
Collaborative Research: Actively Controllable Microfluidics with Film-Confined Redox-Magnetohydrodynamics -- Video and Data

02 Jun 2016

Chaotic Advection-Driven Mixing in Unsteady Three-Dimensional MHD Flows in Microfluidic Devices

Fangping Yuan

Kakkattukuzhy M. Isaac

Missouri University of Science and Technology, isaac@mst.edu

Follow this and additional works at: <https://scholarsmine.mst.edu/cr-acm>

 Part of the [Aerospace Engineering Commons](#), and the [Fluid Dynamics Commons](#)

Recommended Citation

Yuan, Fangping and Isaac, Kakkattukuzhy M., "Chaotic Advection-Driven Mixing in Unsteady Three-Dimensional MHD Flows in Microfluidic Devices" (2016). *Collaborative Research: Actively Controllable Microfluidics with Film-Confined Redox-Magnetohydrodynamics -- Video and Data*. 5.
<https://scholarsmine.mst.edu/cr-acm/5>

This Presentation is brought to you for free and open access by Scholars' Mine. It has been accepted for inclusion in Collaborative Research: Actively Controllable Microfluidics with Film-Confined Redox-Magnetohydrodynamics -- Video and Data by an authorized administrator of Scholars' Mine. This work is protected by U. S. Copyright Law. Unauthorized use including reproduction for redistribution requires the permission of the copyright holder. For more information, please contact scholarsmine@mst.edu.

Chaotic Advection-Driven Mixing in Unsteady Three-Dimensional MHD Flows in Microfluidic Devices

Fangping Yuan
K. M. Isaac

Missouri University of Science & Technology
Aerospace Engineering

229th ECS Meeting, San Diego, May 29-Jun 2, 2016

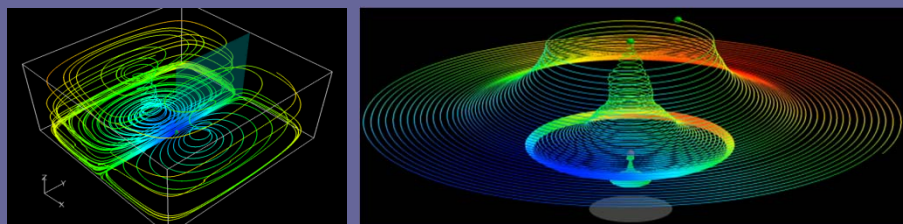
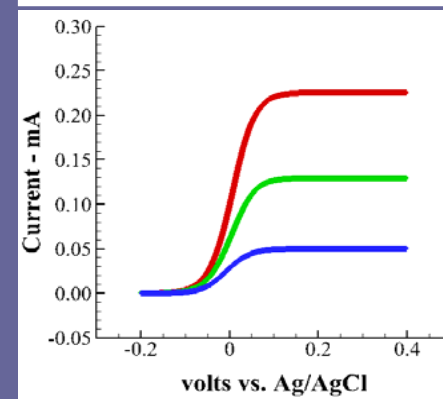
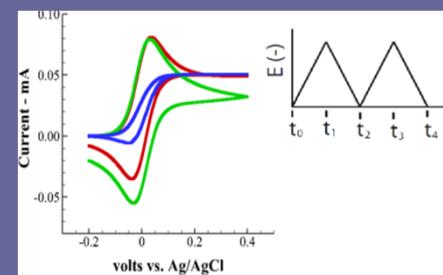
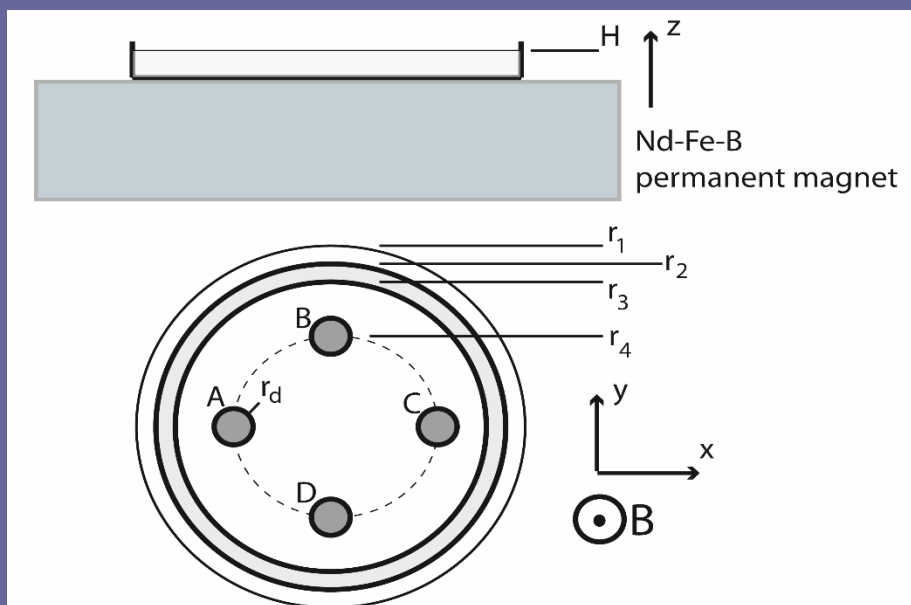


Acknowledgements

National Science Foundation Grants

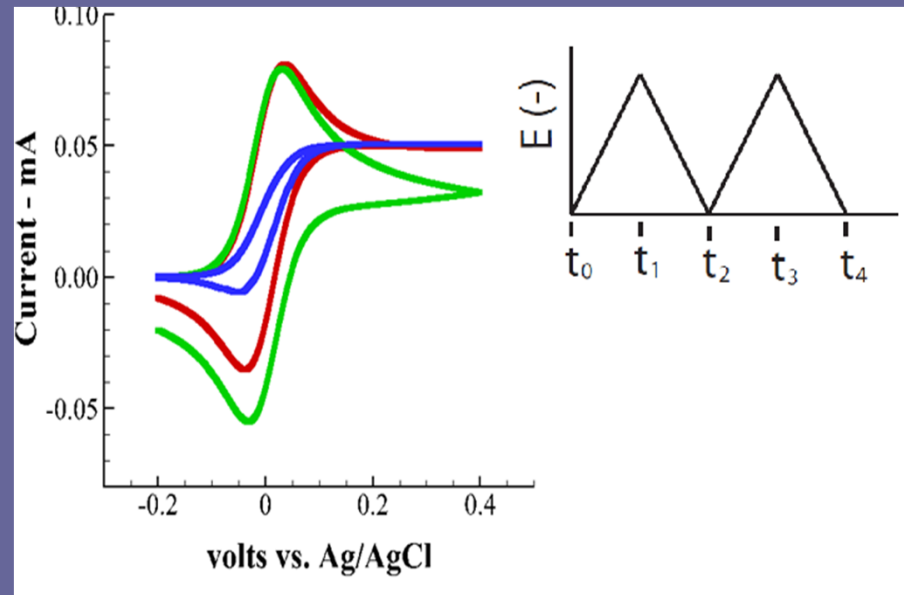
Award number CBET-1336722

Award number CHE-0719097



Outline

- Motivation
- Problem Description
- Technical Highlights
- Results
- Discussion
- Conclusions
- Ongoing and Future Work



Motivation

- Speed of chemical reactions is dictated by mixing efficiency
- Low Reynolds number (~ 1) makes mixing a challenge
- Introducing Turbulent-like features can enhance mixing
- Chaotic advection provides foundation for analysis and control
- Simulations based on Navier-Stokes equations provide many insights
- Powerful post-processing tools help understand underlying mixing phenomena
- Simulation results can help develop new mixing and flow control strategies



Problem Description

- Stirring vs. mixing: An important distinction
- Stirring increases the mean value of gradients
- Mixing decrease the mean value of gradients
- Stirring leads to stretching and folding of the interface
- Chaotic advection is an efficient way of stirring
- Mixing is complete when the gradients disappear

Problem Description

- MHD can introduce a non-intrusive driving force
- Lorentz force $\bar{j} \times \bar{B}$ for MHD is due to ionic current-magnetic field cross product
- Flow field and species concentration fields are obtained by solving unsteady, 3D Navier-Stokes equations
- Extensive post-processing is performed to visualize the results and obtain mixing quality
- Passive numerical particles are tracked by integrating the advection equation

$$\frac{d\bar{X}}{dt} = \bar{V}(\bar{X}, t)$$



Problem Description

Cell geometry

Shallow open cell

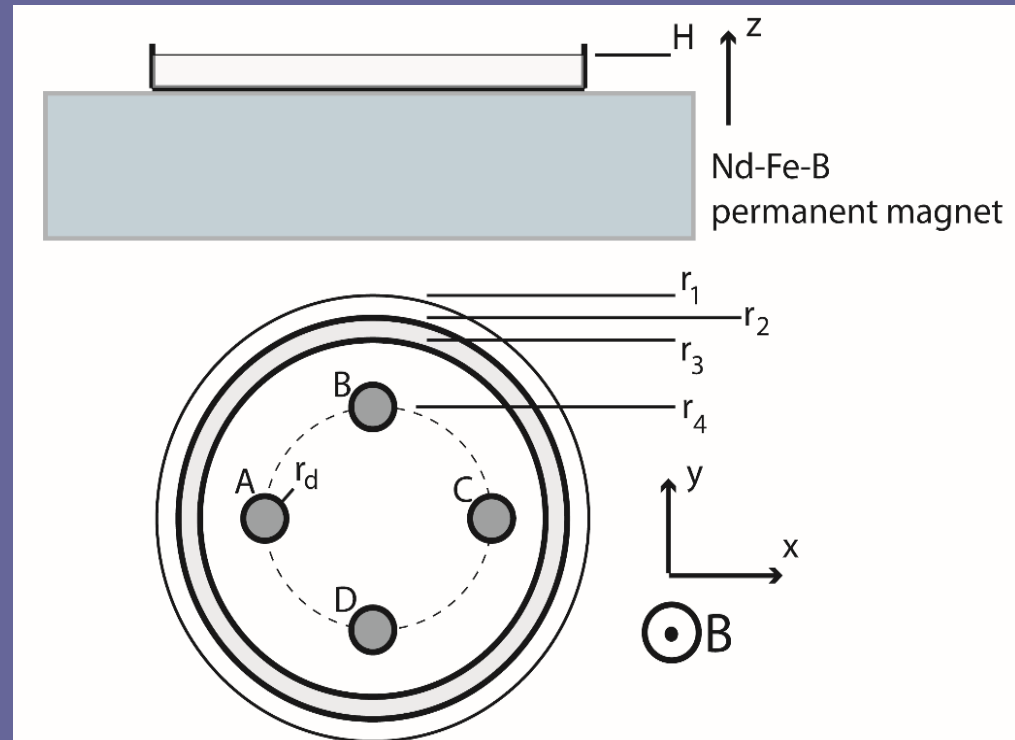
$$r_1 = 3\text{mm}$$

$$r_2 = 2.4\text{ mm}$$

$$r_3 = 2\text{mm}$$

$$r_d = 0.16\text{ mm}$$

$$H = 0.5\text{ mm}$$



NaCl solution: $\rho = 1000\text{ kg/m}^3$, $\mu = 0.001\text{ kg/(m.s)}$, $C^* = 0.1\text{M}$,
 $\phi = 0.04\text{ V}$, $\sigma = 1.29\text{ S/m}$.

Magnetic field: $B_z = 0.36\text{ Tesla}$



Problem Description

Governing Equations

Continuity: $\nabla \cdot \mathbf{V} = 0$

Momentum: $\rho \frac{D\mathbf{V}}{Dt} = -\nabla p + \mu \nabla^2 \mathbf{V} + \mathbf{F}_L$

Lorenz Force: $\mathbf{F}_L = -\sigma \nabla \phi \times \mathbf{B}$

Electric Field Outside Double Layer:

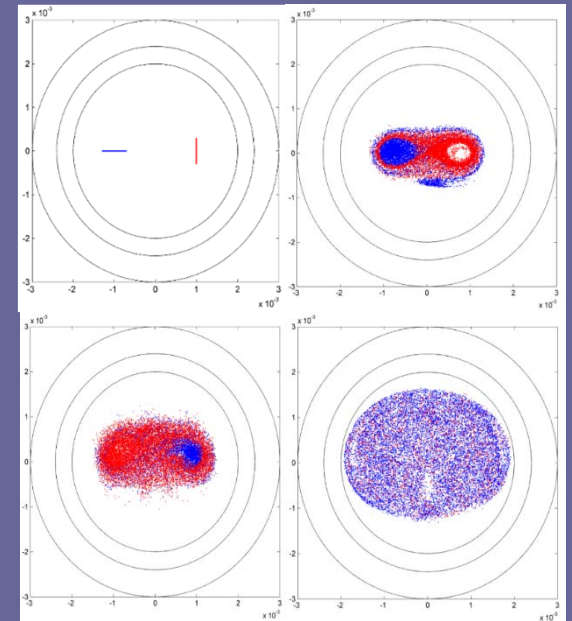
$$\nabla^2 \phi = 0$$

Auxiliary Equations

Passive Particle Trajectory: $\frac{d\mathbf{X}(t)}{dt} = \mathbf{V}(\mathbf{X}, t)$

Mixing Performance: $\alpha(t) = 1 - \frac{\delta^2(t)}{\delta^2(0)}$

$$\delta^2(t) = \iiint_V [C(x, y, z, t) - \bar{C}]^2 dx dy dz$$



2nd order, implicit

$$\frac{3\phi^{n+1} - 4\phi^n + \phi^{n-1}}{2\Delta t} = F(\phi^{n+1})$$



Problem Description

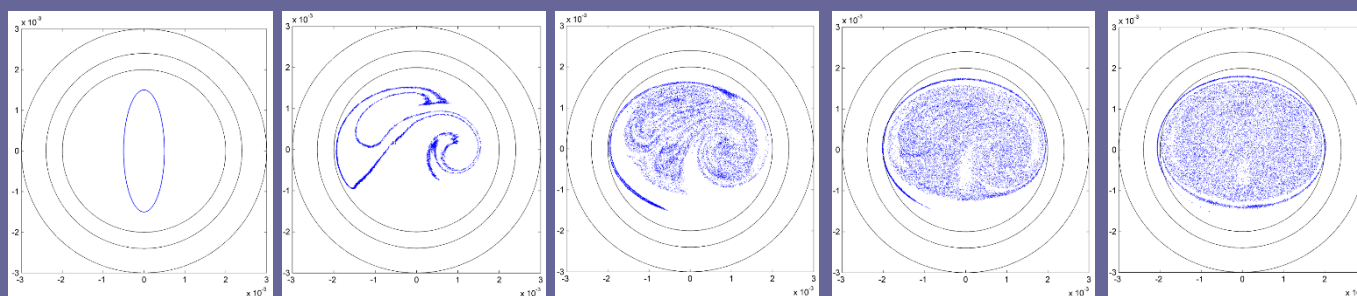
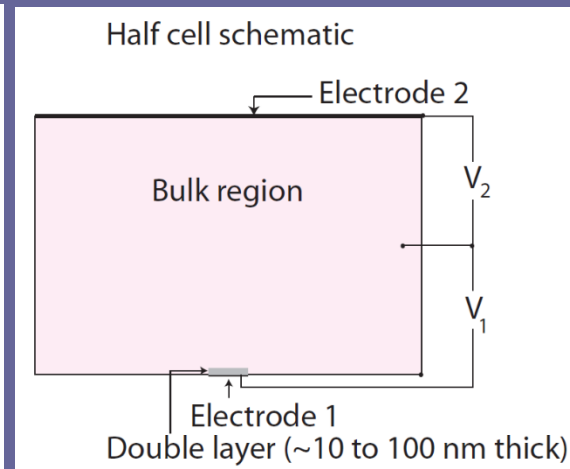
Gouy-Chapman Model for Double Layer

Boltzmann distribution: $C_i = C_i^* \exp\left(-\frac{z_i F \phi}{RT}\right)$

Electric field: $\nabla^2 \phi = -\frac{\rho_e}{\epsilon} = -\frac{F}{\epsilon} \sum_i z_i C_i^* \exp\left(-\frac{z_i F \phi}{RT}\right)$

1D solution for DL: $\frac{\tanh(K\phi)}{\tanh(K\phi_0)} = \exp(-\kappa x)$

$K = Fz / (4RT), \quad \kappa = \left(2F^2 / \epsilon RT\right)^{1/2} z \sqrt{C^*}$



$T = 10s$
 $t = 1T, 5T,$
 $10T, 20T$

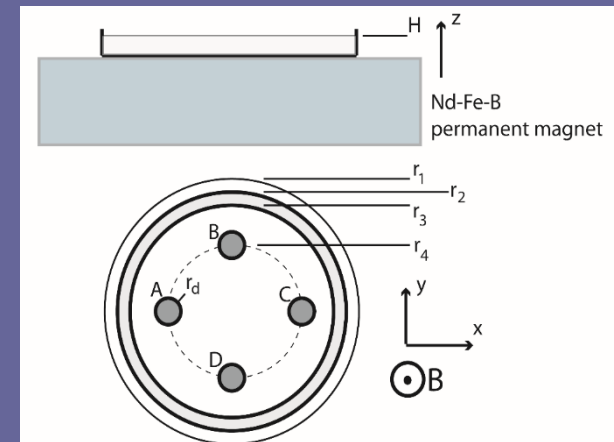


Problem Description

Switching Scheme 1

$$\phi_A = 0.04V \quad \phi_C = 0 \quad kT < t < kT + \frac{T}{2}$$

$$\phi_A = 0 \quad \phi_C = 0.04V \quad kT + \frac{T}{2} < t < (k+1)T$$



Switching Scheme 2

$$\phi_A = 0.04V \quad \phi_C = -0.04V \quad \nabla\phi_B = \nabla\phi_D = 0 \quad kT < t < kT + \frac{T}{2}$$

$$\nabla\phi_A = \nabla\phi_C = 0 \quad \phi_B = 0.04V \quad \phi_D = -0.04V \quad kT + \frac{T}{2} < t < (k+1)T$$

$$k = 0, 1, 2, 3, \dots$$



Results

- **Scheme 1, Scheme 2**

- Velocity vectors
- Velocity profiles
- Potential contours
- Current lines
- Poincaré maps
- Mixing performance

- **Scheme 2**

- ...+
- Particle concentration maps
- Species concentration contours
- Material line deformation

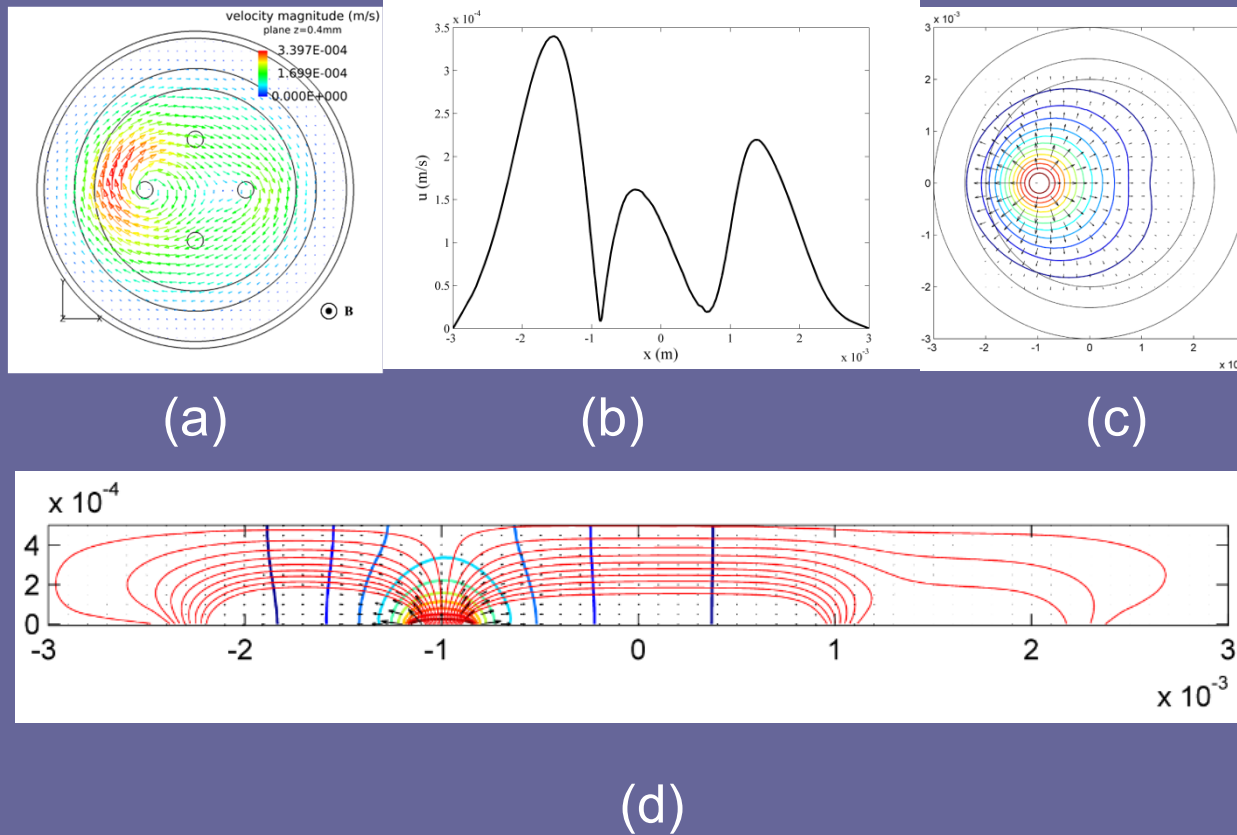
- **Observations**



Results

Scheme 1

- (a) Velocity vectors
- (b) Velocity profile
- (c), (d) Potential contours and current flux lines
(c: $z = 0.4$ mm plane, d: $y = 0$ plane)

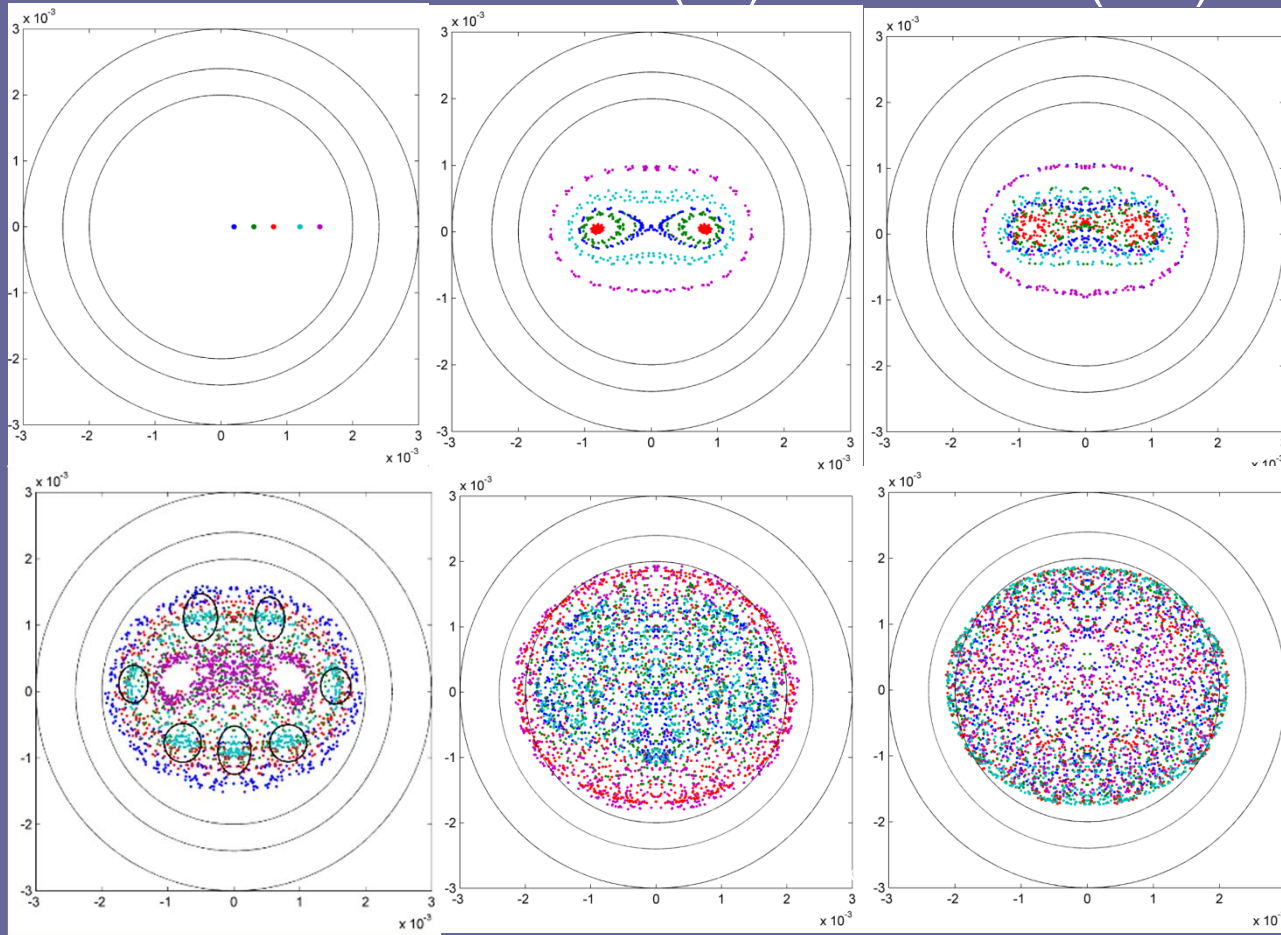


Results

$t = 0$

$T = 1\text{s} (75)$

$T = 2\text{s} (100)$



$T = 4\text{s} (300)$

$T = 8\text{s} (400)$

$T = 10\text{s} (400)$

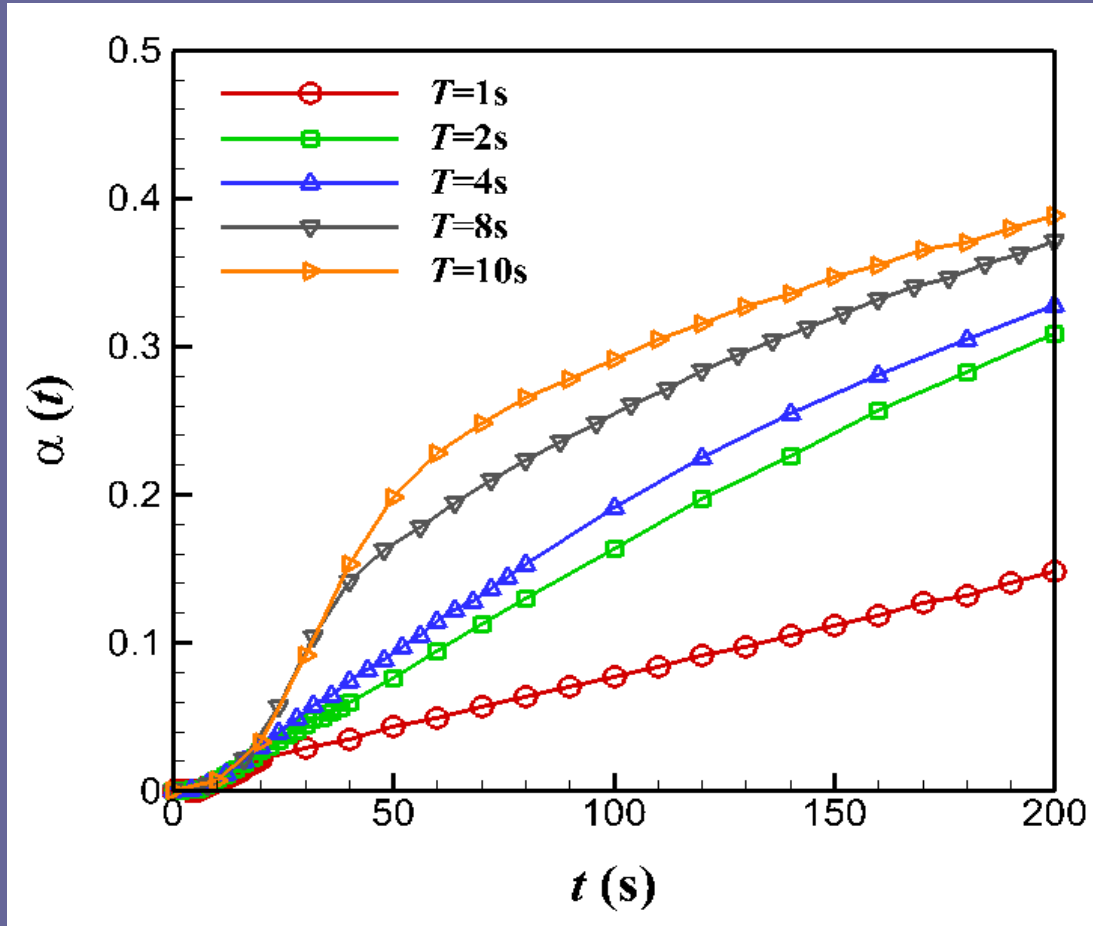
Scheme 1

Poincaré maps
 $z = 0.4\text{ mm}$

Number of
periods are
shown in
parentheses



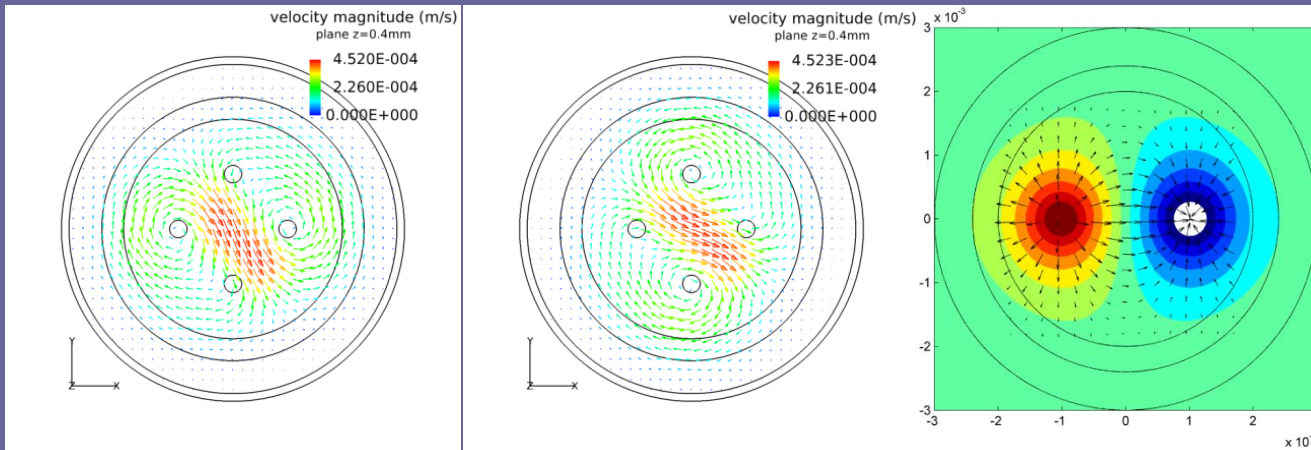
Results



Scheme 1

Evolution of mixing quality vs. time t for 5 values of T in plane $z = 0.4$ mm

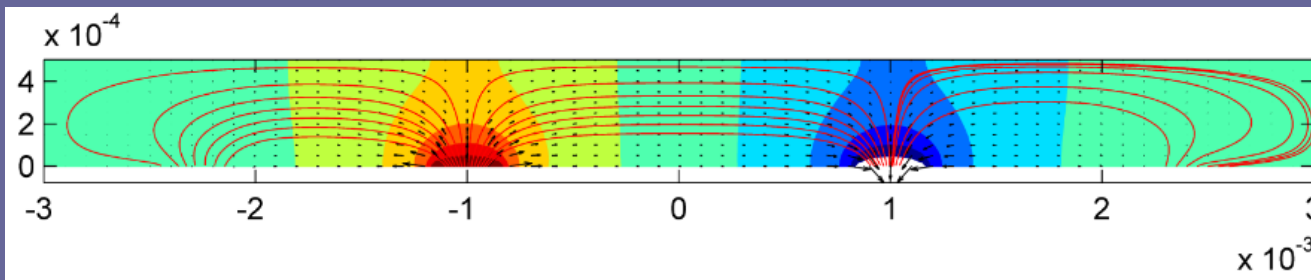
Results



(a)

(b)

(c)

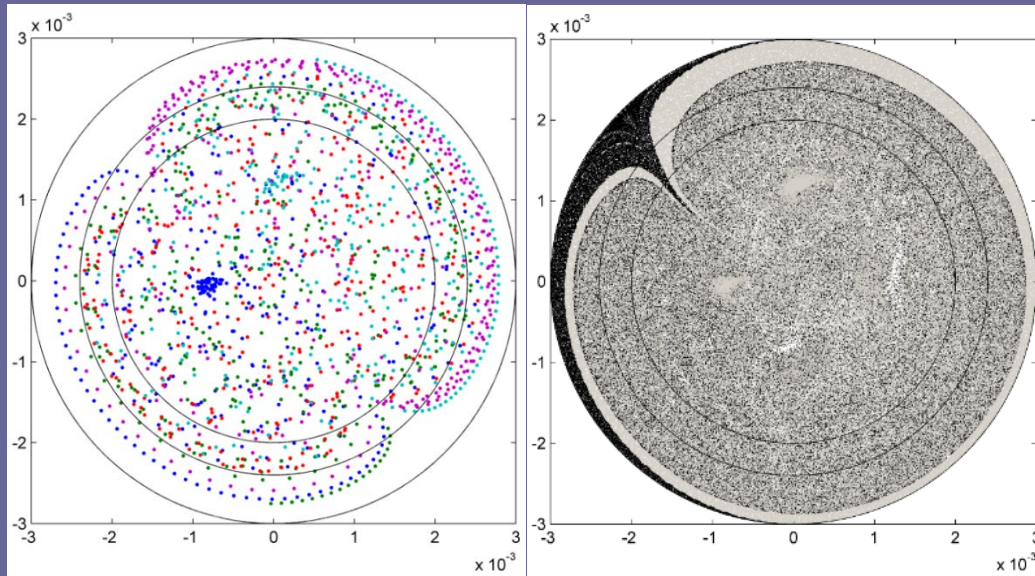


(d)

Scheme 2

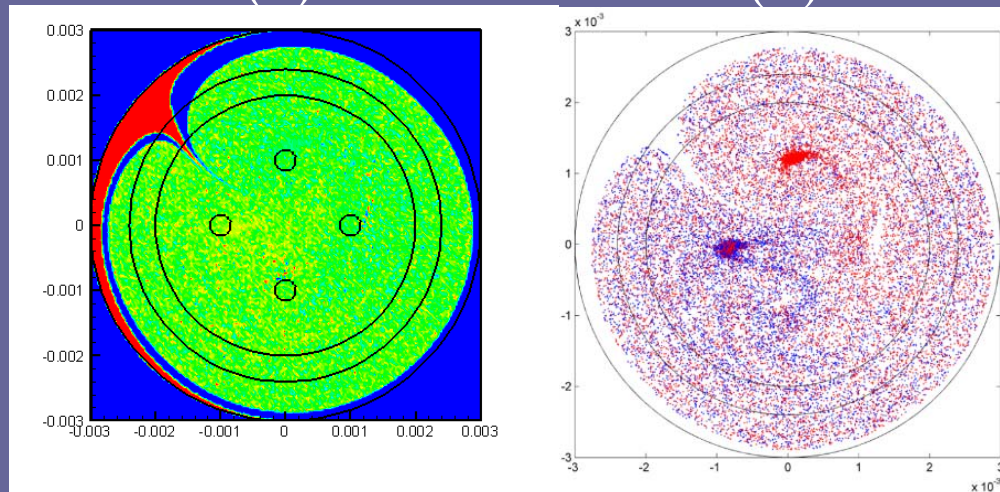
- (a) Velocity vectors
- (b) Velocity vectors
- (c), (d) Potential contours and current flux vectors and lines
(c: $z = 0.4\text{ mm}$ plane, d: $y = 0$ plane)

Results



(a)

(b)



(c)

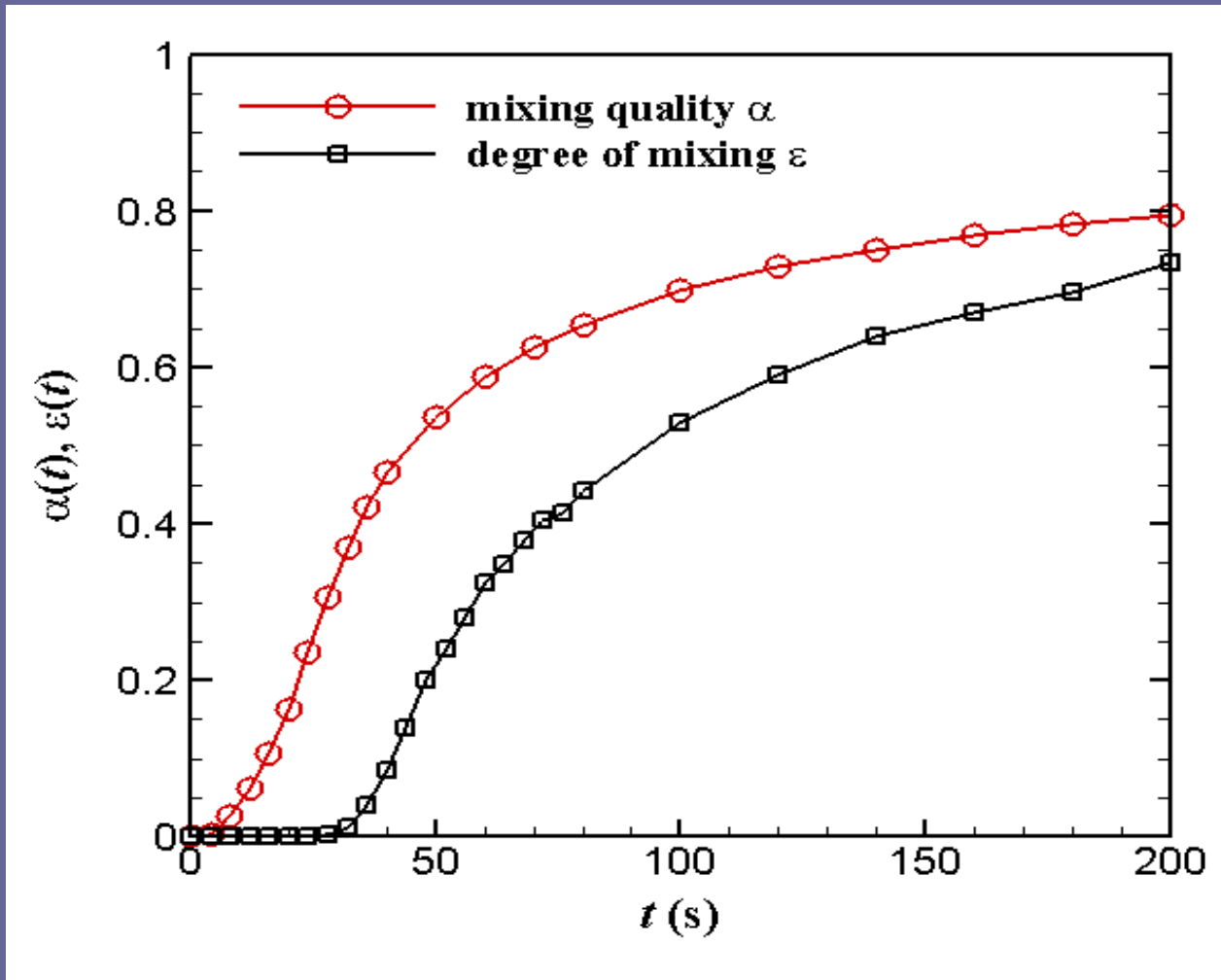
(d)

Scheme 2.

$T = 4s$, $t = 200s$ $z = 0.4$ mm plane.

- (a) Poincaré map
- (b) particle concentration map
- (c) species concentration contours
- (d) Stretching and deformation of two material lines.

Results



Scheme 2
Temporal evolution
of mixing quality
 α and degree of
mixing ϵ . $T = 4$ s.

Results

- Switching Scheme 1 is a practical way to create the blinking vortex model
- Scheme 1 results are similar to those from the Stokes flow model
- Results from particles and material lines show that α increases with period T
- Switching Scheme 2 improves α over Switching Scheme 1
- The islands disappear faster in Scheme 2 compared to Scheme 1
- Each pair of disks is insulated during part of the cycle. *No current on the disk, no vortical flow around the disk to isolate it*



Conclusions and Future Work

- Even better mixing may be possible with more number of electrodes
- Optimal geometric placement and judicious choice of switching schemes may further improve mixing.
- Islands disappear faster with the electrodes insulated during part of the cycle
- The Navier-Stokes simulations show differences compared to the Stokes flow model for $Re_H \gg 1$
- Good post-processing tools are necessary for analysis



Conclusions and Future Work

- Chaotic advection increases with increase with the dimension of the problem (no chaotic advection in 2D, and increases in the order: unsteady 2D and steady 3D, unsteady 3D)
- Length scales and time scales have an influence on chaotic advection
- CFD simulations advance the modeling over potential flow and stokes flow
- Work on Transmission Line Equivalent Circuit (TLEC) model is in progress