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Thermal noise in the highly-reflective mirror coatings is one of the main limitations to the sensitivity of current gravitational-wave detectors. Amorphous silicon (aSi) is an ideal material to reduce thermal noise. Due to high optical absorption at 1064 nm, so far it was mainly considered as a candidate for future, cryogenic detectors using longer wavelengths. This paper summarizes the current state-of-the-art of the optical absorption of aSi at 1064 nm. We show how recent improvements in aSi absorption, and the development of multimaterial coatings, make the use of aSi at 1064 nm realistic, and discuss the possible thermal-noise improvement and corresponding optical absorption in room-temperature gravitational-wave detectors.

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## I. INTRODUCTION

During their first two observing runs, Advanced <sup>53</sup> LIGO [1] and Advanced Virgo [2] detected several <sup>54</sup> gravitational-wave signals from binary black-hole merg- <sup>55</sup> ers and from a neutron-star merger [3–8]. These <sup>56</sup> gravitational-wave detectors are kilometer-scale Michel- <sup>57</sup> son interferometers measuring changes in the separation <sup>58</sup> of suspended, highly-reflective coated mirrors. Thermal <sup>59</sup> noise originating from the mirror coatings will limit the <sup>60</sup> performance of the detectors at their most sensitive fre- <sup>61</sup> quencies once they have reached their design sensitivity. <sup>62</sup> To further increase the sensitivity and therefore the num- <sup>63</sup> ber of detected signals, improved optical coatings are es- <sup>64</sup> sential.

Currently, Advanced LIGO and Advanced Virgo use 66 coatings made from alternating layers of silica (SiO<sub>2</sub>) 67 and tantala doped with titania (TiO<sub>2</sub>:Ta<sub>2</sub>O<sub>5</sub>) [1, 2]. The 68 TiO<sub>2</sub>:Ta<sub>2</sub>O<sub>5</sub> (from now on referred to as Ta<sub>2</sub>O<sub>5</sub>) lay-69 ers dominate coating thermal noise due to a mechanical 70 loss significantly higher than that of SiO<sub>2</sub> at room tem-71 perature [9–13]. Approaches to reduce coating thermal 72 noise include improving Ta<sub>2</sub>O<sub>5</sub> by understanding corre-73 lations between the atomic structure and material prop-74 erties [14–16], elevated temperature deposition [17] and 75 different dopants and doping concentrations [11].

Another option is to replace the  ${\rm Ta_2O_5}$  layers in the  $^{77}$  coating by an alternative amorphous material. However,  $^{78}$  in addition to low mechanical loss, the highly-reflective  $^{79}$  coatings are required to show low optical absorption of,  $^{80}$  ideally, < 1 ppm  $(10^{-6})$  to avoid thermal distortions of  $^{81}$  the mirror due to heating. The optical absorption of the  $^{82}$  best highly-reflective coatings made of  ${\rm SiO_2}$  and  ${\rm Ta_2O_5}$   $^{83}$  is  $\approx 0.27$  ppm [18]. Several potential alternative coating  $^{84}$  materials were investigated over the past years [19–23],  $^{85}$  however, achieving a significant thermal-noise improve-  $^{86}$ 

A completely different approach under investigation is the use of crystalline coatings such as AlGaAs [24]. Crystalline coatings show a significant thermal-noise reduction and low optical absorption, but come with other obstacles such as coating-size limitations and the requirement to remove the coatings from the lattice-matched substrate they are grown on and bond them to the final mirror-substrate used in the detector.

Amorphous silicon (aSi) is a candidate coating material with a low mechanical loss [25] of up to 10 times below that of  ${\rm Ta_2O_5}$ . A high refractive index of 3.5 at 1064 nm makes aSi even more interesting as, due to the high contrast to  ${\rm SiO_2}$  (n=1.45 [26, 27]), fewer layers of both coating materials are required to achieve high reflectivity. This additionally reduces thermal noise, which scales with the coating thickness.

Initially, the optical absorption of aSi was found to be far too high for using it in gravitational-wave detector coatings: At 1064 nm, an extinction coefficient of  $k=7.4\times10^{-3}$  was found after optimum heat treatment (corresponding to  $\approx 4000$  ppm for an highly-reflective coating made of aSi and SiO<sub>2</sub> [28]). However, the past few years have shown a significant absorption reduction to  $k=2.0\times10^{-4}$  [23]. In this paper, we show that at this level of absorption, the use of aSi at 1064 nm is starting to be realistic – when used in a multimaterial design.

Multimaterial coating designs have been proposed to enable the use of materials with high optical absorption, but low mechanical loss [29, 30]. Low-absorption materials are used for the upper coating layers which reflect the majority of the incident laser power. Further down in the coating, where the power is low, higher-absorbing materials with low mechanical loss can be used to reduce the overall thermal noise of the coating. Recently, the multimaterial concept has been experimentally verified: For a proof-of-concept coating, it has been shown that the incorporation of layers with low mechanical loss, but rather high optical absorption, can reduce thermal noise, while

ment and simultaneously keeping the optical absorption low is extremely challenging.

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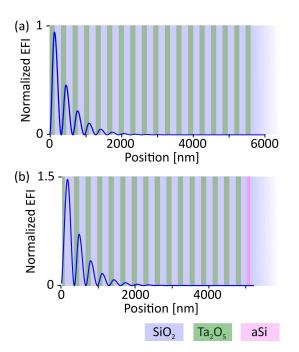


FIG. 1. (a) Schematic of a highly-reflective coating stack<sub>134</sub> made of alternating layers of  $Ta_2O_5$  (green, high n) and  $SiO_2$ <sub>135</sub> (blue, low n) on a  $SiO_2$  substrate. The thickness of all layers is  $\lambda/4$ . The dark-blue line shows the laser light (electric field) intensity (EFI). (b) The  $Ta_2O_5$  layer closest to the substrate was replaced by aSi (pink). The number of  $Ta_2O_5/SiO_2$  layers above was adjusted for coating (b) to have the same reflectiv-139 ity as (a).

low-absorbing layers in the upper layers can keep the total absorption low. Furthermore, it has been shown that heat treatment does not cause damages such as cracks to the multimaterial coating [31].

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In this paper, we give an overview of the current status of aSi absorption at  $1064\,\mathrm{nm}$  and room temperature. We discuss how multimaterial designs – in which as many layers of  $\mathrm{Ta_2O_5}$  as possible are replaced by aSi – can improve the thermal noise of highly-reflective coatings for detectors such as Advanced LIGO, operating at  $1064\,\mathrm{nm}$  at room temperature. Furthermore, we show that the absorption can additionally be reduced by using some split aSi and  $\mathrm{Ta_2O_5}$  layers.

# II. HIGHLY-REFLECTIVE MIRROR COATINGS

This section gives an introduction to the optical $_{151}^{151}$  and thermal-noise performance of highly-reflective mir- $_{152}^{153}$  ror coatings.

# A. Optical Absorption

The optical absorption of a material is proportional to the *extinction coefficient* k. The extinction coefficient is the imaginary part of the complex refractive index  $\mathbf{n} =$  n+ik. We will refer to the real part, n, as refractive index throughout this paper. Highly-reflective mirror coatings are made of stacks of layers with alternating low and high refractive index.

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The reflectivity of the coating depends on the refractive-index contrast at the layer interfaces and on the phase within the layers determined by their optical thickness  $n \times d$ , where d is the geometric thickness. For an optical thickness of  $\lambda/4$  (or odd multiples of this value), the light field within the layers minimizes, resulting in maximum reflectivity. A larger n of one of the coating materials decreases d for all the layers made of this material. If it results in a larger refractive-index contrast between the materials, the total number of (pairs of) layers required to achieve the design reflectivity also reduces. As thermal noise is proportional to d, a thin coating is generally desirable, see Sec. II B.

Figure 1 shows a coating stack made of a total of 17 layers of  $SiO_2$  (blue, n=1.44 at 1064 nm) and 18 layers of  $Ta_2O_5$  (green, n=2.07 at 1064 nm) on a  $SiO_2$  substrate, resulting in a reflectivity of  $99.9994\%^{-1}$ . The amplitude of the blue line, representing the electric-field (light-field) intensity (EFI) within the coating, decreases with every pair of layers – in the example of  $SiO_2$  and  $Ta_2O_5$  by about 50%. The optical absorption in a coating layer is proportional to the spatial integral of the EFI. Consequently, layers contribute less to the total absorption the further down in the coating they are placed.

# B. Coating Thermal Noise

Coating thermal noise amplitude spectral density – from now on referred to as coating thermal noise (CTN) – is given by [30]

$$x(f) = \sqrt{\frac{2k_B T}{\pi^2 f} \frac{1}{w^2} \frac{1 - \sigma_{\text{sub}} - 2\sigma_{\text{sub}}^2}{Y_{\text{sub}}} \sum_{j} b_j d_j \phi_j} . \quad (1)$$

Here,  $k_{\rm B}$  is the Boltzmann constant, T the mirror temperature, f the frequency and w the radius of the laser beam on the coating.  $Y_{\rm sub}$  and  $\sigma_{\rm sub}$  are the Young's modulus and the Poisson ratio of the mirror substrate. d and  $\phi$  are the coating thickness and the mechanical loss with the index j referring to the the  $j^{\rm th}$  layer in the coating (starting from the outermost layer). We assume here that the mechanical losses associated with bulk motion and shear motion [36] are approximately equal  $(\phi_{\rm bulk} \approx \phi_{\rm shear} \approx \phi)$ .

Note that for the gravitational-wave detectors, the design can differ slightly from this simple case in order to optimize coating parameters, e.g. making it reflective at more than one wavelength.

TABLE I. Material parameters of aSi. More detail about the data can be found in [32]. The errors in k and  $\phi$  are taken to be < 10% (see main text for further explanation).

deposition			measured param.			other param.	
facility	method	$T$ [ $^{\circ}$ C]	$\phi$	k	n	Y [GPa]	$\sigma$
MLD	IBS	600	$9^a \times 10^{-5}$	$1.1\times10^{-2}$	$3.34 \pm 0.06$	147 [33]	0.23
ATF	IBS	450	$9 \times 10^{-5} [34]$	$9.0 \times 10^{-3} [28]$	$3.72 \pm 0.12$		
TM	ion plating	500	$2 \times 10^{-5} [35]$	$4.3\times10^{-3}$	$3.82 \pm 0.07$		
UWS	ECR-IBS	400	$2 \times 10^{-5} [23]$	$2(5) \times 10^{-4} [23]$	$3.39 \pm 0.07$		

<sup>&</sup>lt;sup>a</sup> The loss was not measured for MLD coatings. For calculations, the loss of ATF coatings was used.

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In a first approximation, CTN reduces with the square 188 root of the coating thickness [37]. However, the weighting 189 factor  $b_j$  for each coating layer described by

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$$b_{j} = \frac{(1 - 2\sigma_{j})(1 + \sigma_{j})}{(1 - 2\sigma_{\text{sub}})(1 + \sigma_{\text{sub}})} \frac{1}{1 - \sigma_{j}}$$

$$\times \left[ \left( 1 - n_{j} \frac{\partial \theta_{\text{coat}}}{\partial \theta_{j}} \right)^{2} \frac{Y_{\text{sub}}}{Y_{j}} + \frac{(1 - \sigma_{\text{sub}} - 2\sigma_{\text{sub}}^{2})^{2}}{(1 + \sigma_{j})^{2}(1 - 2\sigma_{j})} \frac{Y_{j}}{Y_{\text{sub}}} \right]_{197}^{195}$$

$$(2)_{198}$$

where  $n_i$  is the refractive index of the  $j^{\text{th}}$  layer, shows<sup>200</sup> that the lower coating layers contribute more to the total<sup>201</sup> CTN than the upper layers. The first term in the square<sup>202</sup> brackets of Equ. 2 describes thermal noise arising from 203 fluctuations in coating thickness. It is composed of two204 effects: A fluctuation in a layer changing (1) the optical<sup>205</sup> thickness of this layer, and (2) the position of the front 206 surface of the mirror. These two effects work in opposite<sup>207</sup> directions and partly compensate. The first effect leads to fluctuations in the round-trip phase in each layer,  $\theta_i$ . The sensitivity of the total coating phase,  $\theta_{\rm coat}$ , to these<sup>208</sup> fluctuations, is described by  $\partial \theta_{\rm coat}/\partial \theta_i$ . The magnitude<sup>209</sup> of the latter is proportional to the peak EFI in the  $j^{\rm th}$ layer. As explained in Sec. IIA, the EFI reduces with<sub>210</sub> each double layer resulting in the first effect being smaller<sub>211</sub> for layers positioned further down in the coating, while<sub>212</sub> the second effect is independent of the layer position.213 Therefore, in lower layers, the first effect compensates<sub>214</sub> less for the second effect so that these layers contribute more strongly to  $b_i$ , and consequently to the overall CTN.

## III. AMORPHOUS SILICON ABSORPTION

Future, cryogenic gravitational-wave detectors are<sup>218</sup> likely to use crystalline-silicon mirror substrates and to<sup>219</sup> operate at longer laser wavelengths at which crystalline<sup>220</sup> silicon is transparent [38]. Therefore, to date, most work<sup>221</sup> on aSi has been targeted at coatings for use at 1550 nm<sup>222</sup> and 2000 nm, where it shows an absorption significantly lower than at 1064 nm. Following the success of this work in reducing the absorption, it is on interest to reconsider the use of aSi at 1064 nm.

The absorption of aSi films tends to reduce with heat

treatment, with a minimum occurring at around  $500\,^{\circ}\mathrm{C}^2$ , after which the absorption increases again. Previous measurements on optimally heat-treated, commercially-available, ion-beam sputtered (IBS) aSi films showed that, at  $1064\,\mathrm{nm}$ , k can be as high as  $\approx 1\times 10^{-2}$ . aSi deposited by ion plating can have a significantly lower k, of  $\approx 4\times 10^{-3}$ . An electron cyclotron resonance (ECR) IBS method has recently been shown to produce aSi films with even lower k – as low as  $\approx 2\times 10^{-4}$  at  $1064\,\mathrm{nm}$ . These numbers are summarized in Table I.

The lowest absorption in Table I is still far above the level required for a traditional 2-material coating as shown in Fig. 1(a). However, it is low enough to allow the use of some aSi instead of  $Ta_2O_5$  in the lower layers of a multimaterial coating – see Fig. 1(b).

In the next section, we will look into replacing some of the lower  $Ta_2O_5$  layers (starting from the layer closest to the substrate, see Fig. 1) and investigate the resulting trade-off between optical-absorption increase and thermal-noise reduction.

# IV. MULTIMATERIAL COATING DESIGN: TRADE-OFFS

In this section we discuss the thermal-noise improvement possible by using aSi in multimaterial designs. In Sec. IV A, we introduce the principle we use for our calculations, in Sec. IV B, we introduce different types of aSi, and in Sec. IV C, we discuss the improvement possible.

# A. The Principle

The basic principle of the multimaterial approach is to replace  ${\rm Ta_2O_5}$  layers in the lower part of the coating stack with aSi. This reduces coating thermal noise while the absorption increase is relatively low due to the low laser power at this position.

Figure 1 shows an example: We replace the lowest  ${\rm Ta_2O_5}$  layer in Fig. 1(a) by aSi (with properties for ATF,

<sup>&</sup>lt;sup>2</sup> The exact temperature of the minimum is found to vary with the deposition method and facility used.

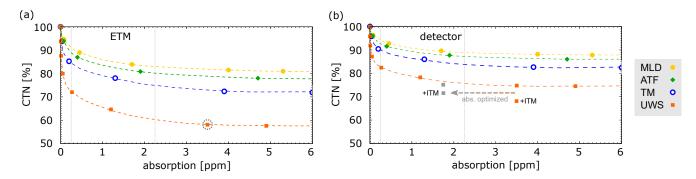


FIG. 2. (a) Relative coating thermal noise of a highly-reflective ETM mirror as a function of additional absorption originating from aSi layers for the four different types of aSi discussed. The 100% mark represents a highly-reflective  $SiO_2/Ta_2O_5$  coating without any aSi, as shown in Fig. 1(a). Each point along the x-axis represents one  $Ta_2O_5$  layer at the bottom replaced by one aSi layer. The number of remaining  $SiO_2/Ta_2O_5$  was adjusted to keep the reflectivity at 99.9994%. A schematic of the coating represented by the point surrounded by the dotted circle is shown in Fig. 3(a). (b) Coating thermal noise of the whole detector taking into account that in the ITM mirror coatings no  $Ta_2O_5$  can be replaced by aSi. The UWS material is an exception with the additional lower orange point showing the improvement possible when also adding aSi to the ITM. A beam radius of 6.2 cm was assumed for the ETM, while for the ITM a smaller beam radius of 5.5 cm was used – as is the case in Advanced LIGO. The vertical, dotted lines indicate the additional tolerable absorption (left lines: current limit, right lines: potential maximum.

see Table I). The aSi layer has a lower loss than Ta<sub>2</sub>O<sub>5259</sub> and it is thinner due to the higher refractive index, both<sub>260</sub> directly reducing CTN. The higher n also results in a reflectivity increase from 99.9994% to 99.9998%. To maintain the reflectivity of (a), we remove layers at the top:261 One complete SiO<sub>2</sub>/Ta<sub>2</sub>O<sub>5</sub> bilayer, plus parts of another bilayer. In this second bilayer we depart from the  $\lambda/4$  design by reducing the thickness of the more lossy Ta<sub>2</sub>O<sub>5</sub> to<sub>263</sub> 0.4 quarters, and increasing the low-loss SiO<sub>2</sub> to 1.6 quarters, of a wavelength, so that we maintain a total optical<sub>264</sub> thickness of  $\lambda/2$  – as otherwise the center wavelength, at<sub>265</sub> which our coating is reflective, would change. This coat- $_{266}$ ing is shown in Fig. 1(b). By replacing one layer of Ta<sub>2</sub>O<sub>5267</sub> by aSi and 'losing' another 1.6 layers, we reduce coating<sub>268</sub> thermal noise by 6% compared to the coating shown in<sub>269</sub> (a). The EFI in the aSi layer at the very bottom of the coating is so low, that the total coating absorption in-270 creases by only 0.07 ppm.

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The number of Ta<sub>2</sub>O<sub>5</sub> layers which can be replaced<sub>273</sub> depends on the tolerable absorption of the final multi-274 material coating: there is a trade-off between the pos-275 sible thermal-noise reduction (thermal noise minimizes<sub>276</sub> when all Ta<sub>2</sub>O<sub>5</sub> layers are replaced by aSi) and the re-<sub>277</sub> sulting absorption (which minimizes when no  $Ta_2O_5$  lay- $_{278}$ ers are replaced). The thermal-compensation system of  $_{279}$ Advanced LIGO is designed for  $0.5\,\mathrm{ppm}$  absorption  $\mathrm{per}_{280}$ coating at maximum laser power [39]. As current coat-281 ings absorb  $0.27\,\mathrm{ppm}$  for the ETMs and  $0.22\,\mathrm{ppm}$  for the ITMs [40], an additional absorption of  $\approx 0.25 \, \mathrm{ppm} \; \mathrm{should}_{283}$ be tolerable without changes. The tolerable absorption may be up to 2.5 ppm with an adjusted compensation system [39, 41], allowing for an additional 2.25 ppm of absorption. These absorption levels are indicated by the vertical, dashed lines in Fig. 2. As possible changes to the system have many implications, in this paper, we analyse the possible thermal noise and absorption performance in a number of cases to point out possibilities for consideration.

#### B. The Materials

We consider four different types of aSi coating material for our evaluation of possible thermal-noise improvement:

- two commercially available ion-beam sputtered coatings produced by MLD and ATF,
- a commercially available ion-plating coating produced by Tafelmaier,
- and a coating produced in a research lab using a novel type of ECR ion-beam sputtering.

The first two materials are 'standard' state-of-the-art aSi, which could be produced in similar quality for use in gravitational-wave detectors without requiring major technological or scientific developments – apart from upscaling to large sizes, an issue all new materials have in common. These two commercial IBS coatings have higher mechanical loss and optical absorption than the other two coatings.

The Tafelmaier (TM) coatings were produced by ion plating <sup>3</sup> and in the past have consistently shown lower mechanical loss and lower optical absorption than the IBS coatings produced by MLD and ATF. This interesting fact shows us that lower loss and absorption could become 'standard' and could give us some hints of how to

<sup>&</sup>lt;sup>3</sup> This technique is unlikely to be used for gravitational-wave detector coatings on short timescales, as it would require several material parameters to be characterized and the development of coating chambers of the required size.

modify other coating procedures such as ion-beam sput-338 tering in order to improve aSi.

The ECR-IBS aSi coatings show an order of magni-339 tude lower absorption than the ion-plating coatings, but 340 this deposition procedure is still in the developmental 341 phase and strong variations in properties have been ob-342 served [23]. A high-quality, high-uniformity, multi-layer 343 optical coating produced using this technique has yet to 344 be demonstrated.

The properties of these four aSi coatings used for our<sup>346</sup> calculations are summarised in Tab. I [32]. In the fol-<sup>347</sup> lowing section, we evaluate how much coating thermal<sup>348</sup> noise in gravitational-wave detectors could be improved<sup>349</sup> by using aSi. We discuss all four materials to show the<sup>350</sup> improvement possible without significant further devel-<sup>351</sup> opment, but also the possibilities the best aSi demon-<sup>352</sup> strated so far would offer.

## C. The Improvement

Figure 2(a) shows the possible coating thermal noise<sub>359</sub> for the four types of aSi introduced in the previous sec-<sub>360</sub> tion, as a function of the additional absorption orig-<sub>361</sub> inating from the aSi layers. The 100 % mark represents a highly-reflective  ${\rm SiO_2/Ta_2O_5}$  coating without any aSi, as shown in Fig. 1(a), with a reflectivity of<sub>362</sub> 99.9994 %. Losses of  $(0.45 \pm 0.03) \times 10^{-4}$  for  ${\rm SiO_2}$  [42] and  $(3.6 \pm 0.1) \times 10^{-4}$  for  ${\rm Ta_2O_5}$  [43] were used for calcu-<sub>363</sub> lations of CTN. The x-axis shows only absorption orig-<sub>364</sub> inating from aSi, and is therefore at 0 ppm for no aSi<sub>365</sub> layers.

For all coatings, the thermal noise reduction converges<sub>367</sub> to a constant level for absorptions of  $\gtrsim 5\,\mathrm{ppm}$ . The rea-<sub>368</sub> son is that for all coating designs we aimed for a re-<sub>369</sub> flectivity of  $R=99.9994\,\%$ . As the sum of absorption,<sub>370</sub> transmission and reflection is defined as  $100\,\%$ , above a<sub>371</sub> certain absorption level, we cannot remove further bilay-<sub>372</sub> ers of  $\mathrm{SiO}_2/\mathrm{Ta}_2\mathrm{O}_5$  or the reflectivity would decrease. The<sub>373</sub> largest error in k was estimated to be  $\sim 10\,\%$  dominated<sub>374</sub> by the measurement reproducibility [32], which directly<sub>375</sub> applies to the x-values in Fig. 2. Error bars were omitted<sub>376</sub> for better legibility.

There has been evidence that the total thermal noise of a multilayer coating can be underestimated by extrapolating from single-layer losses. However, here we do not<sup>378</sup> present any thermal noise numbers, but just relative improvements between coating designs which were all cal-<sup>379</sup> culated based on the same assumptions [42].

An error in mechanical loss affects both the multima- $_{381}$  terial and the  $SiO_2/Ta_2O_5$  coating and, due to the high $_{382}$  loss of  $Ta_2O_5$  dominating the coating loss, largely can- $_{383}$  cels out when looking at the relative CTN improvement $_{384}$  (y-axis). For the coating with the most  $Ta_2O_5$  replaced $_{385}$  with aSi, the relative CTN would change by < 1% for  $_{386}$  an error of  $\le 10\%$  for the mechanical losses of all three $_{387}$  materials (and less for the other coatings).

## 1. ATF

The first green diamond, at a coating thermal noise level of 94% and an absorption of 0.07 ppm, represents a coating with one  $Ta_2O_5$  layer at the bottom replaced by one layer of aSi from ATF and some  $SiO_2/Ta_2O_5$  removed at the top to keep the reflectivity constant, as described in Sec. IV A and shown in Fig. 1(b).

For the second diamond, the second Ta<sub>2</sub>O<sub>5</sub> layer from the bottom was replaced by aSi. One full bilayer of  $SiO_2/Ta_2O_5$  plus the 'modified' (0.4/1.6) bilayer from the first replacement step were removed at the top, reducing coating thermal noise to 87 \%. As we move along the dashed line, each diamond represents a coating in which one more Ta<sub>2</sub>O<sub>5</sub> layer is replaced with aSi. A maximum coating thermal-noise reduction of about 22 % can be achieved with aSi from ATF before the absorption would become too high for a reflectivity of 99.9994 % to be realized. The dashed line is used to indicate that not only full layers of Ta<sub>2</sub>O<sub>5</sub> can be replaced, but also fractions of layers. However, the replacement of fractions does not have a unique solution as we will explain in Sec. VI. Therefore, the dashed lines are only a guide to the eye and partial-layer solutions will not all be positioned exactly on this line.

#### 2. MLD

The MLD coating, represented by the yellow dots and dashed line (top)), shows slightly less thermal noise improvement than the ATF coating. While the mechanical loss is assumed to be identical to ATF, the lower refractive index results in thicker aSi layers and in a slightly lower coating reflectivity, which allows for fewer  ${\rm SiO_2/Ta_2O_5}$  bilayers to be removed at the top of the coating. This results in less coating thermal noise reduction compared to the coating using aSi from ATF. The optical absorption is shifted to lower numbers compared to the ATF coating despite the higher k. This is due to more  ${\rm SiO_2/Ta_2O_5}$  layers remaining on top, which reflect more of the light before it reaches the aSi. With aSi from MLD, a maximum coating thermal-noise reduction of 19% can be achieved.

# 3. Tafelmaier

For the Tafelmaier aSi, both the mechanical loss and the absorption are lower, while the refractive index is higher than for ATF and MLD. Therefore, we win directly in thermal noise from the lower mechanical loss and from the thinner layers reducing d in Equ. 1. Also due to the higher refractive index, fewer  ${\rm SiO}_2/{\rm Ta}_2{\rm O}_5$  layers are required to achieve the design reflectivity, further reducing d. This moves the Tafelmaier results, shown by the blue circles, downwards, while the lower absorption moves the points to the left. When only one  ${\rm Ta}_2{\rm O}_5$ 

layer is replaced by a Si, a significant thermal noise improvement of  $\approx 15\,\%$  can be achieved for an absorption increase of 0.2 ppm. A maximum reduction in coating thermal noise of 28 % would be possible for a coating with the required reflectivity of 99.9994 %.

#### 4. UWS

aSi made by UWS shows an optical absorption about one order of magnitude lower than that of the Tafelmaier coatings. Consequently, coating thermal noise can be improved by almost 30% before the absorption becomes significant. For an absorption of  $\approx 3.5~\rm ppm$ , coating thermal noise reduces by 42% compared to the SiO<sub>2</sub>/Ta<sub>2</sub>O<sub>5</sub> coating shown in Fig. 1(a). For this coating, marked with the dashed circle, only 4.5 bilayers of SiO<sub>2</sub>/Ta<sub>2</sub>O<sub>5</sub> remain on top of 6 bilayers of aSi/SiO<sub>2</sub>, see Fig. 3(a).

## V. ETM AND ITM

So far we have considered highly-reflective coatings suitable for use as the end test-masses (ETMs) in a gravitational-wave detector arm cavity. In the example of Advanced LIGO, the ETM has a reflectivity of 99.9994 % (as shown in Fig. 2(a)). The arm cavity input mirror (input test mass (ITM)) is lower in reflectivity,  $R \approx 98.6$  % in aLIGO. This reflectivity can be achieved using only 6.5 bilayers (if made of ideal quarter layers) of SiO<sub>2</sub>/Ta<sub>2</sub>O<sub>5</sub> plus an additional bilayer made of 1.65 quarter layers of SiO<sub>2</sub> and 0.35 quarter layers of Ta<sub>2</sub>O<sub>5</sub> <sup>4</sup>.

The coating thermal noise of the arm cavity is given by  $^{436}$  the quadrature sum of the contributions from the ITM  $^{437}$  and the ETM. Figure 2(b) shows the total thermal noise improvement of the detector, where the ETMs are improved as in Fig. 2(a), and the ITM design (and thermal  $^{438}$  noise) remains constant, with no aSi layers used.

All coatings using aSi made by MLD, ATF or Tafel-439 maier, shown in Fig. 2(a), require 7.5 or more bilay-440 ers of SiO<sub>2</sub>/Ta<sub>2</sub>O<sub>5</sub> to reduce the EFI before aSi can be<sub>441</sub> used. Therefore, the ITM cannot benefit by using aSi<sub>442</sub> by any of these vendors. However, the optical absorp-443 tion of aSi made by UWS is low enough to allow one<sub>444</sub> bilayer of aSi/SiO<sub>2</sub> to be used underneath 4.5 bilayers of<sub>445</sub> SiO<sub>2</sub>/Ta<sub>2</sub>O<sub>5</sub>, resulting in the ITM design reflectivity of<sub>446</sub> 98.6 % with an absorption of 3.5 ppm. The thermal noise<sub>447</sub> of this coating would be 18 % below that of an ITM made<sub>448</sub> of SiO<sub>2</sub>/Ta<sub>2</sub>O<sub>5</sub> only. As fewer Ta<sub>2</sub>O<sub>5</sub> layers get replaced,<sup>449</sup> less thermal noise improvement is possible than for the<sub>450</sub> ETM.

If we combine this improved ITM with the UWS  ${\rm ETM_{452}}$  of identical absorption, the total thermal noise improve- $_{453}$ 

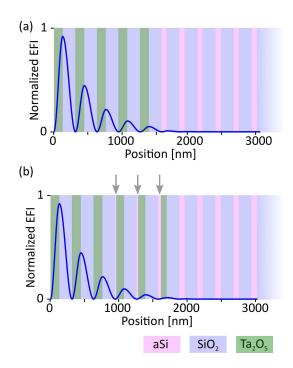


FIG. 3. (a) Schematic of a multimaterial coating using aSi from UWS. The absorption and coating thermal noise of this coating are represented by the orange point surrounded by the dotted circle in Fig. 2(a). (b) Coating with a reflectivity identical to that in (a), but with one aSi layer split and forming 'mixed' layers with  $Ta_2O_5$ . The mixed layers are marked with arrows. As the EFI in the aSi part of the layers is low, the absorption reduces compared to (a).

ment of the detector would be 33 %. This improvement is shown by the orange point labeled '+ ITM' in Fig. 2(b).

## VI. ABSORPTION OPTIMISATION

Further improvement is possible by moving the aSi to positions where the EFI is low. Figure 3(a) shows a multimaterial coating using aSi from UWS with an absorption of about 3.5 ppm. In the coating shown in Fig. 3(b), the same amount of aSi is used as in (a). However, in (b) the outermost aSi layer was split into 0.5, 0.3, and 0.2 quarter layers (adding up to one quarter layer). Together with fractions of  $Ta_2O_5$  quarter layers, they form some of the high refractive-index layers of the coating – marked with grey arrows. The reflectivity of the coating shown in (b) is identical to that in (a), but its absorption has reduced from 3.5 to 1.8 ppm as the aSi layer fractions are positioned where the light field is low.

The total optical layer thickness of the mixed layers is slightly larger than a quarter of a wavelength to keep the reflection maximum and the absorption minimum at the design wavelength, resulting in  $\approx 2.34$  instead of 2 layers of Ta<sub>2</sub>O<sub>5</sub>. Consequently, coating (b) has slightly higher coating thermal noise than coating (a) with an overall reduction of 40.5% (instead of 42%) compared to

<sup>&</sup>lt;sup>4</sup> Note that this is a simplified example and not the actual design<sub>457</sub> of the Advanced LIGO ITMs.

a  $SiO_2/Ta_2O_5$  coating.

The upper, gray square in Fig. 2(b) shows the detector<sup>511</sup> coating thermal noise when the optimized coating is com-<sup>512</sup> bined with a classical ITM (without aSi). For the bot-<sup>513</sup> tom, gray square, the ITM was also optimized <sup>5</sup>. When<sup>514</sup> using mixed layers in both the ITM and ETM, a coating thermal noise improvement of 28.5 % is possible at an absorption of about 1.8 ppm per mirror (compared to<sup>515</sup> 32 % for about 3.5 ppm).

This is an example to illustrate the concept, which<sup>516</sup> we tried to keep as simple as possible while keeping the<sup>517</sup> layers realistically thick (from a deposition perspective).<sup>518</sup> Some other options for incorporating aSi into multimate-<sup>519</sup> rial coatings via stacked-triplet design have recently been<sup>520</sup> investigated by Pinto [44].

#### VII. REFLECTIVITY OF GREEN LIGHT

A way to stabilize the detectors' arm length is the  $_{527}$  use of 532 nm laser light [45, 46]. All coating designs  $_{528}$  discussed so far were optimized for high reflectivity at  $_{529}$  1064 nm only. If, for example, the layer thicknesses of the  $_{530}$  coating shown in Fig. 3(a) were adjusted to 1.27 quarter  $_{531}$  layers for the low, and to 0.71 quarter layers for the high- $_{532}$  index layers, a reflectivity of  $\approx 96\,\%$  at 532 nm could  $_{533}$  be achieved. To maintain the reflectivity at 1064 nm  $_{534}$  at 99.9994 %, in this example, two bilayers of aSi/SiO $_{2535}$  would have to be added at the bottom of the coating.

The additional CTN from these two additional bilayers<sub>537</sub> is low due to the low mechanical loss of these materials.<sub>538</sub> Furthermore, the thickness of the high-loss Ta<sub>2</sub>O<sub>5</sub> layers<sub>539</sub> is reduced and the peak EFI is moved from the layer in-<sub>540</sub> terfaces into the low-loss SiO<sub>2</sub> layers, reducing the overall<sub>541</sub> CTN. This type of design for reflectivity at two wave-<sub>542</sub> lengths is very similar to that previously suggested for<sub>543</sub> CTN optimization [47, 48], as Equ. 1 shows.

In reality, the Advanced LIGO coatings have an opti-545 mized design to enable high reflectivity at both 532 nm546 and 1064 nm which also results in a slight CTN reduc-547 tion. Here, for simplicity, we do not consider thickness optimization for 532 nm reflectivity. However, when optimizing both our reference SiO<sub>2</sub>/Ta<sub>2</sub>O<sub>5</sub> coating and our<sup>548</sup> multimaterial coating, to first order any CTN improvement would be likely to cancel out, as we only consider 549 the relative improvement between coatings.

While from the reflectance perspective there are no ob-551 jections to this, the optical absorption of aSi at 532 nm is552 known to be high. For the UWS coatings, we do not have553 absorption data at this wavelength, but for the Tafel-554 maier coatings, a k of about 0.8 was found, which is in555 agreement with literature 6 and which we will assume556 here as an upper bound for the UWS aSi. At this level557 of absorption, all 532 nm light transmitted into the aSi558

layers would be absorbed. Therefore, it would not be possible to inject the 532 nm light into the interferometer through the ETMs, but another solution would have to be found, e.g. coupling the green light in through (aSi-free) ITMs as in KAGRA [49].

## VIII. CONCLUSION

We have shown that coating thermal noise in gravitational-wave detectors operating at 1064 nm and at room temperatures could be improved by using aSi in a multimaterial coating design. The achievable improvement is a trade-off between thermal noise reduction and absorption increase.

With state-of-the-art commercial coatings, only the ETMs could be improved, while with material in the development phase, an improvement of the ITM would also be possible. By splitting some high-index layers into 'mixed' layers made of aSi and  ${\rm Ta_2O_5}$ , the absorption can additionally be reduced. For the best aSi identified so far, a thermal noise reduction of almost 30 % in the detectors would be possible for an optical absorption of about 1.8 ppm per mirror due to the aSi. (For a higher tolerable absorption, more improvement is possible, for a lower tolerable absorption, it is less.) This approach may therefore be an attractive option.

While the absorption of 532 nm laser light in the aSi layers should not impose a major problem on the coating reflectivity, injecting 532 nm light through the ETMs would very likely not be possible and would need a different solution.

The aSi-based multimaterial coating design is versatile and can easily be used in conjunction with other developments – for example by replacing the remaining  ${\rm Ta_2O_5}$  by a material with lower mechanical loss. Several such materials are currently under development in the community, and we believe that combining these new materials with our multimaterial approach has excellent potential to fully meet the requirements of future room-temperature GW detectors.

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 $0.7 \times 1a_2O_5//51O_2$ ,  $1a_2O_5$ .

<sup>&</sup>lt;sup>5</sup> From the top 2 bilayers of  $SiO_2/Ta_2O_5$ ,  $(0.2 \times aSi + 0.9 \times ^{669} Ta_2O_5)/SiO_2$ ,  $(0.3 \times aSi + 0.8 \times Ta_2O_5)/SiO_2$ ,  $(0.4 \times aSi + ^{670} 0.7 \times Ta_2O_5)/SiO_2$ ,  $Ta_2O_5$ .

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