selected have simple topologies. This is unsurprising given the stringent cut on

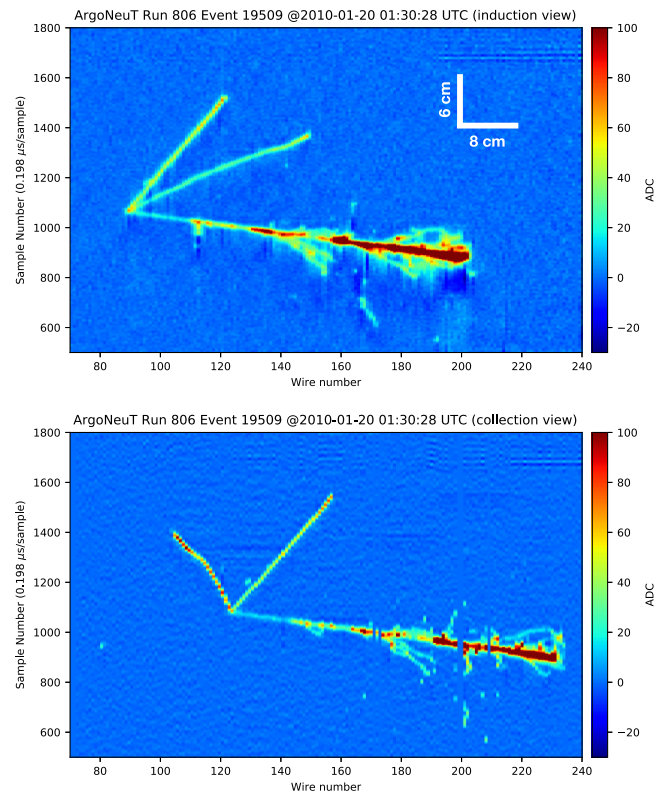


FIG. 4. A candidate electron neutrino interaction. The scale shown applies to both images. The color is proportional to the charge collected.

BDT score and the size of the data sample. While it is difficult to make definitive statements about the initial interaction of the selected events based on final state topology, a hand scan of DIS events in simulation was done to confirm the topologies that are selected could be consistent with DIS interactions. Additionally, there is a non-negligible probability that there is zero or one DIS event in the final selection. A subsequent scan was done of all events in data with BDT score >0.7 , which was also consistent with the expectation and yielded many more clear, high-multiplicity examples of DIS signal-like interactions in data.

The inset of Fig. 3 shows the distribution of vertex dE/dx for all reconstructed simulated electrons in the defined fiducial volume separated according to interaction mechanism (quasielastic, resonant, and DIS). Notably, the tail of the distribution is composed almost entirely of DIS interactions, an important contribution for the few-GeV neutrinos observed by ArgoNeuT.

ArgoNeuT cannot sign select for neutrino type, and significant contributions of both ν_e and $\bar{\nu}_e$ interactions are expected in NuMI low-energy antineutrino mode data. Thus, we define a flux-averaged total cross section such that it is a combination of ν_e and $\bar{\nu}_e$: $\sigma_{\nu_e+\bar{\nu}_e} = (N - B)/(\epsilon N_{\text{Ar}}(\Phi_{\nu_e} + \Phi_{\bar{\nu}_e}))$, where N is the number of events selected in data, B is the number of background events in simulation, ϵ is the selection efficiency, N_{Ar} is the number of argon targets, and Φ is integrated flux. The ν_e and $\bar{\nu}_e$ fluxes can be found in the Supplemental Material of Ref. [10]. Using this convention, we extract a total cross section of $(1.04 \pm 0.38(\text{stat})_{-0.23}^{+0.15}(\text{syst})) \times 10^{-36} \text{ cm}^2$ on argon ($\langle E_{\nu_e} \rangle = 10.5 \text{ GeV}$ and $\langle E_{\bar{\nu}_e} \rangle = 4.3 \text{ GeV}$) consistent with the GENIE expectation (11.2 $_{-1.4}^{+0.4}$ signal events and 3.0 $_{-1.3}^{+2.0}$ total background events). Further, we report a differential cross section in Fig. 5, again as a combination

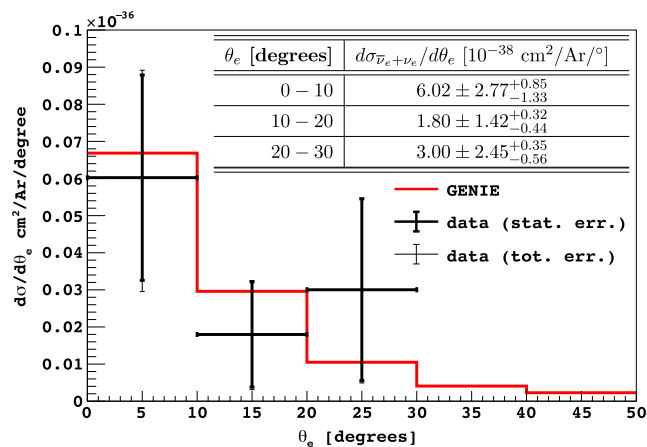


FIG. 5. The ArgoNeuT $\nu_e + \bar{\nu}_e$ CC differential cross section for electron/positron angle with respect to the neutrino beam compared to the GENIE prediction.

of ν_e and $\bar{\nu}_e$ defined in this way: $d\sigma(\theta_{e,i})/d\theta_e = (N_i - B_i)/(\epsilon_i \Delta\theta_{e,i} N_{\text{Ar}}(\Phi_{\nu_e} + \Phi_{\bar{\nu}_e}))$. Uncertainties associated with the GENIE modeling contribute most to the systematic uncertainties. The selection yields 10.5 $_{-0.5}^{+0.6}\%$ efficiency (for events originating inside the fiducial volume) with 78.9 $_{-11.8}^{+8.1}\%$ purity, where the spread in both numbers is due to systematics dominated by GENIE variations. Importantly, the interpretation of these results, for example, in comparisons to model predictions and event generators, *requires* the consideration of both the detailed ν_e and $\bar{\nu}_e$ fluxes simultaneously [10].

There are several notable factors that impact the performance achieved in this analysis. While the efficiency is sufficient for exploring the data-driven classification techniques and performing the measurements reported here, it is limited by ArgoNeuT's intrinsic reconstruction capabilities. First, ArgoNeuT's size is such that EM shower containment is a rarity, which leads to difficulty in event classification. Poor track containment, in general, also affects vertex reconstruction and event classification. Second, the signal selection in ArgoNeuT is necessarily very strict since we cannot move sufficiently far away from the detector walls in the fiducial volume definition to reduce background events produced by interactions external to the active volume of the detector, most notably single gammas. With improvements to these issues anticipated in future detectors like DUNE, we expect a significant increase in inclusive $\nu_e/\bar{\nu}_e$ CC signal selection efficiency.

We have reported a total $\nu_e + \bar{\nu}_e$ cross section and a differential cross section in terms of electron/positron angle with respect to the incoming neutrino using the fully automated reconstruction and analysis framework described above. These are the first measurements of electron neutrino scattering cross sections on argon. The results are statistics limited, further affected by the reconstruction efficiency in ArgoNeuT and the strict selection required to mitigate external backgrounds in the small detector. Furthermore, this is the first measurement of electron neutrino and antineutrino scattering using the same target nucleus and over the same energy range that will be used by the DUNE experiment.

The unique selection techniques presented here are particularly useful for identifying $\nu_e/\bar{\nu}_e$ CC interactions among typical GeV-scale neutrino backgrounds, including events involving gamma-induced showers and/or containing multiple tracks and complicated topologies. We find that obfuscation of key information for electromagnetic shower identification (vertex dE/dx and gap) by multiple hadronic tracks, known to be simulated poorly, is at once extremely important and underemphasized in the literature, and that taking full advantage of the unique topological characteristics of a shower candidate is essential. Further development of calorimetry-based techniques for signal classification is critical to inform and

direct machine learning-based image classification methods currently at the forefront of pattern recognition technology [17–20].

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