

Variational description of Gibbs-non-Gibbs dynamical transitions for the Curie-Weiss model

R. Fernández ¹
F. den Hollander ²
J. Martínez ³

February 22, 2012

Abstract

We perform a detailed study of Gibbs-non-Gibbs transitions for the Curie-Weiss model subject to independent spin-flip dynamics (“infinite-temperature” dynamics). We show that, in this setup, the program outlined in van Enter, Fernández, den Hollander and Redig [3] can be fully completed, namely that Gibbs-non-Gibbs transitions are *equivalent* to bifurcations in the set of global minima of the large-deviation rate function for the trajectories of the magnetization *conditioned* on their endpoint. As a consequence, we show that the time-evolved model is non-Gibbs if and only if this set is not a singleton for *some* value of the final magnetization. A detailed description of the possible scenarios of bifurcation is given, leading to a full characterization of passages from Gibbs to non-Gibbs—and vice versa—with sharp transition times (under the dynamics Gibbsianness can be lost and can be recovered).

Our analysis expands the work of Ermolaev and Külske [7] who considered zero magnetic field and finite-temperature spin-flip dynamics. We consider both zero and non-zero magnetic field but restricted to infinite-temperature spin-flip dynamics. Our results reveal an interesting dependence on the interaction parameters, including the presence of forbidden regions for the optimal trajectories and the possible occurrence of overshoots and undershoots in the optimal trajectories. The numerical plots provided are obtained with the help of MATHEMATICA.

MSC 2010. 60F10, 60K35, 82C22, 82C27.

Key words and phrases. Curie-Weiss model, spin-flip dynamics, Gibbs vs. non-Gibbs, dynamical transition, large deviations, action integral, bifurcation of rate function.

Acknowledgment. FdH is supported by ERC Advanced Grant VARIS-267356. JM is supported by Erasmus Mundus scholarship BAPE-2009-1669. The authors are grateful to A. van Enter, V. Ermolaev, C. Külske, A. Opoku and F. Redig for discussions.

¹Department of Mathematics, Utrecht University, P.O. Box 80010, 3508 TA Utrecht, The Netherlands, R.Fernandez1@uu.nl

²Mathematical Institute, Leiden University, P.O. Box 9512, 2300 RA, Leiden, The Netherlands, denholla@math.leidenuniv.nl

³Mathematical Institute, Leiden University, P.O. Box 9512, 2300 RA, Leiden, The Netherlands, martinez@math.leidenuniv.nl

1 Introduction and main results

Section 1.1 provides background and motivation, Section 1.2 a preview of the main results. Section 1.3 introduces the Curie-Weiss model and the key questions to be explored. Section 1.4 recalls a few facts from large-deviation theory for trajectories of the magnetization in the Curie-Weiss model subjected to infinite-temperature spin-flip dynamics and provides the link with the specification kernel of the time-evolved measure when it is Gibbs. Section 1.5 states the main results and illustrates these results with numerical pictures. The pictures are made with MATHEMATICA, based on analytical expressions appearing in the text. Proofs are given in Sections 2 and 3. Section 1 takes up half of the paper.

1.1 Background and motivation

Dynamical Gibbs-non-Gibbs transitions represent a relatively novel and surprising phenomenon. The setup is simple: an initial Gibbsian state (e.g. a collection of interacting Ising spins) is subjected to a stochastic dynamics (e.g. a Glauber spin-flip dynamics) at a temperature that is *different* from that of the initial state. For many combinations of initial and dynamical temperature, the time-evolved state is observed to become non-Gibbs after a finite time. Such a state cannot be described by any absolutely summable Hamiltonian and therefore *lacks a well-defined notion of temperature*.

The phenomenon was originally discovered by van Enter, Fernández, den Hollander and Redig [2] for *heating dynamics*, in which a low-temperature Ising model is subjected to an infinite-temperature dynamics (independent spin-flips) or a high-temperature dynamics (weakly-dependent spin-flips). The state remains Gibbs for short times, but becomes non-Gibbs after a finite time. Remarkably, heating in this case does not lead to a succession of states with increasing temperature, but to states where the notion of temperature is lost altogether. Furthermore, it turned out that there is a difference depending on whether the initial Ising model has zero or non-zero magnetic field. In the former case, non-Gibbsianness once lost is never recovered, while in the latter case Gibbsianness is recovered at a later time.

This initial work triggered a decade of developments that led to general results on Gibbsianness for small times (Le Ny and Redig [13], Dereudre and Roelly [1]), loss and recovery of Gibbsianness for discrete spins (van Enter, Külske, Opoku and Ruszel [10, 5, 6, 15, 4], Redig, Roelly and Ruszel [16]), and loss and recovery of Gibbsianness for continuous spins (Külske and Redig [12], Van Enter and Ruszel [5, 6]). A particularly fruitful research direction was initiated by Külske and Le Ny [9], who showed that Gibbs-non-Gibbs transitions can also be defined naturally for *mean-field* models, such as the Curie-Weiss model. Precise results are available for the latter, including sharpness of the transition times and an explicit characterization of the conditional magnetizations leading to non-Gibbsianness (Külske and Opoku [11], Ermolaev and Külske [7]). In particular, the work in [7] shows that in the mean-field setting Gibbs-non-Gibbs transitions occur for all initial temperatures below criticality, both for *cooling dynamics* and for *heating dynamics*.

The ubiquitousness of the Gibbs-non-Gibbs phenomenon calls for a better understanding of its causes and consequences. Unfortunately, the mathematical approach used in most references is opaque on the intuitive level. Generically, non-Gibbsianness is proved by looking at the evolving system at two times, the initial and the final time, and applying techniques from equilibrium statistical mechanics. This is an indirect approach that does not illuminate the relation between the Gibbs-non-Gibbs phenomenon and the dynamical effects responsible for

its occurrence. This unsatisfactory situation was addressed in Enter, Fernández, den Hollander and Redig [3], where possible dynamical mechanisms were proposed and a *program* was put forward to develop a theory of Gibbs-non-Gibbs transitions on *purely dynamical grounds*. The present paper shows that this program can be fully carried out for the Curie-Weiss model subject to an infinite-temperature dynamics.

In the mean-field scenario, the key object is the time-evolved single-spin average *conditional* on the final empirical magnetization. Non-Gibbsianness corresponds to a discontinuous dependence of this average on the final magnetization. The discontinuity points are called *bad magnetizations* (see Definition 1.1 below). Dynamically, such discontinuities are expected to arise whenever there is more than one possible trajectory *compatible* with the bad magnetization at the end. Indeed, this expectation is confirmed and exploited in the sequel. The actual conditional trajectories are those minimizing the large-deviation rate function on the space of trajectories of magnetizations. The time-evolved measure remains Gibbsian whenever there is a single minimizing trajectory for every final magnetization, in which case the specification kernel can be computed explicitly (see Proposition 1.4 below). In contrast, if there are multiple optimal trajectories, then the choice of trajectory can be decided by an infinitesimal perturbation of the final magnetization, and this is responsible for non-Gibbsianness.

1.2 Preview of the main results

In the present paper we study in detail the large-deviation rate function for the trajectory of the magnetization in the Curie-Weiss model with pair potential $J > 0$ and magnetic field $h \in \mathbb{R}$ (see (1.1) below). We exploit the fact that, due to the mean-field character of the interaction, this rate function can be expressed as a function of the initial and the final magnetization only (see Proposition 1.2 below), i.e., the trajectories are uniquely determined by the magnetizations at the beginning and at the end (see Corollary 1.3 and Proposition 1.5 below). Here is a summary of the main results (see Fig. 1):

1. If $0 < J \leq 1$ (supercritical temperature), then the evolved state is Gibbs at all times. On the other hand, if $J > 1$ (subcritical temperature) there exists some time Ψ_U at which multiple trajectories appear. The associated non-Gibbsianness persists for all later times when $h = 0$ (zero magnetic field). All these features were already shown by Ermolaev and Külske [7].
2. For $h \neq 0$ there is a time $\Psi_* > \Psi_U$ at which Gibbsianness is restored for all later times.
3. There is a change in behavior at $J = \frac{3}{2}$. For $1 < J \leq \frac{3}{2}$:
 - (a) If $h = 0$, then only the zero magnetization is bad for $t > \Psi_c$.
 - (b) If $h > 0$ ($h < 0$), then there is only one bad magnetization for $\Psi_U < t \leq \Psi_*$. This bad magnetization changes with t but is always strictly negative (strictly positive).

For $J > \frac{3}{2}$:

- (a) If $h = 0$, then there is a time $\Psi_c > \Psi_U$ such that for $\Psi_U < t < \Psi_c$ there are two non-zero bad magnetizations (equal in absolute value but with opposite signs), while for $t \geq \Psi_c$ only the zero magnetization is bad.
- (b) If $h \neq 0$ and small enough, then there are two times $\Psi_T > \Psi_L$ between Ψ_U and Ψ_* such that for $\Psi_U < t \leq \Psi_L$ and $\Psi_T \leq t \leq \Psi_*$ only one bad magnetization occurs, while for $\Psi_L < t < \Psi_T$ two bad magnetizations occur.

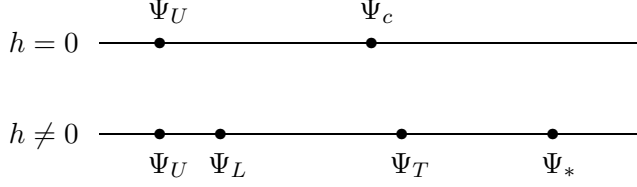


Figure 1: Crossover times for $h = 0$ and $h \neq 0$ when $J > \frac{3}{2}$.

All the crossover times depend on J, h and are strictly positive and finite. Our analysis gives a detailed picture of the optimal trajectories for different J, h and different conditional magnetizations. Among the novel features we mention:

- (1) Presence of *forbidden regions* that cannot be crossed by any optimal trajectory. The boundary of these regions is given by the multiple optimal trajectories when bifurcation sets in. The forbidden regions were predicted in [3] and first found, for $h = 0$, by Ermolaev and Külske [7].
- (2) Existence of *overshoots* and *undershoots* for optimal trajectories for $h \neq 0$.
- (3) Classification of the bad magnetizations leading to multiple optimal trajectories. These bad magnetizations depend on J, h and change with time.

1.3 The model

1.3.1 Hamiltonian and dynamics

The Curie-Weiss model consists of N Ising spins, labelled $i = 1, \dots, N$ with $N \in \mathbb{N}$. The spins interact through a mean-field Hamiltonian —that is, a Hamiltonian involving no geometry and no sense of neighborhood, in which each spin interacts equally with all other spins—. The Curie-Weiss Hamiltonian is

$$H^N(\sigma) := -\frac{J}{2N} \sum_{i,j=1}^N \sigma_i \sigma_j - h \sum_{i=1}^N \sigma_i, \quad \sigma \in \Omega_N, \quad (1.1)$$

where $J > 0$ is the (ferromagnetic) pair potential, $h \in \mathbb{R}$ is the (external) magnetic field, $\Omega_N := \{-1, +1\}^N$ is the spin configuration space, and $\sigma := (\sigma_i)_{i=1}^N$ is the spin configuration. The Gibbs measure associated with H^N is

$$\mu^N(\sigma) := \frac{e^{-H^N(\sigma)}}{Z^N}, \quad \sigma \in \Omega_N, \quad (1.2)$$

with Z^N the normalizing partition sum.

We allow this model to evolve according to an *independent spin-flip* dynamics, that is, a dynamics defined by the generator L_N given by (see Liggett [14] for more background)

$$(L_N f)(\sigma) := \sum_{i=1}^N [f(\sigma^i) - f(\sigma)], \quad f: \Omega_N \rightarrow \mathbb{R}, \quad (1.3)$$

where σ^i denotes the configuration obtained from σ by flipping the spin with label i . The resulting random variables $\sigma(t) := (\sigma_i(t))_{i=1}^N$ constitute a continuous-time Markov chain on Ω_N . We write μ_t^N to denote the measure on Ω_N at time t when the initial measure is μ^N and abbreviate $\mu_t := (\mu_t^N)_{N \in \mathbb{N}}$.

1.3.2 Empirical magnetization

To emphasize its mean-field character, it is convenient to write the Hamiltonian (1.1) in the form

$$H^N(\sigma) = N\bar{H}(m_N(\sigma)) \quad (1.4)$$

where

$$\bar{H}(x) := -\frac{1}{2}Jx^2 - hx, \quad x \in \mathbb{R}. \quad (1.5)$$

and

$$m_N(\sigma) := \frac{1}{N} \sum_{i=1}^N \sigma_i \quad (1.6)$$

is the empirical magnetization of $\sigma \in \Omega_N$, which takes values in the set $\mathcal{M}_N := \{-1, -1 + 2N^{-1}, \dots, +1 - 2N^{-1}, +1\}$. The Gibbs measure on Ω_N induces a Gibbs measure on \mathcal{M}_N given by

$$\bar{\mu}^N(m) := \binom{N}{\frac{1+m}{2}N} \frac{e^{-N\bar{H}(m)}}{\bar{Z}^N}, \quad m \in \mathcal{M}_N, \quad (1.7)$$

where \bar{Z}^N is the normalizing partition sum.

The independent (infinite-temperature) dynamics has the simplifying feature of preserving the mean-field character of the model. In fact, the dynamics on Ω_N induces a dynamics on \mathcal{M}_N , which is a continuous-time Markov chain $(m_t^N)_{t \geq 0}$ with generator \bar{L}_N given by

$$(\bar{L}_N f)(m) := \frac{1+m}{2} N[f(m-2N^{-1}) - f(m)] + \frac{1-m}{2} N[f(m+2N^{-1}) - f(m)], \quad (1.8)$$

for $f: \mathcal{M}_N \rightarrow \mathbb{R}$. Adapting our previous notation we denote $\bar{\mu}_t^N$ the measure on \mathcal{M}_N at time t , and abbreviate $\bar{\mu}_t := (\bar{\mu}_t^N)_{N \in \mathbb{N}}$. Due to permutation invariance, μ_t^N characterizes $\bar{\mu}_t^N$ and vice versa, for each N and t . We write P^N to denote the law of $(m_t^N)_{t \geq 0}$, which lives on the space of càdlàg trajectories $D_{[0, \infty)}([-1, +1])$ endowed with the Skorohod topology.

1.3.3 Bad magnetizations

Non-Gibbsianness shows up through discontinuities with respect to boundary conditions of finite-volume conditional probabilities. For the Curie-Weiss model it is enough to consider the single-spin conditional probabilities

$$\gamma_t^N(\sigma_1 | \alpha_{N-1}) := \mu_t^N(\sigma_1 | \sigma_{N-1}), \quad (1.9)$$

defined for $\sigma_1 \in \{-1, +1\}$ and $\alpha_{N-1} \in \mathcal{M}_{N-1}$, and any spin configuration $\sigma_{N-1} \in \Omega_{N-1}$ such that $m_{N-1}(\sigma_{N-1}) = \alpha_{N-1}$. By permutation invariance, (1.9) does not depend on the choice of σ_{N-1} .

The central definition for our purposes is the following.

Definition 1.1. (Külske and Le Ny [9]) Fix $t \geq 0$.

(a) A magnetization $\alpha \in [-1, +1]$ is said to be good for μ_t if there exists a neighborhood \mathcal{N}_α of α such that

$$\gamma_t(\cdot | \bar{\alpha}) := \lim_{N \rightarrow \infty} \gamma_t^N(\cdot | \alpha_{N-1}), \quad (1.10)$$

exists for all $\bar{\alpha}$ in \mathcal{N}_α and all $(\alpha_N)_{N \in \mathbb{N}}$ such that $\alpha_N \in \mathcal{M}_N$ for all $N \in \mathbb{N}$ and $\lim_{N \rightarrow \infty} \alpha_N = \bar{\alpha}$, and is independent of the choice of $(\alpha_N)_{N \in \mathbb{N}}$. The limit is called the specification kernel. In particular, $\bar{\alpha} \mapsto \gamma_t(\cdot | \bar{\alpha})$ is continuous at $\bar{\alpha} = \alpha$.

(b) A magnetization $\alpha \in [-1, +1]$ is called bad if it is not good.

(c) μ_t is called Gibbs if it has no bad magnetizations.

1.4 Path large deviations and link to specification kernel

The main point of our work is our relation between path large deviations and non-Gibbsianness. For the convenience of the reader, let us recall some basic large deviation results for the Curie-Weiss model. For background on large deviation theory, see e.g. den Hollander [8].

1.4.1 Path large deviation principle

Let us recall that a family of measures ν^N on a Borel measure space satisfies a *large deviation principle* with rate function I and speed N if the following two conditions are satisfied:

$$\liminf_{N \rightarrow \infty} \frac{1}{N} \log \nu^N(A) \geq - \inf_{x \in A} I(x) \quad \text{for } A \text{ open} \quad (1.11)$$

$$\limsup_{N \rightarrow \infty} \frac{1}{N} \log \nu^N(A) \leq - \sup_{x \in A} I(x) \quad \text{for } A \text{ closed} \quad (1.12)$$

The proof of the following proposition is elementary and can be found in many references. The indices S and D stand for *static* and *dynamic*.

Proposition 1.2. (Ermolaev and Külske [7], Enter, Fernández, den Hollander and Redig [3])

(i) $(\bar{\mu}^N)_{N \in \mathbb{N}}$ satisfies the large deviation principle on $[-1, +1]$ with rate N and rate function $I_S - \inf(I_S)$ given by

$$I_S(m) := \bar{H}(m) + \bar{I}(m), \quad \bar{I}(m) := \frac{1+m}{2} \log(1+m) + \frac{1-m}{2} \log(1-m). \quad (1.13)$$

(ii) For every $T > 0$, the restriction of $(P^N)_{N \in \mathbb{N}}$ to the time interval $[0, T]$ satisfies the large deviation principle on $D_{[0, T]}([-1, +1])$ with rate N and rate function $I^T - \inf(I^T)$ given by

$$I^T(\varphi) := I_S(\varphi(0)) + I_D^T(\varphi), \quad (1.14)$$

where

$$I_D^T(\varphi) := \begin{cases} \int_0^T L(\varphi(s), \dot{\varphi}(s)) ds & \text{if } \dot{\varphi} \text{ exists,} \\ \infty & \text{otherwise,} \end{cases} \quad (1.15)$$

is the action integral with Lagrangian

$$L(m, \dot{m}) := -\frac{1}{2} \sqrt{4(1-m^2) + \dot{m}^2} + \frac{1}{2} \dot{m} \log \left(\frac{\sqrt{4(1-m^2) + \dot{m}^2} + \dot{m}}{2(1-m)} \right) + 1. \quad (1.16)$$

Let

$$Q_{t,\alpha}^N(m) := P^N(m_N(0) = m \mid m_N(t) = \alpha), \quad m \in \mathcal{M}_N \quad (1.17)$$

be the conditional distribution of the magnetization at time 0 given that the magnetization at time t is α . The contraction principle applied to Proposition 1.2(ii) implies the following large deviation principle.

Corollary 1.3. *For every $t \geq 0$ and $\alpha \in [-1, +1]$, $(Q_{t,\alpha}^N)_{N \in \mathbb{N}}$ satisfies the large deviation principle on $[-1, +1]$ with rate N and rate function $C_{t,\alpha} - \inf(C_{t,\alpha})$ given by*

$$C_{t,\alpha}(m) := \inf_{\substack{\varphi: \varphi(0)=m, \\ \varphi(t)=\alpha}} I^t(\varphi). \quad (1.18)$$

Note that

$$\inf_{m \in [-1, +1]} C_{t,\alpha}(m) = \inf_{m \in [-1, +1]} \inf_{\substack{\varphi: \varphi(0)=m, \\ \varphi(t)=\alpha}} I^t(\varphi) = \inf_{\varphi: \varphi(t)=\alpha} I^t(\varphi). \quad (1.19)$$

1.4.2 Link to specification kernel

The following proposition provides the fundamental link between the specification kernel in (1.10) and the minimizer of (1.19) when it is unique, and is a straightforward generalization to arbitrary magnetic field of a result for zero magnetic field stated and proved in Ermolaev and Külske [7].

Proposition 1.4. *Fix $t \geq 0$ and $\alpha \in [-1, +1]$. Suppose that (1.19) has a unique minimizing path $(\hat{\varphi}_{t,\alpha}(s))_{0 \leq s \leq t}$. Then the specification kernel equals*

$$\gamma_t(z \mid \alpha) = \frac{\sum_{x \in \{-1, +1\}} e^{x[J\hat{\varphi}_{t,\alpha}(0)+h]} p_t(x, z)}{\sum_{x, y \in \{-1, +1\}} e^{x[J\hat{\varphi}_{t,\alpha}(0)+h]} p_t(x, y)}, \quad z \in \{-1, +1\}, \quad (1.20)$$

where $p_t(\cdot, \cdot)$ is the transition kernel of the continuous-time Markov chain on $\{-1, +1\}$ jumping at rate 1, given by $p_t(1, 1) = p_t(-1, -1) = e^{-t} \cosh(t)$ and $p_t(-1, +1) = p_t(1, -1) = e^{-t} \sinh(t)$.

Remark: Note that the expression in the right-hand side of (1.20) depends on the optimal trajectory only via its initial value $\hat{\varphi}_{t,\alpha}(0)$. Thus, (1.20) has the form

$$\gamma_t(z \mid \alpha) = \Gamma_t(z, J\hat{\varphi}_{t,\alpha}(0) + h), \quad (1.21)$$

where $\hat{\varphi}_{t,\alpha}(0)$ is the unique global minimizer of $m \mapsto C_{t,\alpha}(m)$ and $m \mapsto \Gamma_t(z, m)$ is continuous and strictly increasing (strictly decreasing) for $z = 1$ ($z = -1$).

1.4.3 Reduction

The next proposition allows us to reduce (1.19) to a one-dimensional variational problem. Consider the equation

$$k^{J,h}(m) = l_{t,\alpha}(m) \quad (1.22)$$

with

$$\begin{aligned} k^{J,h}(m) &:= a^J(m) \cosh(2h) + b^J(m) \sinh(2h), \\ l_{t,\alpha}(m) &:= m \coth(2t) - \alpha \operatorname{csch}(2t), \end{aligned} \quad (1.23)$$

where

$$\begin{aligned} a^J(m) &:= \sinh(2Jm) - m \cosh(2Jm), \\ b^J(m) &:= \cosh(2Jm) - m \sinh(2Jm). \end{aligned} \quad (1.24)$$

Proposition 1.5. *Let $C_{t,\alpha}$ be as in (1.18). Then, for every $t \geq 0$ and $\alpha \in [-1, +1]$,*

$$\begin{aligned} C_{t,\alpha}(m) &= I_S(m) \\ &+ \frac{1}{4} \left\{ 4t + \log \left(\frac{1 - \alpha^2}{1 - m^2} \right) + \log \left(\left[\frac{1 - R - 2C_1 \alpha e^{-2t}}{1 + R - 2C_1 \alpha e^{-2t}} \right] \left[\frac{1 + R - 2C_1 m}{1 - R - 2C_1 m} \right] \right) \right. \\ &\left. + 2 \left[\alpha \log \left(\frac{R - C_1 e^{-2t} + C_2 e^{2t}}{1 - \alpha} \right) - m \log \left(\frac{R - C_1 + C_2}{1 - m} \right) \right] \right\} \end{aligned} \quad (1.25)$$

with

$$\begin{aligned} C_1 &= C_1(t, \alpha, m) := \frac{m e^{2t} - \alpha}{e^{2t} - e^{-2t}}, \\ C_2 &= C_2(t, \alpha, m) := \frac{\alpha - m e^{-2t}}{e^{2t} - e^{-2t}}, \\ R &= R(C_1, C_2) := \sqrt{1 - 4C_1 C_2}. \end{aligned} \quad (1.26)$$

Furthermore, the critical points of $C_{t,\alpha}$ are the solutions of (1.22). Hence,

$$\inf_{\varphi: \varphi(t)=\alpha} I^t(\varphi) = \min_{m \text{ solves (1.22)}} C_{t,\alpha}(m), \quad (1.27)$$

and the constrained minimizing trajectories are of the form

$$\hat{\varphi}_{t,\alpha}^{\hat{m}}(s) := \operatorname{csch}(2t) \left\{ m \sinh(2(t-s)) + \alpha \sinh(2s) \right\} \quad 0 \leq s \leq t \quad (1.28)$$

$$\hat{m} = \hat{m}(t, \alpha) = \operatorname{argmin} \left[C_{t,\alpha} \Big|_{\text{solutions of (1.22)}} \right]. \quad (1.29)$$

The identities

$$k^{J,h}(m) = 2 \cosh^2(Jm + h) [\tanh(Jm + h) - m] + m \quad (1.30)$$

and

$$\lim_{t \rightarrow \infty} l_{t,\alpha}(m) = m \quad (1.31)$$

imply that in the limit $t \rightarrow \infty$ (1.22) reduces to $\tanh(Jm + h) = m$. This is the equation for the spontaneous magnetization of the Curie-Weiss model with parameters J, h . This equation has always at least one solution and the value

$$m^\infty = m^\infty(J, h) := \text{the largest solution of the equation } \tanh(Jm + h) = m \quad (1.32)$$

is well known to be strictly positive if $h > 0$ or if $J > 1$. In these regimes, the standard Curie-Weiss graphical argument shows that, for $m > 0$,

$$k^{J,h}(m) \underset{>}{\leq} m \iff m \underset{<}{\geq} m^\infty. \quad (1.33)$$

We also remark that when $t \rightarrow 0$ the function $l_{t,\alpha}$ converges to the line defined by the equation $m = \alpha$. This implies that for short times there is a unique solution of (1.22) and it is close to α .

1.5 Main results

In Section 1.5.1 we state the equivalence of non-Gibbs and bifurcation that lies at the heart of the program outlined in [3] (Theorem 1.6). In Section 1.5.2 we introduce some notation. In Section 1.5.3 we identify the optimal trajectories for $\alpha = 0$, $h = 0$ (Theorems 1.7–1.8). In Section 1.5.4 we extend this identification to $\alpha \in [-1, +1]$, $h \in \mathbb{R}$ (Theorem 1.9). In Section 1.5.5 we summarize the consequences for Gibbs versus non-Gibbs (Corollary 1.10).

1.5.1 Equivalence of non-Gibbs and bifurcation

The following theorem proves the long suspected equivalence between dynamical non-Gibbsianness, i.e., discontinuity of $\alpha \mapsto \gamma_t(\cdot \mid \alpha)$ at α_0 , and non-uniqueness of the global minimizer of $m \mapsto C_{t,\alpha_0}(m)$, i.e., the occurrence of more than one possible history for the same α .

Theorem 1.6. *$\alpha \mapsto \gamma_t(\sigma \mid \alpha)$ is continuous at α_0 if and only if $\inf_{\varphi: \varphi(t)=\alpha_0} I^t(\varphi)$ has a unique minimizing path or, equivalently, $\inf_{m \in [-1, +1]} C_{t,\alpha_0}(m)$ has a unique minimizing magnetization.*

1.5.2 Notation

Due to relation (1.27), our analysis focusses on the different solutions of (1.22) obtained as t, α are varied. In particular, we must determine which of them are minima of the variational problem in (1.19). We write

$$\Delta_{t,\alpha} := \text{the set of global minimizers of } C_{t,\alpha}. \quad (1.34)$$

For brevity, when α is kept fixed and $\Delta_{t,\alpha}$ is a singleton $\{\hat{m}(t, \alpha)\}$ for each t , we write $\hat{m}(t)$ instead of $\hat{m}(t, \alpha)$. When $h, \alpha = 0$, by symmetry we have $\Delta_{t,0} = \{0\}$ or $\Delta_{t,\alpha} = \{\pm \hat{m}(t)\}$, where in the last case we denote by $\hat{m}(t)$ the unique positive global minimizer. If both the initial and final magnetizations are fixed, then there is a unique minimizer that we denote as in (1.28). That is,

$$\hat{\varphi}_{t,\alpha}^m := \underset{\substack{\varphi: \varphi(0)=m, \\ \varphi(t)=\alpha}}{\operatorname{argmin}} I_D^t(\varphi) \quad (1.35)$$

for $m, \alpha \in [-1, +1]$. We emphasize that, by definition, $C_{t,\alpha}(m) = I^t(\hat{\varphi}_{t,\alpha}^m)$ and $\hat{\varphi}_{t,\alpha}(s) = \hat{\varphi}_{t,\alpha}^{\hat{m}(t,\alpha)}(s)$, $s \in [0, t]$. In particular $\hat{m}(t, \alpha) = \hat{\varphi}_{t,\alpha}(0)$.

1.5.3 Optimal trajectories for $\alpha = 0$, $h = 0$

The following theorem refers to a critical time

$$\Psi_c = \Psi_c(J) := \begin{cases} \frac{1}{2} \operatorname{arccoth}(2J - 1) & \text{if } 1 < J \leq \frac{3}{2}, \\ t_* & \text{if } J > \frac{3}{2}, \end{cases} \quad (1.36)$$

where $t_* = t_*(J)$ is implicitly calculable: $t_* = t(m_*)$ where the function $t(m)$ is defined in (2.11) below and $m_* = m_*(J)$ is the solution of (2.18).

Theorem 1.7. (See Fig. 2.) *Consider $\alpha = 0$ and $h = 0$.*

(i) If $0 < J \leq 1$, then

$$\Delta_{t,0} = \{0\}, \quad \forall t \geq 0. \quad (1.37)$$

(ii) If $1 < J \leq \frac{3}{2}$, then

$$\Delta_{t,0} = \begin{cases} \{0\} & \text{if } 0 \leq t \leq \Psi_c, \\ \{\pm \hat{m}(t)\} & \text{if } t > \Psi_c, \end{cases} \quad (1.38)$$

where $t \mapsto \hat{m}(t)$ is continuous and strictly increasing on $[\Psi_c, \infty)$ with $\hat{m}(\Psi_c) = 0$.

(iii) If $J > \frac{3}{2}$, then

$$\Delta_{t,0} = \begin{cases} \{0\} & \text{if } 0 \leq t < \Psi_c, \\ \{\pm \hat{m}(t)\} & \text{if } t \geq \Psi_c, \end{cases} \quad (1.39)$$

where $t \mapsto \hat{m}(t)$ is continuous and strictly increasing on $[\Psi_c, \infty)$ with $\hat{m}(\Psi_c) =: m_* > 0$.

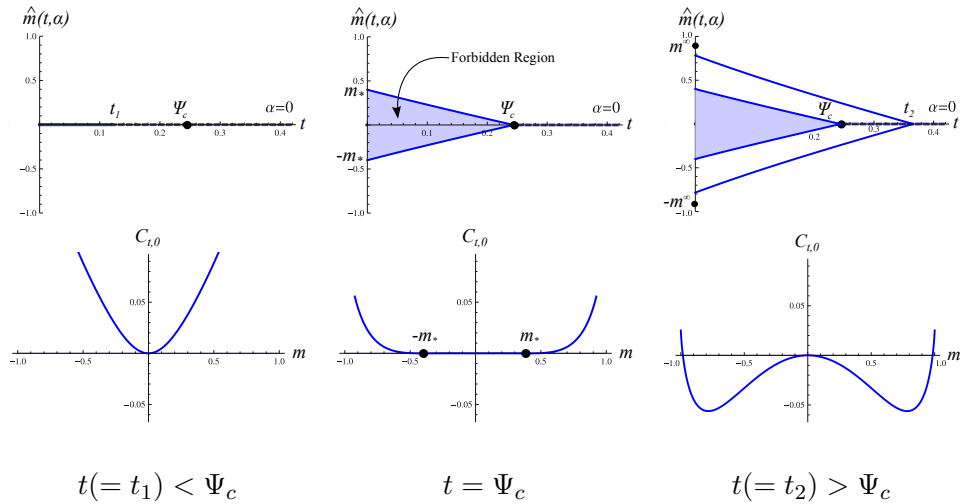


Figure 2: Illustration of Theorem 1.7. *First row:* Time evolution of the minimizing trajectories $\pm(\hat{\varphi}_{t,0}(s))_{0 \leq s \leq t}$ for $t < \Psi_c$, $t = \Psi_c$ and $t > \Psi_c$ for an initial Curie-Weiss model with $(J, h) = (1.6, 0)$ [regime (iii) in the Theorem]. The shaded cone is the forbidden region. *Second row:* Plot of $m \mapsto C_{t,0}(m)$ for the same times and parameter values.

Let $\Lambda_{t,0}(J)$ denote the cone between the trajectories $\pm \hat{\varphi}_{t,0}$. As a consequence of the previous theorem, *no* minimal trajectory conditioned in t' with $t' \geq t$ can intersect the interior of this region. Such a cone corresponds, therefore, to a *forbidden region*. Forbidden regions grow, in a nested fashion, as the conditioning time t grows. There is, however, a distinctive difference between regimes (ii) and (iii) in the previous theorem: In Regime (ii) the forbidden region opens up *continuously* after Ψ_c , while in Regime (iii) it opens up *discontinuously*. These facts are summarized in the following theorem.

Theorem 1.8. *Suppose that $\alpha = 0$ and $h = 0$.*

- (i) $J \mapsto m_*(J)$ is strictly increasing on $(\frac{3}{2}, \infty)$.
- (ii) $J \mapsto \Psi_c(J)$ is strictly decreasing on $(1, \infty)$.
- (iii) $J \mapsto \Lambda_{t,0}(J)$ is left-continuous at $J = \frac{3}{2}$ for all $t > \Psi_c(\frac{3}{2})$.
- (iv) $J \mapsto \Lambda_{\Psi_c(\bar{J}),0}(J)$ is right-continuous at $J = \bar{J}$ for all $\bar{J} > \frac{3}{2}$.
- (v) For every $J \leq 3/2$ the map $t \mapsto \Lambda_{t,0}(J)$ is continuous.

(vi) For every $J > 3/2$ the map $t \mapsto \Lambda_{t,0}(J)$ is continuous except at $t = \Psi_c$ where it exhibits a right-continuous jump.

1.5.4 Optimal trajectories for $\alpha \in [-1, +1]$, $h \in \mathbb{R}$

For fixed (J, h) and α we say that there is (See Fig. 3):

- *No bifurcation* if $\Delta_{t,\alpha} = \{\hat{m}(t, \alpha)\}$, for all $t \geq 0$ and the map $t \mapsto \hat{m}(t, \alpha)$ is continuous on $[0, \infty)$.
- *Bifurcation* when there exists a $0 < t_B < \infty$ such that $t \mapsto \hat{m}(t, \alpha)$ is continuous except at $t = t_B$ and $|\Delta_{t_B, \alpha}| = 2$.
- *Double bifurcation* if there exist times $0 < s_B < t_B < \infty$ such that $t \mapsto \hat{m}(t, \alpha)$ is continuous except at $t = t_B$ and $t = s_B$, and $|\Delta_{s_B, \alpha}| = |\Delta_{t_B, \alpha}| = 2$.
- *Trifurcation* if there exists a $0 < t_T < \infty$ such that $t \mapsto \hat{m}(t, \alpha)$ is continuous except at $t = t_T$ and $|\Delta_{t_T, \alpha}| = 3$.

The bifurcation times t_B and s_B , the trifurcation time t_T and the trifurcation magnetization M_T (defined below) all depend on J, h .

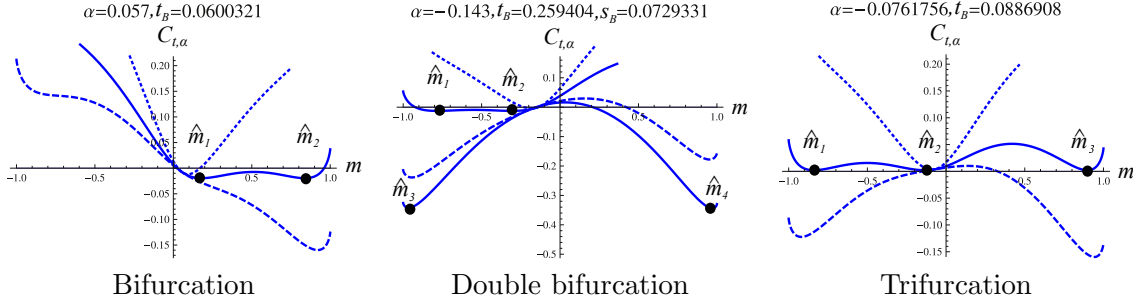


Figure 3: Different scenarios for the evolution in time of $m \mapsto C_{t,\alpha}(m)$. Drawn lines: $t = t_B$, $t = t_B, s_B$, $t = t_T$ (times at which multiple global minima occur or, equivalently, discontinuity points of $t \mapsto \hat{m}(t, \alpha)$). Dotted lines: earlier time. Dashed lines: later time.

The following theorem summarizes the behaviour of $\Delta_{t,\alpha}$ (and therefore of the minimizing trajectories $\hat{\varphi}_{t,\alpha}$) for different t, α . For $J > \frac{3}{2}$, let

$$F(m) := \frac{mk'^{J,h}(m) - k^{J,h}(m)}{\operatorname{csch}[\operatorname{arccoth}(k'^{J,h}(m))]},$$

$$U_B = U_B(J, h) := \max_{m \in [0,1]} F(m), \quad (1.40)$$

$$L_B = L_B(J, h) := \min_{m \in [-1,0]} F(m).$$

Theorem 1.9. (See Figs. 3–4.)

- (1) Suppose that $k^{J,h}(\alpha) \neq 0$.
 - (1a) If $k^{J,h}(\alpha) > 0$ and $\alpha > 0$, then there are $m_R^+ > 0$ and $t_R = t_R(m_R^+) > 0$ (implicitly

calculable from (3.8)) such that $t \mapsto \hat{m}(t)$ is strictly increasing on $[0, t_R]$ and strictly decreasing on $[t_R, \infty)$ with $\hat{m}(t_R) = m_R^+ > m^\infty$.

(1b) If $k^{J,h}(\alpha) < 0$ and $\alpha > 0$, then $t \mapsto \hat{m}(t)$ is strictly decreasing on $[0, \infty)$.

(1c) If $k^{J,h}(\alpha) > 0$ and $\alpha < 0$, then $t \mapsto \hat{m}(t)$ is strictly increasing on $[0, \infty)$.

(1d) If $k^{J,h}(\alpha) < 0$ and $\alpha < 0$, then there are $m_R^- > 0$ and $t_R = t_R(m_R^-) > 0$ (implicitly calculable from (3.9)) such that $t \mapsto \hat{m}(t)$ is strictly decreasing on $[0, t_R]$ and strictly increasing on $[t_R, \infty)$ with $\hat{m}(t_R) = m_R^- < \alpha$.

In all cases $\hat{m}(0) = \alpha$ and $\lim_{t \rightarrow \infty} \hat{m}(t) = m^\infty$.

(2) Suppose that $h = 0$.

(2a) If $0 < J \leq 1$, then there is no bifurcation.

(2b) If $1 < J \leq \frac{3}{2}$, then there is bifurcation only for $\alpha = 0$.

(2c) If $J > \frac{3}{2}$, then there is bifurcation if $\alpha \in (-U_B, U_B)$ and no bifurcation otherwise.

(3) Suppose that $h > 0$.

(3a) If $0 < J \leq 1$, then there is no bifurcation.

(3b) If $1 < J \leq \frac{3}{2}$, then there is bifurcation for $\alpha \in [-1, U_B)$ and no bifurcation for $\alpha \in [U_B, 1]$.

(3c) If $J > \frac{3}{2}$, then there exists a $h_* = h_*(J) > 0$ such that

- for every $0 < h < h_*$ there exists a $M_T \in (L_B, U_B)$ with $M_T < 0$ such that there is

- * no bifurcation for $\alpha \in [U_B, 1]$,
- * bifurcation for $\alpha \in (M_T, U_B)$,
- * trifurcation for $\alpha = M_T$,
- * double bifurcation for $\alpha \in (L_B, M_T)$,
- * bifurcation for $\alpha \in [-1, L_B]$.

- for every $h \geq h_*$ the behavior is the same as in (3b).

In all cases $\alpha \mapsto t_B(\alpha)$ is continuous and decreasing and $\alpha \mapsto s_B(\alpha)$ is continuous and increasing.

Theorem 1.9 gives a complete picture of the bifurcation scenario. Regime (1) —which includes cases with zero and nonzero magnetic field— describes two types of behavior of optimal magnetization trajectories: monotone trajectories [cases (1b) and (1c)] and trajectories with overshoot [cases (1a) and (1d)]. In the latter, $\hat{m}(t)$ increases to some magnetization m_R^+ larger (m_R^- smaller) than m^∞ and afterwards decreases (increases) to m^∞ . Regimes (2) and (3) refer to the existence of bifurcations and trifurcations. We observe that the different bifurcation behaviors —no bifurcation, single and double bifurcation— hold for whole intervals of the conditioning magnetization. In contrast, trifurcation appears at a single final magnetization for each $h \neq 0$.

1.5.5 Gibbs versus non-Gibbs

Theorem 1.6 establishes the equivalence of bifurcation and discontinuity of specifications, as proposed in the program put forward in [3]. Due to this equivalence, the following corollary provides a full characterization of the different Gibbs–nonGibbs scenarios appearing during the infinite-temperature evolution of the Curie-Weiss model. Let

$$0 < \Psi_U := t_B(U_B) < \Psi_T := t_B(M_T) < \Psi_L := t_B(L_B) < \Psi_* := t_B(-1), \quad (1.41)$$

and let M_B be the solution of $t_B(M_B) = \Psi_L$. Denote $\mathcal{D}_t \subseteq [-1, +1]$ the set of α -values for which $\alpha \mapsto \gamma_t(\cdot|\alpha)$ is discontinuous.

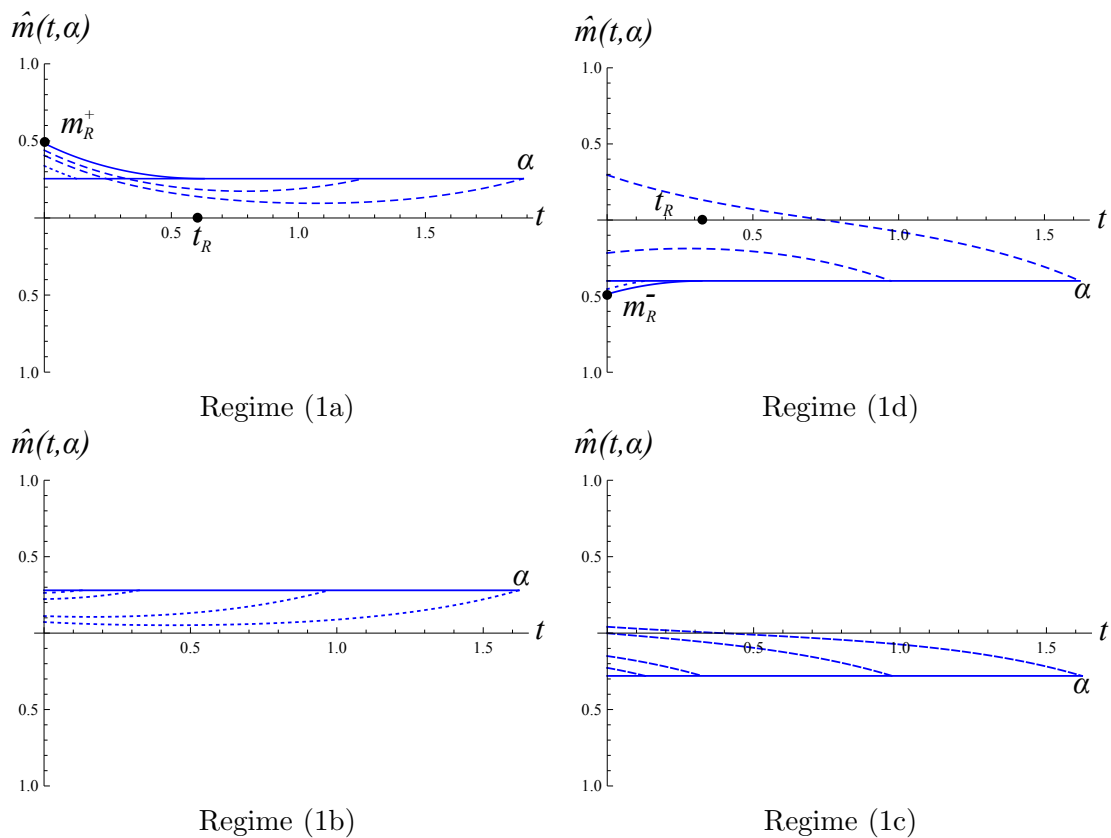


Figure 4: Different regimes of Theorem 1.9. Evolution in time of the minimizing trajectories $\pm(\hat{\varphi}_{t,\alpha}(s))_{0 \leq s \leq t}$ for $t < t_R$ (dotted), $t = t_R$ (drawn), $t > t_R$ (dashed).

Corollary 1.10. (See Fig. 5.)

- (1) *Let $h = 0$.*
- (1a) *If $0 < J \leq 1$, the evolved measure μ_t is Gibbs for all $t \geq 0$.*
- (1b) *If $1 < J \leq \frac{3}{2}$, then μ_t is*
- *Gibbs for $0 \leq t \leq \Psi_c$,*
 - *non-Gibbs for $t > \Psi_c$ with $\mathcal{D}_t = \{0\}$.*
- (1c) *If $J > \frac{3}{2}$, then μ_t is*
- *Gibbs for $0 \leq t \leq \Psi_U$,*
 - *non-Gibbs for $t > \Psi_U$ with*
 - * *$\mathcal{D}_t = \{\pm\alpha\}$ for some $\alpha \in (-U_B, U_B)$ if $\Psi_U < t < \Psi_c$,*
 - * *$\mathcal{D}_t = \{0\}$ if $t \geq \Psi_c$.*
- (2) *Let $h > 0$.*
- (2a) *If $0 < J \leq 1$, then μ_t is Gibbs for $t \geq 0$.*
- (2b) *If $1 < J \leq \frac{3}{2}$, then μ_t is*
- *Gibbs for $0 \leq t \leq \Psi_U$,*
 - *non-Gibbs for $\Psi_U < t \leq \Psi_*$ with $\mathcal{D}_t = \{\alpha\}$ for some $\alpha \in [-1, U_B)$,*
 - *Gibbs for $t > \Psi_*$.*
- (2c) *If $J > \frac{3}{2}$ and $h < h^*$ small enough, then μ_t is*
- *Gibbs for $0 \leq t \leq \Psi_U$,*
 - *non-Gibbs for $\Psi_U < t \leq \Psi_*$ with*
 - * *$\mathcal{D}_t = \{\alpha\}$ for some $\alpha \in [M_B, U_B)$ if $\Psi_U < t \leq \Psi_L$,*
 - * *$\mathcal{D}_t = \{\alpha^1, \alpha^2\}$ for some $\alpha^1, \alpha^2 \in (L_B, M_B)$ if $\Psi_L < t < \Psi_T$,*
 - * *$\mathcal{D}_t = \{\alpha\}$ for some $\alpha \in [-1, M_T]$ if $\Psi_T \leq t \leq \Psi_*$.*
 - *Gibbs for $t > \Psi_*$. If $h \geq h^*$, then the behaviour is as in (2b).*

In all cases $\alpha^1, \alpha^2, \alpha$ depend on (t, J, h) .

2 Proof of Proposition 1.5 and Theorems 1.6–1.8

Proposition 1.5 is proven in Section 2.1, Theorems 1.6–1.8 are proven in Sections 2.2–2.4.

2.1 Proof of Proposition 1.5

Proof. First note that, by (1.14),

$$\inf_{\varphi: \varphi(t)=\alpha} I^t(\varphi) = \inf_{m \in [-1, +1]} \left\{ I_S(m) + \inf_{\substack{\varphi: \varphi(0)=m, \\ \varphi(t)=\alpha}} I_D^t(\varphi) \right\}. \quad (2.1)$$

It follows from (1.14–1.15) and the calculus of variations that the stationary points of the right-hand side of (2.1) are given by the Euler-Lagrange equation, complemented with a free-left-end condition and a fixed-right-end condition:

$$\begin{aligned} \frac{\partial}{\partial s} \frac{\partial L}{\partial \dot{m}}(\varphi(s), \dot{\varphi}(s)) &= \frac{\partial L}{\partial m}(\varphi(s), \dot{\varphi}(s)), & s \in (0, t), \\ \frac{\partial L}{\partial \dot{m}}(\varphi(s), \dot{\varphi}(s)) \Big|_{s=0} &= \frac{\partial I_S}{\partial m}(\varphi(s)) \Big|_{s=0}, \\ \varphi(t) &= \alpha. \end{aligned} \quad (2.2)$$

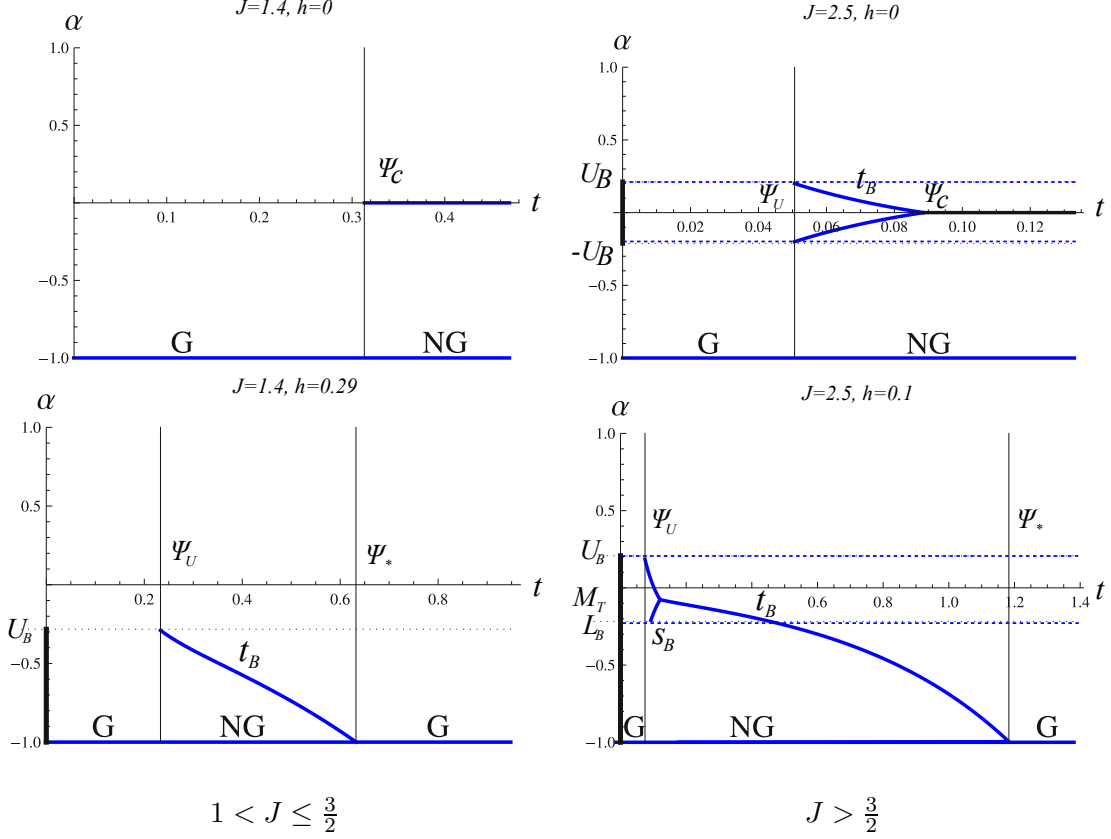


Figure 5: Summary of Corollary 1.10: Time versus bad magnetizations for different regimes. On the vertical α -axis, indicated by a thick line, is the set of bad magnetizations. G=Gibbs, NG=non-Gibbs.

The first and the third equation in (2.2) come from the third infimum in (2.1) and, together with (1.16), determine the form (1.28) of the stationary trajectory. Inserting this form into (1.14) we identify

$$I^t(\hat{\varphi}_{t,\alpha}^m) = C_{t,\alpha}(m), \quad (2.3)$$

as stated in (1.25)–(1.26). This identity reduces (1.19) to a one-dimensional variational problem,

$$\inf_{\varphi: \varphi(t)=\alpha} I^t(\varphi) = \inf_{m \in [-1,+1]} I^t(\hat{\varphi}_{t,\alpha}^m) = \inf_{m \in [-1,+1]} C_{t,\alpha}(m) \quad (2.4)$$

The second equation in (2.2) corresponds to the second infimum in (2.1) or, equivalently, to the rightmost infima in (2.4). It gives a trade-off between the static and the dynamic cost, establishing a relation between the initial magnetization and the initial derivative. After some manipulations this equation can be written in the form

$$-\frac{1}{2}q = a^J(m) \cosh(2h) + b^J(m) \sinh(2h), \quad m = \hat{\varphi}_{t,\alpha}^m(0), \quad q = \dot{\hat{\varphi}}_{t,\alpha}^m(0). \quad (2.5)$$

Differentiating (1.28), we get

$$\dot{\hat{\varphi}}_{t,\alpha}^m(s) = 2 \operatorname{csch}(2t) \left\{ \alpha \cosh(2s) - m \cosh(2(t-s)) \right\}, \quad (2.6)$$

and eliminating q from the last identity in (2.5) in favor of t and α , we conclude that m must be a solution of (1.22). Imposing this restriction to the chain of identities (2.4) we obtain (1.27) and hence (1.28)–(1.29). \square

From now on our arguments rely on the study of (1.22), combined with continuity properties of $C_{t,\alpha}$ as a function of t, α .

2.2 Proof of Theorem 1.6

Equation (1.20) follows in the same way as in the proof of Theorem 2.5 in [7]. Having disposed of this identity, we can now proceed to prove the equivalence. The proof relies on the following lemma.

Lemma 2.1. *For any $t > 0$ and $\alpha_0 \in [-1, +1]$, there exists an open neighbourhood $\mathcal{N}_{\alpha_0} \neq \emptyset$ of the later, such that for all $\alpha \in \mathcal{N}_{\alpha_0} \setminus \{\alpha_0\}$*

1. $C_{t,\alpha}$ has only one global minimum, namely, $\hat{m}(t, \alpha)$.
2. $\tilde{\alpha} \mapsto \hat{m}(t, \tilde{\alpha})$ is continuous at α . If C_{t,α_0} has a unique global minimum, the continuity is also valid at $\alpha = \alpha_0$.
3. If C_{t,α_0} has multiple global minima, for two of them, namely, $\hat{m}_A(t, \alpha_0)$ and $\hat{m}_B(t, \alpha_0)$

$$\lim_{\alpha \downarrow \alpha_0} \hat{m}(t, \alpha) = \hat{m}_A(t, \alpha_0) \quad \text{and} \quad \lim_{\alpha \uparrow \alpha_0} \hat{m}(t, \alpha) = \hat{m}_B(t, \alpha_0).$$

Proof. A straightforward study of (1.22) shows that $C_{t,\alpha}$ has a finite number of critical points, for every fixed choice of J, h, t, α .

Clearly, $\alpha \mapsto C_{t,\alpha}$ and $\alpha \mapsto l_{t,\alpha}$ are continuous with respect to the infinity norm in $C([-1, +1], \mathbb{R})$. This, together with the fact that the left-hand side of (1.22) does not depend on α , implies continuity of any critical point with respect to α .

Let $\hat{m}_i(t, \alpha_0)$, $i = 1, \dots, v$ be the global minima of C_{t,α_0} . By continuity of the critical points, there exists a neighbourhood $\tilde{\mathcal{N}}_{\alpha_0}$ and smooth functions $\tilde{\mathcal{N}}_{\alpha_0} \ni \alpha \mapsto \bar{m}_i(t, \alpha)$, $i = 1, \dots, v$, such that

- i. $\bar{m}_i(t, \alpha)$ are local minima of $C_{t,\alpha}$,
- ii. $\lim_{\alpha \rightarrow \alpha_0} \bar{m}_i(t, \alpha) = \hat{m}_i(t, \alpha_0)$.

This properties proves the lemma if $v = 1$. Otherwise, let

$$B_i(\alpha) := C_{t,\alpha}(\bar{m}_i(t, \alpha)).$$

The minimal cost is attained at the smallest of them:

$$C_{t,\alpha}(\hat{m}(t, \alpha)) = \min_i B_i(\alpha) .$$

Note that there is coincidence at α_0 due to the assumed multiplicity of minima:

$$B(\alpha_0) := B_i(\alpha_0) \quad , \quad i = 1, \dots, v . \tag{2.7}$$

We expand the functions B_i up to first order order

$$B_i(\alpha) = B(\alpha_0) + B_i'(\alpha_0)(\alpha - \alpha_0) + O(\alpha - \alpha_0) , \tag{2.8}$$

and observe that,

$$B'_i(\alpha_0) \neq B'_j(\alpha_0), \quad i \neq j. \quad (2.9)$$

The latter is due to the strict monotonicity of $\frac{\partial C_{t,\alpha}}{\partial \alpha}$ and the fact that each $\bar{m}_i(t, \alpha)$ is a critical points of the function $C_{t,\alpha}(\cdot)$. From (2.2)–(2.9) we conclude that for α in a possibly smaller neighbourhood $\mathcal{N}_{\alpha_0} \subseteq \tilde{\mathcal{N}}_{\alpha_0}$ there is a unique global minimum, and that property 3. holds with

$$a = \arg \min_i B'_i(\alpha_0), \quad b = \arg \max_i B'_i(\alpha_0),$$

and

$$\hat{m}_A(t, \alpha_0) := \hat{m}_a(t, \alpha_0), \quad \hat{m}_B(t, \alpha_0) := \hat{m}_b(t, \alpha_0).$$

□

We are ready to prove Theorem 1.6.

Proof. Suppose that C_{t,α_0} has a unique minimizer, denoted by $\hat{m}(t, \alpha_0)$ and let \mathcal{N}_{α_0} be the neighbourhood of the previous lemma. Then (1.21) holds for every $\alpha \in \mathcal{N}_{\alpha_0}$, and the continuity of $m \mapsto \Gamma_t(z, m)$ for every t, z gives the desired continuity of $\alpha \mapsto \gamma_t(\cdot | \alpha)$ at $\alpha = \alpha_0$.

To prove necessity, assume that C_{t,α_0} has multiple global minima. Consider \hat{m}_A and \hat{m}_B as in the previous lemma. Then, we have that there exist sequences $\alpha_n^- < \alpha_0 < \alpha_n^+$ converging to α_0 and such that $\gamma_t(\cdot | \alpha_n^\pm) = \Gamma_t(\cdot, \hat{m}(t, \alpha_n^\pm))$ and

$$\lim_{n \rightarrow \infty} \hat{m}(t, \alpha_n^-) = \hat{m}_B(t, \alpha_0) \neq \hat{m}_A(t, \alpha_0) = \lim_{n \rightarrow \infty} \hat{m}(t, \alpha_n^+). \quad (2.10)$$

Again using continuity of Γ_t with respect to m , we get

$$\lim_{n \rightarrow \infty} \gamma_t(z | \alpha_n^-) = \Gamma_t(z, \hat{m}_B(t, \alpha_0)) \neq \Gamma_t(z, \hat{m}_A(t, \alpha_0)) = \lim_{n \rightarrow \infty} \gamma_t(z | \alpha_n^+).$$

Hence α_0 is a bad magnetization. □

2.3 Proof of Theorem 1.7

To determine which solutions of (1.22) are global minima of $C_{t,0}$ when $h = 0$, we will pursue the following strategy. Using (1.22) we can write t as a function of m :

$$t(m) := \frac{1}{2} \operatorname{arccoth} \left(\frac{a^J(m)}{m} \right). \quad (2.11)$$

This allows us to determine for which time t the magnetization m can be a *possible minimum* (i.e., a solution of (1.22)).

Lemma 2.2. *Let $A \subseteq [-1, +1]$ be the set of m -values such that m is the solution of (1.22) for some $t > 0$, i.e., $A = \{m \in [-1, +1] : a^J(m)/m > 1\}$. Then, for every $m \in A$,*

$$C_{t(m),0}(m) = \frac{1}{2} J m^2 + \frac{1}{2} \log [1 - m \tanh(Jm)] =: C_M(m). \quad (2.12)$$

In words, (2.12) is the cost for m at the time at which it is a possible minimum.

Proof. Insert (2.11) into (1.25) and use (2.2). □

We now start the proof of Theorem 1.7.

Proof. First note that $l_{t,0}$ is linear with slope $\coth(2t) \in (1, \infty)$ and $l_{t,0}(0) = 0$, and that a^J is antisymmetric. Hence, if m is a solution of (1.22), then also $-m$ is a solution. Further note that

$$\begin{aligned} a^{J'}(m) &= (2J - 1) \cosh(2Jm) - 2Jm \sinh(2Jm), \\ a^{J''}(m) &= 4J(J - 1) \sinh(2Jm) - 4J^2 m \cosh(2Jm). \end{aligned} \quad (2.13)$$

(i) If $0 < J < \frac{1}{2}$, then $a^{J'}(m) < 0$ for all m , and hence $m = 0$ is the unique solution for all $t > 0$. If $\frac{1}{2} \leq J \leq 1$, then $a^{J''}(m) < 0$ for all m , hence a^J is convex, and so it suffices to compare slopes at 0: $k^{J,0'}(0) = a^{J'}(0) = 2J - 1 < 1$ and $l_{t,0}'(0) = \coth(2t) > 1$. Again, $m = 0$ is the unique solution for all $t > 0$ (see Fig. 6).

(ii) As before, $a^{J''}(m) < 0$ for all m , but now the slopes at 0 can be equal, which occurs when $t = \Psi_c$ with Ψ_c defined in (1.36). This proves that $\Delta_t = \{0\}$ for $0 < t \leq \Psi_c$ and $\Delta_t \subseteq \{-\hat{m}(t), 0, \hat{m}(t)\} =$ the set of solutions of (1.22) for $t > \Psi_c$. It is easily seen from Fig. 7(a) that $\hat{m}(t)$ is continuous and strictly increasing on $[\Psi_c, \infty)$ and $\hat{m}(\Psi_c) = 0$. It remains to show that $\{-\hat{m}(t), \hat{m}(t)\}$ are the global minima for all $t > \Psi_c$. This follows from the strategy behind the proof of Lemma 2.2. Since $C_{t,0}(0) = 0$ for all $t > 0$, it suffices to prove that $m \mapsto C_M(m)$ is strictly decreasing. From (2.12) we have

$$C'_M(m) = \frac{\partial C_{t(m),0}}{\partial m}(m) + \frac{\partial C_{t(m),0}}{\partial t}(m) t'(m). \quad (2.14)$$

The first term is zero by the definition of $t(m)$ (each m is a stationary point of $C_{t,0}$ at time $t = t(m)$). The second term is < 0 because $t'(m) > 0$ and

$$\begin{aligned} \frac{\partial C_{t,0}}{\partial t}(m) &= L(\hat{\varphi}_{t,0}^m(t), \dot{\hat{\varphi}}_{t,0}^m(t)) \\ &\quad + \int_0^t \left[\frac{\partial L}{\partial m}(\hat{\varphi}_{t,0}^m(s), \dot{\hat{\varphi}}_{t,0}^m(s)) \frac{\partial \hat{\varphi}_{t,0}^m}{\partial t}(s) + \frac{\partial L}{\partial \dot{m}}(\hat{\varphi}_{t,0}^m(s), \dot{\hat{\varphi}}_{t,0}^m(s)) \frac{\partial \dot{\hat{\varphi}}_{t,0}^m}{\partial t}(s) \right] ds \\ &= L(0, \dot{\hat{\varphi}}_{t,0}^m(t)) + \int_0^t \frac{\partial}{\partial s} \left\{ \frac{\partial L}{\partial \dot{m}}(\hat{\varphi}_{t,0}^m(s), \dot{\hat{\varphi}}_{t,0}^m(s)) \frac{\partial \hat{\varphi}_{t,0}^m}{\partial t}(s) \right\} ds \\ &= L(0, r) - \frac{\partial L}{\partial \dot{m}}(0, r) r \\ &= -\frac{1}{2} \sqrt{4 + r^2} + 1 < 0 \end{aligned} \quad (2.15)$$

with $r = \dot{\hat{\varphi}}_{t,0}^m(t)$, where the second equality uses (2.2). Since $C_M(0) = 0$, this yields the claim (see Figs. 7(a),7(b),7(c)).

(iii) This case is more difficult, because a^J no longer is convex on $(0, 1)$. Let Ψ_c^1 be the first time at which a solution different from 0 exists. To identify Ψ_c^1 , let

$$T_m(x) := (x - m)a^{J'}(m) + a^J(m), \quad (2.16)$$

and let m_1 be the solution of the equation $T_{m_1}(0) = 0 = -m_1 a^{J'}(m_1) + a^J(m_1)$, i.e.,

$$m_1 = \frac{a^J(m_1)}{a^{J'}(m_1)}. \quad (2.17)$$

From m_1 we get t_c^1 by using (2.11): $t_c^1 = t(m_1)$. As before, a solution of (1.22) for $t \geq t_c^1$ is not necessarily a minimum. To find out when it is, we follow the same strategy as in case (ii). Again, $C_{t,0}(0) = 0$ for all $t > 0$, and hence we must look for $m_* > 0$ such that

$$C_M(m_*) = 0. \quad (2.18)$$

Knowing m_* , we are able to compute Ψ_c using (2.11),

$$t_* = t(m_*). \quad (2.19)$$

In words, t_* is the first time at which 0 no longer is a minimum. As in case (ii), it suffices to prove that $m \mapsto C_M(m)$ is strictly decreasing on (m_*, ∞) . Again, we have (2.14). Since

$$t'(m) = \frac{1}{2}(\operatorname{arctanh})' \left(\frac{a^J(m)}{m} \right) \left\{ a^{J'}(m) - \frac{a^J(m)}{m} \right\} \frac{1}{m}, \quad (2.20)$$

it follows that $t'(m) = 0$ if and only if $m = a^J(m)/a^{J'}(m)$, which is the same condition as (2.17). This gives us a graphical argument to conclude that $t'(m) < 0$ for $0 < m < m_1$ and $t'(m) > 0$ for $m > m_1$ (see Figs. 7(d),7(e),7(f)). On the other hand, $m_* > m_1$. \square

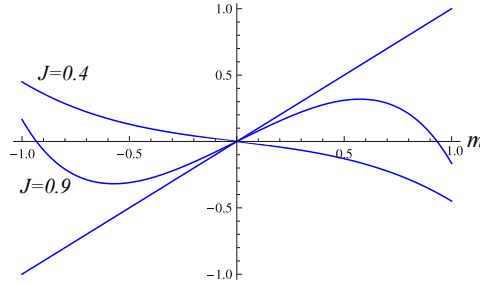


Figure 6: $m \mapsto a^J(m)$, $m \mapsto l_{t,0}(m)$ for Regime (i).

2.4 Proof of Theorem 1.8

Proof. From (1.28) it follows that

$$\tilde{m} > m, \tilde{t} > t \quad \implies \quad \hat{\varphi}_{\tilde{t},0}^{\tilde{m}}(s) > \hat{\varphi}_{t,0}^m(s) \quad \forall 0 \leq s \leq t. \quad (2.21)$$

(i) First note that, because $C_M(m_*) = 0$,

$$\frac{\partial m_*}{\partial J} = - \frac{\frac{\partial C_M}{\partial J}(m_*)}{\frac{\partial C_M}{\partial m}(m_*)}. \quad (2.22)$$

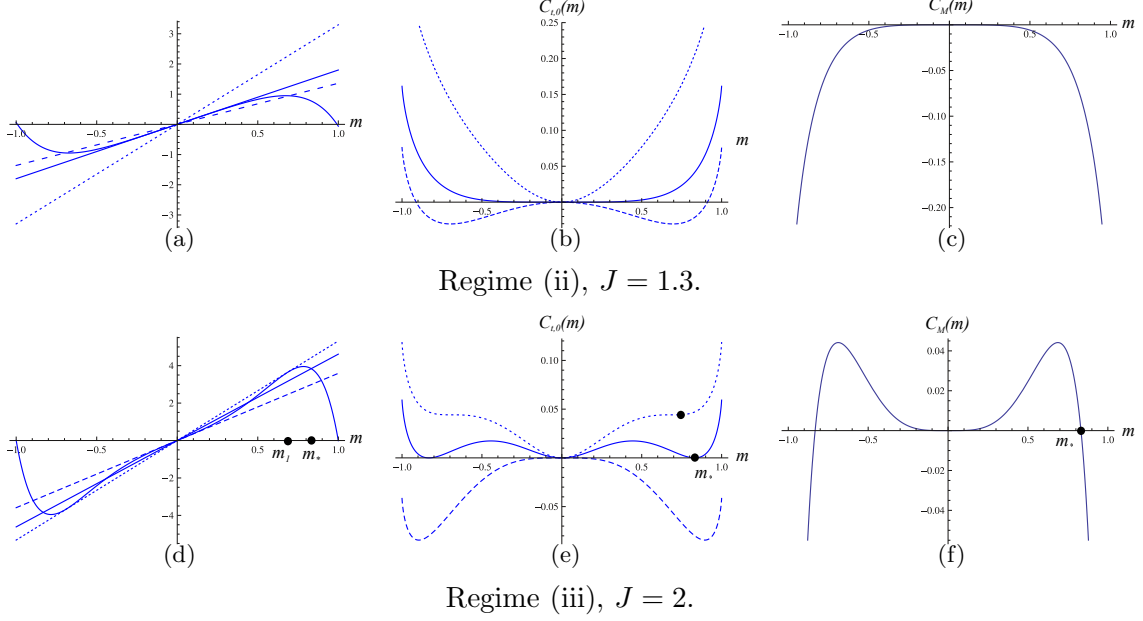


Figure 7: (a+d) $m \mapsto a^J(m)$, $m \mapsto l_{t,0}(m)$; (b+e) $m \mapsto C_{t,0}(m)$, for $0 \leq t < \Psi_c$ (dotted), $t = \Psi_c$ (drawn), $t > \Psi_c$ (dashed); (c+f) $m \mapsto C_M(m)$.

As in Section 2.3, case (iii), we have $(\partial C_M / \partial m)(m_*) < 0$. But

$$\frac{\partial C_M}{\partial J}(m_*) > 0 \iff m_* < \tanh(Jm_*), \quad (2.23)$$

which yields the claim because $m_* < m^\infty$.

(ii) The claim is straightforward for $1 < J \leq \frac{3}{2}$. For $J > \frac{3}{2}$ we need to prove that the function $J \rightarrow a^J(m_*)/m_*$ is strictly increasing. In fact,

$$\begin{aligned} \frac{\partial}{\partial J} \left[\frac{a^J(m_*)}{m_*} \right] &= \frac{1}{m_*^2} \left[-\frac{\partial m_*}{\partial J} \left\{ a^J(m_*) - m_* \frac{\partial a^J}{\partial m}(m_*) \right\} + m_* \frac{\partial a^J}{\partial J}(m_*) \right] \\ &= 1 + \cosh(2Jm_*) - m_* \sinh(2Jm_*). \end{aligned} \quad (2.24)$$

The strict positivity of the last expression is equivalent to the inequality $m_* < \coth(Jm_*)$, which is satisfied for $J > 1$.

(iii) We have $a^J(m) \downarrow a^{\frac{3}{2}}(m)$ as $J \downarrow \frac{3}{2}$ for all $m \in (0, 1)$, with a^J and $a^{\frac{3}{2}}$ continuous. By Dini's theorem, the convergence is uniform.

(iv) Since $\Psi_c(\tilde{J}) < \Psi_c(\frac{3}{2})$, the same argument as in case (iii) can be used.

(v) and (vi) are consequences of parts (ii) and (iii) of Theorem 1.7. \square

3 Proof of Theorem 1.9

In Sections 3.1 and 3.2 we prove that overshoots, respectively, bifurcations, take place in regime (1), respectively (2)–(3), of Theorem 1.9. The analysis of the former regime does not distinguish on whether the initial field h is zero or not.

3.1 Regime (1): Overshoots

The trick is again to write t as a function of m . From (1.22), we have

$$-k^{J,h}(m) + 2m \coth(2t) = m \coth(2t) + \alpha \operatorname{csch}(2t). \quad (3.1)$$

Hence, from (1.23),

$$\begin{aligned} k^{J,h}(m)[-k^{J,h}(m) + 2m \coth(2t)] \\ = [m \coth(2t) - \alpha \operatorname{csch}(2t)] [m \coth(2t) + \alpha \operatorname{csch}(2t)] \end{aligned} \quad (3.2)$$

which implies

$$-k^{J,h}(m)^2 - \alpha^2 = (m^2 - \alpha^2) \coth^2(2t) - 2mk^{J,h}(m) \coth(2t). \quad (3.3)$$

Solving for t we find

$$t_F(m) := \begin{cases} \frac{1}{2} \operatorname{arccoth}\left(\frac{mk^{J,h}(m) + |\alpha| \sqrt{\Phi(m)}}{m^2 - \alpha^2}\right) & \text{if } m \neq \pm\alpha \\ \frac{1}{2} \operatorname{arccoth}\left(\frac{[k^{J,h}(m)]^2 + m^2}{2m k^{J,h}(m)}\right) & \text{if } m = \pm\alpha, m k^{J,h}(m) < 0 \\ 0 & \text{if } m = \pm\alpha, m k^{J,h}(m) > 0 \end{cases} \quad (3.4)$$

$$t_L(m) := \begin{cases} \frac{1}{2} \operatorname{arccoth}\left(\frac{mk^{J,h}(m) - |\alpha| \sqrt{\Phi(m)}}{m^2 - \alpha^2}\right) & \text{if } m \neq \pm\alpha \\ \frac{1}{2} \operatorname{arccoth}\left(\frac{[k^{J,h}(m)]^2 + m^2}{2m k^{J,h}(m)}\right) & \text{if } m = \pm\alpha, m k^{J,h}(m) > 0 \\ 0 & \text{if } m = \pm\alpha, m k^{J,h}(m) < 0 \end{cases} \quad (3.5)$$

with

$$\Phi(m) := k^{J,h}(m)^2 - m^2 + \alpha^2. \quad (3.6)$$

These are times at which m is a stationary point (not necessarily a minimum) both for $C_{t,\alpha}(m)$ and for $C_{t,-\alpha}(m)$ [Equation (3.2), and hence the solutions (3.4) are insensitive to the sign of α]. Overshoots and undershoots occur for values of m satisfying (i) $t_F(m) > 0$ and $t_L(m) > 0$ and (ii) at these times m is a minimum.

The m -dependence of t_F and t_L is depicted in Figs 8 and 9 for cases in which both functions are injective, i.e., for α for which there is only one critical point for each t . In more complicated cases, for instance, when overshoot and bifurcation occur simultaneously, there are two or more stationary points only one of which is a minimum.

We divide the analysis in four steps.

Step A: *Existence of z^+ , z^- and z^{-+} .* We observe that there exists a unique $m^\infty > 0$ such that $k^{J,h}(m^\infty) = m^\infty$. Furthermore,

$$k^{J,h}(m) = 0 \quad \iff \quad \tanh(2h) = \frac{-a^J(m)}{b^J(m)} =: A(m). \quad (3.7)$$

This function A has the features depicted in Figure 10: it is odd, satisfies $A(0) = 0$, $A(1) = 1$ and it is convex with only one global minimum between 0 and 1. We conclude that there

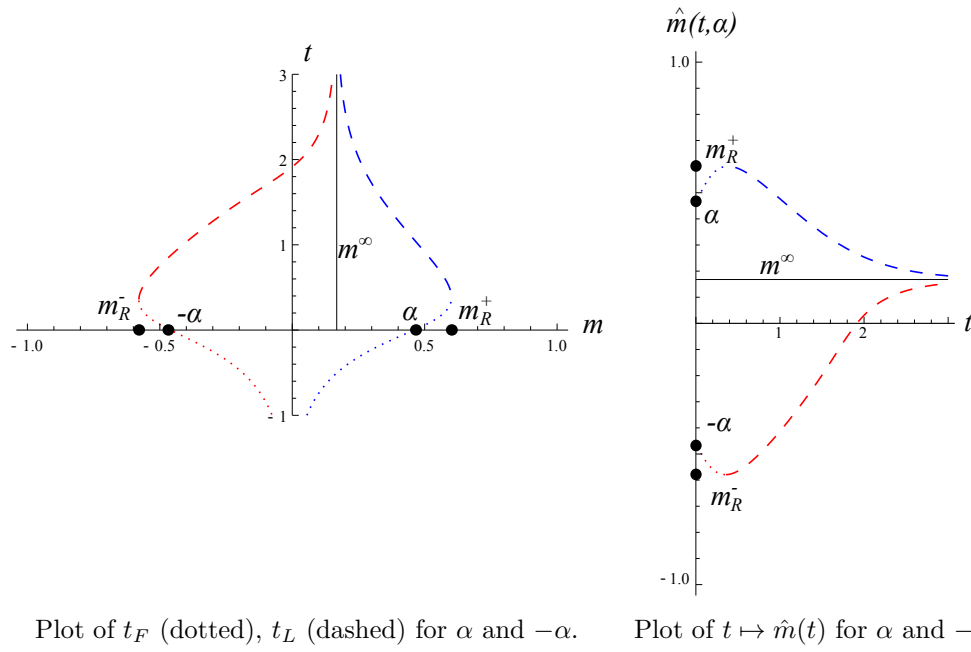


Figure 8: Overshoot for (J, h, α) in Regime (1a) or for $(J, h, -\alpha)$ in Regime (1d). Parameters: $(J, h) = (0.95, 0.01)$, $\alpha = 0.46$.

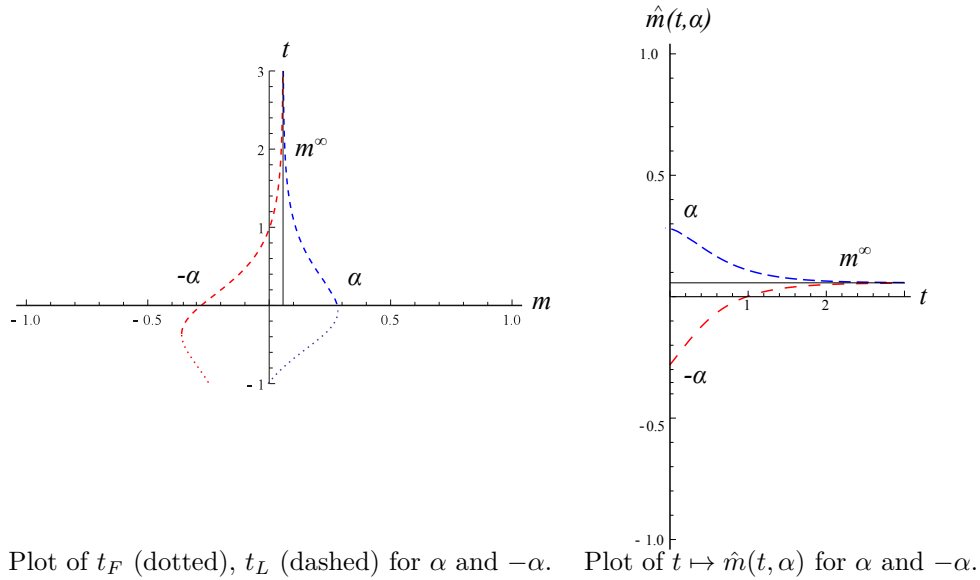


Figure 9: Absence of overshoot for (J, h, α) in Regime (1b) or for $(J, h, -\alpha)$ in Regime (1c). Parameters: $(J, h) = (0.3, 0.04)$, $\alpha = 0.28$.

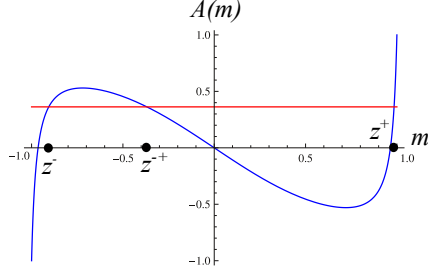


Figure 10: Plot of A and intersection with the constant $\tanh(2h)$.

exists a unique $z^+ = z^+(J, h) > 0$ such that $k^{J,h}(z^+) = 0$ and, in addition, if $\tanh(2h) < \max_{m \in [-1,0]} A(m)$, then there exist $-1 < z^- < z^{-+} < 0$ such that $k^{J,h}(z^-) = k^{J,h}(z^{-+}) = 0$ (see Fig. 10).

Step B: *Existence of m_R^+ , m_R^- and relation between m_∞ and α .*

(Ba) *Existence of m_R^+ :* $k^{J,h}(\alpha) > 0$ together with $\alpha > 0$ and Step A imply $\alpha < z^+$. Since $\Phi(\alpha) > 0$ and $\Phi(z^+) < 0$, it follows that there exists a m_R^+ such that $0 < \alpha < m_R^+ < z^+$ and

$$\Phi(m_R^+) = 0. \quad (3.8)$$

The latter in turn implies that $k^{J,h}(m_R^+)^2 = m_R^{+2} - \alpha^2$. This, together with $k^{J,h}(m_R^+) > 0$, implies that $k^{J,h}(m_R^+) < m_R^+$, which leads to $m_\infty < m_R^+$.

(Bd) *Existence of m_R^- :* As in (Ba), $\Phi(z^-) < 0$ and $\Phi(\alpha) > 0$ imply that there exists a m_R^- such that $z^- < m_R^- < \alpha < 0$ and

$$\Phi(m_R^-) = 0. \quad (3.9)$$

(Bc): $k^{J,h}(\alpha) > 0$ and $\alpha < 0$ imply $\alpha < m_\infty$. This follows from the fact that $k^{J,h}(\alpha) > \alpha$ implies $\alpha < m_\infty$ by (1.33).

(Bb): $k^{J,h}(\alpha) < 0$ and $\alpha > 0$ imply $\alpha > m_\infty$. Again, this is a consequence of (1.33).

Step C: *Consequence of the positivity of times.* Only positive solutions of Equation (3.4) are of interest. This implies the constraints

$$t_F(m) > 0 \iff \eta_F(m) := m k^{J,h}(m) + (\alpha^2 - m^2) + |\alpha| \sqrt{\Phi(m)} \begin{cases} > 0 & \text{if } m^2 > \alpha^2, \\ < 0 & \text{if } m^2 < \alpha^2, \end{cases} \quad (3.10)$$

and

$$t_L(m) > 0 \iff \eta_L(m) := m k^{J,h}(m) + (\alpha^2 - m^2) - |\alpha| \sqrt{\Phi(m)} \begin{cases} > 0 & \text{if } m^2 > \alpha^2, \\ < 0 & \text{if } m^2 < \alpha^2. \end{cases} \quad (3.11)$$

The functions η_F and η_L satisfy

$$\eta_F(\alpha) = \alpha k^{J,h}(\alpha) + |\alpha| |k^{J,h}(\alpha)| = \begin{cases} > 0 & \text{if } \alpha k^{J,h}(\alpha) > 0, \\ = 0 & \text{if } \alpha k^{J,h}(\alpha) < 0, \end{cases} \quad (3.12)$$

$$\eta_L(\alpha) = \alpha k^{J,h}(\alpha) - |\alpha| |k^{J,h}(\alpha)| = \begin{cases} = 0 & \text{if } \alpha k^{J,h}(\alpha) > 0, \\ < 0 & \text{if } \alpha k^{J,h}(\alpha) < 0. \end{cases} \quad (3.13)$$

Also, from (1.33),

$$\begin{aligned} \eta_F(m^\infty) &= 2|\alpha|^2, & \eta'_F(m^\infty) &= 2m^\infty k^{J,h'}(m^\infty), \\ \eta_L(m^\infty) &= 0, & \eta'_L(m^\infty) &= 0. \end{aligned} \quad (3.14)$$

Last line implies that m^∞ is a root of η_L but there is no change of sign around it. Finally, from expressions (3.10)-(3.14) we conclude that:

- The zeros of η_F , η_L are a subset of $\{m^\infty, \pm\alpha\}$.
- The intervals in which η_F and η_L satisfies the constrains (3.10)-(3.11) are:

Regime	(1a)	(1b)	(1c)	(1d)
Condition	$\alpha > 0, k^{J,h}(\alpha) > 0$	$\alpha > 0, k^{J,h}(\alpha) < 0$	$\alpha < 0, k^{J,h}(\alpha) > 0$	$\alpha < 0, k^{J,h}(\alpha) < 0$
$t_F > 0$	$[\alpha, m_R^+]$	\emptyset	\emptyset	$[m_R^-, \alpha]$
$t_L > 0$	$[m^\infty, m_R^+]$	$[m^\infty, \alpha]$	$[\alpha, m^\infty]$	$[m_R^-, m^\infty]$

We observe that in regime (1a) each value of $m \in [\alpha \wedge m_\infty, m_R^+]$ is attained at two different times t_F (“First”) and t_L (“Last”). The same happens in regime (1d) for $m \in [m_R^-, \alpha \vee m_\infty]$. These phenomena correspond respectively, to an over and an under shoot. The proof is completed by showing that the trajectories have the right monotonicity properties.

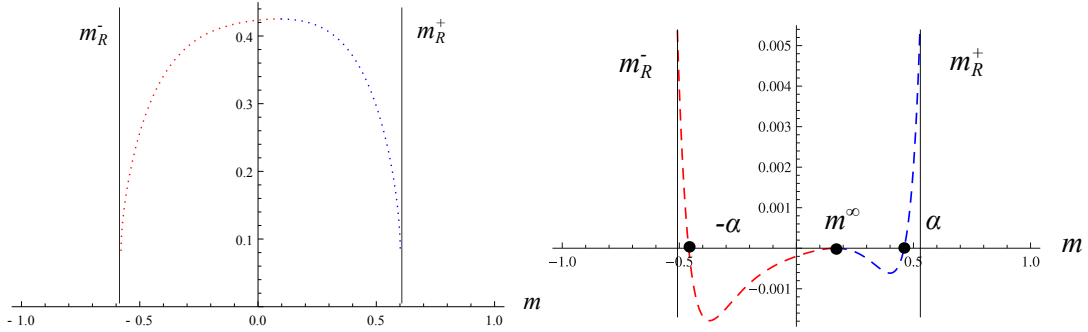


Figure 11: η_F (dotted) and η_L (dashed) for (J, h, α) in regime 1(a) and for $(J, h, -\alpha)$ in regime 1(d). Parameters: $(J, h) = (0.95, 0.01)$, $\alpha = 0.46$ as in Fig. 8.

Step D: Monotonicity.

By using implicit derivation we get

$$\frac{\partial \hat{m}}{\partial t}(t) = \frac{\frac{\partial l_{t,\alpha}}{\partial t}(\hat{m})}{[k^{J,h}]'(\hat{m}) - l'_{t,\alpha}(\hat{m})} = \frac{2 \operatorname{csch}(2t) \left\{ \hat{m} \operatorname{csch}(2t) - \alpha \operatorname{coth}(2t) \right\}}{[k^{J,h}]'(\hat{m}) - \operatorname{coth}(2t)}. \quad (3.15)$$

On the other hand, if $\hat{m} = \hat{m}(t)$ is a critical point $\left(k^{J,h}(\hat{m}) = l_{t,\alpha}(\hat{m}) \right)$, then $\Phi(\hat{m}) = (\hat{m} \operatorname{csch}(2t) - \alpha \operatorname{coth}(2t))^2$. Hence,

$$\frac{\partial \hat{m}}{\partial t}(t) = 0 \iff \Phi(\hat{m}(t)) = 0.$$

Splitting in cases according to different values of \hat{m} , it is not hard to conclude from (3.8) and (3.15) the following monotonicity properties of the trajectories.

Regime	(1a)	(1b)	(1c)	(1d)
\hat{m} incr.	$0 < t < t(m_R^+)$	\emptyset	$0 < t < \infty$	$t(m_R^-) < t < \infty$
\hat{m} decr.	$t(m_R^+) < t < \infty$	$0 < t < \infty$	\emptyset	$0 < t < t(m_R^-)$

This concludes the proof of part 1 of Theorem 1.9.

3.2 Bifurcation

Bifurcation proofs rely on the following facts.

- (B1)** For short times there is a unique critical point, close to α .
- (B2)** Therefore, in order for bifurcation to occur a local maximum and a local minimum must appear in the course of time. Given condition (1.22), usual arguments imply that two (or more) stationary points appear at times larger than \tilde{t} if the curves $l_{\tilde{t},\alpha}$ and $k^{J,h}$ become tangent at a certain magnetization \tilde{m} . The pairs (\tilde{m}, \tilde{t}) are determined by the following two equations (a similar argument was used in the proof of Theorem 1.7(iii)):

$$\begin{aligned} [k^{J,h}]'(\tilde{m}) &= \coth(2\tilde{t}), \\ k^{J,h}(\tilde{m}) &= \tilde{m} \coth(2\tilde{t}) - \alpha \operatorname{csch}(2\tilde{t}). \end{aligned} \tag{3.16}$$

Inserting the first equation into the second, we get

$$F(\tilde{m}) := \frac{\tilde{m} [k^{J,h}]'(\tilde{m}) - k^{J,h}(\tilde{m})}{\operatorname{csch}[\operatorname{arccoth}([k^{J,h}]'(\tilde{m}))]} = \alpha. \tag{3.17}$$

We are left with the task of determining whether or not this equation has solutions. Note that

$$F'(m) = 0 \iff [k^{J,h}]''(m) = 0 \quad \text{or} \quad m = k^{J,h}(m) [k^{J,h}]'(m). \tag{3.18}$$

In what follows all the assertions about F can be checked by using the equivalence in (3.18) and doing a straightforward analysis of $k^{J,h}$.

- (B3)** $t \mapsto C_{t,\alpha}$ is continuous with respect to $\|\cdot\|_\infty$. Hence, when a new minimum appears it cannot be a global one. Both $\alpha \mapsto C_{t,\alpha}$ and $h \mapsto C_{t,\alpha}$ are also continuous.
- (B4)** When $t \rightarrow \infty$ we have two global minima $\pm m^\infty$ if $h = 0$. If $h > 0$ the symmetry is broken and there is only one global minimum $m^\infty > 0$.

Whenever a local maximum/local minimum appears (disappears), we will refer to this behaviour as LMLMA (LMLMD). We proceed by looking at $h = 0$ and $h \neq 0$ separately.

Part (2) ($h = 0$, see Fig. 12). The scenario for $\alpha = 0$ has already been proven in Theorem 1.7. We concentrate on $\alpha > 0$; this is no loss of generality due to the antisymmetry of F .

Claim: Whenever $\alpha > 0$, negative solutions of (3.17) can not cause bifurcations. In fact, let t_{nc} be the time at which the critical point in the negative side emerge. Let also,

$$d(t) := C_{t,\alpha}(\overline{m}_-(t, \alpha)) - C_{t,\alpha}(\hat{m}(t, \alpha)), \quad t \geq t_{nc},$$

where $\overline{m}_-(t, \alpha)$ is the negative *local* [because of **(B3)**] minimum and $\hat{m}(t, \alpha)$ is the global minimum of $C_{t,\alpha}$. The last one is positive due to **(B1)**. and the supposition of $\alpha > 0$. By definition, $d(t_{nc}) > 0$ and by **(B4)** $\lim_{t \rightarrow \infty} d(t) = 0$. Doing calculations similar to (2.15) and using that $\overline{m}_-(t, \alpha)$, $\hat{m}(t, \alpha)$ are both critical points, we get that $d'(t) < 0 \forall t > t_{nc}$. This proves the claim.

In what follows we focus in equation (3.17). Owing to the previous claim, in order to bifurcation to occur a positive solution of (3.17) is needed.

(2a-b) If $0 < J \leq \frac{3}{2}$, then $F'(m) < 0$ for all $m \in (-1, +1)$. Hence, for all $\alpha > 0$ there is only one solution of (3.17). This solution turns out to be negative and hence it can not correspond to a bifurcation.

(2c) If $J > \frac{3}{2}$, F has only one global maximum on the positive side, with value $U_B = U_B(J) > 0$. Combining **B1.-B4.**, we get that there is bifurcation if and only if $\alpha \in [0, U_B] = \text{Im}(F|_{[0,1]})$.

Part (3) ($h > 0$, see Fig. 13).

Remark: If $h > 0$, then **B4.** (Symmetry breaking) allows the appearance of solutions of (3.17) leading to bifurcations.

Once more, let study the different scenarios for F when $h > 0$.

(3a) If $0 < J \leq 1$, then $\text{Im}(F) \subseteq [-1, +1]^c$. Therefore (3.17) has no solution for any $|\alpha| \leq 1$ and there is no bifurcation.

(3b) If $1 < J \leq \frac{3}{2}$, then F has a unique maximum for $m \in [0, 1]$ with value $U_B = U_B(J, h) < 0$, and $[-1, U_B] = \text{Im}(F|_{[0,1]})$. Arguing as in the claim of Part 2 for $\alpha > U_B$, we conclude that there is bifurcation if and only if $\alpha \in [-1, U_B]$.

(3c) Assume $J > \frac{3}{2}$.

1. For $h > 0$ small enough the behaviour is “close” to the $h = 0$ case due to the continuity of $h \mapsto C_{t,\alpha}$ with respect to the infinite norm.

Indeed, there exists $L_B := \min_{[-1,0]} F \approx -U_B(J, 0)$ and $U_B := \max_{[0,1]} F \approx U_B(J, 0)$, with $(-1, U_B] = \text{Im}(F|_{[0,1]})$. There are different regimes for α :

(a) For $\alpha < L_B$, there is a unique solution of (3.17), which is in the positive side, leading to a bifurcation [because of **(B4)**].

(b) For $0 > \alpha \gtrsim L_B$ there are both a negative and a positive solution to (3.17). Both lead to bifurcations, the negative one by continuity **B3.** and the positive due to **B3.-B4.** The negative solution appears earlier in time. We write $s_B = s_B(\alpha)$ for the time of the first (negative side) bifurcation and $t_B = t_B(\alpha)$ for the second bifurcation time.

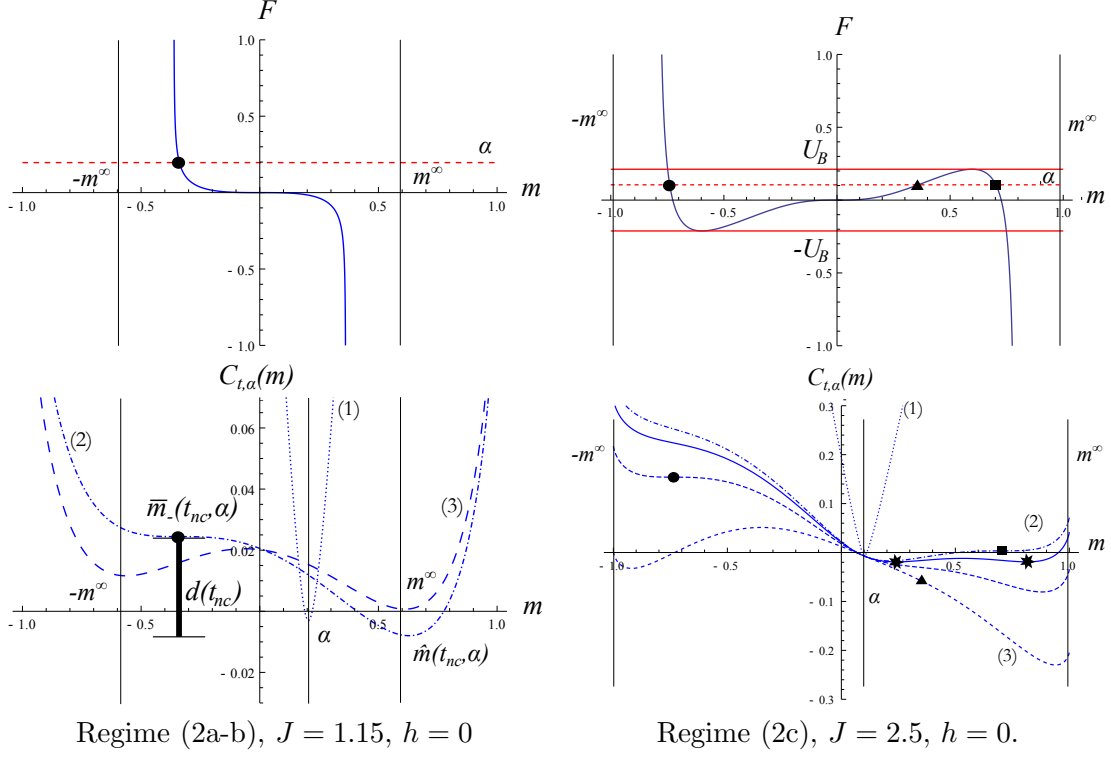


Figure 12: First row: Plot of $m \mapsto F(m)$. \bullet = LMLMA not leading to bifurcation; \blacktriangle = LMLMD; \blacksquare = LMLMA leading to bifurcation; \star = bifurcation. Second row: Plot of $m \mapsto C_{t,\alpha}(m)$ for different times. Short time (dotted) bifurcation (solid), long time (dashed).

- (c) For $0 < \alpha \lesssim U_B$ there is LMLMA on the positive and negative sides. As in the $h = 0$ case, the negative one does not lead to a bifurcation, and thus only one bifurcation occurs, which happens to be in the positive side.
- (d) By the continuity property **(B3)** and the monotonicity of $\alpha \mapsto s_B(\alpha)$ (proved below), the two previous regimes coalesce, leading to an intermediate value $M_T \in (L_B, U_B)$ such that trifurcation occurs at $\alpha = M_T$.
- (e) For $\alpha > U_B$ there is no positive solution to (3.17) with $\alpha \in [0, 1]$. Hence, no bifurcation occurs.

2. The limit $h \rightarrow \infty$ in (3.17) yields

$$\frac{m \left(a^{J'}(m) + b^{J'}(m) \right) - \left(a^J(m) + b^J(m) \right)}{\left(a^{J'}(m) + b^{J'}(m) \right)} = \alpha.$$

Hence, for $h > 0$ large enough we get a behaviour similar to (3b), but with $U_B(J, h) > 0$.

3. The existence of h^* follows from the continuity of the function $h \mapsto C_{t,\alpha}$ commented in **(B3)** with respect to h .

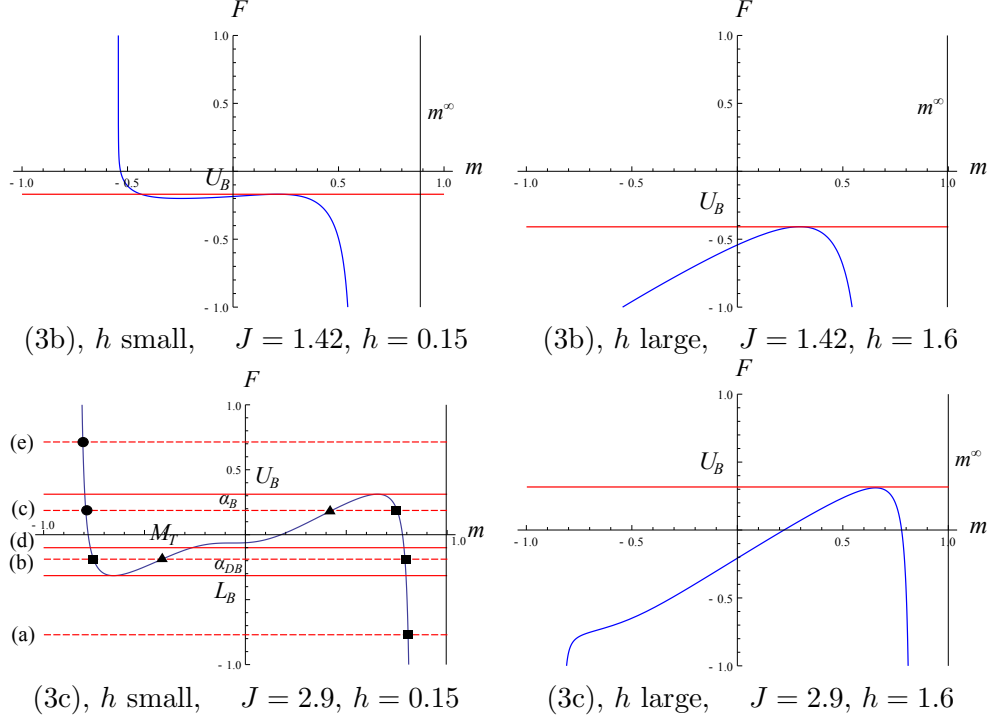


Figure 13: Plot of $m \mapsto F(m)$ for different regimes of J when $h > 0$. \bullet = LMLMA not causing a bifurcation; \blacktriangle = LMLMD; \blacksquare = LMLMA causing a bifurcation.

3.2.1 Monotonicity of the functions $t_B(\alpha)$ and $s_B(\alpha)$

The bifurcation times are characterized by the following equations:

$$\begin{aligned}
 k^{J,h}(\hat{m}_1) &= l_{t_B,\alpha}(\hat{m}_1), \\
 k^{J,h}(\hat{m}_2) &= l_{t_B,\alpha}(\hat{m}_2), \\
 C_{t_B,\alpha}(\hat{m}_1) &= C_{t_B,\alpha}(\hat{m}_2).
 \end{aligned} \tag{3.19}$$

The first two equations say that \hat{m}_1 and \hat{m}_2 are stationary points at the same time t_B , while the third one establishes the equality of costs at this time t_B . Taking the derivative with respect to α of the third equation we get

$$\frac{\partial t_B}{\partial \alpha} = - \frac{\frac{\partial C_{t_B,\alpha}}{\partial \alpha}(\hat{m}_2) - \frac{\partial C_{t_B,\alpha}}{\partial \alpha}(\hat{m}_1)}{\frac{\partial C_{t_B,\alpha}}{\partial t_B}(\hat{m}_2) - \frac{\partial C_{t_B,\alpha}}{\partial t_B}(\hat{m}_1)}. \tag{3.20}$$

A straightforward computation using the first two equations shows that $\frac{\partial t_B}{\partial \alpha} < 0$, which implies that $\alpha \mapsto t_B(\alpha)$ is continuous and decreasing. A similar argument shows that $\alpha \mapsto s_B(\alpha)$ is continuous and increasing.

References

- [1] D. Dereudre and S. Roelly, Propagation of Gibbsianness for infinite-dimensional gradient Brownian diffusions, *J. Stat. Phys.* 121 (2005) 511–551.

- [2] A.C.D. van Enter, R. Fernández, F. den Hollander and F. Redig, Possible loss and recovery of Gibbsianness during the stochastic evolution of Gibbs measures, *Commun. Math. Phys.* 226 (2002) 101–130.
- [3] A.C.D. van Enter, R. Fernández, F. den Hollander and F. Redig, A large-deviation view on dynamical Gibbs-non-Gibbs transitions, *Moscow Math. J.* 10 (2010) 687–711.
- [4] A.C.D. van Enter, C. Külske, A.A. Opoku and W.M. Ruszel, Gibbs-non-Gibbs properties for n -vector lattice and mean-field models, *Braz. J. Prob. Stat.*, 24 (2010) 226–255.
- [5] A.C.D. van Enter and W.M. Ruszel, Loss and recovery of Gibbsianness for XY spins in a small external field, *J. Math. Phys.* 49 (2008) 125208.
- [6] A.C.D. van Enter and W.M. Ruszel, Gibbsianness versus non-Gibbsianness of time-evolved planar rotor models, *Stoch. Proc. Appl.* 119 (2010) 1866–1888.
- [7] V. Ermolaev and C. Külske, Low-temperature dynamics of the Curie-Weiss model: Periodic orbits, multiple histories, and loss of Gibbsianness, *J. Stat. Phys.* 141 (2010) 727–756.
- [8] F. den Hollander, *Large Deviations*, Fields Institute Monographs 14, American Mathematical Society, Providence, RI, 2000.
- [9] C. Külske and A. Le Ny, Spin-flip dynamics of the Curie-Weiss model: Loss of Gibbsianness with possibly broken symmetry, *Commun. Math. Phys.* 271 (2007) 431–454.
- [10] C. Külske and A.A. Opoku, The posterior metric and the goodness of Gibbsianness for transforms of Gibbs measures, *Elect. J. Prob.* 13 (2008) 1307–1344.
- [11] C. Külske and A.A. Opoku, Continuous mean-field models: limiting kernels and Gibbs properties of local transforms, *J. Math. Phys.* 49 (2008) 125215.
- [12] C. Külske and F. Redig, Loss without recovery of Gibbsianness during diffusion of continuous spins, *Probab. Theory Relat. Fields* 135 (2006) 428–456.
- [13] A. Le Ny and F. Redig, Short time conservation of Gibbsianness under local stochastic evolutions, *J. Stat. Phys.* 109 (2002) 1073–1090.
- [14] T.M. Liggett, *Interacting Particle Systems*, Grundlehren der Mathematischen Wissenschaften 276, Springer, New York, 1985.
- [15] A.A. Opoku, *On Gibbs Properties of Transforms of Lattice and Mean-Field Systems*, PhD thesis, Groningen University, 2009.
- [16] F. Redig, S. Roelly and W.M. Ruszel, Short-time Gibbsianness for infinite-dimensional diffusions with space-time interaction, *J. Stat. Phys.* 138 (2010) 1124–1144.