

23 to present the accuracy of the two-phase model for the steam ejector simulation. This
24 demonstrates that the nonequilibrium condensation is essential for the performance
25 analysis of steam ejectors for MED-TVC seawater desalination system.

26 **Keywords:** multi-effect distillation (MED); thermal vapour compression (TVC); steam
27 ejector; seawater desalination; nonequilibrium condensation; wet steam

28 **1. Introduction**

29 The multi-effect distillation with thermal vapour compression (MED-TVC)
30 seawater desalination system provides a solution for freshwater production, which
31 achieves energy saving by upgrading the low-grade energy [1, 2]. The steam ejector, as
32 the main part of the TVC unit, recovers the low-pressure steam from the last effect of
33 evaporators, which accomplish high efficient utilisation of the low-grade energy [3, 4].
34 Thermodynamic and economic analysis has shown that the reduction of the exergy
35 destruction in the TVC was cost-effective for the entire seawater desalination system
36 [5]. The MED-TVC seawater desalination system also contributes to environmental
37 protection by integrating renewable energy [6-8].

38 The studies on steam ejectors for MED-TVC seawater desalination system
39 including the novel design [9], thermodynamic analysis [10] and optimisation studies
40 [11]. Ghaebi and Abbaspour [12] integrated heat recovery steam generators to MED-
41 TVC seawater desalination systems. The thermodynamic analysis illustrated that the
42 novel design improved the exergy efficiency by 57.5%. Tang et al. [13] designed a new
43 steam ejector with pressure regulation and the optimisation improved entrainment ratios
44 by 11.77%. Tang et al. [14] also reported that the multi-optimisation of the entrainment

45 passage can lead to an improvement of the entrainment ratio of 28.75% for the newly
46 designed steam ejector for MED-TVC seawater desalination systems. Gu et al. [15]
47 proposed a steam ejector with a variable geometry by inserting a spindle inside the
48 primary nozzle and averaged entrainment ratios reached 1.39 compared to the normal
49 ejector of 0.69. Xue et al. [16] reported a novel design of two-stage vacuum ejectors
50 and the experimental test demonstrated that the new design provided the vacuum
51 pressure of about 5.3 kPa compared to normal ejectors of 18.6 kPa. Sadeghi et al. [17]
52 proposed a thermodynamic study and optimised power - ejector - desalination systems,
53 and their study illustrated that the decrease of the total exergy destruction resulted in
54 high temperatures at the exit of the steam ejector to produce distilled water.

55 Park [18] optimised steam ejectors for the seawater desalination system by
56 introducing the swirling flow to the primary nozzle using computational fluid dynamics
57 (CFD) modelling, which could obtain high entrainment ratios by changing the strength
58 of the swirling flow. Sharif [19] carried out the 2D axisymmetric and three-dimensional
59 (3D) simulation of a steam ejector and the comparison revealed that the axisymmetric
60 simulation obtained similar results to 3D models considering the overall performance
61 of the steam ejector. Liu et al. [20] computationally investigated the influence of area
62 ratios on steam ejectors for MED-TVC seawater desalination systems and they found
63 that the entrainment ratio increased from 0.025 to 0.8. Khalid et al. [21] optimised the
64 location of steam ejectors for MED-TVC seawater desalination systems and best unit
65 performances could be achieved by integrating the steam ejector at the middle effect
66 regardless of the number of effects. Wang et al. [22] performed the optimising study on

67 primary nozzles within steam ejectors for MED-TVC seawater desalination systems
68 using CFD modelling and they showed that the overall efficiency of the steam ejector
69 increased by 14.41%.

70 The aforementioned CFD and thermodynamic analysis improved the
71 understanding of steam ejectors for MED-TVC seawater desalination systems ignoring
72 nonequilibrium condensation processes in steam ejectors. Fortunately, the importance
73 of the condensation phenomenon has been gradually realized in recent studies. Bonanos
74 [23] proposed a physical model of a steam ejector with assumptions of the one-
75 dimensional and perfect gas model without considering the condensation process, and
76 it was emphasized that the low temperatures due to the acceleration of the fluid may
77 lead to the condensation of the steam, thus invalidating the assumption of perfect gas
78 behaviour. Liu et al. [24] performed a thermodynamic analysis of steam ejectors for
79 MED-TVC seawater desalination units, which investigated the impact of condensation
80 behaviours on steam ejectors based on the homogeneous equilibrium theory assuming
81 that the phase change occurred instantaneously when the steam reached the saturation
82 state. Their study reported that the condensation is a pervasive phenomenon inside
83 steam ejectors. Tang et al. [25] carried out visualization experiments observing
84 condensation behaviours within steam ejectors for MED-TVC seawater desalination
85 systems. The experimental test demonstrated that massive condensing droplets were
86 observed with nonuniform distributions of the droplet size over the cross-area plane
87 within supersonic flows. These studies illustrate that nonequilibrium condensation
88 processes are significant for performance evaluations of steam ejectors for MED-TVC

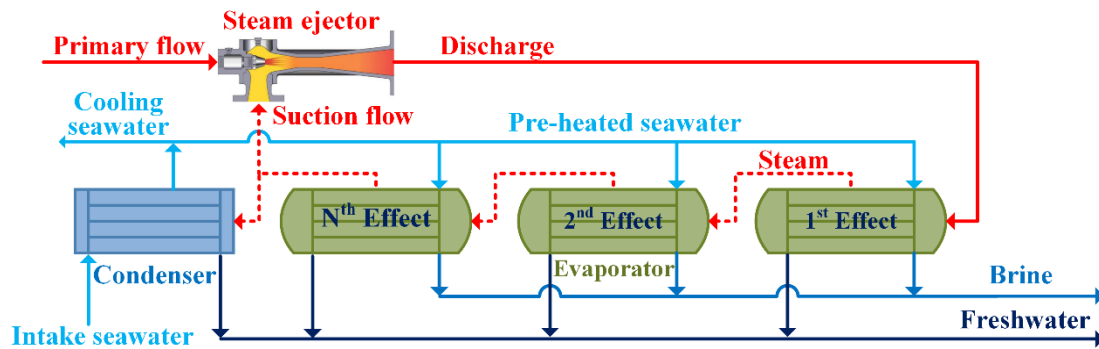
89 seawater desalination systems.

90 The present study fills in the scientific gap including the development of the two-
91 phase wet steam model for nonequilibrium condensation in supersonic flows, the
92 comparison of single-phase flow and two-phase flow models for steam ejectors and the
93 performance evaluation of a steam ejector for MED-TVC seawater desalination
94 systems based on the two-phase flow model. Specifically, the nonequilibrium
95 condensation process is integrated to develop a mathematical model to evaluate
96 performances of steam ejectors for MED-TVC seawater desalination systems. The
97 assessments of the single-phase and two-phase flow models are carried out to describe
98 the flow features within steam ejectors. The nonequilibrium condensation processes are
99 discussed inside steam ejectors for MED-TVC seawater desalination systems.

100 **2. MED-TVC seawater desalination system**

101 The MED-TVC seawater desalination system mainly consists of a seawater supply
102 unit, multi-evaporators, a steam ejector and a condenser, as shown in Fig. 1. Seawater
103 is fed in condensers with heat exchange between seawater and heating steam in the
104 parallel tubes. During this pre-process in the condenser, a small quantity of the feed
105 seawater is used as cooling water, while the main part is sprayed into each effect of the
106 multi-evaporators, where the seawater is heated to evaporate into water vapour. The
107 produced steam then goes into the next effect of multi-evaporators working as heating
108 fluids inside the parallel horizontal tubes which will condense to distilled water due to
109 the heat exchange with the feed seawater. The remaining liquid of the feed seawater is
110 processed as brine to be drained off from the system. These processes are repeated until

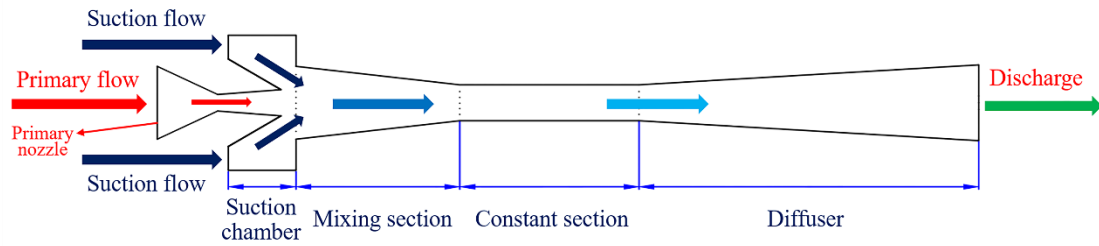
111 the last effect of multi-evaporators, where the main part of produced steam is entrained
112 into an ejector as a recycled steam system while the other part goes into the condenser
113 as the heating fluid to pre-heat the intake seawater.



114

115 Fig. 1 Schematic diagram of a typical MED-TVC seawater desalination system

116 The steam ejector is an important part of MED-TVC seawater desalination
117 systems. The motive primary flow expands in steam ejectors to entrain the low-pressure
118 steam from the last effect of multi-evaporators. The discharge flow of steam ejectors is
119 introduced into the first effect of the multi-evaporators as the heating fluid. This process
120 not only saves steam consumptions for high energy efficiencies but also allows low-
121 pressure operating conditions to eliminate the possibility of fouling and scaling. Figure
122 2 illustrates the schematic diagram of steam ejectors, which mainly includes five parts,
123 namely, a primary nozzle, suction chamber, mixing section, constant section and a
124 diffuser [26]. The dimension of the steam ejector employed for numerical simulations
125 is described in detail in Table 1.



126

127 Fig. 2 Schematic diagram of a steam ejector for MED-TVC seawater desalination

128

systems

129 Table 1 Dimension of the steam ejector for MED-TVC seawater desalination systems

Steam ejector dimension	Size (mm)
Inlet diameter of the primary nozzle	7.76
Throat diameter of the primary nozzle	1.76
Outlet diameter of the primary nozzle	7.00
Length of the converging part of the primary nozzle	28.70
Length of the diverging part of the primary nozzle	33.35
Inlet diameter of the mixing chamber	48.00
Length of the mixing chamber	130.00
Diameter of the constant section	19.00
Length of the constant section	95.00
Outlet diameter of the diffuser	40.00
Length of the diffuser	180.00

130 3. Mathematical model

131 Fundamental equations governing the flow features in a steam ejector are

132 compressible Navier-Stokes equations [27-29]. Considering the complicated flow

133 structure inside steam ejectors, including the supersonic flow, condensing flow, flow

134 separation and shock waves, the shear stress transport (SST) $k-\omega$ turbulence model is

135 employed for performance evaluations of steam ejectors for MED-TVC seawater

136 desalination systems.

137
$$\frac{\partial \rho}{\partial t} + \frac{\partial(\rho u_j)}{\partial x_j} = -\Gamma \quad (1)$$

$$138 \quad \frac{\partial}{\partial t}(\rho u_i) + \frac{\partial}{\partial x_j}(\rho u_j u_i) = -\frac{\partial p}{\partial x_i} + \frac{\partial \tau_{ij}}{\partial x_j} - u_i \Gamma \quad (2)$$

$$139 \quad \frac{\partial}{\partial t}(\rho H) + \frac{\partial}{\partial x_j}(\rho u_j H + p) = -\frac{\partial}{\partial x_j}(\lambda_{eff} \frac{\partial T}{\partial x_j}) + \frac{\partial}{\partial x_j}(u_i \tau_{ij}) - h_v \Gamma \quad (3)$$

$$140 \quad \frac{\partial}{\partial t}(\rho k) + \frac{\partial}{\partial x_i}(\rho k u_i) = \frac{\partial}{\partial x_j} \left(\Gamma_k \frac{\partial k}{\partial x_j} \right) + \bar{G}_k - Y_k + S_k \quad (4)$$

$$141 \quad \frac{\partial}{\partial t}(\rho \omega) + \frac{\partial}{\partial x_j}(\rho \omega u_j) = \frac{\partial}{\partial x_j} \left(\Gamma_\omega \frac{\partial \omega}{\partial x_j} \right) + C_\omega - Y_\omega + D_\omega + S_\omega \quad (5)$$

142 Nonequilibrium condensation processes in steam ejectors are governed by two
143 transport equations[30, 31]:

$$144 \quad \frac{\partial(\rho Y)}{\partial t} + \frac{\partial}{\partial x_j}(\rho Y u_j) = \Gamma \quad (6)$$

$$145 \quad \frac{\partial(\rho N)}{\partial t} + \frac{\partial}{\partial x_j}(\rho N u_j) = \rho J \quad (7)$$

146 where Γ illustrates the condensation rate of the steam:

$$147 \quad \Gamma = \frac{4\pi r_c^3}{3} \rho_l J + 4\pi r^2 \rho_l N \frac{dr}{dt} \quad (8)$$

$$148 \quad J = \frac{q_c}{1+\phi} \frac{\rho_v^2}{\rho_l} \sqrt{\frac{2\sigma}{\pi m_v^3}} \exp\left(-\frac{4\pi\sigma}{3k_B T_v} r_c^2\right) \quad (9)$$

$$149 \quad \frac{dr}{dt} = \frac{\lambda_v (T_s - T_v)}{\rho_l h_v r} \frac{(1 - r_c/r)}{\left(\frac{1}{1+2\beta Kn} + 3.78(1-\nu) \frac{Kn}{Pr} \right)} \quad (10)$$

150 where J is the nucleation rate, and dr/dt is the droplet growth rate.

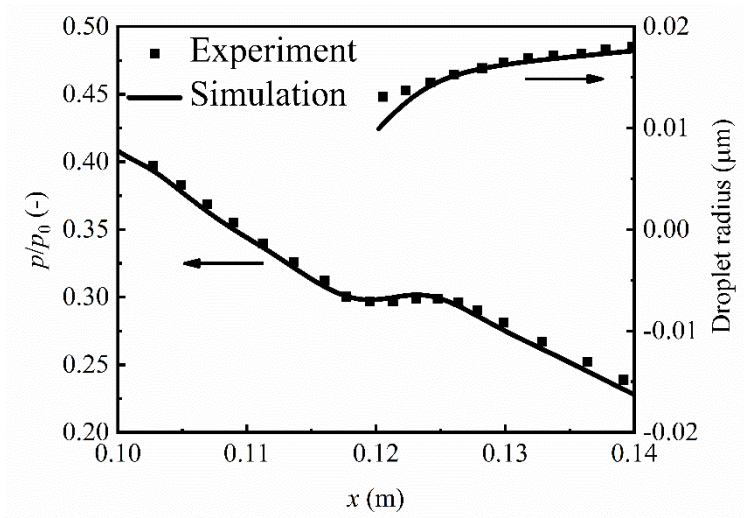
151 ANSYS FLUENT 19 [32] is utilized as a basic computational platform for
152 conservation equations of Eqs. (1)-(5). The additional transport equations of Eqs. (6) -
153 (7) and source terms of Eqs. (8) - (10) are integrated to ANSYS FLUENT by user-
154 defined-scalar and user-defined-function interfaces [33]. The total pressure and total

155 temperature boundary conditions are employed for the entrances and exit of the steam
156 ejector [34].

157 4. Results and discussion

158 4.1. CFD validation of nonequilibrium condensations in supersonic flows

159 The numerical validation of nonequilibrium condensations in supersonic flows is
160 carried out based on the case 252 of the Laval nozzle designed by Moses and Stein [35,
161 36], who performed the experimental measurement on the pressure distribution and the
162 droplet size of the wet steam flow. The supersonic flow condition is assigned to the exit
163 of the Laval nozzle. The computed pressure and droplet radius are compared against
164 experimental data, as shown in Fig. 3. It can be observed that numerical results agree
165 well with experimental data. Specifically, the developed CFD modelling predicts the
166 onset of nonequilibrium condensations while calculates similar droplet radius
167 compared to the values of experimental measurements. This demonstrates that the
168 developed CFD modelling accurately evaluates nonequilibrium condensations in
169 supersonic flows.

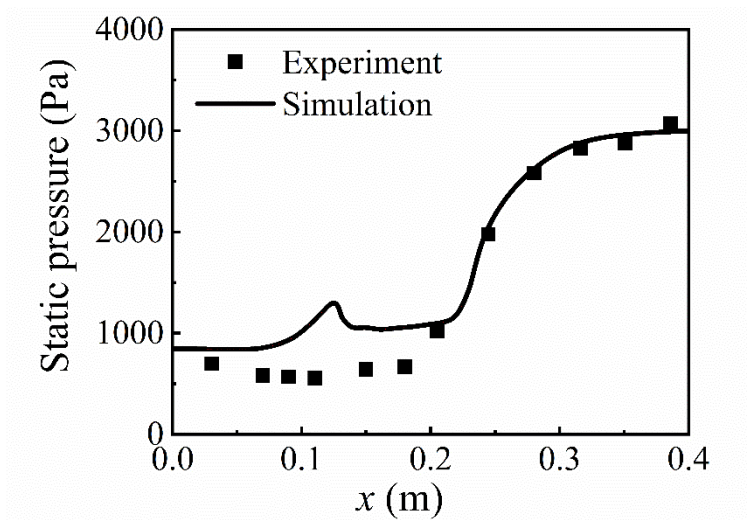


170

171 Fig. 3 CFD validation of nonequilibrium condensations in the supersonic flow

172 4.2. CFD validation of steam ejectors

173 Based on the successful validation of nonequilibrium condensations in the Moses and
174 Stein's Laval nozzle [35, 36], the developed mathematical model is further validated
175 for the steam ejector using experiment data reported by Sriveerakul et al. [37]. The
176 operating conditions include the temperature of the primary flow, $T_p=130$ °C, the
177 temperature of the suction flow, $T_s = 5$ °C, and the back pressure of discharge flow, P_b
178 = 3000 Pa. The numerical and experimental results are described in Fig. 4. The
179 computed pressure agrees well with experimental data in the diffuser while the
180 numerical results are larger than experimental values in the constant section, which is
181 attributed to the difficulty of calibrating the absolute pressure transducer near absolute
182 zero levels [37, 38]. Thus, it is concluded that the developed nonequilibrium
183 condensation model can be accurately used to evaluate a steam ejector for the MED-
184 TVC seawater desalination system.



185

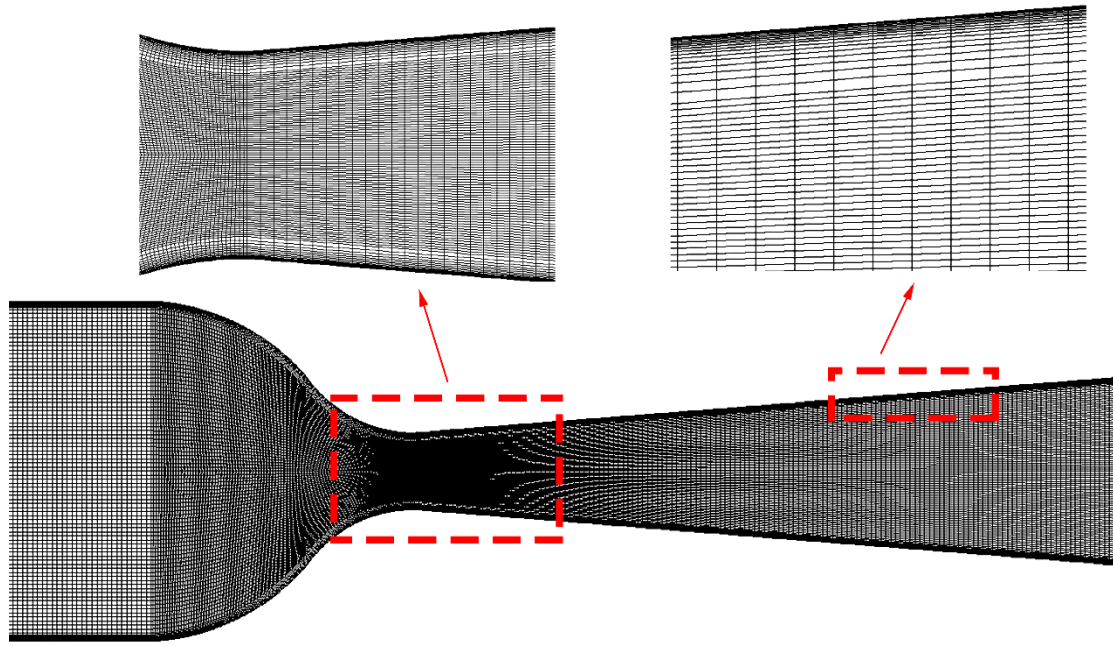
186 Fig. 4 CFD validation of a steam ejector for the MED-TVC seawater desalination

187

system

188 **4.3. Effect of single-phase and two-phase flow models**

189 This section focuses on flow behaviours in a primary nozzle of steam ejectors for
190 MED-TVC seawater desalination systems, which provides an opportunity to show local
191 details on formations and developments of nonequilibrium condensation processes in
192 supersonic flows. The quadrilateral mesh is generated for the primary nozzle, as shown
193 in Fig. 5. Three different grids are employed for analysis of mesh independence on
194 computational results with 16000, 42000 and 88000 cells, respectively. The flow
195 structure is employed to show the impact of grid densities on the supersonic flow. The
196 Mach numbers and nucleation rates in the primary nozzle are illustrated in Figs. 6 - 7.
197 As shown in the contours & profiles of Mach numbers, the difference is observed inside
198 the divergent part where the Mach number decreases because of heat and mass transfer
199 of phase change processes within supersonic flows. Likewise, when looking into the
200 detailed distribution of contours and profiles of nucleation rates, the coarse mesh
201 calculates condensation parameters slightly different from the other two meshes. The
202 coarse mesh predicts the peak nucleation rate of $4.52 \times 10^{25} \text{ m}^{-3} \text{ s}^{-1}$, while medium and
203 fine meshes calculate the maximum value of $1.53 \times 10^{26} \text{ m}^{-3} \text{ s}^{-1}$ and $1.94 \times 10^{26} \text{ m}^{-3} \text{ s}^{-1}$,
204 respectively. Thus, the medium mesh of 42000 quadrilateral cells is utilized to carry
205 out the computations of nonequilibrium condensations within the primary nozzle based
206 on computational accuracies and costs.

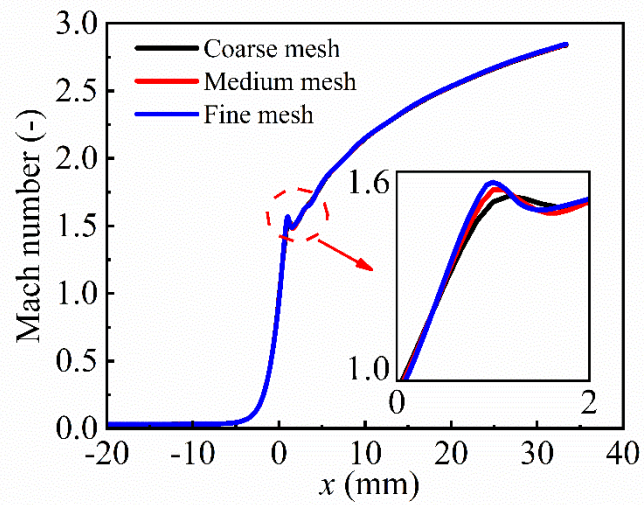


207

208 Fig. 5 Computational grid of the primary nozzle inside a steam ejector for the MED-

209

TVC seawater desalination system

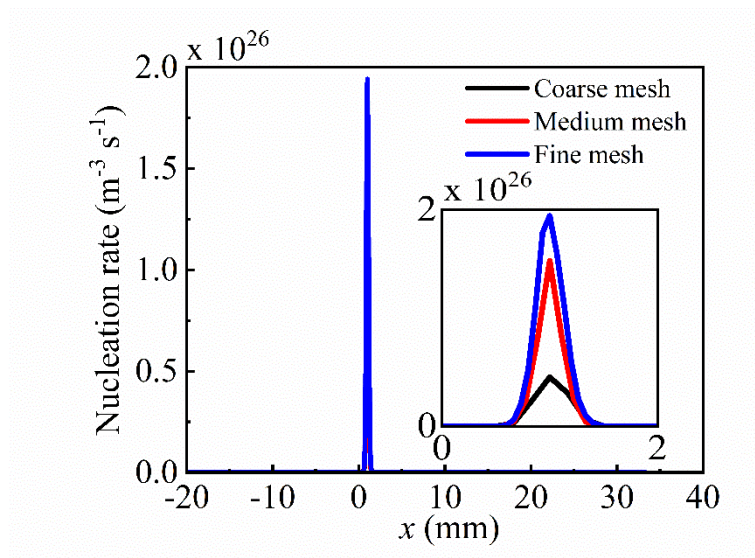


210

211 Fig. 6 Effect of grid density on the Mach number in the primary nozzle of a steam

212

ejector



213

214 Fig. 7 Effect of grid density on the nucleation rate in the primary nozzle of a steam

215

ejector

216

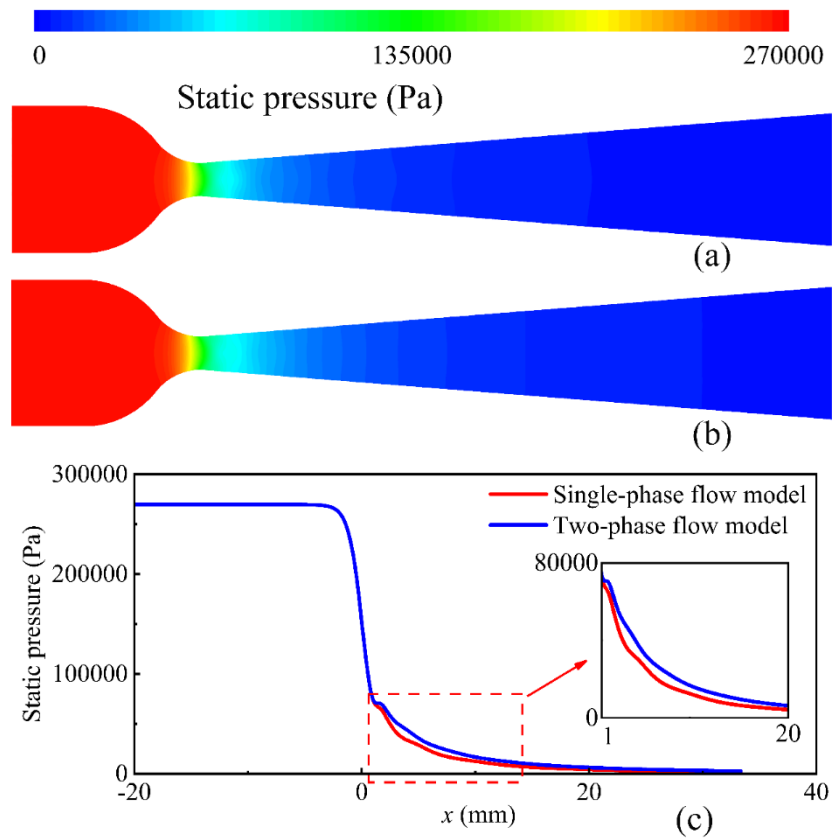
217 Steam flow reaches supersonic flows in the primary nozzle, where the rapid
218 expansion correspondingly induces a sharp decrease of the static pressures and
219 temperatures. The steam flow goes over the saturated-vapour line and reaches the
220 supersaturated state inside diverging parts of primary nozzles. The extreme
221 nonequilibrium state will lead to the occurrence of steam condensation, which is
222 expected to influence the flow behaviour in supersonic flows. After the grid
223 independence test, the single-phase and two-phase flow models are developed for
224 analysing nonequilibrium condensing processes inside supersonic flows within primary
225 nozzles.

225

226 Figures 8-10 illustrate the flow of the steam in the primary nozzle with and without
227 considering nonequilibrium condensations, including static pressures, static
228 temperatures and Mach numbers. It is observed that the single-phase flow represents
dry steam behaviour without considering the phase change. Under the assumption of

229 dry gas flow, the steam pressure and temperature continuously decrease within primary
230 nozzles. For the two-phase steam flow involving the condensation process, the steam
231 expands further to go across the saturated-vapour line, which is far from the equilibrium
232 state to induce condensation nuclei. The latent heat release during phase changes heats
233 the vapour phase and correspondingly results in a jump of the static pressure and static
234 temperature. This means that the steam condensation weakens the expansion
235 characteristics of supersonic flows within primary nozzles. For example, the Mach
236 numbers at the nozzle exit predicted by single- and two-phase flow models reach 3.92
237 and 2.84, respectively. The single-phase flow model overpredicts Mach numbers by 38%
238 compared to the two-phase flow model.

239 Even worse, the single-phase flow model predicts an unphysical decline of the
240 static temperature of the steam in the supersonic flow. In this primary nozzle, the single-
241 phase flow model predicts a static temperature of about 122 K at the nozzle exit, which
242 may raise the query of the formation of the ice in such a low temperature. The pressure-
243 temperature profiles of the water are shown in Fig. 11. It depicts that the pressure-
244 temperature profile calculated by the single-phase flow model crosses the melting line
245 and goes further to the solid region to generate the ice in such conditions. On the
246 contrary, the pressure-temperature crosses the saturation line and stays in the liquid
247 region due to the contribution of the phase change in supersonic flows. This
248 demonstrates that nonequilibrium condensation processes play a significant role in
249 performance evaluations of steam ejectors for MED-TVC seawater desalination
250 systems.



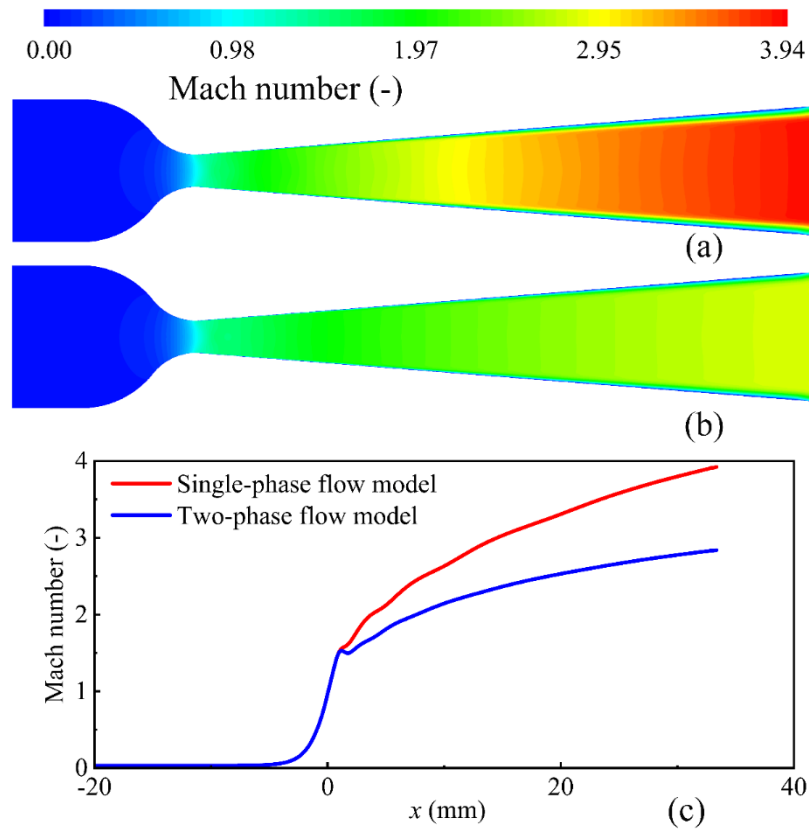
251

252 Fig. 8 Impact of condensing flow on static pressures within primary nozzles of steam

253 ejectors: contours of single- (a) and two-phase (b) flow models, and profiles at the

254 central axis of single- and two-phase flow models (c)

255



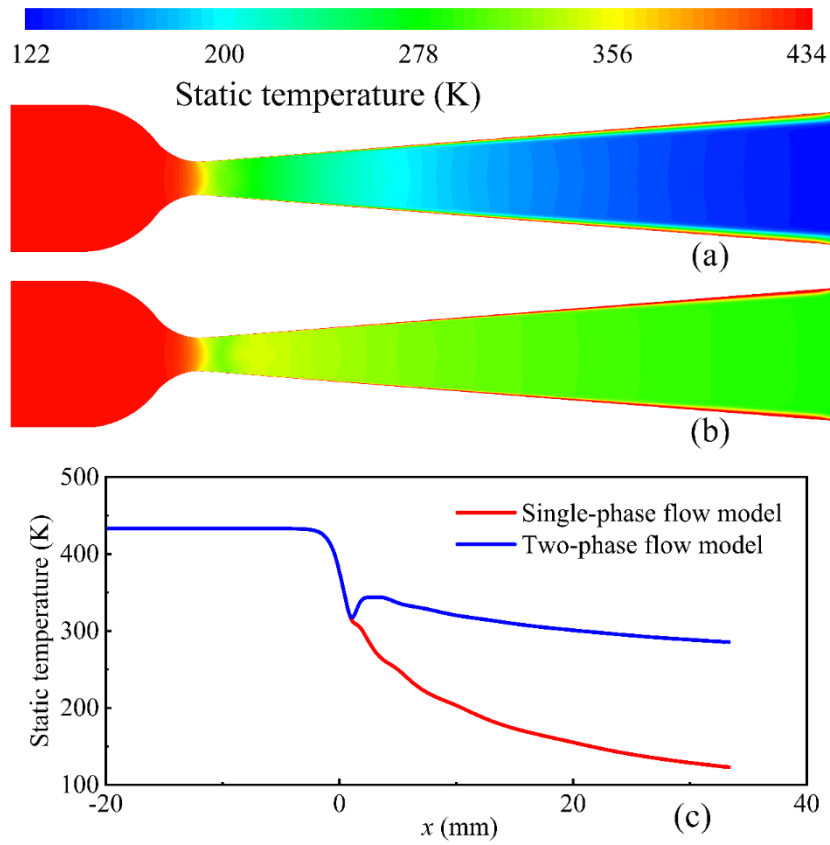
256

257 Fig. 9 Impact of condensing flow on Mach numbers within primary nozzles of steam

258 ejectors: contours of single- (a) and two-phase (b) flow models, and profiles at the

259 central axis of single- and two-phase flow models (c)

260

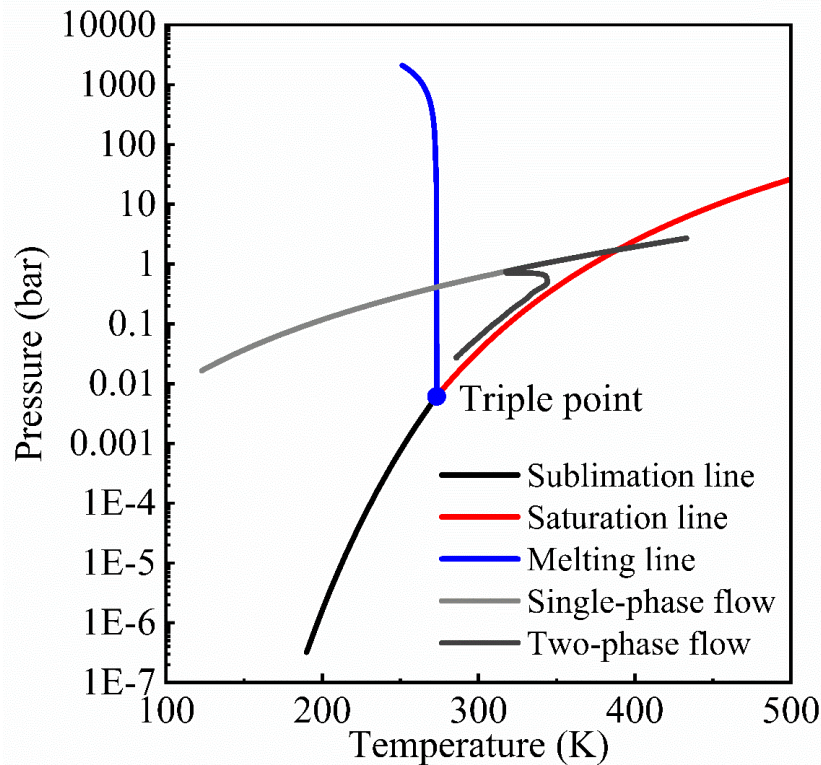


261

262 Fig. 10 Impact of condensing flow on static temperatures within primary nozzles of

263 steam ejectors: contours of single- (a) and two-phase (b) flow models, and profiles at

264 the central axis of single- and two-phase flow models (c)



265

266

Fig. 11 Steam pressure-temperature profiles

267

4.4. Performance analysis of steam ejectors for MED-TVC seawater desalination

268

systems

269

The quadrilateral structured mesh is generated for steam ejectors for MED-TVC

270

seawater desalination systems, as shown in Fig. 12. The impact of grid densities on

271

computational results of steam ejectors is carried out based on three quadrilateral

272

structured meshes including 39000, 73000 and 168000 cells, respectively. Table 2

273

describes the boundary conditions of steam ejectors for MED-TVC seawater

274

desalination systems. The high-pressure steam from the inlet of the primary nozzle is

275

assigned to 270 kPa and 463.15 K, which will expand in the steam ejector to entrain the

276

low-pressure steam from the suction chamber. The low-pressure steam from the last

277

effect of the MED unit will be recirculated to the suction chamber of the steam ejector

278

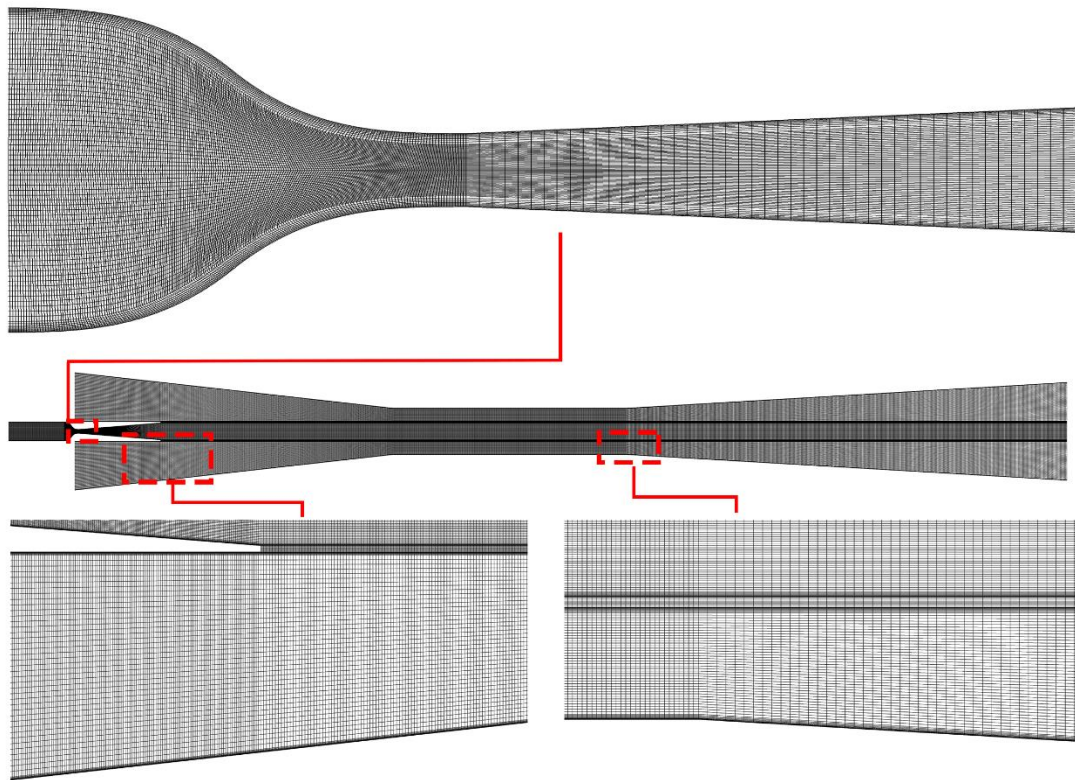
to improve the low-grade heat stream with the total pressure of 1.6 kPa and total

279 temperature of 287.20 K. The steam with the total pressure and total temperature of
 280 3.0kPa and 297.25 K is discharged from the exit of the ejector to the first effect of the
 281 MED unit. Fig. 13 represents the impact of mesh independence on flow behaviours in
 282 steam ejectors including static pressures, nucleation rates and entrainment ratios. It can
 283 be seen that the results based on the coarse mesh are a little bit far from the ones from
 284 medium and fine grids. Specifically, the coarse mesh predicts the maximum nucleation
 285 rate of about $4.53 \times 10^{22} \text{ m}^{-3} \text{ s}^{-1}$, while the medium and fine meshes compute the
 286 maximum value of $4.47 \times 10^{26} \text{ m}^{-3} \text{ s}^{-1}$ and $4.78 \times 10^{26} \text{ m}^{-3} \text{ s}^{-1}$, respectively. Consequently,
 287 the medium mesh is used to evaluate the performances of steam ejectors for MED-TVC
 288 seawater desalination systems.

289 Table 2 Boundary condition of steam ejectors for MED-TVC seawater desalination
 290 systems

Boundary conditions	Inlet of primary nozzle	Inlet of suction chamber	Exit of steam ejector
Total pressure	270 kPa	1.6 kPa	3.0 kPa
Total temperature	463.15 K	287.20 K	297.25 K

291



292

293

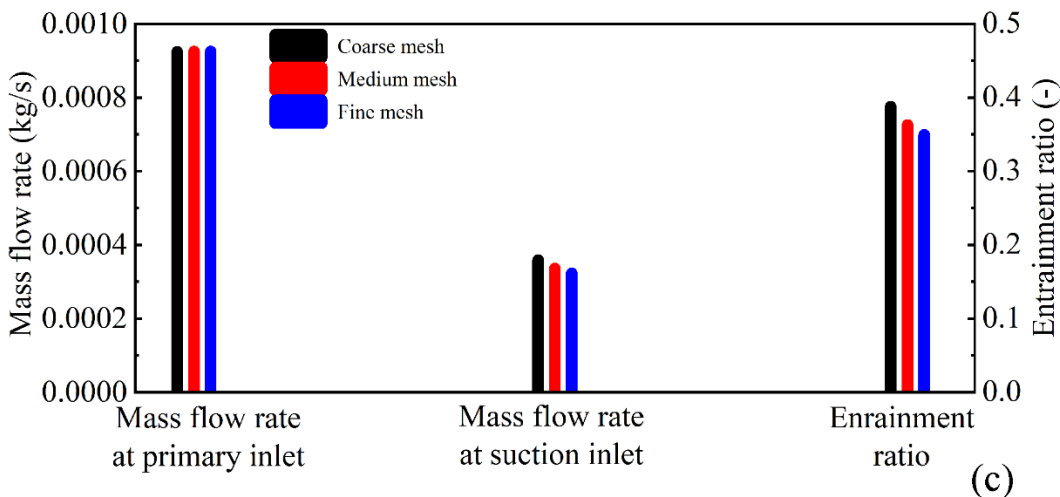
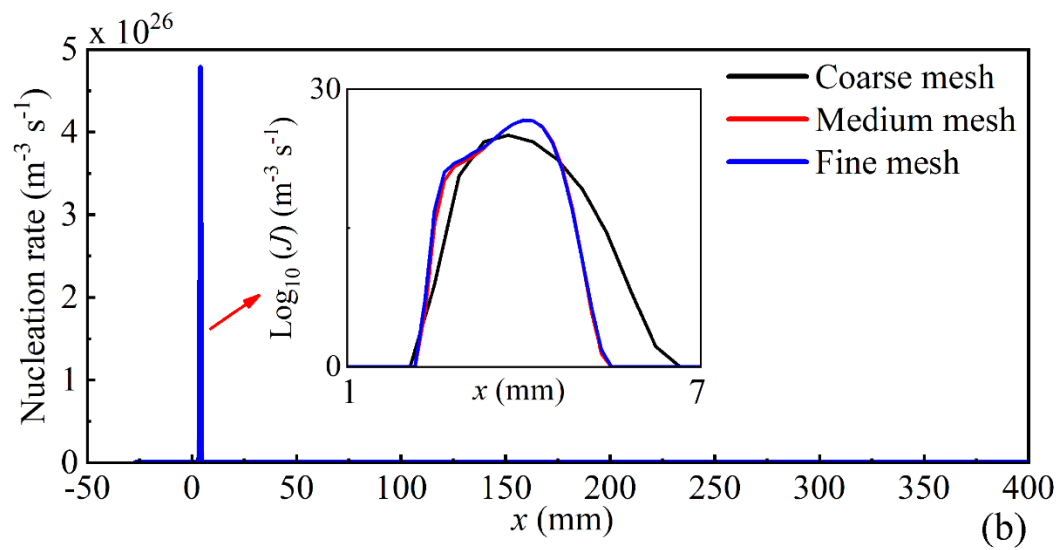
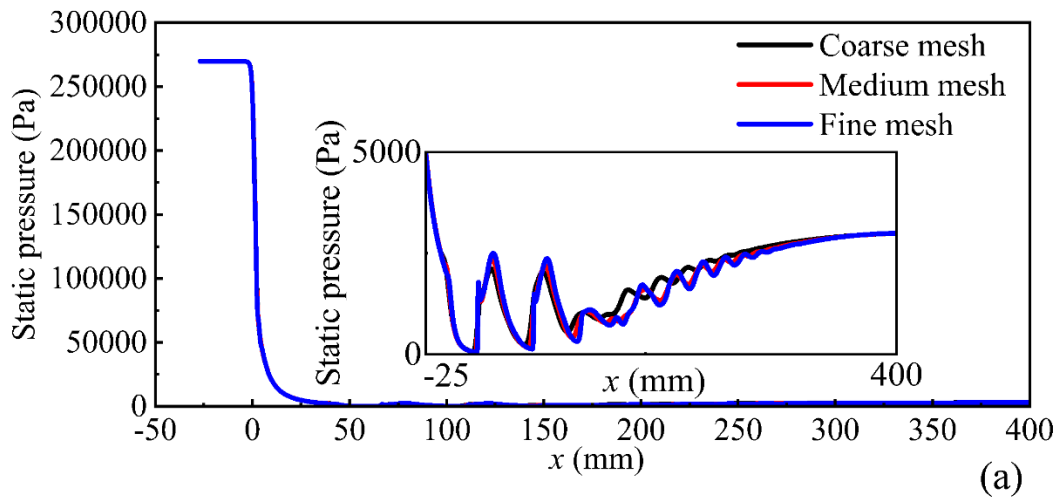
(b)

294

Fig. 12 Computational grids of the steam ejector for the MED-TVC seawater

295

desalination system



296

297 Fig. 13 Impact of the grid density on the steam ejector for the MED-TVC seawater

298

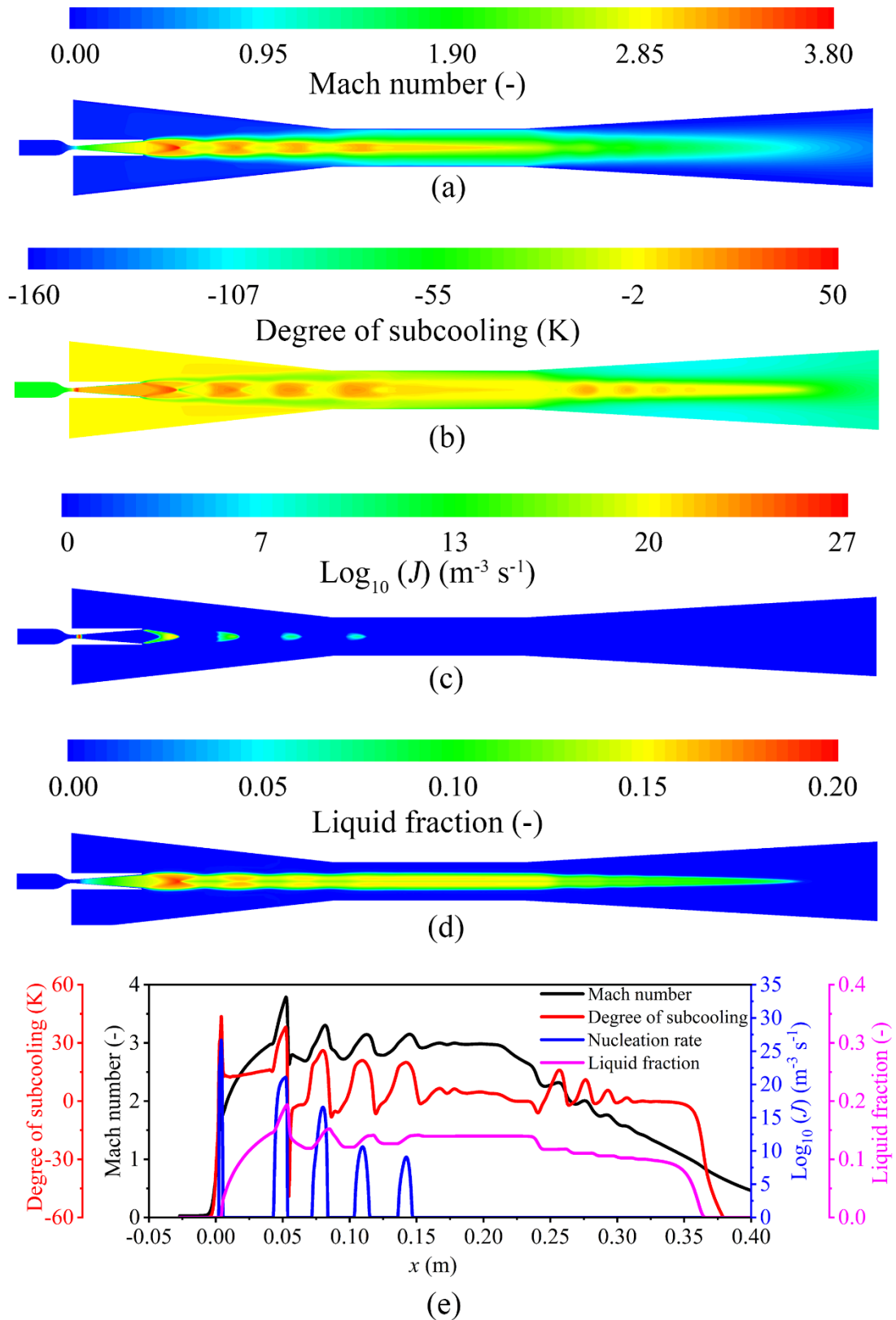
desalination system

299 Figure 14 illustrates flow structures in steam ejectors for MED-TVC seawater
300 desalination systems including Mach numbers, degrees of subcooling and nucleation
301 rates, respectively. The steam achieves supersonic flows inside the divergent section of
302 primary nozzles, which reaches extremely nonequilibrium state downstream the nozzle
303 throat with a maximum degree of subcooling of about 43 K. This induces the first
304 nucleation and condensation process of the steam with a peak nucleation rate of $4.47 \times$
305 $10^{26} \text{ m}^{-3} \text{ s}^{-1}$. The steam still maintains the supersonic flow downstream the condensation
306 region, where the degree of subcooling stays at approximately 12 K till the exit plane
307 of the primary nozzle.

308 The steam undergoes a further expansion-compression process to develop the
309 under-expanded supersonic flow inside mixing and constant sections, which induces
310 the intense change of the degree of subcooling. It can be observed that there are four
311 expansion-compression processes of the steam in mixing and constant sections. The
312 degree of subcooling increases to a peak value of about 38 K and then decreases sharply
313 to the minimum value of -49 K during the first expansion-compression process, which
314 leads to a condensation and re-evaporation process of the steam. Specifically, this
315 further nucleation and condensation process contributes to the maximum liquid fraction
316 of 0.19, which then declines to the trough level of 0.12 due to the re-evaporation process.
317 This alternating condensation and re-evaporation processes occur during the four
318 expansion-compression processes of the steam in mixing and constant sections of steam
319 ejectors for MED-TVC seawater desalination systems.

320 By combining nonequilibrium condensations in the primary nozzle and four

321 condensation and re-evaporation processes in mixing and constant sections of steam
322 ejectors, it is observed that the peak values of degree of subcooling are gradually
323 decreasing, while their trough levels are increasing. For instance, the first to fifth peak
324 values of the degree of subcooling are 43 K, 38 K, 26K, 21 K, 19K, and the first to
325 fourth trough levels are -49 K, -8 K, -6 K, -1 K, respectively. This correspondingly
326 decreases the peak value of the nucleation rate of $4.47 \times 10^{26} \text{ m}^{-3} \text{ s}^{-1}$, $1.22 \times 10^{21} \text{ m}^{-3} \text{ s}^{-1}$,
327 $3.85 \times 10^{16} \text{ m}^{-3} \text{ s}^{-1}$, $4.26 \times 10^{10} \text{ m}^{-3} \text{ s}^{-1}$ and $1.23 \times 10^9 \text{ m}^{-3} \text{ s}^{-1}$ for the first to fifth
328 condensation processes, which finally changes the first to the fourth peak value of the
329 liquid fraction of 0.19, 0.16, 0.143, 0.141, respectively. The degree of subcooling of the
330 steam remains above zero downstream the fourth expansion-compression flow in the
331 constant section of steam ejectors, which makes the liquid fraction stay at around 0.14
332 in this region.



333

334 Fig. 14 Flow structures in the steam ejector for the MED-TVC seawater desalination

335 system: contours of Mach numbers (a), contours of degrees of subcooling (b),

336 contours of nucleation rates (c), contours of liquid fractions (d), and profiles of Mach
337 numbers, degrees of subcooling, nucleation rates and liquid fractions at central lines
338 (e)

339 The flow structures represent the typical gas-liquid two-phase flows inside the
340 divergent part of primary nozzles, mixing and constant sections of steam ejectors for
341 MED-TVC seawater desalination systems, which are expected to affect energy losses
342 within steam ejectors. The entropy loss coefficient is employed here to evaluate the loss
343 of steam ejectors for MED-TVC seawater desalination systems. The entropy loss
344 coefficient is a function of the enthalpy, entropy and temperature, which is defined as
345 [39]:

$$346 \quad \eta = \frac{h_{in} - h_{out}}{h_{in} - h_{out} + T_{v,out} (s_{out} - s_{in})} \quad (11)$$

347 where η is the entropy loss coefficient, h_{in} and h_{out} are the enthalpies at ejector inlet and
348 outlet, s_{in} and s_{out} are the entropy at ejector inlet and outlet, $T_{v,out}$ is the steam
349 temperature at an ejector outlet.

350 Table 3 describes the thermodynamic parameters and entropy loss coefficient of
351 steam ejectors for MED-TVC seawater desalination systems based on single- and two-
352 phase flow models, including the enthalpy, entropy and temperature of the steam at the
353 inlet and out of steam ejectors. It can be found that the single-phase flow model ignoring
354 phase change processes computes larger enthalpy, entropy and temperature than the
355 two-phase flow model at ejector outlets. Subsequently, the single-phase flow model
356 calculates the entropy loss coefficient of 0.084 while the wet steam model predicts a
357 higher value of 0.097. That is, the dry gas model under-predicts the entropy loss

358 coefficient by approximately 15% compared to the two-phase flow model. It
 359 demonstrates that nonequilibrium condensations play a vital role in performance
 360 analyses of steam ejectors for MED-TVC seawater desalination systems.

361

362 Table 3 Entropy loss coefficient and thermodynamic parameters of steam ejectors for

363 MED-TVC seawater desalination systems based on single- and two-phase flow

364 models

	h_{in} (J/kg)	s_{in} (J/(kg.K))	h_{out} (J/kg)	s_{out} (J/(kg.K))	$T_{v,out}$ (K)	η (-)	Error (%)
Dry gas	751761.15	147.59	686752.18	2056.06	368.96	0.084	15%
Wet steam	751761.65	147.59	677778.21	2031.61	364.23	0.097	-

365

366 **5. Conclusions**

367 Performance analyses of a steam ejector are carried out for MED-TVC seawater
 368 desalination systems considering nonequilibrium condensation behaviours. The single-
 369 phase flow model with ignoring phase change processes predicts an unphysical
 370 temperature of the steam in the supersonic flow with the minimum value of
 371 approximately 122 K in the primary nozzle. The two-phase flow model predicts that
 372 the steam expands further and departs from equilibrium states to induce the occurrence
 373 of the nonequilibrium condensation in supersonic flows. The flow structures in steam
 374 ejectors for MED-TVC systems illustrates that the steam reaches extremely
 375 nonequilibrium states in the steam ejector to induce the nucleation and condensation
 376 processes in the primary nozzle, mixing and constant sections. The liquid fraction stays
 377 at around 0.14 in the constant section of the steam ejector due to the nonequilibrium

378 condensation. The single-phase flow model under-predicts the entropy loss coefficient
379 by approximately 15% compared to a two-phase flow model. This demonstrates that
380 nonequilibrium condensations play a crucial role in performance analyses of steam
381 ejectors for MED-TVC seawater desalination systems

382 **Acknowledgements**

383 This project has received funding from the European Union's Horizon 2020
384 research and innovation programme under the Marie Skłodowska-Curie grant
385 agreement No 792876 and the National Natural Science Foundation of China under
386 Grant 51876143.

387 **References**

- 388 [1] Y. Tang, Z. Liu, Y. Li, C. Shi, Combined auxiliary entrainment and structure
389 optimization for performance improvement of steam ejector with consideration of
390 back pressure variation, *Energy Conversion and Management*, 166 (2018) 163-173.
- 391 [2] L. Guimard, A. Cipollina, B. Ortega-Delgado, G. Micale, F. Couenne, P. Bandelier,
392 C. Jallut, New considerations for modelling a MED-TVC plant under dynamic
393 conditions, *Desalination*, 452 (2019) 94-113.
- 394 [3] Y. Tang, Z. Liu, C. Shi, Y. Li, A novel steam ejector with pressure regulation to
395 dredge the blocked entrained flow for performance improvement in MED-TVC
396 desalination system, *Energy Conversion and Management*, 172 (2018) 237-247.
- 397 [4] D. Strušnik, M. Marčič, M. Golob, A. Hribernik, M. Živić, J. Avsec, Energy
398 efficiency analysis of steam ejector and electric vacuum pump for a turbine
399 condenser air extraction system based on supervised machine learning modelling,

- 400 Applied Energy, 173 (2016) 386-405.
- 401 [5] M.L. Elsayed, O. Mesalhy, R.H. Mohammed, L.C. Chow, Exergy and thermo-
402 economic analysis for MED-TVC desalination systems, Desalination, 447 (2018)
403 29-42.
- 404 [6] A. Razmi, M. Soltani, M. Tayefeh, M. Torabi, M. Dusseault, Thermodynamic
405 analysis of compressed air energy storage (CAES) hybridized with a multi-effect
406 desalination (MED) system, Energy Conversion and Management, 199 (2019)
407 112047.
- 408 [7] B. Ghorbani, M. Miansari, S. Zendehboudi, M.-H. Hamed, Exergetic and economic
409 evaluation of carbon dioxide liquefaction process in a hybridized system of water
410 desalination, power generation, and liquefied natural gas regasification, Energy
411 Conversion and Management, 205 (2020) 112374.
- 412 [8] K. Mohammadi, M. Saghafifar, K. Ellingwood, K. Powell, Hybrid concentrated
413 solar power (CSP)-desalination systems: A review, Desalination, 468 (2019)
414 114083.
- 415 [9] Y. Tang, Z. Liu, Y. Li, C. Shi, H. Wu, Performance improvement of steam ejectors
416 under designed parameters with auxiliary entrainment and structure optimization
417 for high energy efficiency, Energy Conversion and Management, 153 (2017) 12-
418 21.
- 419 [10] I.B. Askari, M. Ameri, Techno economic feasibility analysis of Linear Fresnel solar
420 field as thermal source of the MED/TVC desalination system, Desalination, 394
421 (2016) 1-17.

- 422 [11] C. Wang, L. Wang, X. Wang, H. Zhao, Design and numerical investigation of an
423 adaptive nozzle exit position ejector in multi-effect distillation desalination system,
424 Energy, 140 (2017) 673-681.
- 425 [12] H. Ghaebi, G. Abbaspour, Thermo-economic analysis of an integrated multi-effect
426 desalination thermal vapor compression (MED-TVC) system with a trigeneration
427 system using triple-pressure HRSG, Heat and Mass Transfer, 54 (2018) 1337-1357.
- 428 [13] Y. Tang, Z. Liu, C. Shi, Y. Li, A novel steam ejector with pressure regulation to
429 optimize the entrained flow passage for performance improvement in MED-TVC
430 desalination system, Energy, 158 (2018) 305-316.
- 431 [14] Y. Tang, Z. Liu, Y. Li, C. Shi, C. Lv, A combined pressure regulation technology
432 with multi-optimization of the entrainment passage for performance improvement
433 of the steam ejector in MED-TVC desalination system, Energy, 175 (2019) 46-57.
- 434 [15] W. Gu, X. Wang, L. Wang, X. Yin, H. Liu, Performance investigation of an auto-
435 tuning area ratio ejector for MED-TVC desalination system, Applied Thermal
436 Engineering, 155 (2019) 470-479.
- 437 [16] H. Xue, L. Wang, L. Jia, C. Xie, Q. Lv, Design and investigation of a two-stage
438 vacuum ejector for MED-TVC system, Applied Thermal Engineering, 167 (2020)
439 114713.
- 440 [17] M. Sadeghi, M. Yari, S.M.S. Mahmoudi, M. Jafari, Thermodynamic analysis and
441 optimization of a novel combined power and ejector refrigeration cycle –
442 Desalination system, Applied Energy, 208 (2017) 239-251.
- 443 [18] I.S. Park, Numerical investigation of entraining performance and operational

- 444 robustness of thermal vapor compressor having swirled motive steam inflow,
445 Desalination, 257 (2010) 206-211.
- 446 [19] N. Sharifi, Axisymmetric and three dimensional flow modeling within thermal
447 vapor compressors, Heat and Mass Transfer, 49 (2013) 1489-1501.
- 448 [20] J. Liu, L. Wang, L. Jia, X. Wang, The influence of the area ratio on ejector
449 efficiencies in the MED-TVC desalination system, Desalination, 413 (2017) 168-
450 175.
- 451 [21] K.A. Khalid, M.A. Antar, A. Khalifa, O.A. Hamed, Allocation of thermal vapor
452 compressor in multi effect desalination systems with different feed configurations,
453 Desalination, 426 (2018) 164-173.
- 454 [22] K. Wang, L. Wang, L. Jia, W. Cai, R. Gao, Optimization design of steam ejector
455 primary nozzle for MED-TVC desalination system, Desalination, 471 (2019)
456 114070.
- 457 [23] A.M. Bonanos, Physical modeling of thermo-compressor for desalination
458 applications, Desalination, 412 (2017) 13-19.
- 459 [24] J. Liu, L. Wang, L. Jia, H. Xue, Thermodynamic analysis of the steam ejector for
460 desalination applications, Applied Thermal Engineering, 159 (2019) 113883.
- 461 [25] Y. Tang, Z. Liu, Y. Li, H. Wu, X. Zhang, N. Yang, Visualization experimental study
462 of the condensing flow regime in the transonic mixing process of desalination-
463 oriented steam ejector, Energy Conversion and Management, 197 (2019) 111849.
- 464 [26] C. Wen, B. Rogie, M.R. Kærn, E. Rothuizen, A first study of the potential of
465 integrating an ejector in hydrogen fuelling stations for fuelling high pressure

- 466 hydrogen vehicles, *Applied Energy*, 260 (2020) 113958.
- 467 [27] C. Wen, N. Karvounis, J.H. Walther, H. Ding, Y. Yang, Non-equilibrium
468 condensation of water vapour in supersonic flows with shock waves, *International*
469 *Journal of Heat and Mass Transfer*, 149 (2020) 119109.
- 470 [28] H. Ding, Y. Li, E. Lakzian, C. Wen, C. Wang, Entropy generation and exergy
471 destruction in condensing steam flow through turbine blade with surface roughness,
472 *Energy Conversion and Management*, 196 (2019) 1089-1104.
- 473 [29] C. Wen, N. Karvounis, J.H. Walther, Y. Yan, Y. Feng, Y. Yang, An efficient
474 approach to separate CO₂ using supersonic flows for carbon capture and storage,
475 *Applied energy*, 238 (2019) 311-319.
- 476 [30] A. Kantrowitz, Nucleation in very rapid vapor expansions, *The Journal of chemical*
477 *physics*, 19 (1951) 1097-1100.
- 478 [31] J. Young, The spontaneous condensation of steam in supersonic nozzle, *Physico*
479 *Chemical Hydrodynamics*, 3 (1982) 57-82.
- 480 [32] ANSYS Fluent Theory Guide, ANSYS Inc., USA, (2017).
- 481 [33] Y. Yang, J.H. Walther, Y. Yan, C. Wen, CFD modeling of condensation process of
482 water vapor in supersonic flows, *Applied Thermal Engineering*, 115 (2017) 1357-
483 1362.
- 484 [34] Y. Yang, X. Zhu, Y. Yan, H. Ding, C. Wen, Performance of supersonic steam
485 ejectors considering the nonequilibrium condensation phenomenon for efficient
486 energy utilisation, *Applied Energy*, 242 (2019) 157-167.
- 487 [35] C.A. Moses, G.D. Stein, On the growth of steam droplets formed in a Laval nozzle

488 using both static pressure and light scattering measurements, *Journal of Fluids*
489 *Engineering*, 100 (1978) 311-322.

490 [36] J. Starzmann, F.R. Hughes, S. Schuster, A.J. White, J. Halama, V. Hric, M.
491 Kolovratník, H. Lee, L. Sova, M. Št'astný, M. Grübel, M. Schatz, D.M. Vogt, Y.
492 Patel, G. Patel, T. Turunen-Saaresti, V. Gribin, V. Tishchenko, I. Gavrillov, C. Kim,
493 J. Baek, X. Wu, J. Yang, S. Dykas, W. Wróblewski, S. Yamamoto, Z. Feng, L. Li,
494 Results of the International Wet Steam Modeling Project, *Proceedings of the*
495 *Institution of Mechanical Engineers, Part A: Journal of Power and Energy*, 232
496 (2018) 550-570.

497 [37] T. Sriveerakul, S. Aphornratana, K. Chunnanond, Performance prediction of steam
498 ejector using computational fluid dynamics: Part 1. Validation of the CFD results,
499 *International Journal of Thermal Sciences*, 46 (2007) 812-822.

500 [38] T. Sriveerakul, S. Aphornratana, K. Chunnanond, Performance prediction of steam
501 ejector using computational fluid dynamics: Part 2. Flow structure of a steam
502 ejector influenced by operating pressures and geometries, *International Journal of*
503 *Thermal Sciences*, 46 (2007) 823-833.

504 [39] S. Dykas, M. Majkut, K. Smółka, M. Stozik, Study of the wet steam flow in the
505 blade tip rotor linear blade cascade, *International Journal of Heat and Mass*
506 *Transfer*, 120 (2018) 9-17.

507