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 This is a pre print version of the following article:

 Original Citation:

 Availability:

 This version is available http://hdl.handle.net/2318/1768849

 since 2021-01-25T11:01:26Z

 Published version:

 DOI:10.1016/j.carbon.2020.07.033

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4	
5	This is an author version of the contribution published on:
6	Questa è la versione dell'autore dell'opera:
7	Carbon 169, 2020, 357-369 DOI: 10.1016/j.carbon.2020.07.033
8	The definitive version is available at:
9	La versione definitiva è disponibile alla URL:
10	https://www.sciencedirect.com/science/article/pii/S000862232030693X
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How do the graphenic domains terminate in activated carbons and carbon-supported metal catalysts?

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- 9

10 Abstract

11 This work is focused on the determination of the most likely border morphologies in the H-terminated graphenic domains that constitute activated carbons, focusing on four different carbons and six 12 13 supported-metal catalysts. Among the available characterization techniques, Inelastic Neutron Scattering spectroscopy provides detailed vibrational spectra containing the fingerprint of the 14 15 terminal C-H groups of carbon materials. From the experimental spectra, we observed clear differences among carbons having different origin, and a systematic decrease in the integrated area 16 upon metal deposition that linearly scales with the nanoparticles dispersion. Density Functional 17 Theory simulations are fundamental to assign the experimental bands to specific species. Thus, an 18 extended simulation work on both regular and defective aromatic models was carried out, providing 19 the inelastic neutron scattering fingerprint of a large number of different C-H terminations. By fitting 20 the experimental spectra with a linear combination of the simulated spectra, it was possible to quantify 21 22 the concentration of each terminal C-H geometry in the samples, and to identify the species most 23 affected by the metal nanoparticles deposition. Specific benzene rings exposing a single C-H group appear to have a decisive role in the interaction with the metal nanoparticles and their deposition 24 25 procedure.

1 **1. Introduction**

2 Activated carbons are large surface area carbonaceous materials obtained from the pyrolysis and subsequent activation of several natural bio-polymeric precursors (such as wood, peat or coconut 3 4 shells) [1]. They are widely employed as adsorbent or storage materials [2, 3], as supports in 5 heterogeneous catalysis [4] or as catalysts on their own [5-9], and also for innovative applications 6 such as lithium batteries [10]. From a micro-structural point of view, activated carbons are constituted 7 by small and heterogeneously sized and shaped graphenic domains mainly terminated by C-H groups, 8 which also exhibit variable amounts of functional groups containing heteroatoms (such as oxygen or nitrogen) and a certain degree of defectivity. All these properties, as well as the surface area and 9 10 porosity, are largely affected by the nature of the starting precursor and by the activation procedure [11, 12]. Altogether, the textural, morphological and surface properties of an activated carbon have 11 12 dramatic effects on their catalytic performances. As an example, we recently demonstrated that, for 13 Pd-based catalysts supported on activated carbons, oxygen-containing functional groups at the carbon surface favour the hydrogenation of polar substrates, while they are detrimental to the hydrogenation 14 of nonpolar substrates [13]. 15

It emerges that a detailed understanding of the microscopic structure of activated carbons is 16 essential to optimize their performance for a specific application, in particular for catalysis. Among 17 the numerous techniques that can be employed for their characterization [14], Inelastic Neutron 18 Scattering (INS) spectroscopy stands up as one of the most powerful methods to determine the nature 19 of the C-H terminations of the graphenic domains. INS is a vibrational spectroscopy that exploits 20 neutrons to excite molecular and/or atomic vibrations inside materials, and the intensity of the signal 21 22 is proportional to the neutron scattering cross section of the atoms involved in the vibrational mode 23 and to their displacement with respect to their equilibrium position. Since hydrogen displays both the largest neutron scattering cross section and large displacements due to its lightweight, this technique 24 25 is particularly sensitive to the hydrogen-containing species in a material. This is the reason why INS spectroscopy has been successfully employed in the study of several carbon-based materials including 26 27 activated carbons [15-18], carbon blacks [16, 19], coals [16, 20], glassy carbon [16, 21] or reduced graphene oxides [22, 23]. Despite the huge variety of properties displayed by these materials, their 28 29 structure down to the atomic scale is to some extent similar, and so are their main spectral features.

Nevertheless, the bands assignment in an INS spectrum can be tricky. The reason is that this technique has formally no selection rules allowing, at least in principle, the detection of all the vibrational modes. Therefore, INS spectra of complex and heterogeneous samples containing high amounts of hydrogenous species usually show a large number of bands, whose assignment becomes challenging. The employment of computational methods to simulate the INS spectra can be extremely

helpful and, among the available computational chemistry approaches, DFT simulations 1 demonstrated to provide reliable results at reasonable computational costs. Albers et al. were the first 2 authors who adopted a combined INS and computational approach to study activated carbons [15]. 3 Their analysis showed that the fingerprint region of the C-H out-of-plane bending modes (800-1000 4 cm⁻¹) is the one mostly affected by differences in the microstructure of activated carbons. More 5 recently, some of the authors revisited the interpretation of INS spectral features in the fingerprint 6 region of activated carbons [24]. In particular, a more precise assignment of the contributions from 7 8 different C-H terminations of the graphenic domains (including zigzag and armchair edges, as well 9 as corners) was proposed. Still, so far the interpretation of the C-H vibrational features in activated carbons was limited to the comparison of experimental data of mediocre quality with simple 10 11 theoretical models consisting of a few aromatic rings completely terminated by C-H groups, which are not representative of the real complexity of activated carbons. 12

13 Thanks to recent advances in INS instrumentations, the spectra collected on modern neutron spectrometers on well-prepared samples show an unprecedented signal-to-noise ratio, allowing for a 14 more precise interpretation of the spectral features and for a higher reliability for the spectral 15 differences. As an example, Cavallari et al. were recently able to collect high quality INS spectra of 16 regular and defective hydrogenated graphene samples [23]. Combining the experimental spectra and 17 DFT simulations, they identified the INS fingerprints of C-H groups at the periphery of the graphenic 18 planes as well as inside holes. They quantified the two species abundance and they were able to 19 monitor the migration of H atoms to adjacent vacant sites as a function of the temperature. Similarly, 20 in our recent work we investigated a 5 wt.% Pt catalyst supported on an activated carbon in the 21 presence of different amounts of H₂ [18]. The collected INS data allowed for a detailed difference 22 23 analysis between the spectrum of the catalyst and that of the carbon support, as well as between the spectrum of the catalyst in the absence and in the presence of H₂. With this approach, we were able 24 25 to assess more specifically the location of the Pt nanoparticles on the carbon support and to highlight the occurrence of hydrogen spillover from the Pt nanoparticles to the support. Despite the low Pt 26 27 loading, we were also able to detect surface Pt-hydride species forming and modifying upon 28 increasing the H₂ concentration. These examples demonstrate that nowadays we can rely on the fact that differences in the INS spectra of carbon-based samples under different conditions can be 29 30 interpreted as clear structural modifications. At the same time, differences between measured and 31 computed INS spectra can be considered as the result of a lack of specific features in the theoretical 32 model.

In this work, we propose a systematic DFT simulation study of a large number of graphenic
 model systems, aimed at providing the INS fingerprints of a wide range of different C-H termination

geometries, as well as at understanding how they are affected by the presence of defects. The theoretical results are used for interpreting the experimental INS spectra of four physically activated carbons, deriving from wood and peat, and of a series of carbon-supported Pd- and Pt-based catalysts employed for hydrogenation reactions of industrial relevance. In principle, the same methodology can be easily extended to other classes of carbon-based materials. In this respect, we wish that this paper might become a guide for researchers who want to adopt similar approaches in material science.

7 **2. Experimental**

8 2.1. Activated carbon samples

9 The four activated carbons investigated in this work were commercially available carbons and were 10 provided by the Catalyst Division of Chimet S.p.A. All of them were physically activated with steam, but they derive from different raw materials, wood (hereafter labelled as CwA and CwB) and peat 11 12 (labelled as CpA and CpB). The ash content of the four carbons was determined by calcination and 13 is below 2 wt% in the two carbons of wood origin, and of about 3 wt% for CpA and 5.5 wt% for CpB. Table 1 summarizes the specific surface area (SSA_{BET}) and micropore volume (V_{micro}) of the four 14 carbons, as determined by N_2 physisorption measurements. The adsorption isotherms for the four 15 supports are reported in the supporting information, Figure S1. The catalysts supports are activated 16 carbons derived from natural sources and, as such, they are characterized a certain degree of 17 heterogeneity in terms of carbon structure, morphology and porosity intrinsic in the natural source. 18 Representative SEM images of the four activated catalysts at two different magnifications are 19 reported in Figure S2 in Supporting Information to show the structural inhomogeneity of the carbons 20 21 at the μ m scale.

Table 1. Properties of the activated carbons investigated in this work: raw material, activation method, specific surface area (SSA) as determined by N_2 physisorption measurements by applying the BET method and the corresponding volume

of the micropores (V_{micro}) .

Sample	Origin	Activation	$SSA_{BET}(m^2g^{-1})$	V _{micro} (cm ³ g ⁻¹)
CwA	wood	steam	1018	0.63
CwB	wood	steam	1325	0.65
СрА	peat	steam	903	0.49
СрВ	peat	steam	882	0.52

25 **2.2. Catalysts preparation**

All the catalysts were prepared in the Chimet S.p.A. laboratories. Two Pd-based catalysts (5 wt% Pd

27 loading) were prepared using CwA as a support, following the deposition–precipitation method as

reported elsewhere [25, 26]. Na₂PdCl₄ was used as the metal precursor and Na₂CO₃ as the basic agent.

29 The two Pd/CwA catalysts differ in the reduction step: in one case, after the deposition of palladium

as hydroxide, a pre-reduction was carried out by means of HCOONa at 55 °C for 1 h, followed by
water washing in order to remove residual chlorine and drying at 120 °C overnight. In the other case,
no pre-reduction was performed and the Pd deposition step was directly followed by washing and
drying. The two catalysts are labelled as Pd(R)/CwA and Pd(NR)/CwA, where (R) and (NR) stand
for reduced and not-reduced, respectively.

A series of Pt-based catalysts (5 wt% Pt loading) were prepared using the four activated 6 7 carbons as supports, following a deposition-precipitation method similar to the one reported by Kaprielova et al. [27]. A pre-reduction was carried out also in this case by means of HCOONa at 80 8 9 °C for 1 h, followed by water washing and drying at 120 °C overnight. These catalysts are labelled with Pt(R), followed by the name of the carbon employed as support. The main characteristics of all 10 the considered catalysts are summarized in Table 2. In all the cases, the metal dispersion was 11 evaluated by means of CO chemisorption. The measurements were performed at 50 °C using a 12 dynamic pulse method on samples pre-reduced in H₂ at 120 °C [28]. In a typical experiment, the 13 catalyst (200 mg) is loaded inside the U-tube, heated in He up to 120 °C (heating rate of 10 °C min⁻¹), 14 15 reduced in H₂ for 30 min, and finally cooled to 50 °C in He (cooling rate of 10 °C min⁻¹). A CO/metal average stoichiometry of 1 was assumed to calculate the metal dispersion for both the metals, as 16 widely documented [29]. The discussed heterogeneity of the support does not point to a poor quality 17 of the commercial catalysts. To verify that the supported metal nanoparticles are in homogeneous in 18 terms of morphology (spherical-like particles) and display a narrow particle size distribution, we 19 20 recorded TEM micrographs on Pd/CwA and Pt/CwA catalysts (see figure S3 in Supporting information for representative images). 21

Table 2. Properties of the catalysts studied in the current work: the activated carbon support type, the deposited metal, the eventual presence of a pre-reduction step and the metal nanoparticles dispersion (evaluated by means CO chemisorption).

SAMPLE	SUPPORT	METAL	PRE-REDUCED	DISPERSION (%)
PD(NR)/CWA	CwA	Pd	No	24
PD(R)/CWA	CwA	Pd	Yes	24
PT(R)/CWA	CwA	Pt	Yes	72
PT(R)/CwB	CwB	Pt	Yes	77
Pt(R)/CpA	СрА	Pt	Yes	71
Pt(R)/CpB	СрВ	Pt	Yes	69

24 **2.3. INS measurements**

The INS measurements were performed on the LAGRANGE spectrometer at ILL in Grenoble (France) [30-32]. The INS spectra were collected in the 24-1020 cm⁻¹ range, with a resolution of $\Delta E/E$

of about 2%. The spectra were recorded at 20 K to reduce the thermal effects that broaden the spectral

features. The raw data were opportunely reduced to the scattering function $S(\omega)$ and plotted versus

energy transfer (in units of cm⁻¹). In order to ensure an easy comparison of each catalyst with its
respective support, we renormalized the intensity of the former to match the same background
intensity of the latter at 1020 cm⁻¹. The analysis, described in section 3.1, confirmed that contributions
other than the background are negligible in this region.

5 Prior to the measure, all the samples were outgassed at 120 °C in dynamic vacuum 6 (equilibrium pressure below 10⁻⁴ mbar) in order to remove the physisorbed water, and any further 7 manipulation was performed within an Ar-filled glovebox in order to avoid any contamination with 8 air moisture. A weighted amount of each sample (in the order of about 2-3 g) was inserted in a 9 cylindrical Al-cell (4 cm high, 16 mm diameter) sealed with In-wire and then measured.

3. Theoretical methods

11 **3.1. DFT simulations of the INS spectra**

12 The DFT simulations of the INS spectra were performed employing the CRYSTAL17 software [33] in combination with the aClimax software [34]. A good compromise between results reliability and 13 computational cost was found in the combination between the B3LYP hybrid functional [35, 36] and 14 the 6-31G** basis set, as also proposed in the literature for similar systems [37, 38]. Some models 15 presented an open shell electronic configuration, for which the UB3LYP method was chosen. The 16 general models employed for simulating the structure of activated carbons consist in polycyclic 17 aromatic systems displaying different geometries and extensions, with hydrogen atoms used to 18 saturate dangling bonds at edges and corners. Similar models have been already adopted in the 19 literature to simulate the vibrational properties of the graphenic domains in carbon-based materials 20 [15, 24, 38]. A total of 16 regular and 6 defective H-terminated different structures are reported in 21 22 this work (see Figure 3 and Figure 8), while a greater number of them have been effectively simulated. 23 The geometry of all the models was fully optimized, and successively a calculation of the vibrational modes was performed. All the structures were categorized as true energy minima by the absence of 24 imaginary frequencies. The condition for the SCF convergence was set to 10^{-8} and 10^{-10} hartrees 25 during geometry optimization and frequency calculations, respectively. In order to ensure the best 26 accuracy, very thigh values (8 8 8 25 50) have been set for the five CRYSTAL17 integrals screening 27 tolerance. 28

The output of the frequency calculation was used to simulate the INS spectra by employing the aClimax software [34]. Simulations of higher order transitions and phonon wings contributions were also performed, but not included into the final simulated spectra because they generated features that were not experimentally observed. The contribution of specific border fragments to the overall INS spectrum was simulated by manually setting the scattering cross section of all the atoms not 1 included within the interested fragment to 0 barns. For a better comparison with the experimental 2 spectra, the calculated frequencies have been scaled: the employed scaling function was obtained 3 through a least squares linear regression of the experimental and calculated frequencies for the 4 benzene molecule, as proposed in reference [39], resulting in the scaling function $v_{scaled}=22.1 + 0.9543v_{calculated}$.

6 **3.2. Linear combination analysis**

7 To evaluate the weight of each simulated C-H terminal geometry in the experimental spectra, a linear 8 combination fit analysis was performed. Each theoretical C-H contribution was normalised to the 9 intensity of a single C-H group, averaged over similar models and then multiplied for a weight 10 coefficient, whose final value in the best-fit solution is proportional to the concentration of that species in the sample. To compare among the values obtained for different samples, the optimised 11 values of the sum of these coefficients were normalised to 1 for each support, while the catalysts were 12 normalised respect to the corresponding support. A small multiplicative scaling factor s (0.99 < s <13 1.01) was introduced on the frequency to adjust for small shifts along the horizontal axis. 14

15

16 **4. Results and discussion**

17 4.1. Experimental INS spectra

In the analysis of the experimental INS spectra, we decided to focus our attention on the 690 - 102018 cm⁻¹ range, characteristic of the out-of-plane C-H bending vibrations [15-18, 24]. This is the region 19 of the spectra that exhibits the greatest variations among the samples, and thus can be considered as 20 the INS fingerprint of C-H terminations in the activated carbon samples. The spectra in the extended 21 range 200-1020 cm⁻¹ are reported in Figure S4, S5, S6 and S7 in the supporting information, while 22 an overall assignment of all the observed bands is described elsewhere [17]. Figure 1 shows the INS 23 24 spectra of the four activated carbons and of the corresponding Pd- and Pt- based catalysts in the 690 -1020 cm⁻¹ range. All the spectra are characterized by three main bands centred at similar positions 25 (ca. 800, 880 and 955 cm⁻¹, with a tail extending at 750 cm⁻¹). Differences in the relative intensity of 26 the three bands are observed between carbons derived from different raw materials (Figure 1a-d), but 27 also between carbons having the same origin (e.g. for CwB, Figure 1b, the band at ca. 880 cm⁻¹ is 28 distinctively more intense than for CwA, Figure 1a). These variations are expected to be caused by 29 changes in the relative concentration of different C-H terminal geometries of the graphenic domains 30 within the investigated samples. 31



1

Figure 1. Experimental INS spectra of the four activated carbons (black) and of the corresponding Pd- and Pt-based
catalysts, as measured at the LAGRANGE spectrometer at ILL. The spectra are shown in the 690 – 1020 cm⁻¹ range,
characteristic of the out-of-plane C-H bending vibrations. The intensity of the catalysts' spectra is normalised to the one
of the support to the value of the background.

Variations are also observed when comparing the INS spectra of the Pd and Pt catalysts (red, 6 7 orange and yellow spectra in Figure 1a-d) with those of the parent carbons (black spectra). The 8 observed changes are very similar irrespective of the carbon type, of the metal and of the metal 9 deposition method. A decrease in intensity of all the three bands is observed, more pronounced for the one centred at 880 cm⁻¹ than for the others. This behaviour could be the result of the interaction 10 of Pt and Pd nanoparticles with the C-H terminations of graphenic domains in the carbon support, 11 and it highlights that some types of C-H terminations are more affected than others. Similar variations 12 13 were recently reported for the Pt(R)/CwA catalyst compared to the bare carbon support, and interpreted as an evidence that the Pt nanoparticles can be located at the edges of the sp² graphenic 14 15 domains [18].

16 The absolute intensity of INS spectra is directly proportional to the hydrogen concentration 17 within each sample [15, 17]. In this respect, CwA is the activated carbon displaying the highest

hydrogen concentration, closely followed by CwB and CpA, while CpB shows a much lower amount 1 of hydrogen in respect to the others. Albers et al. [15] demonstrated that the integrated area of the 2 spectra in the C-H bending region linearly scales with the H concentration in the sample. We have 3 thus decided to use the integrated area of the spectra in the out-of-plane bending region (690-1020) 4 cm⁻¹) to evaluate the H concentration within each of the samples. The variation in the integrated area 5 of the spectra of all the catalysts with respect to the bare carbon support was calculated and plotted 6 against the metal nanoparticles dispersion, as shown in Figure 2. For the Pt catalysts, the decrease in 7 8 the integrated area (i.e. in the H concentration) is directly proportional to the dispersion of the metal 9 nanoparticles, as shown in Figure 2b. This observation suggests that smaller the nanoparticles, higher is their interaction with the borders of the graphenic platelets. The two Pd catalysts instead exhibit 10 large variations in H content albeit their low dispersion, as shown in Figure 2a, suggesting that Pd 11 nanoparticles interact more with the C-H terminations than the Pt ones. 12



13

Figure 2: a) Difference in the integrated area of the INS spectra in the out-of-plane bending region (690-1020 cm⁻¹)
 between the bare carbon support and the corresponding catalysts versus the metal dispersion. b) Magnification of the Pt
 catalysts points: a linear trend can be observed.

17 4.2. DFT simulations of the INS spectra for regular C-H terminations

The three bands observed in the experimental spectra shown in the previous paragraph are attributed to vibrations of specific C-H terminal geometries. The observed changes in relative intensity reflect a change in their concentration. A qualitative bands assignment was done previously on the basis of DFT calculations performed on very simple aromatic models [24], that cannot account for the heterogeneity of real samples. For this reason, we underwent a large and systematic theoretical study aimed at exploring a much wider range of possible sizes, border geometries and terminations, with the aim to perform a quantitative analysis of the experimental spectra.

The models initially employed for the new DFT simulation of the INS spectra consist in 1 regular aromatic structures where all the carbon atoms are sp^2 hybridized and all the terminations are 2 saturated by H atoms. Their geometry can be described in terms of extended regular borders and 3 corners. Extended regular borders can show a regular sequence of benzene rings exposing a single C-4 H group named zigzag edge (Z), or an alternation of benzene rings exposing two adjacent C-H groups 5 named armchair edge (A). The intersection of two regular borders generates a corner site that can 6 7 present two, three or four adjacent C-H groups and will be labelled as *duo* (D), *trio* (T) or *quatro* (Q), respectively, following the nomenclature proposed by Zander for substituted benzenes [40, 41] 8

9 The models have been categorized in Figure 3 as a function of their shape, i.e. the relative extension of Z and A edges. Figure 3a summarizes the models having a rectangular shape, which is 10 defined by two zigzag and two armchair edges and by four T corners. Moving in the horizontal 11 direction, the Z edge increases in length while keeping constant the dimension of the A edge; moving 12 13 in the vertical direction, the A edge increases, instead. Figure 3b displays models having a parallelogram shape, which is defined by a single type of edge, either Z (C₃₈H₁₈, C₅₈H₂₀, C₅₀H₂₀), or 14 A ($C_{54}H_{22}$). For the $C_{38}H_{18}$ and $C_{50}H_{20}$ models, an aromatic ring was added in order to break the 15 symmetry, resulting in the addition of a corner of Q type. In this category, the typologies of corners 16 17 are more varied, resulting in an alternation of D and T corners in the models featuring Z edges and of T and Q corners for the model featuring A edges. Finally, Figure 3c shows two hexagonal models, 18 the latter being also the largest model adopted in this work, showing only Z edges alternating with D 19 corners. The relative abundance of Z and A edges, as well as of D, T, and Q corners is reported in 20 Figure 3 for all the models in parenthesis, following the notation (Z/A) for the edges, and [D/T/Q] for 21 22 the corners.





2 Figure 3. Overview of the models simulated in this work, categorized in terms of their shape. For all the models, the C/H ratio is indicated, as well as the relative abundance of zigzag and armchair edges (Z/A), and of duo, trio and quatro

3 4 corners [D/T/Q].

Figure 4a shows three representative simulated INS spectra for three models having different 1 size and shape, in comparison with the experimental spectrum of CwA, taken in the following as a 2 reference. As expected, all the models contribute with some bands in the 690 - 1020 cm⁻¹ range. 3 However, different models clearly show distinct bands in terms of number, position and relative 4 intensity. The simulated INS spectrum of $C_{54}H_{18}$ is the closest to the experimental one in terms of 5 both bands positions and relative intensity, while the INS spectra of $C_{66}H_{22}$ and $C_{38}H_{18}$ evidently 6 show bands not present in the experimental one. This demonstrated that the model geometry has a 7 8 dramatic effect on the result.

9 In order to better analyse the characteristic spectral features of each C-H termination portion in the models, their contributions to the overall simulated spectra were separated and analysed 10 individually. For example, model C₅₄H₁₈ includes two distinct terminal C-H geometries (Z edges and 11 D corners), whose spectral contributions are reported in Figure 4Figure 4b. It is evident that the two 12 13 C-H terminations display significantly different spectral features: the D corners are responsible for the appearance of the two bands at 800 and 955 cm⁻¹, while the C-H at the Z edges are mainly 14 responsible for the band at 880 cm⁻¹. However, D corners belonging to two different models are 15 characterized by very similar spectral features as shown in Figure 4c. Analogous considerations apply 16 17 for the other models and terminations, making it clear that the INS spectra of each C-H termination in the 690 - 1020 cm⁻¹ range mostly depend on local geometry and are only marginally affected by 18 the general structure of the model. This does not hold for the rest of the spectral range, where the 19 vibrational modes often display collective features, resulting in spectra reflecting the geometry of the 20 whole model more than the single C-H terminations. [24] 21







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Figure 4. Part a): Comparison between the simulated INS spectra of three models having different shape and the experimental spectrum of CwA (black). Different model geometries result in different simulated spectra. Part b): The partial contribution of the two types of C-H termination characterizing the $C_{54}H_{18}$ model in comparison to the experimental spectrum of CwA (black). Part c): Simulated spectral contribution of two D corners belonging to two different models, compared to the experimental spectrum of CwA (black).

- On these bases, we systematically investigated the INS contributions of specific C-H terminal
 geometries in the 690 1020 cm⁻¹ range. Five different classes of C-H terminations were identified,
 and their partial INS spectra were simulated and compared.
- **Extended zigzag borders** give rise to a couple of INS signals centred at ca. 800 cm⁻¹ and 880 9 cm⁻¹, which are attributed to the out-of-plane bending mode, δ (C-H), in which adjacent C-H sites 10 vibrate out-of-phase and in-phase, respectively. Being a correlated mode, the exact position of these 11 12 two bands and their relative intensity depend on the extension of the border or, in other words, on the distance of the C-H groups from the corners. For the sites located at the extremities, it also depends 13 14 on the geometry of adjacent sites. Figure 5a and b show some representative spectra of zig-zag borders belonging to different models, whereby each C-H group has been labelled as Z, followed by a number 15 indicating the distance from the corner (e.g. Z2 is the second C-H belonging to the regular zigzag 16 border starting counting from the corner). For Z1 sites (i.e. those adjacent to a corner) another letter 17 is added to describe the adjacent corner (D = duo, T = trio). Trying to generalize, the spectra of Z 18 sites far from the corners tend to converge, as shown in Figure 5a (e.g. the spectra of Z3 and Z4 sites 19 in $C_{76}H_{26}$ are very similar). In these cases, usually the band at ca. 800 cm⁻¹ is more intense than that 20 at 880 cm⁻¹. Z2 sites tend to show the same two δ (C-H) bands, shifted at higher energy transfer. The 21 spectra of Z1 sites are more scattered depending on the model (Figure 5b), and mainly depend on the 22 type of the adjacent corner. For the Z1D sites (i.e. Z1 sites close to a *duo* corner) the band at ca. 800 23 cm⁻¹ is largely reduced in intensity and the spectra are dominated by the band at 880 cm⁻¹. In the 24 spectra of Z1T sites the band at ca. 870 cm⁻¹ is usually split in two components and more intense than 25 that at ca. 800 cm^{-1} . 26

The spectra of **extended armchair borders** (Figure 5c) are generally characterized by two bands centred at ca. 810 cm⁻¹ and 920 cm⁻¹, corresponding to the out-of-plane δ (C-H) modes of the two adjacent C-H sites that vibrate in-phase and out-of-phase, respectively. In this case, the final spectrum is almost the same irrespective of the length of the border and the geometry of the platelet.



Figure 5. Part a): Representative contributions of different types of C-H terminations at zig-zag borders (Z sites) far from the corners and belonging to different models. Part b): representative contributions of Z1 C-H terminations differing in the type of adjacent corner (D = duo, T = trio) and belonging to different models. Part c): Representative contributions of different C-H terminations at armchair borders (A sites).

A *duo* corner (Figure 6a) consists in a benzene ring at a corner position exposing two adjacent C-H groups. The INS spectra of duo corners are all characterized by two bands centred at ca. 795 and 960 cm⁻¹, corresponding to the out-of-plane δ (C-H) modes of the two adjacent C-H groups that vibrate inphase and out-of-phase, similarly to the extended armchair border case. However, the spectra of D 10 and A sites can be easily distinguished on the bases of the position of the two bands, evidently more 11 spaced for D corners than for A edges. A *trio* corner (Figure 6b) consists in a benzene ring exposing 12 three adjacent C-H groups. The vibrational modes of the three C-H groups are strongly correlated, 13 and usually four bands are observed, the first two at ca. 750 and 800 cm⁻¹ (in some cases the 800 cm⁻¹ 14 ¹ band appears as a shoulder of the former), the third at 850 and the forth at 950 cm⁻¹. The relative 15 intensities of the bands appear to be affected by the typology of adjacent extended borders (either Z 16 or A). Finally, a quatro corner (Figure 6c) consists in a benzene ring exposing four adjacent C-H 17 sites. The spectra of Q corners display four bands centred at about 740, 840, 930 and 960 cm⁻¹ and 18 are almost invariant with the model. 19

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1 Figure 6. Representative contributions of different types of C-H corners, belonging to different models: *duo* (part a), *trio*

2 (part b) and *quatro* (part c).

3 4.3. Linear combination fit with regular C-H components

To quantify, at least in relative terms, the abundance of the different C-H terminal geometries in a 4 5 certain sample, we fitted the experimental INS spectra shown in Figure 1 with a linear combination 6 of the simulated spectra discussed in section 4.2. Due to the limited number of experimental points 7 compared to the number of possible C-H contributions from models, we had to limit the number of 8 included functions. Consequently: 1) all the simulated spectra belonging to the same C-H termination 9 category were averaged, reducing the number of possible contributions to eight cases; 2) T and Q corner sites were excluded from the fit analysis because they display too strong contributions in the 10 730-800 cm⁻¹ region, where the experimental spectra are very weak; 3) the contributions of extended 11 zigzag borders (Z1T, Z2 and Z3-4) always resulted in negligible values (much lower than 1%) when 12 performing preliminary linear combination fit tests, and were thus excluded. Therefore, the 13 14 experimental spectra were fitted with a linear combination of the simulated functions for Z1D, A and 15 D terminations. The intensities of the spectra introduced in the linear combination were then 16 normalised to account for a single C-H group.



17

Figure 7. Result of the linear combination fit for the experimental spectrum of CwA (black) obtained by including the simulated functions of Z1D, A and D terminations. The best fit solution is reported in yellow, and the percentages of each component of the fit in the final result are reported in table form. The goodness of fit is evaluated by the RMSE.

Figure 7 shows the best fit (yellow) for the experimental spectrum of CwA (black) and the three different C-H contributions, complemented by a table summarising the resulting relative abundance of the C-H sites. Its Root Mean Square Error (RMSE) expresses the goodness of the fit. Although the three main bands characterizing the experimental spectrum are well reproduced in terms of position, they are not reproduced in terms of relative intensities. In particular, if we renormalize the fitted and experimental spectra to the intensity of the band at 800 cm⁻¹, the bands at 850 and 955 cm⁻¹ are missing some intensity as well as additional contributions around 830 and 930 cm⁻¹ (i.e. where the experimental spectrum shows two minima) are not accounted for. Analogous results were obtained for all the other samples. Hence, it becomes clear that the fingerprint region of the INS experimental spectrum of a physically activated carbon cannot be fully reproduced by considering only regular C-H terminations.

8 4.4. DFT simulations of the INS spectra for physical defects

9 Since activated carbons are by definition disordered and defective materials, the missing components 10 in the fit reported in Figure 7 might be associated with defects. Defects may arise either because of 11 the presence of heteroatoms (i.e. functional groups), or because of ruptures of the regularity in the 12 aromatic domains (physical defects). Since physically activated carbons display a very low 13 concentration of functional groups [17, 42], we focused on the latter class of defects. Although the 14 number of possible physical defects is extremely high, we decided to focus on those containing C-H 15 groups and hence detectable by INS.

First, the aromatic models can simply include some disordered C-H terminations, hardly 16 17 classifiable as any of the regular borders and corners classes described in section 4.2. For these disordered C-H terminations, the greatest variations in the INS spectra are observed for single zigzag 18 sites, for which some examples are reported in Figure 8a, first plot. These terminations give a quite 19 broad band centred at ca 830 cm⁻¹, in correspondence with one of the minima of the experimental 20 spectra. Disordered duo terminations instead (second plot in Figure 8a) give two main bands, 21 similarly to the spectra of D and A terminations (Figure 5c and Figure 6a), but in this case the two 22 bands fall at ca. 800 and 935 cm⁻¹, about midway between the A and D regular cases. Generalizing, 23 we can conclude that the INS fingerprint of terminations consisting in two C-H exposed at the same 24 25 benzene ring fall within a continuum of cases, of which D and A sites correspond to the two extremes. Finally, T and Q corners are always inserted in peripheral positions in these models and hence they 26 are not affected by disorder and/or defectivity of the model. 27

We also considered holes within the aromatic platelets and saturated by H atoms, since they have been recently demonstrated to occur in defective graphene [23]. Due to the limited size of the platelets in the activated carbons under investigation [17], only small holes can be imagined. Thus, we have focused on holes created by removing a single C atom from the model forming a five member ring and a C-H group (Figure 8b), which we labelled as a single vacancy. A distinct band in the fingerprint region at 950 cm-1, characterizes the simulated INS spectrum of this termination inside the hole. A third kind of possible irregular termination consists in a single C-H site completely isolated from other C-H groups on adjacent benzene rings (Figure 8c). In this case, the considered C-H groups are isolated from their neighbours by the presence of four pentagonal rings. In this case, the out-ofplane bending mode of the C-H group has a very little tendency to combine with the vibrations of other C-H groups, resulting in a single and sharp band centred at ca 880 cm⁻¹.

Activated carbons were also proposed to feature topological defects that introduce local 6 7 curvature and sphericity within the graphenic planes [43]. We thus decided to simulate two additional models featuring the same kind of C-H terminations as model C₅₄H₁₈ but with different curvatures, 8 9 designed by introducing either a five or a seven-membered ring at the centre of the model. The analysis of their simulated spectra, shown in Figure S8 in the supporting information, immediately 10 shows that the spectral features are almost undistinguishable from the corresponding flat model. This 11 indicates that this region of the INS spectra is mostly sensible to the local geometry of the C-H 12 13 terminations, and much less on defects not directly involving H atoms.

Three further simulations were run on models in which the source of the curvature (i.e. introduction 14 15 of a 5-member ring) is located closer to the borders of the model, and so where it can directly affect the nearby C-H terminations. The characteristic INS spectral features of the three models are reported 16 17 in Figure S9 in supporting information. The first two models where build starting from the hexagonal C₅₄H₁₈ model, by adding five member rings at two distinct positions close to the border. For these 18 models (Figure S9A and B), the Z1 and D terminations in close proximity with the five member rings 19 show bands in the same frequency ranges and with similar intensities as the previously introduced 20 disordered Z1 and D terminations on flat models of Figure 8a. Consequently, these curved models 21 contribute to the disordered terminations similarly to the flat ones and we cannot distinguish in 22 between them. The third model contains a five membered ring exposing one C-H termination. Its 23 characteristic INS spectrum (Figure S9C) shows the most intense features in the 600-750 cm⁻¹ range, 24 in correspondence with one of the minima of the experimental spectra. This observation makes us 25 conclude that this specific termination is not expected to be significantly present in our samples. 26



Figure 8. Disordered and defective models considered in this work and the characteristic contributions to the INS spectra of their defective C-H terminations. The sites that behave differently than the regular ones are indicated by a circle. The reported simulated INS spectra refer to the species highlighted with a full green circle. Part a): Models containing disordered terminations, where disordered Z1 sites are circled in dark green and disordered D sites in light green. A representative spectrum for both disordered Z1 and D sites is reported. Part b): Models containing a Single Vacancy defect saturated by one H atom, and a representative spectrum of their INS contribution. Part c): Model containing an isolated Z1 site, obtained by introducing four five membered rings and curvature in the model, and characteristic INS spectrum of this site.

10 **4.5. Introduction of physical defects in the fit**

With the simulations of the physical defects in the hands, we came back to the linear combination 11 analysis and added these new contributions to improve the fit result. To ensure the reliability of the 12 result we limited again as much as possible the number of functions included in the fit, and thus the 13 number of parameters to be optimised. At first, we decided to describe the D, A and disordered D 14 terminations with a single function (D-all), obtained by averaging their characteristic spectra. Then, 15 three new functions were added to the fit, accounting for disordered Z1, isolated Z1 and single 16 vacancy defects, obtained by averaging all the spectra belonging to the same category of C-H 17 18 terminations.

The new linear combination fit was therefore performed by including five functions Z1D, Dall, disordered Z1, isolated Z1 and single vacancy, and imposing the following constrains: 1) the concentration of regular terminations should be higher than that of irregular and defective sites; 2) for the catalysts, the values of the coefficients in the fit should be equal or lower than those obtained

1 in the fit of the parent carbon. These choices were justified by the observation that the three main peaks in the experimental spectra show a better matching with the spectra of regular terminations, 2 and that the intensity of the INS spectra for the catalysts is always less intense than for the parent 3 carbons. The result of this fit for CwA is shown in Figure 9. The same procedure was followed for 4 5 all the other activated carbons and catalysts, resulting in the final values reported in Table 3. For the four activated carbons, all the coefficients obtained from the fit were normalised so that their sum 6 gives 100%, while all the results of the catalysts were normalised with respect to the parent support, 7 8 in order to allow a direct evaluation of the loss of hydrogen at the graphenic terminations as a 9 consequence of the metal deposition. The plots showing the results of the linear combination fit analysis for the other samples are reported in Figure S10, S11, S12 and S13 in the supporting 10 information. 11



12

Figure 9: Result of the linear combination fit for the sample CwA. The best result was obtained by introducing in the linear combination the functions for Z1D, D-all, disordered Z1, isolated Z1 and single vacancy sites. An example of a C-H termination for each function in the fit is reported for clarity.

16 Table 3: Summary of the results of the linear combination fit on the experimental INS spectra of all the activated carbons 17 and catalysts. The values for the activated carbon supports are normalised so that their sum gives 100%, while those of 18 the catalysts were normalised with respect to the parent support. The goodness of the fit is evaluate by the RMSE.

	Z1D %	D-all %	Disordered Z1 %	Single vacancy %	Isolated Z1 %	тот	RMSE
CwA	25	38	19	12	5	100	0.001
Pt(R)/CwA	20	36	19	12	5	92	0.001
Pd(NR)/CwA	15	32	17	11	4	78	0.001
Pd(R)/CwA	13	32	19	11	4	79	0.002
CwB	28	33	23	11	5	100	0.001
Pt(R)/CwB	15	29	11	9	4	67	0.001
СрА	24	42	22	10	2	100	0.001
Pt(R)/CpA	16	35	18	9	2	78	0.001
СрВ	27	41	20	11	2	100	0.001

Pt(R)/CpB	18	36	19	11	2	86	0.001
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2 As far as the bare carbons are concerned, the main C-H contributions are relatively similar for all the considered samples. In all the cases the experimental spectra are well reproduced by using 3 mostly Z1D, D-all and disordered Z1 sites, followed by lower amounts of single vacancies, while the 4 contribution of isolated Z1 sites is almost negligible. For all the four carbons, the main contribution 5 is given by D-all sites, followed by Z1D and disordered Z1. The D-all sites are much more frequent 6 in the activated carbons of peat origin than in those of wood origin. Among carbons with the same 7 origin, the greatest difference is observed between CwA and CwB, and it is mostly reflected by the 8 9 different D-all/Z1D ratio (1.5 and 1.2, respectively).

10 The deposition of Pd or Pt nanoparticles causes a small decrease in the total intensity of the spectra, as already observed in our previous experiment [5]. The result of the linear combination fit 11 12 analysis allows us to evaluate which are the most affected C-H species: in fact, albeit the amount of all C-H sites decreases with respect to the parent carbon, the greatest variation is always observed for 13 14 Z1D sites. For CwA, the perturbation caused by the deposition of Pt nanoparticles (sample Pt(R)/CwA) is relatively modest, with the greatest change observed for Z1D sites which decreased 15 16 by about 5%. The deposition of Pd nanoparticles, instead, causes a much larger reduction of the regular C-H terminations, mostly evident for Z1D sites which decrease by about 10% in 17 18 Pd(NR)/CwA, and by 12% for Pd(R)/CwA. These data also clearly indicate that the cause of the 19 observed H loss in the catalysts is the deposition of the metal, and not the successive pre-reduction step. The most significant change was observed when depositing Pt nanoparticles on the CwB 20 support, which resulted in a decrease of about 13% for Z1D sites and of about 12% for disordered Z1 21 sites. The effect of the Pt deposition is instead quite similar for the two carbons of peat origin, where 22 the Z1D sites decreased of ca. 8-9% and the D-all decrease by about 5-7 %. 23

Even if the matching between the fitted solution and the experimental spectra increased significantly with the introduction of physical defects, it is evident for all the samples that the fits are still missing some contribution at about 920 cm⁻¹, in correspondence with one of the minima of the experimental spectra. None of the defective models that we have simulated so far seems able to cover that region. Further simulations and experiments in this direction are planned.

29 **5. Conclusions**

In this work, we present the high-resolution INS spectra of four different activated carbons and of six industrial catalysts obtained by depositing Pd or Pt nanoparticles on the same carbons. The INS spectra contain information on the type and amount of the C-H terminations at the graphenic domains constituting the carbon. Albeit the investigated carbon supports show a certain degree of structural, porous and morphological heterogeneity, the unpreceded high signal-to-noise ratio of our spectra in comparison with past measurements [17], allowed us to highlight tiny differences among carbons having different origins, and between the catalysts and the corresponding carbon supports. In particular, we observed that the deposition of the metal nanoparticles causes a general decrease of the C-H signals, which for Pt scales linearly with the metal dispersion: smaller the metal nanoparticles, larger is the amount of C-H affected by the metal deposition.

7 In order to interpret the differences observed in the experimental spectra, an extensive DFT 8 study was performed. At first, we simulated the INS spectra of several aromatic models characterized 9 by regular C-H terminations and/or by the presence of some physical defects. Successively, we determined the characteristic INS fingerprint of different C-H terminations. The experimental INS 10 spectra were systematically fitted with a linear combination of five functions corresponding to: 1) C-11 H at regular zigzag borders and vicinal to a duo corner (Z1D); 2) C-H at regular armchair borders, at 12 13 duo corners or at disordered duo terminations (D-all); 3) C-H at disordered zigzag borders (disordered Z1); 4) isolated C-H (isolated Z1); and 5) C-H in single vacancies inside the graphenic platelets. We 14 found out that, although regular aromatic models are already able to reproduce the position of the 15 experimental features, physical defects and disordered C-H species are necessary to obtain a good 16 17 match between experimental and fitted spectra. We were not able to identify a general spectral fingerprint for the curvature and not planarity of the graphenic domains, but two of the defective 18 contributions identified by linear combination analysis (Isolated Z1 and Single Vacancy) were indeed 19 given by not planar models. Thus, it does not appear possible to describe in full the platelets structure 20 within activated carbons considering flat geometries only. 21

The results of the fits indicate that the most frequent C-H terminations in all the activated 22 carbons are D-all, followed by Z1D and disordered Z1, while the other two types of defects are much 23 less frequent. The relative abundance of Z1D terminations is lower for carbons of peat origin (Cp) 24 25 than for those of wood origin (Cw). Since the metal dispersion for catalysts prepared on Cp carbons is systematically lower than that on Cw carbons, these results suggest that the metal dispersion is 26 improved on carbons having a large percentage of Z1D terminations. The decisive role of Z1D sites 27 28 is confirmed by the observation that the metal nanoparticles deposition mainly affects this typology of C-H termination. These outcomes are particularly relevant for catalysis, since they offer for the 29 30 first time an explanation of the relationship between the type of the carbon support and the metal 31 dispersion. On a more general ground, this study demonstrates that high resolution INS spectra 32 complemented by a systematic DFT study are invaluable tools for characterizing carbon-based 33 materials.

34

1 Acknowledgments

- 2 This work was strongly encouraged from the beginning by our beloved Carlo Lamberti and Giuseppe
- 3 Leofanti, to whom we will be infinitely grateful for having guided us with passion and enthusiasm.
- 4 We are also grateful to ILL for the time we have accounted on the computational cluster.

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