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Intense Reactivity in Sulfur-Hydrogen Mixtures at High Pressure under X-ray Irradiation

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Abstract

Superconductivity near room-temperature in the sulfur-hydrogen system arises from a se-3 quence of reactions at high pressures, with x-ray diffraction experiments playing a central role 4 in understanding these chemical-structural transformations and the corresponding S:H stoi-5 chiometry. Here we document x-ray irradiation acting as both a probe and as a driver of chem-6 ical reaction in this dense hydride system. We observe reaction between molecular hydrogen 7 (H_2) and elemental sulfur (S_8) under high pressure, induced directly by x-ray illumination, at 8 photon energies of 12 keV using a free electron laser. The rapid synthesis of hydrogen sulfide q (H_2S) at 0.3 GPa was confirmed by optical observations, spectroscopic measurements, and 10 micro-structural changes detected by x-ray diffraction. These results document x-ray induced 11 chemical synthesis of superconductor-forming dense hydrides, revealing an alternative produc-12 tion strategy and confirming the disruptive nature of x-ray exposure in studies on high-pressure 13 hydrogen chalcogenides, from water to high-temperature superconductors. 14

J5 Graphical TOC Entry



The behavior of simple hydrides under extreme conditions is of fundamental importance to 17 achieving superconductivity at high temperatures, ^{1–6} with chemical bonds between hydrogen atoms 18 (H) and heavier elements favoring high phonon frequencies. Recently, the high-pressure forma-19 tion of H₂S (and H₂Se) from constituent elements has been of substantial interest, following early 20 predictions ^{7,8} and then observation of a high superconducting transition temperature (T_c) of 203 21 K in the S-H system at pressures above 155 GPa,^{2,9-13} with similar properties predicted in the 22 analogous selenium-hydrogen (Se-H) system.^{14,15} The experimentally observed high T_c in the 23 hydrogen-sulfur system is considered to occur in H₃S, a high-pressure phase containing a high 24 content of hydrogen.^{2,8,16,17} H₃S ($T_c = 204$ K) is predicted to form above 111 GPa from the re-25 action between H₂ and H₂S, upon compression of the stoichiometrically equivalent, molecular 26 van-der-Waals compound (H₂S)₂H₂.^{8,18} This compound (and structurally analogous (H₂Se)₂H₂) 27 itself forms above 4 GPa upon room-temperature compression of H₂ and H₂S (H₂Se).^{11–13,18} For-28 mation of the initial H₂ and H₂S (H₂Se) mixtures is facilitated by directly reacting the constituent 29 elements S (Se) and excess H₂ at high pressures within diamond anvil cells, a sluggish process 30 near room temperature.^{11–13} 31

Hard x-ray experimental probes, in particular x-ray diffraction (XRD), have been crucial in the characterization of the structures, reactions, stoichiometries and high-pressure phase transformations of S-H and Se-H systems.^{9,10,12,16,18,19} However, there are notable inconsistencies in the results of the XRD studies within both the S-H^{9,10,16} and Se-H^{12,13} systems, though the reasons for these differences are currently unclear.

In fact, x-ray exposure itself is observed to cause material changes within the electronically analogous oxygen-hydrogen (O-H) system, where irradiation of H₂O by hard x-rays of intermediate energy (\sim 10 keV) induces a chemical transformation to an O₂-H₂ alloy.^{20,21} This raises the prospect that x-ray exposure could have a chemical effect on other chalcogen (group 16 element) hydride systems under pressure, yet this question remains unexplored. Additionally, the effect of the x-ray source intensity on reaction rates must be explored: while prolonged synchrotron x-ray irradiation of the O-H system induced a reaction over several hours (order 10³ s),²⁰ the high intensity of modern sources such as x-ray free electron lasers (XFELs) may induce more rapid reactions. This may be an essential consideration when taking advantage of the high-brilliance of x-ray pulses from XFEL sources, such as for high-resolution diffraction^{22–24} and chemical dynamics²⁵ studies that are not possible at synchrotrons. XFELs have been demonstrated to be advantageous for the study of very high pressure conditions, thus far using dynamic compression methods,^{22–24} but have not yet been used in conjunction with static compression techniques.



Figure 1: Chemical reactivity induced by XFEL exposure under high pressure. An initial crystalline (orthorhombic) sulfur sample (2 unit cells shown, viewed diagonally between the *a* and *b* axes) adjacent to fluid hydrogen compressed to 0.3 GPa (a) is exposed to a burst of hard x-rays from a free electron laser (b), causing photoelectron and secondary electron damage across the sample disrupting 1-10% of S₈ molecules, leading to rapid etching of the sulfur and formation of phase separated H₂S just after the exposure (c).



Figure 2: Representation of the experimental set-up used at the XSS beamline of the X-ray Free Electron Laser at the Pohang Accelerator Laboratory (PAL-XFEL).^{26,27} The beam is directed horizontally through the axis of the diamond-anvil-cell (DAC). The diffraction detector is placed behind and to the side of the sample at $\sim 45^{\circ}$ to the incident direction of the x-ray beam, to avoid damage from direct contact. The microscope provides *in situ* microscopy of the sample chamber viewed along the DAC axis using an angled mirror, via which the beam is transmitted through a small hole.

Here we present the results of XFEL irradiation of S and H₂ under static compression, and 50 demonstrate an unambiguous high-pressure chemical reaction that is induced (Fig. 1), and probed, 51 by the femtosecond (10^{-15} s) x-ray exposures. The experimental configuration is shown in Figure 2 52 and described in detail in the Supporting Methods. A diamond anvil cell (DAC) was prepared with 53 pieces of crystalline sulfur surrounded by fluid hydrogen at 0.3 GPa. XFEL pulses were passed 54 through the sample cavity, diffracting from the solid portions of the sample. The x-ray intensity 55 was gradually increased; higher powers were sufficient to initiate a visible, rapid reaction between 56 S and H₂, and alter the diffraction patterns indicating physical changes in the sulfur component. 57 The short duration of each pulse means that any physical changes to the sample induced by a given 58 XFEL pulse are observed in the diffraction pattern produced by the following pulse.²⁸ 59 The diffraction patterns showed that the orthorhombic $S_8(\alpha)$ phase persisted from shot to

⁶⁰ The diffraction patterns showed that the orthorhombic $S_8(\alpha)$ phase persisted from shot to ⁶¹ shot (Fig. 3). However, above a given x-ray power, the Debye-Scherrer ring azimuthal intensity



Figure 3: Rietveld refinement of the orthorhombic α -S structure to the integrated profile obtained from the sum of the four diffraction images shown in Figure 4, showing the observed (points) and calculated (line) diffraction patterns, the calculated reflection positions (vertical lines), and the difference profile (lower line). The four 2D diffraction images were summed to help reduce the effects of sample texture. The refined parameters for space group *Fddd* are *a* = 10.380(3) Å, *b* = 12.770(1) Å, and c = 24.421(3) Å (density of 2.105 g cm⁻³); see Supporting Methods for further details.

distributions shifted, evidencing microstructural changes in sulfur. The most pronounced changes were observed at the highest tested power (Fig. 4). The visual appearance of the solid S_8 also changed coincident with the pulses (Fig. 5). Over the duration of the experiment, the initially precompressed translucent flake became a scattered powder at the irradiation point (Fig. 5b).

Immediately upon exposure of the sample to the XFEL pulse, we observed a small vesicle 66 forming at the point of irradiation, which moved steadily downward through the fluid hydrogen 67 medium, before disappearing upon reaching the edge of the sample chamber (see Supplementary 68 Video). This was interpreted as a rapidly-produced bubble of the reaction product H_2S , the stable 69 hydride of S at these pressures. The motion is consistent with a dense H₂S vesicle $(1.03 \text{ g cm}^{-3})^{29}$ 70 sinking through the lower density molecular hydrogen medium (0.10 g cm^{-3}) ,³⁰ where isothermal 71 equations of state are used as heat associated with the reaction should dissipate within millisec-72 onds,³¹ before the bubble moves. Following Stokes law for a spherical particle falling through a 73 fluid,³² the velocity v is related to the medium viscosity μ , particle radius R, and densities of the 74

⁷⁵ particle (ρ_p) and medium (ρ_m), as

$$v = \frac{2gR^2(\rho_p - \rho_m)}{9\mu}\gamma\tag{1}$$

⁷⁶ where γ is a correction factor, of order 0.1-1.0, that depends on the nature of the motion and ⁷⁷ proximity to a wall (diamond culet). With a measured bubble velocity of ~160 μ m/s and diameter ⁷⁸ of ~8 μ m, and considering the viscosity of hydrogen³⁰ ($\mu = 2.8 \times 10^{-5}$ Pa s), the bubble's motion ⁷⁹ is consistent with a fall of a pure H₂S vesicle immediately adjacent to the culet ($\gamma \sim 0.1$),³² as ⁸⁰ would be expected for a reaction occurring in the vicinity of the S₈ flake attached to the culet.

Raman analysis before and after the XFEL exposures confirmed that a chemical reaction be-81 tween S₈ and H₂ had been induced, from the appearance of the characteristic fluid H₂S stretching 82 mode at 2610 cm⁻¹ (Fig. 5a).^{33,34} The H₂S signal was ubiquitous throughout the fluid H₂ region 83 of the sample, indicating that the concentration was low enough for the phases to be mixed (ra-84 tio of $H_2S:H_2 \ll 1$), as at higher concentrations (~ 1:1) H_2S and H_2 separate into distinct fluid 85 regions.^{12,18} In contrast, the appearance of a small bubble of the product immediately after the 86 reaction is attributed to temporary phase separation made possible by an intense localized reaction, 87 that quickly enhanced the local concentration of product before dissolution could proceed. Treating 88 the bubble as pure H₂S, and assuming it comprises all or most of the product from a single x-ray 89 pulse, the product mass is estimated as 2.8×10^{-13} kg per pulse. Given the x-ray beam diameter 90 of 14 μ m, this implies that the volume of solid S₈ which reacts as a result of each x-ray pulse is 91 equivalent to a $\sim 0.8 \ \mu m$ thickness of the solid S₈ layer in the area of the beam, accounting for its 92 elevated density under pressure (Fig. 3). The initial S₈ layer thickness is $\sim 15 \ \mu$ m, therefore about 93 5% of the exposed S_8 is reacted with each exposure. This explains the disintegration of the S_8 flake 94 (Fig. 5b) in the region of irradiation after ~ 17 pulses above the detectable reaction threshold. 95

These observations collectively indicate that an intense chemical reaction between S_8 and H_2 is induced by x-ray illumination on the femtosecond timescale. In contrast, at room temperature without x-ray excitation there is no evidence of H_2S formation from samples of S_8 in fluid H_2 at ⁹⁹ similar pressures on a timescale of two weeks (the time between loading and final observation in ¹⁰⁰ this study). Synthesis of H_2S from S_8 and H_2 was therefore attributed exclusively to exposure to ¹⁰¹ x-ray photons.



Figure 4: The microstructure changes induced by XFEL exposure. a) Detector images showing the textured Debye-Scherrer rings arising from S₈ upon four successive exposures to XFEL pulses at the highest pulse energy used in this experiment (0.120 mJ/pulse). All four images were obtained from the region of the sample indicated in Fig. 5b, and were obtained at ~ 1 s intervals. Each diffraction image shows the state of the sample produced by the previous XFEL pulse. The azimuthal angle ϕ is shown in yellow in (i), and the black rectangular areas in the images cover an area of damage to the detector. Plots b) and c) show the azimuthal variations in the intensity of the (222) (red) and (026) (blue) Debye-Scherrer rings (identified in (a) by the red and blue arrows), respectively, in each of the four images. The changes in intensity distribution arise from the texture changes induced by each XFEL pulse. d) Azimuthally integrated S diffraction patterns as a function of increasing XFEL energy. Each pattern is labelled by shot number (#) followed by fraction of maximum pulse energy in brackets. The intensities of the three patterns obtained at lower pulse energies are scaled up by a factor of four; the patterns obtained at maximum pulse energy are integrated profiles of 4a) i - iv.

The absorption of an x-ray pulse by the sample initially disrupts the electrons, while the atoms remain fixed in place. The energy absorbed in the sample per volume (H) is estimated as

$$H = \Lambda \alpha \tag{2}$$

where Λ is the x-ray beam energy density within the sample, and α is the local absorption coeffi-104 cient, i.e. of either the sulfur or hydrogen component. The number of x-ray absorption events in the 105 sample, per volume, is H/E_{ph} , where E_{ph} is the photon energy (12 keV). Considering the number 106 of atoms in the same volume, the number of absorption events in hydrogen is very low, or about 107 one event per 10^9 atoms at peak power, whereas for S₈ there is one absorption event for every $\sim 10^5$ 108 atoms; we thus concentrate on the S₈ component. If we also account for electron-impact ionization 109 cascades (affecting a further $10^2 - 10^3$ atoms²⁸), and the disruption of one S₈ molecule of α -S per 110 ionization (Fig. 1), we can expect a direct electronic disruption of bonding for roughly 1-10% of S 111 atoms following each x-ray pulse. The dominant process will thus involve ionization of atoms by 112 electron collisions, with core electron expulsion and subsequent Auger decay rapidly ionising and 113 fragmenting the S₈ rings (e.g. $S_8 \rightarrow S + S_7^{2+} + 2e^-$). The disrupted S could then rapidly react with 114 adjacent, or plausibly interstitial, molecular hydrogen in an attempt to fill valence p-orbitals; unsta-115 ble ring-fragments could break down further in similar interactions. We do not see any evidence, 116 from Raman and XRD measurements, of other allotropes of sulfur forming; therefore the S₈ rings 117 must break down completely or re-form. The fraction of irradiated S atoms directly affected by 118 x-ray energy deposition and ionization is similar to that consumed by the reaction in a single pulse, 119 estimated on the basis of optical observations to be of order 5%. Therefore, it is plausible that the 120 intense reaction observed is driven primarily by the chemical disruption of the sample during and 121 immediately after the femtosecond irradiation, which then undergoes a rapid reaction to form H₂S. 122 Previous evidence of photochemistry in both S-H and Se-H systems has been identified from 123 the partial formation of H₂S (H₂Se) upon prolonged exposure to laser light of sufficient energy (2.3 124 eV) and substantial flux.^{11,13} Our study shows a considerably more intense and complete photore-125 action between S and H₂ resulting from the delivery of a large dose of keV energy photons. In such 126 a photochemical process, the total dose of radiation can be critical in driving a reaction. A simi-127 lar process may thus occur under lower power (i.e. synchrotron) hard x-ray irradiation on longer 128 but routinely accessed timescales. Indeed, photochemically driven molecular reorganization was 129 reported in the H-O system using synchrotron x-rays of comparable photon energy over a much 130

longer period. ²⁰ While these qualitative comparisons suggest that the total dose and the photon energy are critical parameters, nonlinear effects of high beam intensity including multiple ionization,
bulk chemical disruption, and sample heating may also be enhancing the reaction intensity.



Figure 5: a) Raman spectra obtained within H_2 media before (black) and after (blue) exposure to high intensity x-rays. The black spectrum shows pure H_2 whilst the blue spectrum shows a mixture of H_2S dissolved in fluid H_2 . b) photomicrographs of the DAC sample chamber before and after x-ray exposure. The position of the Raman laser-spot is indicated on both images; the right image shows the area exposed to XFEL pulses.

¹³⁴ Concomitant with electronic chemical disruption is the rapid thermalization of the hot electrons ¹³⁵ on a timescale of $\sim 10^{-12}$ seconds,^{28,31} leading to an equilibrium heated state in S₈ (direct heat-¹³⁶ ing of H₂ is negligible). This hot state in the S₈ persists on timescales of $\sim 10^{-6}$ seconds before ¹³⁷ quenching and will additionally heat the surrounding hydrogen to high temperatures.³¹ Temper-¹³⁸ ature induced reaction and subsequent formation of hydrides under pressure is well documented, ¹³⁹ for example as a result of laser-heating S in H₂,^{9,10,19} and gentle heating of Se in H₂ at 473 K.¹² ¹⁴⁰ The equilibrium peak temperature (before heat dissipation) in these experiments can be estimated 141 as³¹

$$T - T_0 = \frac{H}{\rho C_P} \tag{3}$$

where ρ is the density and C_P the heat capacity at constant pressure of the sample. The temperature achieved in S₈ during the pulse is thus estimated to peak at ~1000 K. As similarly elevated temperatures can themselves enhance a chemical reaction between S and H₂, it is possible that this contributes partially to the total reaction progression, but as the heating timescale is extremely short compared to existing studies, ^{9,10,19} this may not have as strong an effect as the direct electronic disruption.

It is worth observing that such a peak temperature is well in excess of the melting point of solid 148 S under pressure (477 K at 0.3 GPa³⁵). Therefore, the successive changes to the microstructure of 149 the crystalline S_8 (Fig. 4) must be related both to the sample breakup associated with the reaction 150 (Fig. 5b), and also to the local melt and recrystallization of bulk S, as these processes occur well 151 before the arrival of the subsequent x-ray pulse and diffraction measurement. In particular, the 152 change in the texture of the Debye-Scherrer rings from smooth arcs (Fig 4a.i), to spotty points 153 of intensity (Fig 4a.ii-iii), then back to smooth arcs (Fig 4a.iv) is suggestive of both crystallite 154 displacement and melt recrystallization. In summary, we find that the reaction mechanism is one 155 where the XFEL pulse induces flash localized heating and massive chemical disruption in the S, 156 some of which reacts with adjacent or interstitial H₂ thus forming H₂S, while unreacted molten S 157 recrystallizes back into orthorhombic α -S₈. 158

In conclusion, we find that an intense high-pressure chemical reaction between sulfur and hy-159 drogen can be initiated by hard x-ray pulses from a free electron laser. Delivery of a large x-ray 160 photon dose (10¹¹ photons) in a short time (fs) produces rapid chemical damage as well as sam-161 ple heating, primarily localized to the higher-Z reactant (crystalline S). This reactivity in the S-H 162 system, induced by the x-rays, shows interesting differences from those in the high-pressure O-H 163 system,²⁰ with formation rather than dissociation of hydride phases. This suggests a potentially 164 rich x-ray induced reaction behavior in chalcogen hydride systems at high pressures. This im-165 portant observation is relevant to high pressure experiments concerning the superconducting S-H 166

¹⁶⁷ system, where the reaction behavior and results of x-ray experiments have been in dispute.^{9,10,16}

Our study shows that x-ray irradiation at high photon energies may not act simply as a noninvasive probe of the structural and molecular state of this system, as is commonly assumed, but may in fact cause reactivity that would otherwise not be observed. These effects must be further investigated as functions of pressure, in both S-H and Se-H systems, and may also provide a novel synthesis route in superconductor forming hydrides.

Our study finally benchmarks a new application for XFEL radiation in high-pressure chem-173 istry studies under static compression, demonstrating a strategy for non-destructive XFEL probing 174 and excitation using robust sample confinement in a diamond anvil cell, and the ability to obtain 175 high-quality diffraction data from such samples in femtoseconds. With new experimental regimes 176 expected to be accessible with advances in capabilities at 4th generation x-ray sources, such as 177 improved brilliance and faster pulse-repetition rates, it is important to gain a comprehensive un-178 derstanding of their effect on high pressure materials systems in order to fully take advantage of 179 these modern light sources. 180

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¹⁹⁴ Supporting Information Available: Supporting methods and supplementary video recording.

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