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<b>Author(s)</b>	<b>Tang, WH; Gao, J</b>
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# Compatibility of Nd and Ba in $\text{YBa}_2\text{Cu}_3\text{O}_y$ Superconductor

W.H. Tang and J. Gao

Department of Physics, The University of Hong Kong, Pokfulam Road, Hong Kong

**Abstract** — A series of samples with nominal compositions of  $\text{YBa}_{2-x}\text{Nd}_x\text{Cu}_3\text{O}_y$  were prepared by solid state reactions and investigated by X-ray diffraction, magnetic and electrical measurements. The solubility,  $x$ , was determined to be  $x < 0.25$ . It is found that the solubility of rare earth Re ions at the Ba-sites in  $\text{Re}_{1+x}\text{Ba}_{2-x}\text{Cu}_3\text{O}_y$  is strongly related to a geometric parameter, the tolerance factor  $t$ . For  $x < 0.25$ ,  $T_{c0}$  are all above 92K. For  $x > 0.3$   $T_{c0}$  drops sharply to about 84K, and finally for  $x=0.5$   $T_{c0}$  falls to 30K and  $T_{c0}$  is below 10K. Our results suggest that the small compatibility of Nd and Ba in  $\text{YBa}_2\text{Cu}_3\text{O}_y$  system does not seriously suppress the superconductivity.

## I. INTRODUCTION

It is well known that  $\text{ReBa}_2\text{Cu}_3\text{O}_y$  (ReBCO) superconductors show almost the same superconducting transition temperature ( $T_c \sim 90\text{K}$ ) [1].  $\text{YBa}_2\text{Cu}_3\text{O}_y$  (YBCO) is a well-studied electrode material for high- $T_c$  superconductor-normal metal-superconductor (SNS) junctions. Barrier materials with suitable physical properties and good interface character are key factors for the fabrication of reliable and high quality SNS junctions. Various barriers such as  $\text{PrBa}_2\text{Cu}_3\text{O}_y$  (PBCO) [2],  $\text{Pr}_{1-x}\text{Y}_x\text{Ba}_2\text{Cu}_3\text{O}_y$  [3] and doped YBCO [4] have been used for high- $T_c$  SNS junctions.  $\text{NdBa}_2\text{Cu}_3\text{O}_y$  (NBCO) possesses the highest  $T_c$  in the Re123 family [5], however, its  $T_c$  depends strongly on the fabricating conditions [6]. Since the lattice constants  $a$  and  $b$  of NBCO match very well with those of YBCO, it could be another good barrier material for SNS junctions. Oh et al [7] have reported on a SNS ramp junction with a NBCO barrier. The interface between the YBCO and NBCO is fairly clean, and the interface resistance is very small. To understand the interface between the YBCO and NBCO layers, it is important to investigate the compatibility of YBCO and NBCO. However, it has been suggested recently that the magnetic rare earth ions at the Ba-sites are fatal to the superconductivity [8]. It is even predicted that PBCO without Pr-at-Ba-site defects should be a superconductor with a  $T_c \sim 90\text{K}$ .  $\text{Y}^{3+}$  ion is nonmagnetic, its solubility at the Ba-sites is very limited.  $\text{Nd}^{3+}$  ion, however, is magnetic ( $3.54\mu_B$ ) and its solubility at the Ba-sites is very wide. Thus the  $\text{YBa}_{2-x}\text{Nd}_x\text{Cu}_3\text{O}_y$  system is suitable for as investigation into the effects of magnetic Re ions at the Ba-sites on superconductivity. In this work, we aim to check the compatibility of Nd and Ba and its influence on crystal structure and superconductivity in  $\text{YBa}_2\text{Cu}_3\text{O}_y$ .

## II. EXPERIMENTAL

The samples of  $\text{YBa}_{2-x}\text{Nd}_x\text{Cu}_3\text{O}_y$  with  $x=0.0, 0.05, 0.10, 0.15, 0.20, 0.25, 0.30, 0.35, 0.40, 0.50$  were prepared by solid state reactions. The mixture of the powders of  $\text{Y}_2\text{O}_3$  (99.9%),  $\text{Nd}_2\text{O}_3$  (99.9%),  $\text{BaCO}_3$  (99.5%) and  $\text{CuO}$  (99.9%) was ground and calcined in air at  $970^\circ\text{C}$  for 12h. The resultant powder was reground and sintered in air at  $960^\circ\text{C}$  for 12h. After this sintering process the powder samples were examined by x-ray diffraction with  $\text{CuK}_\alpha$  radiation. After being reground and pelletized, the pellet was fired in air at  $960^\circ\text{C}$  for 12h and at  $500^\circ\text{C}$  in flowing oxygen (1l/min) for 10h. Finally, the samples were cooled down to room temperature at a slow cooling rate ( $1^\circ\text{C}/\text{min}$ ). The final samples were examined by the x-ray diffraction again.

The resistance vs temperature curves were measured by a standard DC four-probe method. The samples were cooled by liquid nitrogen or a closed-cycled cryogenic system. The cooling rate was controlled to be  $2\text{K}/\text{min}$  by a computer program. The temperature was measured by a Pt resistance thermometer. The temperature error was estimated to be less than  $0.5\text{K}$ . The susceptibility was measured by an ac mutual inductance method, with an operating frequency of  $120\text{Hz}$ .

## III. RESULTS AND DISCUSSION

All samples were examined by x-ray powder diffraction. The results indicate that samples with  $x < 0.3$  are nearly single Re123 phase. The impurity phase content was estimated to be less than 1wt.%. Fig. 1 shows the x-ray powder diffraction patterns for the samples of  $x=0.05, 0.15$  and  $0.25$ . All the intensities of the (001) diffraction peaks were enhanced strongly, indicating an obvious (001) preferred direction for the bulk samples. The x-ray powder diffraction patterns can be indexed by a tetragonal Re123 lattice. The orthorhombic splitting of (200), (020) and (006) diffraction peaks at  $2\theta \sim 46.5^\circ$  is very small (see Fig.2). The solubility of Nd at the Ba-sites in  $\text{YBa}_{2-x}\text{Nd}_x\text{Cu}_3\text{O}_y$  is less than 0.25, which is smaller than that in  $\text{Nd}_{1+x}\text{Ba}_{2-x}\text{Cu}_3\text{O}_y$  ( $x \leq 0.6$ ) [9]. The size effect of the rare earth ions is responsible for the reduction in the solubility of Nd.

It is well known that the stability of the perovskite structure can be scaled by a geometrical parameter, i.e. tolerance factor,  $t$ . For a  $\text{ABO}_3$  perovskite structure the tolerance factor,  $t$ , is defined as[10]:

$$t = (r_A + r_O) / \sqrt{2} (r_B + r_O) \quad (1)$$

where  $r_A$ ,  $r_B$  and  $r_O$  are the radii of the A, B and O ions, respectively. For  $0.8 < t < 1.0$ , the perovskite structure can

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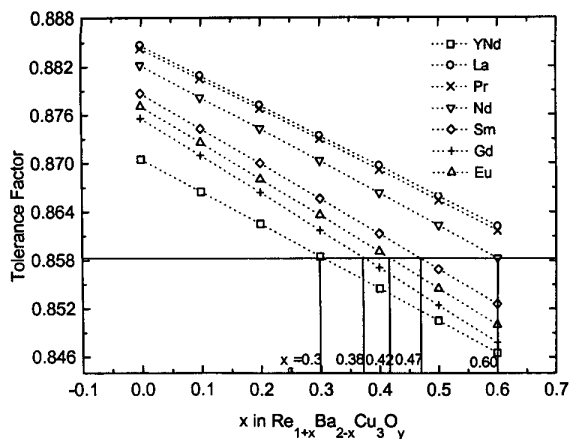


Fig. 3. Calculated tolerance factor of  $\text{Re}_{1+x}\text{Ba}_{2-x}\text{Cu}_3\text{O}_y$

TABLE II  
SUPERCONDUCTING TEMPERATURE OF  $\text{YBa}_{2-x}\text{Nd}_x\text{Cu}_3\text{O}_y$

x	$T_c^{\text{on}}$ (K)	$T_c^{\text{zero}}$ (K)	$T_c^*$ (K)
0.00	95.5	92.0	93.0
0.05	95.0	92.0	92.5
0.10	95.5	92.5	93.0
0.15	95.0	91.5	92.0
0.20	95.0	92.5	92.0
0.25	96.0	94.0	93.0
0.30	95.0	83.0	85.0
0.35	95.0	86.0	88.0
0.40	95.0	85.0	86.0
0.50	30.0	<10.0	

$T_c^{\text{on}}$  is the onset transition temperature of the resistance curve;  
 $T_c^{\text{zero}}$  the zero resistance temperature;  
 $T_c^*$  the transition temperature of the ac susceptibility.

The substitution of  $\text{Nd}^{3+}$  for  $\text{Ba}^{2+}$  probably gives rise to three kinds of possible effects. The first one is the size effect from the difference of the radii of  $\text{Nd}^{3+}$  (0.995Å) and  $\text{Ba}^{2+}$  (1.34Å). It has little effect on the superconductivity, but has an obvious effect on the solubility of Nd at the Ba-sites and the structural transition from orthorhombic to tetragonal as discussed above. The second one is the magnetic pair breaking from the magnetic  $\text{Nd}^{3+}$  ions ( $3.54\mu_B$ ) due to the spin-dependent exchange scattering of mobile holes in the  $\text{CuO}_2$  planes by the magnetic  $\text{Nd}^{3+}$  ions at the Ba-sites. According to the theory of Abrikosov and Gor'kov, a linear suppression of  $T_c$  with the magnetic impurity concentration  $x$  would be expected. That does not agree with our experimental results. From the view point of the crystal structure, the rare earth and Ba sites originate from the same crystallographic sites of the disordered  $(\text{Re},\text{Ba})\text{CuO}_3$  perovskite. In the ordered distorted perovskite  $\text{Re}123$  structure, the distances from the Re- and Ba-sites to the  $\text{CuO}_2$  plane are slightly different. It is well known that the magnetic interaction is strongly dependent on the interaction distance. In fact, the distance from Ba-sites to the  $\text{CuO}_2$  plane is longer than that from Re-sites to  $\text{CuO}_2$  plane. Thus the magnetic depairing effect from the magnetic Re ions at Ba-sites should be smaller than that at Re-sites. It is believed that magnetism of Re ions at Re-sites does not effect the superconductivity in

$\text{ReBCO}$  superconductors. The third one is due to the difference in the valences of  $\text{Nd}^{3+}$  and  $\text{Ba}^{2+}$ . The superconductivity of p-type doped superconductors is dependent on the holes predominantly on the oxygen sites in the  $\text{CuO}_2$  planes, and substitution of  $\text{Re}^{3+}$  for  $\text{Ba}^{2+}$  would require donor electrons, whence canceling partial holes on the  $\text{CuO}_2$  planes. The superconducting transition temperature,  $T_c$ , should then decrease with increasing the content of  $\text{Re}^{3+}$ . Our results show that  $T_c$  is almost unchanged for  $x \leq 0.25$ . That might suggest that the charge balance is maintained by the extra oxygen and the total hole concentration should not change. Neutron diffraction indicates that the increase of Nd at the Ba-sites results in an increase of the occupation of the  $\text{O}_5$  sites [9]. The  $\text{O}_5$  does not affect the  $\text{CuO}_2$  plane directly, but affects on the  $\text{CuO}_2$  plane indirectly through  $\text{O}_4$ . Due to the screening of  $\text{O}_4$  it is possible for  $\text{O}_5$  to not suppress the superconductivity in a certain occupation of  $\text{O}_5$ . The superconducting transition temperatures  $T_c$  of  $\text{Nd}_{1+x}\text{Ba}_{2-x}\text{Cu}_3\text{O}_y$  solid solution, reported by different groups, seems very scattered [9][15-19]. Yossefov et al. reported that  $T_c$  was insensitive to Nd content,  $x$ , a ranged between 98.7 and 94K for  $x < 0.25$  [19]. That is very similar to our result (See Fig.4). The scattered results for the  $T_c$  of  $\text{Nd}_{1+x}\text{Ba}_{2-x}\text{Cu}_3\text{O}_y$  solid solutions indicate that the oxygen content and the oxygen occupation are key factors for the superconductivity. We believe that the oxygen content and charge transfer due to the difference in the valences of  $\text{Nd}^{3+}$  and  $\text{Ba}^{2+}$  may be the key factor in the superconductivity of  $\text{YBa}_{2-x}\text{Nd}_x\text{Cu}_3\text{O}_y$ . The magnetic pair breaking of magnetic rare earth ions at the Ba-sites seems not to be a major factor in the  $T_c$ -depression of  $\text{ReBa}_2\text{Cu}_3\text{O}_y$  superconductors. This was supported by our recent results on  $\text{YBa}_{2-x}\text{Pr}_x\text{Cu}_3\text{O}_y$  [20].

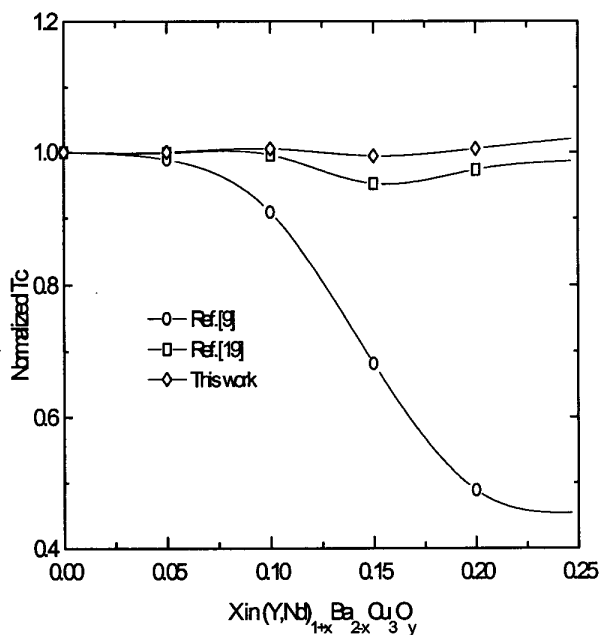


Fig. 4 Normalized  $T_c$  for  $\text{Nd}_{1+x}\text{Ba}_{2-x}\text{Cu}_3\text{O}_y$  and  $\text{YBa}_{2-x}\text{Nd}_x\text{Cu}_3\text{O}_y$ .

## IV. CONCLUSIONS

The solubility of Nd on the Ba-sites in  $\text{YBa}_{2-x}\text{Nd}_x\text{Cu}_3\text{O}_y$  was determined to be  $x=0.25$  by x-ray powder diffraction. The solubility of Nd and the structural transition from orthorhombic to tetragonal are related to a geometric parameter, the tolerance factor,  $t$ . The onset superconducting transition temperature,  $T_c^{\text{on}}$ , determined from the curves of resistance vs temperature, exhibits a weak dependence on the Nd content,  $x$ , for  $x < 0.4$ . The variation in the zero resistance transition temperatures,  $T_c^{\text{zero}}$ , vs the Nd content,  $x$ , exhibits two steps. For  $x \leq 0.25$   $T_c^{\text{zero}}$  values are all above 92K, and  $T_c^{\text{zero}}=94\text{K}$  for  $x=0.25$ . For  $x \geq 0.3$   $T_c^{\text{zero}}$  drops sharply to about 84K, and finally for  $x=0.5$   $T_c^{\text{on}}$  falls to 30K and  $T_c^{\text{zero}}$  is below 10K. The small compatibility of Nd and Ba in  $\text{YBa}_2\text{Cu}_3\text{O}_y$  system does not seriously suppress the superconductivity. The oxygen content and charge transfer due to the unbalance between  $\text{Nd}^{3+}$  and  $\text{Ba}^{2+}$  might be key factors in the superconductivity of  $\text{YBa}_{2-x}\text{Nd}_x\text{Cu}_3\text{O}_y$ . The magnetic pair breaking of magnetic rare earth ions at the Ba-sites seems not to be a major factor in the  $T_c$ -depression of  $\text{REBa}_2\text{Cu}_3\text{O}_y$  superconductors.

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