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Dynamic Nuclear Polarization Enhancement of 200 at 21.15 T Enabled by 65 kHz Magic Angle Spinning

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For value in analytical chemistry, sensitivity and spectral
resolution are cornerstone parameters of solid-state
nuclear magnetic resonance (SSNMP) experiments Dynamic nuclear magnetic resonance (SSNMR) experiments. Dynamic nuclear polarization (DNP) is a powerful technique to dramatically increase the sensitivity of NMR by transfer of electron spin polarization to nearby nuclear spins upon microwave (μw) irradiation.^{[1](#page-3-0)} In solid-state NMR, the combination of high magnetic field and fast magic angle spinning (MAS) is often a requirement to allow the study of complex systems and to implement sophisticated multidimensional experiments. State-of-the-art commercial NMR instrumentation uses magnetic fields of up to 28 T (1.2 GHz) and MAS probes that can spin samples to 120 kHz at room temperature using 0.7 mm diameter rotors. Such spinning regimes significantly improve ¹H NMR resolution on samples from small crystalline pharmaceuticals to fully protonated protein assemblies. 2^{-4} 2^{-4} 2^{-4} 2^{-4} There is therefore tremendous interest in combining these approaches with DNP. DNP MAS experiments are more demanding as they usually require cryogenic temperatures $(\approx 100 \text{ K})$ to obtain reasonable saturation of the electron spin transition. In 2016, Chaudhari et al. reported DNP MAS experiments at 40 kHz MAS rates using 1.3 mm rotors, 5 where faster MAS led to significantly longer coherence lifetimes in ²⁹Si spectra. They then discovered that, very importantly, under certain conditions DNP enhancements increased significantly as the sample spinning rate increased. $6,7$ $6,7$ $6,7$ They also showed that fast spinning yielded ²⁷Al spectra that were free of sidebands or spectral distortions. Wang et al. have shown the possibility of using $^1\mathrm{H}$

detection in this context.⁸ In particular, ¹H-detected ¹H-⁸⁹Y DNP HETCOR at 40 kHz MAS surpassed the sensitivity of DNP HETCOR recorded using 3.2 mm rotors at a lower spinning rate.

Here we present the first DNP MAS experiments using a 0.7 mm MAS probe. This allows us to reach spinning frequencies of 65 kHz, with μ w irradiation, at 100 K and at a field of 21.15 T. We report the performance, at very fast MAS, of BDPA in 95% deuterated orthoterphenyl (OTP- $d_{95\%}$) and HyTEK2 in 1,1,2,2-tetrachloroethane (TCE). In both, we find that enhancements continue to increase significantly with increasing spinning rates, and we obtain enhancement factors of up to 200 for HyTEK2. Beyond the good DNP performance, we also report the highest directly detected ¹H resolution in DNP MAS SSNMR spectra observed so far.

In 2014, Can et al. used the carbon-centered radical BDPA in a polystyrene matrix to obtain significant proton DNP enhancements (ε_H) using either Overhauser effect (OE) or solid effect (SE). Interestingly, OE DNP scaled positively with the external magnetic field from 9.4 T (ε_{H} = 14) to 18.8 T (ε_{H}) = 20),^{[9](#page-4-0)} and later enhancements of ε_{H} = 65 at 9.4 T and ε_{H} =

Received: August 15, 2020 Accepted: September 10, 2020 Published: September 22, 2020

Figure 1. $^1\mathrm{H}$ DNP enhancements measured for a 60 mM frozen solution of BDPA in OTP- $d_{95\%}$ at MAS rates of up to 65 kHz. Measurements were taken with polarization delays of 10 s (blue) and 75 s (orange) at each spinning frequency. The error bars are calculated from the estimated error based on the signal-to-noise ratio of the μ w ON and OFF spectra (see the [Supporting Information](http://pubs.acs.org/doi/suppl/10.1021/acs.jpclett.0c02493/suppl_file/jz0c02493_si_001.pdf) for details). The blue and orange lines correspond to the simulated DNP enhancements in both cases, based on the source-sink relayed DNP model (see the [Supporting Information](http://pubs.acs.org/doi/suppl/10.1021/acs.jpclett.0c02493/suppl_file/jz0c02493_si_001.pdf) for details regarding the simulations).^{[6](#page-4-0)} The OTP and BDPA structures are shown.

Figure 2. ¹H DNP enhancements (blue) and ¹H DNP buildup times in seconds (purple) for a 32 mM solution of HyTEK2 in 1,1,2,2 tetrachloroethane (TCE) as a function of the MAS rate. The error bars of the ¹H enhancement are estimated from the signal-to-noise ratio of the μ w ON and OFF spectra, and the error bars of the buildup times are the error of the fit (see the [Supporting Information](http://pubs.acs.org/doi/suppl/10.1021/acs.jpclett.0c02493/suppl_file/jz0c02493_si_001.pdf) for details). HyTEK2 and TCE molecules are represented in the figure.

75 at 18.8 T were reported by using a 95% deuterated orthoterphenyl $(OTP-d_{95\%})$ matrix.^{[10](#page-4-0)} Later, an enhancement ε_{H} = 73 was reported using a BDPA/OTP- $d_{95\%}$ mixture at 21.15 T.[11](#page-4-0) All of these experiments were performed at moderate sample spinning rates using 3.2 mm rotors at between 8 kH $z^{9,10}$ and 12.5 kHz.^{[11](#page-4-0)} In 2017, OE DNP with BDPA in OTP- $d_{95\%}$ was investigated at fast magic angle spinning using 1.3 mm rotors spinning at ≤40 kHz in an 18.8 \hat{T} magnet,⁷ where it was reported that the ¹H enhancement ramped up from 23 at 0.5 kHz MAS to 105 at 40 kHz MAS. Thus, OE DNP with BDPA in OTP- $d_{95\%}$ at both high field and fast MAS appeared to be a system of choice for performing

DNP measurements using very fast MAS in 0.7 mm rotors at 21.15 T. At 100 K, the 0.7 mm rotors can stably spin at rates of up to 65 kHz. Although 0.7 mm rotors can spin stably up to 120 kHz at room temperature, at 100 K fluid dynamics of the cold nitrogen gas limit the maximum rate. Notably, the density of N_{2(g)} is ~3 times higher at 100 K than at 298 K, and the speed of sound is slower by a factor $2.^{12}$ (In principle, this could be mitigated by, for example, spinning using helium gas, but this is technically not straightforward.) Spinning stability at 65 kHz and 100 K is usually very reliable and stable in the range of \pm 9 Hz (see [Figure S4\)](http://pubs.acs.org/doi/suppl/10.1021/acs.jpclett.0c02493/suppl_file/jz0c02493_si_001.pdf).

Figure 1 shows the ${}^{1}H$ DNP enhancement of a 60 mM solution of BDPA in OTP- $d_{95\%}$. The sample was prepared following the method described in ref [10.](#page-4-0) Following on from the observations made at \leq 40 kHz, here we see that the 1 H

Figure 3. (a) 900 MHz DNP enhanced ¹H spectra at different MAS rates between 10 and 60 kHz of $[U^{-13}C_1^{15}N]$ histidine·HCl·H₂O impregnated with 32 mM HyTEK2 in TCE. The spectra were all recorded with a 3 ms total spin echo prior to acquisition for solvent suppression.^{[45](#page-5-0)} Note that the spectrum at 10 kHz is dominated by the solvent signal. (b) $\mathrm{^{1}H-^{13}C}$ and (c) $\mathrm{^{1}H-^{15}N}$ $\mathrm{^{1}H-detected}$ DNP HETCOR spectra of the same sample spinning at 62.5 kHz. For the sake of clarity, the signal-free region between 70 and 150 ppm of the $^{15}{\rm N}$ dimension of panel c is not shown. H2/3 signals, clearly visible on panel b, are not visible in panel a as they do not survive the spin echo used for solvent suppression. See the [Supporting Information](http://pubs.acs.org/doi/suppl/10.1021/acs.jpclett.0c02493/suppl_file/jz0c02493_si_001.pdf) for detailed acquisition parameters.

enhancement continues to increase as the MAS frequency increases to 65 kHz, reaching a maximum of $\varepsilon_{\rm H}$ = 106. Although [Figure 1](#page-1-0) reports a maximum enhancement of 106, we have been able to reach ε_H = 128 at 65 kHz MAS with a different sample. We attribute this difference to the quality of glass formation upon freezing the sample, which is a recurrent feature in many DNP matrices, including OTP.[13](#page-4-0)−[16](#page-4-0) At 65 kHz MAS, we measure a buildup time $(T_{b,on})$ of 58 s, and the enhancements were determined with a polarization (recycle) delay of 1.3 \times T_{b,on} = 75 s (orange dots in [Figure 1\)](#page-1-0). We also measured the $^1\mathrm{H}$ enhancements with a polarization delay of 10 s and observed enhancements consistently lower than for the longer recycle delay (blue dots in [Figure 1\)](#page-1-0). This is a clear signature of H ¹H $-$ ¹H spin diffusion relaying hyperpolarization among the ${}^{1}H$ nuclei in the OTP matrix.^{[17,18](#page-4-0)} It is possible to rationalize this behavior using the source-sink spin diffusion model proposed in ref [7.](#page-4-0) This model hypothesizes that the detected DNP behavior results from the flow of hyperpolarization from the BPDA/OTP- $d_{95\%}$ source matrix to dilute sinks (e.g., paramagnetic impurities). Using this model, the enhancements with polarization delays of 10 and 75 s were simulated across the whole range of spinning rates from 0 to 65 kHz. The simulated enhancements are reported in [Figure 1](#page-1-0) (plain lines) and fit very well to the experimental enhancement data. The details of the simulation model are given in the [Supporting Information.](http://pubs.acs.org/doi/suppl/10.1021/acs.jpclett.0c02493/suppl_file/jz0c02493_si_001.pdf) The spin diffusion model explains both why $\varepsilon_{\rm H}$ increases when the MAS rate increases at a constant polarizing delay and why for a given MAS rate $\varepsilon_{\rm H}$ increases when the polarization delay increases from 10 to 75 s. The results give additional very strong support to the original source-sink model.

The temperature of the sample in the rotor was monitored at each MAS rate using the T_1 ⁽⁷⁹Br) of a KBr crystal added to the rotor.^{[19](#page-4-0)} The internal temperature under μ w irradiation increases from 95 K at 10 kHz MAS to \approx 105 K at 65 kHz MAS, indicating that the frictional heating induced by MAS in the 0.7 mm DNP LTMAS probe is only around 10 K. The contribution factor has not been remeasured here as it was previously found to be 1 up to 40 kHz MAS.^{[7](#page-4-0)}

Although OE DNP using BDPA provides a large ¹H DNP enhancement, the long $^1\mathrm{H}$ buildup time prevents it from being an efficient polarizing system. In some cases, flip-back $11,20$ or frequency selective^{[21](#page-4-0)} approaches should be applicable. The rational design of radicals performing OE DNP is still very challenging, which slows development of better OE $DNP²²$ $DNP²²$ $DNP²²$ While there is great interest in \overline{OE} DNP,^{[23](#page-4-0)} it does not appear here to be the polarizing system of choice under these conditions.

On the contrary, cross-effect (CE) DNP has been extensively studied and many design parameters have been discovered^{[15](#page-4-0),[24](#page-4-0)−[29](#page-4-0)} and theorical models proposed,^{[30](#page-4-0)−[38](#page-4-0)} allowing us today to tailor the structures of CE polarizing agents and thus optimize their performance. Unfortunately, CE polarizing agents that perform very well at moderate magnetic fields (e.g., AMUPOL) scaled unfavorably when going to higher magnetic fields (see [Table S1](http://pubs.acs.org/doi/suppl/10.1021/acs.jpclett.0c02493/suppl_file/jz0c02493_si_001.pdf)).^{[6](#page-4-0),[25,](#page-4-0)[39](#page-5-0)} Capitalizing on the work of Hu et al.,^{[32](#page-4-0)} Mathies et al. introduced the concept of mixed biradicals in 2015 in which a TEMPO-like moiety was tethered to a trityl radical.[40](#page-5-0) In particular, the performance of the resulting TEMTriPol-1 scaled favorably with the magnetic field between 5 and 14.1 T, yielding ${}^{1}H$ enhancements of 60 at 5 T, 87 at 14.1 T, and 65 at 18.8 T, which was attributed to the strength of the exchange interaction between the trityl radical and the

nitroxide.[40](#page-5-0) Capitalizing on the previous concept of decorated binitroxides to provide longer electron relaxation times, Wisser et al. proposed a series of mixed BDPA−nitroxide biradicals and evaluated their performance at both high field (up to 21.15 T) and fast MAS (up to 40 kHz using a 1.3 mm rotor). 6 In 6 In particular, HyTEK2 yields ¹H enhancements of 185 at 18.8 T and 40 kHz MAS in a 1.3 mm rotor which was, so far, the highest DNP enhancement reported at high magnetic field and fast MAS.⁶

[Figure 2](#page-1-0) plots the 1 H DNP enhancement and the 1 H DNP build-up time as a function of the MAS rate obtained with a 32 mM HyTEK2 solution in TCE. We measured $\varepsilon_H = 147$ at 10 kHz, which then increases almost linearly to 200 at the maximum MAS rate of 65 kHz.

At the same magnetic field shown in a 3.2 mm rotor, enhancements of 64 at 10 kHz were previously reported. We attribute the improvement of the performance between 0.7 and 3.2 mm rotors at the same spinning speed to better penetration of the μ w. An analogous difference was observed previously between 3.2 and 1.3 mm rotors at the same MAS rate.^{5,[6](#page-4-0)} We also made the same observation with a 10 mM AMUPOL in a $6/3/1$ (v/v/v) glycerol- $d_6/D_2O/H_2O$ solution, and the results are listed in [Table S1](http://pubs.acs.org/doi/suppl/10.1021/acs.jpclett.0c02493/suppl_file/jz0c02493_si_001.pdf).

As reported in ref [6](#page-4-0), the combination of fast MAS and high magnetic field with the mixed radical nitroxide-BDPA HyTEK2 allows us to reach unprecedented DNP performance at this magnetic field; the increase with MAS rate is again attributed to decoupling of the polarizing bath from polarization sinks, as for the BDPA/OTP- $d_{95\%}$ example mentioned above. We note that prior to measurement, the sample was degassed of dissolved O_2 using freeze−thaw cycles in the probe.^{[41,42](#page-5-0)} This is always necessary when using TCE as a DNP solvent; degassing here allows us to increase the ¹H DNP enhancement (measured at 20 kHz MAS) by 66% after four freeze−thaw cycles. The temperature of the sample was monitored on the basis of longitudinal relaxation rate T_1 ⁽⁷⁹Br) of a KBr crystal incorporated in the sample¹⁹ (see [Figure S2](http://pubs.acs.org/doi/suppl/10.1021/acs.jpclett.0c02493/suppl_file/jz0c02493_si_001.pdf) for details). Similarly, to what we observed above with BDPA/ OTP- $d_{95\%}$, very fast MAS induces sample heating of only ~10 K, and here in addition μ w irradiation induces another 10 K heating due to the use of TCE that absorbs some μ w. The contribution factor has not been remeasured here; it was previously found to be constant at 0.71 between 5 and 40 kHz MAS ^{[6](#page-4-0)}, and we believe this value can be extrapolated to 65 kHz MAS. In total, our data demonstrate that the HyTEK2 radical is a very efficient polarizing agent for very fast MAS DNP at high field. We were able to reach a ¹H DNP enhancement of 200, which is a new record for both a magnetic field of 21.15 T and the HyTEK2 polarizing agent.

One of the most promising outcomes for NMR with fast MAS (60 kHz and higher) is the achievable ${}^{1}H$ resolution, which allows one, for example, to record directly or indirectly ¹H NMR spectra without the use of sophisticated ¹H homonuclear decoupling schemes.^{[43](#page-5-0)} In DNP MAS, reaching such regime was not possible until now as hardware was limited to \leq 40 kHz. Here, as a proof of concept, we present ¹H-detected DNP MAS experiments at 65 kHz MAS on a microcrystalline powder of $\overline{[U^{-13}C, ^{15}N]}$ histidine hydrochloride impregnated with 32 mM HyTEK2 in TCE, at 100 K. In the impregnation approach, the histidine does not dissolve in TCE and the hyperpolarization diffuses from the frozen HyTEK2/ TCE phase to the histidine crystal through spontaneous spin diffusion, as is now well established.^{17,[44](#page-5-0)} HyTEK2/TCE was

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chosen between the two DNP systems benchmarked here as it provides the highest sensitivity gain (in terms of DNP enhancements and build-up times). [Figure 3a](#page-1-0) shows DNPenhanced ¹H spectra at different MAS rates. The strong TCE signal was partially suppressed using a spin echo of 3 ms prior to acquisition.⁴⁵ At 10 kHz MAS, this method of solvent suppression was not efficient enough, and the spectrum is actually dominated by the TCE signal. One can observe both the significant increase in sensitivity and spectral resolution as the spinning frequency is increased. In addition to the increase in the DNP enhancement with HyTEK2 at higher MAS rates, these observations are related to the increase in the $^1\mathrm{H}$ coherence lifetime, the decrease in the ¹H line width, and the decrease in the number of sidebands as the spin rate increases.[43,46](#page-5-0) The duration of the spin echo for solvent suppression was kept constant (at 3.0 ms) between the different MAS rates to compare the intensity between the spectra. The enhancement of the histidine signals in the 20 kHz spectrum is 41 (note that the sample was not degassed). The enhancement was not measured at 65 kHz for this sample but was previously measured to be ε_H = 25 at 35 kHz for the histidine signals and ε_H = 35 at 65 kHz. The slight asymmetry of the ¹ H signals is attributed to shimming issues, probably due to the current probe design, which includes an extension of the waveguide into the stator to improve overall performance.⁴

[Figure 3b](#page-1-0) shows the two-dimensional (2D) DNP-enhanced H-detected ¹ H−13C HETCOR spectrum and [Figure 3](#page-1-0)c ¹H−¹⁵N HETCOR spectrum of impregnated [U-¹³C,¹⁵N]histidine at 62.5 kHz MAS. The HETCOR spectra are fully resolved and allow all of the resonances to be assigned.

In conclusion, we have reported DNP MAS experiments using 0.7 mm rotors spinning at 65 kHz at a magnetic field of 21.15 T. The enhancements measured for the samples studied here increase with increasing MAS rates. This allows us to obtain the highest enhancements measured at this magnetic field so far. In particular, we found that HyTEK2 can yield a ¹H DNP enhancement of 200 at the maximum spinning frequency. We then showed how reaching these spin rates under DNP conditions opens up the possibility of using $^1\mathrm{H}$ detected spectra more widely in MAS DNP, and as an example, we obtained resolved DNP-enhanced ¹H-detected 1D and 2D spectra.

■ ASSOCIATED CONTENT

6 Supporting Information

The Supporting Information is available free of charge at [https://pubs.acs.org/doi/10.1021/acs.jpclett.0c02493](https://pubs.acs.org/doi/10.1021/acs.jpclett.0c02493?goto=supporting-info).

Details of experimental parameters, spin diffusion simulation details, and the DNP performance of the AMUPOL polarizing agent ([PDF\)](http://pubs.acs.org/doi/suppl/10.1021/acs.jpclett.0c02493/suppl_file/jz0c02493_si_001.pdf)

The MatLab code used to perform simulations and the NMR raw data have been deposited at: [http://doi.org/](http://doi.org/10.5281/zenodo.4032323) [10.5281/zenodo.4032323.](http://doi.org/10.5281/zenodo.4032323)

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Notes

The authors declare no competing financial interest.

■ ACKNOWLEDGMENTS

P.B., A.B., and G.S. acknowledge Innosuisse, the NCCR Chemical Biology, and the Marie-Skłodowska-Curie program (Grant 796904), respectively, for support. This work was supported by Swiss National Science Foundation Grant 200020_178860.

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