

# A Spectroscopic Follow-up Program of Very Massive Galaxies at 3 < z < 4: Confirmation of Spectroscopic Redshifts, and a High Fraction of Powerful AGNs

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Received 2016 June 3; revised 2017 April 13; accepted 2017 May 5; published 2017 June 8

#### Abstract

We present the analysis and results of a spectroscopic follow-up program of a mass-selected sample of six galaxies at 3 < z < 4 using data from Keck-NIRPSEC and VLT-Xshooter. We confirm the z > 3 redshifts for half of the sample through the detection of strong nebular emission lines, and improve the  $z_{phot}$  accuracy for the remainder of the sample through the combination of photometry and spectra. The modeling of the emission-line-corrected spectral energy distributions (SEDs) adopting improved redshifts confirms the very large stellar masses of the sample  $(M_* \sim 1.5-4 \times 10^{11} M_{\odot})$  in the first 2 Gyr of cosmic history, with a diverse range in stellar ages, starformation rates, and dust content. From the analysis of emission-line luminosities and widths, and far-infrared (FIR) fluxes, we confirm that  $\gtrsim 80\%$  of the sample are hosts to luminous hidden active galactic nuclei (AGNs), with bolometric luminosities of  $\sim 10^{44-46}$  erg s<sup>-1</sup>. We find that the MIPS 24  $\mu$ m photometry is largely contaminated by AGN continuum, rendering the SFRs derived using only 24  $\mu$ m photometry to be severely overestimated. By including the emission from the AGN in the modeling of the UV-to-FIR SEDs, we confirm that the presence of the AGN does not considerably bias the stellar masses (<0.3 dex at  $1\sigma$ ). We show evidence for a rapid increase of the AGN fraction from  $\sim$ 30% to  $\sim$ 60%–100% over the 1 Gyr between  $z \sim$  2 and  $z \sim$  3. Although we cannot exclude some enhancement of the AGN fraction for our sample due to selection effects, the small measured [O III] contamination to the observed K-band fluxes suggests that our sample is not significantly biased toward massive galaxies hosting AGNs.

*Key words:* cosmology: observations – galaxies: evolution – galaxies: formation – galaxies: high-redshift – galaxies: stellar content – infrared: galaxies

## 1. Introduction

One of the surprising findings in observational studies of galaxies in the early universe is that, in contrast to the bottomup assembly of dark matter halos inferred from simulations of ACDM cosmologies, in which low-mass halos form early and subsequently grow via continued accretion and merging to form more massive halos at later times, the stellar component of halos appear to have assembled in an anti-hierarchical, topdown manner, with low-mass galaxies assembling most of their mass and forming most of their stars at later times compared to more massive systems (Thomas et al. 2005; Gallazzi et al. 2006; Pérez-González et al. 2008; Wiklind et al. 2008; Marchesini et al. 2009, 2010; Domínguez Sánchez et al. 2011; Muzzin et al. 2013b; Stefanon et al. 2015).

Closely related to this issue is the intriguing finding that the number density of the most massive galaxies (i.e., those with stellar masses  $M_* > 3 \times 10^{11} M_{\odot}$ ) seems to evolve slowly from  $z \sim 4$  to  $z \sim 1.5$  (Pérez-González et al. 2008; Marchesini et al. 2009, 2010; Muzzin et al. 2013b), suggesting that very massive galaxies were already in place at  $z \sim 3.5$  and implying that their stellar content was assembled rapidly in the first ~1.5 Gyr of cosmic history (Caputi et al. 2015; Stefanon et al. 2015). While theoretical models of galaxy formation and evolution have significantly improved in matching observations in recent years, they are still unable to reproduce the observed number

density of galaxies at 3 < z < 4 with log  $(M_*/M_{\odot}) > 11.5$ (Fontanot et al. 2009; Marchesini et al. 2009, 2010; Ilbert et al. 2013), unless the model-predicted stellar mass functions (SMFs) are convolved to account for the effect of Eddington bias by ~0.3–0.4 dex (Henriques et al. 2015), arguably much larger than the typical random uncertainties on  $M_*$  at these redshifts (e.g., Marchesini et al. 2010; Ilbert et al. 2013; Muzzin et al. 2013b). Furthermore, it remains to be determined to what extent the inferred number density measurements at high *z* are affected by blending of sources in ground-based images.

In recent years, robust evidence for the existence of massive  $(M_* \gtrsim 10^{11} M_{\odot})$  galaxies at z > 3 was enabled by the wellsampled spectral energy distributions (SEDs) and accurate photometric redshifts delivered by near-infrared (NIR) imaging surveys adopting medium-bandwidth filters in the NIR, i.e., NEWFIRM Medium-Band Survey (NMBS; Whitaker et al. 2011, and the FourStar Galaxy Evolution Survey (zFOURGE; Spitler et al. 2014). The wide-field NMBS robustly revealed, a population of "monster" galaxies at 3 < z < 4 with log  $(M_*/M_{\odot}) > 11.4$  (Marchesini et al. 2010), whereas the deeper but pencil beam zFOURGE survey extended the study of massive galaxies at z > 3 down to the characteristic stellar mass at  $z \sim 3-4$ , i.e.,  $\log(M_*/M_{\odot}) \sim 10.6$  (Spitler et al. 2014; Straatman et al. 2014). Stefanon et al. (2015) used the

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UltraVISTA DR1 (Muzzin et al. 2013a) to search for very massive galaxies at 4 < z < 7, finding a robust candidate of a monster galaxy at  $z \sim 5.6$  with  $\log(M_*/M_{\odot}) \sim 11.6$  and quenched star-formation activity.

In contrast to the local universe where massive galaxies are predominantly quiescent and have old stellar populations with red colors (Blanton et al. 2003), the massive galaxy population at  $z \sim 3$  is dominated by dusty, star-forming galaxies (~40–60% of massive population, Marchesini et al. 2010; Spitler et al. 2014), and includes a significant population of already quiescent galaxies (~30%–40%) with ages consistent with the age of the universe at the targeted redshifts. In addition, from the X-ray and radio detections, as well as the ubiquitous very bright fluxes in MIPS 24  $\mu$ m (rest-frame 5–6  $\mu$ m), it appears that massive galaxies in the early universe commonly host active galactic nuclei (AGNs), both at z > 3 (Marchesini et al. 2010; Stefanon et al. 2015), as well as at z < 3 (Barro et al. 2013; Föster Scheriber et al. 2014).

Despite the well-sampled SEDs delivered by the NMBS and zFOURGE surveys, there is still ambiguity in the photometric redshift solutions for the rest-frame UV faint massive galaxies at z > 3. Quantitatively, among the population of very massive galaxies with  $\log(M_*/M_{\odot}) > 11.4$  at 3 < z < 4 selected from the NMBS by Marchesini et al. (2010), ~50% could have a photometric redshift  $z_{phot} < 3$  solution if their stellar populations were characterized by both an evolved (i.e., stellar age  $\gtrsim 1$  Gyr) and very dusty (i.e.,  $\langle A_V \rangle \sim 3$  mag) component. This elusive population of high-*z* old and dusty galaxies, originally proposed in Marchesini et al. (2010) and effectively non-existent at z < 1, appears to play a major role among high-*z* massive galaxies (Marchesini et al. 2014), and has only recently been targeted for detailed investigations (e.g., Bedregal et al. 2017, G. Brammer et al. 2017, in preparation).

Because of the ambiguity in the photometric redshift solutions for as much as half of the population, it is of paramount importance to spectroscopically measure the redshifts of very massive galaxies at z > 3. Recently, Marsan et al. (2015) presented the first spectroscopic confirmation of an ultra-massive ( $M_* \sim 3 \times 10^{11} M_{\odot}$ ) galaxy at z = 3.35, dubbed "The Vega Galaxy," due to the incredible similarity of the integrated SED to an A0V-star such as Vega. The detailed analysis of its UV-to-FIR SED reveals that most of its stars formed at z > 4 in a highly dissipative, intense, and short burst of star formation, and that it is transitioning to a post-starburst phase while hosting a powerful AGN.

In this paper, we present the results of a spectroscopic follow-up program of a stellar-mass-complete sample of six galaxies at 3 < z < 4 originally selected from the NMBS (Marchesini et al. 2010), aimed at spectroscopically confirming their redshifts and further investigating their stellar populations. The brightest galaxy within this sample has been extensively studied and the results were presented in Marsan et al. (2015).

We confirm z > 3 redshifts for half of the sample through the detection of nebular emission lines, and improve the  $z_{phot}$ accuracy for the remainder of the sample with the combination of photometry and spectra. We present the SED modeling of the combined NMBS photometry and spectroscopy and discuss the properties of the mass complete sample at z > 3.

This paper is organized as follows. In Section 2, we describe the target selection, and present the ground-based spectroscopy and data reduction of the mass-selected sample in Section 3. In Section 4, we present the analysis of the spectra, update the redshifts for our targets ( $z_{\text{spec}}$ , if emission lines are detected; otherwise  $z_{\text{cont}}$ , obtained through including binned spectra in the modeling of the observed SEDs), the modeling of the SEDs with updated redshifts, and the investigation of the AGNs content of the sample. The results are summarized and discussed in Section 5. We assume  $\Omega_{\text{M}} = 0.3$ ,  $\Omega_{\Lambda} = 0.7$ ,  $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ , and a Kroupa (2001) initial mass function (IMF) throughout the paper. All magnitudes are in the AB system.

## 2. Target Selection

The targets presented in this paper are primarily selected from the stellar mass complete sample of massive galaxies log  $(M_*/M_{\odot}) > 11.4$  at 3 < z < 4 presented in Marchesini et al. (2010), constructed using the NMBS (Whitaker et al. 2011). The NMBS is a moderately wide, moderately deep nearinfrared imaging survey (van Dokkum et al. 2009) targeting the COSMOS field (Scoville et al. 2007) and AEGIS strip (Davis et al. 2007). The NMBS photometry presented in Whitaker et al. (2011) includes deep optical ugriz data from the CFHT Legacy Survey, deep Spitzer-IRAC and MIPS imaging (Sanders et al. 2007), Galaxy Evolution Explorer (GALEX) photometry in the FUV (150 nm) and NUV (225 nm) passbands (Martin et al. 2005), NIR imaging with NEWFIRM using the five medium-band filters  $J_1$ ,  $J_2$ ,  $J_3$ ,  $H_1$ ,  $H_2$ , broadband K (van Dokkum et al. 2009) for both fields. The COSMOS field additionally includes deep Subaru images with the  $B_I V_I r^+ i^+ z^+$ broadband filters (Capak et al. 2007), Subaru images with 12 optical intermediate-band filters from 427 to 827 nm (Taniguchi et al. 2007), and JHK<sub>S</sub> broadband imaging from the WIRCam Deep Survey (McCracken et al. 2010).

The medium-bandwidth NIR filters in NMBS allows the Balmer/4000 Å break of galaxies at 1.5 < z < 3.5 to be more finely sampled compared to the standard broadband NIR filters providing more accurate photometric redshift estimates (Whitaker et al. 2011). The public NMBS catalog uses photometric redshifts determined with the EAZY code (Brammer et al. 2008) modeling the full FUV-8  $\mu$ m SEDs. Stellar mass and other stellar population parameters were determined using Fitting and Assessment of Synthetic Templates (FAST; Kriek et al. 2009), adopting the redshift solution output by EAZY (i.e.,  $z_{\text{peak}}$ ), the stellar population synthesis models of Bruzual & Charlot (2003), the Calzetti et al. (2000) reddening law with  $A_V = 0-4$  mag, solar metallicity, and an exponentially declining star-formation history (SFH). These SED-modeling assumptions are the same used in the analysis presented in Marchesini et al. (2010)

Table 1 lists the properties of the six targets selected for spectroscopic follow-up. Four objects are the brightest candidates of very massive galaxies at z > 3 selected from Marchesini et al. (2010), namely C1-15182, C1-19536, A2-15753, and C1-23152 (i.e., the Vega galaxy, already studied in detail in Marsan et al. 2015). The spectroscopic program targeted two additional galaxies, C1-2127 and C1-19764, consistent with being at  $z \gtrsim 3$  but not originally included in the sample presented in Marchesini et al. (2010) because they did not strictly satisfy the completeness limit in stellar mass (though both massive with log  $(M_*/M_{\odot}) \sim 11.2-11.3$ ). Table 1 also lists C1-21316 from the sample of Marchesini et al. (2010) but not targeted by our spectroscopic follow-up program, since it has a previously confirmed spectroscopic redshift of  $z_{\rm spec} = 3.971$  (Capak et al. 2010; Smolčić et al. 2012; but see also Miettinen et al. 2015).

	Ground-based Spectroscopic Observations										
				NMBS		NIRSPEC			X-shooter		
ID	$\alpha$ (J2000)	δ(J2000)	K (mag)	Zpeak	$\log M^* \ (M_\odot)$	$\lambda_{ m range} \ (\mu)$	t <sub>exp</sub> (minutes)	FWHM (")	$\lambda_{ m range} \ (\mu)$	$t_{exp}$ (minutes)	FWHM (")
C1-2127	10h00m12.96	$+02^{d}12^{m}11^{s}.5$	$21.69\pm0.07$	$3.17_{-0.12}^{+0.12}$	$11.26\substack{+0.13\\-0.15}$	1.94-2.37	60	0.5			
C1-15182	09 <sup>h</sup> 59 <sup>m</sup> 24 <sup>s</sup> .39	$+02^{d}25^{m}36^{s}.5$	$21.62\pm0.09$	$3.56\substack{+0.11\\-0.11}$	$11.54\substack{+0.04\\-0.05}$	2.04-2.46	60	0.5	0.3-2.4	60	0.5
C1-19536	09h59m31.82	$+02^{d}30^{m}18.2$	$21.65\pm0.06$	$3.19_{-0.08}^{+0.07}$	$11.55_{-0.03}^{+0.03}$	1.94-2.37	60	0.5	0.3-2.4	60	0.8
C1-19764	10h00m21s12	+02d30m33.5	$21.81\pm0.12$	$3.04_{-0.19}^{+0.17}$	$11.22\substack{+0.03\\-0.06}$	1.94-2.37	60	0.5	0.3-2.4	60	1.1
C1-21316 <sup>a</sup>	10 <sup>h</sup> 00 <sup>m</sup> 19 <sup>s</sup> .74	+02d32m04.3	$22.29\pm0.16$	$3.68^{+0.12}_{-0.11}$	$11.52_{-0.61}^{+0.01}$						
C1-23152 <sup>b</sup>	10 <sup>h</sup> 00 <sup>m</sup> 27 <sup>s</sup> .81	$+02^{d}33^{m}49^{s}.3$	$20.31\pm0.02$	$3.29_{-0.06}^{+0.06}$	$11.42_{-0.01}^{+0.01}$	1.48-2.37	75	0.7	0.3-2.4	60	0.5
A2-15753	14h18m30:83	$+52^{d}40^{m}24.6^{s}6$	$22.25\pm0.06$	$3.14\substack{+0.10\\-0.09}$	$11.40\substack{+0.02 \\ -0.09}$	1.94-2.37	210	0.6	•••		

Table 1

Notes. Listed ID's are from the NMBS catalog (v4.4). "C1-" and "A2-" refer to the COSMOS and AEGIS fields, respectively. The listed redshifts are the best-fit EAZY redshifts (z<sub>peak</sub>), quoted stellar masses are the bestfit FAST stellar masses used for the selection of the targets (from the NMBS catalog v4.4, Marchesini et al. 2010).  $\lambda_{range}$  is the wavelength range covered by the instrumental setup;  $t_{exp}$  is the on-source exposure time in minutes; FWHM is the average seeing in arcseconds of the observations. <sup>a</sup> C1-21316 was not included in the spectroscopic follow-up program, but it has a spectroscopic redshift present in Capak et al. (2010). <sup>b</sup> A detailed study of C1-23152 was already presented in Marsan et al. (2015).

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C1-21316 is also a submillimeter galaxy listed as AzTEC 5 in the AzTEC millimeter survey catalog in COSMOS (Scott et al. 2008).

# 3. Observations and Data Reduction

In this section, we briefly describe the ground-based spectroscopic data obtained from Keck-NIRSPEC and VLT-X-shooter as part of follow-up programs aimed at spectroscopically confirming the existence of very massive galaxies at z > 3. Along with the observational techniques, we provide a summary of the data reduction. We refer the reader to Geier et al. (2013) and Marsan et al. (2015) for a detailed description of the spectroscopic data reduction. Table 1 summarizes the wavelength coverage, total exposure time, and seeing of the spectroscopic observations. We note that several targets were observed with both NIRSPEC and X-Shooter (C1-15182, C1-19536, C1-19764, and C1-23152).

## 3.1. NIRSPEC Spectroscopy

We used NIRSPEC (McLean et al. 1998) on the Keck II telescope for H- and K-band spectroscopy of massive z > 3targets primarily in search of their [O III] emission lines. The observations were carried out on the nights of 2011 February 11-13 as part of the NOAO program 2011A-0514 (PI: Marchesini) with a typical seeing of 0."7, which worsened throughout the sequence of observations due to cloudy variable weather. Observations were conducted following an ABA'B' on-source dither pattern. The orientation of the slit was set to include a bright point source when possible to serve as reference when analyzing and combining the two-dimensional rectified frames. Targets were acquired using blind offsets from a nearby bright star. The alignment of the offset star in the slit was checked before each individual 900 s science exposure and corrected when necessary. Before and after each observing sequence, a spectrophotometric standard and an AV0 star was observed for the purpose of correcting for telluric absorption and detector response.

The data reduction for the NIRSPEC observations used a combination of custom IDL scripts and standard IRAF tasks.<sup>8</sup> Bad pixel masks were created by flagging outlier pixels in dark and flat frames. The cosmic rays on the science frames were removed using L.A.Cosmic (van Dokkum 2001). Each spectrum was sky subtracted using an adjacent spectrum with the IDL routines written by G. Becker (2017, private communication). The sky subtracted frames and the sky spectra were rotated and rectified to a linear wavelength scale using a polynomial to interpolate between adjacent pixels. An absolute wavelength dispersion solution was applied with the use of OH skylines. The standard star frames used to correct for atmospheric absorption and detector response were rectified and reduced in the same manner as the science frames. A onedimensional spectrum was extracted for each telluric standard star (before and after science observations) by summing all the rows (along the spatial direction) with a flux greater than 0.1 times that of the central row. The average of the onedimensional telluric star spectra was used to correct the twodimensional science and spectrophotometric star frames for telluric absorption. The one-dimensional spectrum of the spectrophotometric star was extracted in the same manner, and used to create a response function for flux calibration. The position of an on-slit bright source was used to determine the necessary shifts to align target continuum. The two-dimensional rectified and reduced science frames were combined by weighting according to their signal-to-noise ratio (S/N).

#### 3.1.1. VLT-X-shooter

X-shooter is a single-object, medium-resolution echelle spectrograph with simultaneous coverage of wavelength range  $0.3-2.5 \,\mu m$  in three arms (UVB, VIS, NIR; D'Odorico et al. 2006). The observations of our targets were carried out in queue mode as part of the ESO program 087.A-0514 (PI: Brammer) in 2011 May, following an ABA'B' on-source dither pattern using the  $11'' \times 1.''0$  and  $11'' \times 0.''9$  slits for the UVB and VIS/NIR arms, respectively. This instrumental setup resulted in a spectral resolution of R = 4200, 8250, and 4000for the NIR, VIS, and UVB arms, respectively. For calibration purposes, telluric and spectrophotometric standard stars were observed in the same setup as science observations. We refer the reader to Table 1 for the seeing conditions (FHWM, ") and exposure times for each object. The data reduction for the X-shooter observations used custom scripts based on the standard X-shooter reduction pipeline (Modigliani et al. 2010). The calibration steps (master darks, order prediction, flat fields, and the two-dimensional maps for later rectification of the spectra) were run with the default parameters in the pipeline (Goldoni 2011). After these five calibration steps, the echelle spectra were dark-subtracted, flat-fielded, and rectified, and the orders stitched (12, 15, and 16 for the UVB, VIS, and NIR arms, respectively). The sky was then subtracted using adjacent exposures. Standard star observations were reduced with the same calibration data as the science frames and used to correct for telluric absorption and detector response. A final spectrum was created by mean stacking all exposures. We refer to Geier et al. (2013) for a detailed description of reduction steps of X-shooter spectra.

## 3.1.2. Extraction of One-dimensional Spectra

Following Horne (1986), one-dimensional spectra were extracted by summing all adjacent lines (along the spatial direction) using weights corresponding to a Gaussian centered on the central row with a full width at half maximum equal to the slit width used in each observation. To correct for slit losses and obtain an absolute flux calibration, spectroscopic broad/ medium-band fluxes were obtained by integrating over the corresponding filter curves, and a constant scaling was applied to each spectra individually. A binned, lower resolution spectrum with higher S/N was extracted for each spectra using optimal weighting, excluding parts of spectra contaminated by strong sky emission or strong atmospheric absorption. The resulting spectral resolutions of the binned X-shooter spectra were  $R \approx 10-45$ , 15–40, and 25 – 50 for the UVB, VIS, and NIR arms, respectively, while the spectral resolutions of the binned NIRSPEC spectra were  $R \approx 30-100$  and 30-300for the *H* and *K* bands, respectively.

# 4. ANALYSIS and Results

In this section, we built and improved upon previously derived and published stellar population parameters from Marchesini et al. (2010) and Whitaker et al. (2011) by

<sup>&</sup>lt;sup>8</sup> IRAF is distributed by the National Optical Astronomy Observatory, which is operated by the Association of Universities for Research in Astronomy (AURA), Inc., under cooperative agreement with the National Science Foundation.



Figure 1. Observed one-dimensional and two-dimensional spectra in the regions around considered spectral features for C1-19536 and C1-15182. Red solid curves indicate the best-fit single Gaussian profiles to the emission lines in 1D spectra. The gray shaded regions indicate regions of the spectra significantly affected by telluric sky-lines.

modeling the emission-line-corrected SEDs constructed from the combination of the UV-to-8  $\mu$ m NMBS photometry and the binned spectra. We adopted the measured spectroscopic redshift,  $z_{\text{spec}}$ , when available, or  $z_{\text{cont}}$ , the improved photometric redshift estimated by modeling the binned spectrum in combination with the to-8  $\mu$ m NMBS photometry. [O III] nebular emission lines were detected in the two brightest targets, C1-19536 and C1-15182, allowing for the measurement of the spectroscopic redshift. For the remaining targets with continuum detections (C1-2127, A2-15753, and C1-19764), we combined the binned spectrum and the NMBS photometry to model the finely sampled SEDs with EAZY to derive  $z_{cont}$ . In addition to the EAZY template set adopted in Whitaker et al. (2011), we also modeled the SEDs using a template set augmented by an "old-and-dusty" template consisting of a 1 Gyr old single stellar population with  $A_V = 3$  mag of dust extinction. Marchesini et al. (2010) found that the inclusion of such a template caused approximately half of the massive galaxies population at z > 3to be consistent with a somewhat lower redshift in the range 2 < z < 3. For a detailed description of the properties of this template and investigation of its effects on the estimated photometric redshifts, we refer the reader to G. Brammer et al. (2017, in preparation).

The results of the detailed analysis of the spectroscopic data of the brightest *K*-band candidate from Marchesini et al. (2010), namely C1-23152, were presented in Marsan et al. (2015). This paper presented the first spectroscopic

confirmation of an ultra-massive galaxy at redshift z > 3, along with a detailed investigation of the stellar population and structural properties of a progenitor of local most massive elliptical galaxies when the universe was less than 2 Gyr old, showing its ultra-compact nature, the presence of an obscured powerful AGN, and discussing its evolutionary path to z = 0. We refer the reader to Marsan et al. (2015) for the derived stellar population properties of this source.

### 4.1. Emission-line Features and Spectroscopic Redshift

The extracted one-dimensional spectra were used to measure spectroscopic redshifts ( $z_{spec}$ ) when nebular emission lines were detected. The nebular emission lines that we set out to measure were primarily Ly $\alpha$  and [O III] $\lambda\lambda$ 4959, 5007. The [O III] doublet was identified in the spectra of C1-19536 and C1-15182, along with a weak Ly $\alpha$  detection in the spectrum of C1-19536. A symmetric Gaussian profile plus a local continuum was used to fit these lines in the extracted onedimensional spectra. Identical redshifts, FWHM, and a ratio of 1:3 for the amplitudes of the [O III]4959, 5007 doublet were assumed. Figure 1 shows the observed one-dimensional spectra around the regions of the considered spectral lines. The regions of the spectra significantly affected by skylines, indicated in gray regions, were excluded from the line profile modeling. The best-fit Gaussian profiles are indicated in red.

Spectral Line Properties						
Feature	$\lambda_{ m lab} \ ({ m \AA})$	Z,	FWHM (km s <sup>-1</sup> )	EW <sub>obs</sub> (Å)	$\frac{L}{(10^{42} \mathrm{erg  s}^{-1})}$	
C1-19536						
Lyα	1215.24	$3.1544_{-0.0027}^{+0.0017}$	$1535.4^{+1845.6}_{-905.8}$	$280.4{\pm}442.4$	$11.5^{+4.9}_{-11.9}$	
O III(Xsh)	5008.240	$3.1396\substack{+0.0007\\-0.0008}$	$896.53_{-87.51}^{+89.67}$	$520.63^{+142.86}_{-104.16}$	$26.55^{+4.49}_{-3.83}$	
O III(NIRSPEC)	5008.240	$3.1373_{-0.0008}^{+0.0007}$	$871.82^{+212.09}_{-182.63}$	$636.31^{+254.28}_{-142.78}$	$17.62^{+5.93}_{-4.64}$	
C1-15182						
O III(Xsh)	5008.240	$3.3708^{+0.0014}_{-0.0021}$	$734.0^{+479.4}_{-365.9}$	$54.8^{+2.2}_{-18.3}$	$16.0^{+6.8}_{-5.3}$	
O III(NIRSPEC)	5008.240	$3.3692\substack{+0.0025\\-0.0021}$	$767.9^{+127.7}_{-152.2}$	$221.7^{+108.6}_{-76.3}$	$7.9^{+5.9}_{-2.8}$	

 Table 2

 ctral Line Properties

Note.  $\lambda_{lab}$  is the laboratory rest-frame wavelength of the targeted spectral lines; z is the derived redshift; FWHM is the intrinsic velocity width of the spectral lines; EW<sub>obs</sub> is the observed equivalent widths of the spectral lines; and L is the integrated line luminosity calculated using the adopted systemic redshift  $z_{spec}$ .

A Monte Carlo approach was used to measure the uncertainties in the centroid, flux, and width of the fitted emission lines. For each 1D spectrum, 1000 simulated spectra were created by perturbing the flux of the true spectrum at each wavelength by a Gaussian random amount with the standard deviation set by the level of the  $1\sigma$  error at that specific wavelength. Line measurements were obtained from the simulated spectra in the same manner as the actual data. We computed the formal lower and upper confidence limits by integrating the probability distribution of each parameter (centroid, width, continuum, and emission-line flux) from the extremes until the integrated probability is equal to 0.1585. We calculated the fluxes of the observed emission lines by integrating the fitted Gaussian profiles, with uncertainties determined using the  $1\sigma$  distribution of integrated fluxes from Monte Carlo simulations as described above. The equivalent widths were calculated individually for the observed and simulated spectra by dividing the integrated line flux by the continuum of the fit, with uncertainties derived as above.

The  $[O III]\lambda\lambda4009$ , 5007 doublet emission is detected in both NIRSPEC and Xshooter 1D spectra for C1-19536 and C1-15182; however, the lower S/N of the Xshooter spectrum is evident in the panels of Figure 1. This is true especially for C1-15182, where a strong continuum is not detected when fitting the emission lines with the Xshooter spectrum, deeming the calculated equivalent width of [O III] emission highly uncertain. We therefore fixed the continuum level when fitting the [O III] lines in the Xshooter spectrum for C1-15182 to the median flux values of the spectral region ( $\pm 450$  Å) around the [O III] lines and not contaminated by skylines. We used the [O III] emission-line fitting results from the higher S/N NIRSPEC spectrum to analyze the emission-line properties of C1-15182 and C1-19536.

The best-fit redshifts, line velocities corrected for the instrumental profile (determined by the width of skylines in each spectra), and the observed equivalent widths are listed in Table 2, with the quoted uncertainties corresponding to the  $1\sigma$  errors estimated from the Monte Carlo simulations. We estimated the line luminosities by adopting the redshifts of the [O III] emission lines as the systemic redshifts, listed in Table 2. [O III] line luminosities are  $L_{[O III]} = 1.76^{+0.6}_{-0.5} \times 10^{43}$  erg s<sup>-1</sup> and  $7.9^{+5.9}_{-2.8} \times 10^{42}$  erg s<sup>-1</sup> for C1-19536 and C1-15182, respectively, at the high end of [O III] luminosities found in galaxies harboring AGNs (Maschietto et al. 2008; Kuiper et al. 2011; Harrison et al. 2016).

The resulting systemic redshifts for C1-15182 and C1-19536 are  $z_{\text{spec}} = 3.369 \pm 0.002$  and  $z_{\text{spec}} = 3.1373 \pm 0.0008$ , respectively,

spectroscopically confirming their redshifts to be at z > 3. Compared to the photometric redshifts derived in Marchesini et al. (2010) and listed in Table 1, we note that the photometric redshift of C1-19536 is in very good agreement with the spectroscopic redshift. For C1-15182, the photometric redshift is smaller by ~4% in  $\Delta z/(1+z)$  compared to the spectroscopic redshift, but consistent with it at the  $\sim 2\sigma$  level. For C1-21316, the photometric redshift is smaller by ~6% in  $\Delta z/(1 + z)$  compared to the spectroscopic redshift provided by Capak et al. (2010), but still consistent with it at the  $\sim 2.4\sigma$  level. We note, however, that the corresponding spectrum from which  $z_{spec}$  was calculated for C1-21316 has been deemed to be of poor quality (Miettinen et al. 2015). The top panel of Figure 2 shows the comparison between the photometric redshifts from Marchesini et al. (2010) and the spectroscopic redshifts or the improved  $z_{\text{cont}}$ . The [O III] line widths listed in Table 2 are  $\approx 768^{+123}_{-152}$  km s<sup>-1</sup> and  $871^{+212}_{-183}$  km s<sup>-1</sup> for C1-15182 and C1-19536, respectively, typical of AGNs of similar  $L_{\rm [O III]}$  at low-z (Hao et al. 2005) and intermediate-z (Harrison et al. 2016).

#### 4.2. Improved Photometric Redshifts

When no emission lines are detected, we used the 1D binned spectra combined with the full FUV-8  $\mu$ m broad- and mediumband photometry from the NMBS to determine more accurate photometric redshifts ( $z_{cont}$ ). Following Muzzin et al. (2013a), the default template set used in this work consists of nine templates: the six templates taken from the optimized template set of EAZY, but augmented with emission lines; a template of a 12.5 Gyr old single stellar population constructed using the Maraston (2005) models; a 1 Gyr old post-starburst template; and a slightly dust-reddened Lyman break template. The SEDs of the galaxies in our sample are shown in Figure 3 (left panels), with the EAZY redshift probability functions plotted on the right panels along with the spectroscopic redshift (when available; red), the  $z_{peak}$  from Marchesini et al. (2010; gray), and the  $z_{cont}$  (black).

Marchesini et al. (2010) found that up to ~50% of the massive 3 < z < 4 galaxy sample could be contaminated by a previously unrecognized population of massive, old, and very dusty galaxies at z < 3. As in Marchesini et al. (2010), we considered the case of adding an additional template representing an old (1 Gyr;  $\tau = 100$  Myr) and very dusty ( $A_V = 3$  mag) galaxy and repeating the redshift analysis. The resulting redshift probability distributions and best fit  $z_{\text{cont}}$  are overplotted in the right panels of Figure 3 (dotted–dashed blue). The resulting photometric redshift estimates of only two galaxies in our sample (A2-15753 and C1-21316) are



**Figure 2.** Comparison of the best-fit stellar masses and redshifts for sources with continuum detections and spectroscopic confirmation (C1-23152, red; C1-19536, aquamarine; C1-15182, orange; C1-21316, yellow; C1-2127, magenta; C1-19764, green and A2-15753, blue).  $z_{peak}$  and log  $M_{*}(z_{peak})$  are from the NMBS catalog v4.4 (Marchesini et al. 2010). Stars indicate the spectroscopic redshifts when available. Error bars indicate  $1\sigma$  limits output from EAZY (Brammer et al. 2008) and FAST. Gray dotted–dashed line is the 1–1 relation for redshifts and steller mass.

influenced by the addition of the old and dusty template in the analysis, yet both estimates still formally lie at z > 3. Only the photometric redshift of C1-19764, not included in the Marchesini et al. (2010) sample, moves below z = 3 when the old and dusty template is including in EAZY, with a formal solution of  $z_{\text{cont}} = 2.72^{+0.26}_{-0.29}$ .

We note that the spectroscopic redshifts of the galaxies in our sample tend to be different from the photometric redshifts at more than  $1\sigma$ . Specifically, we can see from the right panels of Figure 3 that EAZY overestimates the redshift of twothirds of our targets by ~0.17 and ~0.21 for C1-15182 and C1-19536, corresponding to ~4% and ~5% in  $\Delta z/(1 + z)$ , respectively.

# 4.3. SED Modeling and Stellar Population Properties

In order to robustly constrain the stellar population parameters, all spectra and photometry must be corrected for contamination from nebular emission lines. For sources with detected nebular emission lines, we corrected the observed photometry for emission-line contamination by comparing the observed-frame equivalent width of each emission line to the bandwidth of the corresponding filter. We find that the contribution due to [O III] emission is ~15% and ~5% for C1-19536 and C1-15182, respectively, for the NMBS *K* and *K*<sub>S</sub> filters. Due to the highly uncertain equivalent width calculated for the Ly $\alpha$  emission line of C1-19536, we chose to remove the contaminated filter (Subaru, IA505) when proceeding with the SED fittings.

We estimated the stellar population properties by fitting the binned X-shooter and NIRSPEC spectra in combination with the broadband and medium-band photometry with stellar population synthesis (SPS) models. We used FAST (Fitting and Assessment of Synthetic Templates; Kriek et al. 2009) to model and fit a full grid in metallicity, dust content, age, and star-formation timescale. Stellar population synthesis models of Bruzual & Charlot (2003) were adopted assuming a Kroupa (2001) IMF, an exponentially declining SFH and the Calzetti et al. (2000) extinction law. The age of the stellar population's fit ranged between 10 Myr and the maximum age of the universe at the redshift of sources with a step size of 0.1 dex. We adopted a grid for  $\tau$  between 3 Myr and 10 Gyr in steps of 0.10 dex and allowed the dust attenuation  $(A_V)$  to range from 0 to 7 mag with a step size of 0.01 mag. We initially modeled the observed SED with the metallicity as a free parameter (Z = 0.004, 0.08, 0.02, 0.05) and repeated the modeling by treating the metallicity as a systematic uncertainty by fixing it to solar metallicity ( $Z_{\odot}$ ). Figure 4 shows the observed SEDs of our targets. The results of the SED modeling and the corresponding  $1\sigma$  errors are listed in Table 3. The bottom panel of Figure 2 shows the comparison between the stellar masses from Marchesini et al. (2010) and the stellar masses derived from the improved redshifts (either  $z_{\text{spec}}$  or  $z_{\text{cont}}$ ) and better sampled SEDs.

As shown in the lower panel of Figure 2, the stellar masses obtained when adopting the improved redshifts and better sampled SEDs are consistent with those based only on the NMBS photometry (Marchesini et al. 2010), spectroscopically confirming the existence of ultra-massive (i.e.,  $\log(M_*/M_{\odot})$ ) >11.4) galaxies at z > 3, when the universe was younger than 2 Gyr. We note that, although error bars on both axes in the panels of Figure 2 correspond to  $1\sigma$  limits, the stellar population parameter space over which the SEDs are modeled are not identical (the observed SEDs are modeled in this work with a finer grid in parameter space than Marchesini et al. 2010, resulting in larger  $1\sigma$  errors). For C1-19536 and C1-15182, which have detected nebular emission lines, we find that removing the emission-line contamination to the observed SEDs decreases the stellar masses by  $\sim 0.1$  dex (still consistent with being above the mass completeness threshold in Marchesini et al. 2010). The stellar mass of A2-15753 is affected the most when modeling the observed SED including the binned 1D spectra, resulting in a stellar mass of log  $(M_*/M_{\odot}) = 11.14^{+0.13}_{-0.23}$ , i.e., ~0.25 dex smaller than what was derived in Marchesini et al. (2010).

The stellar mass of C1-2127 and C1-19764, not part of the sample of ultra-massive galaxies at z > 3 from Marchesini et al. (2010), remain effectively the same, i.e., log  $(M_*/M_{\odot}) = 11.23^{+0.12}_{-0.44}$  and  $11.12^{+0.20}_{-0.24}$ . C1-21316 is the only source with a larger stellar mass when remodeling with improved redshift. This is not surprising because the updated



**Figure 3.** Left: observed SEDs from the combination of the medium- and broadband NMBS photometry (black filled circles) and the binned spectra (NIRSPEC, green; X-shooter, purple). The gray curves represent the best-fit EAZY templates using the default EAZY template set adopted in this work (light gray when modeling only the medium- and broadband photometry; dark gray when also including the 1D binned spectra). The blue dotted–dashed curves represent the best-fit EAZY template when the additional old and dusty template is included in the estimate of the photometric redshifts. Right: the EAZY redshift probability distributions. The gray curves represent the distribution calculated using only medium- and broadband photometry (NMBS v4.4), whereas the black curves represent the resulting redshift distribution including the 1D binned spectra in EAZY. The blue dotted–dashed curves show the resulting redshift probability distributions when the old and dusty SED is included in the template set of EAZY to estimate the photometric redshifts. Vertical lines indicate the  $z_{peak}$  output by EAZY for different template sets with 1 $\sigma$  uncertainties. Red vertical lines indicate spectroscopic redshifts when available.



Figure 4. Left: observed SEDs from the combination of the medium- and broadband NMBS photometry (black filled circles) and the binned spectra (NIRSPEC in green; X-shooter in purple). The gray curves represent the best-fit models adopting the BC03 population model, with an exponentially declining SFH and free metallicity. Best-fit stellar population properties and redshifts are indicated in each panel.

spectroscopic redshift of the source is  $\Delta z \sim 0.3$  greater than the  $z_{\text{phot}}$  quoted in Marchesini et al. (2010).

The time elapsed since the onset of star formation for our targets ranges from ~400 Myr to ~ 1.25 Gyr, consistent with having formation redshifts of  $z_{\rm form} \sim 5-9$ . For roughly half of the sample, the best-fit SFR-weighted mean ages  $\langle t \rangle_{\rm SFR}$  (as defined in Föster Scheriber et al. 2004) are  $\approx$ 800 Myr. Two of the massive galaxies (C1-19764, with  $z_{\rm cont} < 3$  and C1-2127) in our sample have best-fit SFR estimates from SED modeling consistent with low SFR ( $<1 M_{\odot} \text{ yr}^{-1}$ ), while the remaining galaxies have SFRs on the order of few tens to hundreds of solar masses per year. We stress that the estimated SFRs, stellar ages, and dust extinction values have large uncertainties.

Figure 5 shows the rest-frame U - V versus V - J diagram (hereafter, UVJ diagram) with the rest-frame colors of our targets (color scheme identical to Figure 2) overlayed on the distribution of the rest-frame colors of all galaxies at 3.0 < z < 4.0 from the  $K_S$ -selected UltraVISTA catalog of Muzzin et al. (2013a) (grayscale representation). The restframe UVJ diagram is a powerful diagnostic to separate starforming and quiescent galaxy populations in color–color space (Labbé et al. 2005; Wuyts et al. 2007; Brammer et al. 2011; Whitaker et al. 2011; Muzzin et al. 2013b). We calculated the rest-frame U - V and V - J colors of the targets using EAZY (Brammer et al. 2008). A Monte Carlo approach was used to measure the uncertainties in the U - V and V - J colors. Specifically, 1000 photometry catalogs were created by perturbing each flux by a Gaussian random number with the standard deviation set by the level of each flux error. The simulated catalogs were each fit with EAZY separately, and the formal upper and lower limits were obtained in a similar manner as for the emission-line fits. The rest-frame colors and  $1\sigma$  confidence limits for the galaxies in our sample are plotted in Figure 5 in solid colored circles (color scheme identical to Figure 2). Figure 5 also shows color-color evolution tracks of Bruzual & Charlot (2003) models stellar population synthesis models for different SFHs and dust extinction values. The orange solid and dashed tracks represent the evolution of an exponentially declining SFH with  $\tau = 100$  Myr and  $A_{\rm V} = 0$ and 2 mag, respectively. The blue solid and dashed tracks represent the evolution for a constant SFH with  $A_{\rm V} = 0$  and 2 mag, respectively. The evolution of an exponentially declining SFH with  $\tau = 500$  Myr and  $A_V = 1$  is indicated by the dashed green track.

Based on rest-frame colors, only one (C1-19764, only galaxy with best-fit  $z_{cont} < 3$ ) of the targets falls firmly in the quiescent

ID	Z	log $ au$ (yr)	Metallicity	log(Age) (yr)	A <sub>V</sub> (mag)	$\log(M^*) \ (M_{\odot})$	SFR $(M_{\odot} \text{ yr}^{-1})$	$\chi^2$		
C1-2127	$3.19\substack{+0.08\\-0.07}$	$\begin{array}{c} 6.70\substack{+3.30\\-0.20}\begin{pmatrix}+3.30\\-0.20\end{pmatrix}\\ 8.00\substack{+0.24\\0.000}\begin{pmatrix}+2.00\\0.000\end{pmatrix}\end{array}$	$\begin{array}{c} 0.008\substack{+0.042\\-0.004}(\substack{+0.042\\-0.004})\\ 0.02\end{array}$	$7.90^{+0.82}_{-0.90}(^{+1.20}_{-0.90}) \\ 8.60^{+0.19}_{-0.18}(^{+0.35}_{-1.60})$	$2.00^{+1.10}_{-1.10}(^{+1.18}_{-1.70}) \\ 1.40^{+0.37}_{-0.47}(^{+1.60}_{-0.90})$	$11.23^{+0.16}_{-0.44} ( {}^{+0.26}_{-0.54} ) \\ 11.35^{+0.06}_{-0.04} ( {}^{+0.14}_{-0.60} )$	$\begin{array}{c} 0.0067\substack{+7244.\\-0.006}(\substack{+8128.\\-0.006})\\ 53.7\substack{+69.3\\-35.9}(\substack{+6402.\\-49.53})\end{array}$	1.12 1.17		
C1-15182	$3.3703\substack{+0.0013\\-0.0011}$	$\begin{array}{c} 8.00\substack{+0.92\\-1.50}(\substack{+2.00\\-1.50})\\ 8.30\substack{+0.33\\-1.14}(\substack{+0.76\\-1.80})\end{array}$	$\begin{array}{c} 0.004^{+0.046}_{0.000}(^{+0.046}_{0.000})\\ 0.02\end{array}$	$\begin{array}{c} 8.60\substack{+0.60\\-0.79}(\substack{+0.60\\-1.60})\\ 8.80\substack{+0.39\\-0.71}(\substack{+0.40\\-1.03})\end{array}$	$\frac{1.80^{+0.59}_{-0.73}(^{+1.52}_{-1.26})}{1.60^{+0.50}_{-0.61}(^{+0.72}_{-1.04})}$	$11.45^{+0.22}_{-0.16} \begin{pmatrix} +0.28 \\ -0.52 \end{pmatrix}$ $11.56^{+0.09}_{-0.16} \begin{pmatrix} +0.18 \\ -0.29 \end{pmatrix}$	$\begin{array}{c} 69.2^{+367.3}_{-69.18}(^{+10895}_{-69.2}) \\ 102.3^{+137.5}_{-96.3}(^{+398.8}_{-102.3}) \end{array}$	1.86 1.85		
C1-19536	$3.1385\substack{+0.0005\\-0.0005}$	$\begin{array}{c} 8.30\substack{+0.14\\-0.28}\\ 8.50\substack{+0.15\\-0.11}\\ (-0.34) \end{array}$	$\begin{array}{c} 0.05\substack{+0.000\\-0.014}(\substack{+0.000\\-0.046})\\ 0.02\end{array}$	$\begin{array}{c}9.00\substack{+0.14\\-0.36}\\9.10\substack{+0.20\\-0.10}\begin{pmatrix}+0.20\\-0.37\end{pmatrix}\end{array}$	$\begin{array}{c} 0.60\substack{+0.10\\-0.15}(\substack{+1.13\\-0.31})\\ 0.90\substack{+0.10\\-0.51}(\substack{+0.72\\-0.61})\end{array}$	$11.41^{+0.03}_{-0.05} \begin{pmatrix} +0.17 \\ -0.15 \end{pmatrix}$ $11.45^{+0.09}_{-0.10} \begin{pmatrix} +0.18 \\ -0.15 \end{pmatrix}$	$11.5^{+1.4}_{-1.5}(^{+117.3}_{-5.7}) \\ 22.9^{+4.0}_{-15.7}(^{+89.3}_{-16.6})$	1.35 1.44		
C1-21316	3.971	$\begin{array}{c} 8.70^{+1.30}_{-2.20}(^{+1.30}_{-2.20})\\ 9.30^{+0.70}_{-2.46}(^{+0.70}_{-2.80})\end{array}$	$\begin{array}{c} 0.05\substack{+0.00\\-0.046}(\substack{+0.00\\-0.046})\\ 0.02\end{array}$	$\begin{array}{c}9.10\substack{+0.00\\-1.66}\begin{pmatrix}+0.00\\-2.10\end{pmatrix}\\9.10\substack{+0.00\\-1.51}\begin{pmatrix}+0.00\\-2.10\end{pmatrix}\end{array}$	$\begin{array}{c} 1.10^{+1.21}_{-0.51}(^{+1.84}_{-0.91}) \\ 1.60^{+0.74}_{-0.70}(^{+1.33}_{-1.02}) \end{array}$	$11.61^{+0.04}_{-0.55} \begin{pmatrix} +0.07 \\ -0.94 \end{pmatrix}$ $11.58^{+0.06}_{-0.52} \begin{pmatrix} +0.09 \\ -0.91 \end{pmatrix}$	$95.5^{+795.7}_{-95.3} \begin{pmatrix} +9024. \\ -95.5 \end{pmatrix} \\ 281.8^{+1416.}_{-230.5} \begin{pmatrix} +7846. \\ -281.8 \end{pmatrix}$	0.64 0.68		
C1-19764	$2.72\substack{+0.26\\-0.29}$	$\begin{array}{c} 8.20\substack{+0.42\\-1.70}(\substack{+0.81\\-1.70})\\ 8.10\substack{+0.30\\-1.60}(\substack{+0.79\\-1.60})\end{array}$	$\begin{array}{c} 0.008\substack{+0.042\\-0.004}(\substack{+0.042\\-0.004})\\ 0.02\end{array}$	$\begin{array}{c}9.10\substack{+0.30\\-0.54}\begin{pmatrix}+0.30\\-0.83\end{pmatrix}\\9.00\substack{+0.25\\-0.46}\begin{pmatrix}+0.40\\-0.81\end{pmatrix}\end{array}$	$\begin{array}{c} 0.20\substack{+1.13\\-0.20}(\substack{+2.19\\-0.20})\\ 0.20\substack{+1.02\\-0.20}(\substack{+2.12\\-0.20})\end{array}$	$11.12^{+0.20}_{-0.24} \begin{pmatrix} +0.26 \\ -0.39 \end{pmatrix}$ $11.20^{+0.10}_{-0.13} \begin{pmatrix} +0.18 \\ -0.40 \end{pmatrix}$	$\begin{array}{c} 0.4^{+3.8}_{-0.4}(^{+28.4}_{-0.4}) \\ 0.6^{+3.7}_{-0.6}(^{+28.9}_{-0.6}) \end{array}$	0.81 0.83		
A2-15753	$3.13\substack{+0.07\\-0.07}$	$\begin{array}{c} 8.90\substack{+1.10\\-0.51}(\substack{+1.10\\-2.40})\\ 10.0\substack{+0.00\\-0.98}(\substack{+0.00\\-3.50})\end{array}$	$\begin{array}{c} 0.05\substack{+0.00\\-0.033}(\substack{+0.00\\-0.046})\\ 0.02\end{array}$	$\begin{array}{c}9.10\substack{+0.30\\-0.43}\begin{pmatrix}+0.30\\-1.80\end{pmatrix}\\9.40\substack{+0.00\\-0.41}\begin{pmatrix}+0.00\\-1.79\end{pmatrix}\end{array}$	$\begin{array}{c} 1.20\substack{+0.40\\-0.39}(\substack{+1.36\\-0.80})\\ 1.40\substack{+0.31\\-0.27}(\substack{+0.82\\-0.68})\end{array}$	$11.14^{+0.13}_{-0.23}(^{+0.26}_{-0.73}) \\ 11.17^{+0.13}_{-0.15}(^{+0.23}_{-0.60})$	$57.5^{+74.3}_{-32.4}(^{+965.7}_{-57.5}) \\ 69.2^{+62.6}_{-28.4}(^{+328.9}_{-69.2})$	0.91 1.05		

 Table 3

 Best-fit Stellar Population Parameters

**Note.** Estimated stellar population parameters from the modeling of the binned UV-to-NIR spectra in combination with the broad- and medium-bandwidth UV-to-8  $\mu$ m photometry from NMBS derived assuming the Bruzual & Charlot (2003) stellar population synthesis models and an exponentially declining SFH. Spectroscopic redshifts are listed when available, otherwise best-fit EAZY outputs are listed. Quoted errors are 1 $\sigma$  confidence intervals output by FAST and EAZY. The values in parantheses correspond to  $3\sigma$  confidence intervals. A Kroupa (2001) IMF and a Calzetti et al. (2000) extinction law is assumed in all cases.



**Figure 5.** Rest-frame U - V vs. V - J color–color diagram. The grayscale representation indicates the distributions of 3.0 < z < 4.0 galaxies above the 95% stellar mass completeness limits from the  $K_{\rm S}$ -selected UltraVISTA catalog (Muzzin et al. 2013a). The cuts used to separate star-forming from quiescent galaxies from Muzzin et al. (2013b) are shown as the solid black lines. Galaxies are indicated with the identical color scheme as in Figure 2. Color evolution tracks of Bruzual & Charlot (2003) models are also shown: exponentially declining SFHs with  $\tau = 100$  Myr in orange (A = 0 mag, solid;  $A_V = 2$  mag, dashed) and  $\tau = 500$  Myr in green with  $A_V = 1$ , and constant SFHs in blue (A = 0 mag, solid;  $A_V = 2$  mag, dashed). The dust vector indicates an extinction of  $A_V = 1$  mag for a Calzetti et al. (2000) extinction curve.

region. Two galaxies with confirmed  $z_{spec}$ , C1-19536 and C1-15182, that are also hosts to powerful AGNs (based on the X-ray detections see Section 4.4.4), are on the border of the quescient—star-forming galaxy rest-frame color separation. The quenched Vega galaxy (C1-23152, Marsan et al. 2015), has rest-frame colors consistent with transitioning into the quiescent region. The positions of the remaining three targets (C1-2127, C1-21316, and A2-15753) indicate that they are dusty, star-forming galaxies, in line with the derived stellar population parameters based on SED modeling.

#### 4.4. Active Galactic Nuclei Content

### 4.4.1. Infrared Spectral Energy Distribution

We found that five-sixths of our targets have significant  $(>3\sigma)$  MIPS 24  $\mu$ m detections (four already studied in Marchesini et al. 2010), consistent with the high fraction of MIPS-detected sources in the sample of IRAC-selected massive galaxies at z > 3.5 over GOODS-North (Mancini et al. 2009). At the redshifts of our targets, the observed 24  $\mu$ m probes the rest-frame  $\sim$ 4.8–7.8  $\mu$ m. Emission at these wavelengths can arise from warm/hot dust and polycyclic aromatic hydrocarbons (Draine & Li 2007), heated by dust-enshrouded star formation or AGNs.

Figure 6 shows the observed far-infrared (FIR) SEDs of our targets with significant MIPS 24  $\mu$ m detections. We include the Herschel PACS 100 and 160  $\mu$ m photometry from the PACS Evolutionary Probe (PEP) survey DR1 public data release9 (Lutz et al. 2011; Magnelli et al. 2013) and Herschel SPIRE 250, 350, and 500 µm photometry from the Herschel Multi-tiered Extragalactic Survey (HerMES) DR3/4 public data release<sup>10</sup> (Roseboom et al. 2010, 2012; Oliver et al. 2012; Hurley et al. 2017). These fluxes are listed in Table 4. For C1-21316 (a.k.a. AzTEC 5), we also included the de-boosted photometry from SMA at 890  $\mu$ m  $(S_{890 \ \mu m} = 9.3 \pm 1.3 \text{ mJy}; \text{ Younger et al. 2007}), \text{ JCMT AzTEC}$ at 1.1 mm ( $S_{1.1 \text{ mm}} = 6.5 \pm 1.4 \text{ mJy}$ ; Scott et al. 2008), ASTE AZTEC at 1.1 mm ( $S_{1.1mm} = 4.8 \pm 1.1 \text{ mJy}$ ; Aretxaga et al. 2011), and JCMT SCUBA2 at 450 and 850  $\mu$ m ( $S_{450 \ \mu m} =$  $25.35 \pm 6.04$  mJy and  $S_{850 \ \mu m} = 11.42 \pm 1.38$  mJy; Casey et al. 2013). Most of our sources are not detected in any of the

<sup>&</sup>lt;sup>9</sup> http://www.mpe.mpg.de/ir/Research/PEP/DR1

<sup>&</sup>lt;sup>10</sup> http://hedam.lam.fr/HerMES/index/dr4



**Figure 6.** FIR SED with *Spitzer* MIPS 24  $\mu$ m, *Herschel* PACS 100, 160  $\mu$ m, SPIRE 250, 350, 500  $\mu$ m with IR templates overplotted. For C1-21316 (AzTEC-5), the plotted photometry also includes the de-boosted fluxes from SMA at 890  $\mu$ m (Younger et al. 2007), JCMT AzTEC at 1.1 mm (Scott et al. 2008), ASTE AzTEC at 1.1 mm (Aretxaga et al. 2011), and JCMT SCUBA2 at 450 and 850  $\mu$ m (Casey et al. 2013). Filled black circles represent observations with detections S/N > 1. Photometric points with no detection are indicated with their 1, 2, and 3 $\sigma$  upper limits. The red dashed SED represents the average SED of Dale & Helou (2002) templates used in Marchesini et al. (2010). Dashed (blue) templates represent the mean SEDs for  $\alpha$ =1,..., 2.5 from the template set of Dale et al. (2014) for varying degrees of AGN contribution (dark blue: 0% to light blue: 100% AGN contribution in uniform 20% steps). The thick sold curve represents the template from Dale et al. (2014) with the minimum amount of AGN contamination allowed by the IR fluxes or 3 $\sigma$  upper limits. The minimum amount of AGN contribution is also specified in blue. Two blue solid curves and AGN values are provided for C1-21316, depending on the IR data used to constrain the AGN contamination (the SCUBA 450  $\mu$ m or the fluxes at  $\lambda$  > 800  $\mu$ m). Violet SEDs are the IR color-based templates from Kirkpatrick et al. (2015). Each panel specifies whether the galaxy is also an X-ray AGN or a radio-loud AGN. Bottom right panel: distribution of our sources in IR colorspace using the observed frame colors (colors same as in Fig 2). The black line represents the empirical separation between the AGNs and SFGs defined in Kirkpatrick et al. (2013), the gray dashed lines show the dividing lines between color regions from which average templates are calculated in Kirkpatrick et al. (2015). The  $\langle f (AGN)_{IR} \rangle$  for each colorspace is indicated in the corresponding region (Kirkpatrick et al. 2015).

*Herschel* bands. Only C1-21316, which is a submillimeter galaxy and a radio-loud AGN, is detected in SPIRE and at longer wavelengths; C1-2127 is detected in the blue and green SPIRE bands, but not at 500  $\mu$ m. No source is detected in PACS despite all being robustly detected in the MIPS 24  $\mu$ m band. Figure 6 shows the IR SEDs of the targeted sources, along with the 1, 2, and  $3\sigma$  upper limits indicated for the bands without detections.

The 24  $\mu$ m emission has been widely used in the literature to estimate total infrared luminosities ( $L_{IR} = L_{8-1000 \ \mu\text{m}}$ ; see, e.g., Papovich et al. 2007; Rigby et al. 2008), which can then be transformed into dust-enshrouded SFRs (Kennicutt 1998). In this section, we will show that this approach cannot be blindly adopted in ultra-massive galaxies at z > 3 given the almost ubiquitous presence of obscured AGNs in this population and the resulting AGN contamination of the 24  $\mu$ m emission.

First, we derived dust-enshrouded SFRs from the MIPS 24  $\mu$ m fluxes using the approach presented in Wuyts et al. (2008), which has become one of the most adopted approaches in the literature. Specifically, this method uses the mean log*L*<sub>IR, $\alpha=1,...,2.5$ </sub> of the infrared SEDs of star-forming galaxies provided by Dale & Helou (2002) to calculate SFR<sub>IR</sub>. This method was validated out to  $z \sim 3.5$  for lower mass (i.e.,  $\log(M_*/M_{\odot}) < 11$ ) star-forming galaxies using Herschel data (Wuyts et al. 2011; Tomczak et al. 2016). The calculated total 8–1000  $\mu$ m rest-frame luminosities and the corresponding SFR<sub>IR</sub> adapted for a Kroupa (2001) IMF from Kennicutt (1998) are listed in Table 5, third and fourth columns. The corresponding errors are also listed, with and without the uncertainties due to random photometric redshift uncertainties when spectroscopic redshifts are not available. As shown in

ID	$\frac{S_{\nu} (100 \ \mu \text{m})}{(\text{mJy})}$	$S_{\nu} (160 \ \mu \text{m}) \\ (\text{mJy})$	$\frac{S_{\nu} (250 \ \mu \text{m})}{(\text{mJy})}$	$S_{ u}$ (350 $\mu$ m) (mJy)	$\frac{S_{\nu} (500 \ \mu \text{m})}{(\text{mJy})}$			
C1-2127	(1.8)	(3.4)	$13.4^{+2.2}_{-2.3}$	$22.2_{-4.7}^{+4.7}$	(10.9)			
C1-15182	(1.7)	(3.4)	(3.6)	(2.1)	(2.9)			
C1-19536	(1.7)	(3.4)	(2.7)	(3.6)	(5.1)			
C1-19764	(1.7)	(3.4)	(2.7)	(3.6)	(5.1)			
C1-21316	(2.0)	(3.4)	$29.1^{+2.5}_{-3.2}$	$38.3^{+3.1}_{-3.4}$	$31.6^{+6.6}_{-29.2}$			
A2-15753	(1.2)	(2.5)	(3.6)	(5.9)	(5.6)			

 Table 4

 FIR Flux Density Detections and Limits

Note. The Herschel PACS 100, 160  $\mu$ m, SPIRE 250, 350, 500  $\mu$ m detections and limits. Photometric points with detections are listed with the corresponding 1 $\sigma$  errors, whereas non-detections are represented with the 1 $\sigma$  error limits in parenthesis.

Table 5, the random uncertainties on  $L_{IR}$  and SFR<sub>IR</sub> as derived using the Wuyts et al. (2008) approach are dominated by the contribution from random photometric redshift errors. We note that a significant scatter of  $\sim 0.2-0.3$  dex was found by Wuyts et al. (2011) and Tomczak et al. (2016) between the MIPS 24  $\mu$ m derived SFRs and the SFRs derived including robust Herschel PACS detections. Therefore, we stress that the total error budget of  $L_{IR}$  and SFR<sub>IR</sub> listed in Table 5 as derived using the Wuyts et al. (2008) approach is dominated by this scatter and can be as large as a factor of a few. When the Wuyts et al. (2008) approach is adopted, the  $L_{\rm IR}$  of our targets range from  $\sim 5 \times 10^{12} L_{\odot}$  to  $\sim 5 \times 10^{13} L_{\odot}$ , typical of Ultra Luminous IR galaxies (ULIRGs) and Hyper Luminous IR galaxies (HLIRGs; Murphy et al. 2011), corresponding to SFR  $\sim 500-5000 M_{\odot} \text{ yr}^$ of obscured star formation. We used the rest-frame 2800 Å luminosity ( $\nu L_{2800 \text{ Å}}$ ), determined via the best-fit templates in EAZY (using the same methodology for deriving rest-frame colors described in Brammer et al. 2011), as a proxy for  $L_{\rm UV}$ , reflecting the contribution of unobscured star formation. We used the approach detailed in Bell et al. (2005) to convert  $\nu L_{2800 \text{ Å}}$  to SFR<sub>UV</sub>. The calculated  $L_{\rm UV}$  and SFR<sub>UV</sub> are listed in Table 5.

The SFRs derived by assuming that all the flux at the observed 24  $\mu$ m is associated with dust-enshrouded star formation are  $\sim 2-3$  orders of magnitude larger than the SFRs estimated from SED modeling (although values for C1-2127 and A2-15753 are consistent within  $1\sigma$  uncertainties). We can see in Figure 6 that the median pure-SF IR SED template derived from the library of Dale & Helou (2002), and adopted in the Wuyts et al. (2008) approach, is not an adequate fit to the observed detections and limits. For each object, we overplot the median Dale & Helou (2002) template scaled to the observed MIPS 24  $\mu$ m detections as dashed red curves. It can be seen that a model in which the IR emission is due only to obscured star formation cannot reproduce the observed IR SEDs. Specifically, the detections and upper limits in *Herschel* bands put strong constraints on the overall shape of the IR SEDs, rejecting the scenario in which the observed IR properties are due entirely to obscured star-formation activity, and pointing to the presence of a component with warmer dust temperature (arguably originating from the dusty torus of the AGN) not present in the average template from Dale & Helou (2002). This result was already presented in Marsan et al. (2015) for the Vega galaxy at z = 3.35, and we now show that it appears to be a common property of ultra-massive galaxies at z > 3.

To quantitatively assess the amount of AGN contribution to the IR SEDs, we used the IR template set of Dale et al. (2014), which includes fractional AGN contributions to the mid-

infrared (MIR) radiation. We calculated the mean  $\log(L_{IR})$  for  $\alpha = 1,..., 2.5$  from this template set for varying AGN contributions. For visual ease in comparison, we only plot the templates with AGN contributions of 0%, 20%, 40%, 60%, 80%, and 100% scaled to the observed MIPS 24  $\mu$ m photometry (dashed blue curves in Figure 6). The Dale et al. (2014) templates that minimize the  $\chi^2$  to detections and are consistent with upper limits for each target are plotted as solid blue curves, with  $f(AGN)_{MIR}$  indicated in each panel. We note that the specified f(AGN)<sub>MIR</sub> refers to the fraction of AGN light to the integrated mid-infrared (i.e., 5–20  $\mu$ m) emission. A significant amount of AGN contribution to the MIR SEDs is required by all ultra-massive galaxies at z > 3 to explain the IR SEDs. If  $3\sigma$  upper limits for the *Herschel* photometry are adopted, we find a lower limit on the fraction of AGNs to the MIR SEDs ranging from 15% to 50% depending on the galaxy. If the  $1\sigma$  upper limits are adopted, we find a lower limit on the fraction of AGNs to the MIR SEDs ranging from 40% to 75%. Columns 5–8 of Table 5 list the values or upper limits of  $L_{\rm IR}$ and SFR<sub>IR</sub>, and the values or lower limits of  $f_{Bol}(AGN)$  and  $f_{\rm MIR}(\rm AGN)$  as obtained from adopting the best-fit templates from Dale et al. (2014). The upper limits were derived adopting the  $3\sigma$  upper limits in the *Herschel* photometry. We stress that these newly derived  $L_{IR}$  and SFRs should be considered as upper limits, given that we only have lower limits on the AGN contribution to the IR. We note that the three galaxies with X-ray detections, namely A2-15753, C1-15182, and C1-19536 (see Section 4.4.3 below), are undetected in Herschel while being detected at several  $\sigma$  in the MIPS 24  $\mu$ m band, f(AGN)<sub>MIR</sub> > 40% (75%), 50%(80%), and 15%(60%), respectively, when adopting the  $3\sigma$  (1 $\sigma$ ) upper limits in *Herschel*. One target, C1-21316, which has  $f (AGN)_{MIR} \sim 0.35-0.50$  is detected in radio data, indicative of hosting a radio-loud AGN (Section 4.4.4).

We also used the templates from the observed FIR colorbased library of Kirkpatrick et al. (2015) to investigate the AGN continuum contribution to the IR SEDs. In this study, libraries of empirical IR SED templates were created from a large sample of high-redshift ULIRGS with rest-frame mid-IR spectroscopy from the *Spitzer Space Telescope* Infrared Spectrograph (IRS). We display their color diagnostic in the bottom right panel of Figure 6, with the FIR colors of our targets calculated using photometry from *Herschel* SPIRE and *Spitzer* MIPS/IRAC plotted in solid colored circles (color scheme identical to Figures 2 and 5), with arrows indicating  $2\sigma$ limits. Kirkpatrick et al. (2015) used mid-IR features to determine the strength of the AGN contribution to IR SEDs by jointly fitting the mid-IR spectrum of a prototypical starburst

	Spitzer-24 $\mu$ m Fluxes and the Derived $L_{IR}$ and SFR								
ID	S <sub>24</sub> (μJy)	$L_{ m IR}^{ m a}$ ( $10^{12}L_{\odot}$ )	$\frac{\text{SFR}_{\text{IR}}^{\text{a}}}{(M_{\odot} \text{ yr}^{-1})}$	$L_{\rm IR}^{\rm b}$ (10 <sup>12</sup> $L_{\odot}$ )	$\frac{\text{SFR}_{\text{IR}}^{\text{b}}}{(M_{\odot} \text{ yr}^{-1})}$	$f (AGN)_{bol}^{c}$ %	$f (AGN)_{MIR}^{c}$ %	$L_{\rm UV} \ (10^9 L_{\odot})$	$\frac{\text{SFR}_{\text{UV}}}{(M_{\odot} \text{ yr}^{-1})}$
C1-2127	$110.6\pm10.7$	$9.6 \pm 0.9(^{+2.0}_{-1.5})$	$1049 \pm 101(^{+220}_{-168})$	$7.2 \pm 0.7(^{+0.8}_{-0.8})$	$761 \pm 291(^{+295}_{-294})$	2.7±1.0	~15	5.3	1.7
C1-15182	$78.6\pm7.3$	$10.3\pm1.0$	$1119 \pm 104$	${<}2.6\pm0.24$	${<}259\pm96$	$>9.2\pm3.2$	>50	5.6	1.8
C1-19536	$61.6\pm6.7$	$4.7\pm0.5$	$512\pm56$	$<3.6\pm0.4$	$<384\pm148$	$>4\pm1$	>15	4.8	1.6
C1-21316	$177.1\pm7.8$	$46.0 \pm 1.2$	$4993 \pm 64$	$9.1\pm0.414.9\pm0.7$	$900 \pm 325  1516 \pm 540$	4 - 12	$\sim 35 - 50$	5.4	1.8
A2-15753	$165.7 \pm 4.2$	$13.3 \pm 0.6(^{+2.5}_{-2.1})$	$1357 \pm 34(^{+276}_{-227})$	$<5.3\pm0.1(^{+0.3}_{-0.3})$	$<533\pm188(^{+190}_{-190})$	$> 10 \pm 4$	>40	8.5	2.7

Table 5

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**Notes.** The *Spitzer* 24  $\mu$ m fluxes and derived  $L_{IR}$  and SFR<sub>IR</sub> assuming <sup>a</sup> the mean Dale & Helou (2002) IR template

<sup>b</sup> the best-fit mean Dale et al. (2014) IR template in agreement with *Herschel*  $3\sigma$  detection limits; these values are strictly upper limits to the  $L_{IR}$  and SFRs allowed by the observed FIR fluxes and limits <sup>c</sup> lists the fraction of AGN luminosity contribution to the total bolometric (5–1100 µm) and MIR (5–20 µm) luminosity (values in paranthesis) for the best-fit mean Dale et al. (2014) templates; these values are lower limits constrained by FIR fluxes and limits. For C1-21316, the two values provided in columns 5 to 8 correspond to values derived assuming the two Dale et al. (2014) templates constrained by either the SCUBA 450  $\mu$ m or the  $\lambda > 800 \,\mu$ m fluxes (solid blue curves in Figure 6). The errors listed for  $L_{IR}$  and SFR<sub>IR</sub> were computed using just the 24  $\mu$ m photometric errors (values not in parentheses) and the combination of the 24  $\mu$ m photometric errors and the photometric redshift errors (values in parentheses).

(M82) and a pure power law for the AGN component to the rest-frame mid-IR spectroscopy. FIR color-based templates were created by stacking the SEDs of galaxies that fall within each  $S_{250}/S_{24}$  (observed) versus  $S_8/S_{3.6}$  (observed) color box  $\langle (f_{AGN}) \rangle$  of galaxies that fall into each color box is indicated in the lower right panel of Figure 6). The Kirkpatrick et al. (2015) color-based templates consistent with SPIRE, MIPS, and IRAC fluxes and detection limits of our targets are displayed in violet in Figure 6, scaled to the MIPS 24  $\mu$ m flux of each galaxy. Using the appropriate color-based templates and the listed fraction of  $L_{\rm IR}$  attributable to star formation for each template (Table 3 in Kirkpatrick et al. 2015), we re-calculated the dustenshrouded SFRs. The SFRs calculated accounting for the AGN contribution to the FIR SEDs this way are broadly consistent with the observed UV-8  $\mu$ m SED derived SFRs. Specifically, we calculated the dust-enshrouded SFRs for C1-15182, C1-19536, and A2-15753 to be  $<100 M_{\odot} \text{ yr}^{-1}$ ,  $\sim 100-200 M_{\odot} \text{ yr}^{-1}$  for C1-21316, and  $\sim 200 M_{\odot} \text{ yr}^{-1}$  for C1-2127, in line with SED derived SFRs within uncertainties. It should be noted that the Kirkpatrick et al. (2015) library is more representative of our targets because the templates were created using a large sample of ULIRGs at higher redshifts  $(z \sim 0.3-2.8)$  than compared to Dale et al. (2014) templates, which used local galaxies.

To summarize, the investigation of the FIR SEDs of the massive galaxy sample reveals that none of them can be adequately described as purely star-forming systems. In fact, a MIR AGN contribution of at least  $\sim$ 40%, on average, is necessary to explain the marginal detections in the Herschel bands. It is promising that estimates of AGN contribution using the template sets of Dale et al. (2014) and Kirkpatrick et al. (2015) are consistent with each other. Although in both libraries the fractional AGN contribution is calculated over the rest-frame MIR spectrum, the quasar and starburst templates differ. Dale et al. (2014) uses a single quasar template (median PG quasar spectrum of Shi et al. 2013) to linearly mix with a suite of local normal star-forming galaxies spanning a range in  $\alpha$  to characterize different heating levels ( $\alpha = 1$  and 2.5 for active and quiescent galaxies, respectively) over the 5–20  $\mu$ m wavelength range. In Kirkpatrick et al. (2015), the mid-IR spectrum of the prototypical starburst M82 is used to represent the star-formation component, and the AGN component is characterized by a pure (free-slope) power law.

## 4.4.2. [O III] Luminosities and Bolometric Correction

We use the luminosity-dependent O[III] bolometric correction factor,  $C_{\rm O\,III}$  from Lamastra et al. (2009) for the two galaxies with [O III] emission-line detections (namely C1-19536 and C1-15182),  $L_{\rm bol} \approx 454 \times L_{\rm O\,III}$ . Because both C1-19536 and C1-15182 have non-negligible best-fit  $A_{\rm V}$  values from SED fitting, we assume that the observed [O III] line luminosities are a lower limit to the intrinsic luminosity of [O III]. Assuming all the observed [O III] emission is due to AGN radiation, we find  $L_{\rm bol} > 8 \times 10^{45} \, {\rm erg \, s}^{-1}$  and  $> 3.5 \times 10^{45} \, {\rm erg \, s}^{-1}$  for C1-19536 and C1-15182, respectively, consistent with them hosting powerful hidden AGNs.

# 4.4.3. Continuum Emission from AGNs

The stellar population parameters derived in Section 4.3 may be influenced by the presence of strong AGN continuum contamination to the observed photometry, most significantly biasing the stellar mass estimates. We investigated this by using the publicly available fully Bayesian Markov Chain Monte Carlo SED fitting algorithm of active galaxies, AGNfitter<sup>11</sup> (Calistro Rivera et al. 2016), to model the observed SEDs of the 24  $\mu$ m detected targets. AGNfitter simultaneously fits the total active galaxy SED by decomposing it into four physically motivated components: the stellar population of the host galaxy, cold dust emission from starforming regions, the accretion disk emission (Big Blue Bump, BBB) and hot dust emission surrounding the accretion disk (torus). We used the fiducial library of templates that are supplied with AGNFITTER, which has been tested to perform well in classifying Type1 and Type 2 AGNs purely based on observed broadband photometry. We briefly describe the modeling assumptions (as adopted in Calistro Rivera et al. 2016) and the fiducial library of templates used in AGNFITTER, and highlight important distinctions compared to the modeling assumptions employed in previous sections.

The BBB describing the thermal radiation emitted from the accretion disk surrounding the central supermassive black hole was modeled using a single modified template based on the composite spectrum of Sloan Digial Sky Survey Type 1 QSOs (Richards et al. 2006). Extinction to the emitted BBB spectrum was modeled assuming a Small Magellanic Cloud reddening law (Prevot et al. 1984). Emission from the dusty torus was modeled by using a library of SEDs with varying hydrogen column densities ( $N_{\rm H} = 21-25$ ) created based on the set of empirical templates of Silva et al. (2004). The BC03 stellar population synthesis models are used to construct the library of templates to describe the host galaxy's stellar emission. Although the stellar population synthesis model assumptions are similar to those adopted in FAST as presented in Section 4.3, the explored parameter space is somewhat reduced. Specifically, a similar exponential declining SFH is adopted in AGNFITTER, of the form SFR(t)  $\propto e^{-t/\tau}$ , with a mildly more restricted range for the timescale  $\tau$  from 0.1 to 10 Gyr. A model with constant SFR is also included in AGNFITTER. The grid of stellar population ages is created from 0.2 Gyr to the age of the universe at the target's redshift in step of  $\Delta t \sim 0.1$  dex at fixed solar metallicity. AGNFITTER adopts a Chabrier (2003) and the Calzetti et al. (2000) reddening law to model the extinction to the stellar light. The library of templates for modeling the emission from cold dust in star-forming regions includes the template set of Dale & Helou (2002) (also used in Section 4.4.1) and Chary & Elbaz (2001). For a more detailed description of the fitting algorithm, modeling assumptions, and template libraries, we refer the reader to Calistro Rivera et al. (2016).

Figure 7 illustrates the results when modeling the observed SEDs with the full parameter space range of the fiducial templates of AGNFITTER and Table 6 lists the best-fit stellar masses converted to the Kroupa (2001). The fractions of AGN emission in the rest-frame UV–optical (0.1–1  $\mu$ m) and MIR (1–30  $\mu$ m) for each object, calculated by comparing the relative contribution of AGN luminosities in the considered wavelenghts are also listed in Table 6. The amount of AGN contribution to the rest-frame MIR estimated by AGNFITTER is remarkably consistent with the values or lower limits derived in Section 4.4.1. The AGN contribution in the rest-frame UV–optical is found to range between a few percent (for A2-15753)

<sup>&</sup>lt;sup>11</sup> https://github.com/GabrielaCR/AGNfitter



Figure 7. SEDs output when modeling the observed UV-FIR photometry of the  $24 \mu m$  detected targets with four components using AGNFITTER (Calistro Rivera et al. 2016). The orange and green curves correspond to the stellar emission and starburst components, respectively. The accretion disk and the surrounding hot dusty torus components are represented as blue and purple curves, respectively. In each panel, the SED components corresponding to 10 randomly selected realizations from the posterior probability density functions are overplotted in order to visualize the range of parameter space. The total SEDs from the realizations are plotted in red.

 Table 6

 Best-fit Stellar Mass and AGN Contributions with AGNfitter

ID	$\log M_*$	$\Delta \log M_*$	$f_{\rm AGN}$ (0.1–1 $\mu {\rm m}$ )	$f_{\rm AGN}~(1{-}30~\mu{\rm m})$
C1-2127	$11.24_{-0.24}^{+0.17}$	$-0.11^{+0.17}_{-0.25}$	$0.37\substack{+0.07\\-0.07}$	$0.09^{+0.18}_{-0.05}$
C1-15182	$11.61\substack{+0.05\\-0.09}$	$+0.05\substack{+0.17\\-0.13}$	$0.09\substack{+0.12\\-0.07}$	$0.96\substack{+0.04\\-0.17}$
C1-19536	$11.37\substack{+0.03\\-0.04}$	$-0.08\substack{+0.10\\-0.10}$	$0.10\substack{+0.10\\-0.07}$	$0.98\substack{+0.02\\-0.21}$
C1-21316	$11.34_{-0.68}^{+0.23}$	$-0.24^{+0.57}_{-0.68}$	$0.28^{+0.35}_{-0.23}$	$0.39_{-0.05}^{+0.04}$
A2-15753	$11.21\substack{+0.01 \\ -0.03}$	$+0.04\substack{+0.15\\-0.13}$	$0.04\substack{+0.03\\-0.02}$	$0.99\substack{+0.01\\-0.01}$

**Note.** Results obtained from the modeling of the binned UV–NIR spectra, NMBS UV-8  $\mu$ m photometry, and FIR detections and upper limits with AGNFITTER. The best-fit stellar masses were adjusted to the Kroupa (2001) IMF in order to compare with results listed in Table 3.  $\Delta M_*$  lists the effect on the derived stellar mass when including AGN templates in the modeling of the panchromatic SEDs ( $\Delta \log M_* = \log (M_*)_{(AGNfitter)} - \log (M_*)_{(FAST)}$ ) at fixed metallicity ( $Z = Z_{\odot}$ ).  $f_{AGN}$  is calculated as the relative contribution of the AGN templates in the considered wavelengths (i.e., the Big Blue Bump emission describing the accretion disk at 0.1–1  $\mu$ m and the hot dusty torus emission at 1–30  $\mu$ m).

to  $\sim$ 40% (for C1-2127), whereas the AGN contribution in the rest-frame MIR is found to range between  $\sim$ 10% (for C1-2127) to almost 100% (for C1-15182, C1-19536, and C1-21316).

The third column of Table 6 lists the difference between the stellar mass estimated by AGNFITTER and the stellar mass derived using FAST (i.e., assuming no AGN contamination to the observed SEDs). It is interesting to note that, while the naive expectation would be that not accounting for AGN contamination to the observed photometry would lead to overestimating the stellar masses, two galaxies (C1-15182 and A2-15753) have AGNFITTER derived stellar masses marginally

larger (by  $\sim 0.04-0.05$  dex) than when no AGN contamination is assumed. This appears to be caused by slightly older best-fit ages of the stellar population (and hence slightly larger massto-light ratios) when the SEDs are modeled accounting for the AGN contribution. We note, however, that these differences are smaller than (or comparable to) the  $1\sigma$  error on the stellar mass. For the other three galaxies, the stellar masses derived by accounting for the AGN contamination are found to be smaller (by  $\sim$ 0.1–0.3 dex on average) than the stellar masses derived when no AGN contamination is assumed. The best-fit stellar masses derived assuming AGN contamination ranges from  $1.5 \times 10^{11} M_{\odot}$  to  $4 \times 10^{11} M_{\odot}$ . All galaxies except C1-21316 are found to be more massive than  $10^{11} M_{\odot}$  within  $1\sigma$  when modeled by AGNFITTER. For all galaxies but C1-21316, the systematic effect to the derived stellar masses from AGN contamination is found to be smaller than a factor of  $\sim 2$  at  $1\sigma$ . This is because the rest-frame UV-optical SEDs of these galaxies are dominated by the stellar light despite the rest-frame MIR SEDs being dominated by the obscured AGN. For C1-21316, the stellar mass is found to be overestimated by a factor of  $\sim 2$  on average, although the  $1\sigma$  uncertainty allows for solutions that are smaller by as much as a factor of  $\sim$ 8. Finally, we note that the stellar masses of all studied galaxies derived from AGNFITTER are consistent with the stellar masses derived using FAST within the  $1\sigma$  uncertainties.

## 4.4.4. X-Ray and Radio Detections

As previously done in Marchesini et al. (2010), we used the publicly available *Chandra* X-ray catalogs (Laird et al. 2009 and Elvis et al. 2009 for the AEGIS and COSMOS fields, respectively)

and the improved redshifts to revise the X-ray luminosities of the three detected sources (C1-15182, C1-19536, and A2-15753). Assuming a power-law photon index  $\Gamma = 1.9$  (Nandra & Pounds 1994), the X-ray luminosities  $L_{2-10 \text{ keV}}$  are  $(6.4 \pm 1.7) \times 10^{45}$  for C1-15182,  $(16.6 \pm 2.6) \times 10^{45}$  for C1-19536 and  $(9.6 \pm 1.5) \times 10^{45}$  erg s<sup>-1</sup> for A2-15753 (redshift uncertainties are included in quoted errors). We used the empirically derived observed  $L_{\rm [O III]} - L_{2-10 \text{ keV}}$  relation of local AGNs from Ueda et al. (2015) to investigate the expected X-ray flux of the sources with detected [O III] emission (C1-19536 and C1-15182). The estimated  $L_{2-10 \text{ keV}}$  for these galaxies using the observed  $L_{[O \text{ III}]}$  are an order of magnitude less than the luminosities calculated using X-ray detections (using  $\Gamma = 1.4$  lowers this factor to  $\sim 6-8$ ). We calculated upper limits to the rest-frame  $L_{2-10 \text{ keV}}$  for the targets not detected (C1-2127, C1-19764, and C1-21316) using the  $3\sigma$ detection limits with  $\Gamma = 1.9$ . The  $3\sigma$  upper limits are in the range of  $L_{2-10 \text{ keV}} \sim 5-20 \times 10^{44} \text{ erg s}^{-1}$ , comparable to only the very powerful X-ray AGN, and therefore we cannot conclude or rule out whether these galaxies harbor AGNs.

From the sample of galaxies presented here, only C1-21316 (AzTEC-5) is detected at 1.4 GHz (Schinnerer et al. 2010), with a flux of  $0.126 \pm 0.015$  mJy. We calculated the rest-frame flux densities assuming the canonical value of the radio spectral index of  $\alpha = 0.8$  (Condon 1992). The calculated rest-frame flux density is  $L_{1.4 \text{ GHz}} = 10.6 \pm 1.3 \times 10^{27} \text{ W Hz}^{-1}$ , above the threshold to be classified as a radio-loud AGN (log( $L_{1.4}$  GHz) > 25; Schinnerer et al. 2007). The radio spectral index is not observed to evolve significantly over cosmic time (Magnelli et al. 2015), and is consistent with values derived for different samples of galaxies (Ibar et al. 2009, 2010; Ivison et al. 2010; Bourne et al. 2011).

#### 4.4.5. Redshift Evolution of the AGN Fraction

The fraction of AGNs  $(f_{AGN})$  within a galaxy population represents a probe of the level of AGN activity and of the growth of supermassive black holes at the center of galaxies. The energetic output from AGNs is considered to be one of the dominant feedback processes that can regulate, and even halt altogether, the infalling of gas and star formation in massive galaxies. The measurement of the level of AGN activity in massive galaxies as a function of redshift is therefore an indirect probe of the impact of AGN feedback at a given cosmic time.

Figure 8 shows the evolution of  $f_{AGN}$  as a function of redshift for galaxies with stellar mass  $M_{\star} > 10^{11} M_{\odot}$ . In our sample of very massive galaxies at 3 < z < 4, AGNs appear almost ubiquitous, with a very large AGN fraction of  $f_{\rm AGN} pprox 0.8 \pm 0.2$  (in which we assumed that C1-2127 does not host an AGN). Figure 8 also shows the fraction of AGN in massive galaxies at 0.6 < z < 3.0 from Cowley et al. (2016), at 1.7 < z < 2.7 from Kriek et al. (2007), and in the local universe from Kauffmann et al. (2003). Our measurement is consistent with the very high AGN fraction of  $f_{AGN} = 0.86^{+0.14}_{-0.22}$  at 2.6 < z < 3.0 as measured by Cowley et al. (2016) using the zFOURGE data set. On the contrary, the AGN fraction in massive galaxies at  $z \sim 0$  is much smaller,  $f_{\rm AGN} \approx 0.3$ , and it appears to remain relatively constant all the way out to  $z \sim 2.5$ . Figure 8 clearly shows a dramatic transition in the fraction of AGNs in galaxies with  $M_{\star} > 10^{11} M_{\odot}$  happening at  $z \sim 2.5$ , i.e., when the universe was  $\sim 3$  Gyr old. We quantified the significance of the transition in the AGN fraction at  $z \sim 2.5$  by assuming that



**Figure 8.** Evolution of the fraction of AGNs in galaxies with stellar mass  $M_{\star} > 10^{11} M_{\odot}$ . The red star represents the measurement from our work. The green, pink, and gray points represent the fraction of AGNs at  $z \sim 0$  (SDSS; Kauffmann et al. 2003), at 1.7 < z < 2.7 from Kriek et al. (2007), and at 0.6 < z < 3.0 from Cowley et al. (2016). The dotted–dashed curve represents a fit to all plotted points using the empirical model  $f_{\text{AGN}} = \frac{1}{C + e^{-(z-A)/B}} + D$  with the  $1\sigma$  uncertainties on the model fit indicated by the light blue hatched region.

the fraction of AGN is constant with redshift, and equal to the bi-weight mean of the observed fraction at z < 2.5 (i.e.,  $f_{\rm AGN} \approx 0.29$ ) and testing this hypothesis with the chi-squared statistics of the two points at z > 2.5. If each point at z > 2.5 is considered by itself, we find a probability p < 0.0006 of obtaining each z > 2.5 measurement at least this discrepant from the no evolution assumption. If the two z > 2.5measurements of the AGN fraction are considered together, we find a probability  $p < 10^{-7}$  of obtaining both z > 2.5points at least this discrepant from the no evolution assumption. At earlier times, most (all) supermassive black holes at the center of massive galaxies appear to be actively accreting, whereas at later times most ( $\sim 2/3$ ) of them are found dormant. Although the small sample does not allow us to robustly quantify a potential enhancement of the AGN fraction in our spectroscopically confirmed sample of very massive galaxies at z > 3 due to [O III] emission-line contamination to the stellar continuum. However, we note that the [O III] line emissions detected in the spectra are found to contribute at most 15%, and more generally  $\leq 5\%$ , to the observed K-band fluxes. The limited observed [O III] emission-line contamination suggests that our sample is not significantly biased toward strong [O III] emitters, i.e., toward massive galaxies hosting AGNs, and that the derived AGN fraction in our sample is not considerably enhanced. Future wide area surveys will be able to constrain more robustly the fraction of AGN in massive galaxies at z > 2.5, whereas deeper surveys will be able to investigate the evolution of the AGN fraction in lower-mass galaxies.

## 5. Summary and Discussion

In this paper, we investigated the observed-frame UV and NIR spectra of a sample of six  $3 \leq z_{\text{phot}} < 4$  massive galaxies selected from the NMBS. We confirmed the spectroscopic redshift of two galaxies (C1-15182,  $z_{\text{spec}} = 3.371$ ; C1-19536,  $z_{\text{spec}} = 3.139$ ) through nebular emission lines (Figure 1). We re-modeled the medium- and broadband photometry in combination with the binned one-dimensional spectra of detected sources to derive improved photometric redshifts with the EAZY code ( $z_{\text{cont}}$ ). The best-fit  $z_{\text{cont}}$  of only one of our

targets (C1-19764, not included in Marchesini et al. 2010 sample to calculate the high-mass end SMF at  $3 \le z_{\text{phot}} < 4$ ) is z < 3 (although the redshift solution extends to z > 3).

We used FAST in conjunction with the BC03 stellar population synthesis models, the Calzetti et al. (2000) extinction law, the Kroupa (2001) IMF, and an exponentially declining SFH to model the observed SEDs (Figure 4, Table 3). We included an additional source in our analysis from the Marchesini et al. (2010) sample, C1-21316, which has a spectroscopic redshift of  $z_{\text{spec}} = 3.971$ . We find that the updated stellar population parameters are consistent with those previously derived using only medium- and broadband photometry. The stellar masses of galaxies with detected optical emission lines decrease by  $\sim 0.1 \, \text{dex}$  when emission-line contamination to the observed photometry is accounted for. From the SED modeling, our sample of massive galaxies show ranges in stellar population properties, in accordance with previous studies at  $3 \le z < 4$ (Muzzin et al. 2013b; Marchesini et al. 2014; Spitler et al. 2014). We find that  $\sim$  50% of sources in this sample are characterized by average stellar ages of  $\approx$ 800 Myr, suggesting that the bulk of the stellar mass in these systems was formed early in the universe  $(z_{\rm form} \sim 4.5)$  (these also have ongoing star formation, with SFR  $\sim \text{few} \times 10\text{--}100 \, M_{\odot} \, \text{yr}^{-1}$ , but these values are highly dependent on model parameters).

All but one of the sources in this sample have significant MIPS 24  $\mu$ m detections. The total IR luminosities estimated from the observed 24  $\mu$ m assuming IR templates representative of starburst galaxies are  $\sim 2-3$  orders of magnitude greater than the SFRs estimated from SED modeling. Inspecting the observed FIR SEDs including Herschel-PACS and SPIRE detections and upper limits effectively rules out dustenshrouded star formation as the only source of the observed 24  $\mu$ m flux. In fact, using two separate IR libraries that include different levels of contribution from an AGN (Dale et al. 2014; Kirkpatrick et al. 2015) reveals that >15% of the MIR emission must originate from obscured AGNs. However, despite the significant AGN emission contribution to the MIR SEDs, we find that the rest-frame UV-optical light is largely dominated by stellar emission. For all but one galaxy, the presence of the AGN does not considerably bias the stellar mass estimates, with typical systematic differences within  $\sim 0.1$  dex, and stellar masses at most a factor of  $\sim 2$  smaller at the  $1\sigma$  level. Even after accounting for the presence of the (mostly obscured) AGN, the inferred stellar masses of all but one galaxy are found to be larger than  $10^{11} M_{\odot}$  within the  $1\sigma$  uncertainties. The large fraction (>80%) of massive galaxies hosting AGNs at 3 < z < 4 inferred from strong [O III] emission lines, IR SED properties, and X-ray and radio detections implies that 3 < z < 4 must be an extremely active time in the cosmic history for the growth of the supermassive black holes at the massive end of the galaxy population.

The authors thank the anonymous referee for useful comments and suggestions, which have improved this manuscript. Z.C.M. gratefully acknowledges support from the John F. Burlingame and the Kathryn McCarthy Graduate Fellowships in Physics at Tufts University. D.M. and Z.C.M. acknowledge the support of the Research Corporation for Science Advancement's Cottrell Scholarship, the support from Tufts University Mellon Research Fellowship in Arts and Sciences. We thank Gabriela Calistro Rivera for her support on the implementation of AGNFITTER. This study is partially

based on data products from observations made with ESO Telescopes at the La Silla Paranal Observatory under program ID 087.A-0514. Some of the data presented in this study were obtained at the W. M. Keck Observatory, which is operated as a scientific partnership among the California Institute of Technology, the University of California, and the National Aeronautics and Space Administration. The Observatory was made possible by the generous financial support of the W. M. Keck Foundation. Keck telescope time was granted by NOAO, through the Telescope System Instrumentation Program (TSIP). TSIP is funded by NSF. The authors wish to recognize and acknowledge the very significant cultural role and reverence that the summit of Mauna Kea has always had within the indigenous Hawaiian community. We are most fortunate to have the opportunity to conduct observations from this mountain. This study makes use of data from the NEWFIRM Medium-Band Survey, a multi-wavelength survey conducted with the NEWFIRM instrument at the KPNO, supported in part by the NSF and NASA. We thank the NMBS and COSMOS collaborations for making their catalogs publicly available. The combination of NEWFIRM Medium-Band Survey and Herschel data was supported by the National Aeronautics and Space Administration (NASA) under grant number NNX13AH38G issued through the NNH12ZDA001N Astrophysics Data Analysis Program (ADAP).

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