

# 1 Submesoscale eddies in the South China Sea

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12

## 13 **Key Points**

- 14 • Submesoscale eddies are detected automatically from ocean colour data and are  
15 analyzed statistically in the SCS
- 16 • The surface structure of submesoscale eddies shows the classical ‘cat’s-eye’  
17 pattern
- 18 • Submesoscale eddies can significantly modulate surface tracer distribution

19

## 20 **Abstract**

21 Submesoscale eddies are often seen in high-resolution satellite-derived ocean  
22 colour images. To efficiently identify these eddies from surface chlorophyll data, here  
23 we develop an automatic submesoscale eddy detection method and apply it to the  
24 South China Sea (SCS). The detected submesoscale eddies are found to have a radius  
25 of  $13\pm 5$  km and an aspect ratio of  $0.5\pm 0.2$ , with a notable predominance of cyclones.  
26 Further investigation reveals that the surface structure of these eddies displays a  
27 unique ‘cat’s-eye’ pattern and the eddies become more circular with increasing eddy  
28 radius. Submesoscale eddies can strongly regulate surface chlorophyll via horizontal  
29 advection while they have less coherent signatures in sea surface temperature. These  
30 findings may help to improve submesoscale parameterizations in Earth system  
31 models.

## 32 **Plain Language Summary**

33 Ubiquitous ocean eddies play a crucial role in the upper ocean dynamics. Using  
34 high-resolution satellite remote sensing data, we have developed an automatic method  
35 to detect small elliptical eddies in the SCS over a 10-year period. The results show  
36 that these ‘submesoscale’ eddies of the order of 10 km appear to have a unique  
37 ‘cat’s-eye’ structure with significant effect on the surface tracer distribution. This  
38 study therefore improves our understanding of oceanic submesoscale dynamics and  
39 contributes to parameterizing the impact of submesoscale eddies in climate and ocean  
40 models.

## 41 **1. Introduction**

42 Submesoscale spiral eddies of the order of 10 km have been frequently observed  
43 in different regions over the world ocean since they were first seen in the sun-glitter  
44 from the Apollo Mission in 1968 (e.g., Munk et al., 2000; Shen and Evans, 2002;  
45 Buckingham et al., 2017). Although submesoscale eddies are believed to be important

46 for upper ocean dynamics and biogeochemical processes (Haine and Marshall, 1998;  
47 Munk et al., 2000; McWilliams, 2010; Mahadevan, 2016), progress in characterizing  
48 and understanding them has been slow, because the resolutions of in-situ ocean  
49 measurements and satellite altimetry observations are typically too coarse to resolve  
50 these small-scale and short-lifetime eddies. One way to overcome this obstacle is to  
51 utilize other satellite remote sensing data, such as sea surface temperature (SST) and  
52 near-surface chlorophyll, which is available at high resolution and wide coverage  
53 (Munk et al., 2000; Liu et al., 2014; Buckingham et al., 2017). However, to our  
54 knowledge, no methods exist yet that are able to extract submesoscale spiral eddies  
55 from the remote sensing images in an automatic and systematic way. In this study, we  
56 first develop an automatic submesoscale eddy detection method and then apply it to  
57 the South China Sea (SCS), the largest marginal sea in the western Pacific that is rich  
58 in submesoscale eddies.

59 The SCS is characterized by varying seafloor topography, a seasonal upper ocean  
60 circulation, a complex upwelling-front system and active mesoscale eddies, which  
61 facilitate the generation of submesoscale phenomena (Wang et al., 2003; Hu and  
62 Wang, 2016; Lin et al., 2020). Although submesoscale eddies have been seen a few  
63 times in remote sensing data in the northern and western SCS (e.g., Su, 2004; Liu et  
64 al., 2014; Yu et al., 2018), the statistical properties of these eddies in the SCS (e.g.,  
65 size, polarity and shape) have not been determined. In a seminar paper on spiral  
66 eddies, Munk et al. (2000) proposed that the surface structure of submesoscale spiral  
67 eddies can be described by an extension of the classical Stuart (1967) solution, which  
68 yields the well-known ‘cat’s eye’ configuration (Thomson, 1880; Fig. 1a). However,  
69 this cat’s-eye surface structure proposed for submesoscale eddies is yet to be  
70 observationally confirmed and the key parameter in the Stuart solution to be  
71 determined. Automatic submesoscale eddy detection enables composite analyses of  
72 chlorophyll and SST anomalies associated with these eddies and as such is a useful  
73 tool for analyzing the surface structure of submesoscale eddies as well as their impact  
74 on surface tracer distributions.

## 75 **2. Data**

76 The daily Moderate Resolution Imaging Spectroradiometer (MODIS)  
77 chlorophyll and SST data from the National Aeronautics and Space Administration  
78 (NASA) Ocean Colour project are analyzed in this study for a 10-year period from  
79 January 2006 to December 2015. Both the chlorophyll and SST data are level-2  
80 products provided with a spatial resolution of  $\sim 1$  km. Because of the log-normal  
81 distribution of chlorophyll concentration, we follow Chelton et al. (2011) and  $\log_{10}$   
82 transform the chlorophyll field before compositing chlorophyll anomalies associated  
83 with submesoscale eddies.

## 84 **3. Results**

### 85 **3.1. Statistical Features**

86 We first develop an automatic submesoscale eddy detection method based on the  
87 curvature of contours extracted from high-resolution chlorophyll data. The  
88 chlorophyll images are first processed to fill small blank patches due to clouds (Oram  
89 et al., 2008). The extracted chlorophyll contours are then broken into segments  
90 according to the contour curvature direction. The clustering segments that curl in the  
91 same direction are regarded as different parts of the same submesoscale eddy if they  
92 further satisfy a number of criteria. The type, edge and center of a submesoscale eddy  
93 are defined as the type, convex hull and geometric center of the segments of the eddy,  
94 respectively. A detailed description of the automatic submesoscale eddy detection  
95 method is provided in the Supporting Information (Fig. S1). For example, based on  
96 this method, two cyclonic submesoscale eddies are identified in the western SCS  
97 during the summer of 2012 (Fig. 1b) and an anticyclonic submesoscale eddy is  
98 detected in the eastern SCS during the winter of 2012 (Fig. 1c). Overall, about 5983  
99 (4372) snapshots of cyclonic (anticyclonic) submesoscale eddies are identified in the  
100 entire SCS over the 10-year study period. The elevated number of cyclonic  
101 submesoscale eddies over their anticyclonic counterparts is consistent with the

102 findings of previous theoretical and numerical studies that anticyclonic submesoscale  
103 eddies are subject to inertial instability while cyclonic submesoscale eddies are not  
104 (Munk et al., 2000; Shen and Evans, 2002; Dong et al., 2007; Hasegawa et al., 2009).  
105 Note that in weakly-stratified waters anticyclonic eddies are found to be more stable  
106 than cyclonic eddies (Buckingham et al. 2020). Submesoscale eddies in the SCS are  
107 frequently detected in the coastal regions (Fig. 1d), including the northern SCS  
108 shelf-slope region, both sides of the Luzon strait and the coastal waters off Vietnam,  
109 where submesoscale eddies have been reported before (e.g., Su, 2004; Zheng et al.,  
110 2008; Liu et al., 2014). In these boundary regions, enhanced along-slope velocity  
111 shear, strong coastal front instability and vortex stretching due to tidal flow over  
112 shallow waters are known to be able to generate submesoscale eddy activity (Munk et  
113 al., 2000; Gula et al., 2015; Li et al., 2020). A recent high-resolution modelling study  
114 by Lin et al. (2020) confirms that submesoscale processes are particularly active in  
115 these coastal regions of the SCS. Furthermore, the large chlorophyll gradients near the  
116 coast (Fig. S2a) facilitate identification of submesoscale eddies via our detection  
117 method which is based on chlorophyll contours. For both types of submesoscale  
118 eddies, they are more frequently detected in winter and summer while less in spring  
119 and autumn (Fig. S3), which is probably related to the strongly seasonally-varying  
120 upper ocean circulation in the SCS driven by the monsoon (Wang et al., 2003; Su,  
121 2004; Liu et al., 2014).

122 Here we define the radius of a submesoscale eddy as the radius of a circle that  
123 has the same area as the eddy. Statistical analysis shows that the radii of submesoscale  
124 eddies in the SCS range from about 3 km to more than 30 km, with a mean value of  
125 14.2 km (13.4 km) and a standard deviation of 5.2 km (4.5 km) for cyclones  
126 (anticyclones) (Table 1; Fig. 2a). The eddy radii estimated in this study are  
127 comparable in magnitude to those estimated from various data in previous research  
128 (Liu et al., 2014; Xu et al., 2015; Yu et al., 2018). When it comes to characterizing  
129 eddy shape, one useful metric is eddy aspect ratio, which is defined as the ratio  
130 between the minor and major radius of the fitted ellipse. The probability density

131 function of the aspect ratios of submesoscale eddies contains a skewed distribution  
132 (Fig. 2b), with an average of 0.48 (0.49) and a standard deviation of 0.18 (0.18) for  
133 cyclones (anticyclones) (Table 1). Interestingly, the eddy aspect ratio is found to be a  
134 function of the eddy radius, irrespective of the eddy polarity (Fig. 2c); the larger the  
135 submesoscale eddies, the more circular they are.

### 136 **3.2. Horizontal Structure**

137 The identified eddy edges are also used to investigate the horizontal structure of  
138 submesoscale eddies. We first create a rotated coordinate system for the eddies, where  
139 the coordinate center is defined as the center of each eddy, with the major (minor)  
140 axis of the eddy on the x-axis (y-axis) (Supporting Information; Fig. S4). After that,  
141 we project the edges of cyclonic and anticyclonic submesoscale eddies separately  
142 onto the rotated eddy coordinate (Figs. 3a, b and S5). The average edges of cyclonic  
143 and anticyclonic submesoscale eddies are found to be almost identical, revealing a  
144 nearly perfect ‘cat’s-eye’ structure as shown in previous theoretical and numerical  
145 studies (Munk et al., 2000; Shen and Evans, 2002). We then compare the observed  
146 mean edges of submesoscale eddies with the Stuart solution  
147  $\psi = -U/k \cdot \log(\cosh(ky) - \alpha \cdot \cos(kx))$ , where  $U = \pm 0.3 \text{ m s}^{-1}$  is the background  
148 shear flow,  $k \approx 0.0003 \text{ m}^{-1}$  is the ratio between  $2\pi$  and eddy length scale, and  $\alpha$  is  
149 an unknown parameter between 0 and 1 that needs to be determined (following Munk  
150 et al., 2000). The Stuart solution yields parallel shear flows when  $\alpha = 0$  and  
151 concentrated point vortices as  $\alpha$  approaching 1. By adjusting  $\alpha$  to obtain a best fit  
152 of the Stuart solution to the observed eddies, both cyclonic and anticyclonic, we find  
153  $\alpha = 0.6$  gives a good agreement. Our result therefore provides the first statistical  
154 observational evidence in support of the ‘cat’s-eye’ horizontal structure proposed by  
155 Munk et al. (2000) for submesoscale eddies.

156 Given that the submesoscale eddy aspect ratio depends on eddy radius (Fig. 2c),  
157 the value of  $\alpha$  in the Stuart solution may also vary with the radius of submesoscale  
158 eddies. To test this conjecture, we divide the identified eddies into five bins, at an

159 interval of 5 km from 5 km to 30 km, according to the eddy radius. Then, we average  
160 all the fitted ellipse edges of submesoscale eddies in each bin to estimate the  
161 best-fitting  $\alpha$  for each bin. The value of  $\alpha$  is indeed found to vary with the  
162 submesoscale eddy radius, increasing from over 0.4 to around 0.7, with slightly  
163 smaller values for cyclones (Fig. 3c). Moreover, binning of  $\alpha$  as a function of the  
164 radius of cyclonic (anticyclonic) submesoscale eddies displays a nearly linear  
165 relationship, with  $\alpha = 0.015r + 0.322$  ( $\alpha = 0.015r + 0.344$ ) where  $r$  is the  
166 radius of submesoscale eddies. The relationship between the eddy radius and  $\alpha$   
167 found in this study can be used to improve the Stuart solution to better describe the  
168 surface structure of submesoscale eddies which may have implications for  
169 submesoscale eddy parameterizations.

### 170 **3.3. Composite chlorophyll and SST**

171 To examine the impact of submesoscale eddies on surface tracer distributions,  
172 the  $\log_{10}$ -transformed chlorophyll and SST data of the 10-year study period are first  
173 high-pass filtered using a Gaussian filter (Ni et al., 2020) and then are projected and  
174 averaged onto the rotated submesoscale eddy coordinate (Supporting Information; Fig.  
175 S4). Note that the flank of an eddy with positive chlorophyll anomalies is taken as the  
176 positive y-axis. Fig. 4a (b) shows the resulting composite chlorophyll anomalies  
177 inside and around cyclonic (anticyclonic) submesoscale eddies detected in the SCS.  
178 On average, the magnitude of  $\log_{10}$ -transformed chlorophyll anomalies induced by  
179 submesoscale eddies is on the order of  $\pm 0.1 \text{ mg m}^{-3}$ , which is comparable to the  
180 magnitude of seasonal variations of surface chlorophyll anomalies averaged over the  
181 SCS (Fig. S2b) but several times larger than that associated with mesoscale eddies  
182 (Chelton et al., 2011; Gaube et al., 2014; He et al., 2019). We also note that the  
183 composite chlorophyll anomalies indicate a ‘cat’s-eye’ shape and display a distinct  
184 dipole pattern which consists of two rotational anomalies of opposite sign. Similar  
185 dipole structure has been seen in the composite maps of tracer anomalies (i.e.,  
186 chlorophyll and SST) induced by mesoscale eddies, which is known to result from

187 lateral eddy advection of background tracer gradients (Chelton et al., 2011; Hausmann  
188 and Czaja, 2012; Gaube et al., 2015). In regions of significant background chlorophyll  
189 gradient, the effect of horizontal eddy rotation is to advect high (low) chlorophyll  
190 concentration to the side of low (high) chlorophyll concentration and thereby result in  
191 positive (negative) chlorophyll anomalies. Indeed, the composite maps of Figs. 4a and  
192 b indicate the existence of distinct chlorophyll fronts at  $y \approx 0$ .

193 The composite SST anomalies associated with the identified cyclonic and  
194 anticyclonic submesoscale eddies are shown in Figs. 4c and d, respectively. One  
195 outstanding feature is that positive (negative) SST anomalies on the flanks of  
196 submesoscale eddies are collocated with negative (positive) chlorophyll anomalies,  
197 consistent with the fact that near the coast the chlorophyll concentration is higher  
198 while the SST is colder. Furthermore, the signatures of submesoscale eddies in the  
199 composite SST anomaly images tend to be more obscure when compared to  
200 chlorophyll. One possible explanation is that there exist various formation  
201 mechanisms for submesoscale eddies. For the mechanism of frontal instability, the  
202 pattern of chlorophyll anomalies is expected to be similar to that of SST anomalies  
203 (Munk et al., 2000; Klein and Lapeyre, 2009). For the mechanism of shear instability,  
204 however, a different picture occurs. For example, submesoscale eddies caused by  
205 flow-island interaction may occur in a relatively homogeneous temperature field (Fig.  
206 S1f; Yu et al., 2018), and as a result the imprint of submesoscale eddies in the SST  
207 anomalies are less pronounced. Previous research indeed found greater chlorophyll  
208 variance at submesoscales than SST (Mahadevan, 2016). This is why we choose  
209 chlorophyll rather than SST to identify submesoscale eddies in our method. The  
210 difference between submesoscale eddy signatures in chlorophyll and SST maps also  
211 reflects the degree of conservativeness in their behaviour, which may need to be  
212 accounted for when parameterizing the effect of submesoscale eddies in the tracer  
213 equations.



#### 214 **4. Conclusions**

215 In this work we have developed an automatic submesoscale spiral eddy  
216 identification method based on high-resolution chlorophyll data and then applied it to  
217 the SCS which is a marginal sea rich in submesoscale eddies. The detected  
218 submesoscale eddies in the SCS are found to have a radius of  $13\pm 5$  km and an aspect  
219 ratio of  $0.5\pm 0.2$ , with a notable predominance of cyclones. We have shown that the  
220 surface structure of submesoscale eddies displays the classical ‘cat’s-eye’ pattern and  
221 further determined the key unknown parameter in the Stuart solution that describes  
222 the shape of the cat’s-eye pattern. Submesoscale eddies are found to induce dipole  
223 surface chlorophyll and SST anomalies via horizontal advection of background  
224 chlorophyll and SST gradients.

225 The widespread existence of submesoscale eddies is believed to be important in  
226 tracer transport, energy cascade, re-stratification and biological processes in the upper  
227 ocean (Ubelmann and Fu, 2011; McWilliams, 2010; Haine and Marshall, 1998;  
228 Mahadevan, 2016). However, the present global ocean and climate models have too  
229 coarse spatial resolutions to resolve submesoscale processes and as such would rely  
230 on parameterizing the effect of submesoscale eddies for the foreseeable future (e.g.,  
231 Fox-Kemper et al., 2011). The submesoscale eddy structure and statistics found in this  
232 study may provide observation-based guidance for future development of  
233 submesoscale eddy parameterizations. For example, anisotropy in submesoscale eddy  
234 length scales, i.e., shorter length scale in the cross-front direction than along-front  
235 direction, implies anisotropic submesoscale eddy diffusivity if the parameterization  
236 scheme employs a mixing length approach.

237 The high-resolution Surface Water and Ocean Topography (SWOT) satellite  
238 altimeter is scheduled to launch in 2021 (Qiu et al., 2017), which aims at resolving sea  
239 level variability at submesoscales. Combining the chlorophyll-based submesoscale  
240 eddy detection method developed in this study with SWOT-derived submesoscale sea  
241 level anomalies should have potential to further improve our understanding of the

242 surface pattern, dynamics and impact of submesoscale eddies. Nevertheless, in  
243 addition to satellite remote sensing, we still need in-situ observing technologies with  
244 high-enough spatiotemporal resolution to reveal the three-dimensional structure of  
245 these eddies.

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254 SST data are available at <https://oceandata.sci.gsfc.nasa.gov/>.

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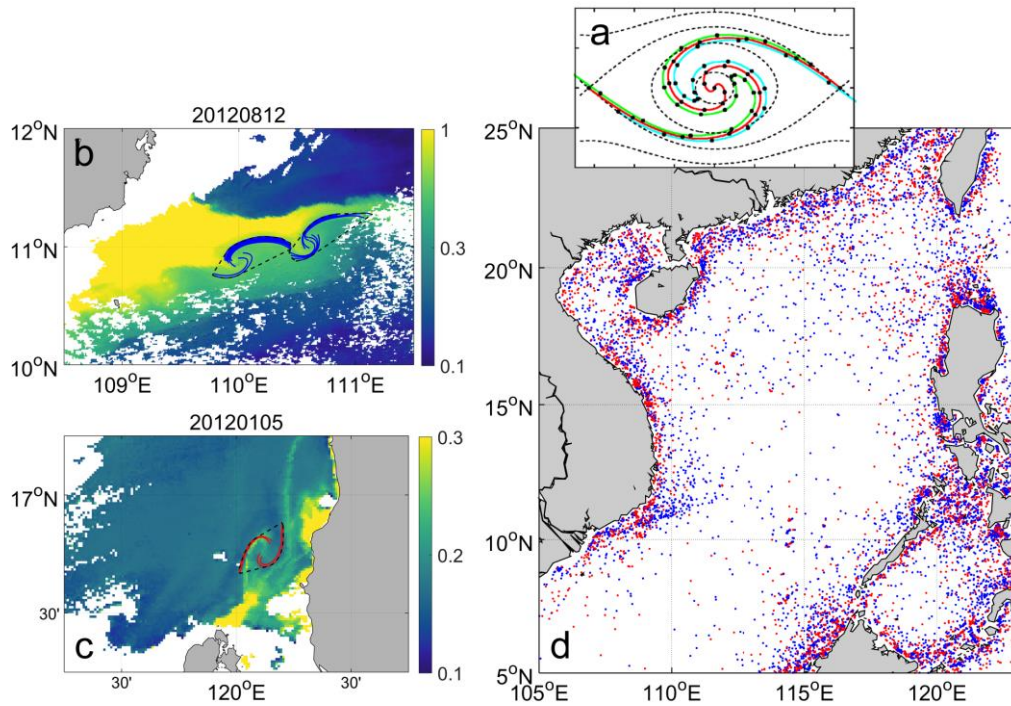
341

342 **Table**

343 Table 1. Statistical features of submesoscale eddies detected in the South China Sea  
 344 from 2006 to 2015

Polarity	r (km)		$r_{min}/r_{maj}$	
	Mean	STD	Mean	STD
Cyclonic	14.2	5.2	0.48	0.18
Anticyclonic	13.4	4.5	0.49	0.18

345

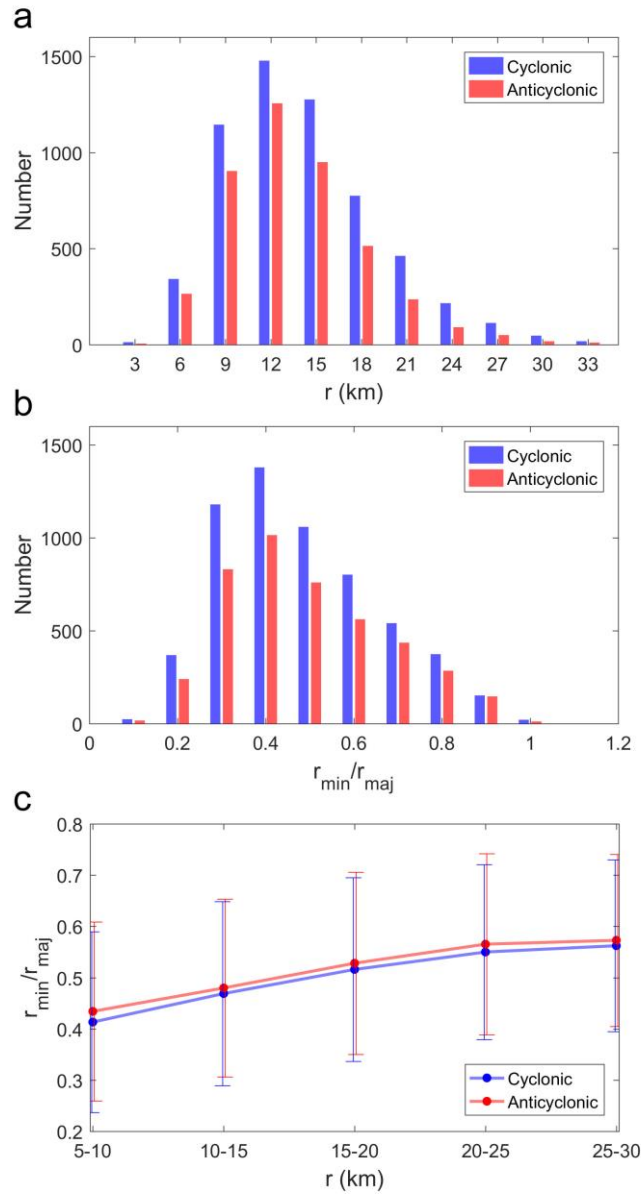


347

348 Figure 1. (a) Particle distribution (black dots and colour curves) in a Stuart spiral eddy  
 349 (black dashed contour) that shows a ‘cat’s-eye’ pattern. Adapted from Munk et al.  
 350 (2000). (b) One-day snapshot of cyclonic submesoscale eddies (blue curves)  
 351 identified from high-resolution chlorophyll data (colour shading;  $\text{mg m}^{-3}$ ). The eddy  
 352 edges are denoted by black dashed curves. (c) Same as Fig. 1b but for an anticyclonic  
 353 submesoscale eddy (red curves). (d) Distributions of cyclonic (blue dots) and  
 354 anticyclonic (red dots) submesoscale eddies identified in the South China Sea (SCS)  
 355 from 2006 to 2015.

356

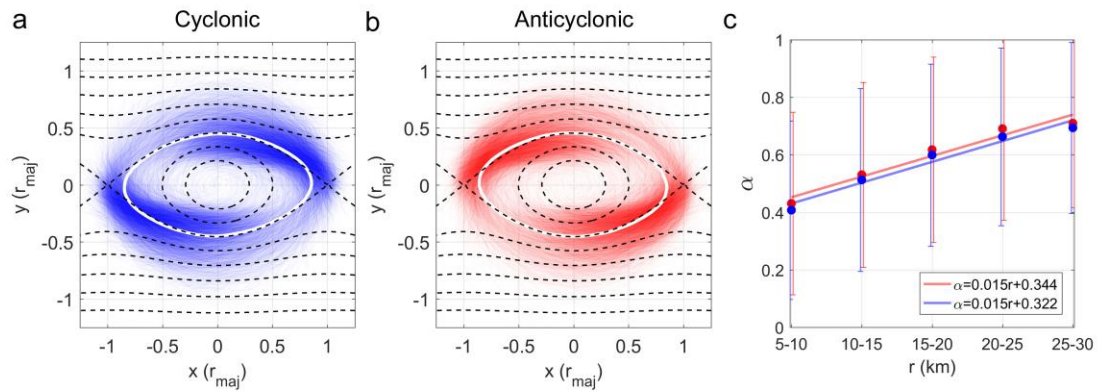




357

358 Figure 2. (a) Histogram of the radius of submesoscale eddies in the SCS. (b) Same as  
 359 Fig. 2a but for the eddy aspect ratio that is defined as the ratio between the minor and  
 360 major radius of a submesoscale eddy. (c) Variations of eddy aspect ratio with eddy  
 361 radius (averaged in an eddy-radius bin of 5 km). Vertical lines denote one standard  
 362 deviation.

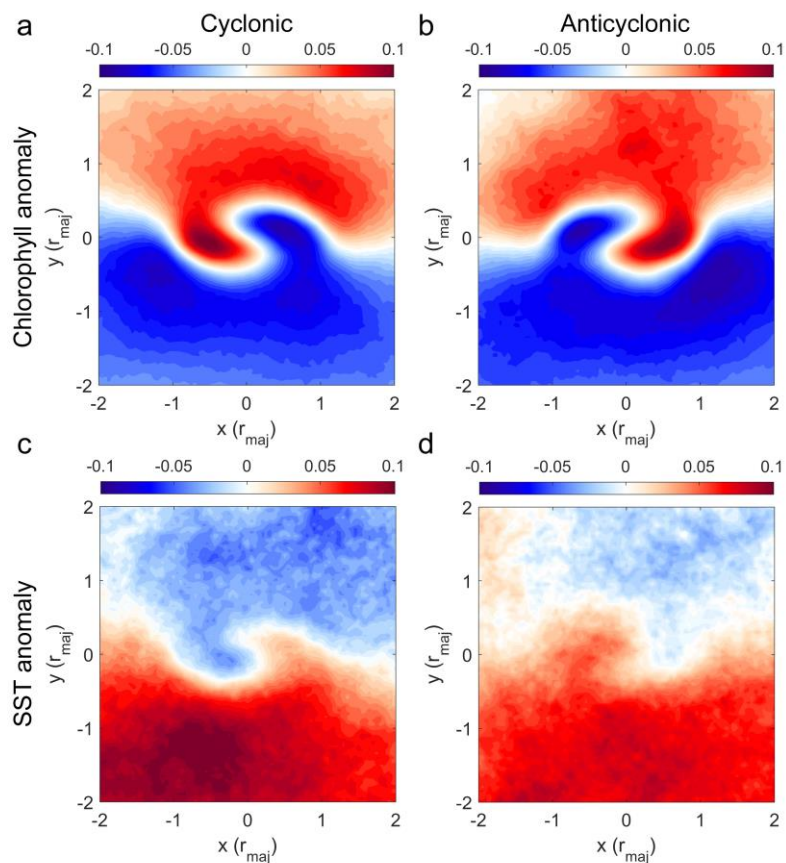
363



364

365 Figure 3. Horizontal structure of submesoscale eddies in the SCS. (a) Edges of  
 366 cyclonic eddies (blue curves) and their average (white curve) on a rotated  
 367 submesoscale eddy coordinate system (Supporting Information). Black dashed  
 368 contours are the horizontally normalized streamfunction contours derived from the  
 369 Stuart solution  $\psi = -U/k \cdot \log(\cosh(ky) - \alpha \cdot \cos(kx))$ , where  $U = \pm 0.3 \text{ m s}^{-1}$ ,  
 370  $k \approx 0.0003 \text{ m}^{-1}$ , and  $\alpha = 0.6$ . (b) Same as Fig. 3a but for anticyclonic eddies (red  
 371 curves). (c) Values of  $\alpha$  as a function of the radius of cyclonic (blue dots) and  
 372 anticyclonic (red dots) submesoscale eddies and the corresponding linear fitting  
 373 results (lines).

374



375

376 Figure 4. (a, b) Composite  $\log_{10}$ -transformed chlorophyll anomalies ( $\text{mg m}^{-3}$ ) on the  
 377 rotated submesoscale eddy coordinate. (c, d) Same as Fig. 4a, b but for SST anomalies  
 378 ( $^{\circ}\text{C}$ ).